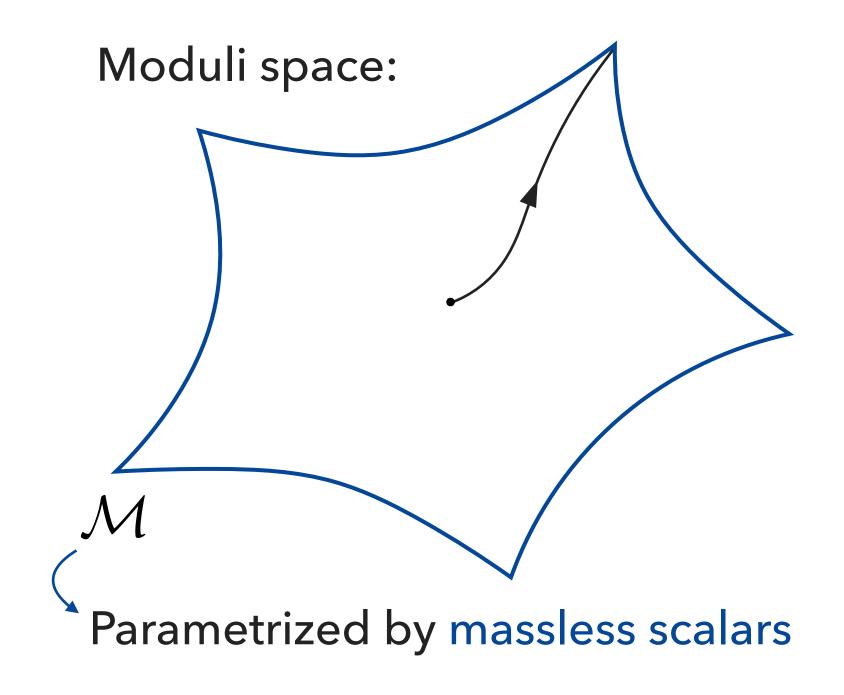
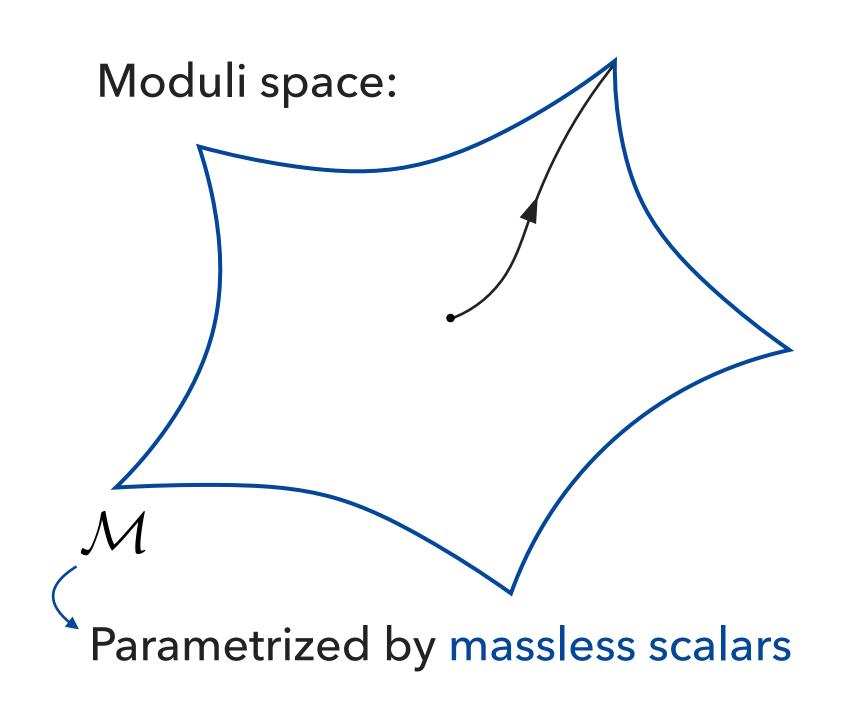
Conserved Currents at Infinite Distance in the Conformal Manifold José Calderón Infante Based on 2305.05693 with Florent Baume

String Phenomenology 2023, IBS Korea, 06/07/2023





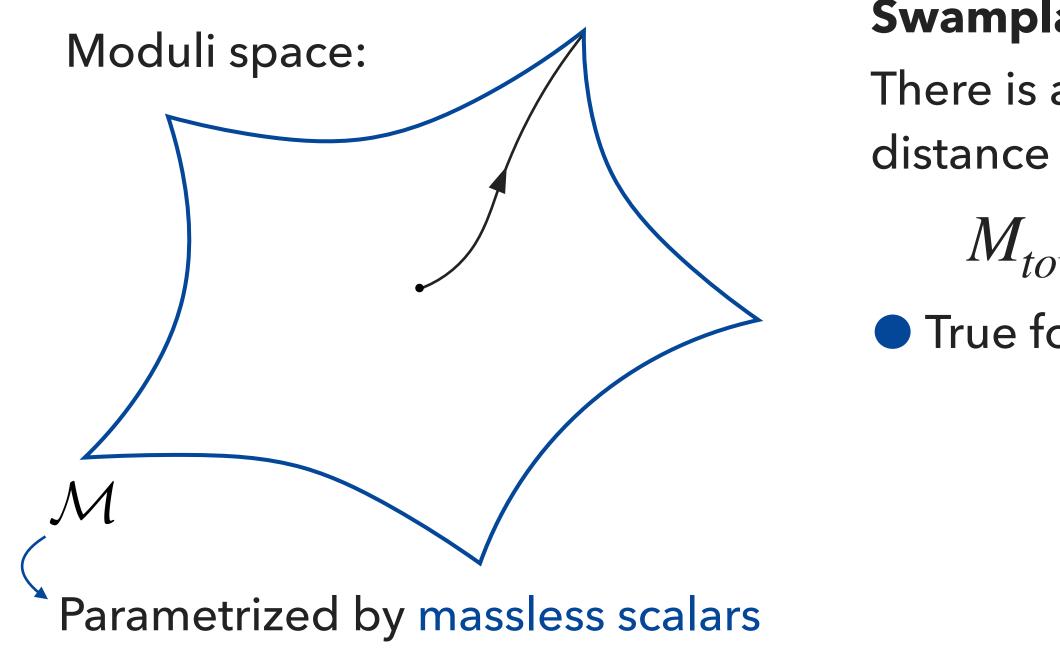


Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

 M_{to}

wer
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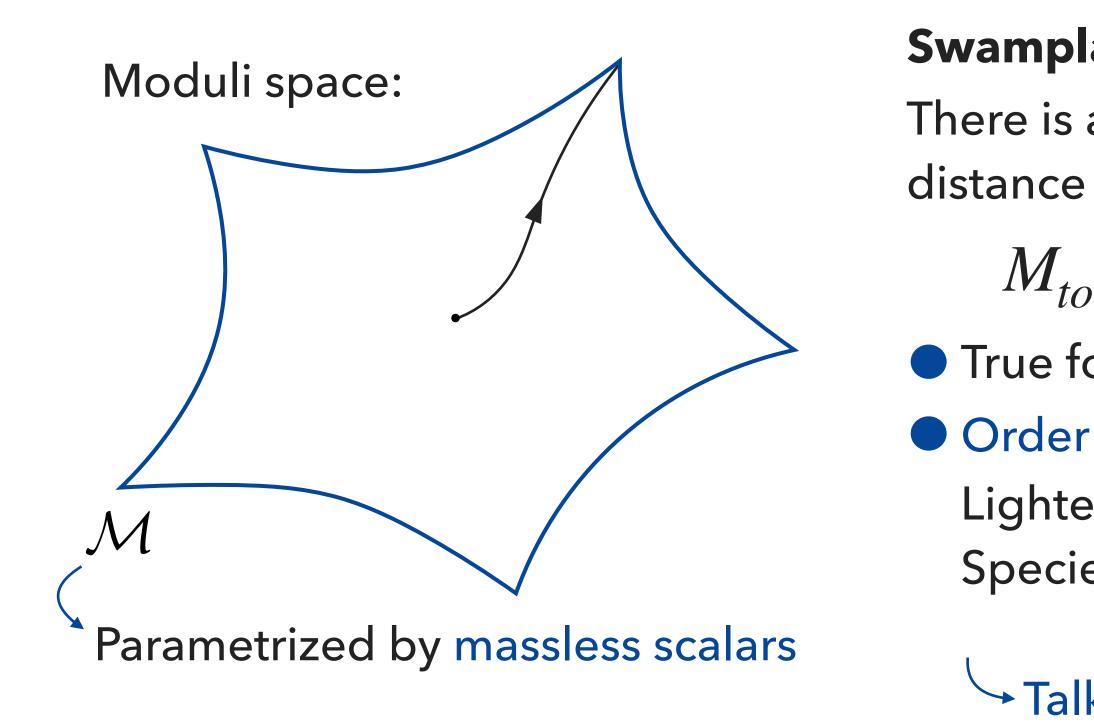


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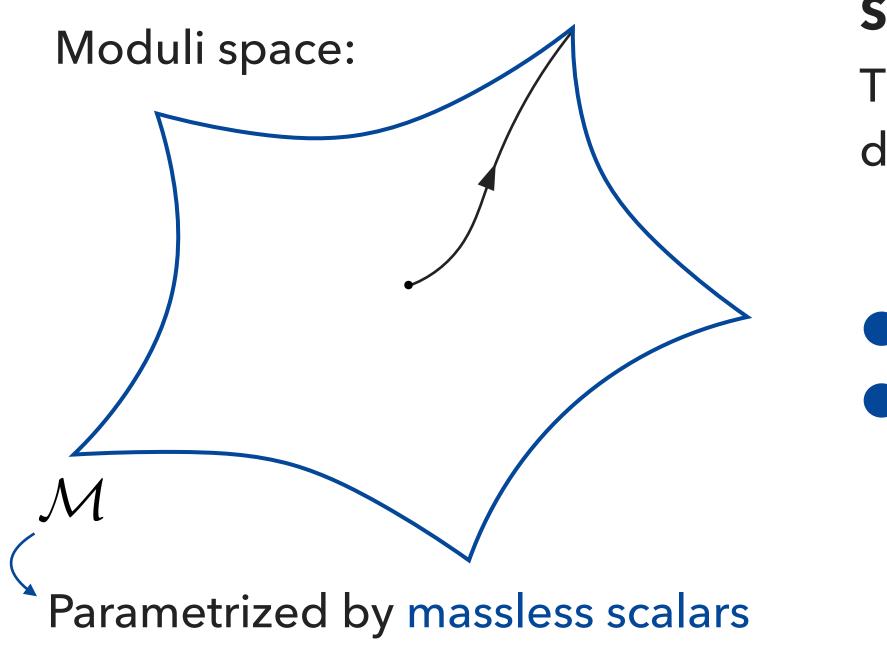
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→ Talks by Alberto, Álvaro, Ben, Ignacio, Irene and Max



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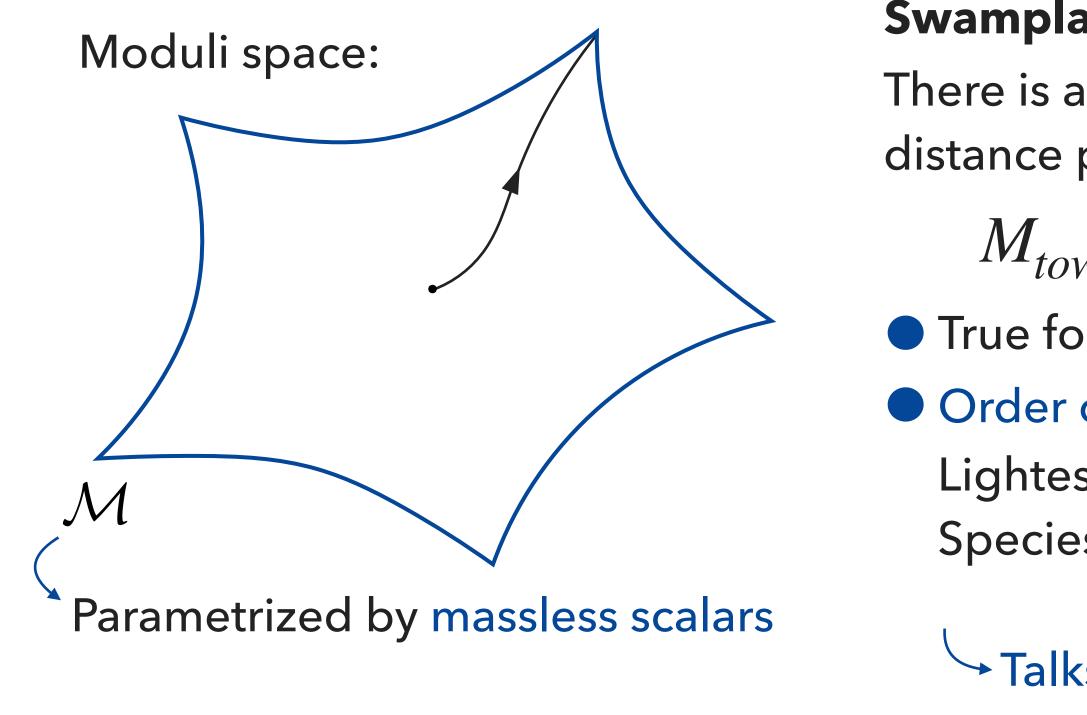
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Lots of top-down evidence!

- String theory flat space vacua: [Grimm, Palti, Valenzuela '18] [Lee, Lerche, Weigand '18-'19] + many many more!
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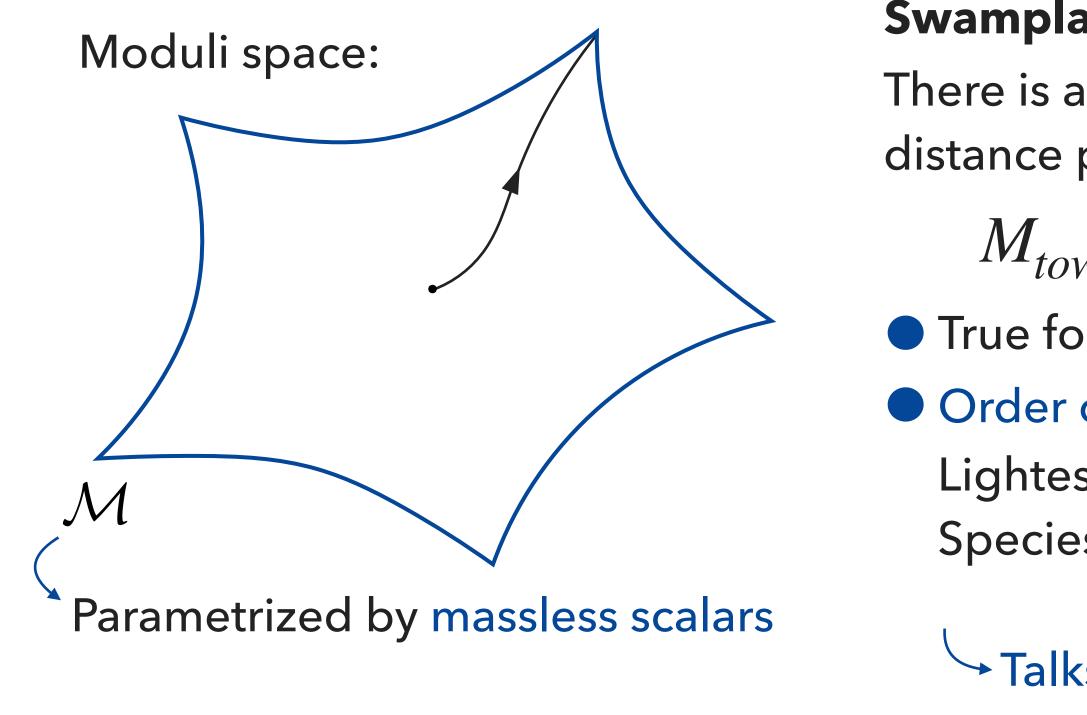
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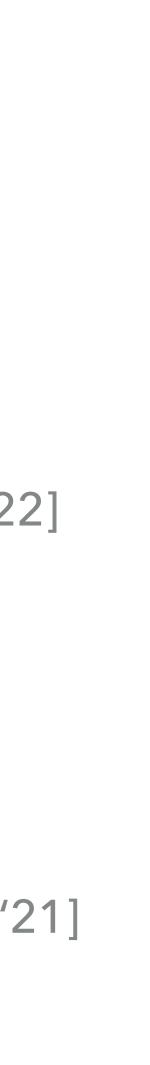
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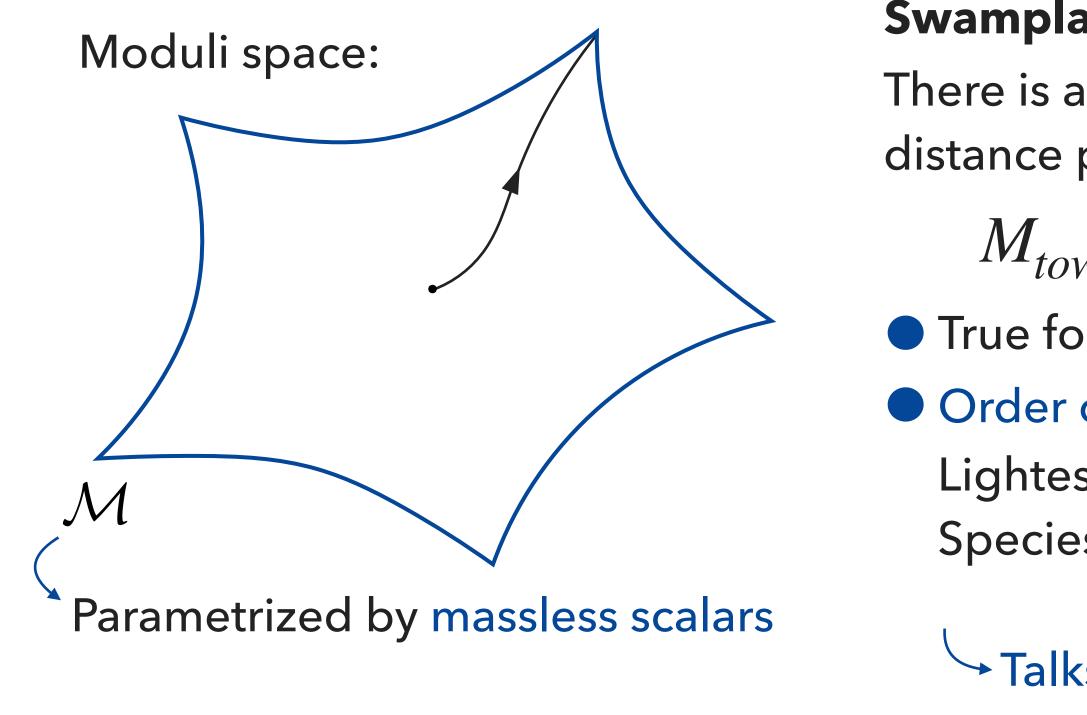
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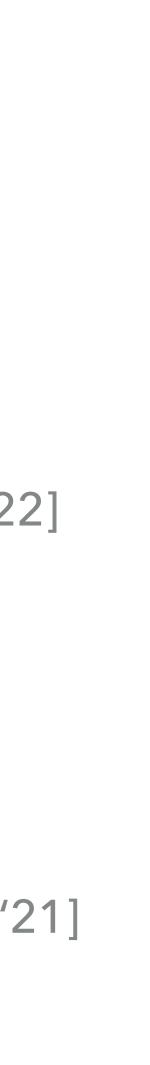
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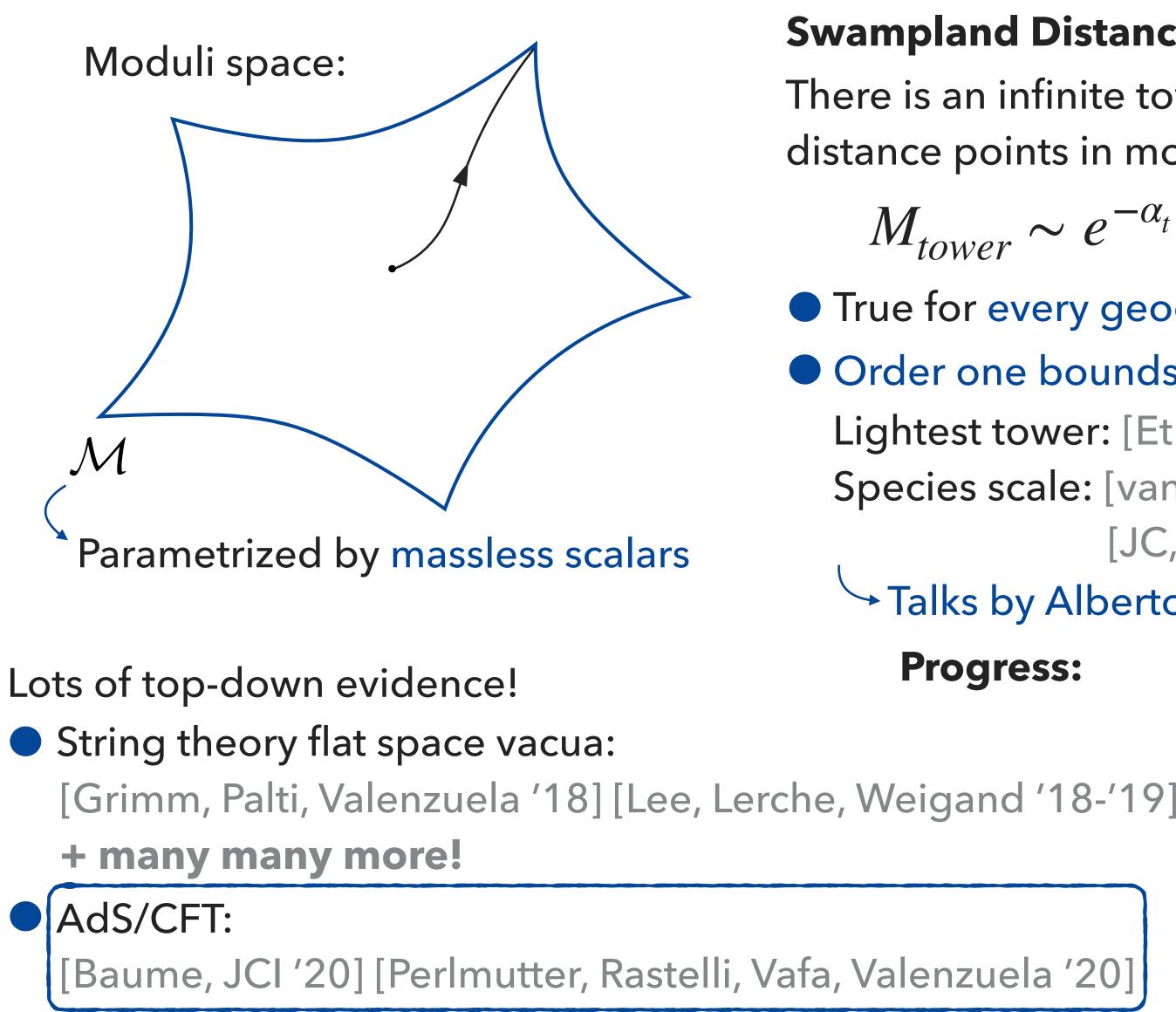
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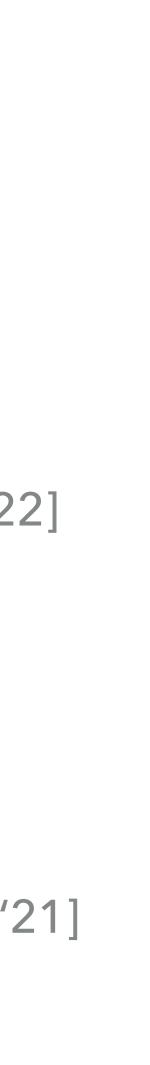
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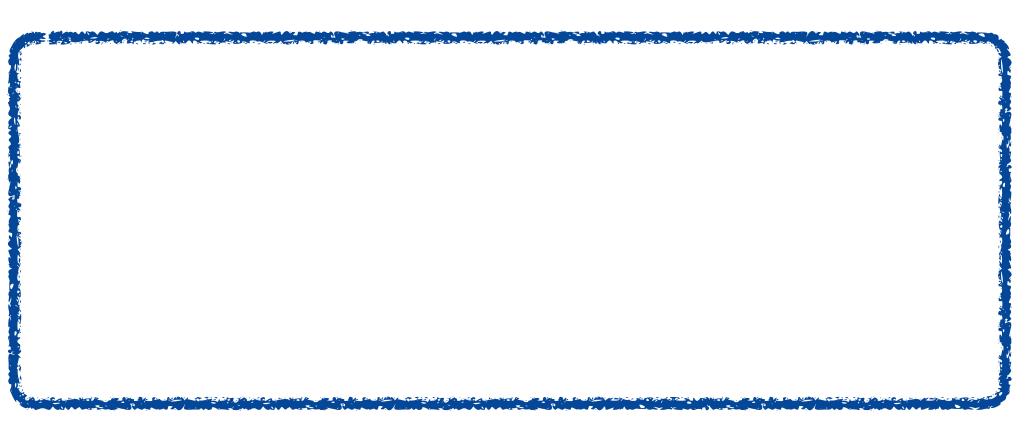


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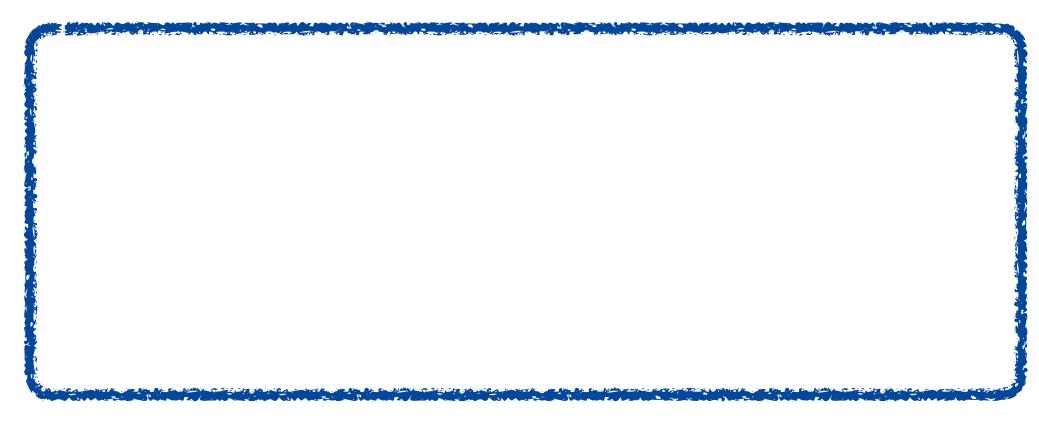
Conformal manifold of local CFT in d>2

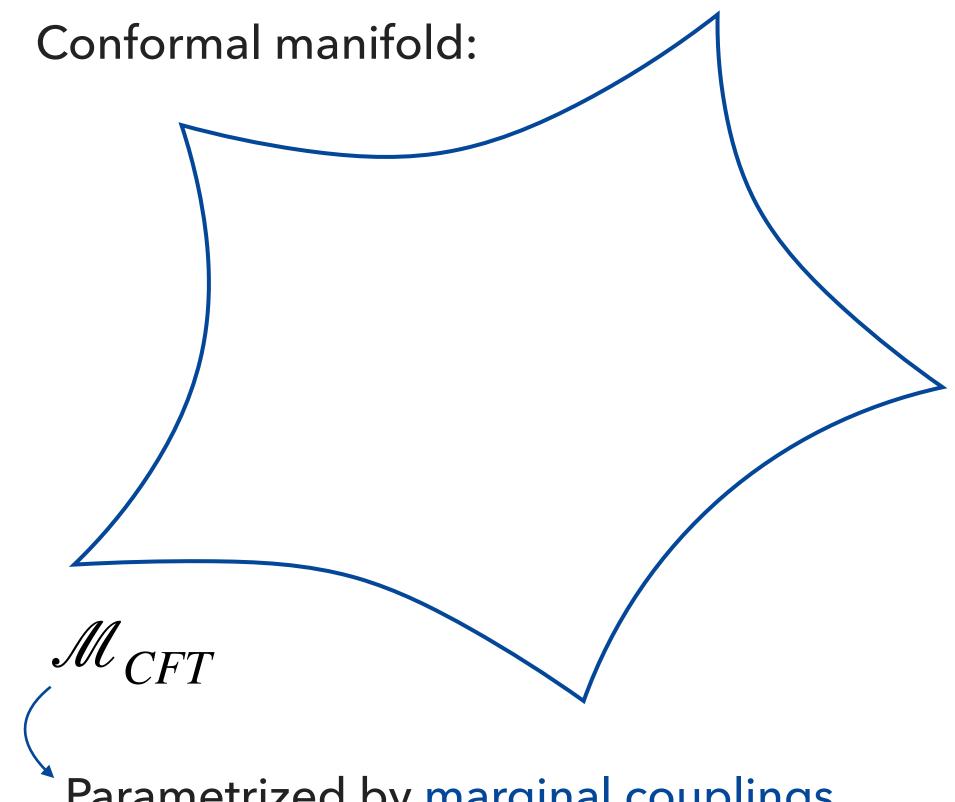


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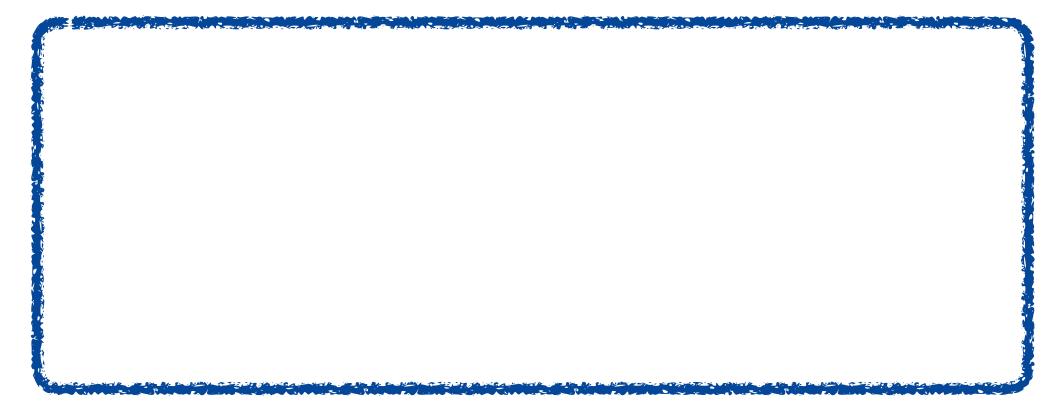




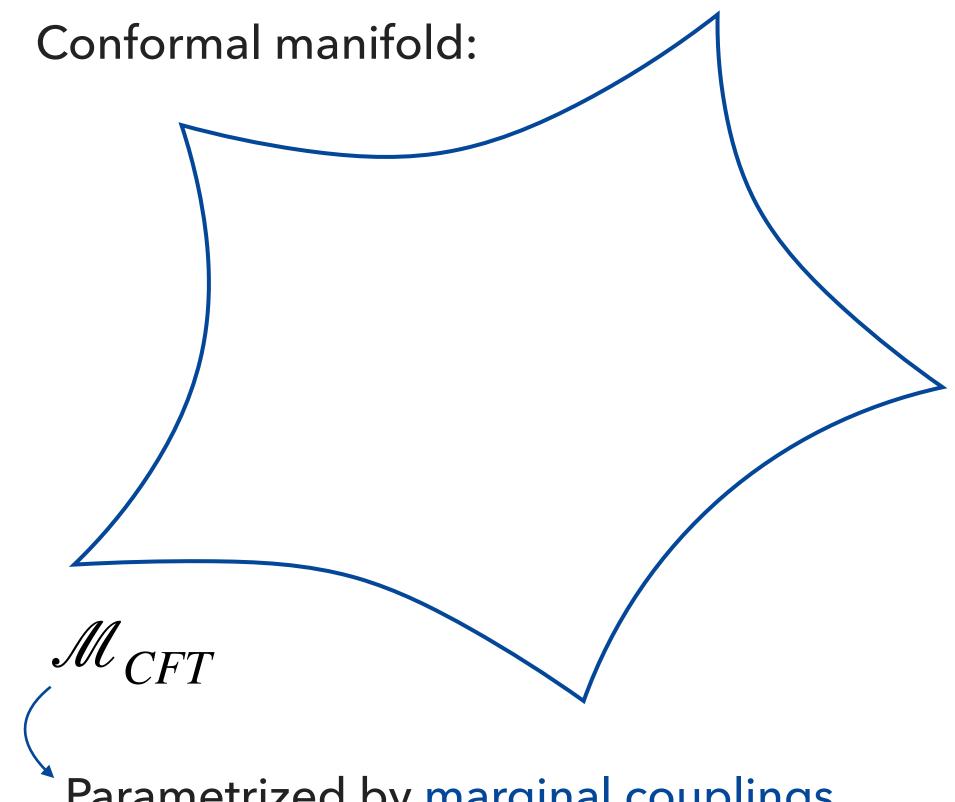
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Local CFT: Posses energy-momentum tensor



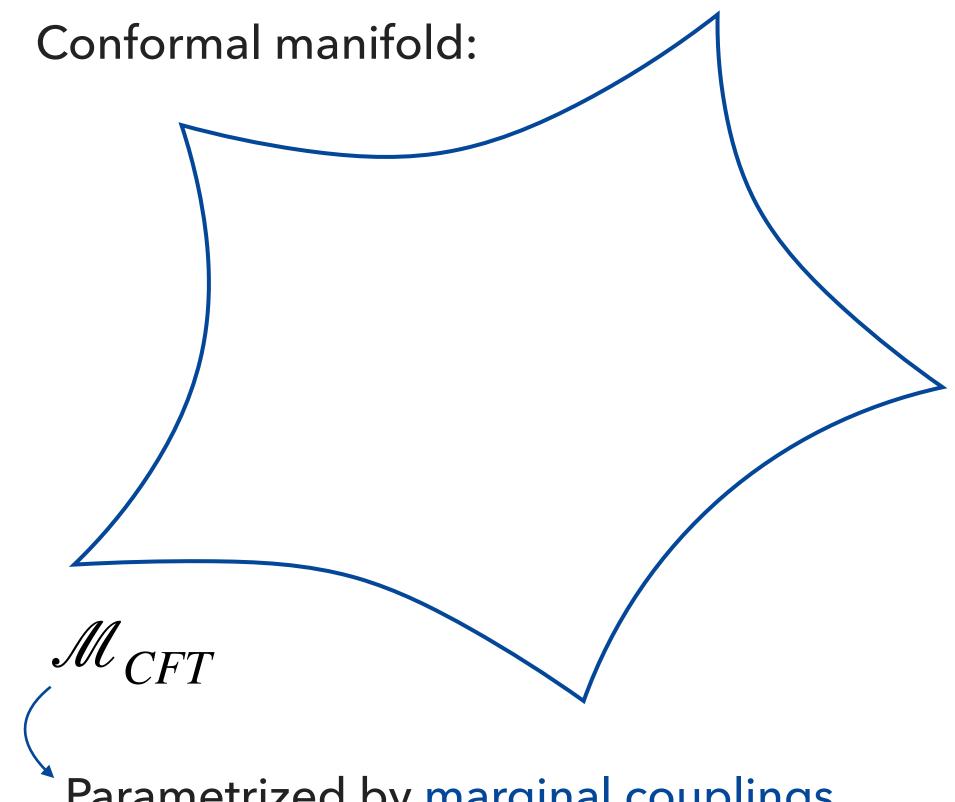
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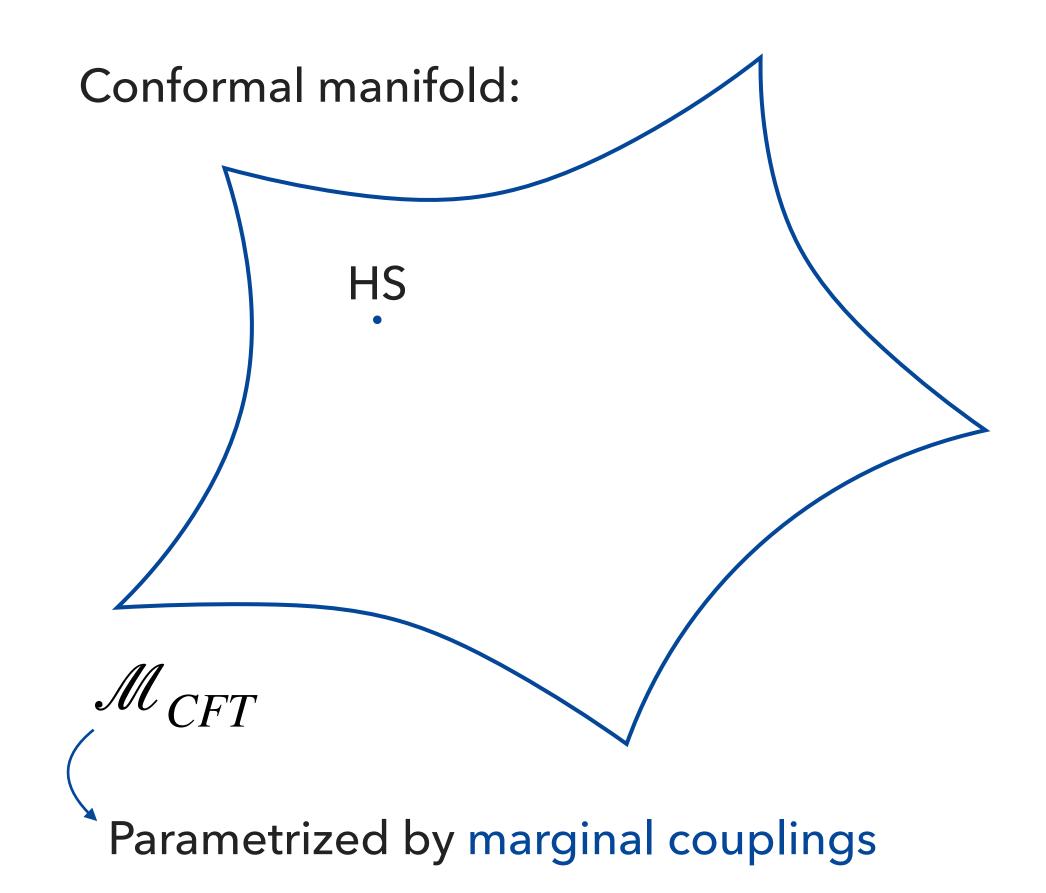
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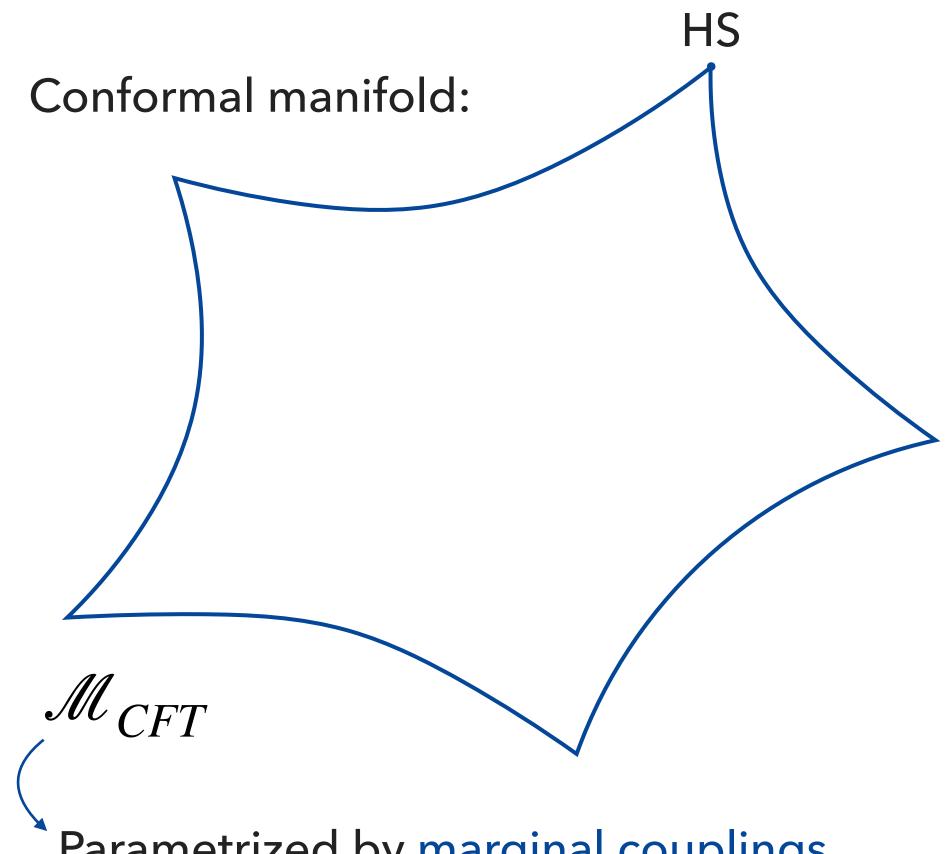
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Dynamical gravity in the bulk!



Parametrized by marginal couplings

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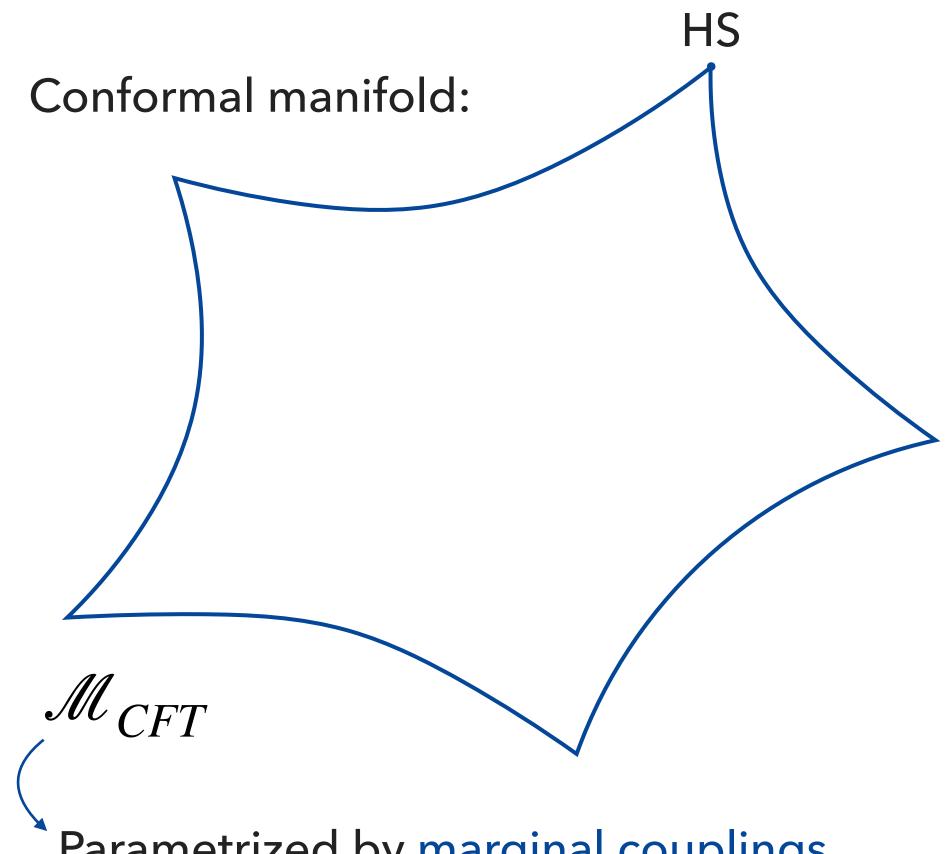
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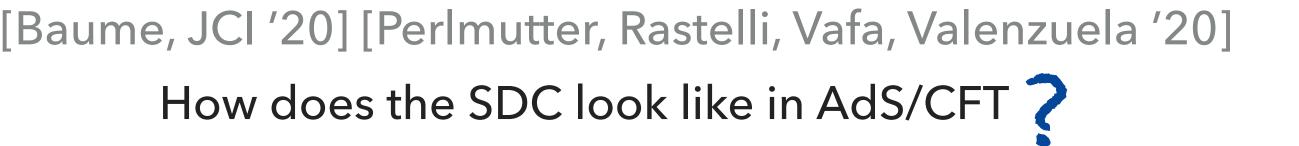
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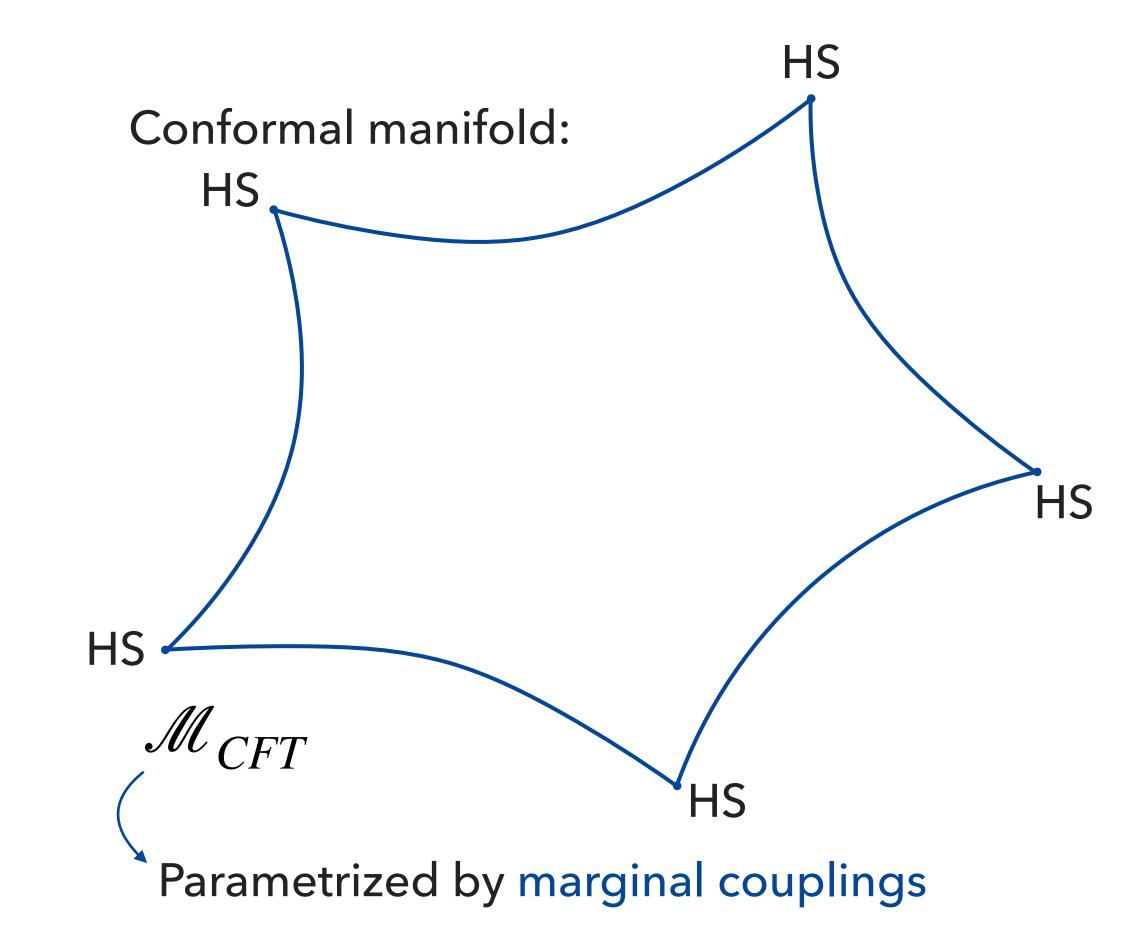
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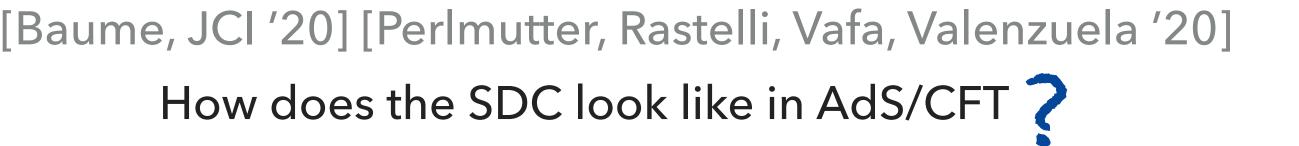


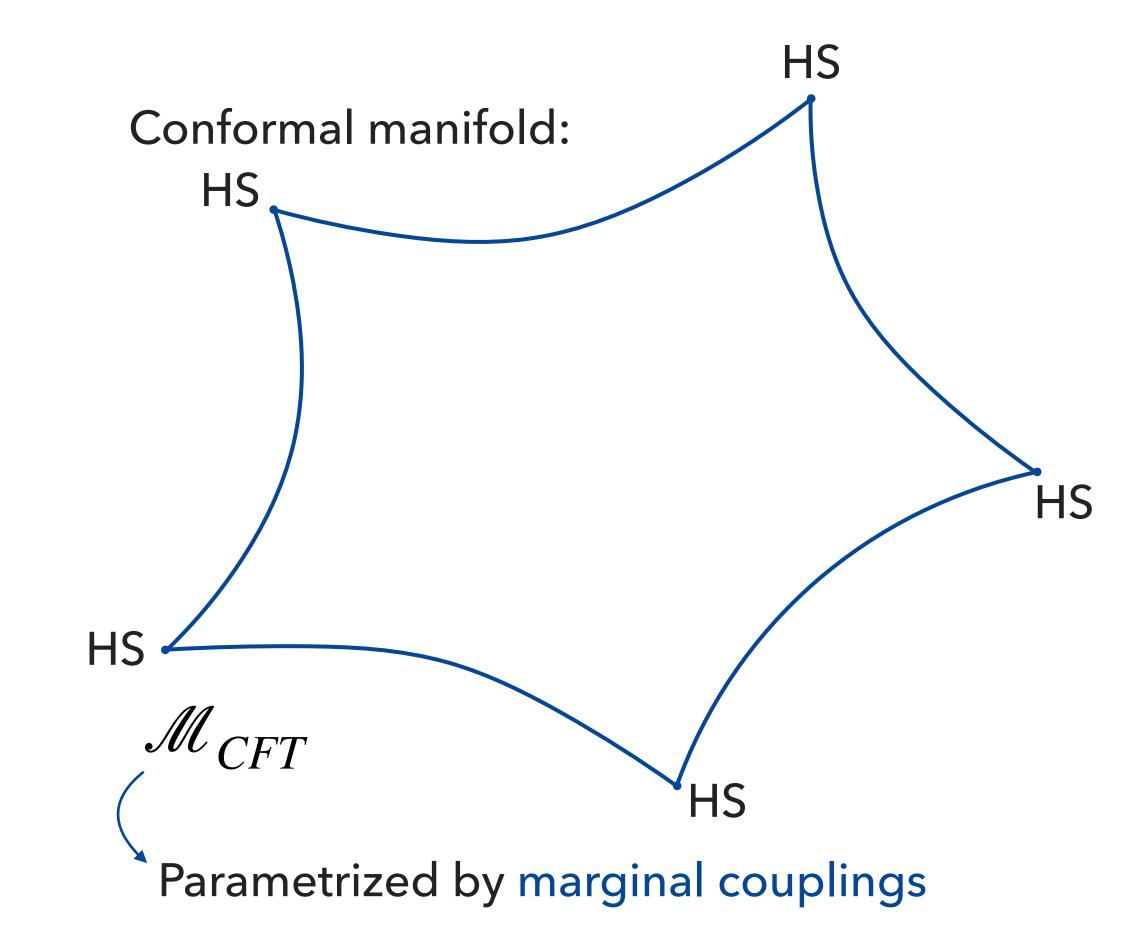


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Zamolodchikov distance

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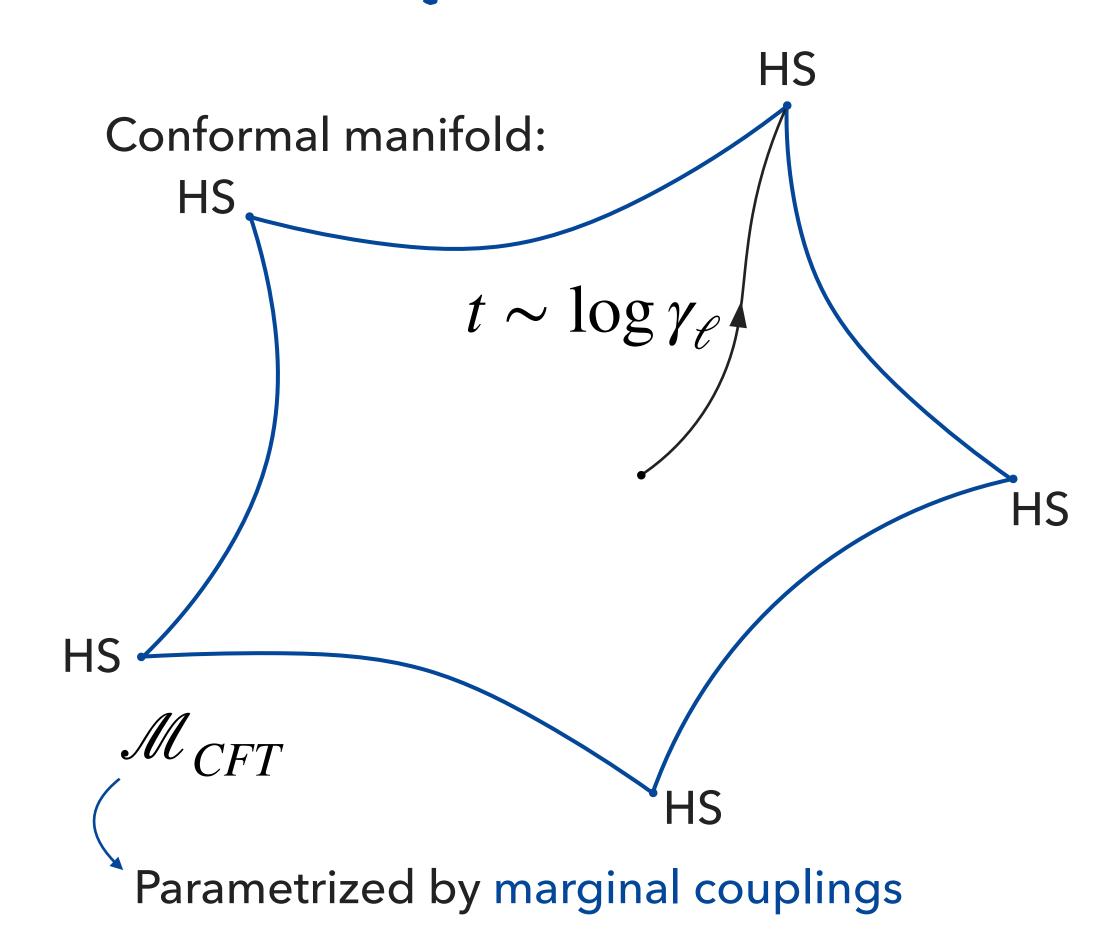


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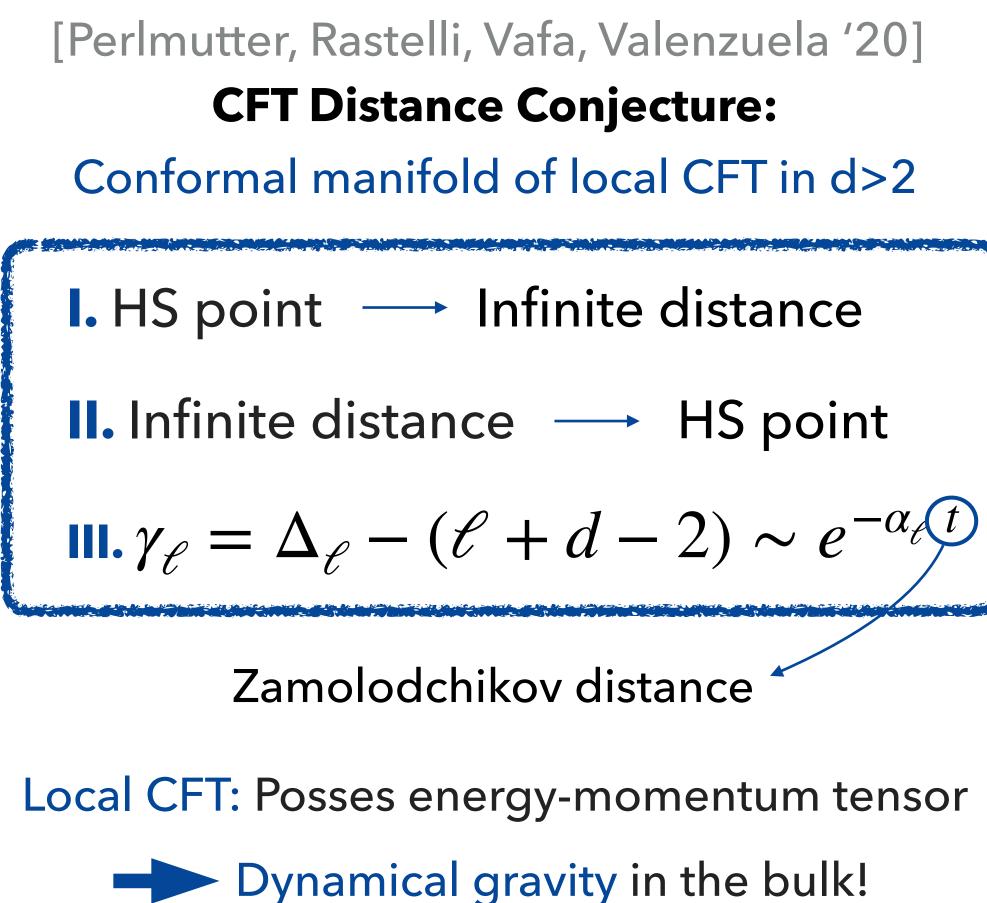
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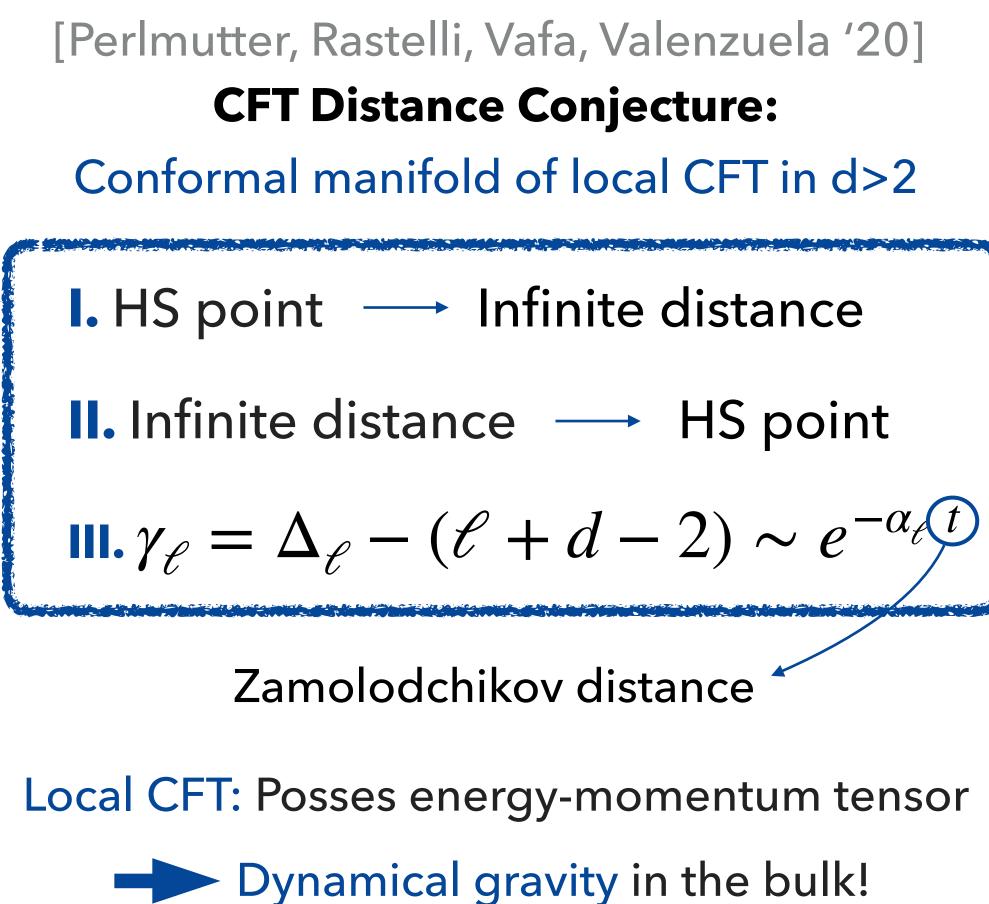


Can we prove this using CFT techniques **?**





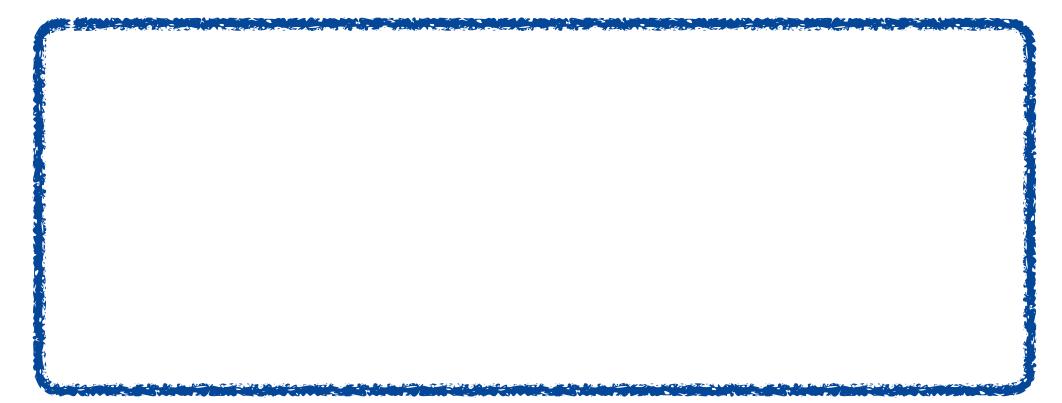
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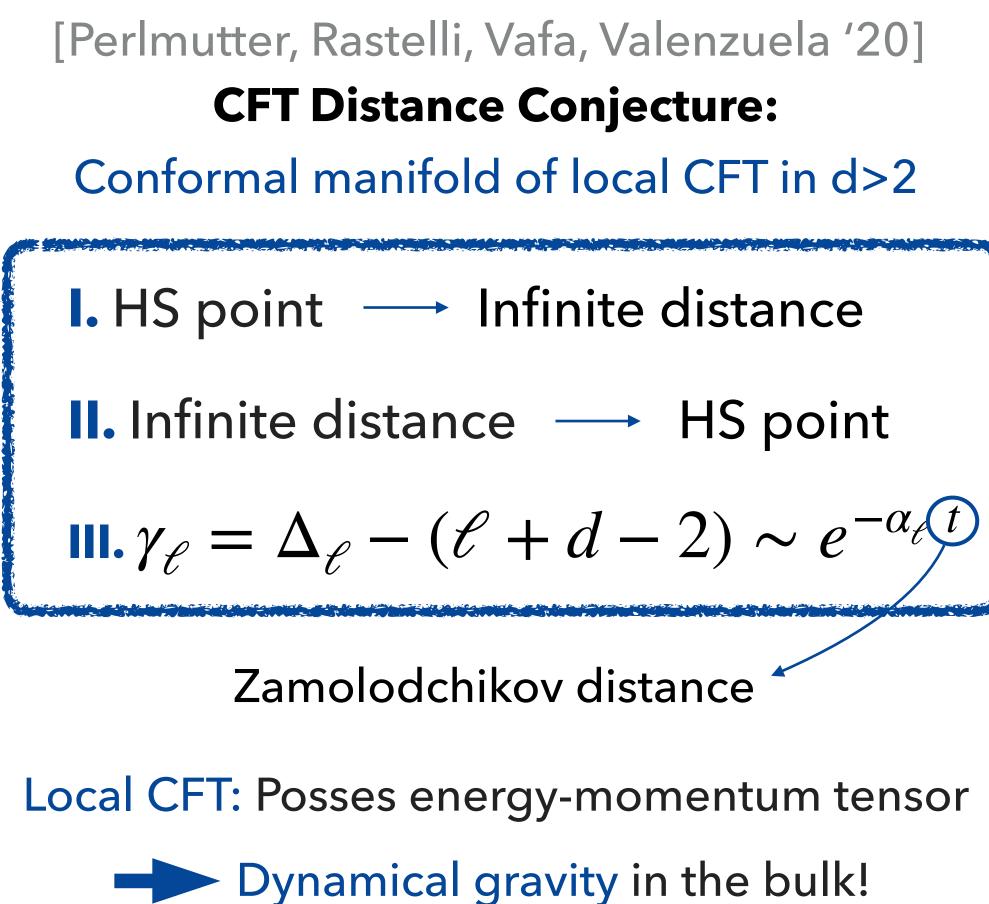
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Towards a CFT Distance Theorem:

Conformal perturbation theory + HS symmetry



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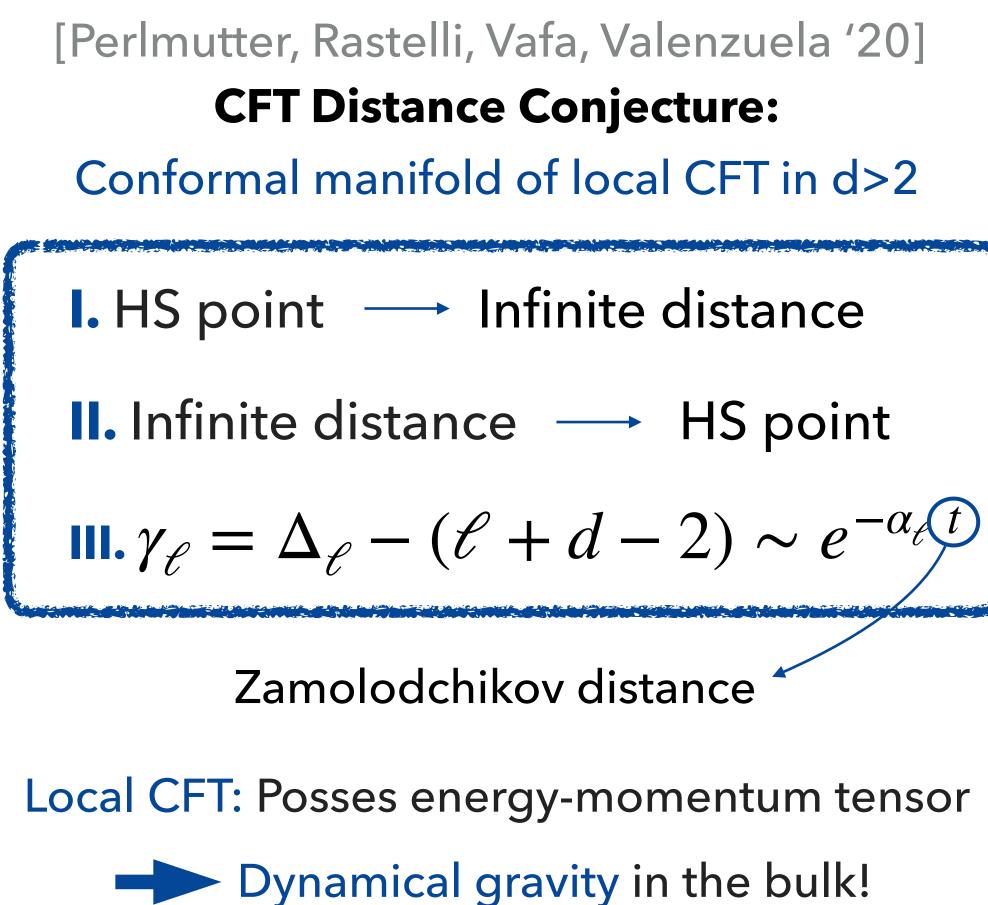
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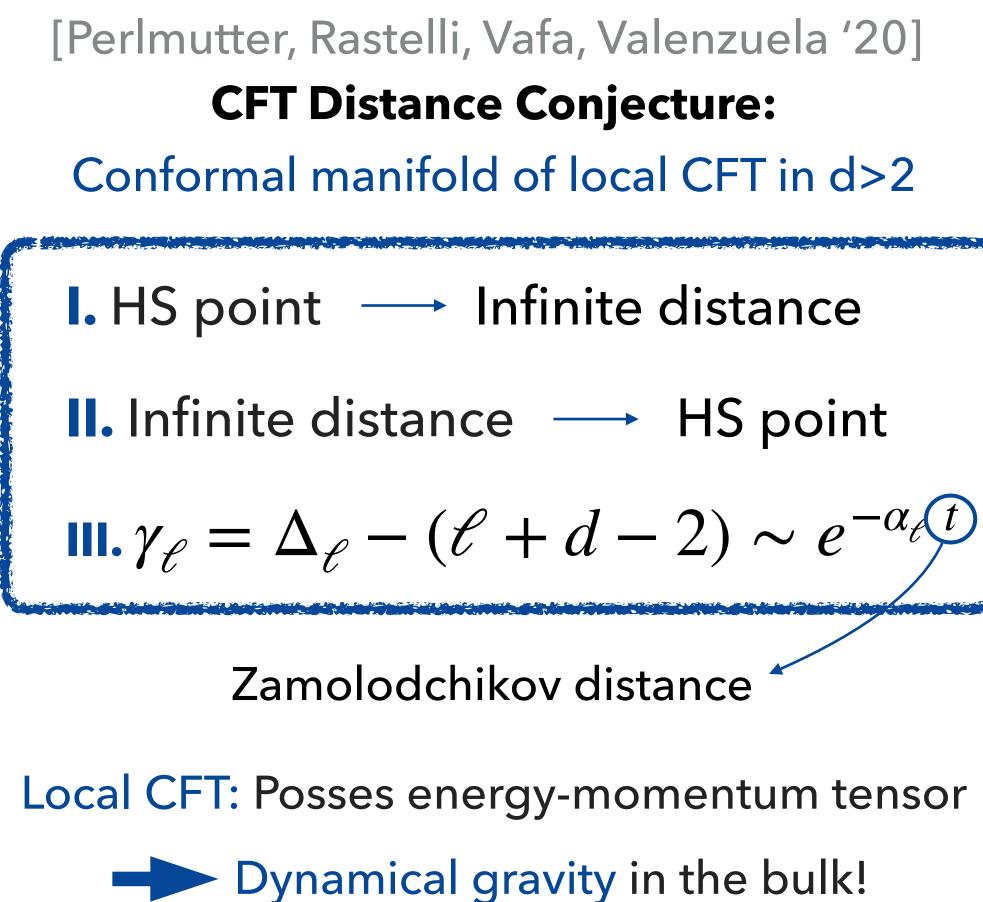
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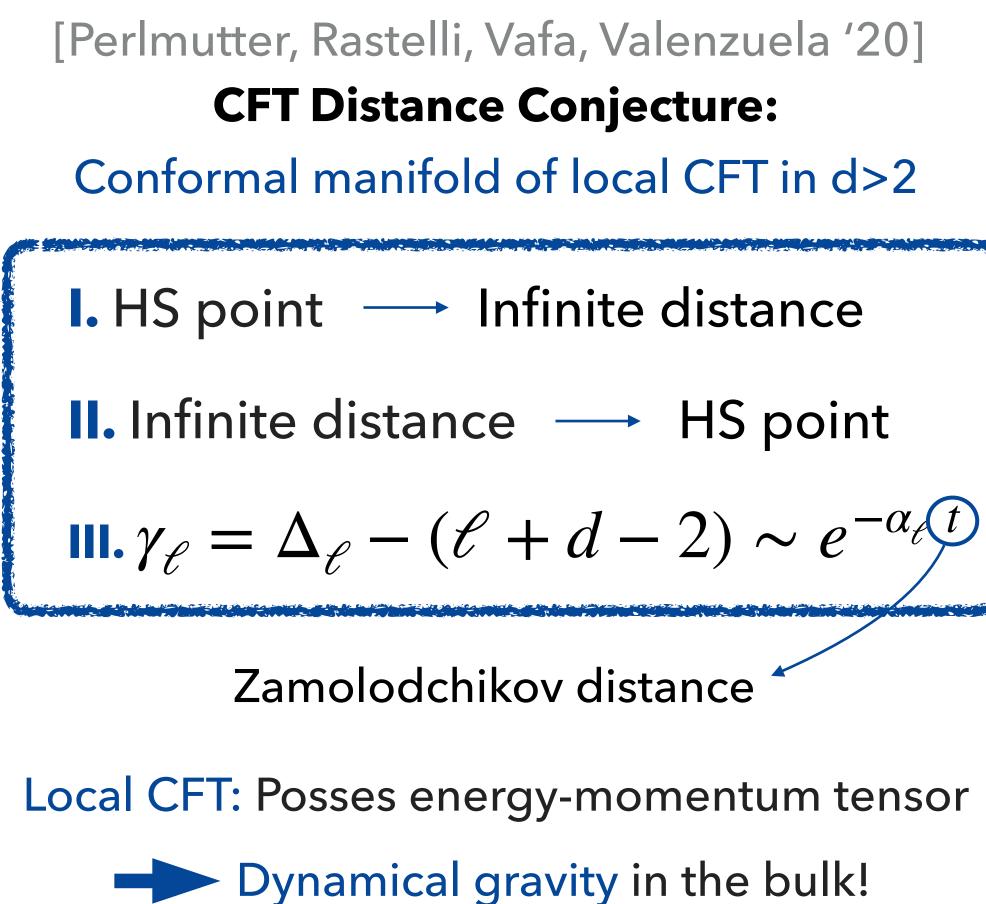
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$$\blacksquare \leftrightarrow \alpha_{\ell} \sim \left\langle K_{\ell-1} K_{\ell-1} \mathcal{O} \right\rangle_{HS} \neq 0$$

Evaluated at HS point! -

Can we prove this using CFT techniques **?**



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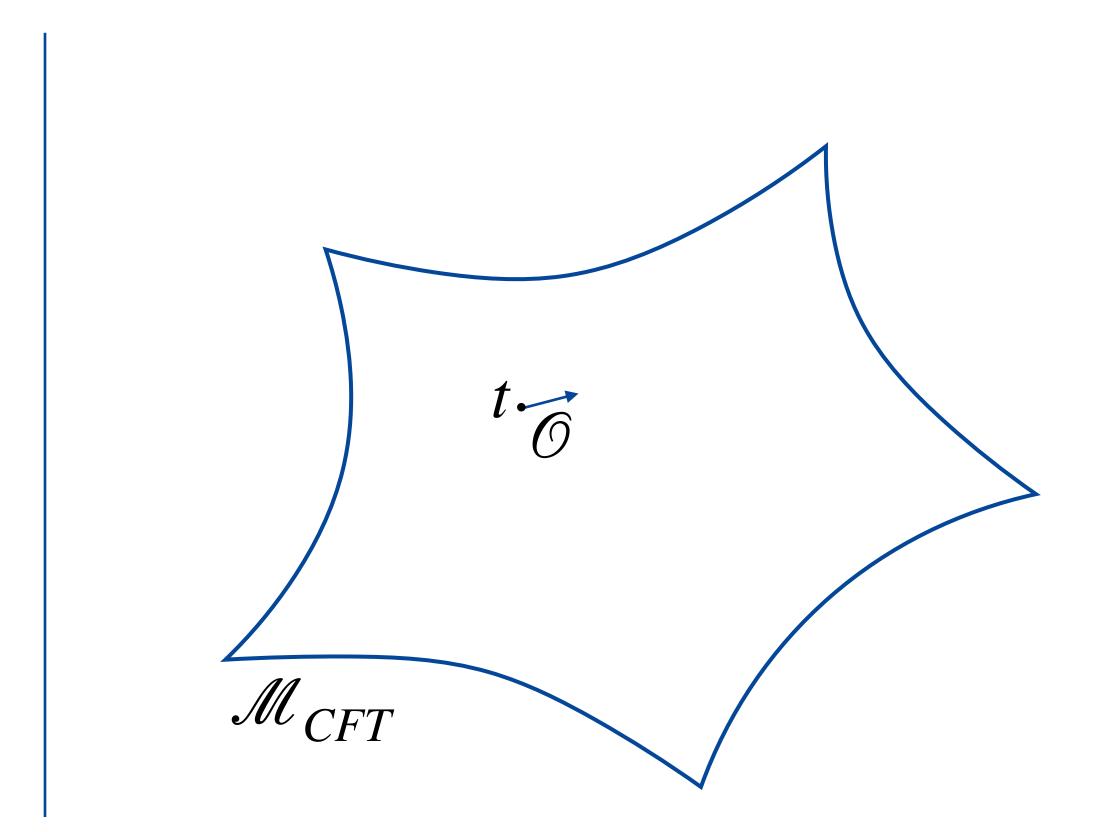


No extra assumption, e.g., no supersymmetry + existence of energy-momentum is crucial!



 $\delta \left\langle J_{\ell} J_{\ell} \right\rangle_{t} = \delta t \int \left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle_{t}$

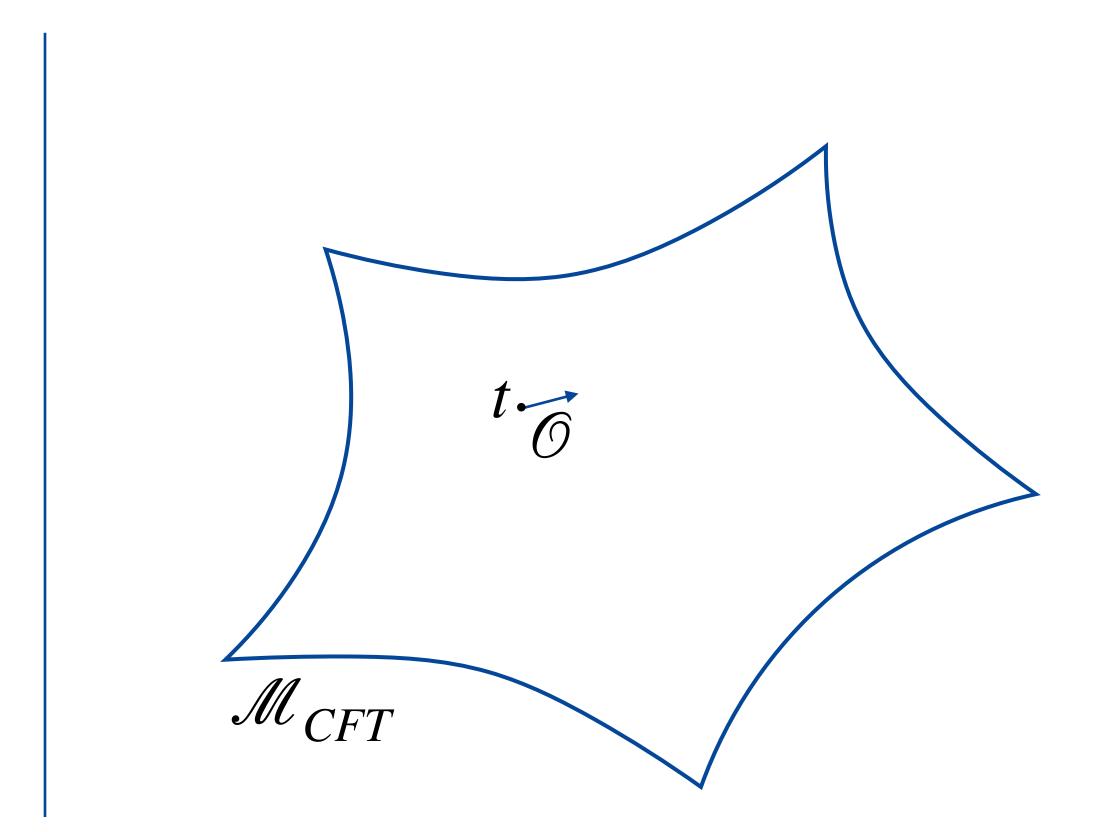
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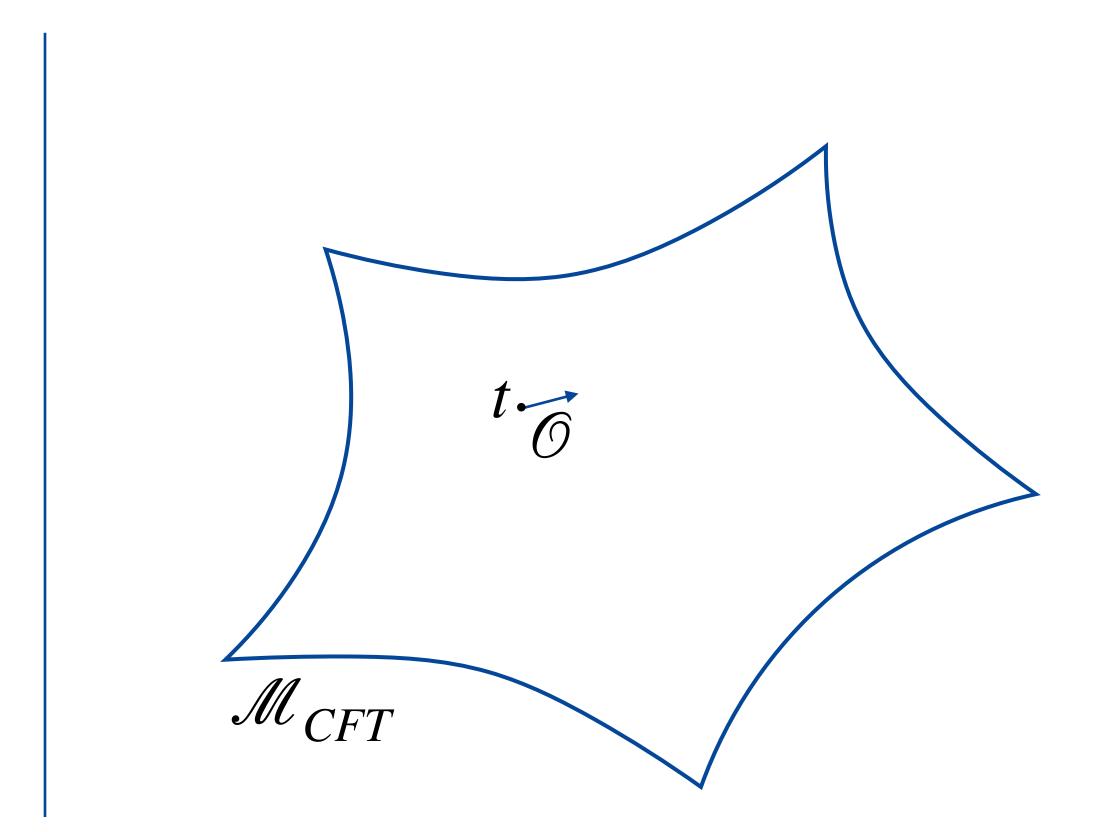
$$\delta \left\langle J_{\ell} J_{\ell} \right\rangle_{t} = \delta t \int \left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle_{t}$$

$$\downarrow$$

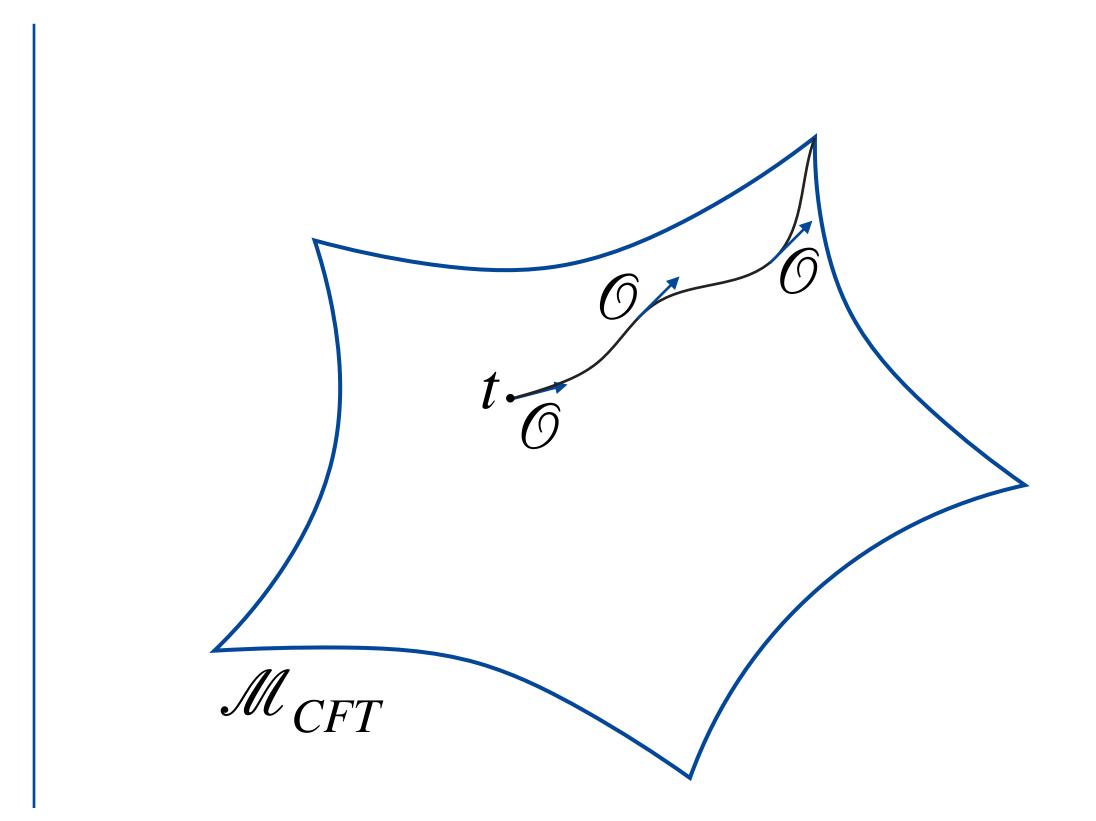
$$\delta \gamma_{\ell} = -C_{JJ\mathcal{O}}(t) \, \delta t$$



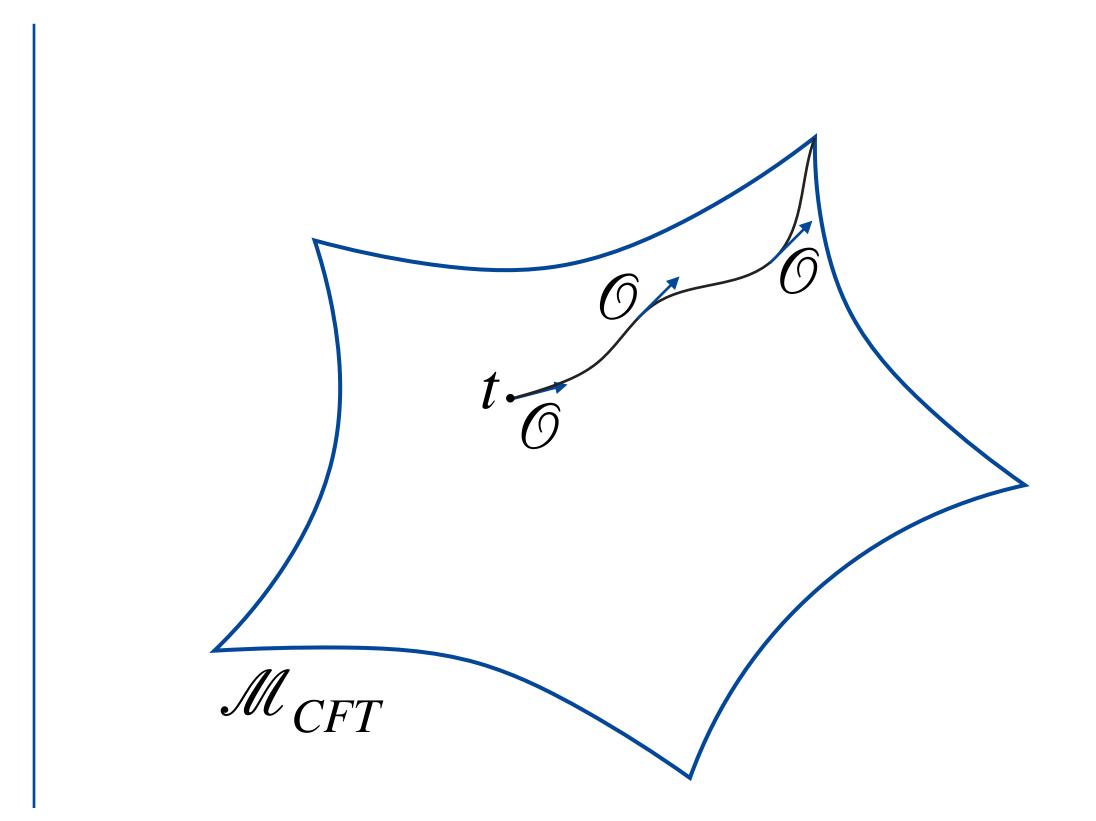
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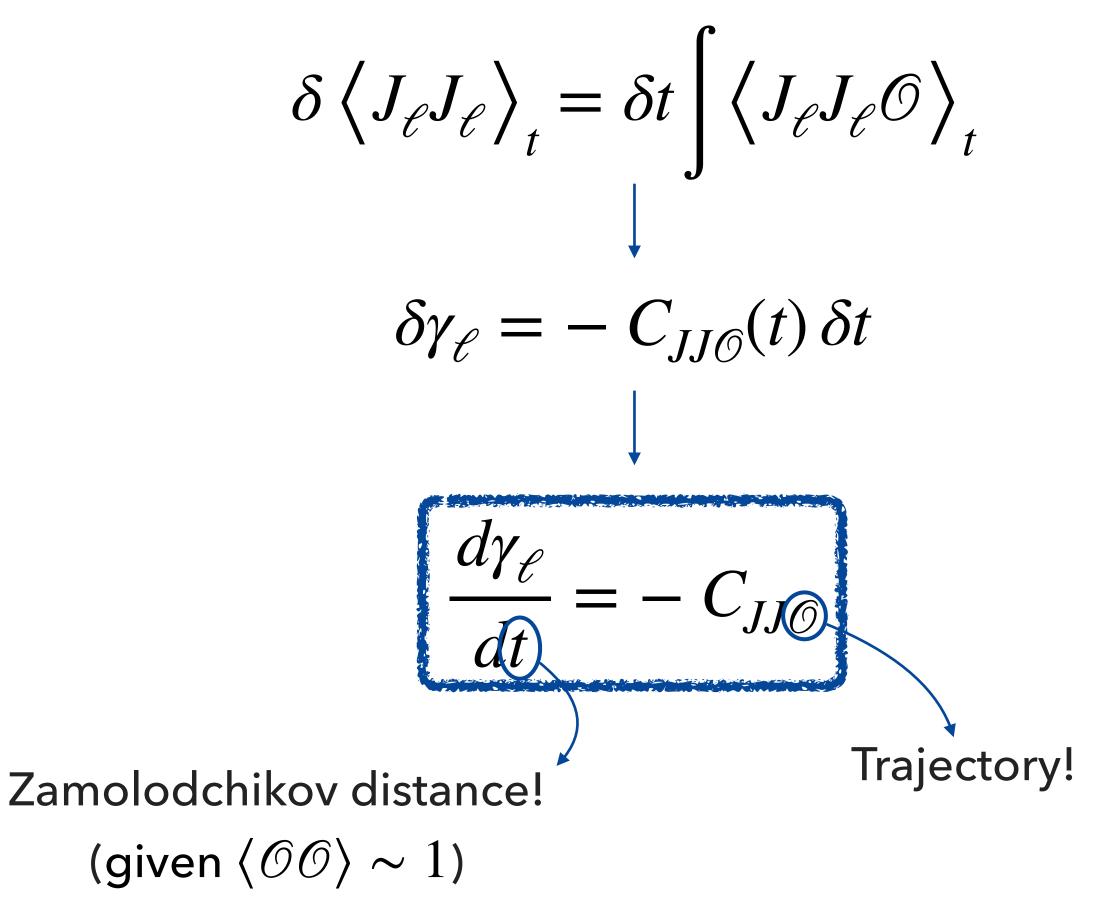
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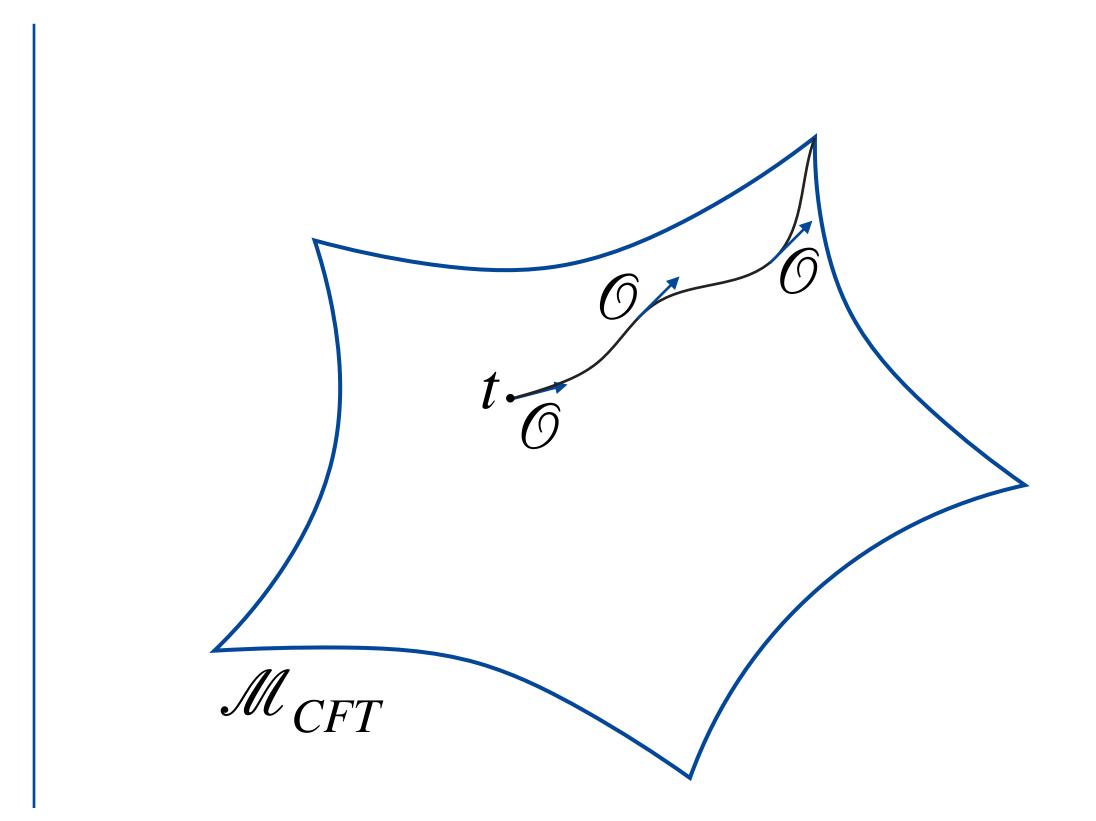
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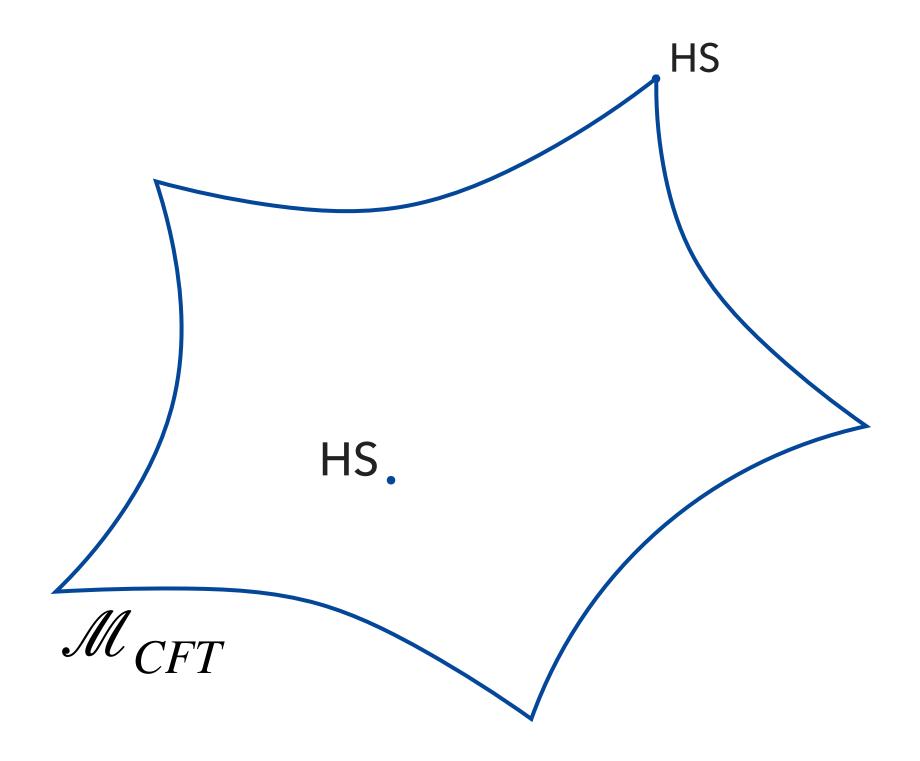
Conformal perturbation theory



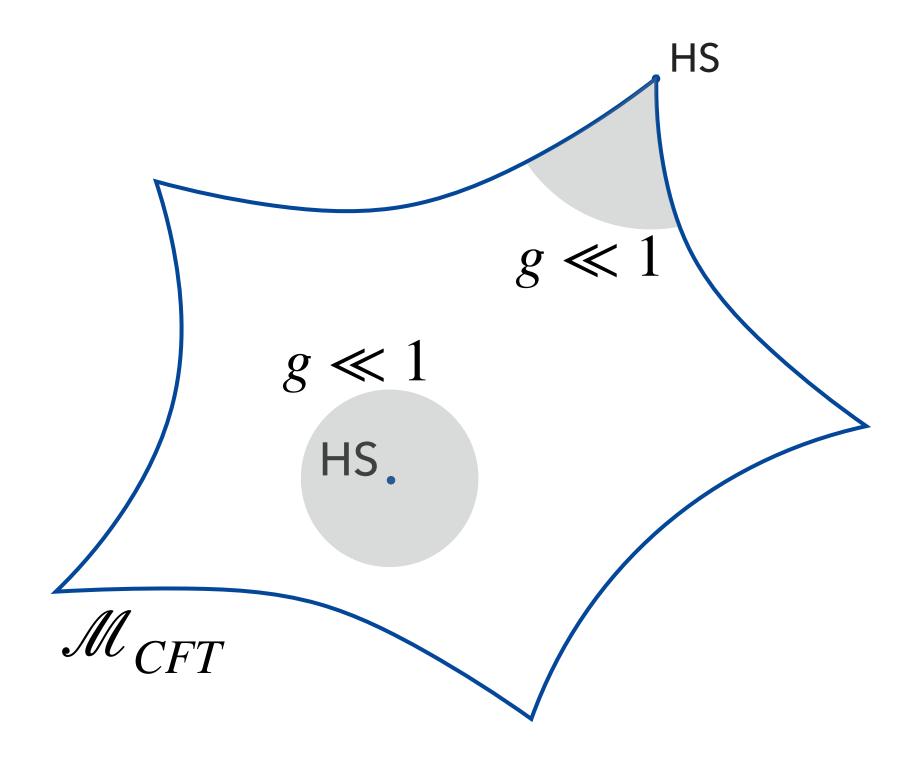
Sketch of the Proof



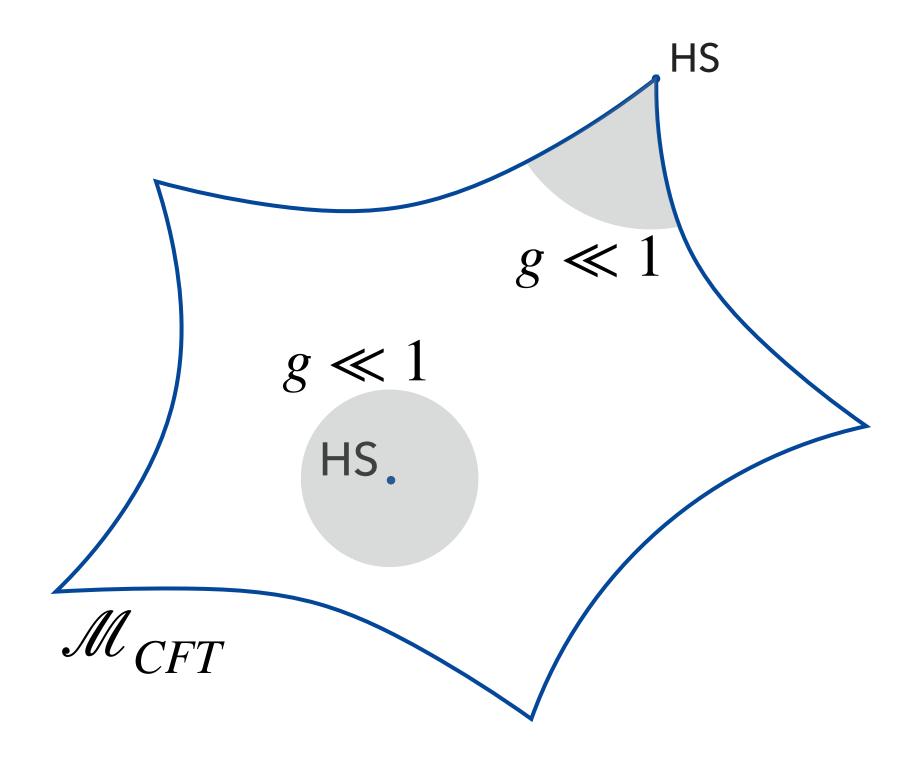
$$\partial \cdot J_{\ell} = g \, K_{\ell-1}$$



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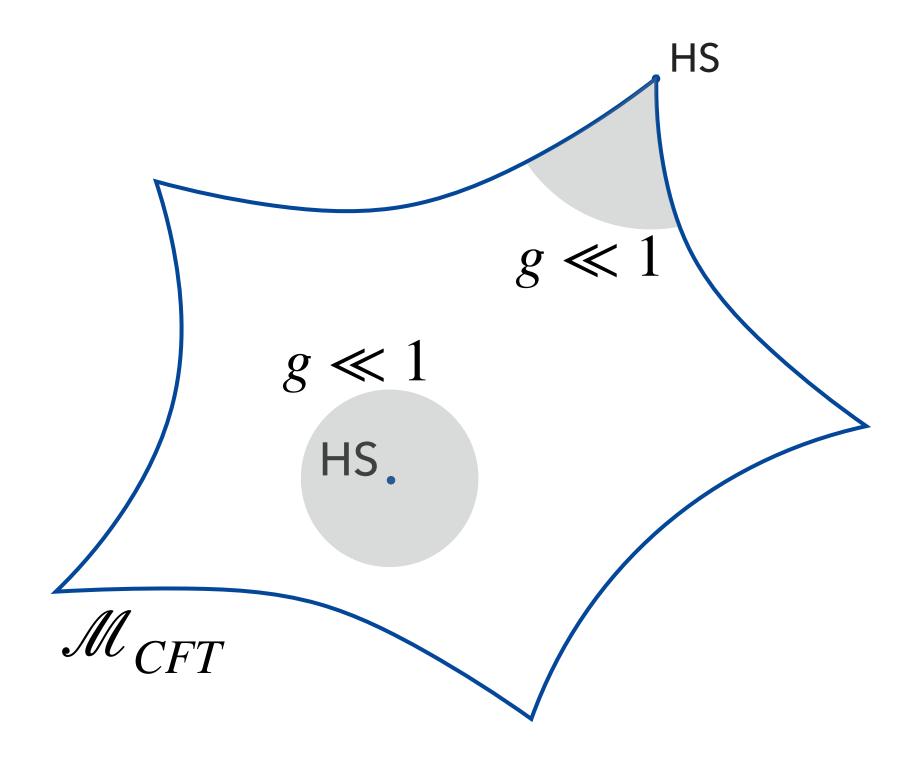
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(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = g \, K_{\ell-1}$$

 $C_{JJ\mathcal{O}} \simeq C_{JJ\mathcal{O}}^{HS} + C_{JK\mathcal{O}}^{HS} g + C_{KK\mathcal{O}}^{HS} g^2 + \cdots$



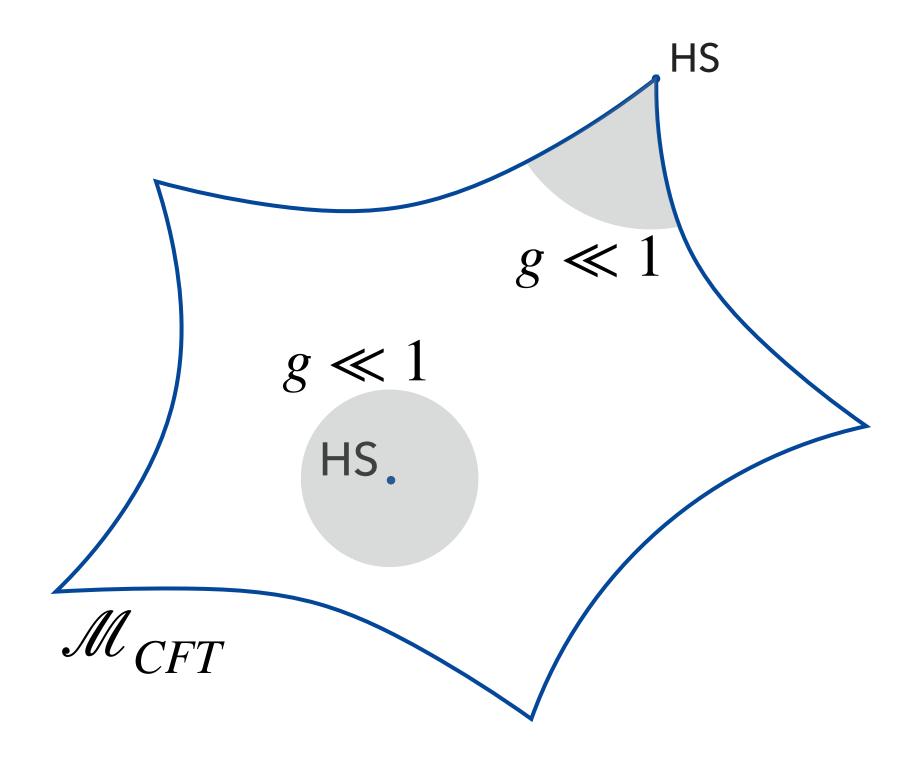
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 $C_{JJ\mathcal{O}} \simeq C_{JI\mathcal{O}}^{HS} + C_{JK\mathcal{O}}^{HS}g + C_{KK\mathcal{O}}^{HS}g^2 + \cdots$

Essentially, as $g \rightarrow 0$:

 $\partial \cdot \left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle = g \left\langle J_{\ell} K_{\ell-1} \mathcal{O} \right\rangle$



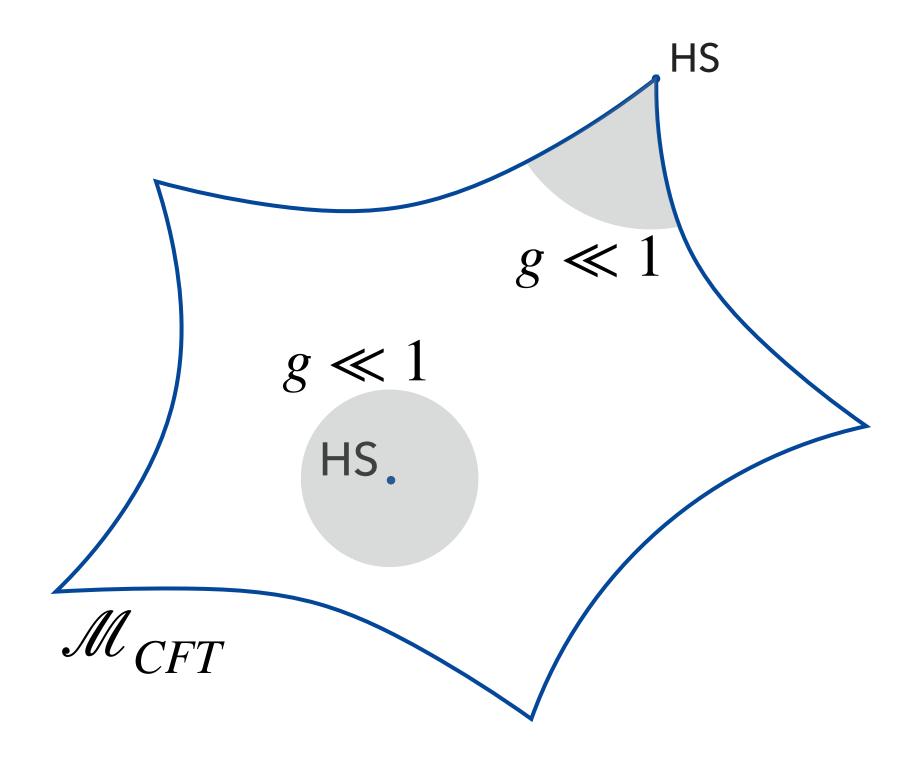
(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = g \, K_{\ell-1}$$

 $C_{IIO} \simeq C_{IIO}^{HS} + C_{IKO}^{HS}g + C_{KKO}^{HS}g^2 + \cdots$

Essentially, as $g \rightarrow 0$:

 $\partial \cdot \left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle = g \left\langle J_{\ell} K_{\ell-1} \mathcal{O} \right\rangle$ $C_{JJ\mathcal{O}} \quad C_{JK\mathcal{O}}$

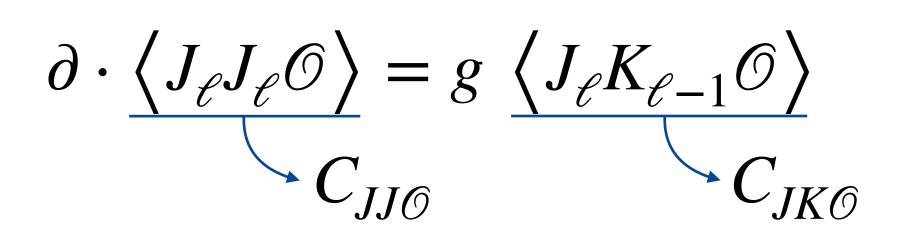


(Weakly-broken) HS symmetry

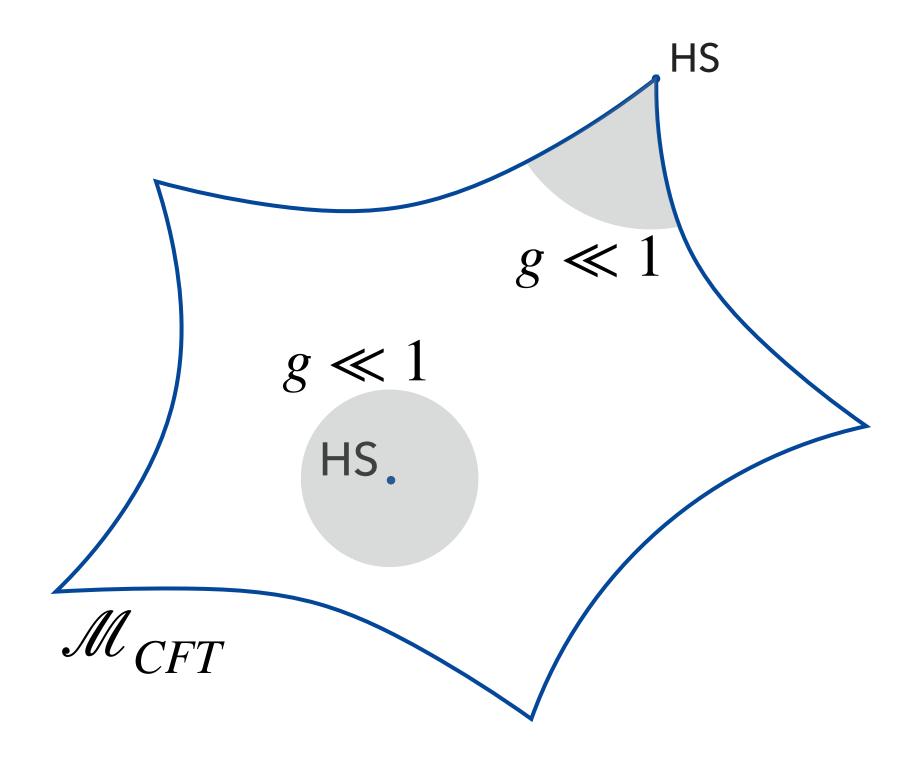
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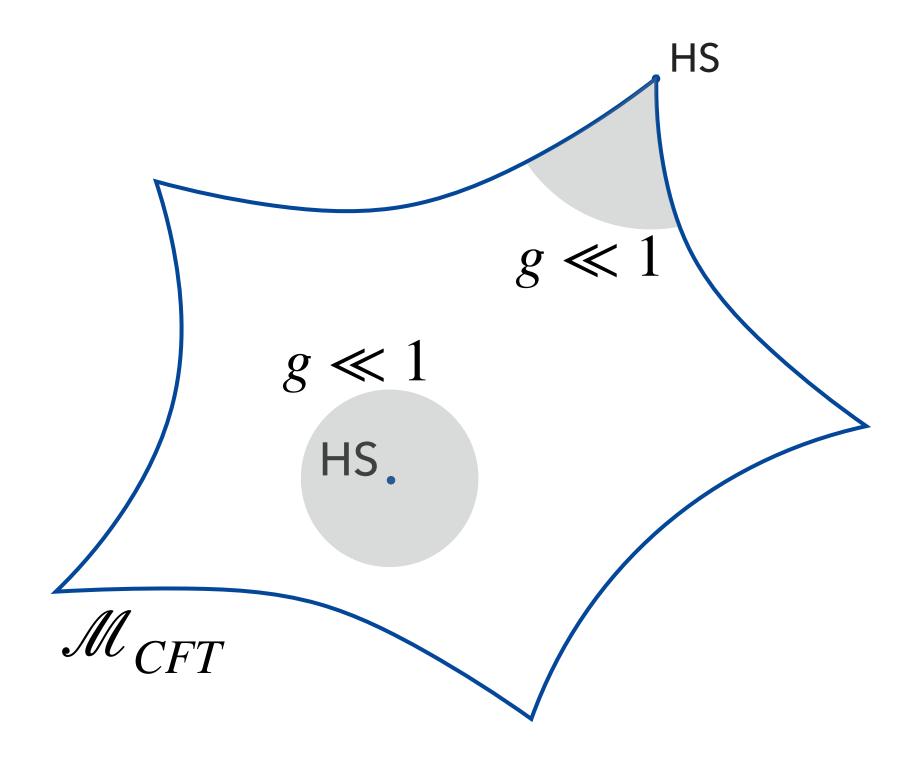


A bit more complicated than this... → More details in [Baume, JC '23]

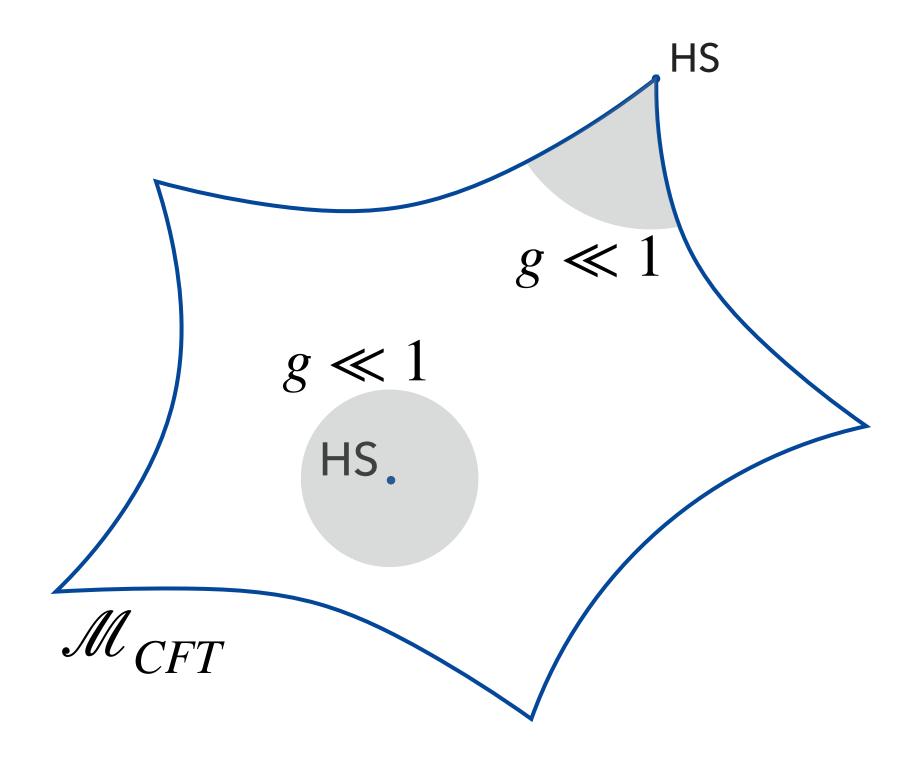


 $\partial \cdot J_{\ell} = g K_{\ell-1}$

 $C_{JJO} \simeq C_{JIO}^{HS} + C_{JKO}^{HS}g + C_{KKO}^{HS}g^2 + \cdots$



 $\partial \cdot J_{\ell} = g K_{\ell-1}$ $C_{JJO} \simeq C_{JJO}^{HS} + C_{JKO}^{HS} g + C_{KKO}^{HS} g^2 + \cdots$ HS symmetry constraints



(Weakly-broken) HS symmetry

 $\partial \cdot J_{\ell} = g K_{\ell-1}$

 $C_{JJ\mathcal{O}} \simeq C_{JJ\mathcal{O}}^{HS} + C_{JK\mathcal{O}}^{HS}g + C_{KK\mathcal{O}}^{HS}g^2 + \cdots$ HS symmetry constraints $\gamma_{\ell} \sim g^2$ $C_{JJO} \lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} \to 0$

Conformal perturbation theory

 $\delta \left\langle J_{\ell} J_{\ell} \right\rangle_{t} = \delta t \left[\left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle_{t} \right]$ $\delta \gamma_{\ell} = -C_{IIO}(t)\,\delta t$ $\frac{C_{II}}{L} = -C_{JJO}$

Sketch of the Proof

(Weakly-broken) HS symmetry

 $\partial \cdot J_{\ell} = g K_{\ell-1}$

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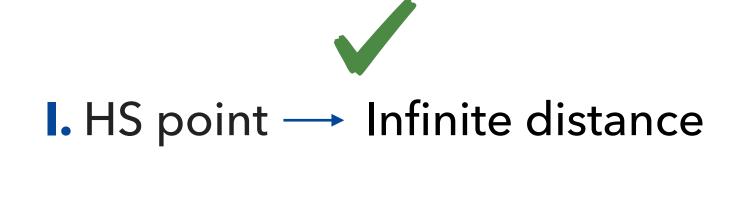


Sketch of the Proof

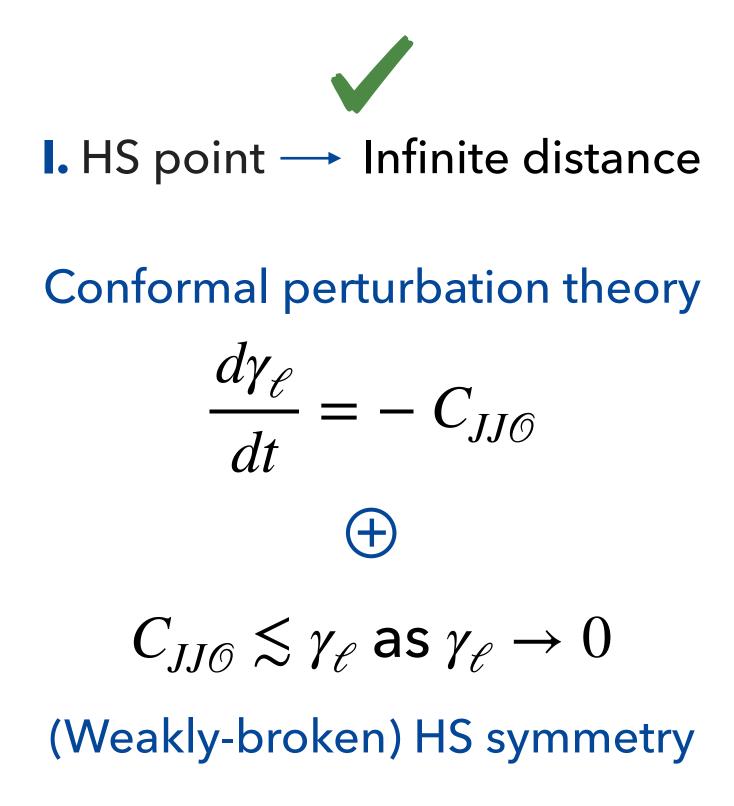
(Weakly-broken) HS symmetry $\partial \cdot J_{\ell} = g K_{\ell-1}$ $C_{JJO} \simeq C_{JJO}^{HS} + C_{JKO}^{HS}g + C_{KKO}^{HS}g^2 + \cdots$ HS symmetry constraints $\gamma_{\ell} \sim g^2$ $C_{JJO} \lesssim \gamma_\ell \, \mathrm{as} \, \gamma_\ell o 0$

CFT Distance Conjecture

CFT Distance Conjecture



CFT Distance Conjecture



CFT Distance Conjecture

II. Infinite distance → HS point

I. HS point → Infinite distance

Conformal perturbation theory

$$\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$$

$$C_{JJO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$$

(+)

CFT Distance Conjecture

I. HS point → Infinite distance

Conformal perturbation theory

$$\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$$

(+)

$$C_{JJO} \lesssim \gamma_\ell \text{ as } \gamma_\ell o 0$$

(Weakly-broken) HS symmetry

- II. Infinite distance → HS point
 - No HS symmetry



 $\exists \mathscr{O}: C_{JJ\mathscr{O}} \neq 0$

CFT Distance Conjecture

CFT Distance Sufficient but Criterion

I. HS point → Infinite distance Conformal perturbation theory $\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$ (+) $C_{IIO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$ (Weakly-broken) HS symmetry

- II. Infinite distance → HS point
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 - ?
 - $\exists \mathcal{O}: C_{JJ\mathcal{O}} \neq 0$
 - not necessary
 - Finite distance

CFT Distance Conjecture

Criterion

I. HS point → Infinite distance Conformal perturbation theory $\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$ (+) $C_{IIO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$ (Weakly-broken) HS symmetry

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- **CFT Distance** Sufficient but not necessary
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 $\blacksquare \cdot \gamma_{\ell} \sim e^{-\alpha_{\ell} t}$

CFT Distance Conjecture

Criterion

I. HS point → Infinite distance Conformal perturbation theory $\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$ (+) $C_{IIO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$ (Weakly-broken) HS symmetry

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$$\blacksquare \gamma_{\ell} \sim e^{-\alpha_{\ell} t}$$

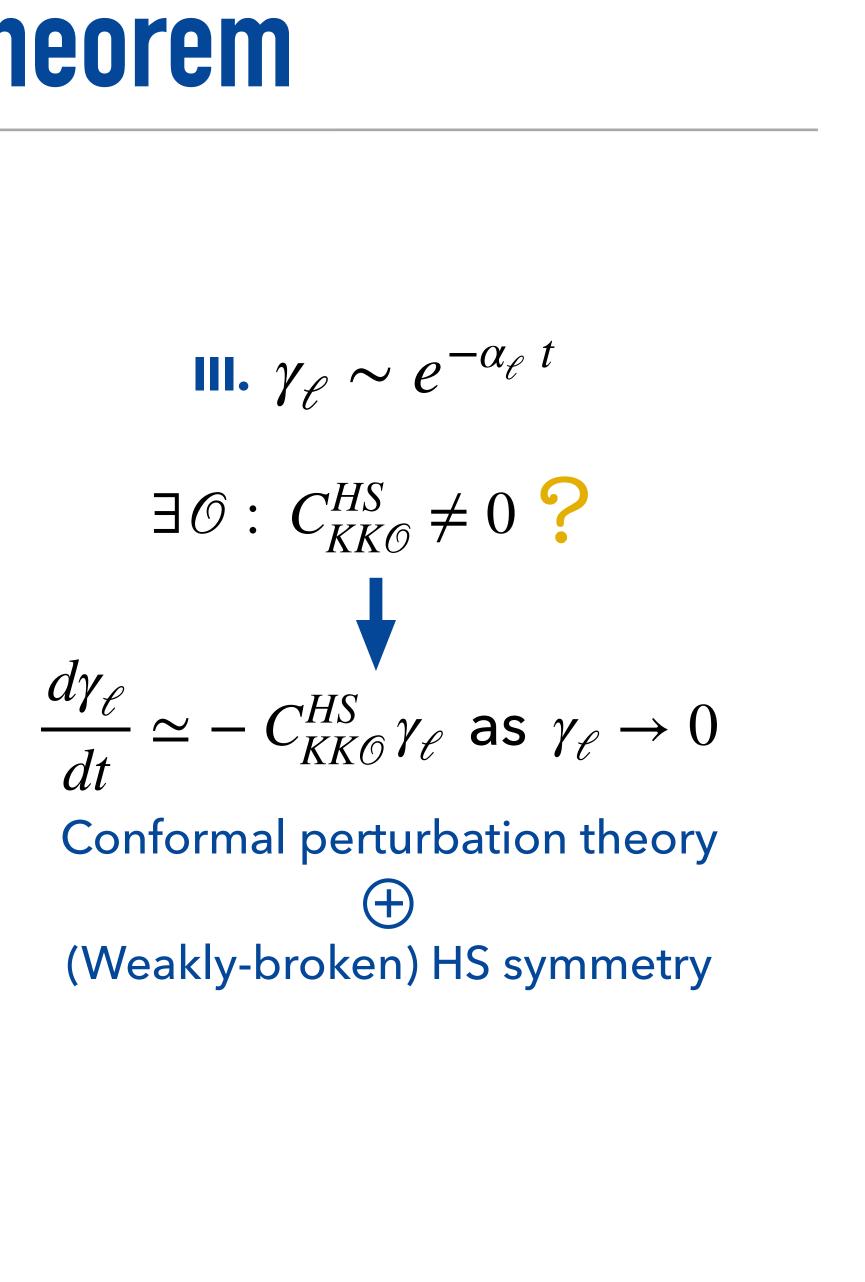
 $\exists \mathscr{O}: C_{KK\mathscr{O}}^{HS} \neq 0 ?$

CFT Distance Conjecture

CFT Distance Sufficient but Criterion

I. HS point \rightarrow Infinite distance **Conformal perturbation theory** $\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$ (+) $C_{IIO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$ (Weakly-broken) HS symmetry

- II. Infinite distance → HS point
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CFT Distance Conjecture

I. HS point → Infinite distance

Conformal perturbation theory

 $\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$

(+)

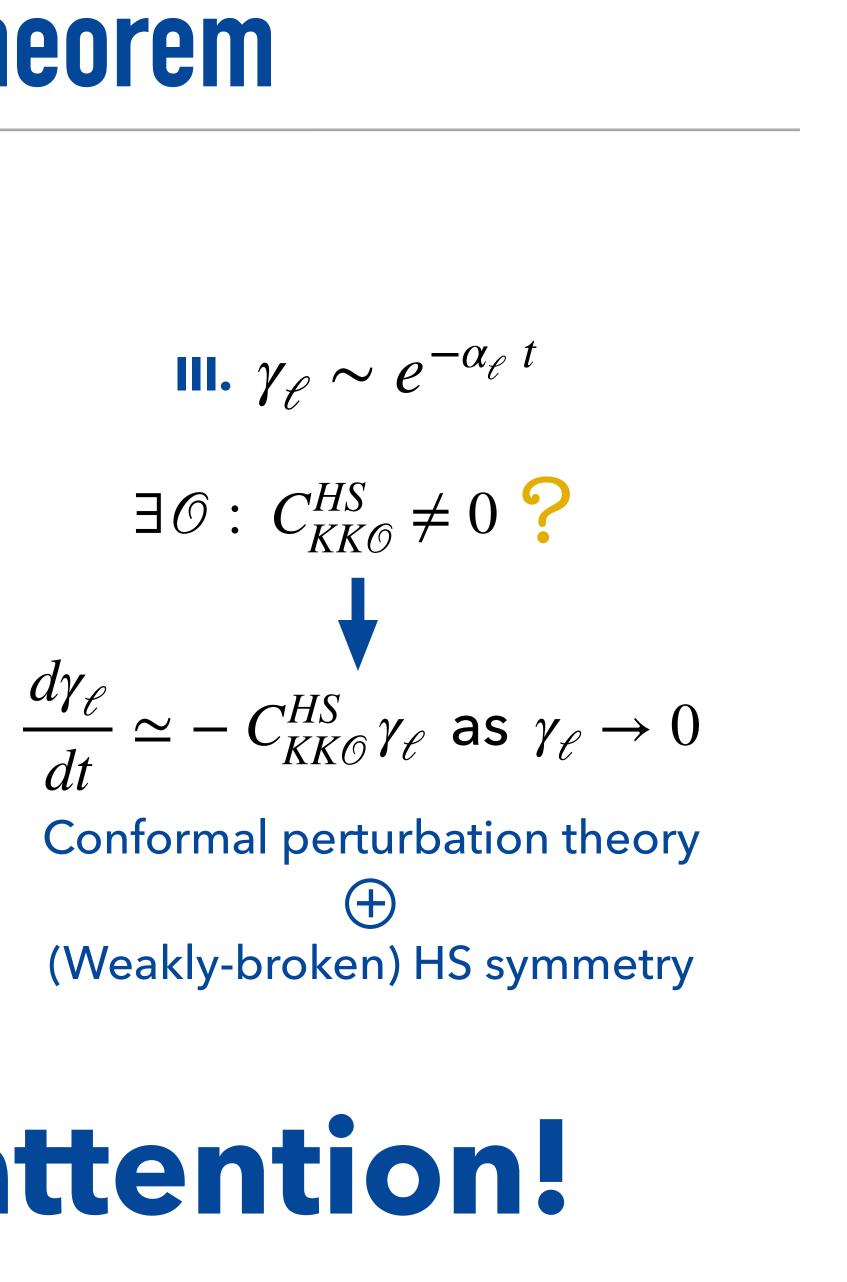
 $C_{IIO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$

(Weakly-broken) HS symmetry

Criterion

Thank you for your attention!

- II. Infinite distance → HS point
 - No HS symmetry
 - $\exists \mathscr{O}: C_{JJ\mathscr{O}} \neq 0$
- **CFT Distance** Sufficient but not necessary
 - Finite distance



Backup slides



Conformal Perturbation Theory for HS operators

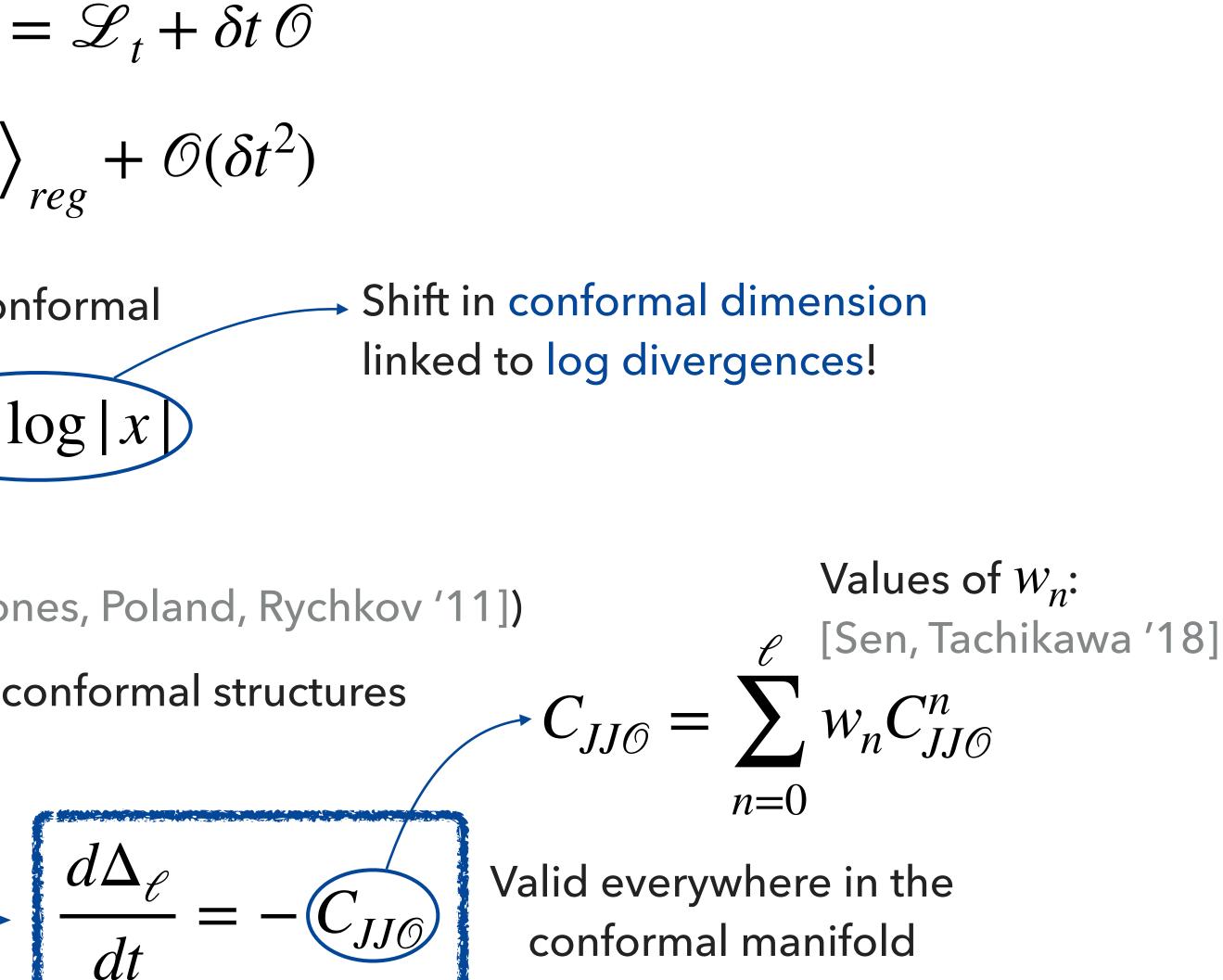
Deform the theory with an operator: $\mathscr{L}_{t+\delta t} = \mathscr{L}_t + \delta t \mathcal{O}$

$$\rightarrow \delta \left\langle J_{\ell} J_{\ell} \right\rangle = \delta t \int d^{d} y \left\langle J_{\ell} J_{\ell} \mathcal{O}(y) \right\rangle_{\mu}$$

• \mathcal{O} marginal operator \longrightarrow Theory remains conformal $\rightarrow \delta \left\langle J_{\ell} J_{\ell} \right\rangle = -\frac{H_{12}^{\ell}}{|x|^{2\Delta_{\ell} - 2\ell}} 2\delta \Delta_{\ell} \log |x|$

Conformal structures (see e.g. [Costa, Penedones, Poland, Rychkov '11])

$$\searrow \left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle = \sum_{n=0}^{\ell} C_{JJ\mathcal{O}}^{n} \Theta_{n}^{\ell} \mathcal{C}^{\ell+1} \cos \theta_{n}^{\ell} \right\rangle$$
$$\delta \Delta_{\ell} = -\delta t C_{JJ\mathcal{O}} + \mathcal{O}(\delta t^{2}) \longrightarrow 0$$





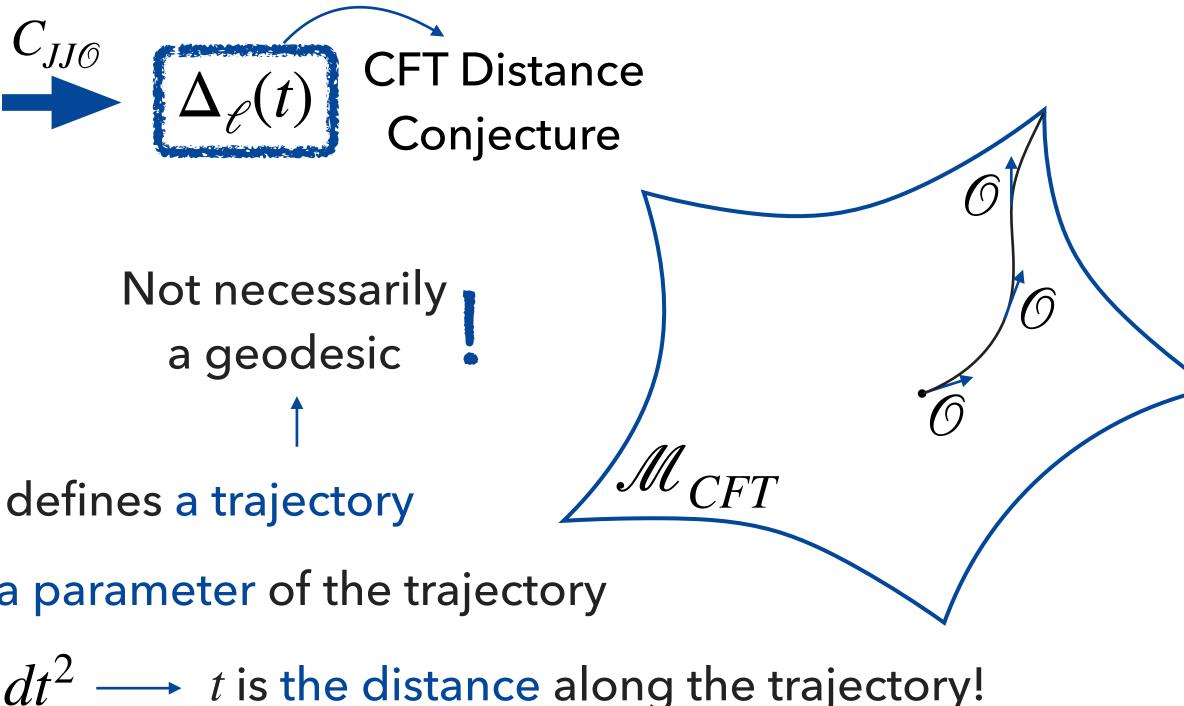
Distances from Conformal Perturbation Theory

$$\frac{d\Delta_{\ell}}{dt} = -C_{JJO}$$

Two important questions:

• What is the meaning of t? **1.** \mathcal{O} defines a direction in $\mathcal{M}_{CFT} \longrightarrow \mathcal{O}(t)$ defines a trajectory **2.** $t + \delta t$ represents how we move $\longrightarrow t$ is a parameter of the trajectory **3.** Take $\langle OO \rangle \sim 1 \longrightarrow ds^2 \sim \langle OO \rangle dt^2 \sim dt^2 \longrightarrow t$ is the distance along the trajectory! • How do we learn about C_{JJO} ? Usually: CPT for $\langle JJ O \rangle$ -

Here: Something different! In particular, weakly-broken HS symmetry



Complicated equations involving all spectrum of operators and conformal blocks

(See e.g. [Behan '18])



A CFT Distance Criterion

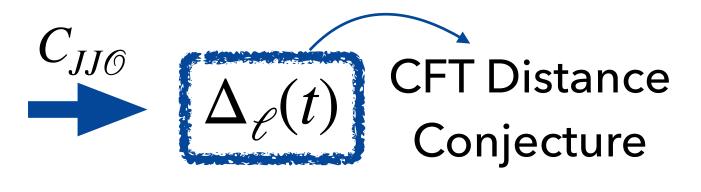
$$\frac{d\Delta_{\ell}}{dt} = -C_{JJO}$$

Consider a point with $\Delta_{\ell} = \Delta_{\ell}^*$ and $\Delta_{\ell}(t)$ inver

Take parametrization:
$$C_{JJO} \propto \left(\Delta_{\ell} - \Delta_{\ell}^{*}\right)^{a+1}$$
 $a \in \mathbb{R}$
 $\rightarrow t \propto \left(\Delta_{\ell} - \Delta_{\ell}^{*}\right)^{-a}$ Finite distance $\rightarrow a < 0$
Infinite distance $\rightarrow a \ge 0$

CFT Distance criterion:

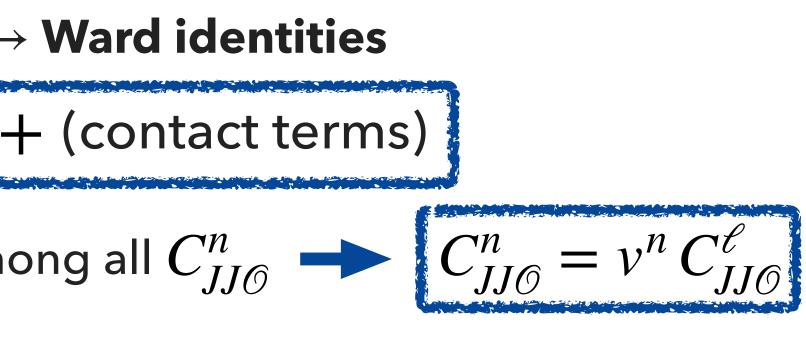
Finite distance $\longleftrightarrow \exists \mathcal{O}$ (marginal operator) such that a < 0Infinite distance $\rightarrow a \ge 0 \quad \forall \mathcal{O} \ (marginal operator)$ • Notice! Any point with $C_{IIO} \neq 0$ any automatically at finite distance!



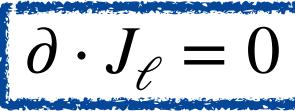
rtible in a neighbourhood:
$$C_{JJO}(t) \longrightarrow C_{JJO}(t)$$

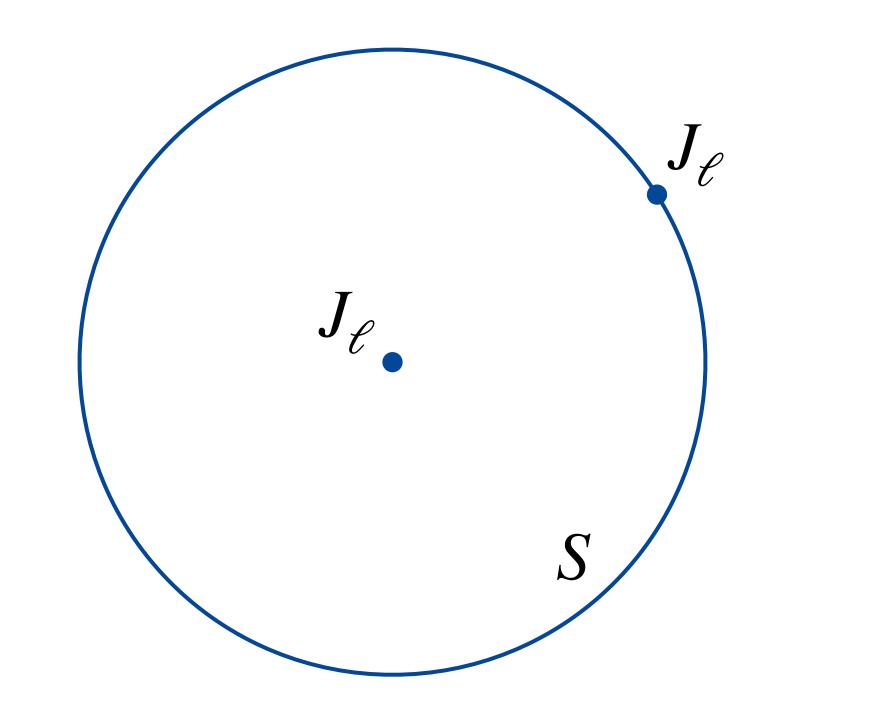
$$C_{JJO}(\Delta_{\ell})$$

HS point
$$\rightarrow$$
 Ward identities
 $\partial \cdot J_{\ell} = 0 + (\text{contact terms})$
• $\langle \partial J_{\ell} J_{\ell} \mathcal{O} \rangle = 0 \rightarrow \text{Recursive relation among all } C_{JJ\mathcal{O}}^n \rightarrow C_{JJ\mathcal{O}}^n = v^n$
 $\rightarrow C_{JJ\mathcal{O}}^{HS} = C_{JJ\mathcal{O}}^{\ell} \sum_{n=0}^{\ell} w_n v^n = 0$? No :(...Need to work a bit
• Integrated Ward identity! Twist conservation: $\Delta_s - s = d - 2$
1. Define conserved charges: $Q_{\ell} = \int_S J_{\ell} \rightarrow [Q_{\ell}, J_{\ell}] \sim \sum_s J_s$
2. Use: $\int_S \langle J_{\ell} J_{\ell} \mathcal{O} \rangle \propto \langle [Q_{\ell}, J_{\ell}] \mathcal{O} \rangle$



it harder



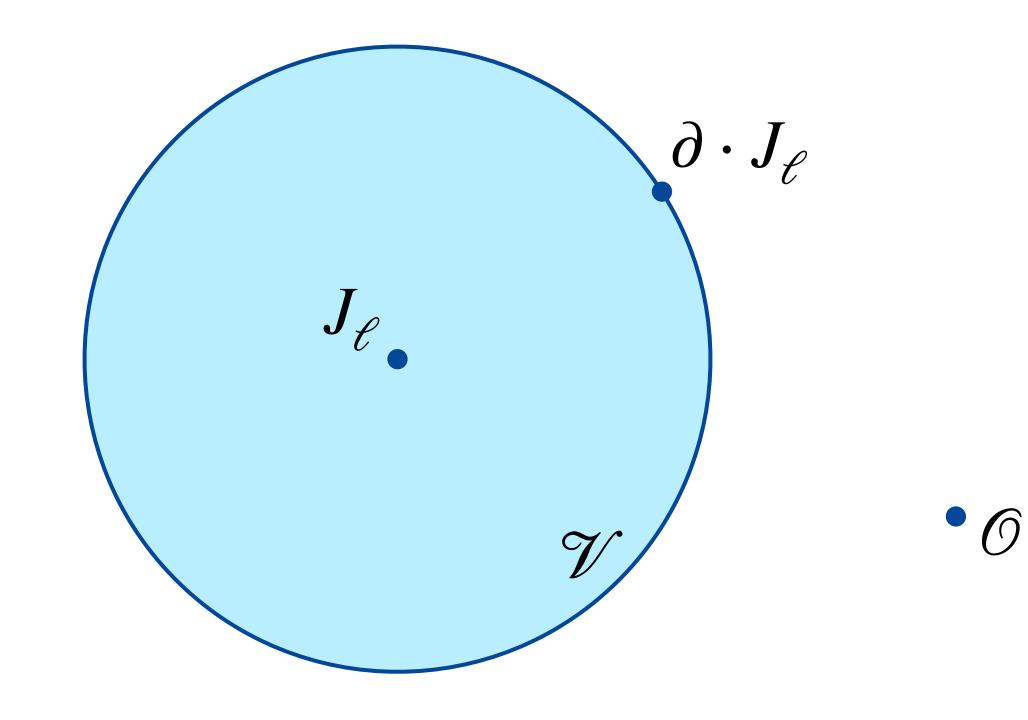


HS point \rightarrow **Ward identities**

 $\partial \cdot J_{\ell} = 0 + (\text{contact terms})$

Surface integral



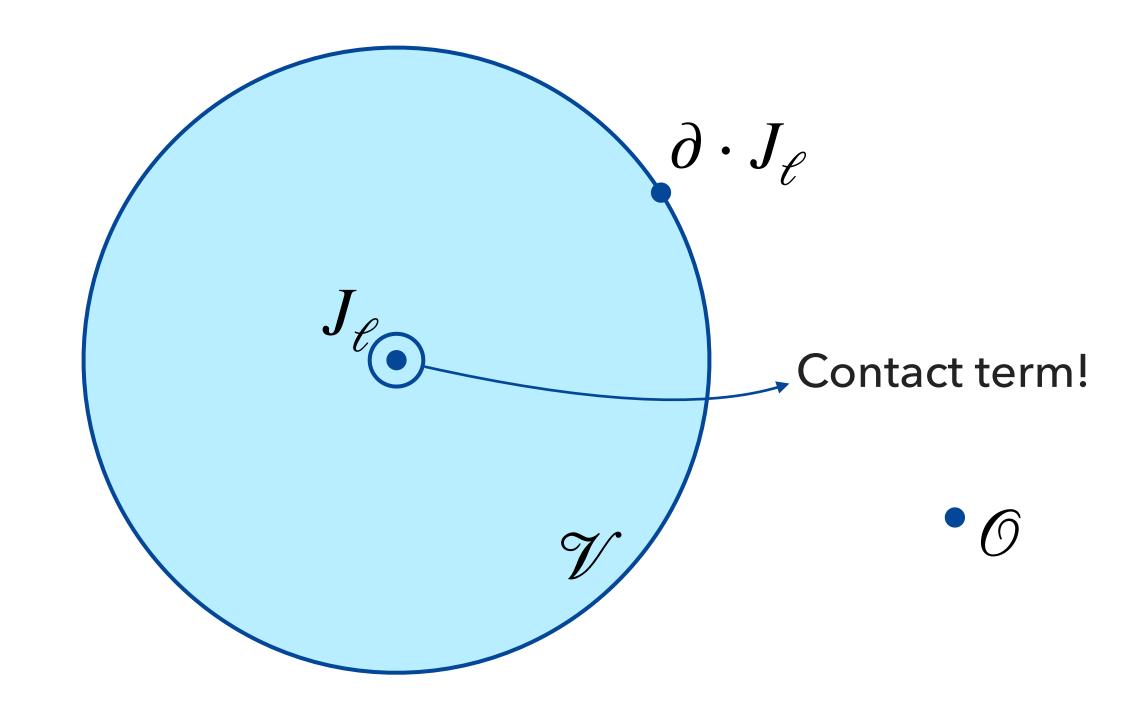


HS point \rightarrow **Ward identities**

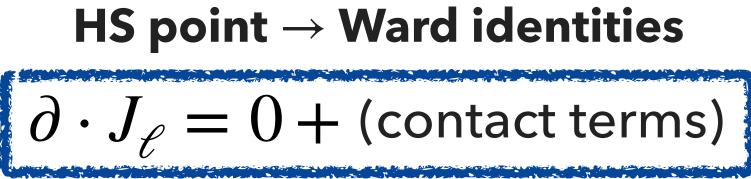
 $\partial \cdot J_{\ell} = 0 + \text{(contact terms)}$

Surface integral + Stokes theorem

HS point \rightarrow **Ward identities** $\partial \cdot J_{\ell} = 0 + \text{(contact terms)}$



Surface integral + Stokes theorem + Ward identity

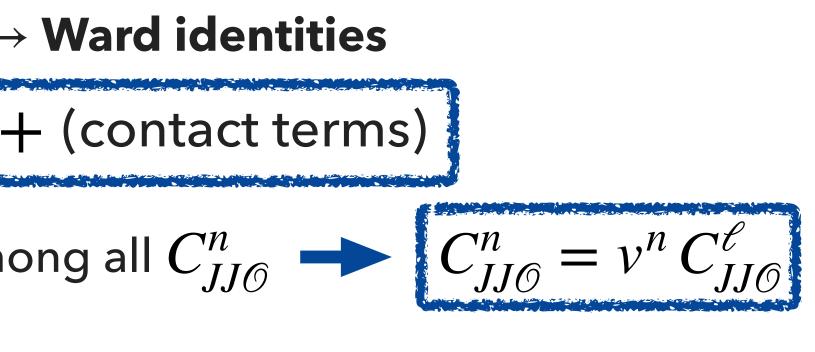


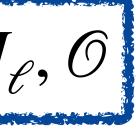
 $[Q_{\ell}, J_{\ell}]$

Surface integral + Stokes theorem + Ward identity

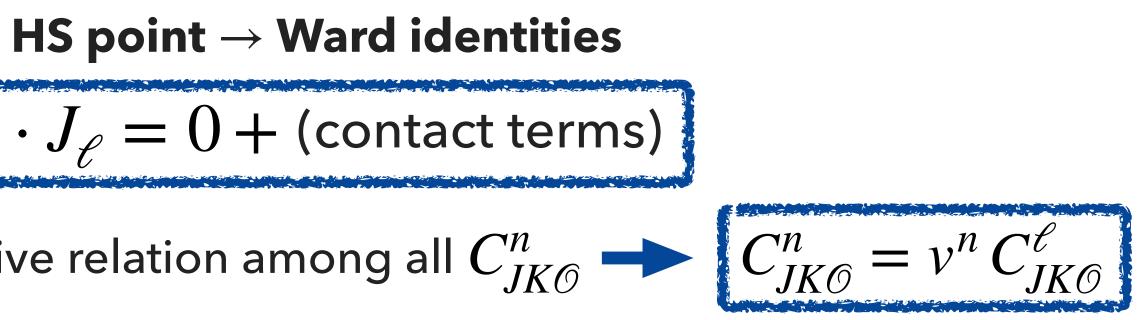


HS point
$$\rightarrow$$
 Ward identities
 $\partial \cdot J_{\ell} = 0 + (\text{contact terms})$
At generic points!
 $\langle \partial J_{\ell} J_{\ell} \mathcal{O} \rangle = 0 \rightarrow \text{Recursive relation among all } C_{JJ\mathcal{O}}^n \rightarrow C_{JJ\mathcal{O}}^n = v^n C_{JJ\mathcal{O}}^{\ell}$
 $\rightarrow C_{JJ\mathcal{O}}^{HS} = C_{JJ\mathcal{O}}^{\ell} \sum_{n=0}^{\ell} w_n v^n = 0$? No :(...Need to work a bit harder
Integrated Ward identity!
 $\text{Integrated Ward identity!}$
 $\text{Integrated Ward identity!}$
 $\text{Integrated charges: } Q_{\ell} = \int_{S} J_{\ell} \rightarrow [Q_{\ell}, J_{\ell}] \sim \sum_{s} J_{s} \qquad \langle J_{\ell} \mathcal{O} \rangle = 0 \quad \forall J_{\ell}$
 $2. \text{ Use: } \int_{S} \langle J_{\ell} J_{\ell} \mathcal{O} \rangle \propto \langle [Q_{\ell}, J_{\ell}] \mathcal{O} \rangle \rightarrow \langle k C_{JJ\mathcal{O}}^{\ell} \propto \sum_{\ell} \langle J_{\ell} \mathcal{O} \rangle = 0 \rightarrow C_{JJ\mathcal{O}}^{\ell} = 0$
Non-vanishing!





At generic points!
•
$$\langle \partial J_{\ell} K_{\ell-1} \mathcal{O} \rangle = 0$$
 \longrightarrow Recursive relation among all $C_{JK\mathcal{O}}^n \rightarrow C_{JK\mathcal{O}}^n = v^n C_{JK\mathcal{O}}^{\ell}$
 $\langle J_{\ell}J_{\ell}\mathcal{O} \rangle$: ℓ equations for $\ell + 1$ variables
 $\langle J_{\ell}K_{\ell-1}\mathcal{O} \rangle$: ℓ equations for ℓ variables
• Integrated Ward identity!
• Integrated Ward identity!
• Integrated Ward identity!
• Define conserved charges: $Q_{\ell} = \int_{S} J_{\ell} \rightarrow [Q_{\ell}, K_{\ell-1}] \sim \sum_{s} K_{s}$
• $Q_{\ell} = \int_{S} J_{\ell} \rightarrow [Q_{\ell}, K_{\ell-1}] \sim \sum_{s} K_{s}$
• $L_{\ell} = K_{\ell} =$





At generic points! • $\langle \partial J_{\ell} K_{\ell-1} \mathcal{O} \rangle = 0 \longrightarrow$ Recursive relation among all $C_{JK\mathcal{O}}^n \longrightarrow$ $\langle J_{\ell} J_{\ell} \mathcal{O} \rangle$: ℓ equations for $\ell + 1$ variables $\langle J_{\ell} K_{\ell-1} \mathcal{O} \rangle$: ℓ equations for ℓ variables $\langle J_{\ell} K_{\ell-1} \mathcal{O} \rangle$: ℓ equations for ℓ variables

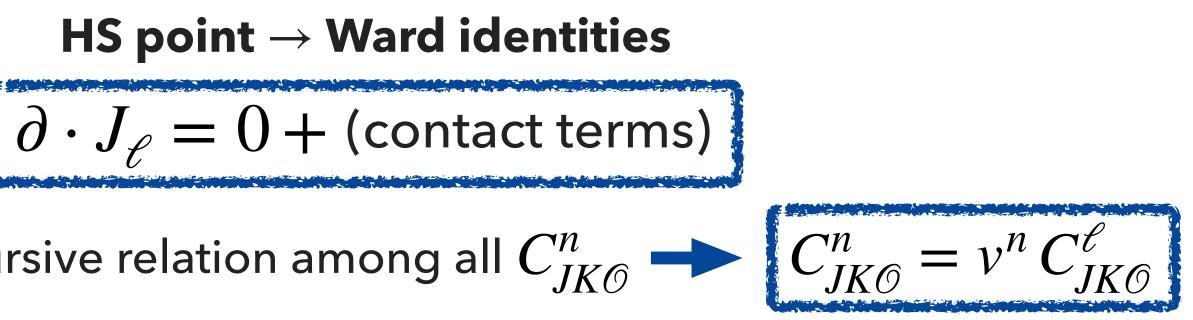
Charge conservation identity! [Maldacena, Zhiboedov '11]
 Involves presence of energy-momentum tensor

 Dynamical gravity in the bulk

$$\left\langle [Q_{\ell}, T_2 K_{\ell-1} \mathcal{O}] \right\rangle = 0 \longrightarrow \mathcal{B}_{\ell} \left\langle J_{\ell} K_{\ell-1} \right\rangle$$

Non-vanishing! [Maldacena, Zhiboedov '11]

Goal: Show that only consistent solution is such that this contribution vanishes



Enough? No:(... Need to work a bit harder

