

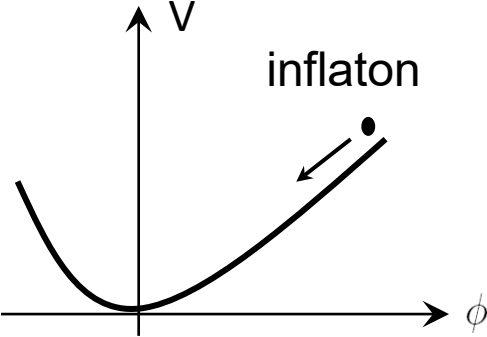
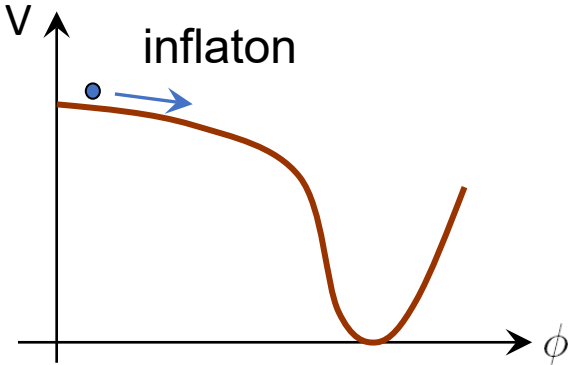
# Primordial Features as Probes of Inflation Models and High Energy Physics

Xingang Chen

CfA Harvard

2106.07546, 2108.10110, with Braglia, Hazra;  
2205.01107, with Ebadi, Kumar;  
2210.07028, with Braglia, Hazra, Pinol;  
2303.03406, with Fan, Li.

# Simplest Single Field Inflation Models



## Primordial Features

Motivations:

- Fine-tuning problem of potential
- Inflationary landscape

Low energy trajectory may not be smooth, i.e., sharp features, periodic features  
Orthogonal directions are lifted by potentials and represented by massive fields



Small (or large) components that significantly depart from scale-invariance, in power spectrum or non-Gaussianities,  
Induced by scale-dependent physics

However, no known model-independent lower bound on the amplitudes

- Sharp feature: Temporary deviation from attractor, due to a sharp feature

E.g., step, bump, kink, interactions, ... ..

(Starobinsky, 92, ...)

$$\frac{\Delta P_\zeta}{P_{\zeta 0}} \propto 1 - \cos(2k/k_f)$$

With a highly model-dependent envelop



- Resonance feature: (Semi)-periodic background oscillations with frequency larger than Hubble  
(XC, Easter, Lim, 06; Flauger et al, 09, ...)

Inflation  
example:

Background oscillation with fixed frequency  
Inflaton oscillation with time-dependent (decrease) frequency



Resonate mode by mode

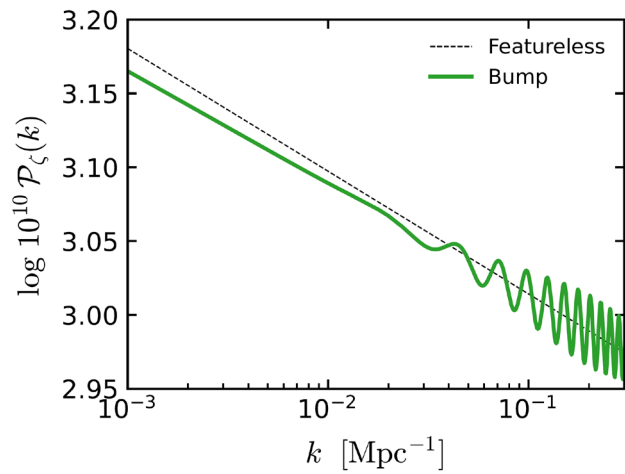
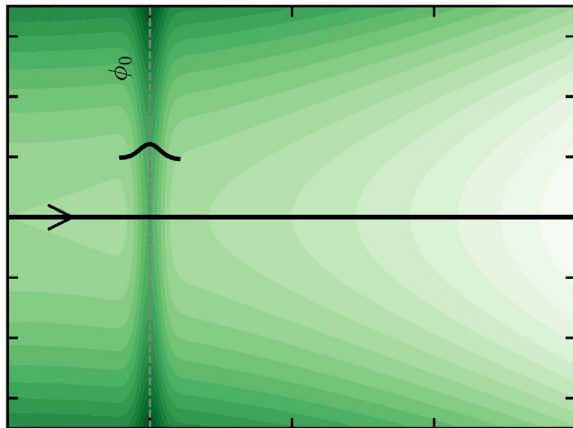
$$\frac{\Delta P_{\zeta}}{P_{\zeta 0}} \propto \sin(\Omega \log(2k) + \phi)$$

With a model-dependent envelop

# Feature Model Examples

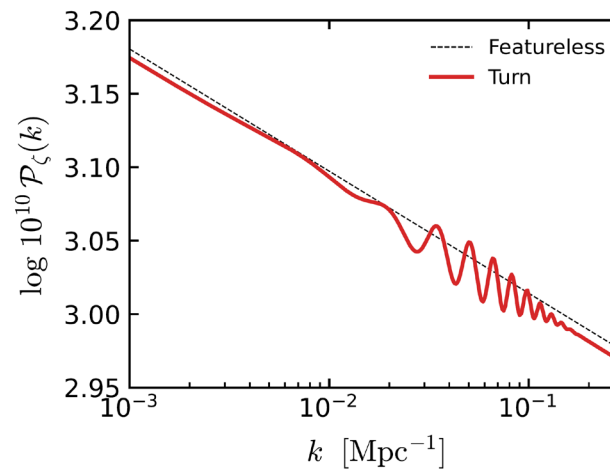
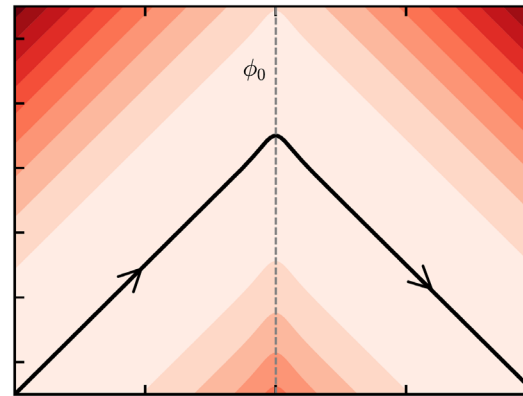
Sharp feature in potential  
e.g. step, bump, dip, ...

(Starobinsky, 92; Adam, Cresswell, Easter, 01, ...)



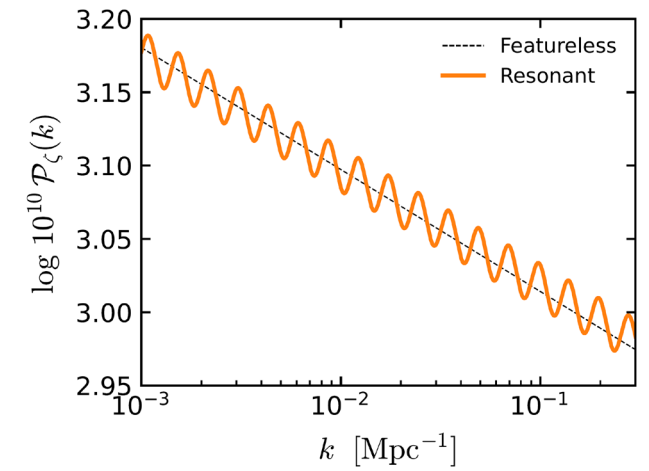
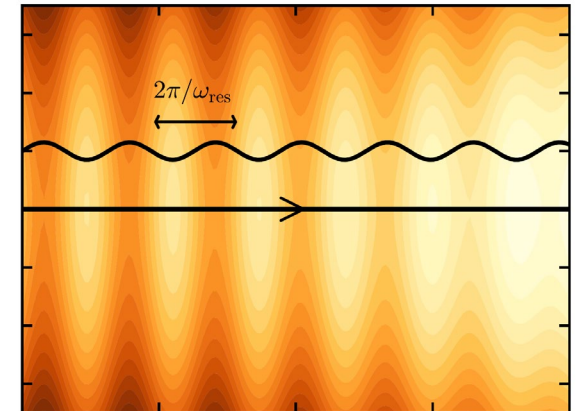
Multifield sharp-turn

(Achucarro et al, 10)



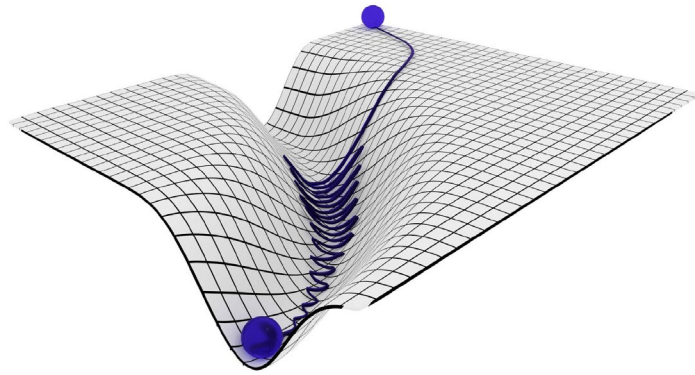
Period features in potential

(XC, Easter, Lim, 06; Flauger et al, 09, ...)

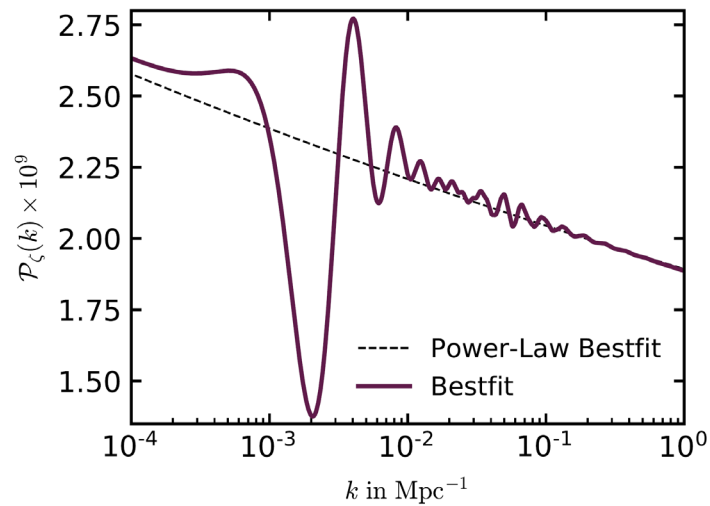
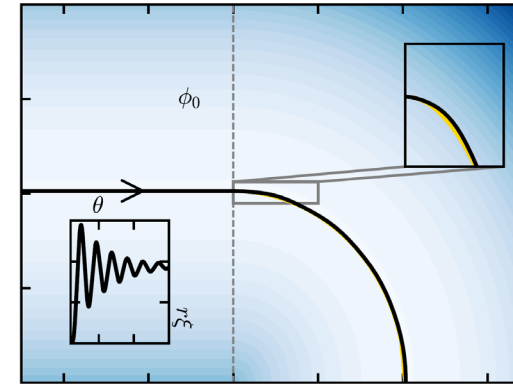


## Feature Model Examples

Tachyonic falling inducing massive field oscillation



Bending trajectory inducing massive field oscillation



Sophisticated combination  
of sharp and resonance features

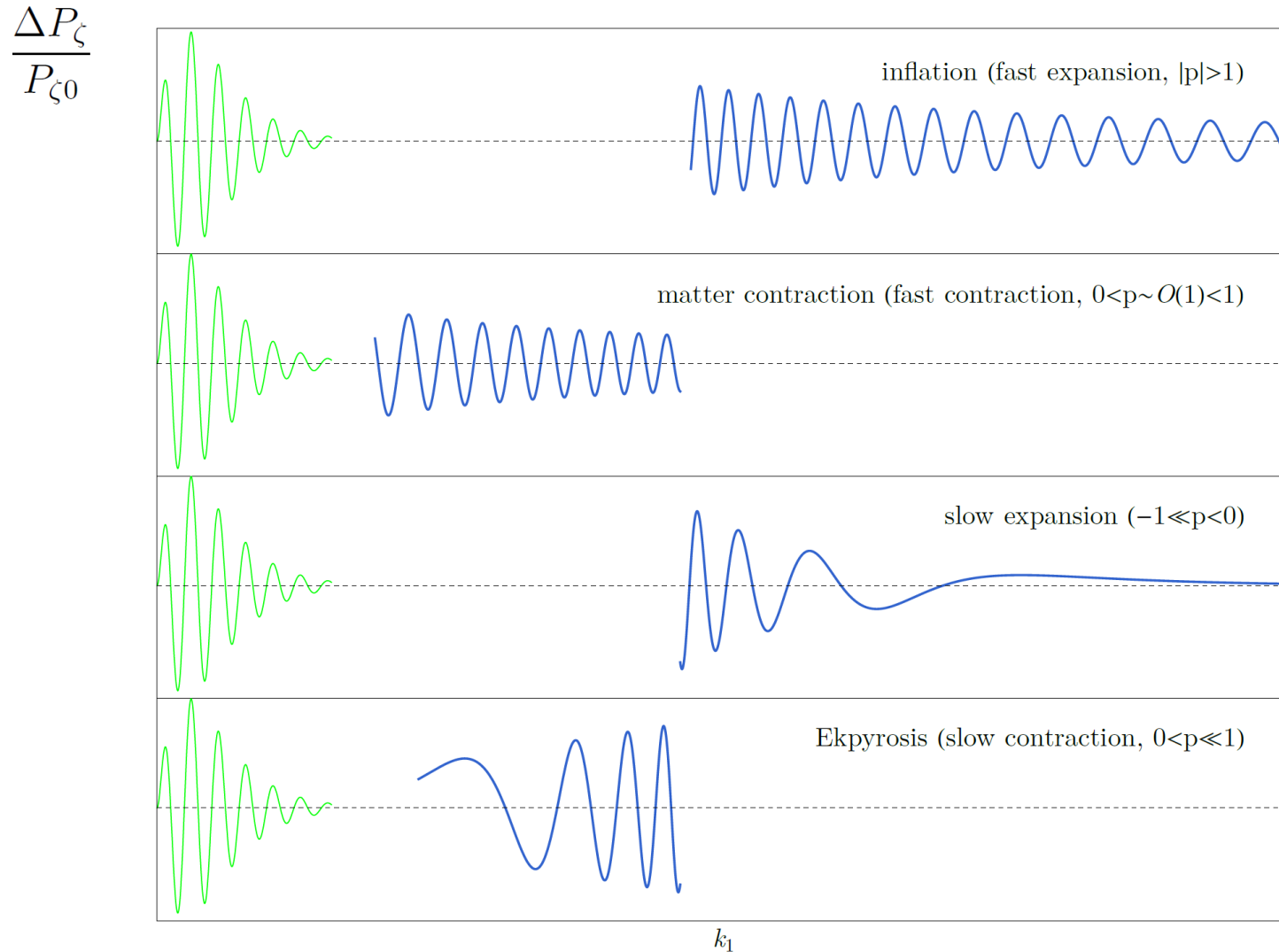
## **(Classical) Primordial Standard Clocks**

# Measuring $a(t)$ of Different Scenarios

(XC, 11, XC, Namjoo, Wang, 14)

In both power spectra (as corrections) and non-Gaussianities

Qualitative sketch



$$a(t) = a_0 \left( \frac{t}{t_0} \right)^p$$

## The Clock Signal in Classical PSC

The background oscillation resonates with curvature fluctuations mode by mode

The clock signal:  $\sim \sin \left[ p \frac{m}{m_{h,0}} \left( \frac{K}{k_r} \right)^{1/p} + \varphi \right]$   $K \equiv k_1 + k_2 = 2k_1$  for power spectrum

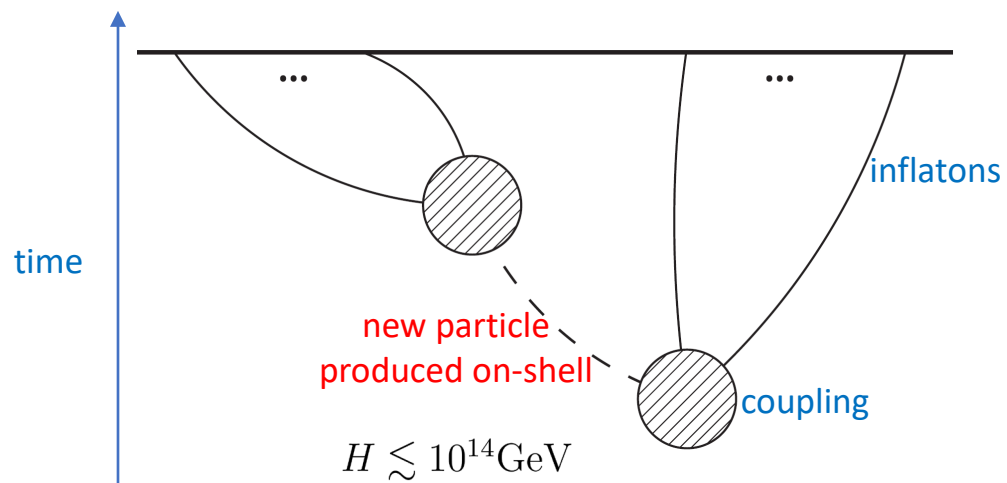
horizon mass  
at time of sharp feature

**Inverse function of  $a(t)$**

This phase pattern is a direct measure of  $a(t)$

## Cosmological Collider Physics

## Quantum production of massive particles



New particles, with mass up to  $O(H)$ , are produced on-shell due to quantum fluctuations,  
They couple to inflatons, and leave imprints in inflaton correlation functions



What are the observational signatures of these particle states?

**Mass and spin spectra** of intermediate state are encoded in **soft limits** of non-Gaussianities:

E.g. **Squeezed limit bispectrum**

$$S \xrightarrow[\text{limit}]{\text{squeezed}} e^{-\pi\mu} \left( \frac{k_{\text{long}}}{k_{\text{short}}} \right)^{\frac{1}{2} \pm i\mu} P_s(\cos \theta)$$

↓ Boltzmann suppression
 ↓ mass
 ↓ spin

$$\mu = \sqrt{\frac{m^2}{H^2} - \frac{9}{4}}$$

**Cosmological Collider Physics**

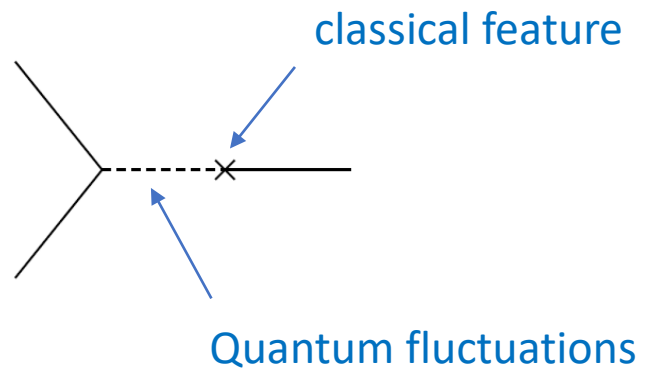
(XC, Wang, 09  
Arkani-Hamed, Maldacena, 15)

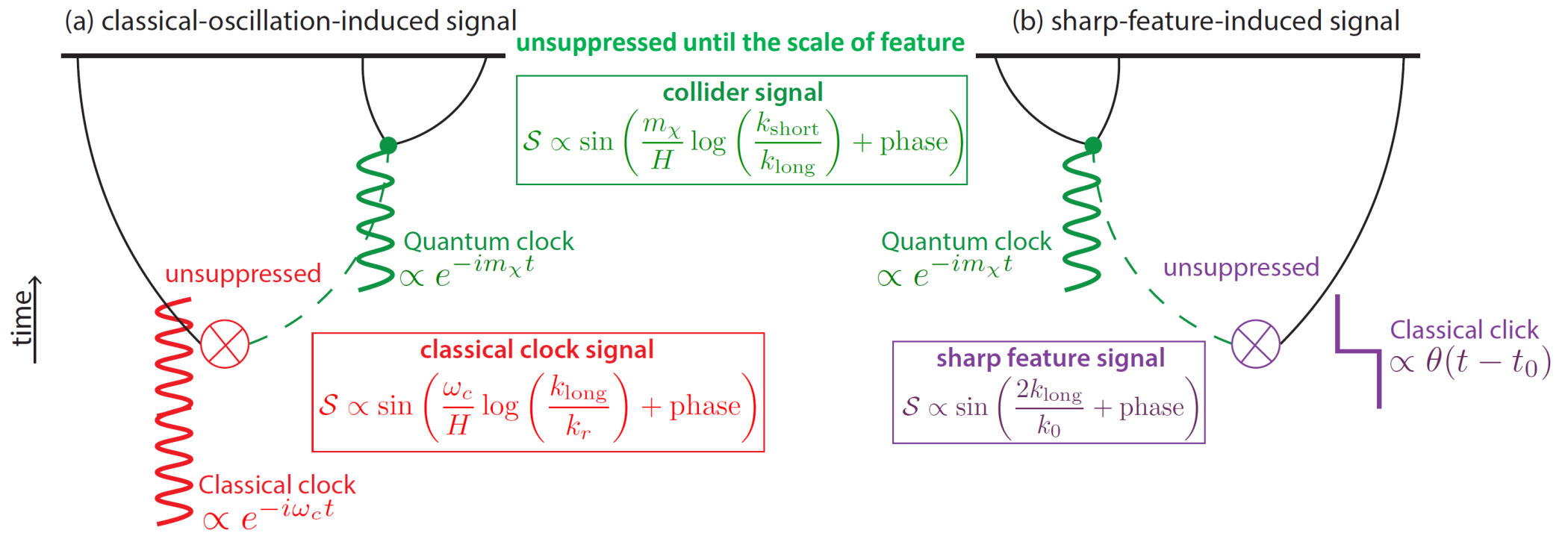
- Amplitude of non-G is very model-dependent, open for predictions from model building

## Classical Cosmological Collider Physics

(XC, 11; XC, Ebadi, Kumar, 22)

Primordial features inject extra energy to inflation models  
and  
boost the energy of the cosmological collider by orders of magnitude





## **An Example of Particle Dark Matter Model**

## Axion Model as Example of Massive Field

(XC, Fan, Li, 23)

$$\mathcal{L} = -\frac{(\partial_\mu\phi)^2}{2} - |\partial_\mu\chi|^2 - V_\phi(\phi) - V_\chi(\chi)$$

$$V_\chi(\chi) = \frac{\lambda}{2} \left( |\chi|^2 - \frac{f_a^2}{2} \right)^2$$

Add a coupling:  $-\frac{c}{\Lambda^2}(\partial\phi)^2|\chi|^2$

A sharp feature in inflaton potential can induce oscillation of radial mode



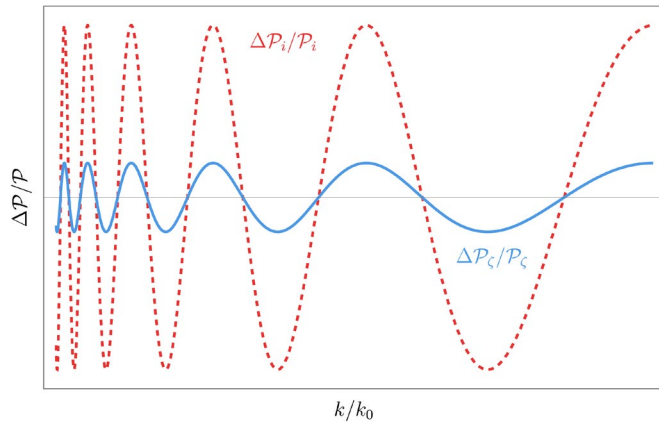
Imprint in both inflaton and axion fields



Correlated clock signals in both curvature and dark matter isocurvature perturbations



- A signature of Peccei-Quinn field during inflation
- Feature signal dominates in DM isocurvature perturbation



$$\left| \frac{\Delta P_\zeta}{P_\zeta} \right|_{\text{clock;amp}} \approx 0.019 \left( \frac{q}{0.02} \right)^2 \left( \frac{bV_{\phi 0}}{0.3\dot{\phi}_0^2} \right) \left( \frac{\dot{\phi}_0}{(60H)^2} \right)^2 \left( \frac{40H}{f_I} \right)^{7/2} \left( \frac{1}{\lambda} \right)^{3/4}$$

$$\left| \frac{\Delta P_i}{P_i} \right|_{\text{clock;amp}} \approx 0.96 \left( \frac{q}{0.02} \right) \left( \frac{bV_{\phi 0}}{0.3\dot{\phi}_0^2} \right) \left( \frac{\dot{\phi}_0}{(60H)^2} \right)^2 \left( \frac{40H}{f_I} \right)^{7/2} \left( \frac{1}{\lambda} \right)^{3/4}$$

- Break the degeneracy in measurement of  $f_I$  and  $H$

Dark matter physics  
(amplitude of axion isocurvature):

$$\beta \propto \theta_i^2 H^2 f_I^{1/3}$$

Cosmological collider physics

$$\mu_\rho = \sqrt{\lambda} (f_I/H)$$

Degeneracy directions in  $f_I$ - $H$  plane  
are quite orthogonal to each other

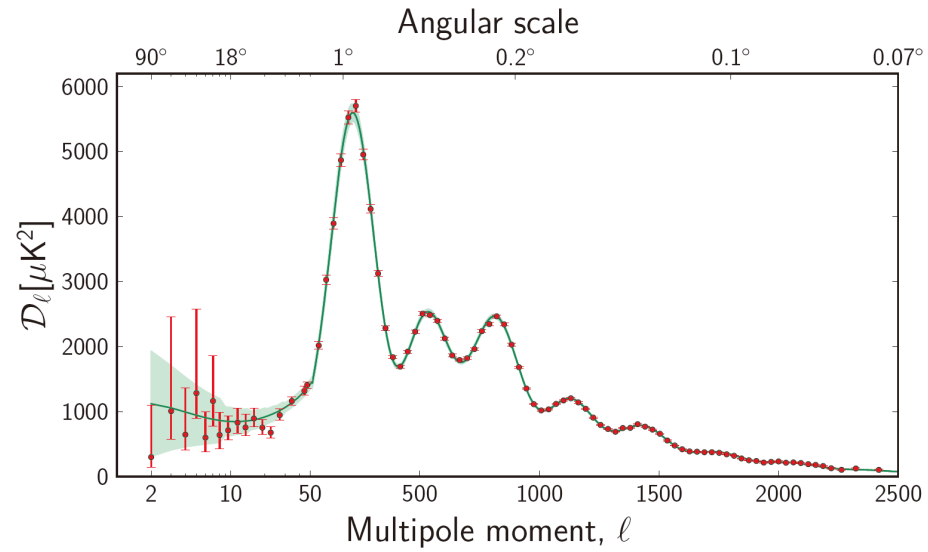
**Summary: What may we learn from signals of primordial features?**

- Inflation or not inflation
- Details of (non-) inflationary potential or internal space
- High energy physics at super-Hubble energy scales

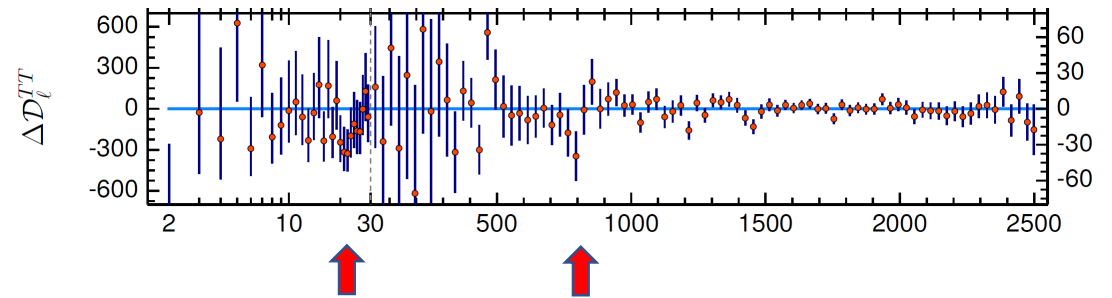


## CMB Power Spectra Residuals

There are some interesting, statistically marginal, anomalies in CMB residual data



CMB TT residuals (also has counterparts in TE and EE)

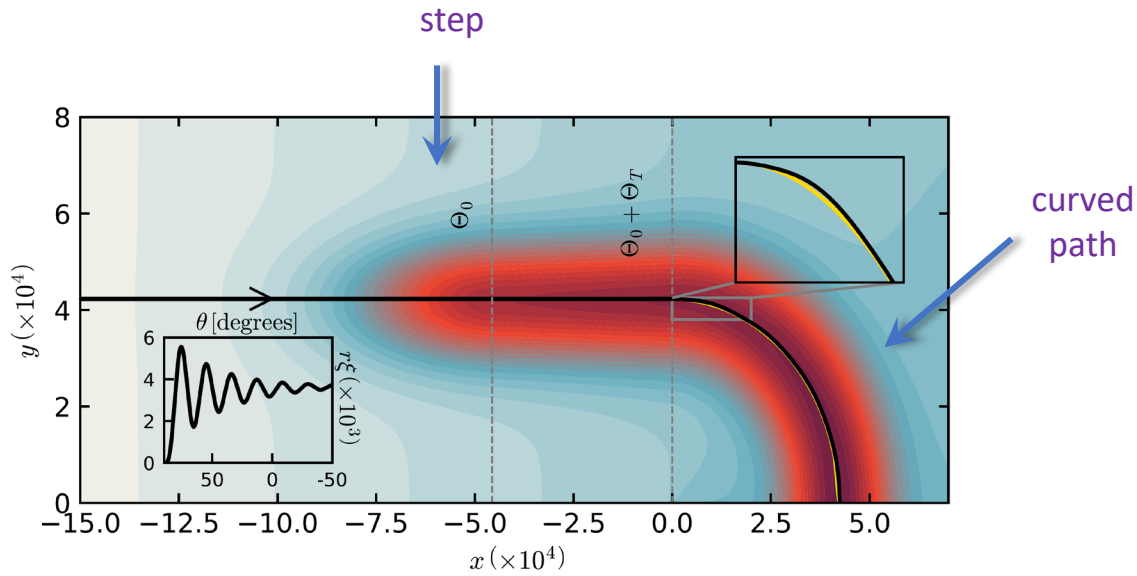


(Planck, 13, 18)

The two well-separated features in CMB may be connected by the Standard Clock effect; or explained by other feature models.

# A Classical Primordial Standard Clock Model

(Braglia, XC, Hazra, 21)



$$\mathcal{L} = -\frac{1}{2} [1 + \Xi(\Theta)\sigma]^2 (\partial\Theta)^2 - \frac{1}{2} (\partial\sigma)^2 - V(\Theta, \sigma)$$

$$\Xi(\Theta) = \xi \text{Heav}(\Theta - \Theta_0 - \Theta_T)$$

$$V(\Theta, \sigma) = V_{\text{inf}} \left\{ 1 - \frac{1}{2} C_\Theta \Theta^2 + C_\sigma \left[ 1 - \exp \left( -\frac{(\Theta - \Theta_0)^2}{\Theta_f^2} \text{Heav}(\Theta - \Theta_0) - \frac{\sigma^2}{\sigma_f^2} \right) \right] \right\}$$

(6 more parameters than standard model)

Birdseye view of the potential

1<sup>st</sup> stage inflation → Rolling off a cliff → Fast-roll, entering a curved path of valley

→ Overshoots bottom of valley and climbs on cliff → Excites classical oscillation of a massive field

→ Oscillation decays, settles down to 2<sup>nd</sup> stage inflation

Use effective parameters and nested sampling.

To compare models, we use **Bayes Factor**:

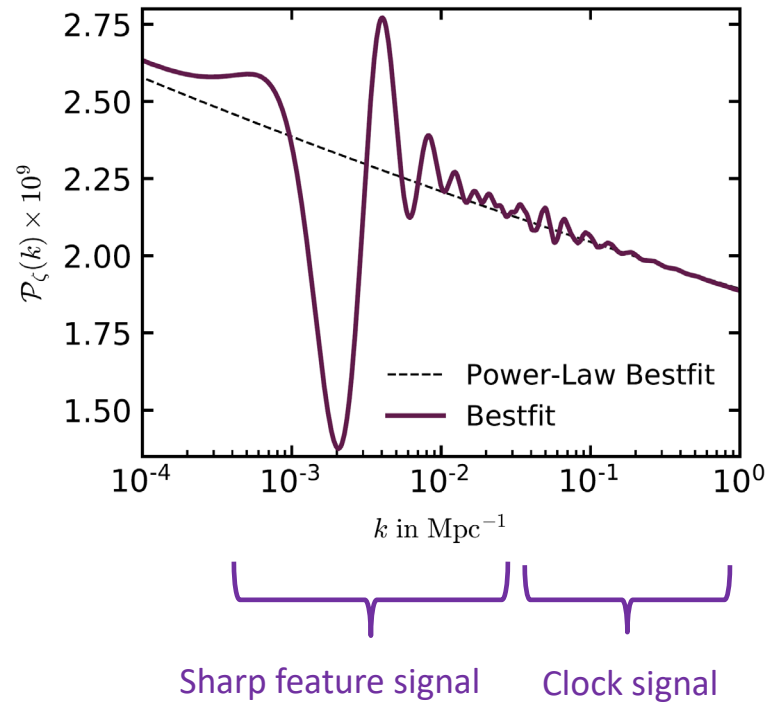
$$B_{12} \equiv \frac{Z_1}{Z_2} = \frac{\int d\boldsymbol{\theta}_1 \pi(\boldsymbol{\theta}_1 | \mathcal{M}_1) \mathcal{L}(\mathbf{x} | \boldsymbol{\theta}_1, \mathcal{M}_1)}{\int d\boldsymbol{\theta}_2 \pi(\boldsymbol{\theta}_2 | \mathcal{M}_2) \mathcal{L}(\mathbf{x} | \boldsymbol{\theta}_2, \mathcal{M}_2)}$$

Z's are marginalized likelihoods for different models:  
probability of data given the model

### Evidence Categories for Bayes Factor given by Jeffreys (1961)

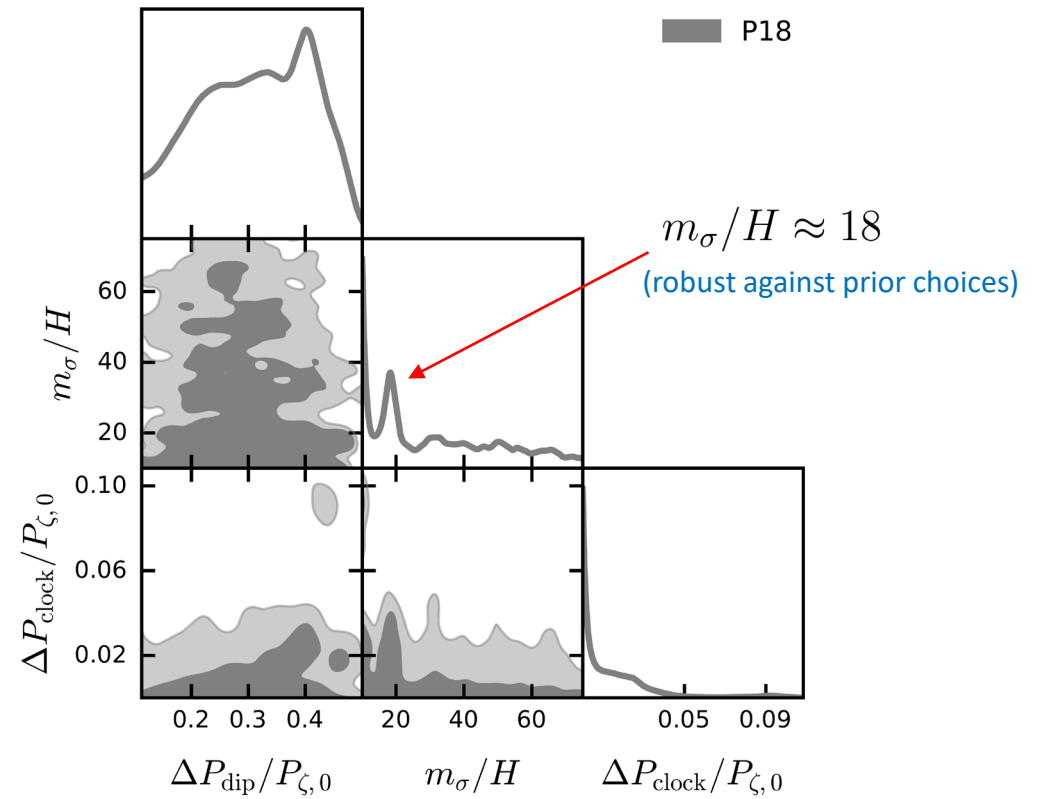
$BF_{ij}$	$\ln(BF_{ij})$	Interpretation
$>100$	$>4.61$	Decisive evidence for $H_i$
$30-100$	$3.40$ to $.61$	Very strong evidence for $H_i$
$10-30$	$2.30$ to $3.40$	Strong evidence for $H_i$
$3-10$	$1.10$ to $2.30$	Substantial evidence for $H_i$
$1-3$	$0$ to $1.10$	Not worth more than a bare mention
$1/3-1$	$-1.10$ to $0$	Not worth more than a bare mention
$1/10-1/3$	$-2.30$ to $-1.10$	Substantial evidence for $H_j$
$1/30-1/10$	$-3.40$ to $-2.30$	Strong evidence for $H_j$
$1/100-1/30$	$-4.61$ to $-3.40$	Very strong evidence for $H_j$
$<1/100$	$<-4.61$	Decisive evidence for $H_j$

## Best-fit, Posteriors, and Bayes Evidence



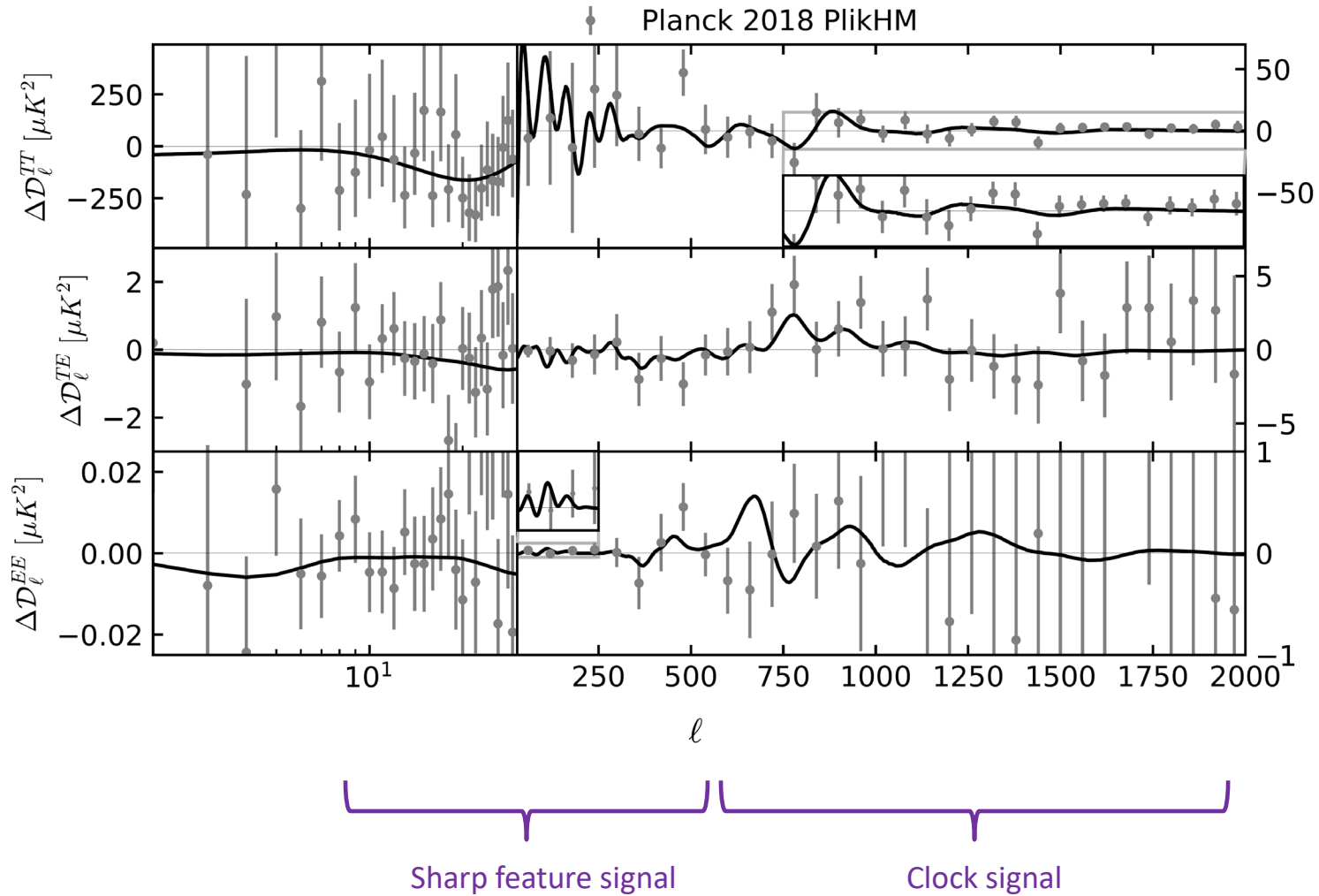
Best-fit:  $\Delta\chi^2 = 19.8$  (with 6 extra parameters)

Bayes factor:  $\ln B \equiv \ln Z_{\text{feature}} - \ln Z_{\text{featureless}} = -0.13 \pm 0.38$

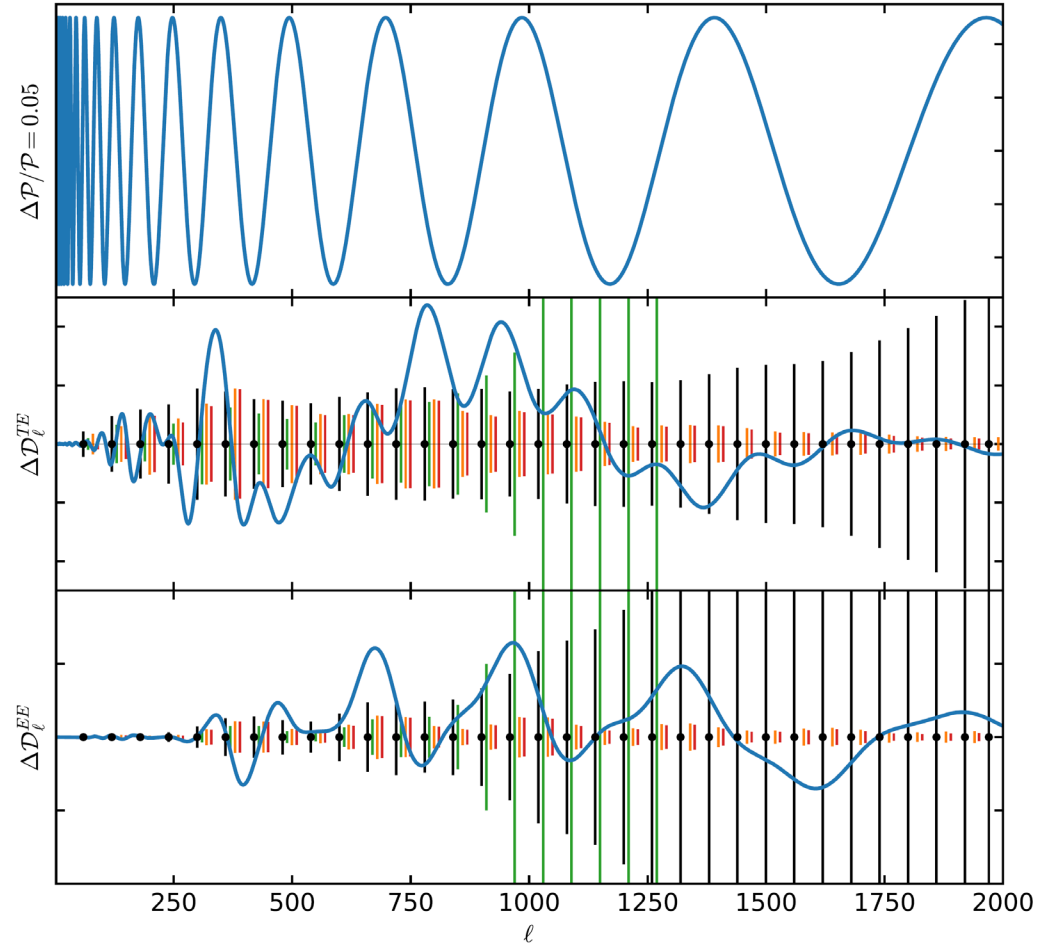
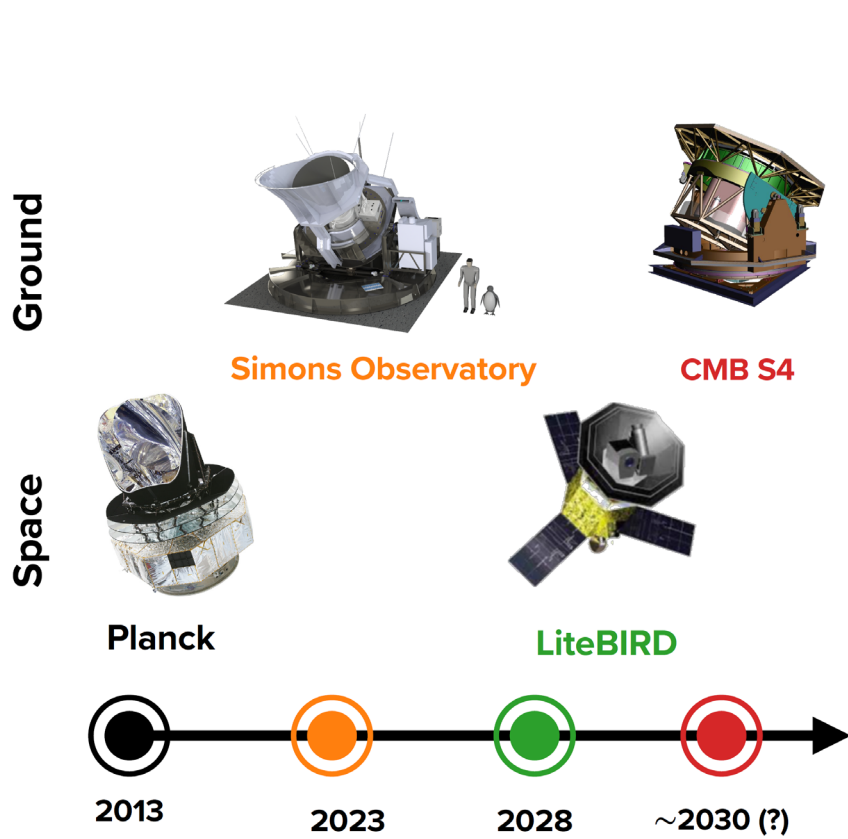


This feature model is currently indistinguishable from Standard Model

# Best-fit v.s. binned residuals (Planck 18)

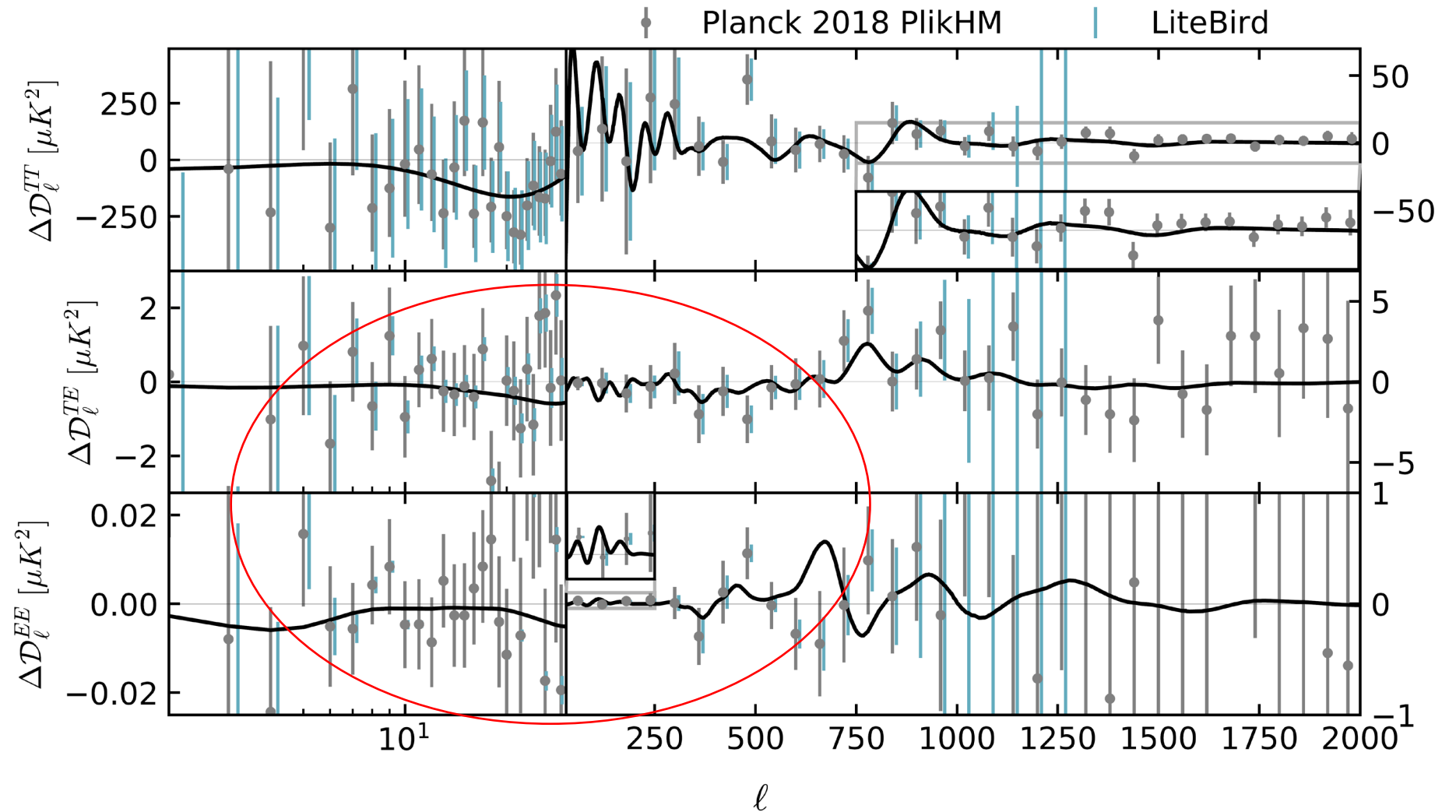
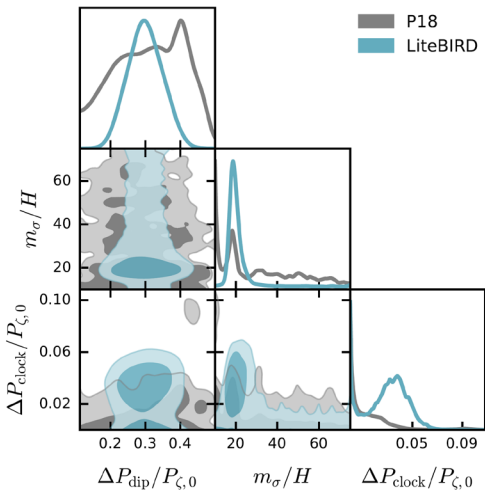


## Future CMB Polarization Experiments

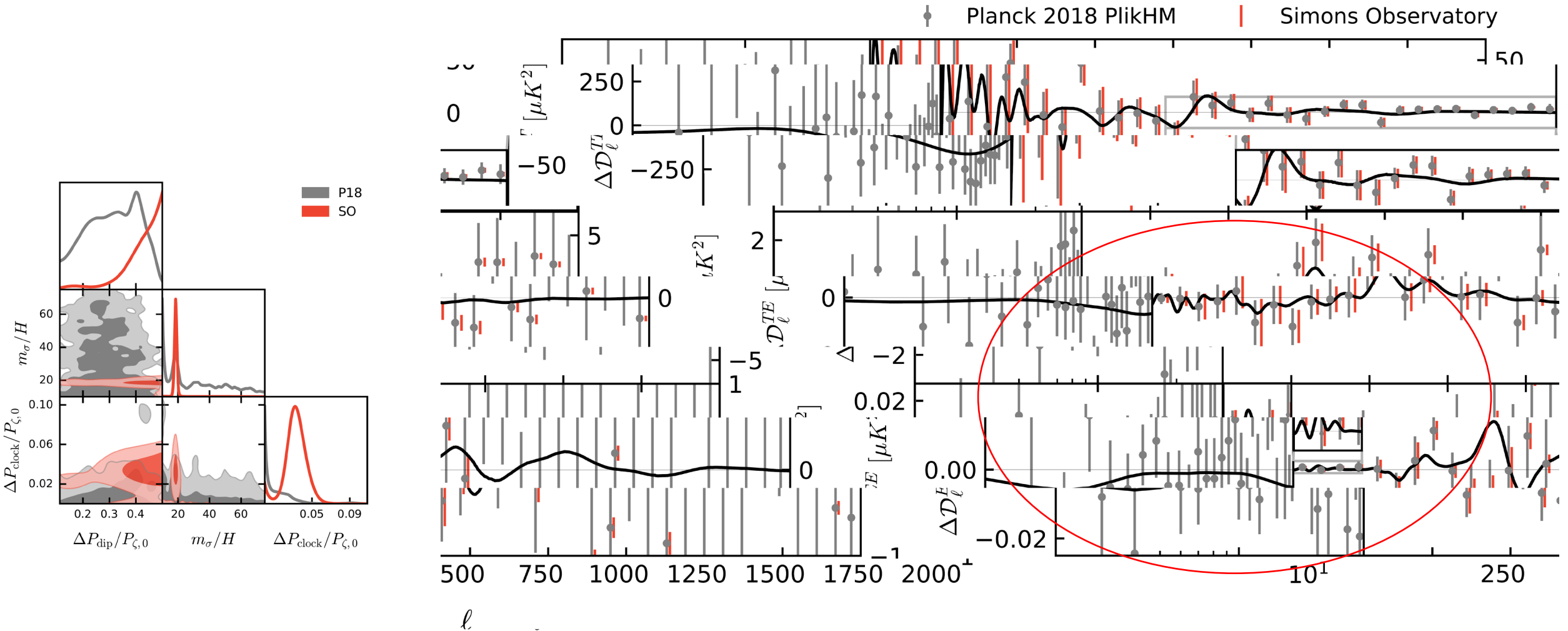


Error bar sizes in the example of resonance model

## Forecast with the best-fit model: LiteBIRD



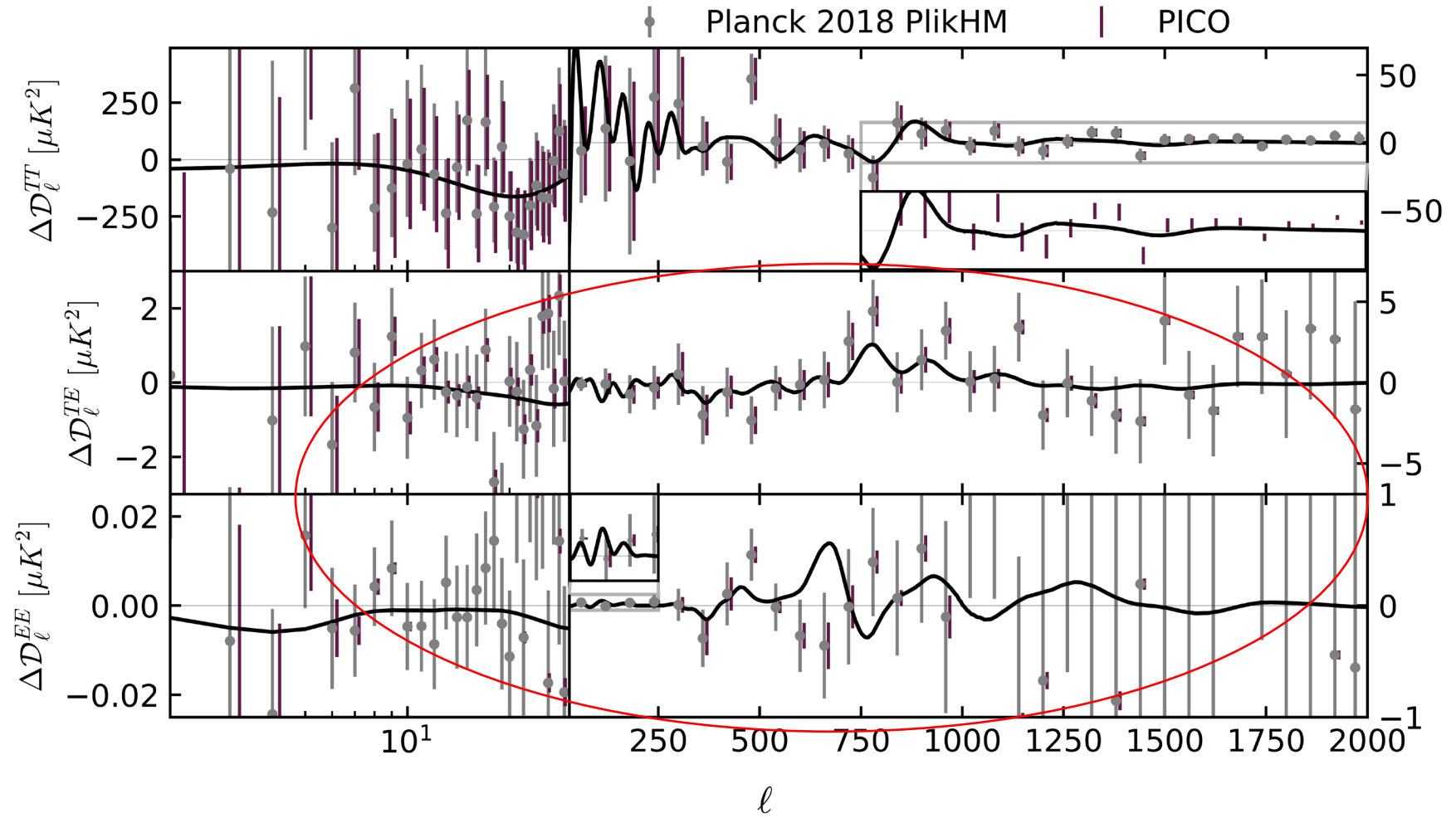
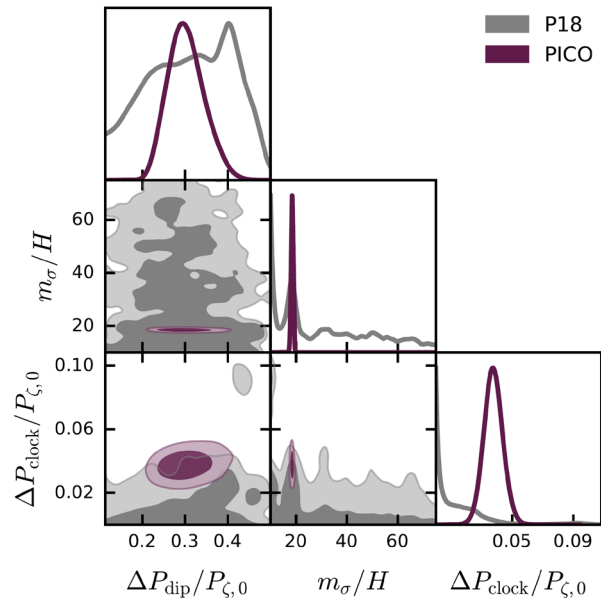
## Forecast with the best-fit model: **Simons Observatory**



The clock signal carries information of  $a(t)$ , which can be used to distinguish the inflation and alternatives in a model-independent fashion.

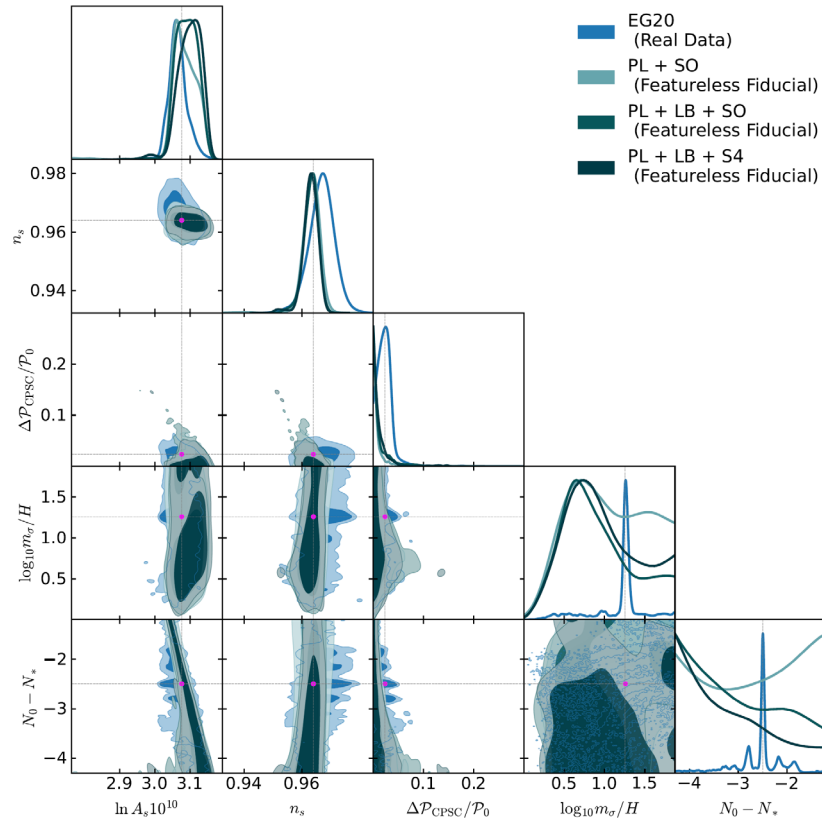


## Forecast with the best-fit model: PICO

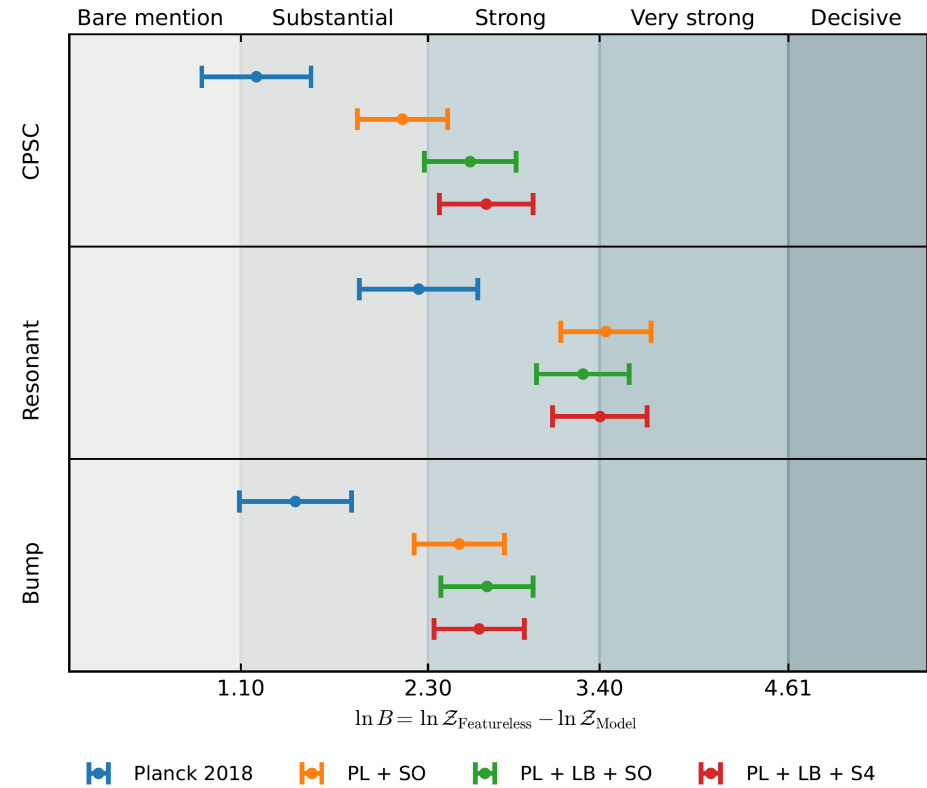


## Constraints on feature model using featureless as fiducial

Example: a CPSC model



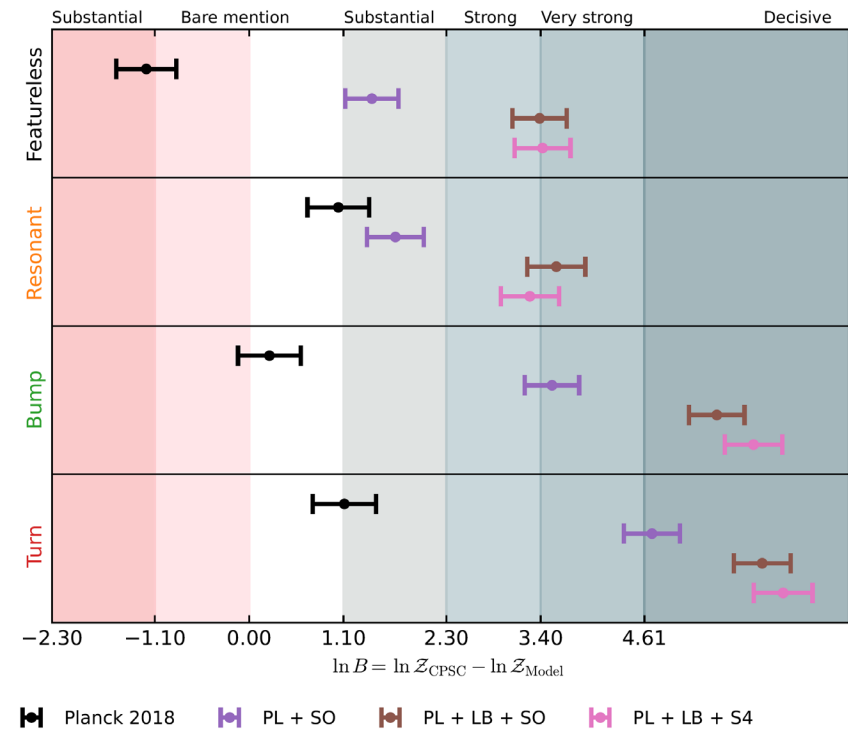
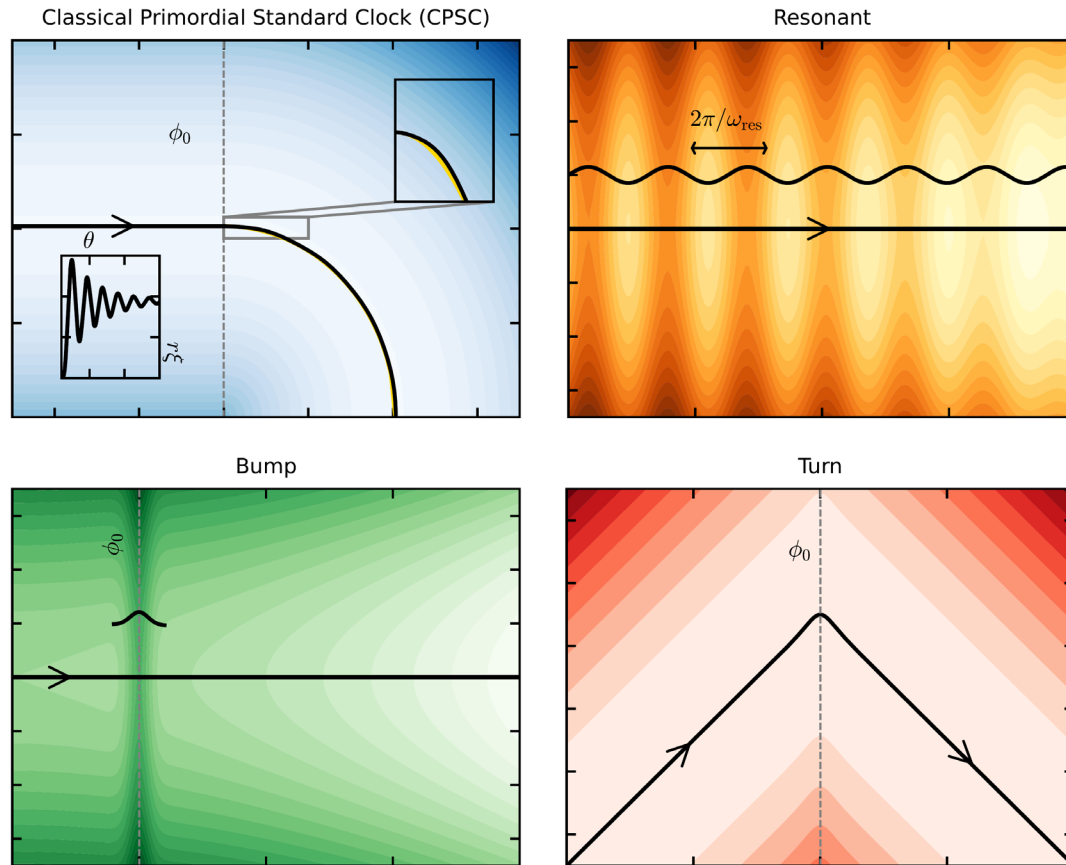
Summary for all example models:  
in terms of Bayes factors



Evidences are generally strong, but not decisive -- feature amplitudes can be zero

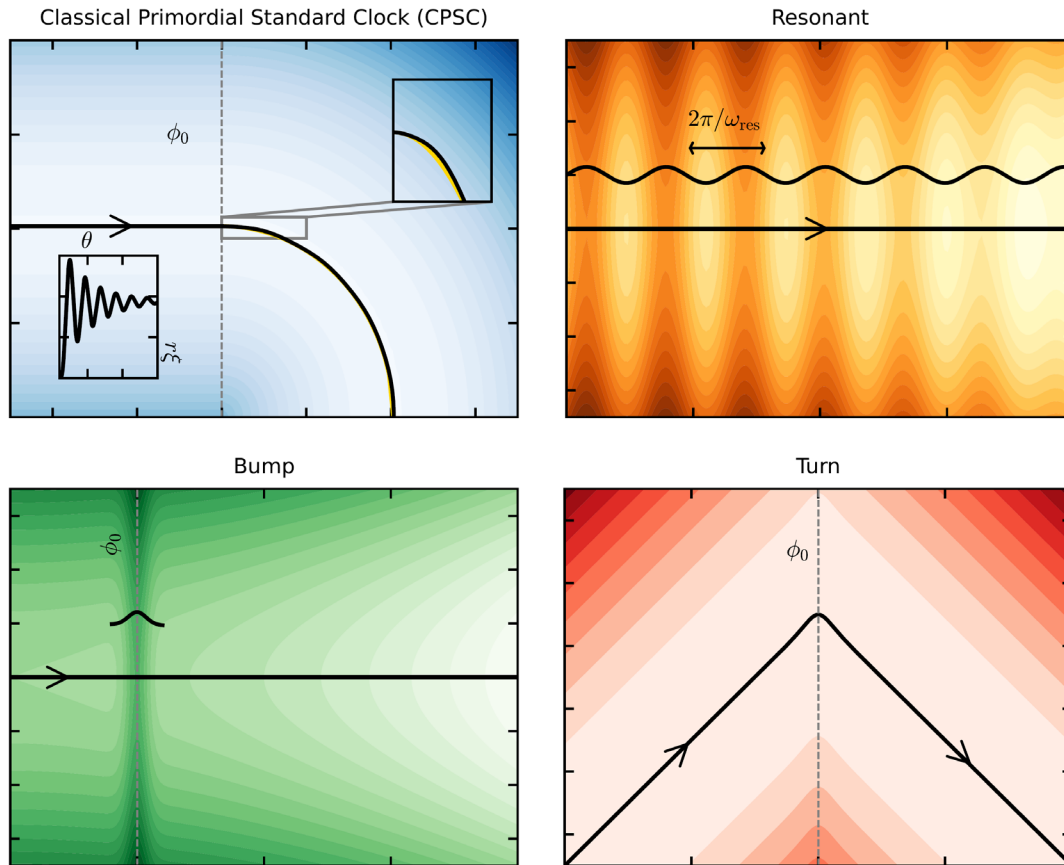
Even if a primordial feature signal were discovered, at initial stages it will likely have multiple possible explanations.

### 1) How confident can we detect the feature if the Universe has such a feature?

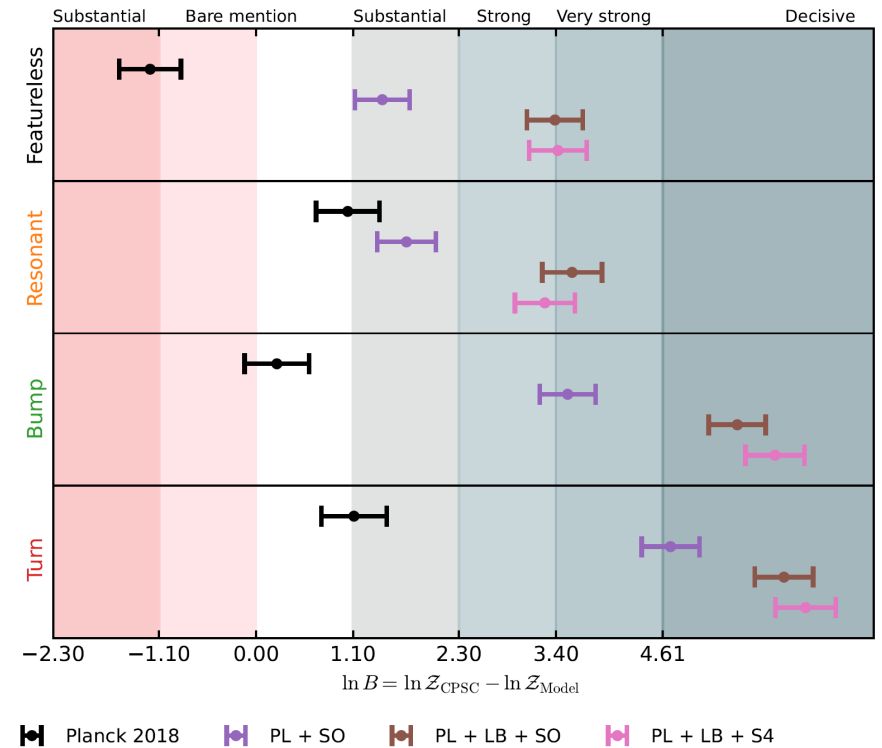


Even if a primordial feature signal were discovered, at initial stages it will likely have multiple possible explanations.

2) To what extent can we distinguish different feature models if the Universe has a feature?



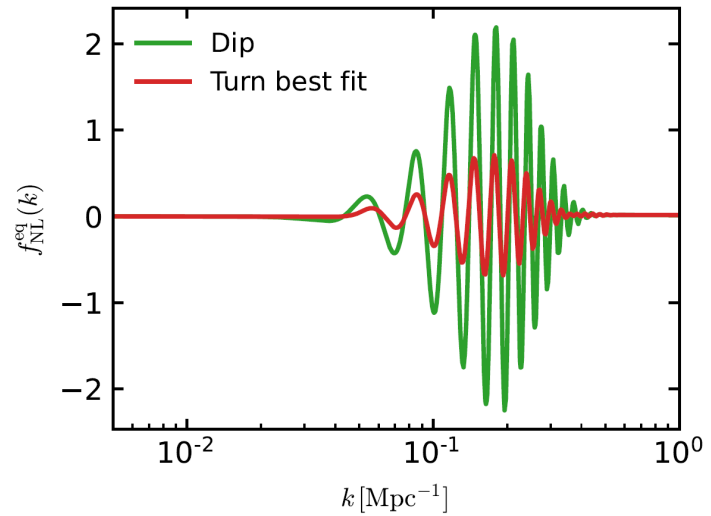
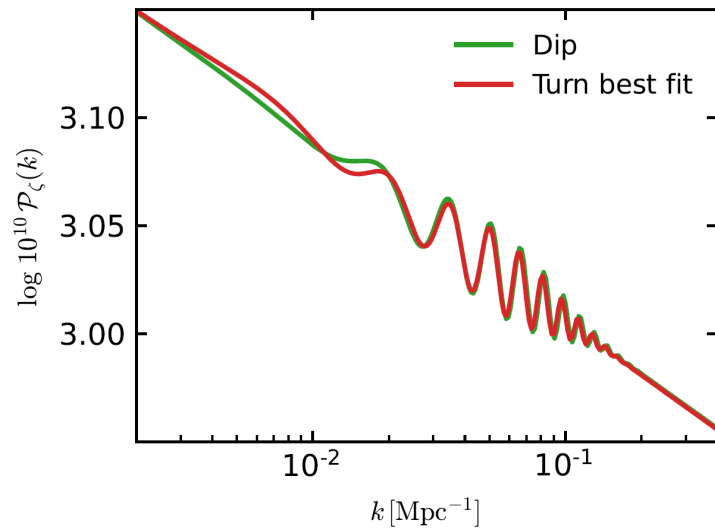
E.g. using the CPSC bestfit as fiducial



Strong to decisive evidences

In some special cases, the two feature models have very similar 2pt.

We may need higher order correlation functions.



## Correlated feature signals should also appear in other observables

- Galaxy surveys

(Huang,Verde,Vernizzi,12; XC, Dvorkin,Huang,Namjoo,Verde,16;  
Ballardini,Finelli,Fedeli,Moscardini,16; Palma,Sapone,Sypsas,17;  
L'Huilier,Shafieloo,Hazra,Smoot,Starobinsky,17;  
Beutler,Biagetti,Green,Slosar,Wallisich,19, .....)

- 21 cm from atomic hydrogen

(XC,Meerburg,Munchmeyer,16; Xu,Hamann,Chen,16 )

- Stochastic gravitational wave background

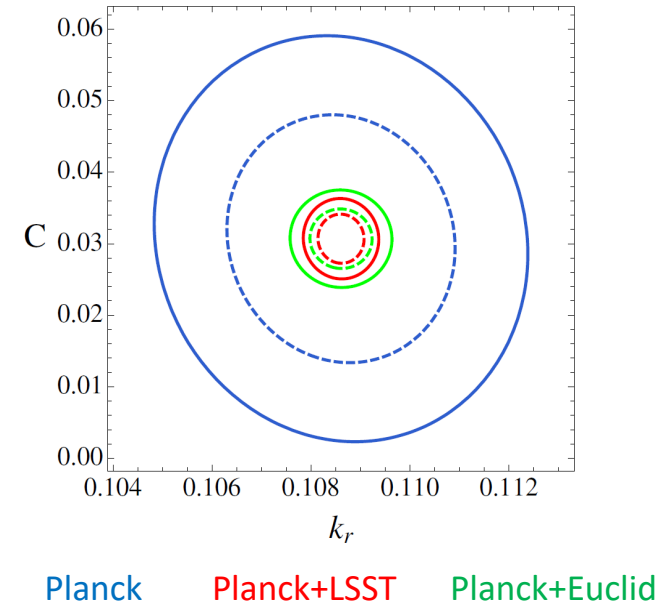
(Fumagalli, Renaux-Patel, 20; Braglia, XC, Hazra, 20; Bodas, Sundrum, 22)

- Non-Gaussianities

(Fergusson, Shellard, Liguori, 14, ... )

- Isocurvature perturbations

(XC, Fan, Lin, 23)



*Thank You !*