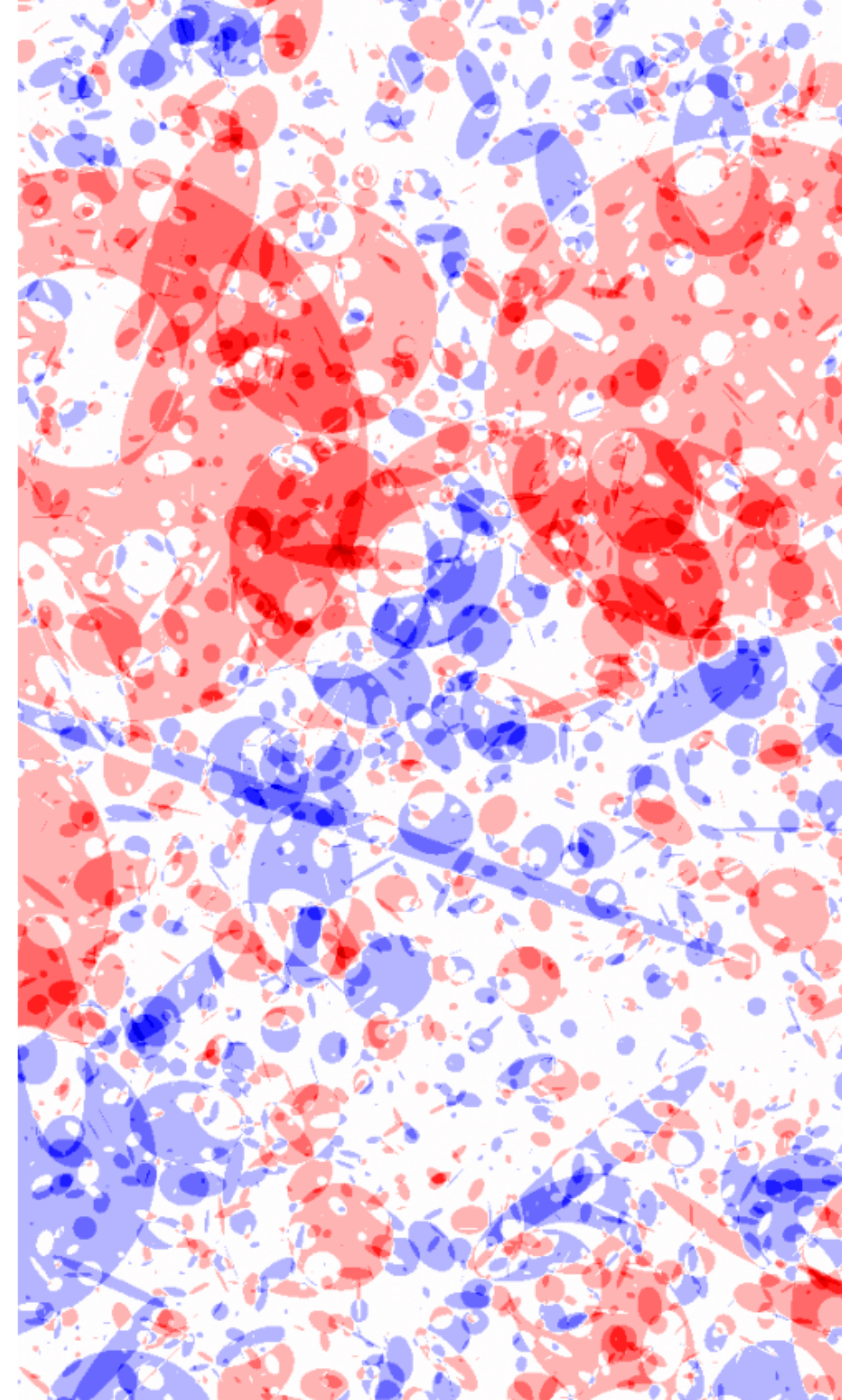


Charge quantisation, Axion strings, and Cosmic birefringence

arXiv:2305.02318 (2023)

arXiv:2111.12741 (2022)

Winston Yin, 2023-04-28



Charge quantisation

Charges beyond the Standard Model

- Electric charges in Standard Model (SM) are multiples of $1/3$
- What are the charge assignments in beyond-the-SM theories?
- Any new fermions with charges less than $1/3$?
- Axion-photon coupling can help us answer this question
- Observable: cosmic birefringence (in CMB) induced by axion strings
- P. Agrawal, A. Hook, J. Huang (2020)

Ultralight axion (-like particles)

- Axions are pseudo-scalar fields, i.e. their values are periodic $a \in [0, 2\pi f_a)$
- Generic product of breaking of global $U(1)$ symmetry [Peccei-Quinn (PQ) symmetry] in models beyond SM
- Ultralight axions with mass $m_a \lesssim H_{\text{cmb}} \simeq 3 \times 10^{-29} \text{ eV}$
 - ~~CP problem~~
 - ~~Dark matter~~
 - Potential as dark energy
 - Predicted in large numbers in string theory "axiverse" scenarios

Axion-photon coupling

After PQ symmetry breaking

- Induces a Chern-Simons (topological) axion-photon coupling

The diagram shows the Lagrangian term for axion-photon coupling: $\mathcal{L} \supset \frac{\mathcal{A} \alpha_{\text{em}}}{4\pi f_a} a F \tilde{F}$. Blue arrows point from text labels to parts of the equation: 'PQ-EM anomaly coefficient' points to \mathcal{A} ; 'fine structure constant' points to α_{em} ; 'axion' points to a ; 'EM field strength' points to $F \tilde{F}$; and 'axion periodicity / PQ symmetry breaking scale' points to f_a .

PQ-EM anomaly coefficient

fine structure constant

axion

EM field strength

axion periodicity /
PQ symmetry breaking scale

$$\mathcal{L} \supset \frac{\mathcal{A} \alpha_{\text{em}}}{4\pi f_a} a F \tilde{F}$$

Anomaly coefficient

$$\mathcal{L} \supset \frac{\mathcal{A} \alpha_{\text{em}}}{4\pi f_a} a F \tilde{F}$$

- f_a is subject to renormalisation
- But \mathcal{A} is not, so its value is fixed on all energy scales

$$\mathcal{A} = \sum_f Q_f^{\text{PQ}} \left(Q_f^{\text{EM}} \right)^2$$

all beyond-SM fermions \rightarrow f \leftarrow PQ charges, integers \leftarrow electric charges

- \mathcal{A} is integer multiple of square of the smallest electric charge beyond SM
- Beyond-the-SM theories predict different $\mathcal{A} = \mathcal{O}(1)$, e.g. 4/3 for minimal GUT
- Axion strings will allow us to measure \mathcal{A} directly, unaffected by f_a

Cosmic birefringence

Induced by axion-photon coupling

$$\mathcal{L} \supset \frac{\mathcal{A} \alpha_{\text{em}}}{4\pi f_a} a F \tilde{F}$$

- Polarisation of CMB photons is rotated by intervening axion field

$$\Delta\Phi = \frac{\mathcal{A} \alpha_{\text{em}}}{2\pi f_a} \Delta a$$

- Rotation angle is proportional to net change in axion value along photon path
- Typically $a \ll 2\pi f_a$, so effect is very weak (naively)
- With axion strings, $\Delta a \approx n 2\pi f_a$ for integers n for any CMB photon, possible due to axion periodicity
- $\Delta\Phi \approx n \mathcal{A} \alpha_{\text{em}} = n \mathcal{O}(\text{deg})$, rotation angle is macroscopic and quantised!

Axion strings

Axion strings

Topological defects in axion field

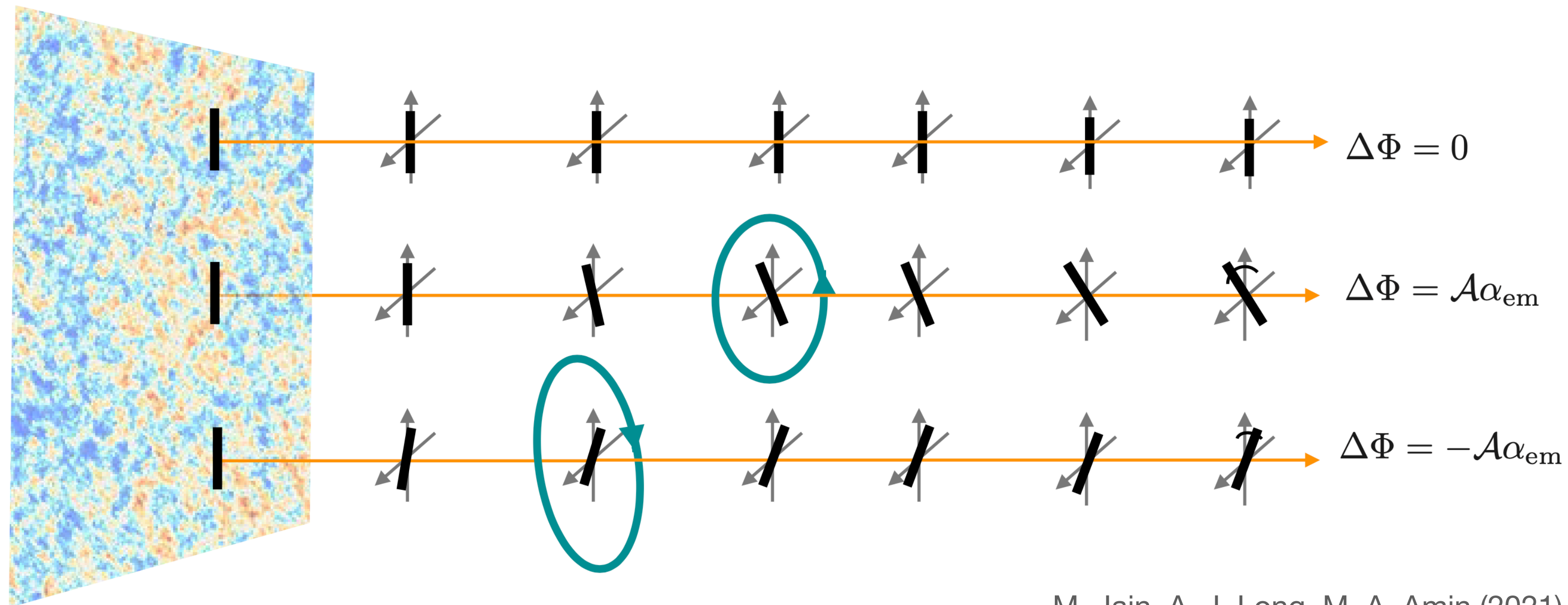
- a changes by exactly one period $2\pi f_a$ around a string
- Topologically stable (cannot be continuously deformed into vacuum)
- Formed in large numbers by Kibble mechanism if PQ symmetry breaking occurred after inflation
- Ultralight axion strings have no detectable gravitational effect

Cosmic birefringence

Induced by axion strings

$$\Delta\Phi = \frac{\mathcal{A}\alpha_{\text{em}}}{2\pi f_a} \Delta a$$

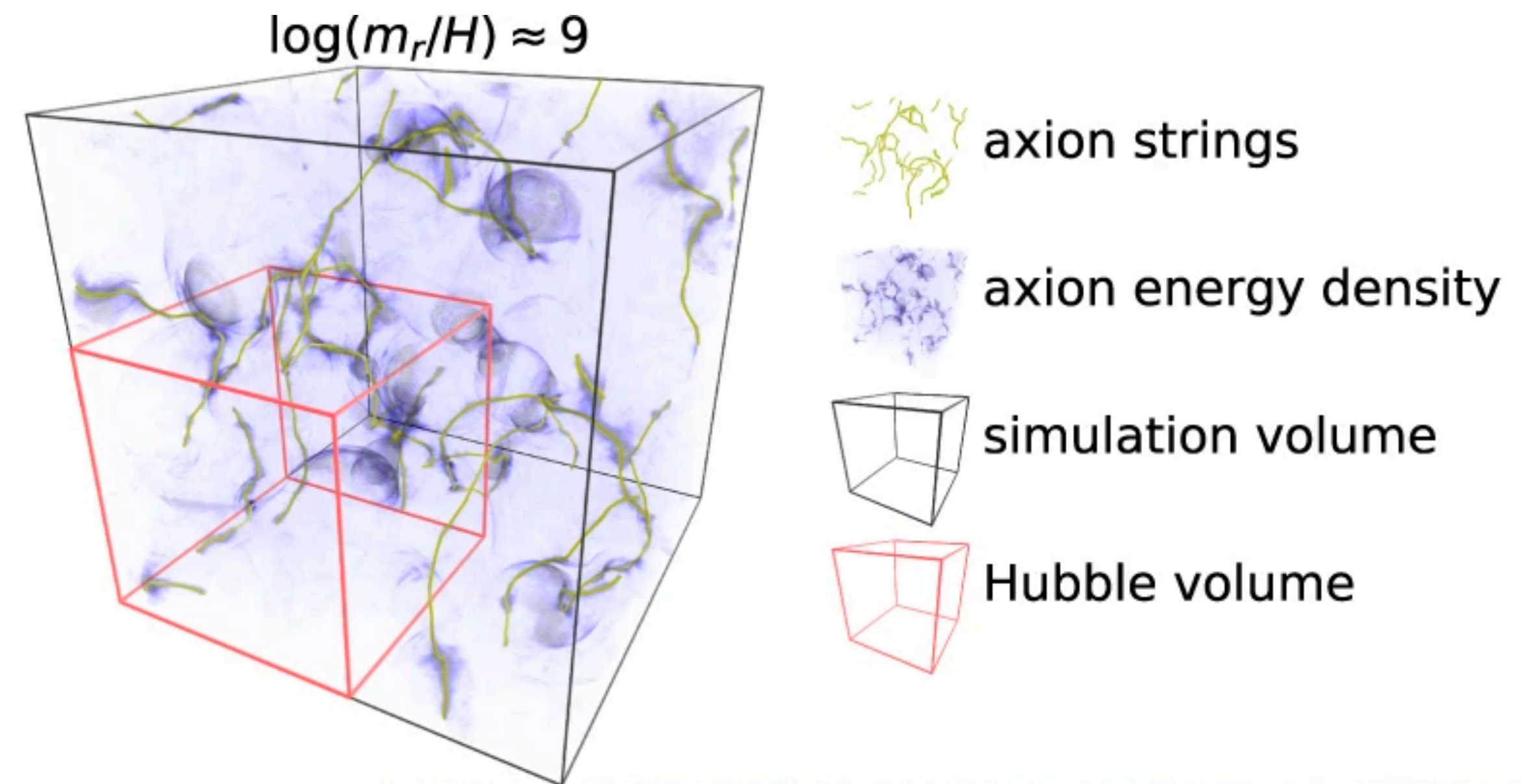
- CMB polarisation rotates by $\Delta\Phi = \pm \mathcal{A}\alpha_{\text{em}}$ if photon passes through a loop
- Observable: anisotropies in CMB polarisation rotation field



Simulations

Of axion string networks

- String dynamics leads to the same loop length distribution
- Most strings ($\sim 80\%$) are Hubble or super-Hubble scale
- The rest are logarithmically distributed sub-Hubble scale



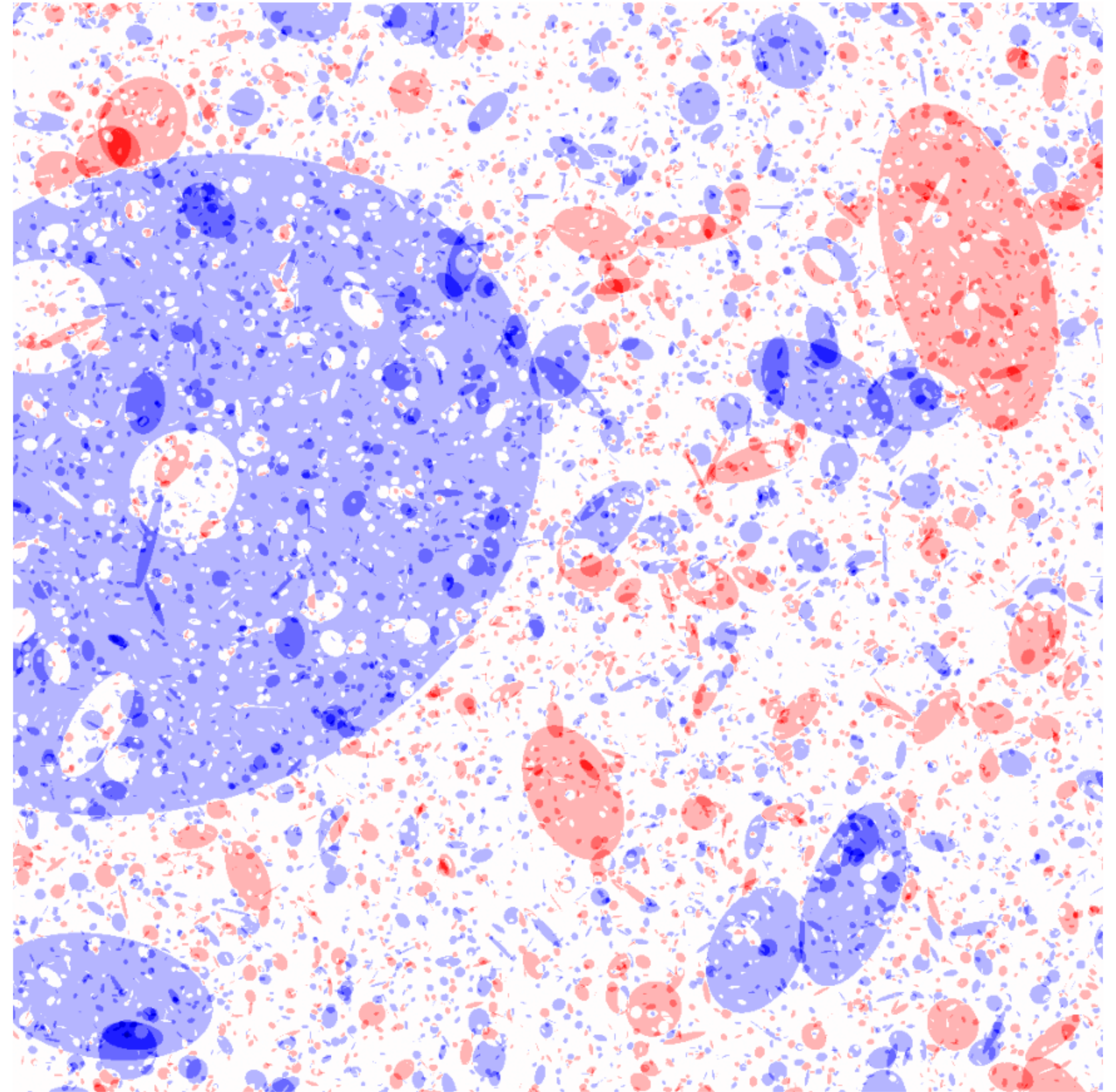
M. Buschmann, et al. (2022)

Loop-crossing model

A phenomenological string network model

- M. Jain, A. J. Long, M. A. Amin (2021)
- Circular string loops in random orientations scattered throughout the universe
- Loop radius distribution specified for one redshift then evolves via scaling law
- Each loop "paints" an ellipse on polarisation rotation field filled with $\pm \mathcal{A} \alpha_{\text{em}}$

Simulated polarisation rotation field (red+, blue-)



Parameters

Of the loop-crossing model

- $\mathcal{A} = 0.1 \sim 1$
Overall scaling of string-induced CMB polarisation rotation signal
- $\xi_0 = 1 \sim 100$
Effective number of (Hubble-scale) string loops per Hubble volume
- Other parameters that control loop radius distribution at any given redshift

Cosmic birefringence

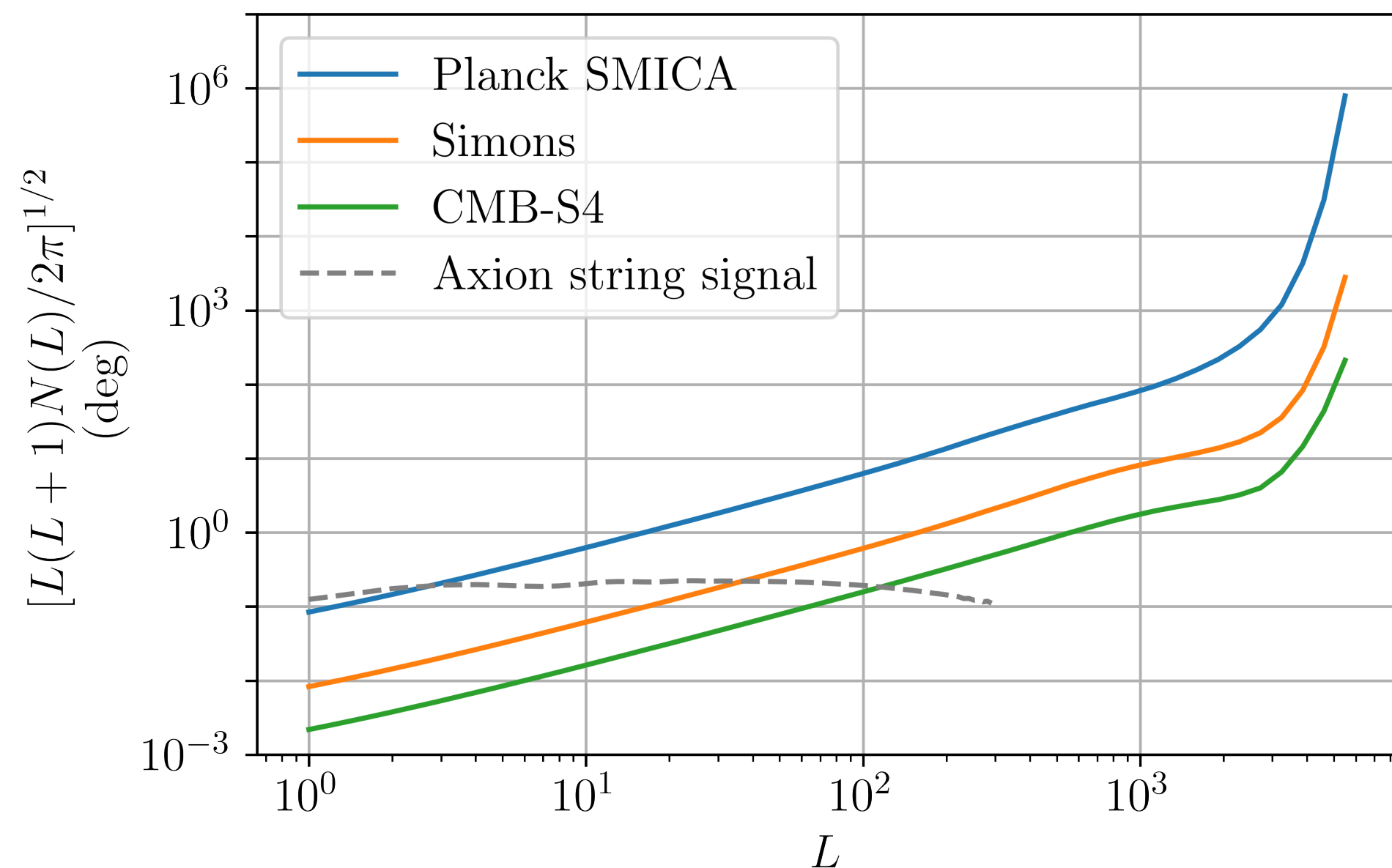
Quadratic estimators

- CMB polarisation rotation field must be estimated from cross-correlations between primary CMB observables T , E , and B
- Quadratic estimators (QE), well established for weak lensing, can be applied
- Lensing potential and polarisation rotation field can be simultaneously estimated via QEs
- Cannot resolve individual strings, need statistical detection of many strings
- Power spectrum of QEs well understood, use as summary statistics

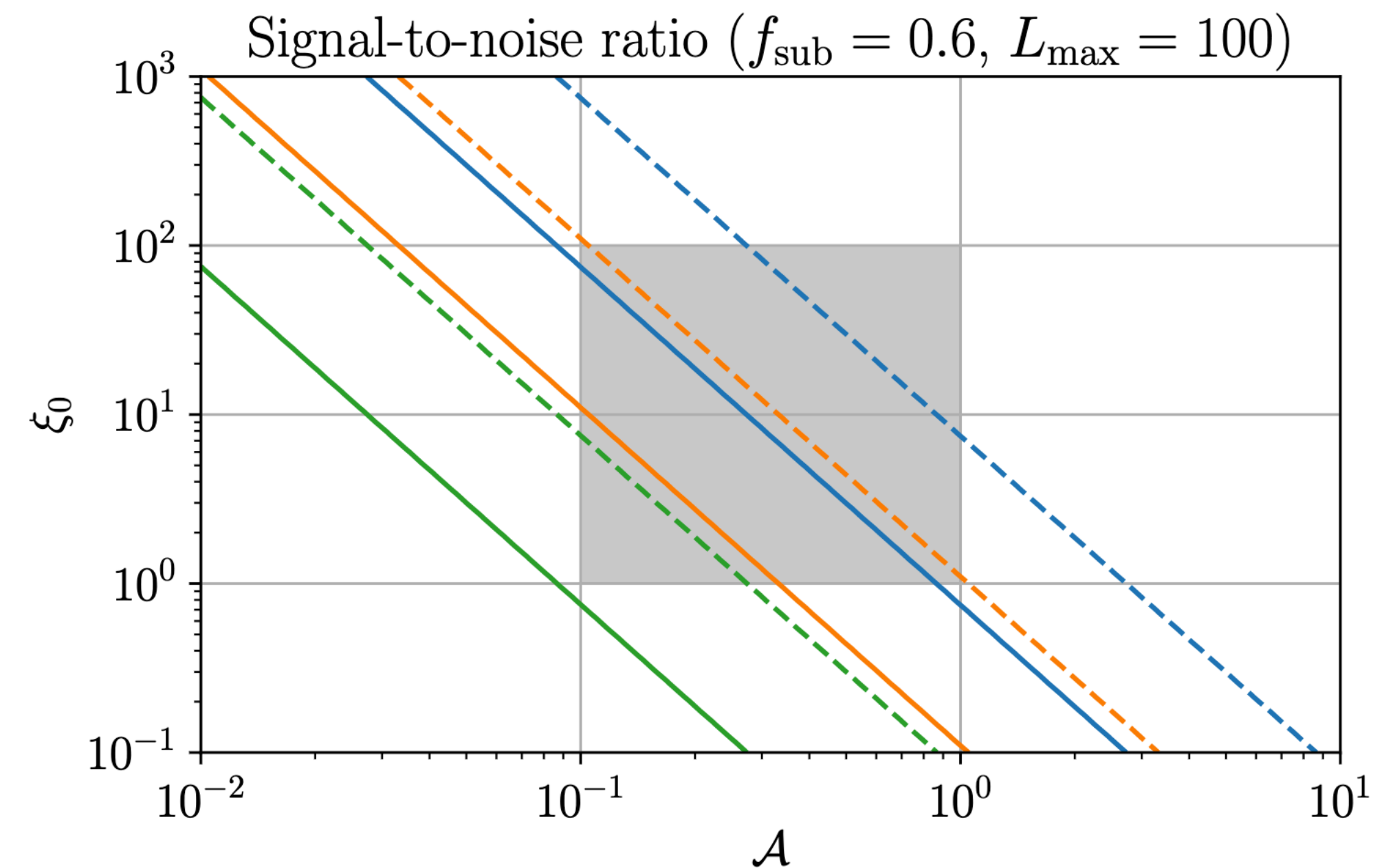
QE sensitivity

To axion string signal

- Signal is dominated by $L \lesssim 100$ modes in the rotation field power spectrum
- CMB Stage III, IV will discover or falsify axion string-induced anisotropic polarisation rotation



W. W. Yin, L. Dai, S. Ferraro (2022)



W. W. Yin, L. Dai, S. Ferraro (2022)

Planck constraint

- For the loop-crossing model in which a fraction of string loops have logarithmically distributed sub-Hubble sizes
- Planck 2015 data gives constraint $\mathcal{A}^2 \xi_0 < 0.93$ at 95% confidence
- Consistent with absence of axion strings

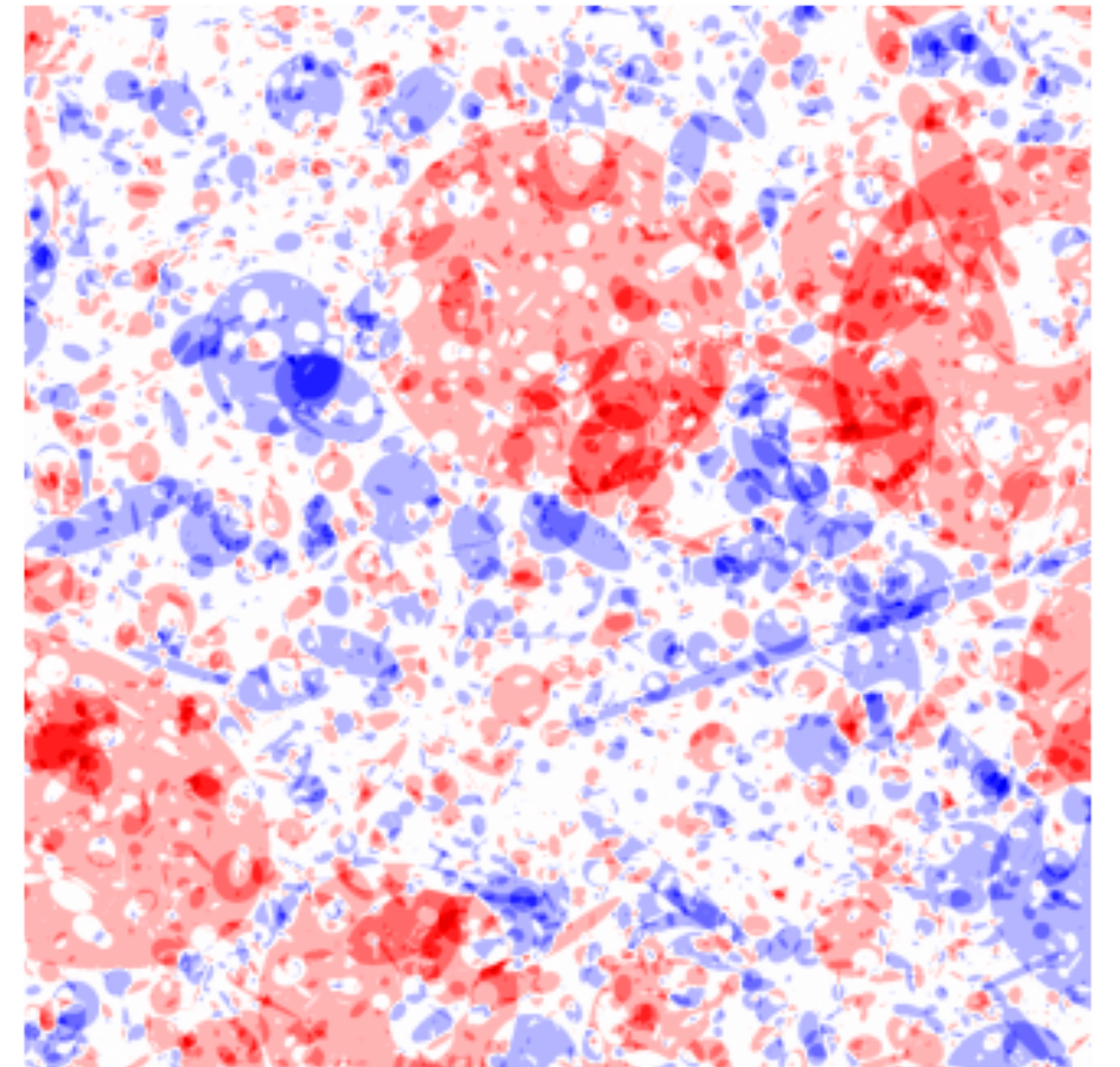
Beyond the power spectrum

Limitations

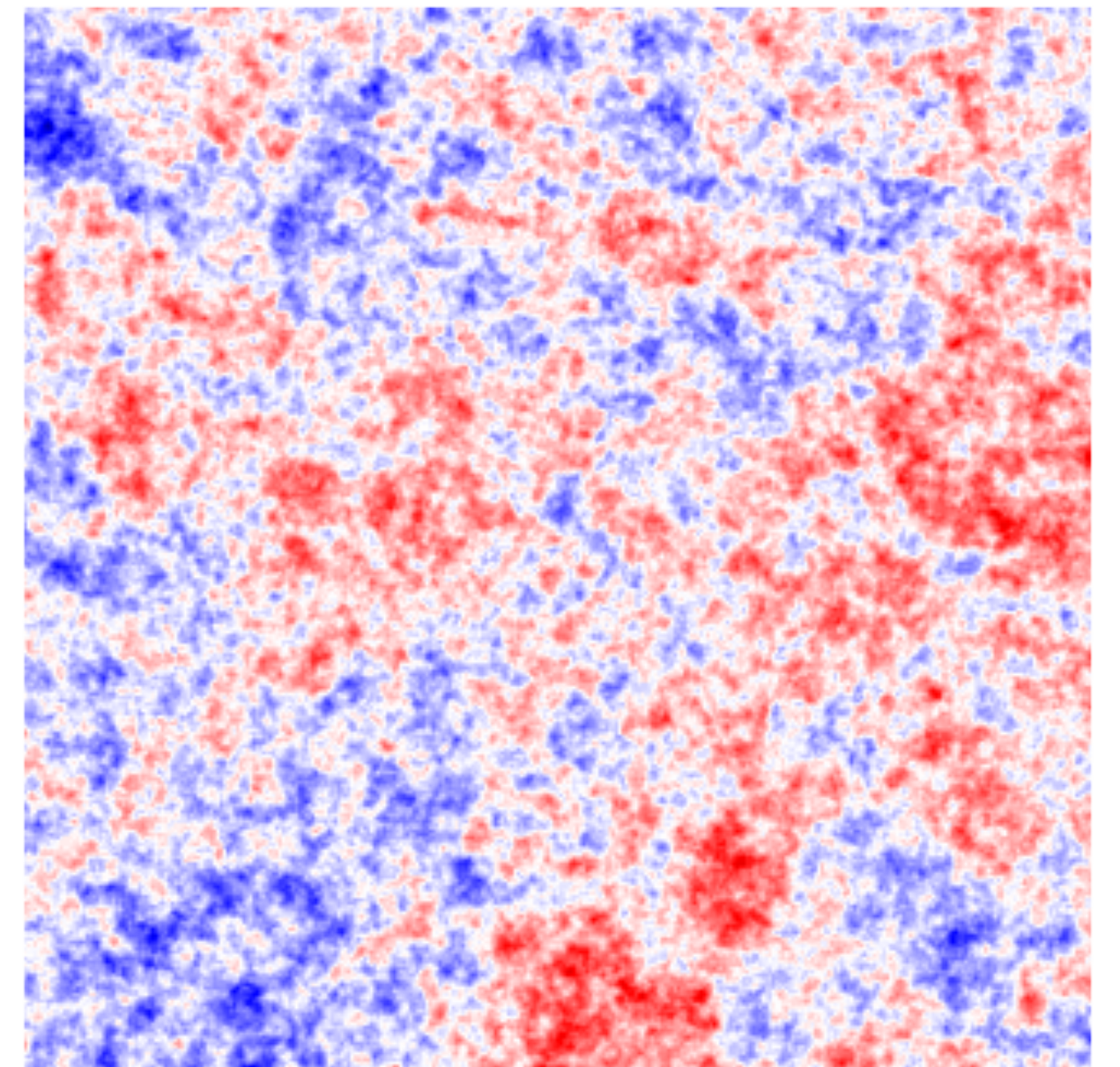
Of power spectrum

- Only captures Gaussian information
- Unable to distinguish between string networks with the same $\mathcal{A}^2 \xi_0$ but different \mathcal{A}
- Need to go beyond power spectrum
 - Bi-/trispectrum? CNN? Scattering transform? 🎉

Few strong strings

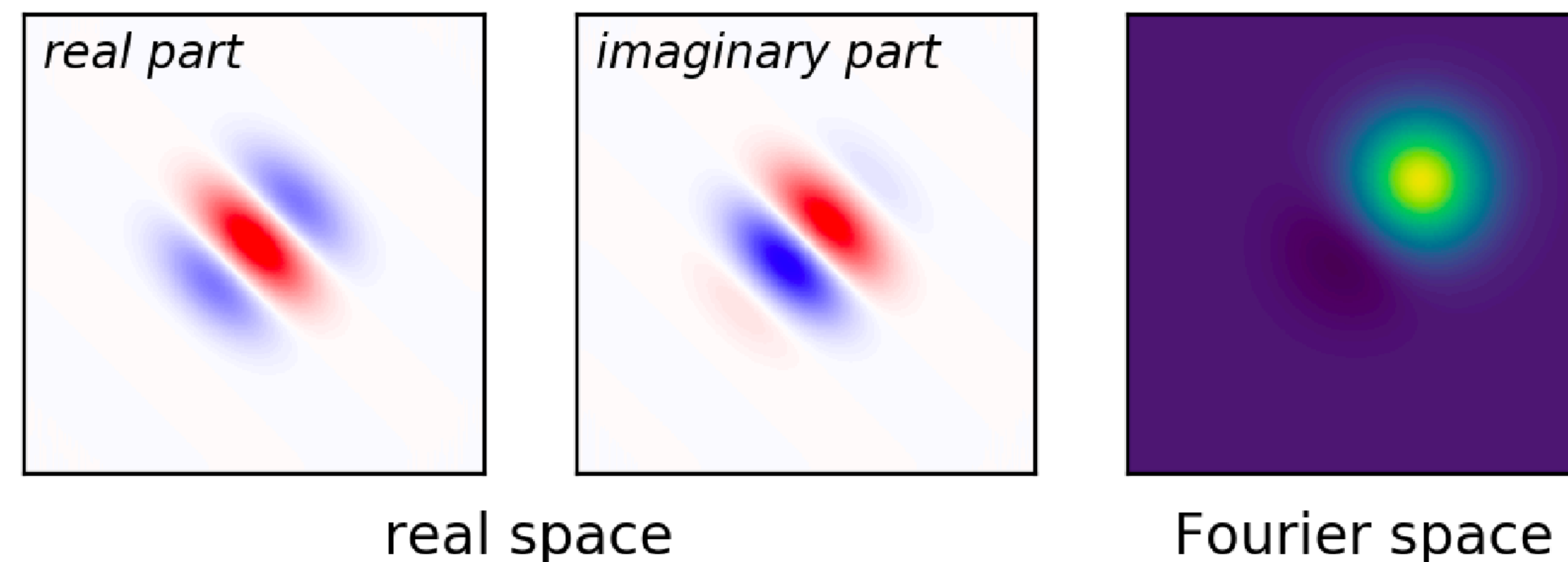


Many weak strings



Scattering transform

- Convolve input field with Morlet wavelets (localised in real and Fourier space) of different scales and orientations



S. Cheng, et al. (2020)

- Similar to convolutional neural network but requires no training
- Non-linear transform (absolute value) after convolution ensures non-Gaussian info is extracted
- Can reduce to translationally and rotationally symmetric summary statistics

Advantages

Of scattering transform (ST)

- Over bi-/trispectra
 - Packs non-Gaussian info in small number of coefficients: for 128^2 input field, only 21 ST coefficients
 - Higher-order ST does not suffer from the increased sample variance from $(\text{input field})^{\text{high power}}$
- Over convolutional neural network (CNN)
 - Requires no training
 - Reduced coefficients are inherently translationally and rotationally symmetric
 - Individual coefficients have interpretable meaning

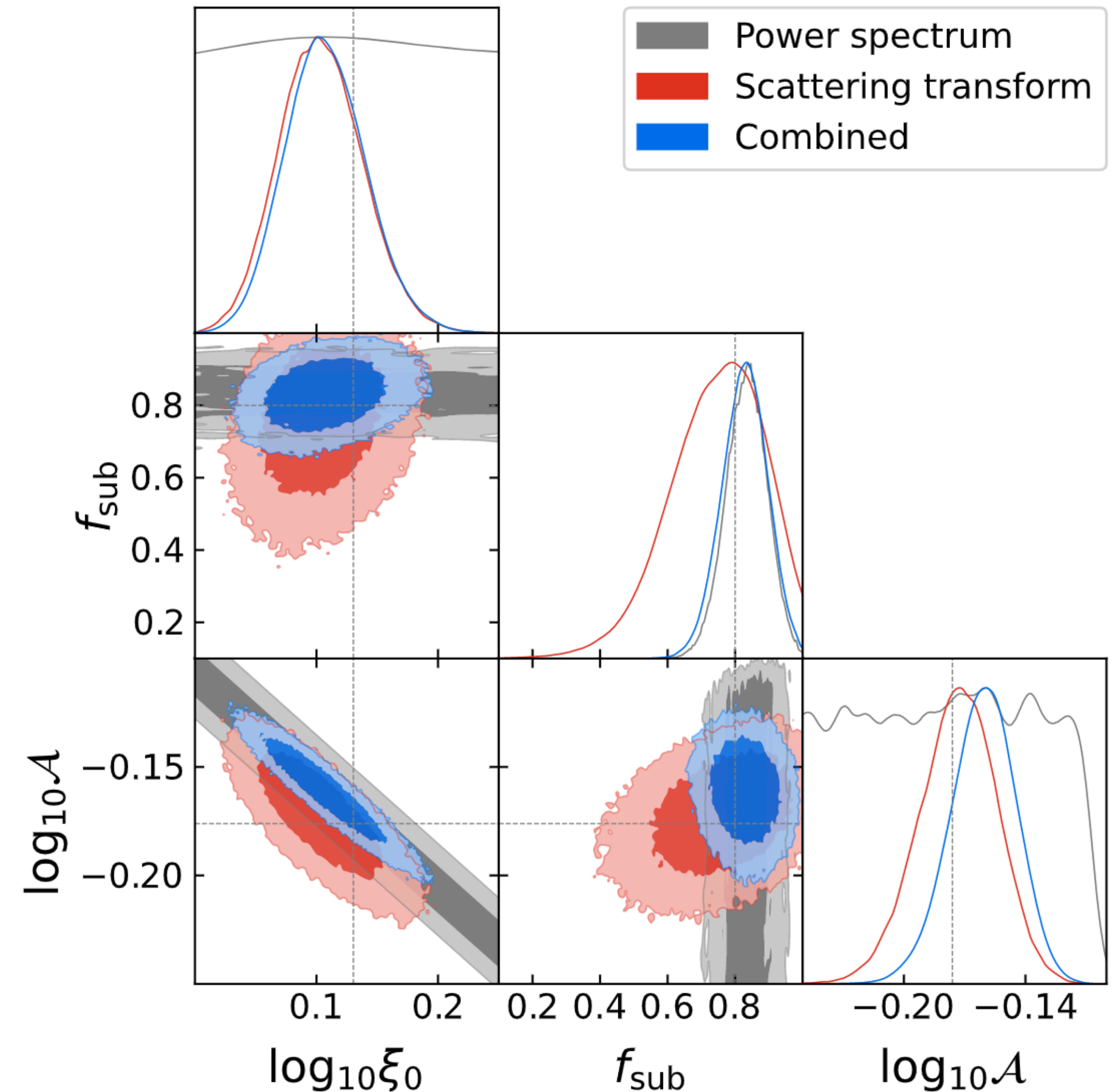
Parameter inference

Using scattering transform

- Generate a large number of polarisation rotation fields on the discretised parameter space of the loop-crossing model
- Compute their scattering transform coefficients (Kymatio Python package)
- Compute sample mean and covariance matrix at each parameter grid point
- Interpolate to obtain the "theory" against which actual CMB polarisation rotation field is compared
- Likelihood maximisation by MCMC

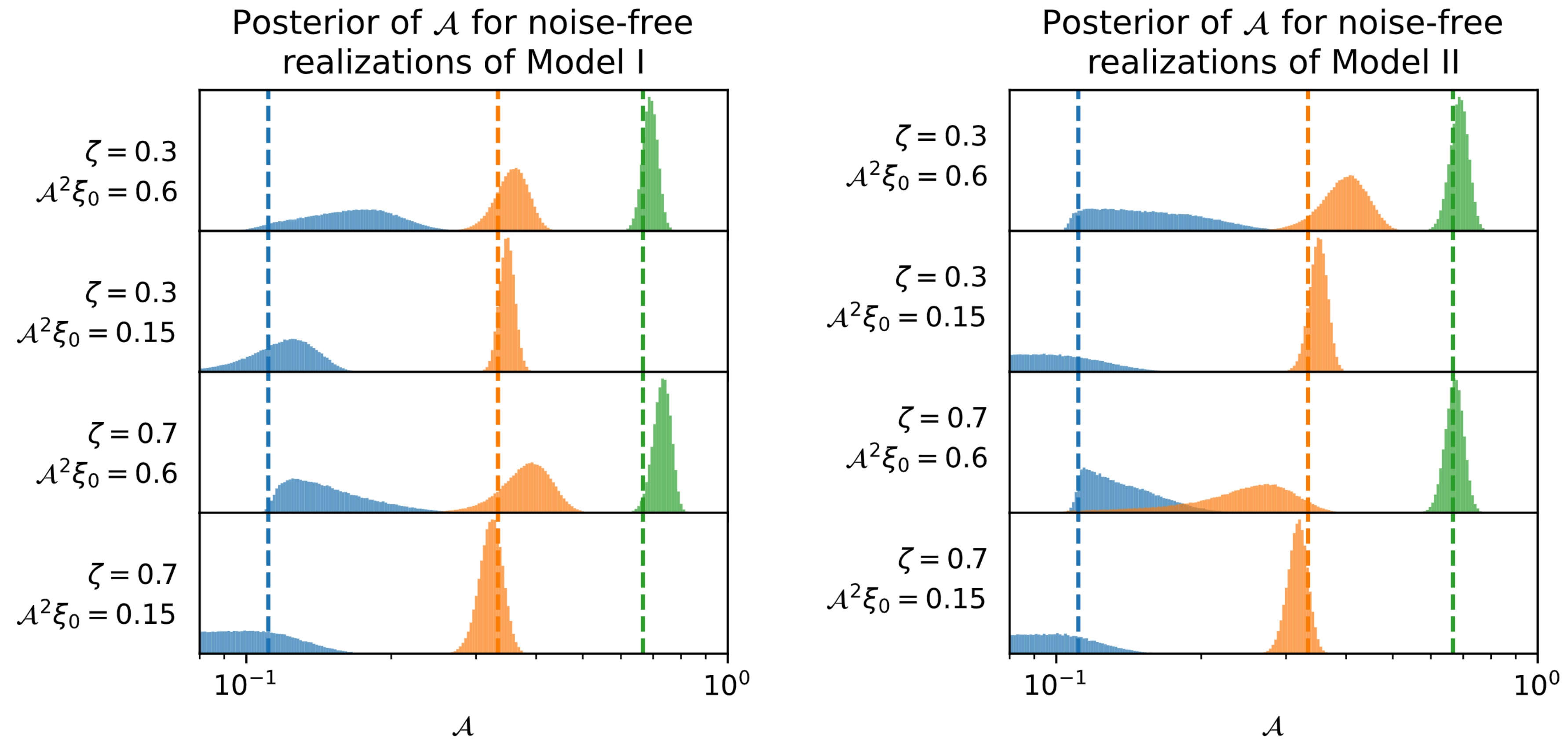
Evaluation

- Testing is done using mock polarisation rotation fields with known parameters as input fields
- Procedure repeated for both ideal noise-free case and QE reconstruction noise at future CMB-HD level
- Compared with power spectrum analysis



Ideal noise-free case

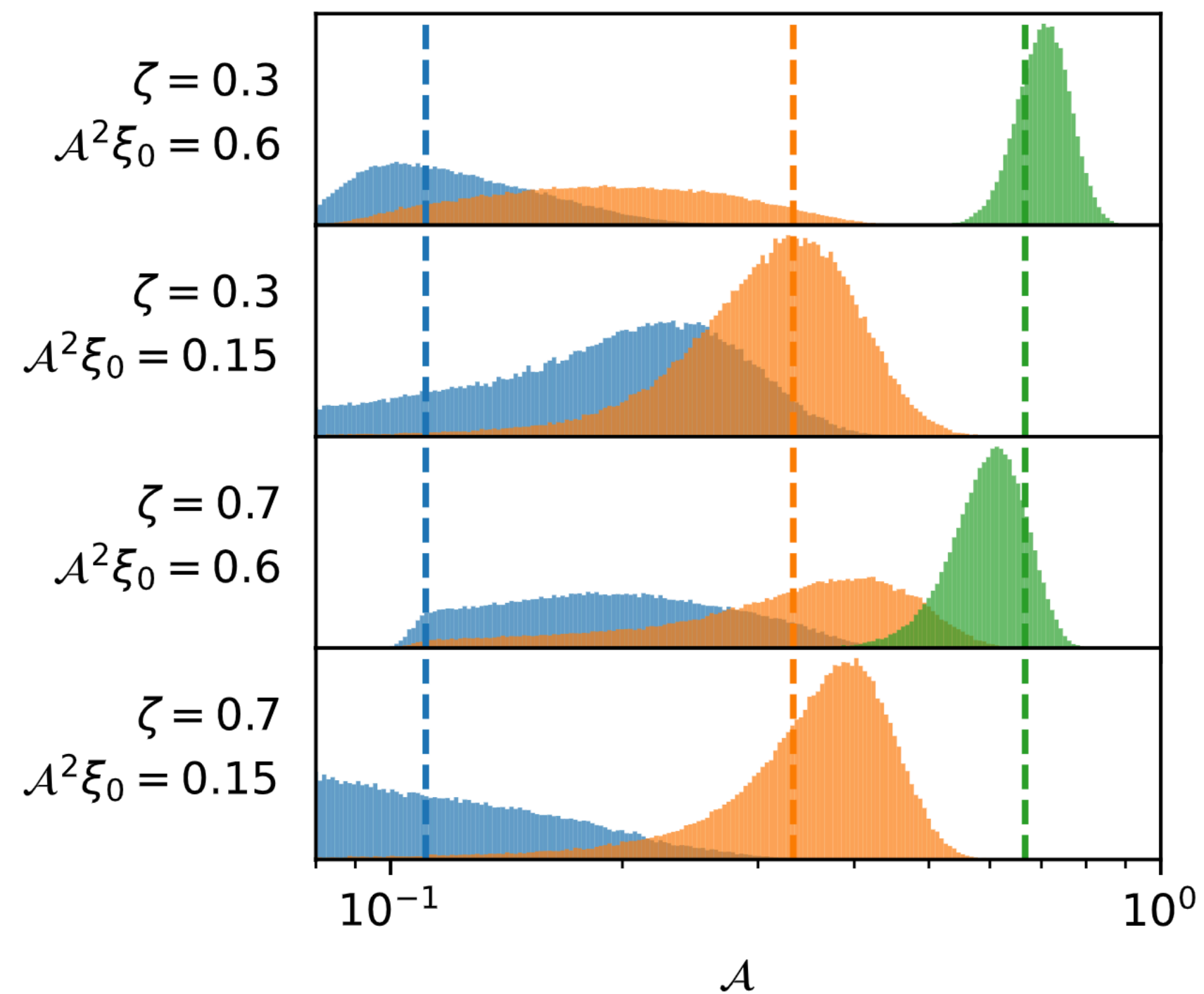
- Able to clearly distinguish between $\mathcal{A} = 1/9, 1/3, 2/3$



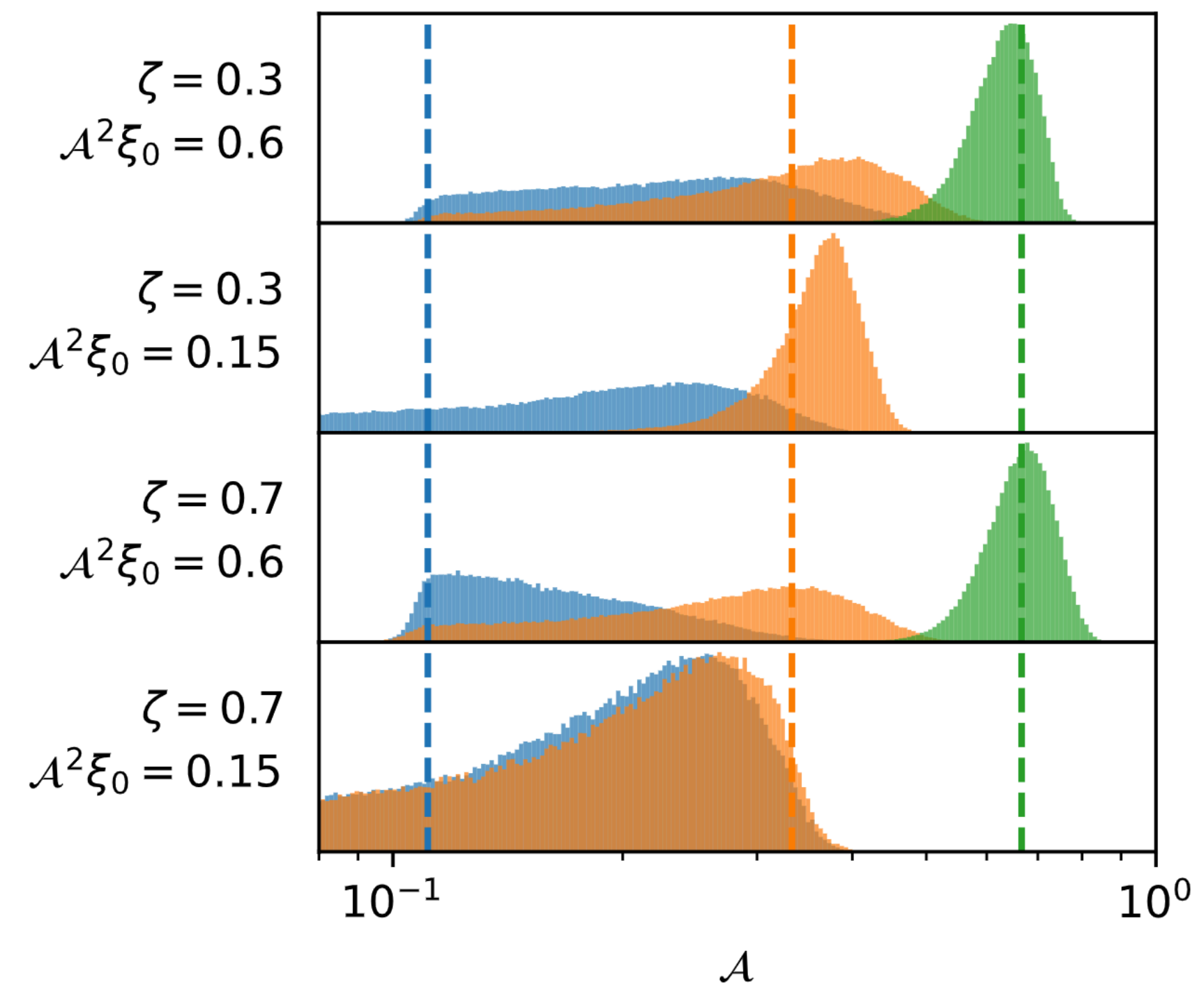
CMB-HD noise level

- Able to clearly distinguish $\mathcal{A} = 2/3$ from $\mathcal{A} = 1/9, 1/3$ and marginally between $\mathcal{A} = 1/9, 1/3$

Posterior of \mathcal{A} for realizations of Model I with CMBHD reconstruction noise



Posterior of \mathcal{A} for realizations of Model II with CMBHD reconstruction noise



Summary

- Axion-photon coupling is proportional to anomaly coefficient \mathcal{A}
- \mathcal{A} reveals charge assignments beyond the SM
- Axion strings induce quantised anisotropic rotation of CMB polarisation $\propto \mathcal{A}$
- CMB Stage III, IV will give us a conclusive answer on axion strings through power spectrum of QE
- Power spectrum analysis suffers from $\mathcal{A}^2 \xi_0$ degeneracy
- If axion strings are discovered, scattering transform can measure \mathcal{A} (at CMB-HD noise level) and rule out certain beyond-the-SM theories

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