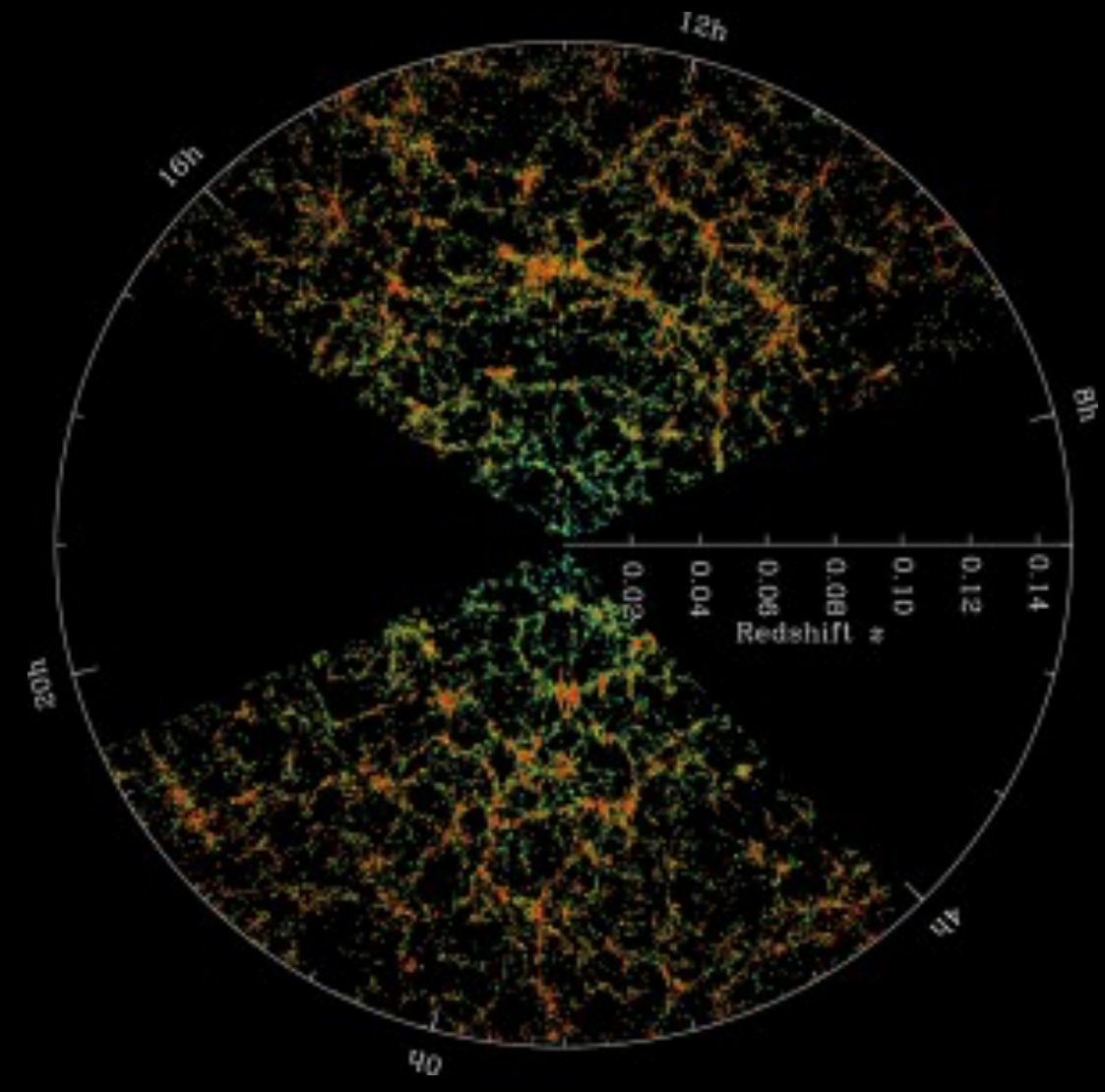
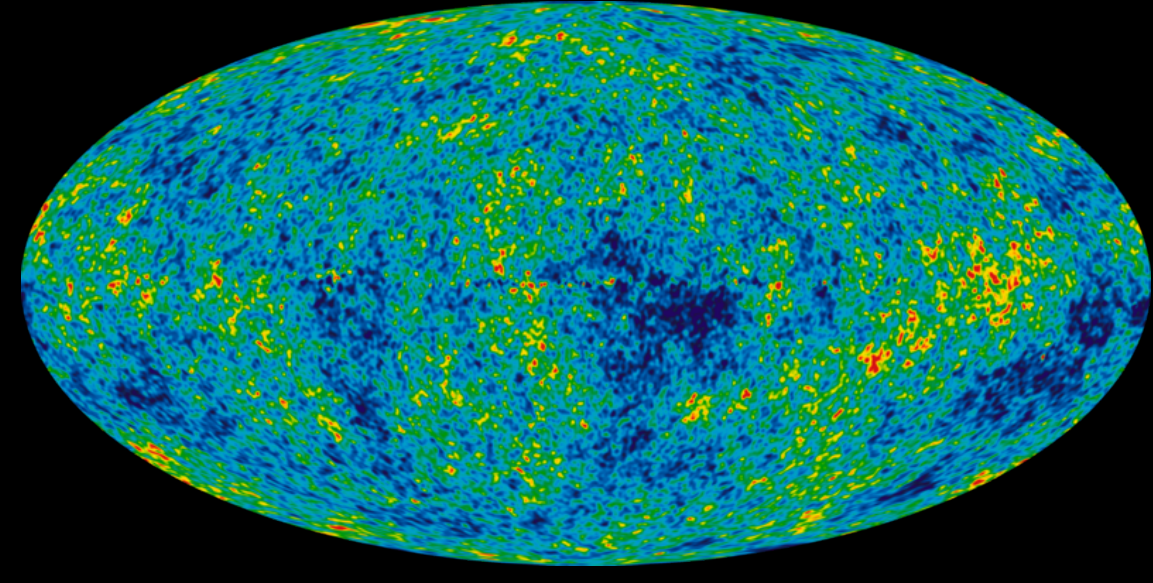


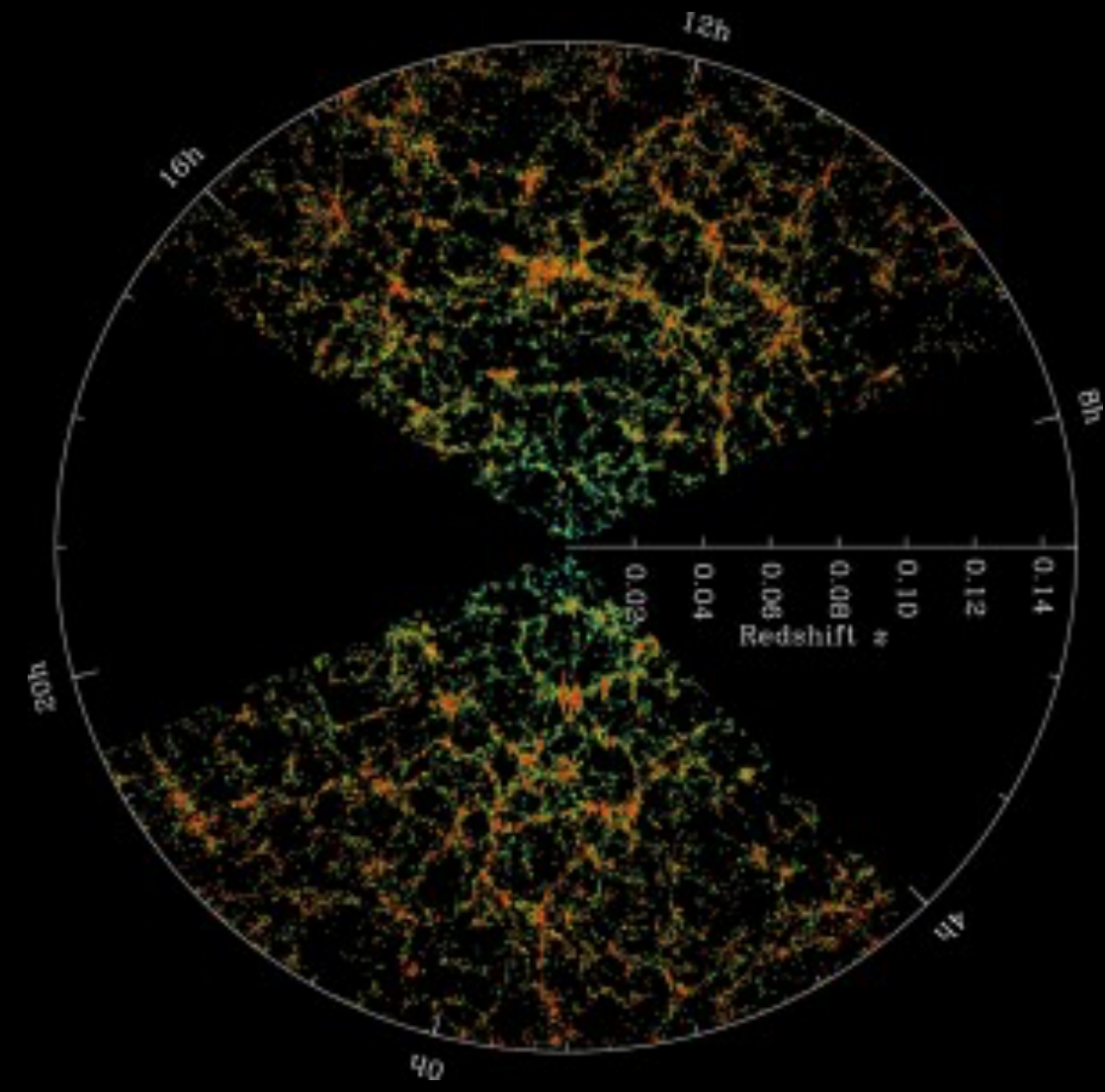
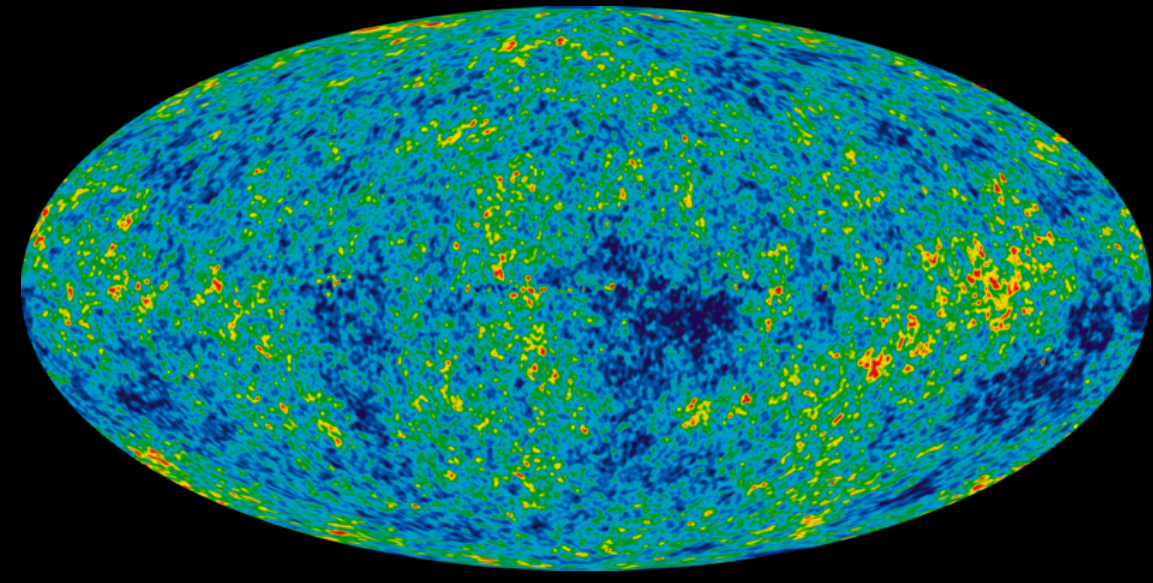
Seeing highly anisotropic gravitational wave backgrounds from phase transitions

Arushi Bodas

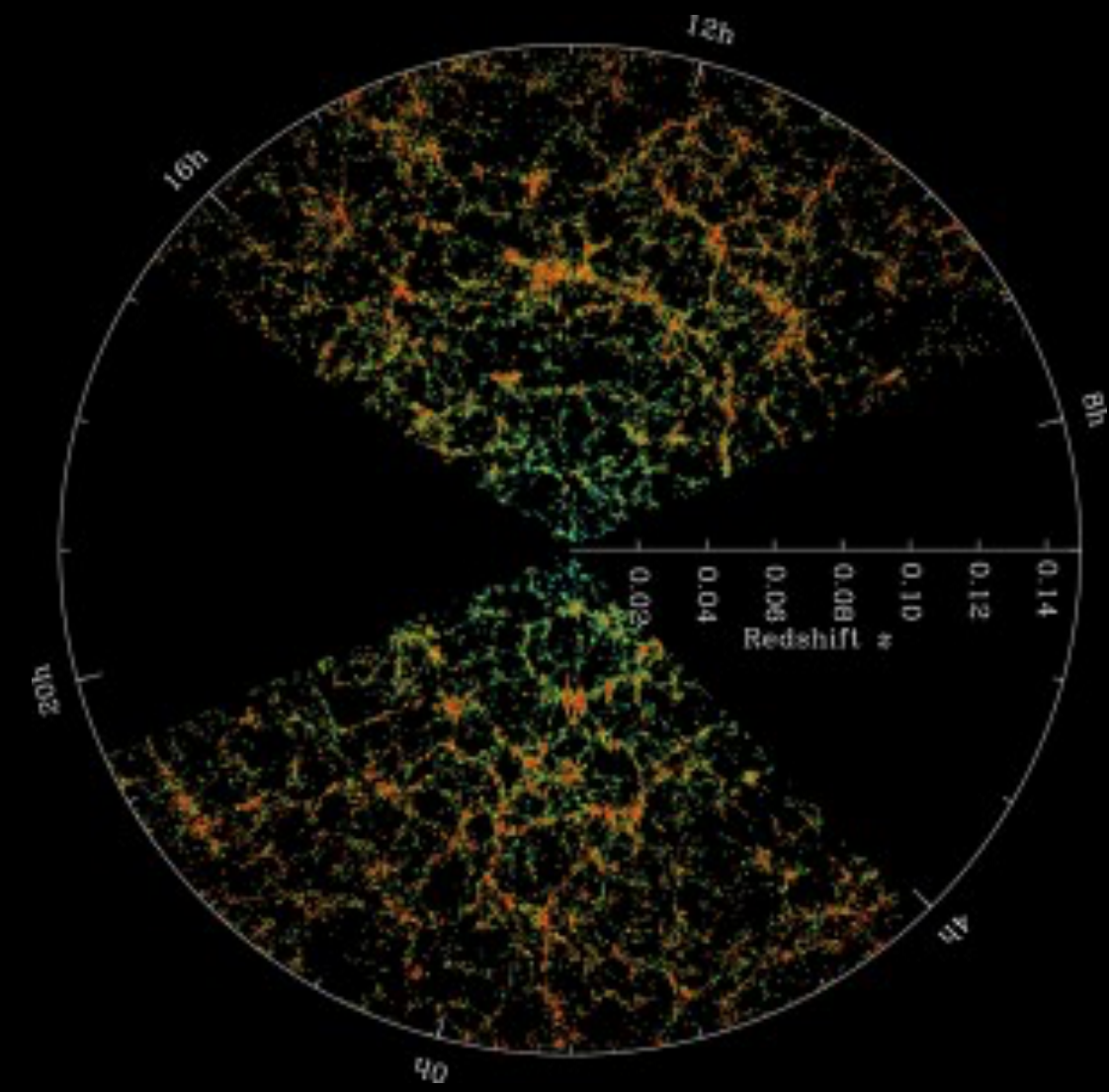
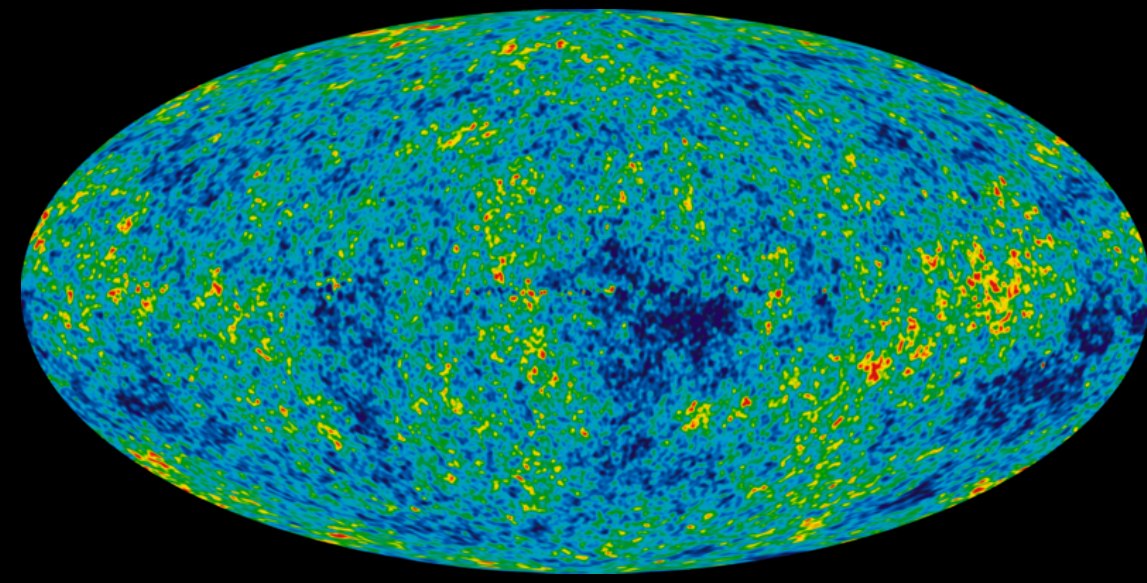
University of Chicago and Fermilab

Based on work with R. Sundrum at University of Maryland, *JHEP* 06 (2023) 029





Quantum fluctuations of a scalar field + inflation



Quantum fluctuations of a scalar field + inflation

Anisotropy maps are windows into inflationary physics

Adiabatic nature of known datasets

- Current observations tell us that density perturbations in all Standard Model (SM) species and Dark Matter (DM) originate from the same inflationary fluctuation (e.g. the inflaton).
[Planck,1807.06211](#)
- Future datasets based SM and DM species (line-intensity maps, $C_{\nu B}$) will also be adiabatic.
- Detection of new inflationary physics in these future maps is already constrained from non-detection in CMB.
- Could there be cosmological anisotropy maps very different from the CMB (dominantly isocurvature)?

Something that is...

- Copiously produced in the early universe
- Free-streaming (does not thermalize with SM and DM)
- Could be detectable with the technology that we have today

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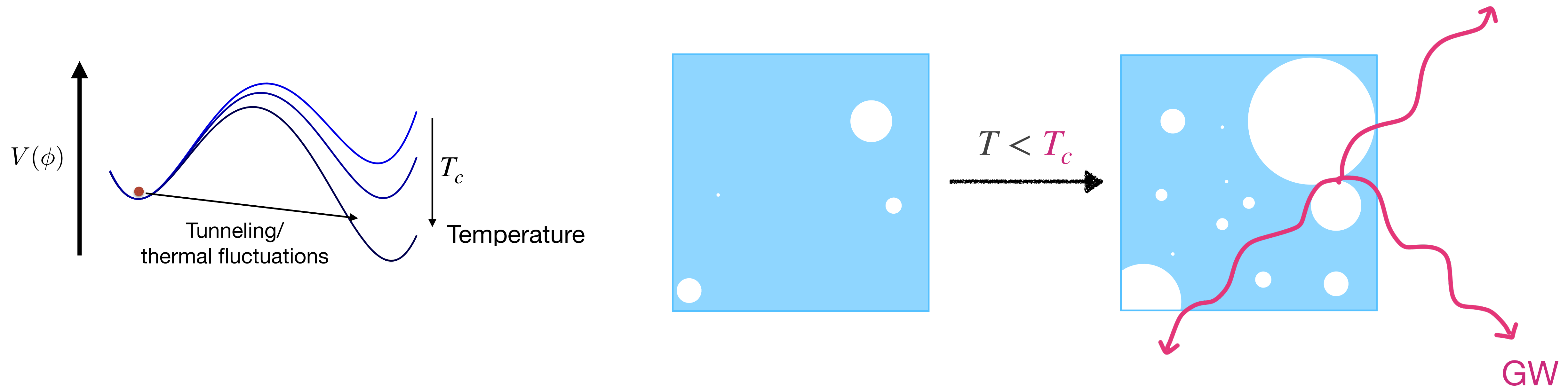
Gravitational waves!

Outline

- Background: Source and anisotropy of gravitational wave background (GWB)
- A simple model that can produce isocurvature GWB
- **Tradeoff**: large isocurvature leads to a suppression of the GWB signal
- Our model: weakens this tradeoff, thereby improving detection prospects of highly isocurvature GWB

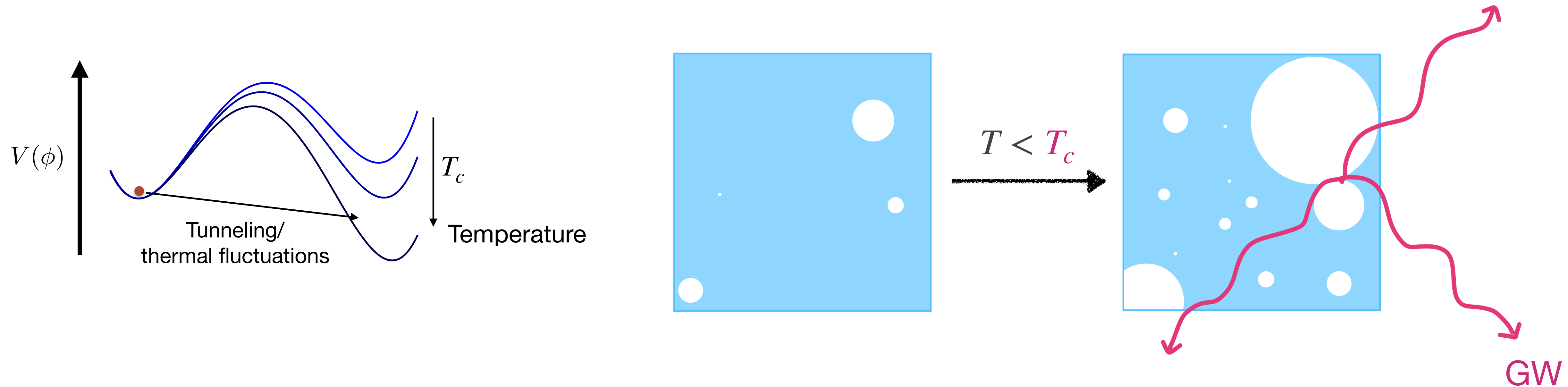
Gravitational waves from the early universe

- We focus on first-order phase transitions (PT) as sources of GW.



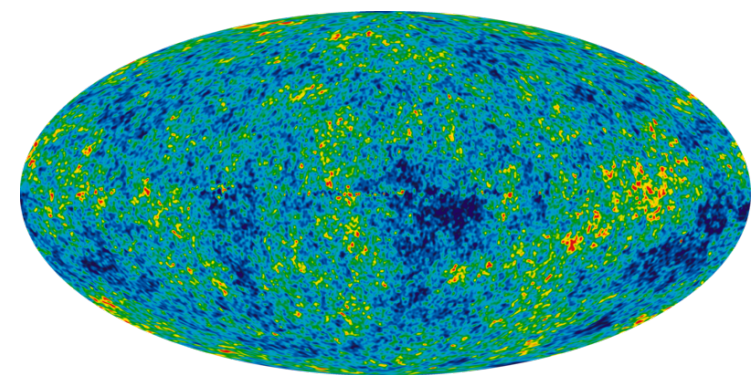
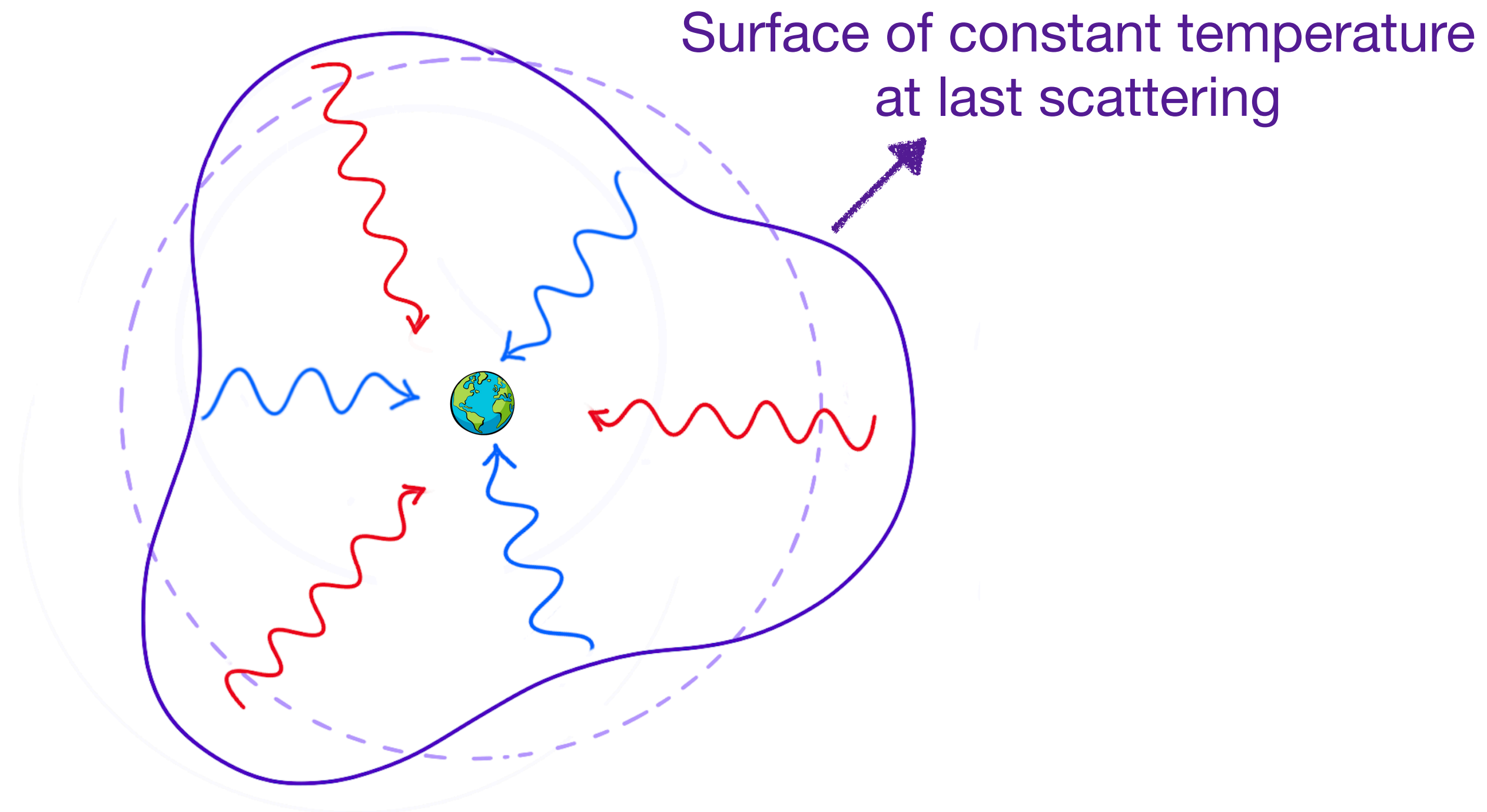
Gravitational waves from the early universe

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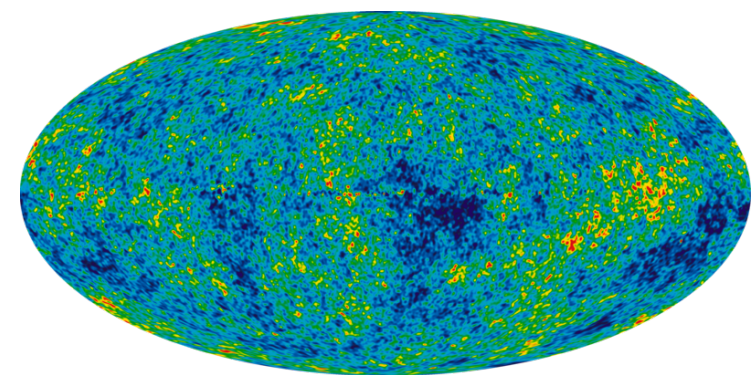
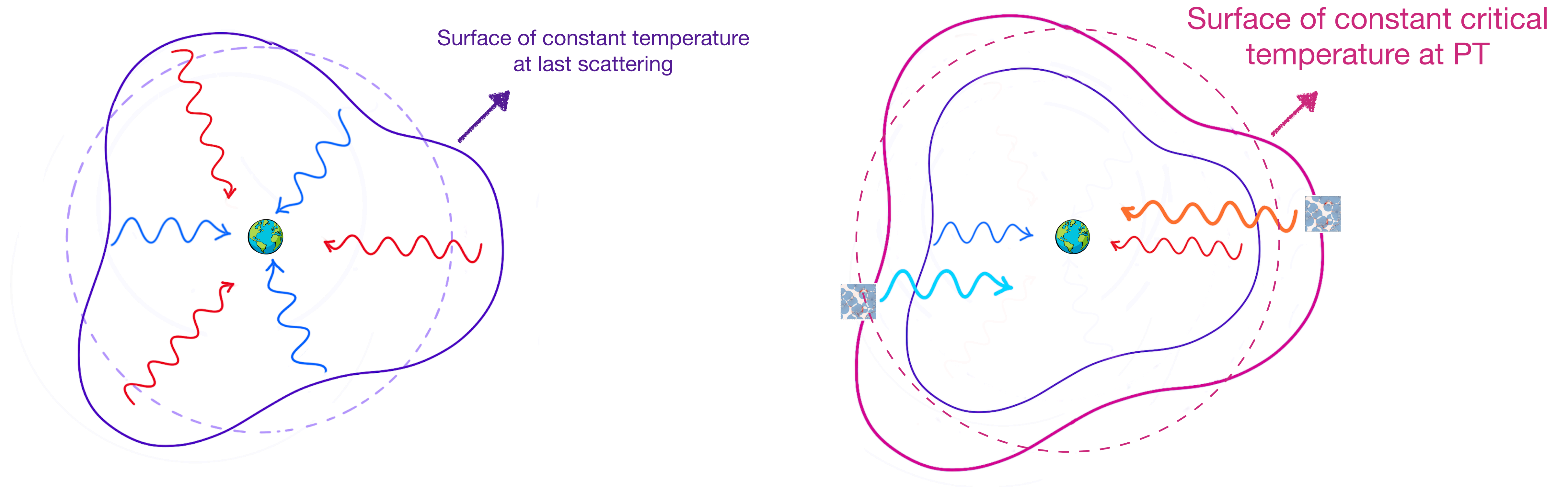
- First order phase transitions are dramatic.
- Phase transitions are “instantaneous”, giving a 2D anisotropy map (analogous to CMB)

Recap: CMB anisotropies

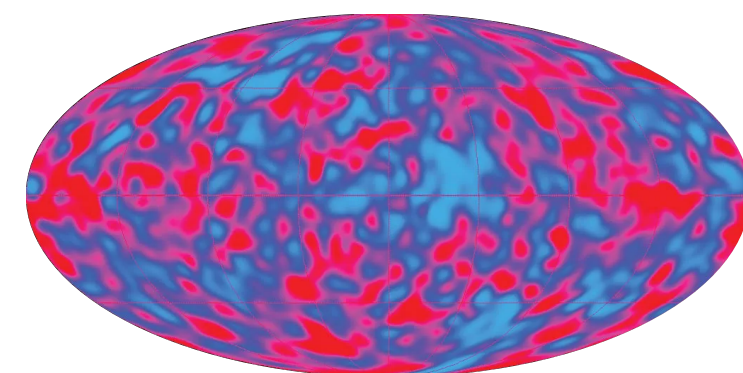


$$\delta_{\text{CMB}} \equiv \frac{\Delta T}{\bar{T}} \sim 10^{-5}$$

GWB anisotropies

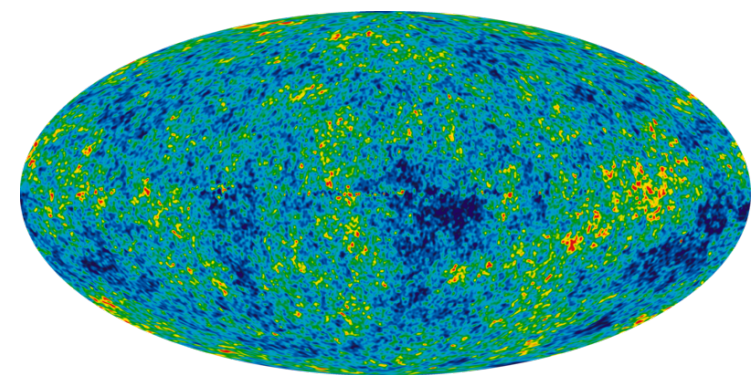
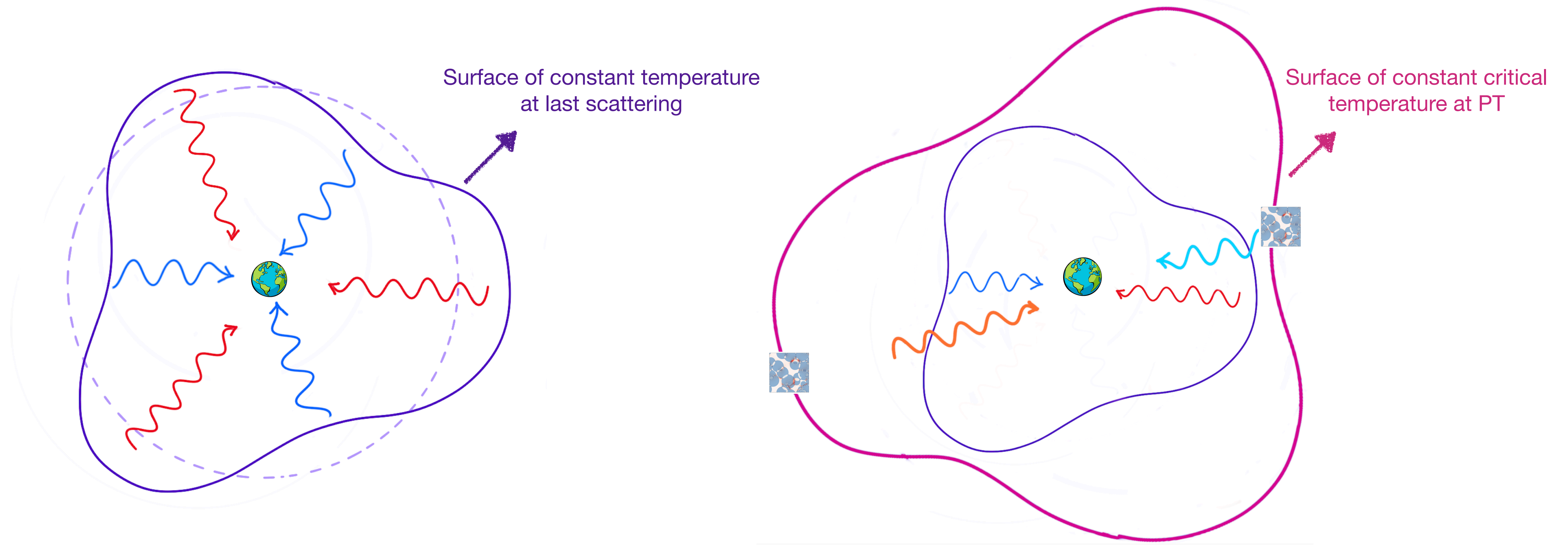


$$\delta_{\text{CMB}} \equiv \frac{\Delta T}{\bar{T}} \sim 10^{-5}$$

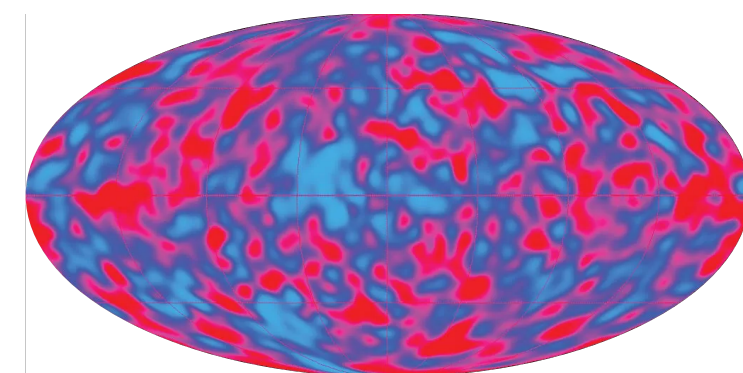


$$\delta_{\text{GW}} \equiv \frac{\Delta \rho_{\text{GW}}}{\bar{\rho}_{\text{GW}}} \sim 10^{-5}$$

Isocurvature GWB anisotropy

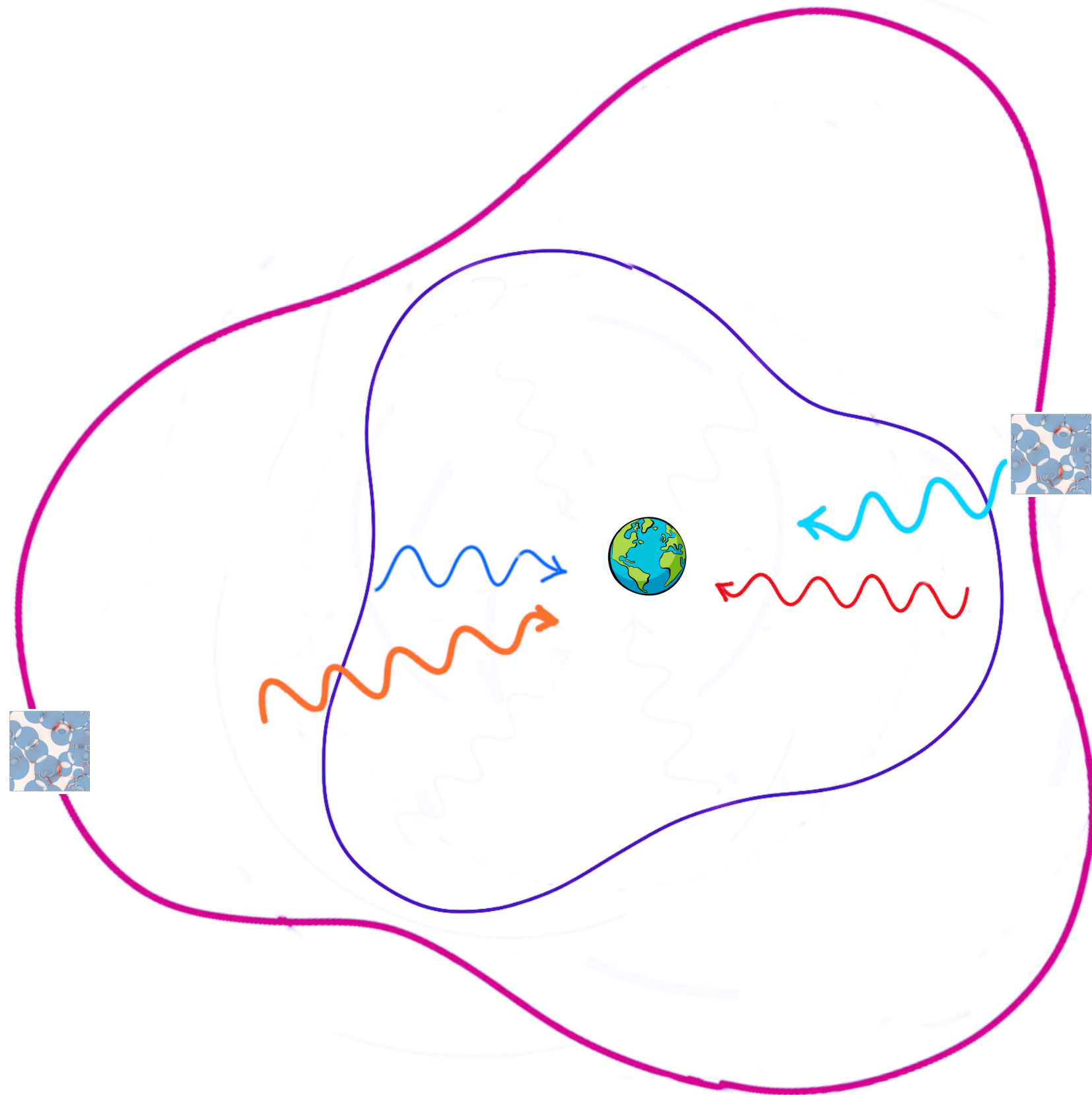


$$\delta_{\text{CMB}} \equiv \frac{\Delta T}{\bar{T}} \sim 10^{-5}$$



$$\delta_{\text{GW}} \equiv \frac{\Delta \rho_{\text{GW}}}{\bar{\rho}_{\text{GW}}} \neq 10^{-5}$$

Isocurvature GWB anisotropy



Uncorrelated large-scale fluctuations → another **light field** from inflation

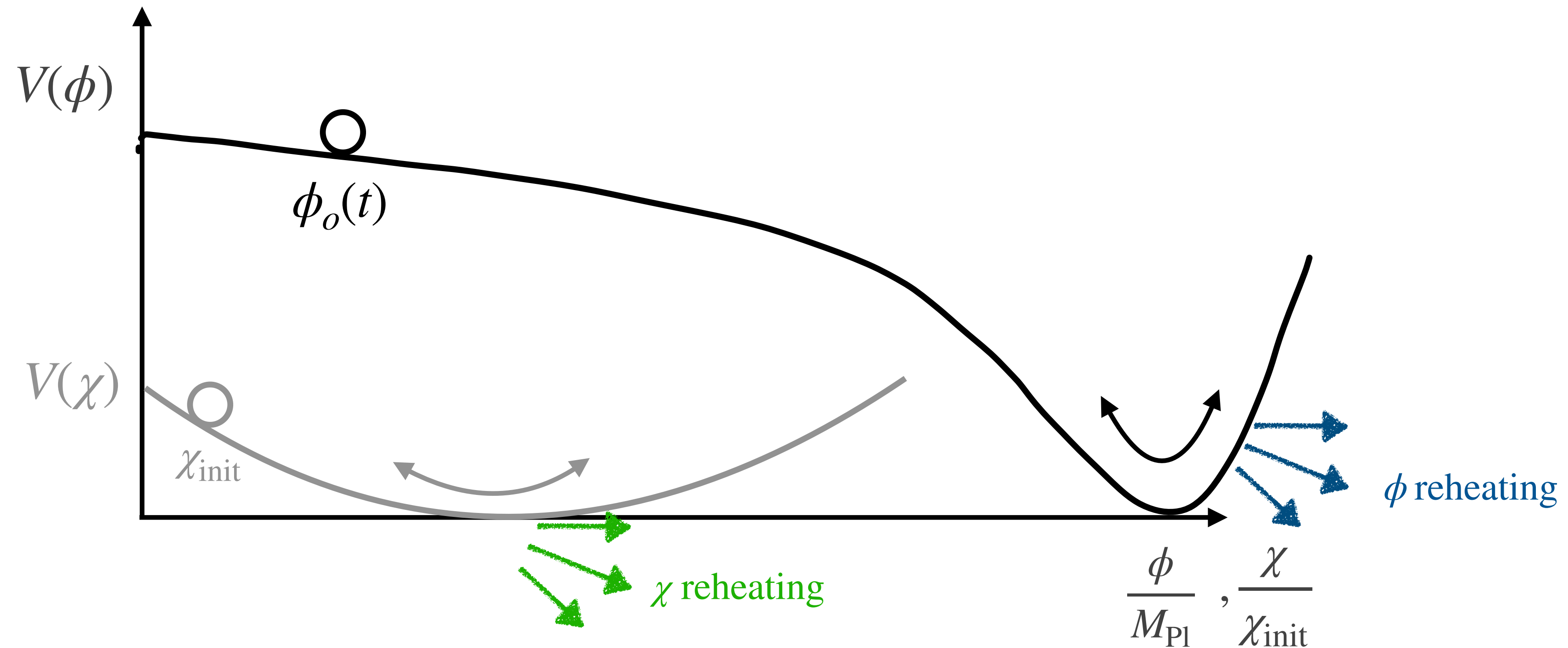
Interesting features such as tilt, scale-invariance breaking feature, non-Gaussianity → **new signatures** of physics from inflation

L. Valbusa Dall'Armi et al 2021

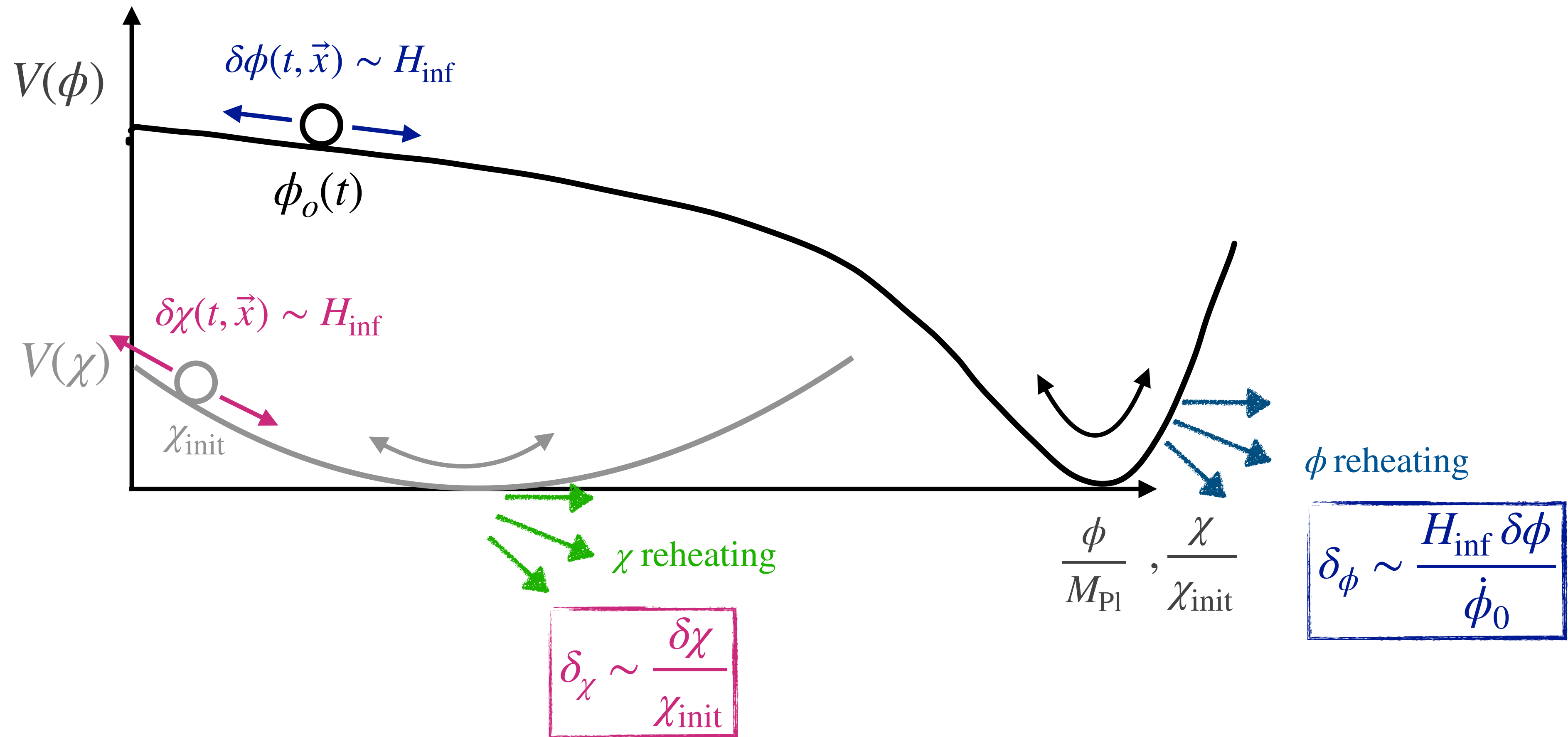
S. Kumar, R. Sundrum and Y. Tsai 2021

AB, R. Sundrum 2022

A new light field from inflation: ALP (χ)

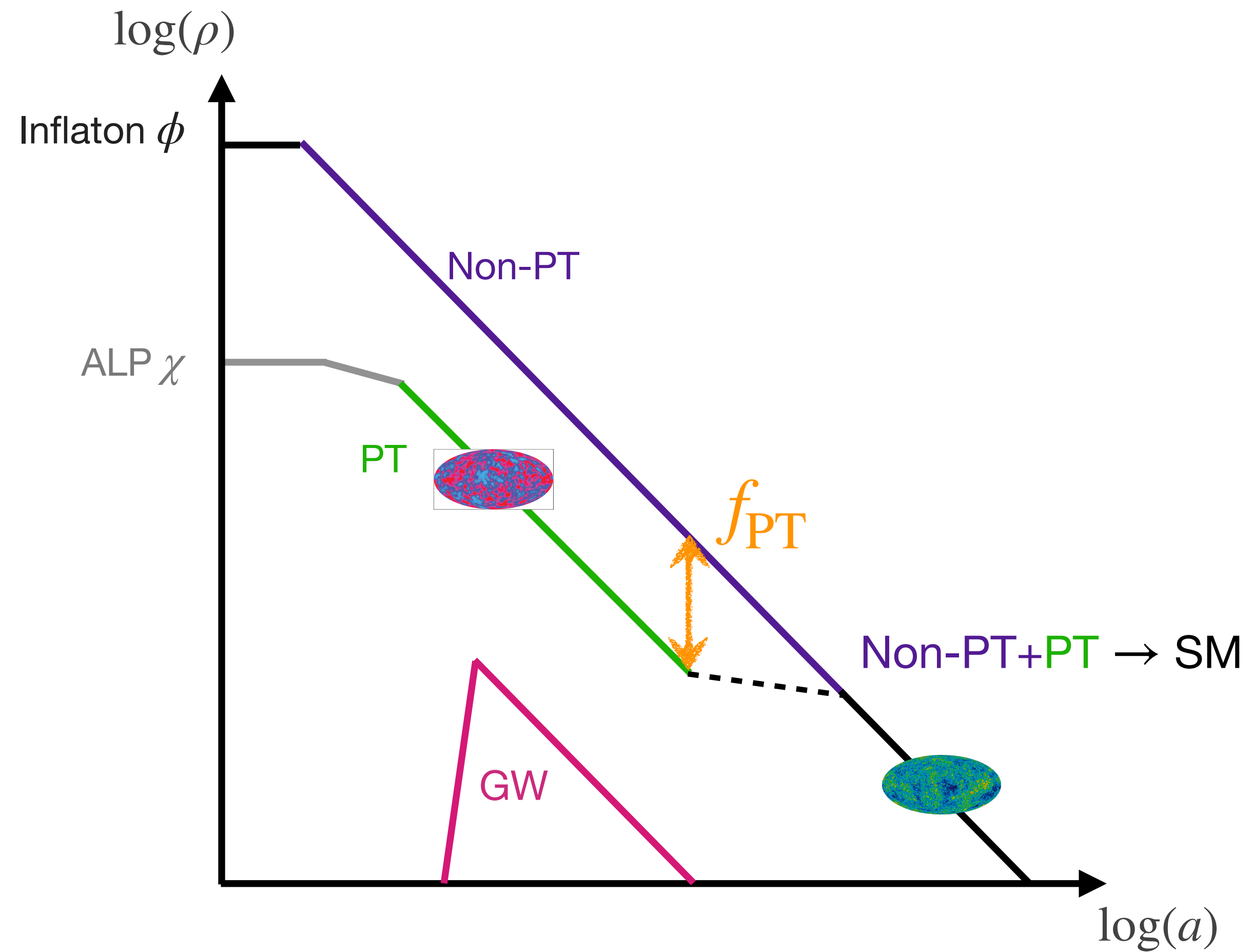


A new light field from inflation: ALP (χ)



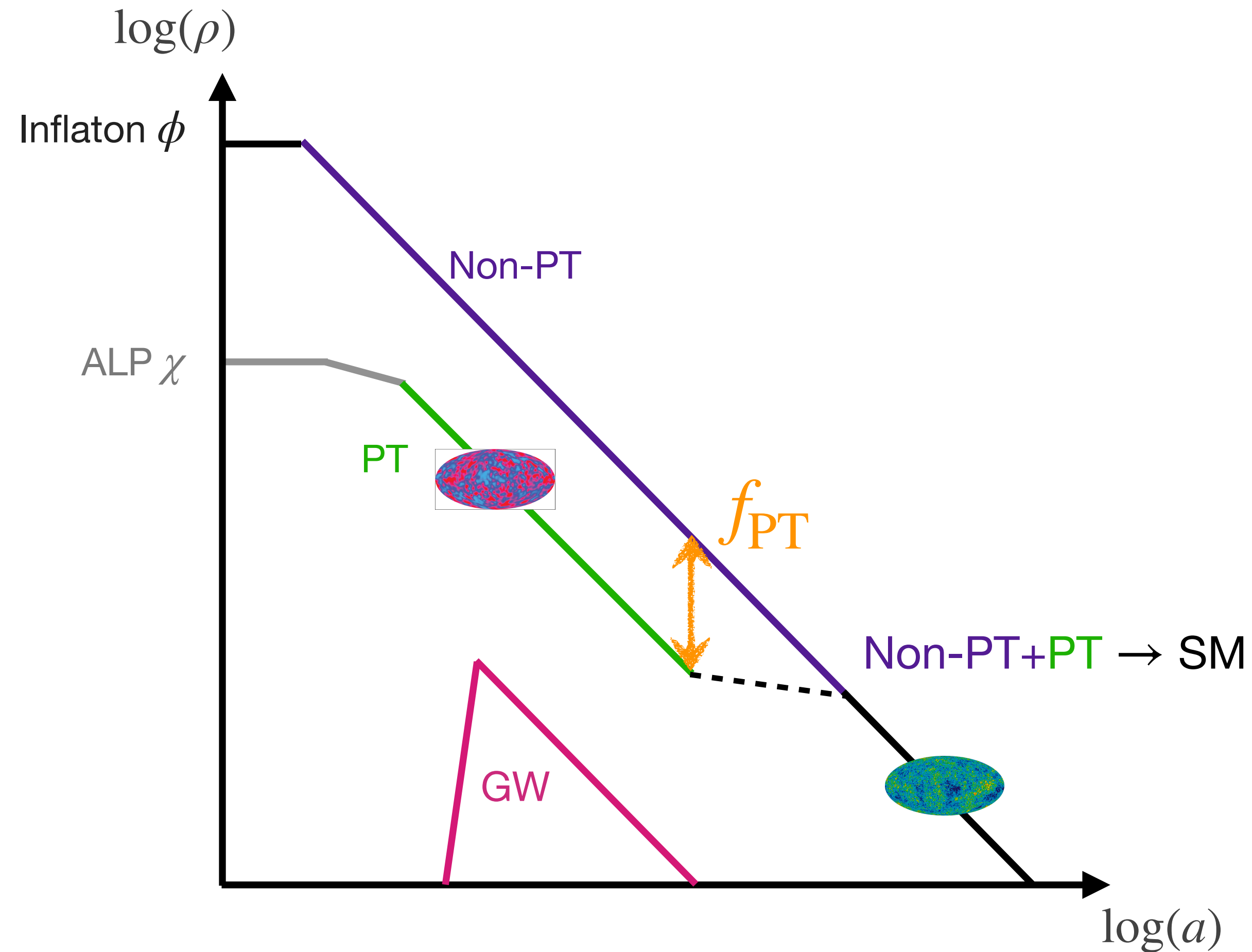
A model of isocurvature GWB

Geller, Hook, Sundrum, Tsai 1803.10780



A model of isocurvature GWB

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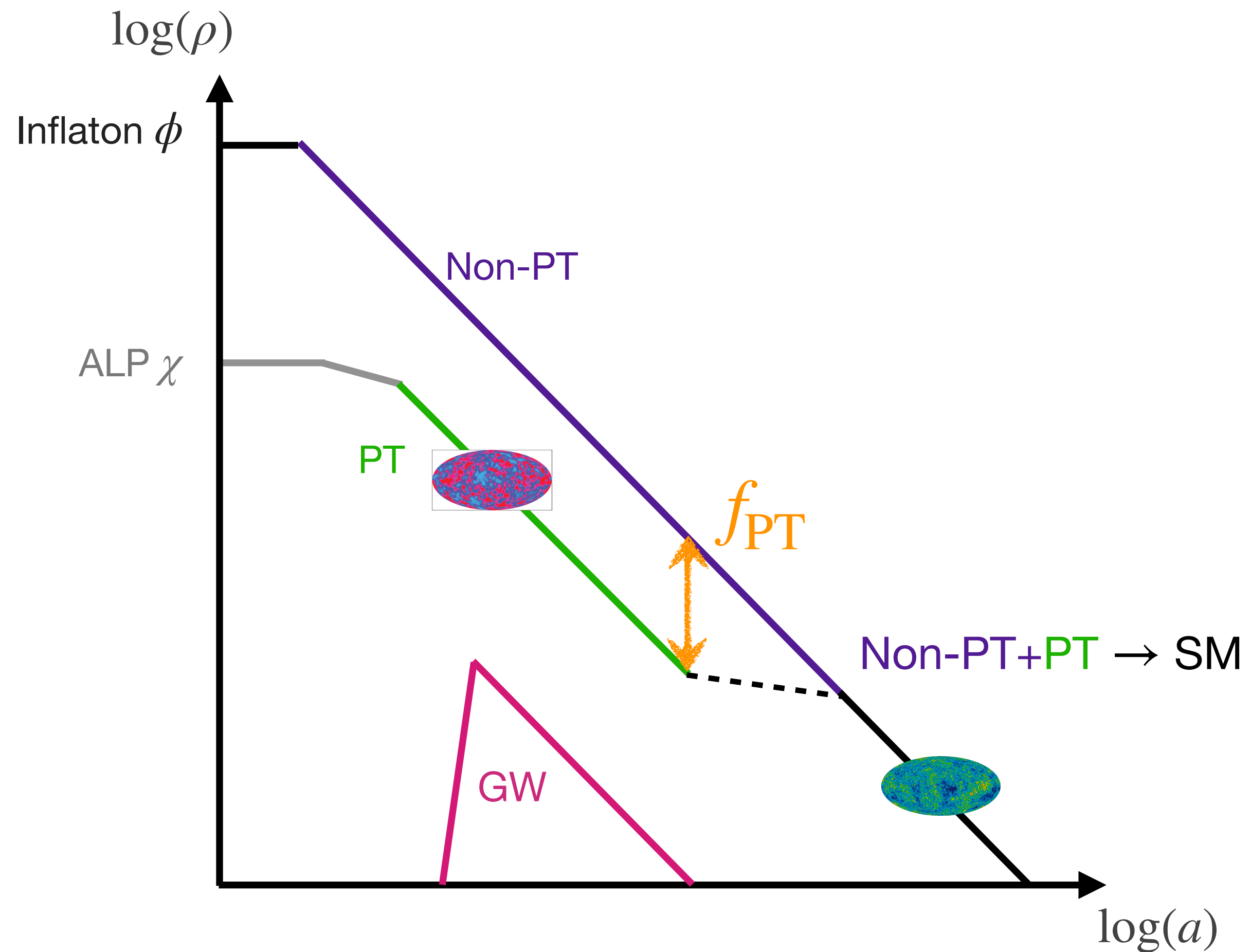
$$\delta_{GW} \sim \delta_\chi + \delta_\phi$$

Sachs-Wolfe contribution $\sim 10^{-5}$

$$\delta_{CMB} \sim \delta_\phi + f_{PT} \delta_\chi \sim 10^{-5}$$

$$f_{PT} = \frac{\rho_{PT}}{\rho_{total}} \ll 1$$

Large isocurvature is more interesting

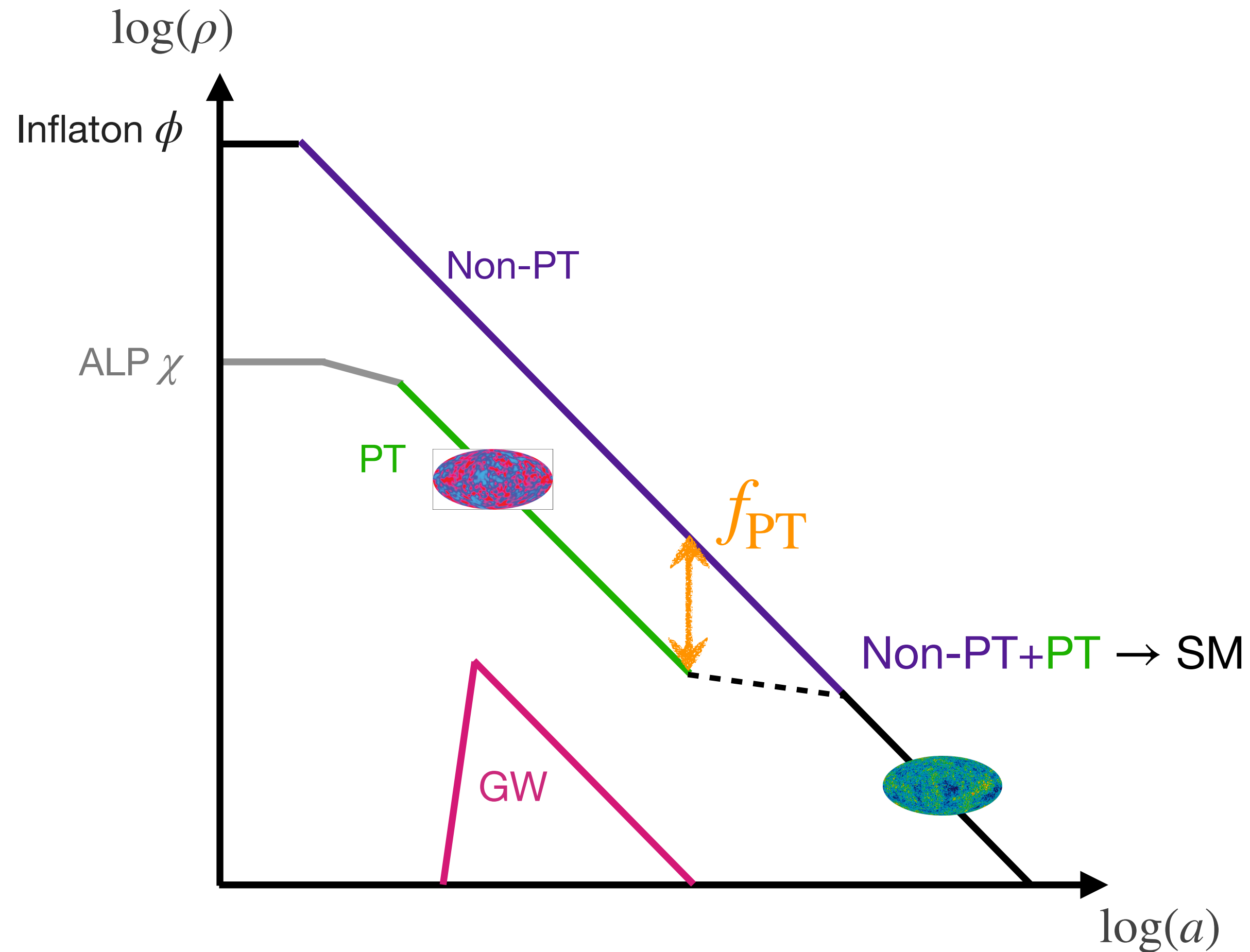


$$\delta_{GW} \sim \delta_{\chi} + \delta_{\phi}$$

$\sim 10^{-5}$

- $\delta_{\chi} > 10^{-5}$ is more distinct

Large isocurvature is more interesting



$$\delta_{\text{GW}} \sim \delta_{\chi} + \delta_{\phi} \rightarrow \sim 10^{-5}$$

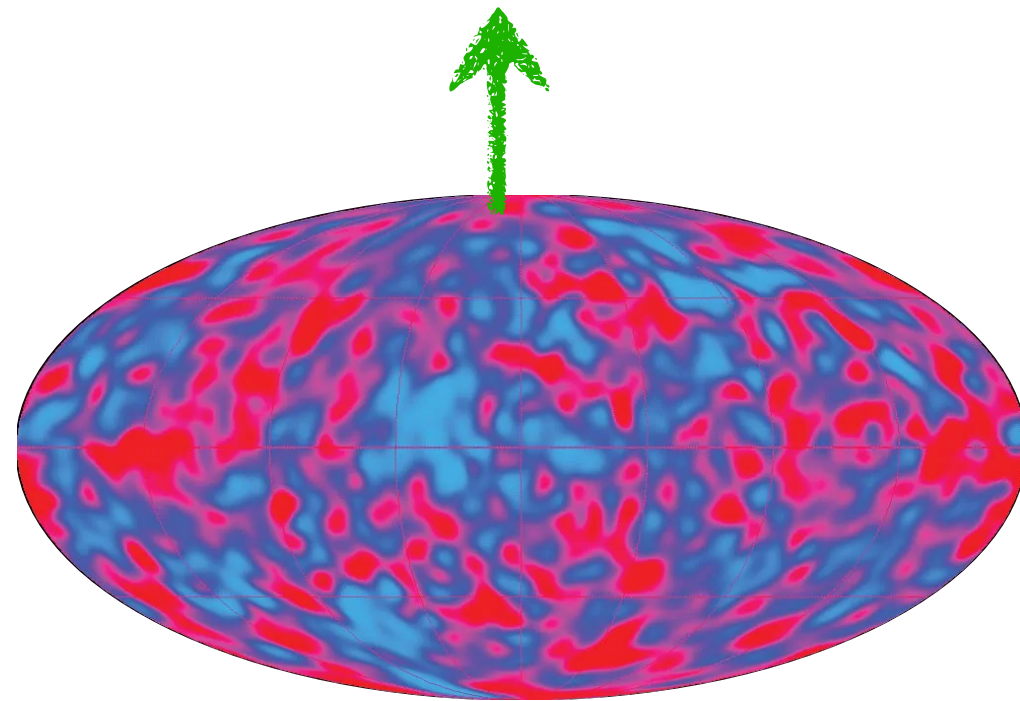
- $\delta_{\chi} > 10^{-5}$ is more distinct

- $\delta_{\chi} \sim \frac{H_{\text{inf}}}{\chi_{\text{init}}} \rightarrow$ in high-scale inflation

$H_{\text{inf}} \sim 10^{-5} M_{\text{Pl}}$, any sub-Planckian misalignment χ_{init} gives $\delta_{\chi} > 10^{-5}$

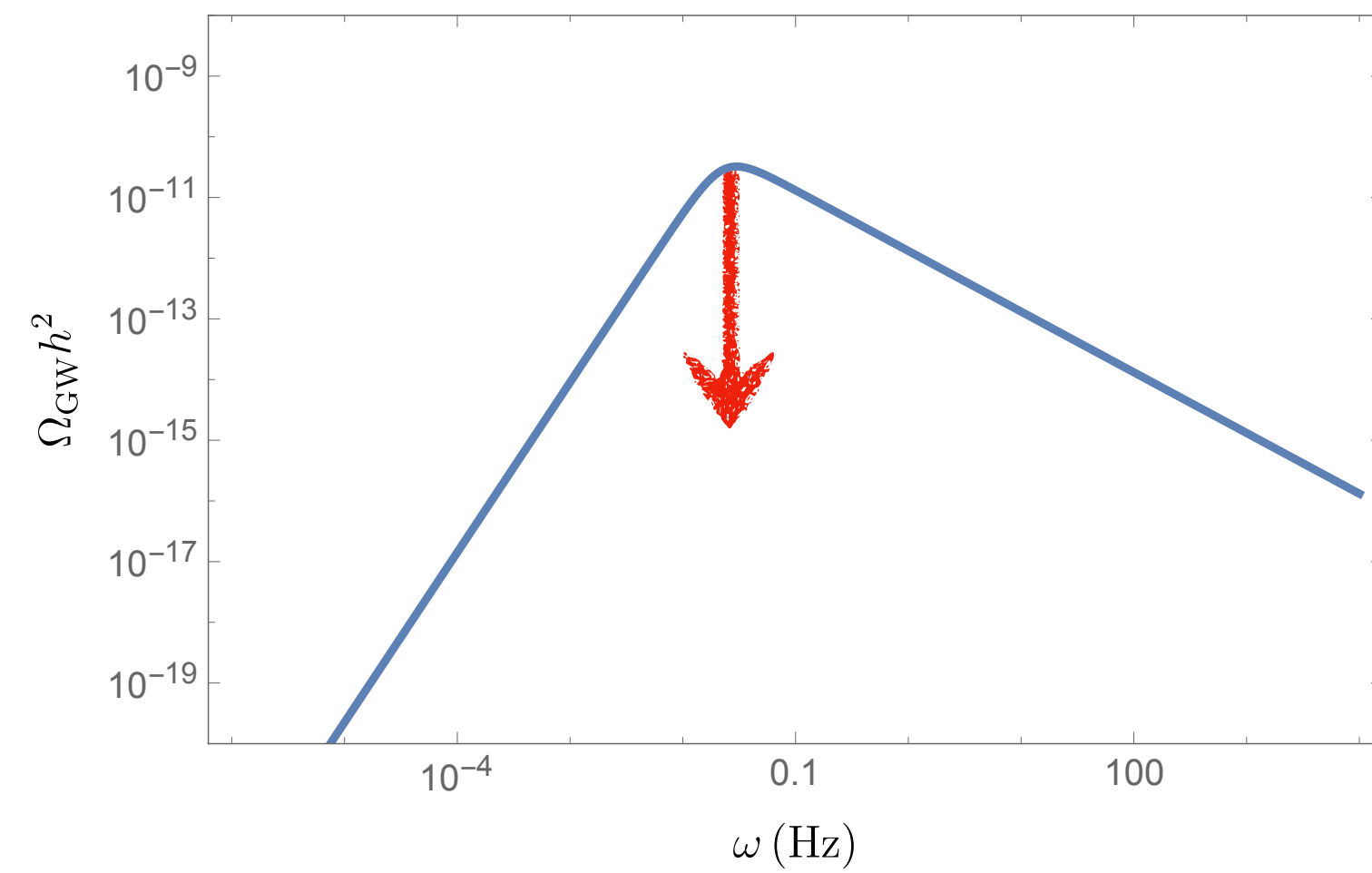
The tradeoff

Geller et al. 1803.10780
M. Breitbach et al. 1811.11175
M. Fairbairn et al. 1901.11038



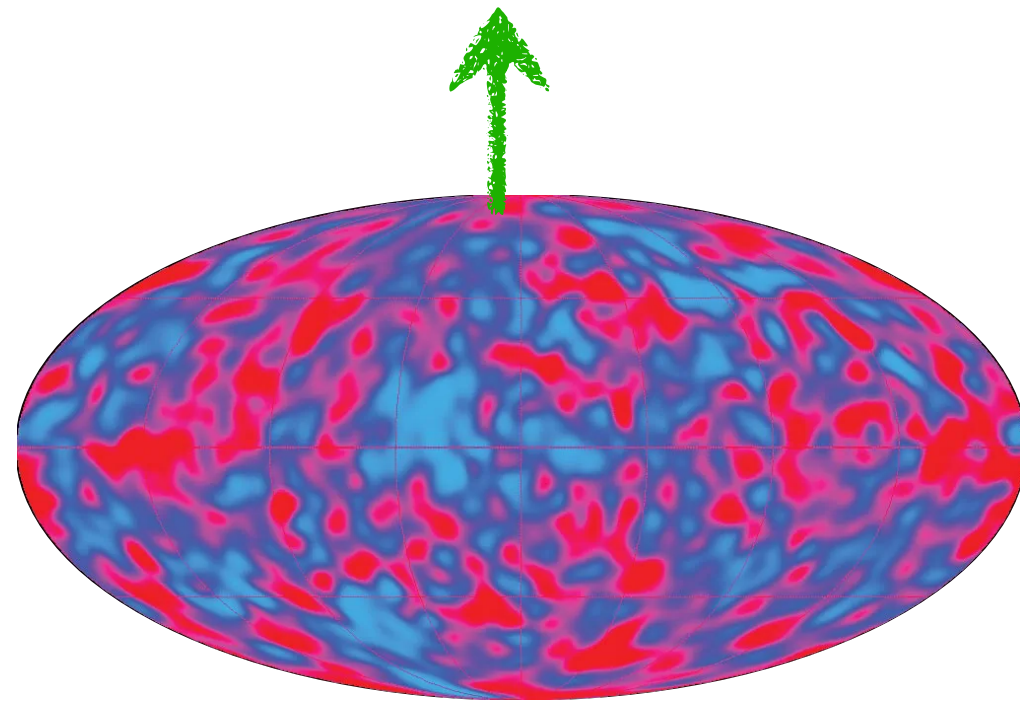
$$\delta_{\text{GW}} \sim \delta_{\chi} + \delta_{\phi}$$

A green arrow points upwards from the δ_{χ} term in the equation.



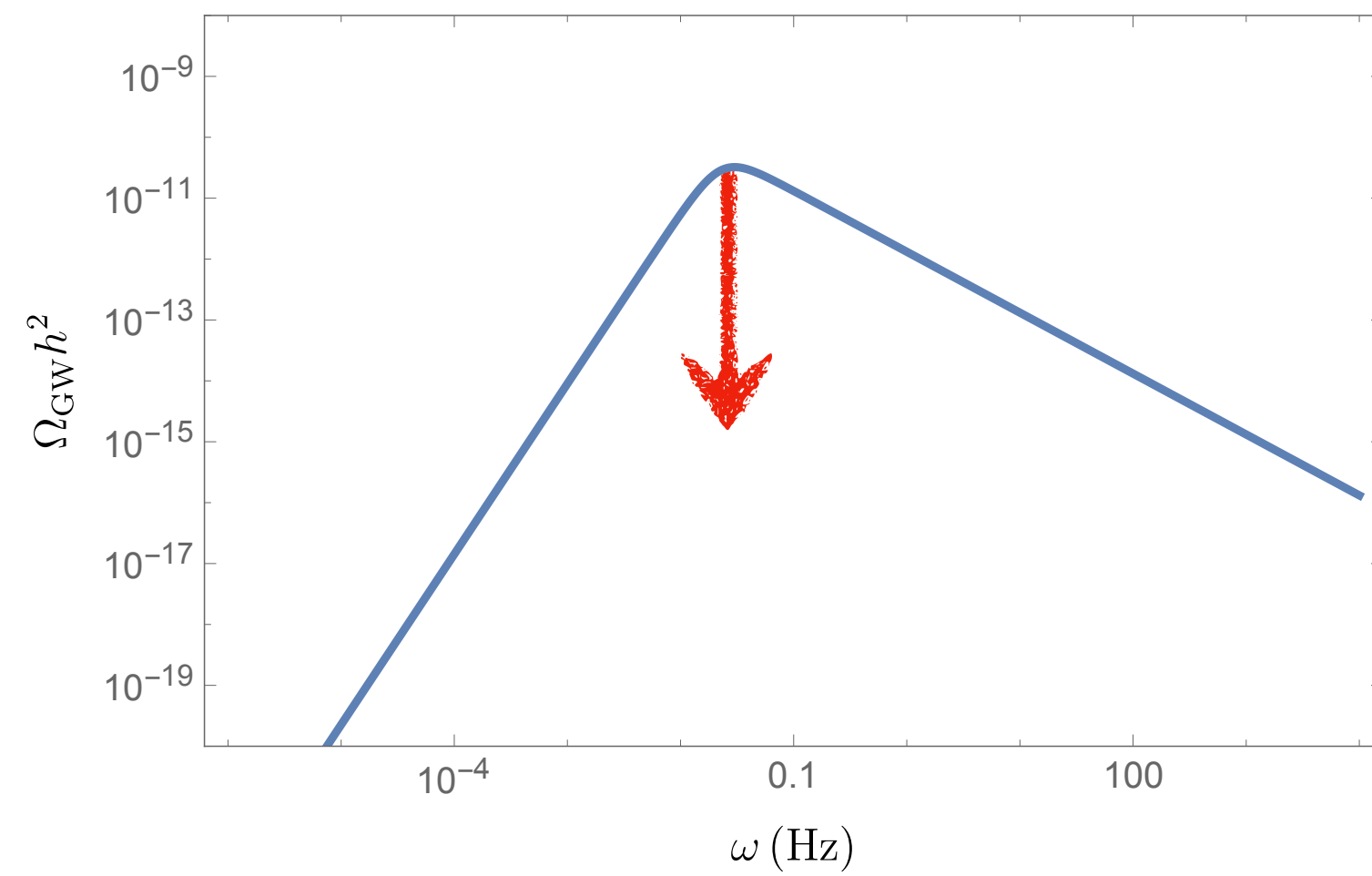
The tradeoff

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$$\delta_{\text{GW}} \sim \delta_{\chi} + \delta_{\phi}$$

A green arrow points upwards from the δ_{χ} term in the equation.

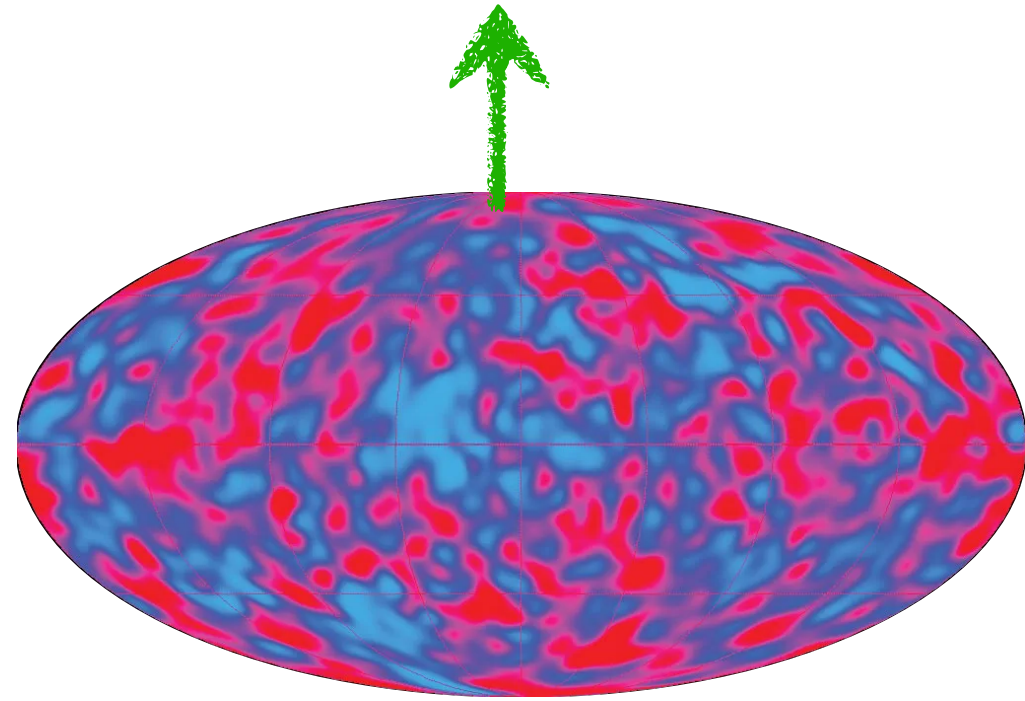


$$\delta_{\text{CMB}} \sim \delta_{\phi} + f_{\text{PT}} \delta_{\chi} \sim 10^{-5}$$

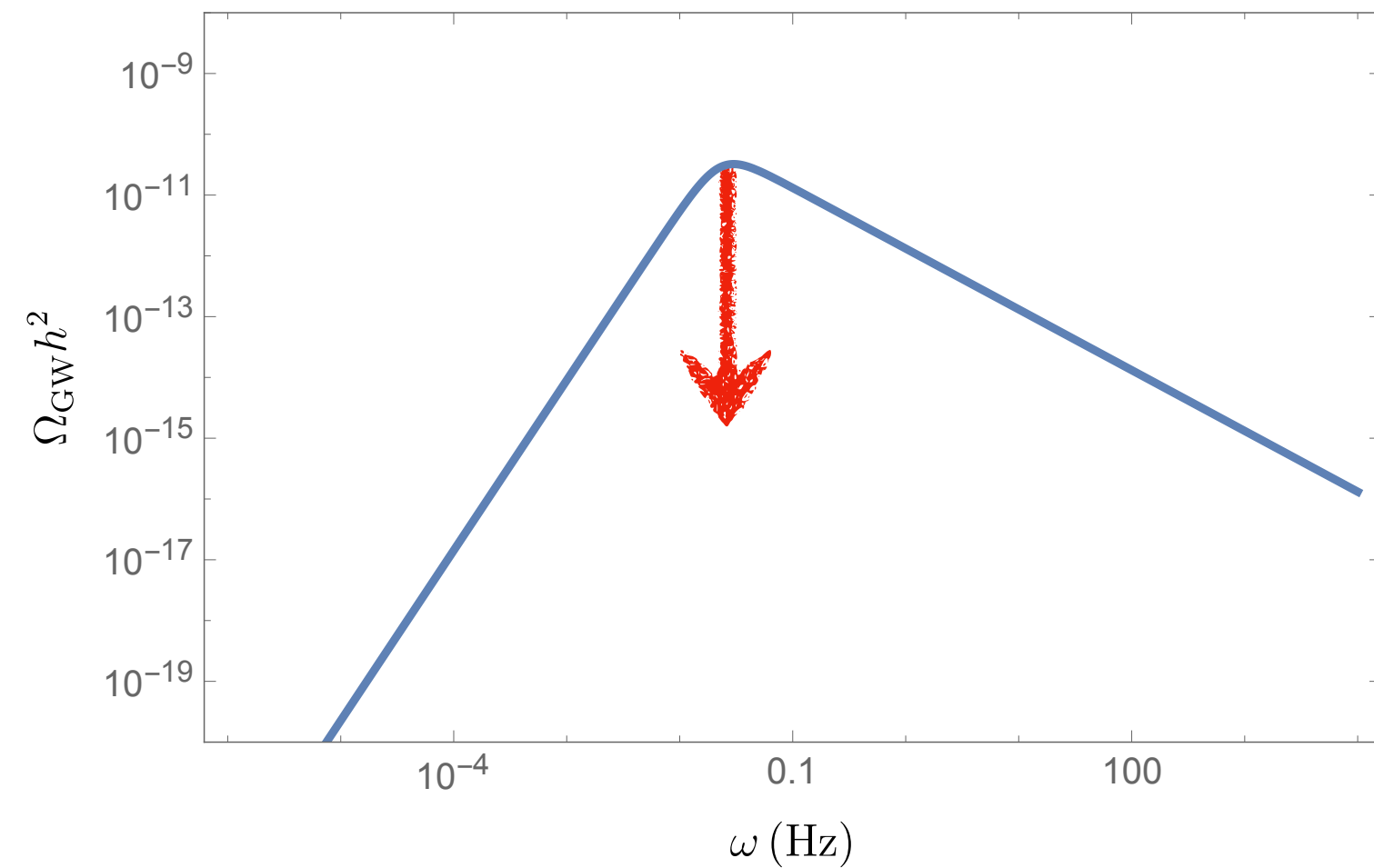
A green arrow points upwards from the δ_{χ} term, and a red arrow points downwards from the f_{PT} term.

The tradeoff

Geller et al. 1803.10780
 M. Breitbach et al. 1811.11175
 M. Fairbairn et al. 1901.11038



$$\delta_{\text{GW}} \sim \delta_{\chi} + \delta_{\phi}$$

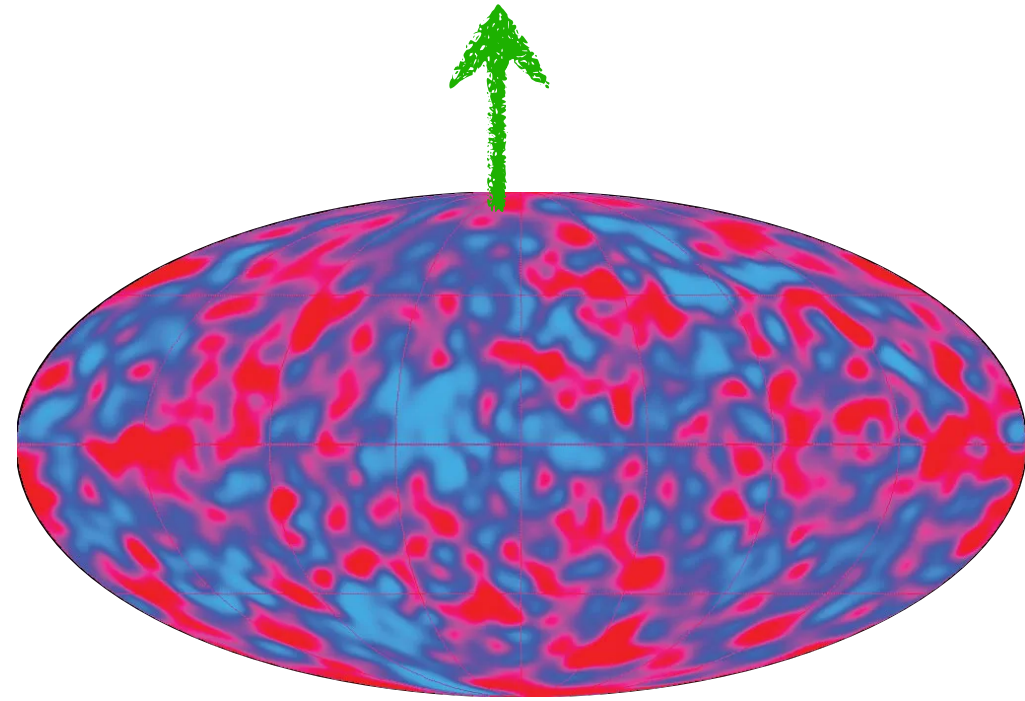


$$\delta_{\text{CMB}} \sim \delta_{\phi} + f_{\text{PT}} \delta_{\chi} \sim 10^{-5}$$

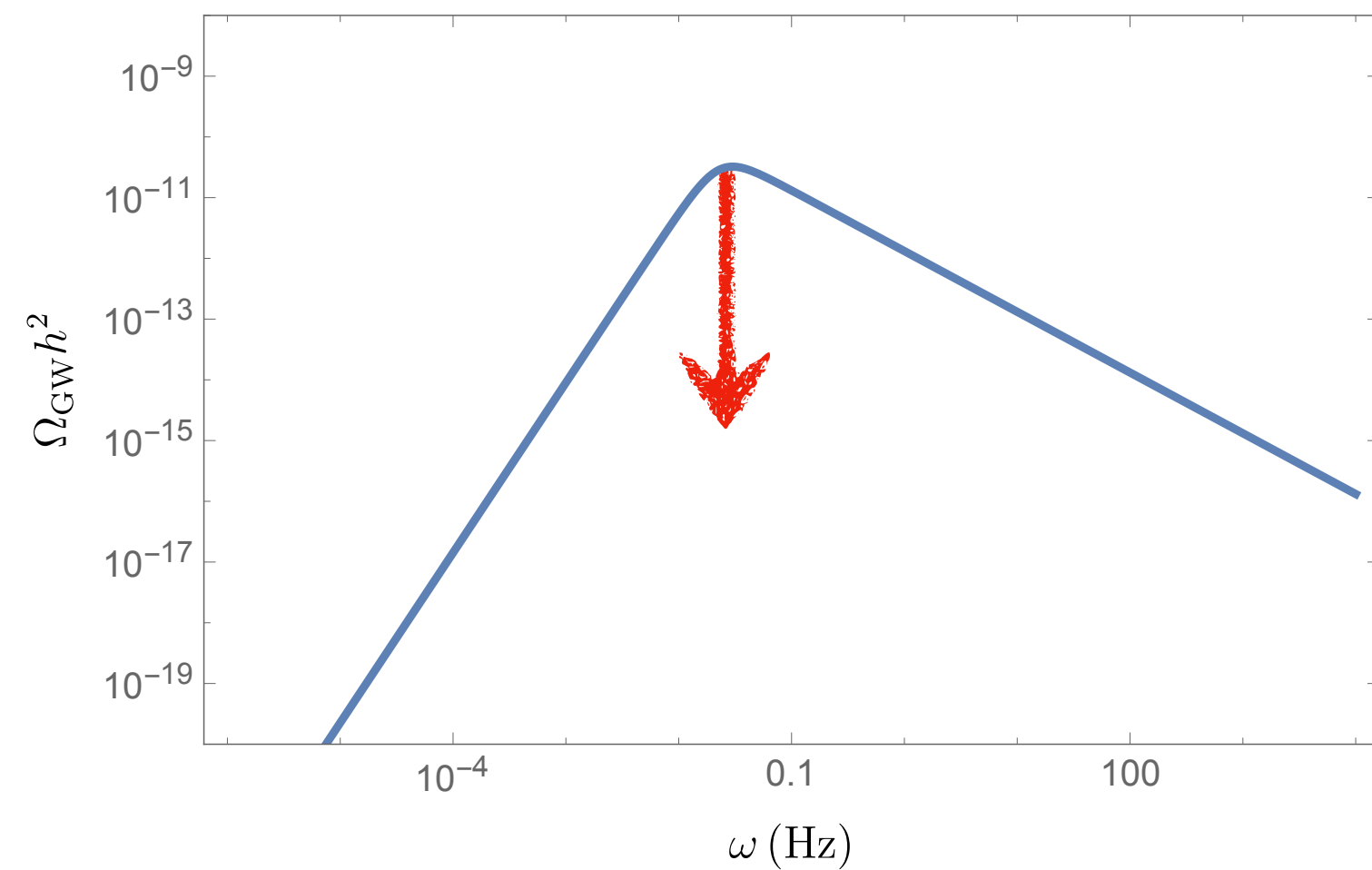
$$\Omega_{\text{GW}} := \frac{\rho_{\text{GW}}}{\rho_{\text{total}}} \propto f_{\text{PT}}^2$$

The tradeoff

Geller et al. 1803.10780
 M. Breitbach et al. 1811.11175
 M. Fairbairn et al. 1901.11038



$$\delta_{\text{GW}} \sim \delta_{\chi} + \delta_{\phi}$$



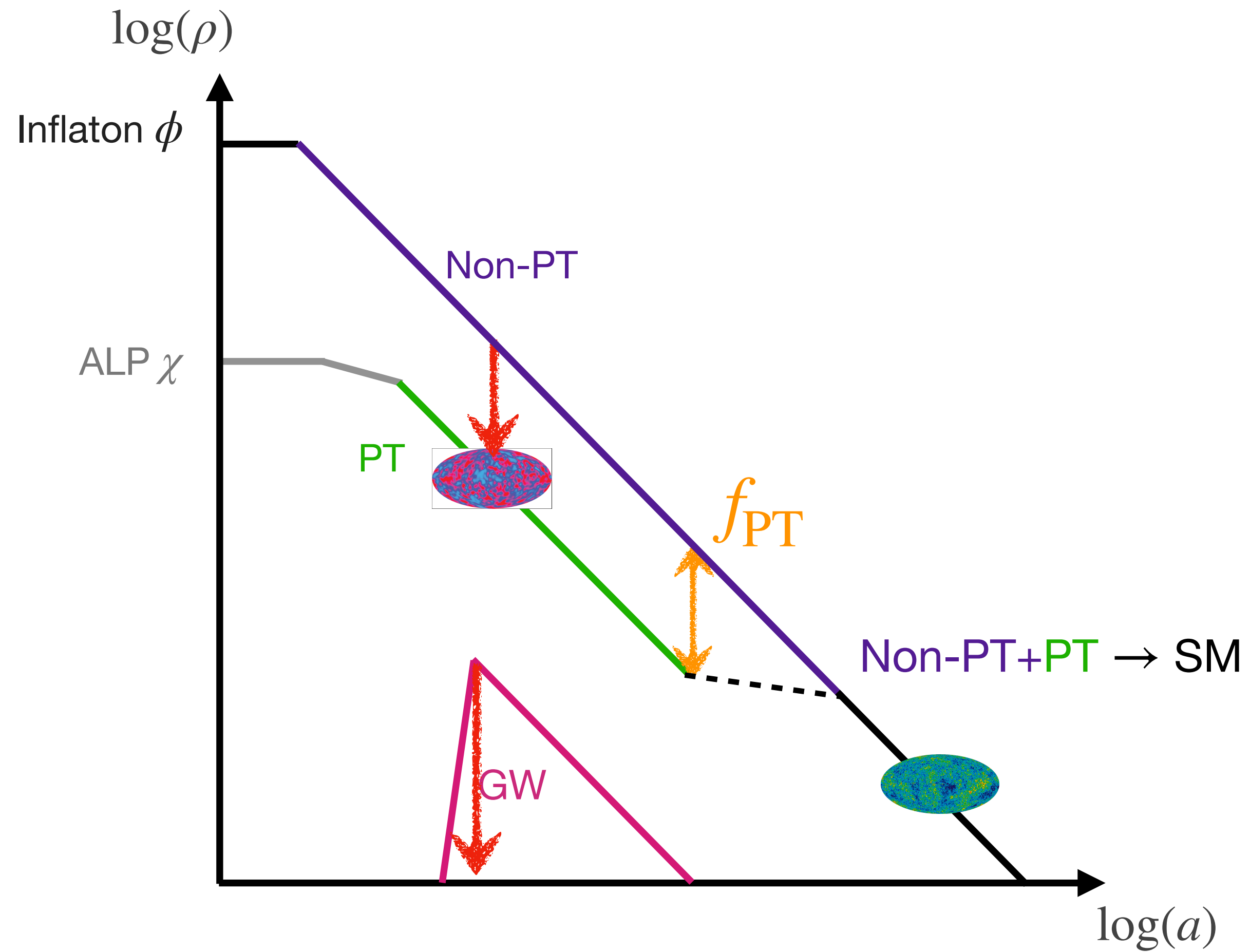
$$\delta_{\text{CMB}} \sim \delta_{\phi} + f_{\text{PT}} \delta_{\chi} \sim 10^{-5}$$

$$\Omega_{\text{GW}} := \frac{\rho_{\text{GW}}}{\rho_{\text{total}}} \propto f_{\text{PT}}^2$$

$$h^{\text{TT}} \propto \rho_{\text{PT}}$$

$$\rho_{\text{GW}} \propto (h^{\text{TT}})^2 \propto \rho_{\text{PT}}^2$$

The solution



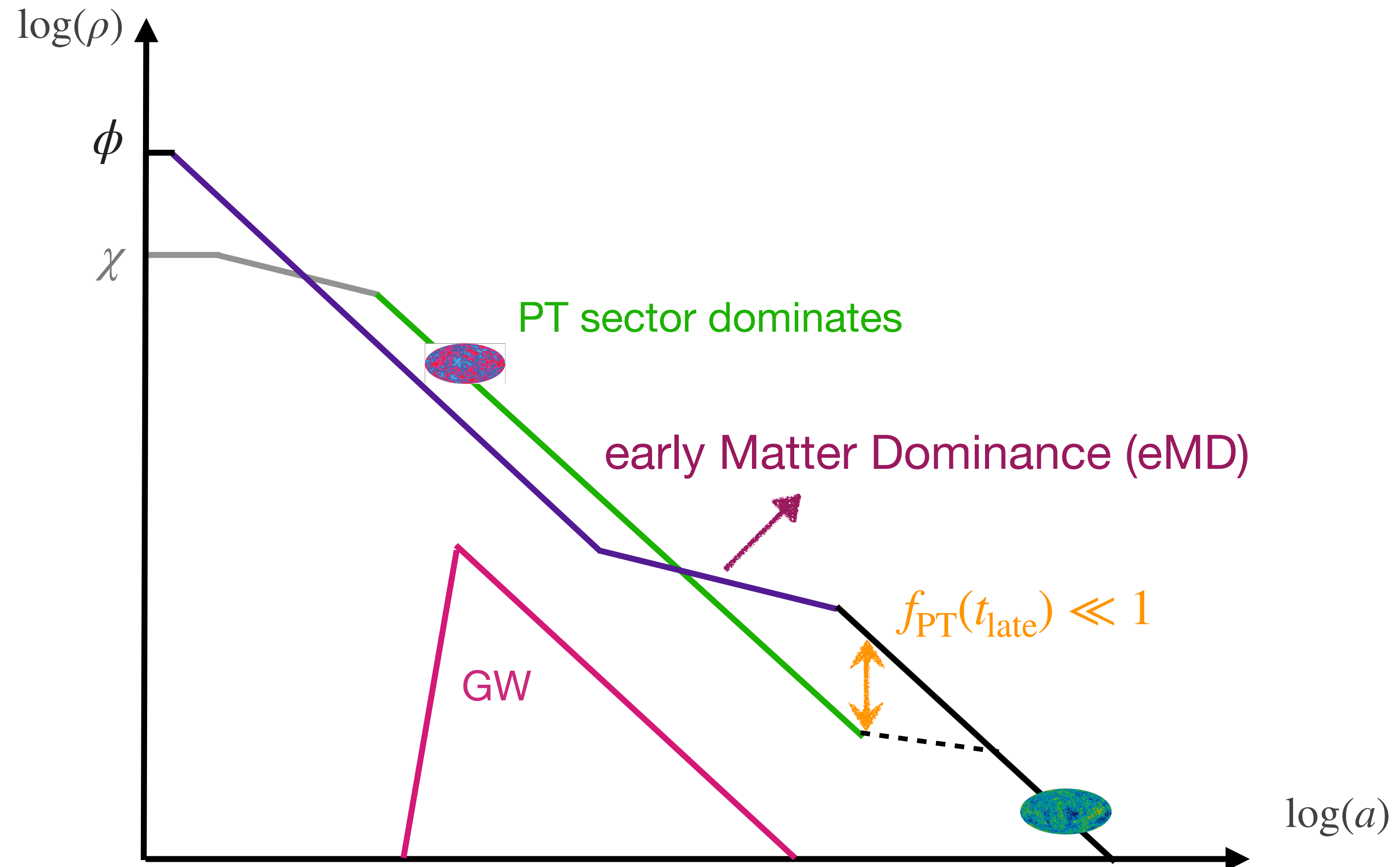
Make f_{PT} time-dependent such that:

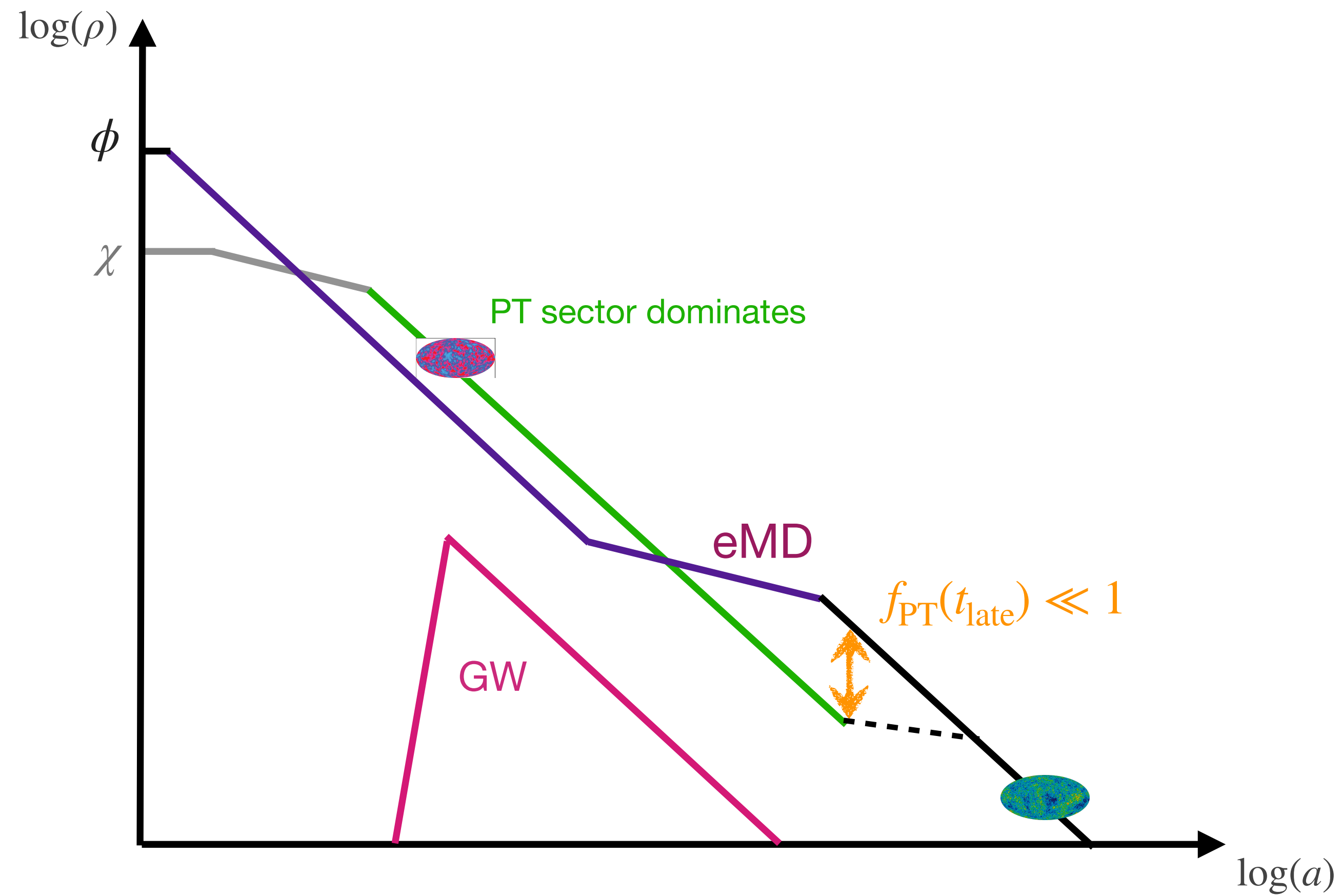
$$f_{PT} \sim 1 \text{ at PT}$$

$$f_{PT} \ll 1 \text{ at CMB}$$

Utilize slower dilution of matter compared to radiation

AB, R. Sundrum: JHEP 06 (2023) 029



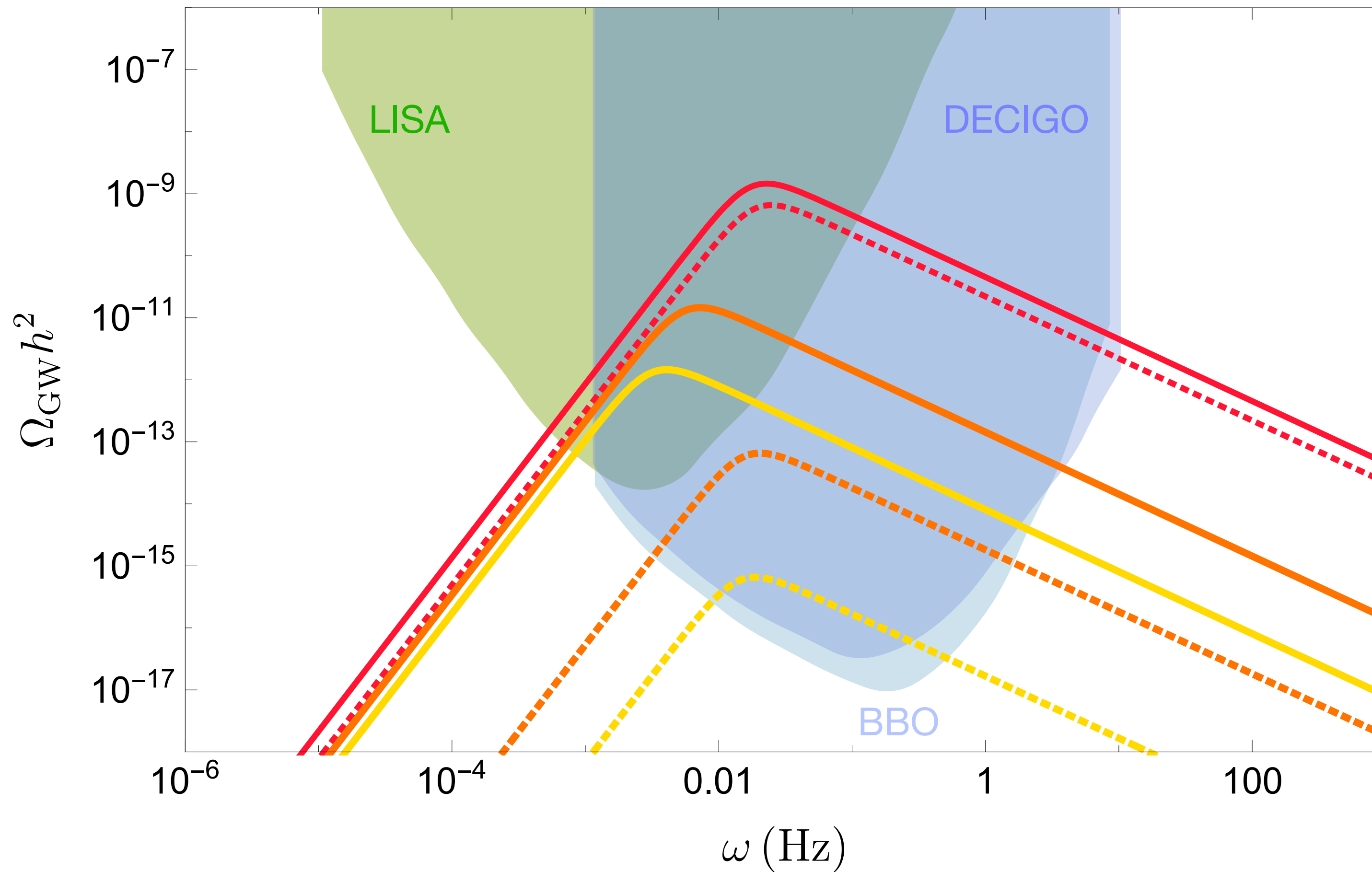


No suppression at production

Relative dilution from eMD

$$\Omega_{\text{GW}}^{\text{today}} \propto f_{\text{PT}} \text{ instead of } f_{\text{PT}}^2$$

Improvement in the GWB signal



Dashed lines: First model

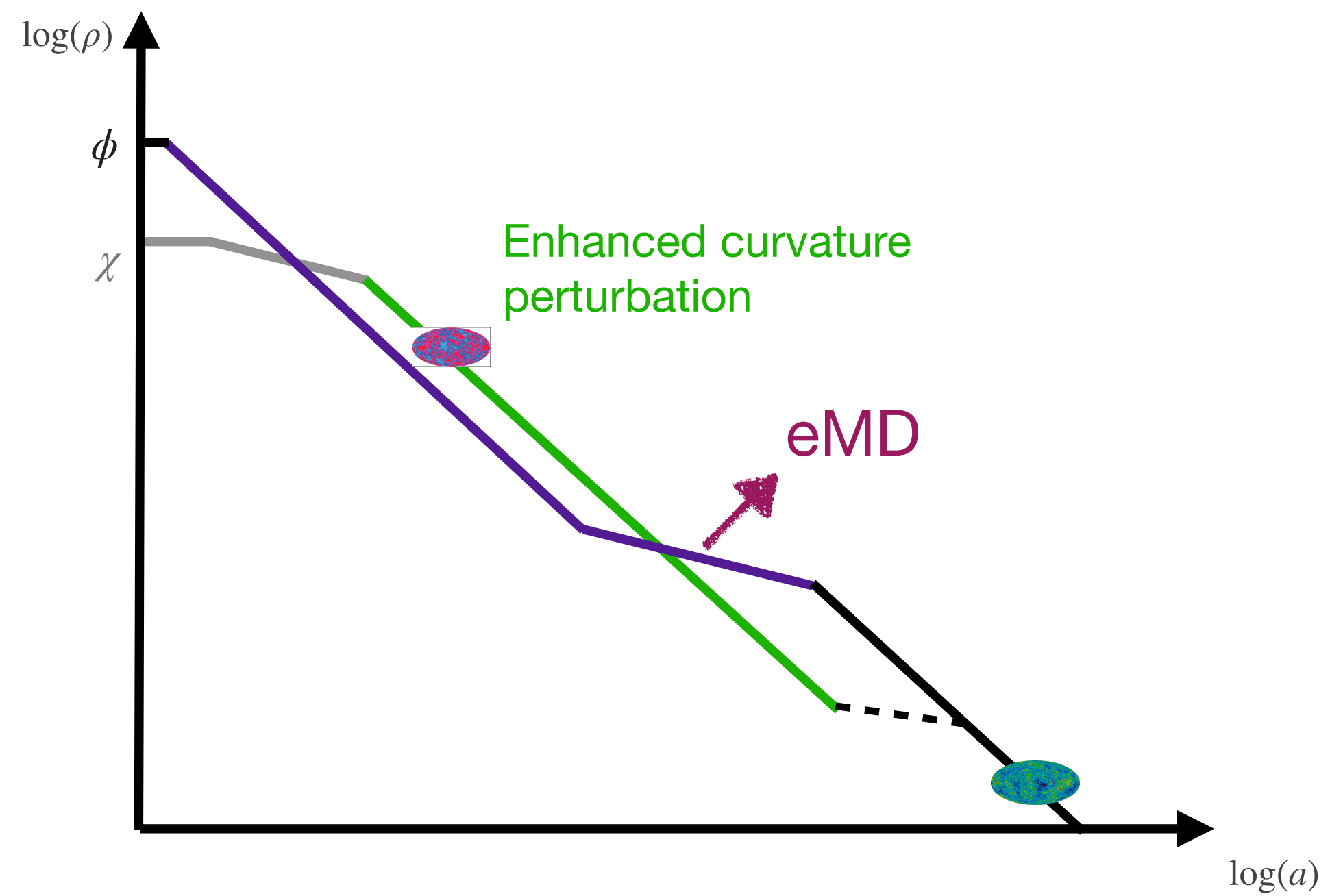
Solid lines: eMD model

— $\delta_{\text{GW}} \sim 10^{-4}$, $f_{\text{PT}} \sim 10^{-1}$

— $\delta_{\text{GW}} \sim 10^{-2}$, $f_{\text{PT}} \sim 10^{-3}$

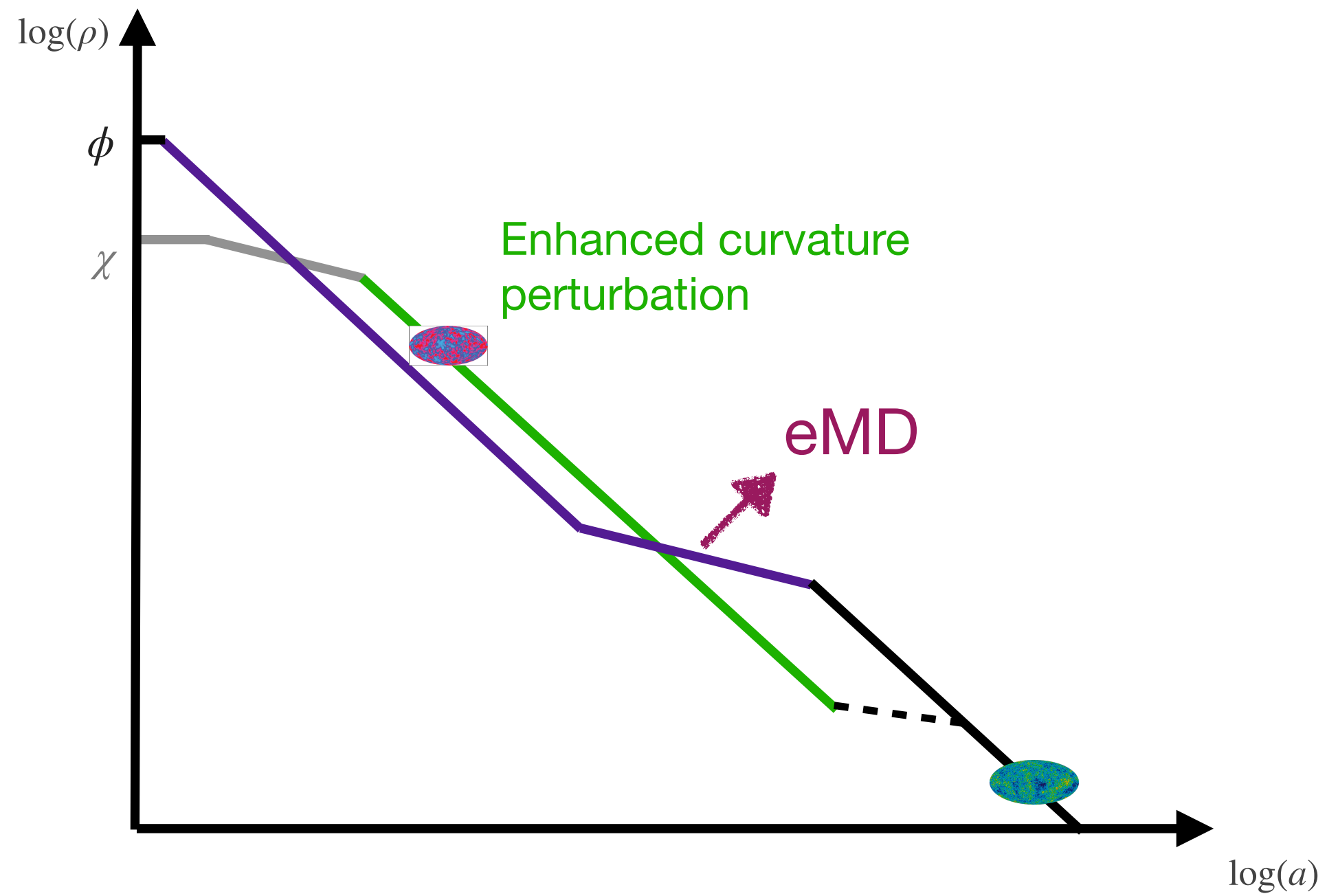
— $\delta_{\text{GW}} \sim 10^{-1}$, $f_{\text{PT}} \sim 10^{-4}$

Primordial black holes

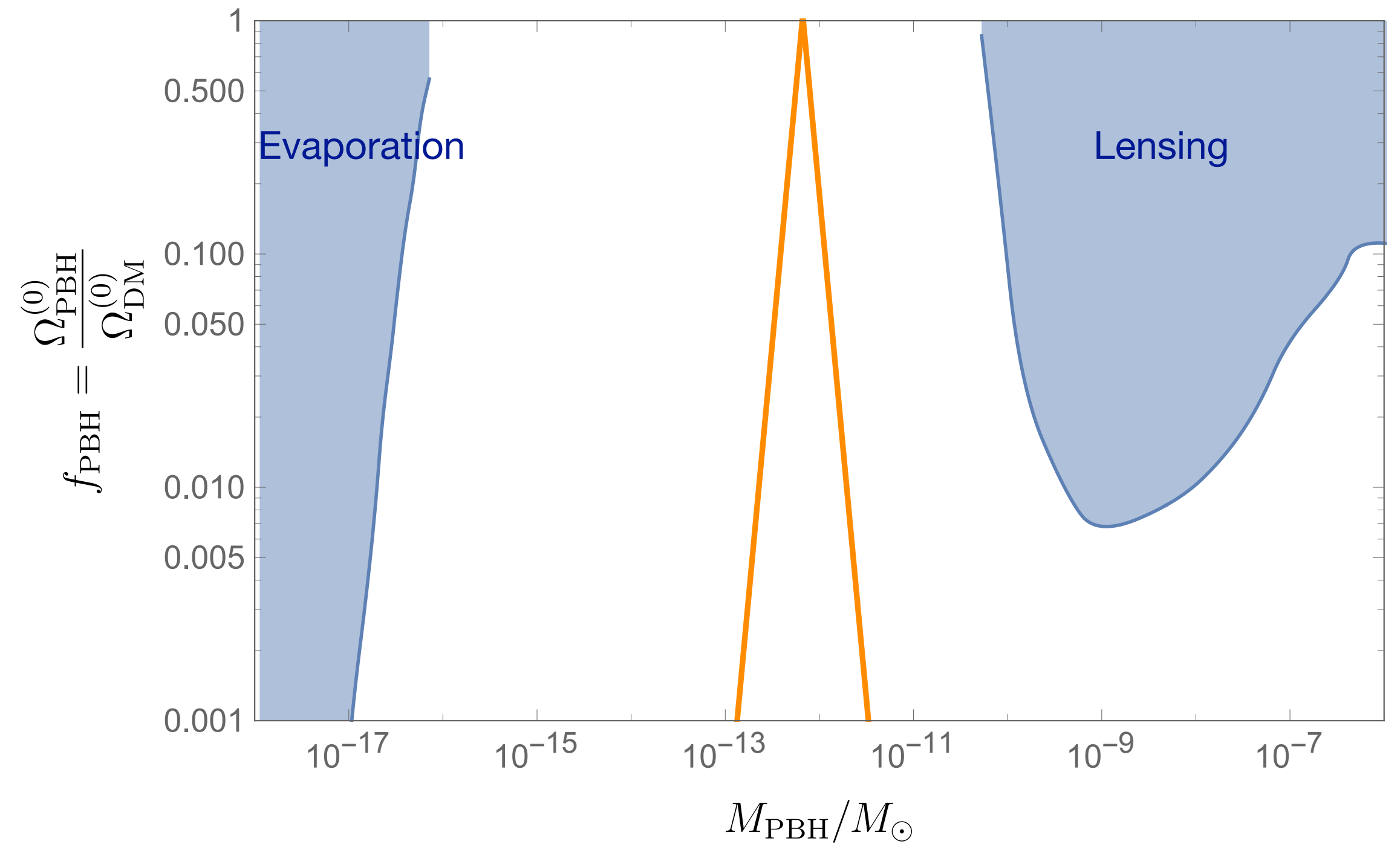


Require $\delta_\chi < 0.01$ on small scales

Primordial black holes



Require $\delta_\chi < 0.01$ on small scales



— $\delta_\chi \sim 10^{-2}, f_{PT} \sim 10^{-3}$

Summary and future directions

- GWB from phase transitions could be highly anisotropic.
- Contrary to the previous belief, such highly anisotropic GWB could be within the reach of detectors like LISA.
- In the future, it would be interesting to extend the study of isocurvature GWB from other sources.
- I invite you to think about other kinds of isocurvature maps, which may reveal new physics from the inflationary era.

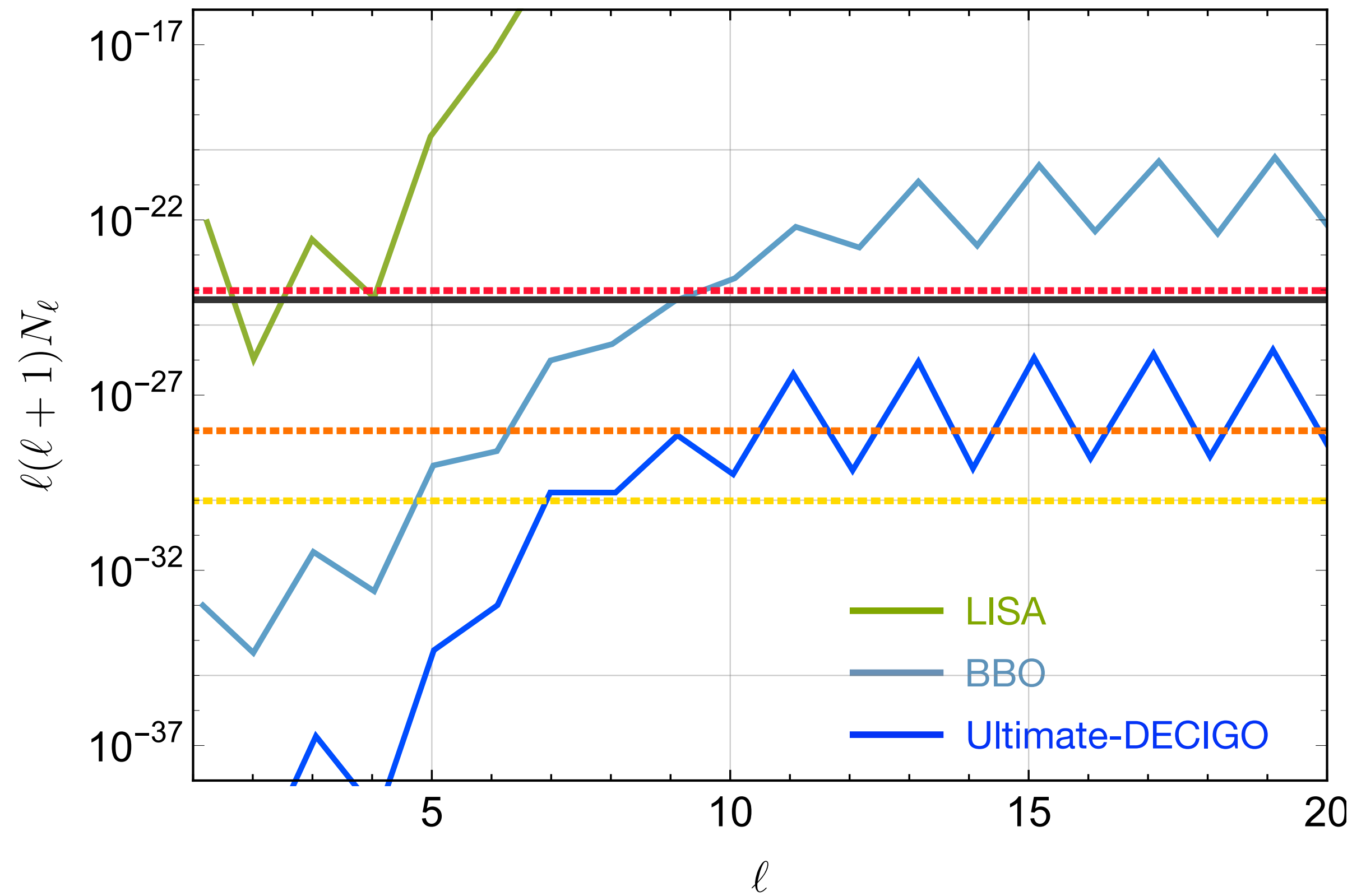
Summary and future directions

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Thank you!

Back-up slides

Detectability of large anisotropy



Dashed lines: Simple model

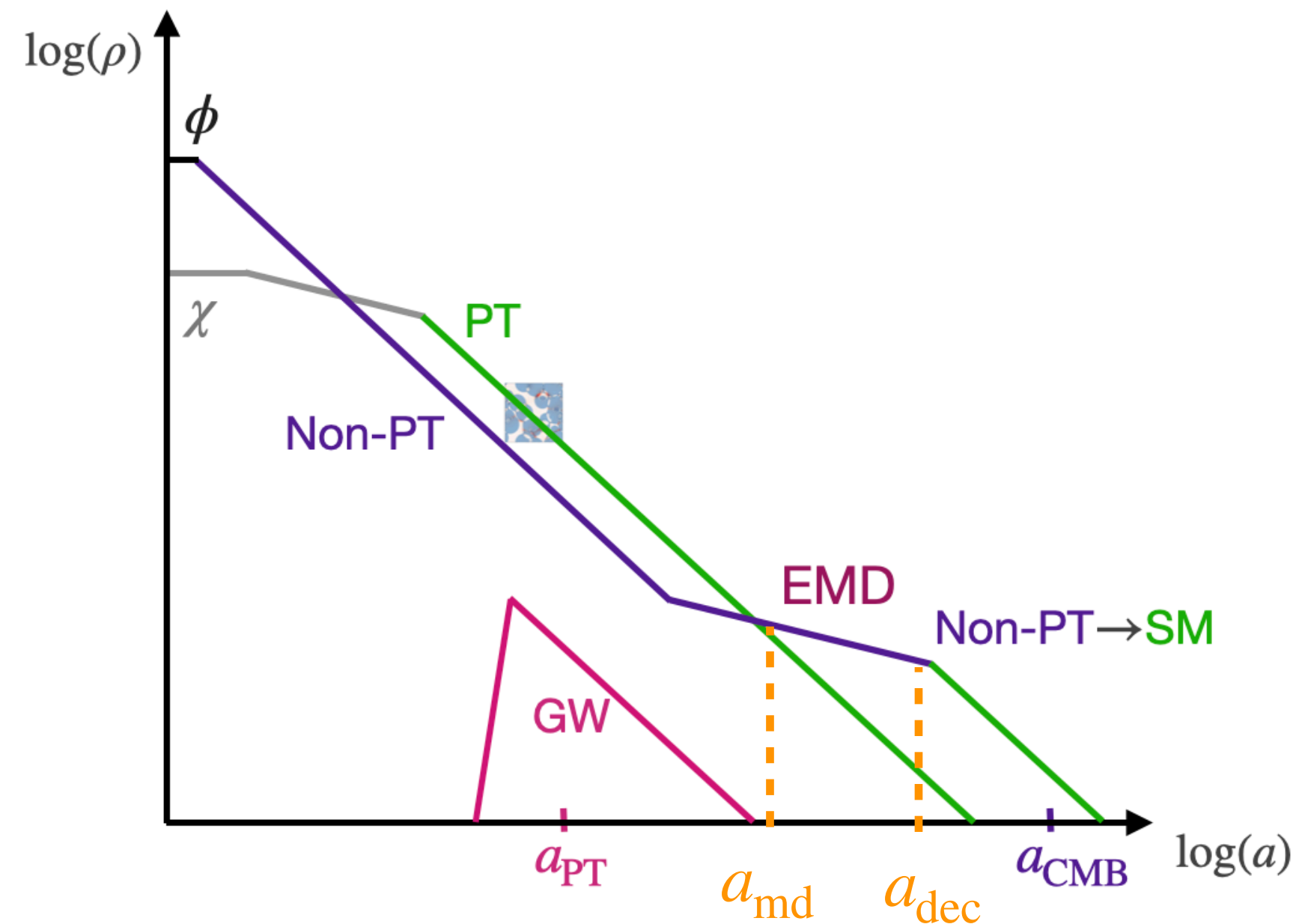
Solid lines: eMD model

— $\delta_{\text{GW}} \sim 10^{-4}, f_{\text{PT}} \sim 10^{-1}$

— $\delta_{\text{GW}} \sim 10^{-2}, f_{\text{PT}} \sim 10^{-3}$

— $\delta_{\text{GW}} \sim 10^{-1}, f_{\text{PT}} \sim 10^{-4}$

Constraints on eMD model: small scale structure



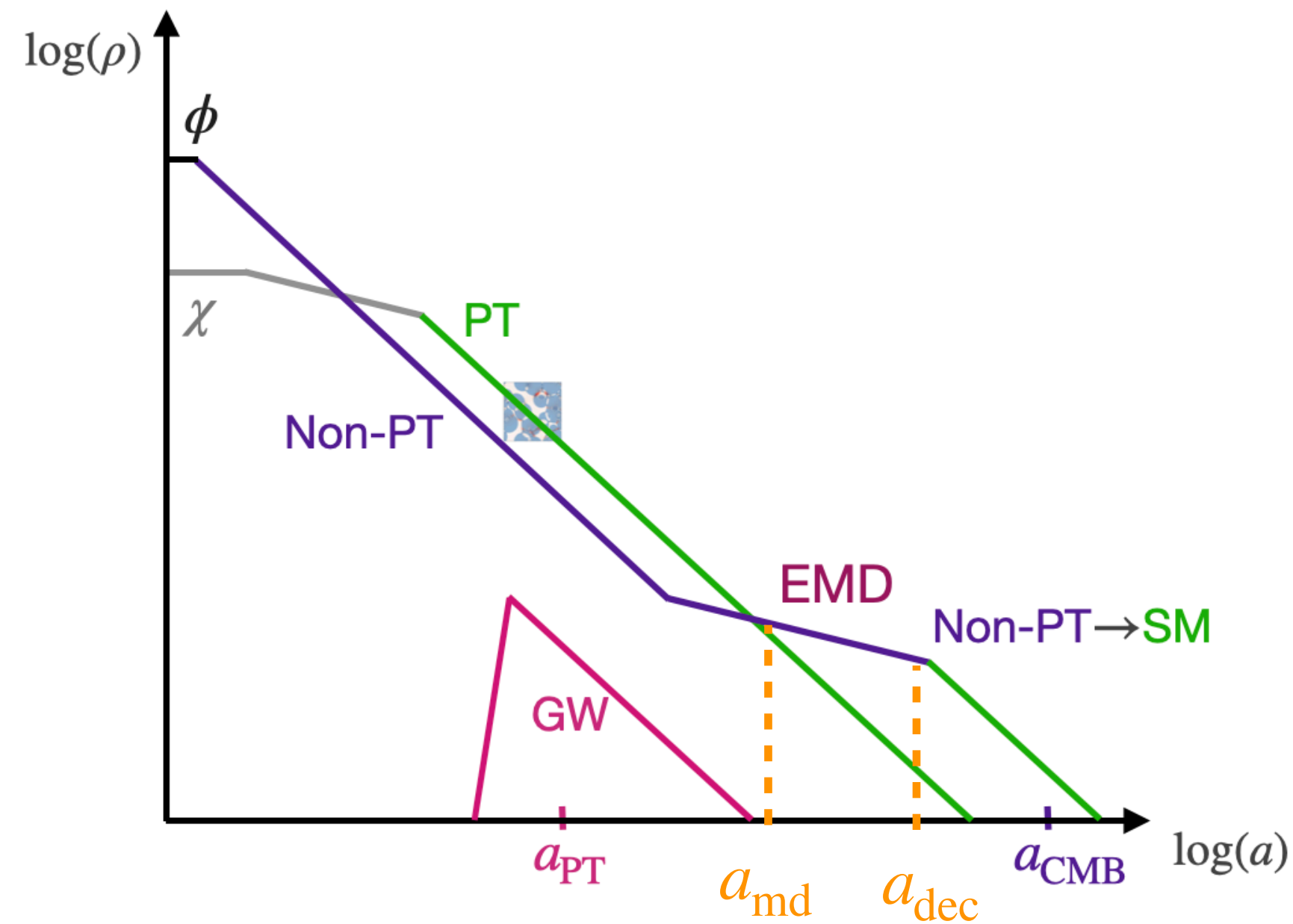
During eMD, $\delta_{\text{mat}} \propto a$, could form structures on small scales with potential observational constraints

However, several mechanism erase these large fluctuations after the decay of non-PT sector into SM:

- Damping of radiation perturbation during the decay
Fan, Özsoy, and Watson '14
- Frictional damping during kinetic decoupling of DM
Loeb, Zaldarriaga '05, Bertschinger '06
- Free-streaming of DM after kinetic decoupling
Loeb, Zaldarriaga '05, Bertschinger '06
- Silk damping in photons and baryons
Silk '68, Hu, Sugiyama '96

Decay before DM decoupling $\rightarrow T_{\text{dec}} \gtrsim 100 \text{ GeV}$

Secondary (induced) GW



$$\frac{a_{\text{dec}}}{a_{\text{md}}} = f_{\text{PT}}^{-1}$$

GW sourced by scalar perturbations at second order, larger contribution from early matter domination

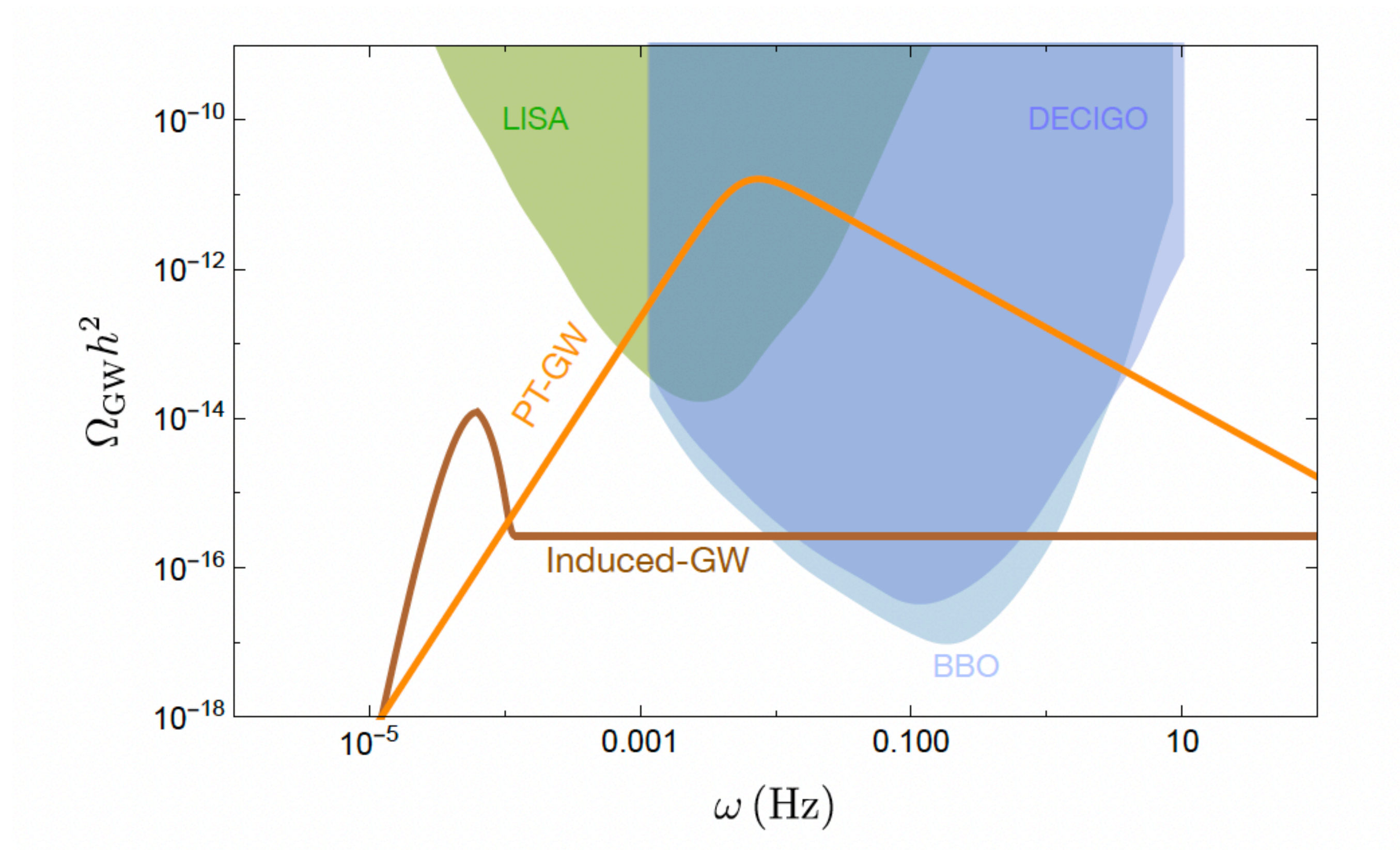
$$(\Omega_{\text{GW}}^{\text{sec}})_{\text{peak}} \sim 10^{-5} \frac{a_{\text{dec}}}{a_{\text{md}}} \delta_{\chi}^4$$

Baumann et al 0703290,
Assadullahi, Wands 0901.0989
Kohri, Terada 1804.08577

eMD imprints unique features in the frequency spectrum of secondary GW.

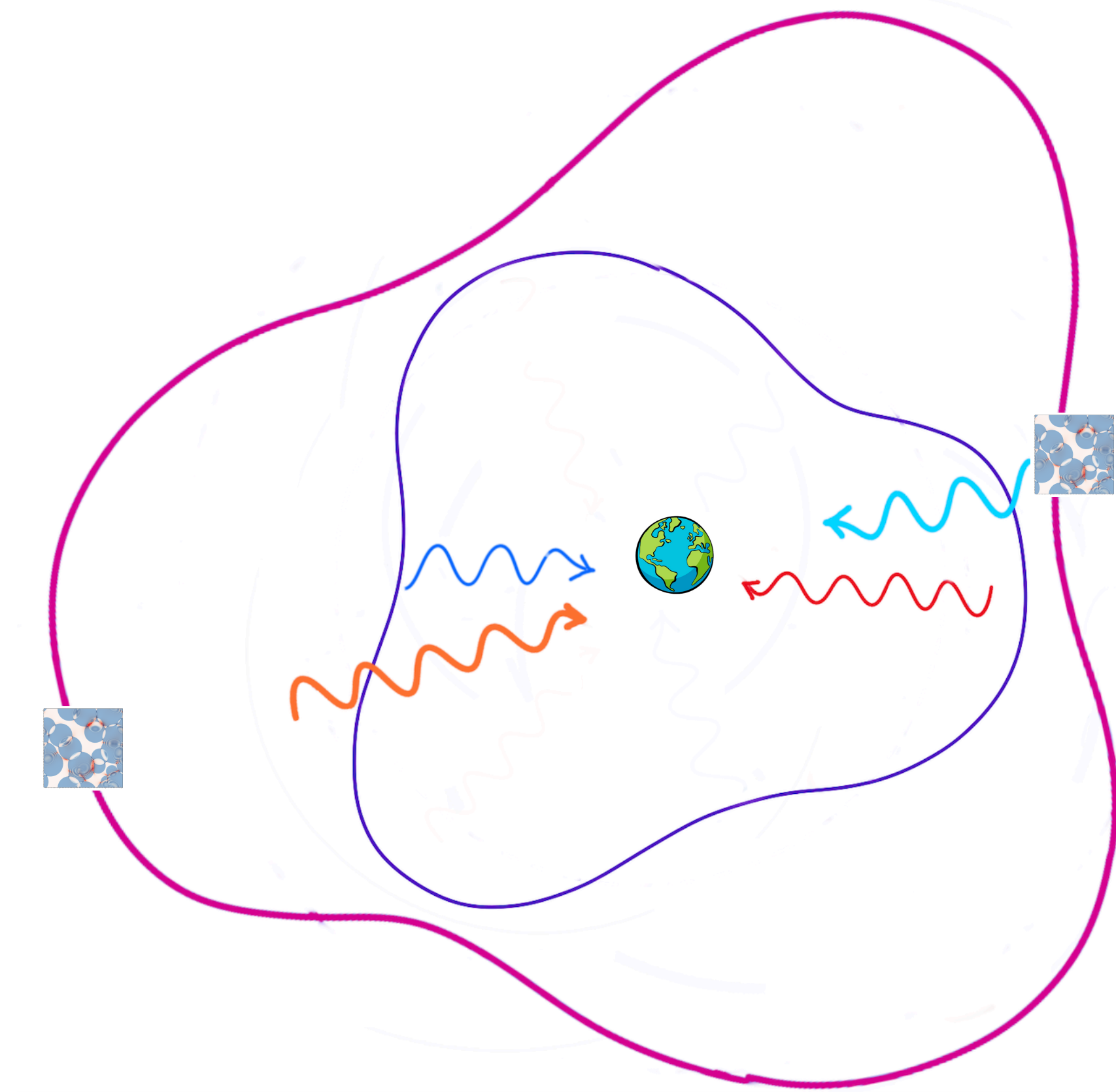
Gouttenoire, Servant, Simakachorn 2111.01150

Secondary (induced) GW



— $\delta_\chi \sim 10^{-2}, f_{\text{PT}} \sim 10^{-3}$

Bubbles are unresolvable

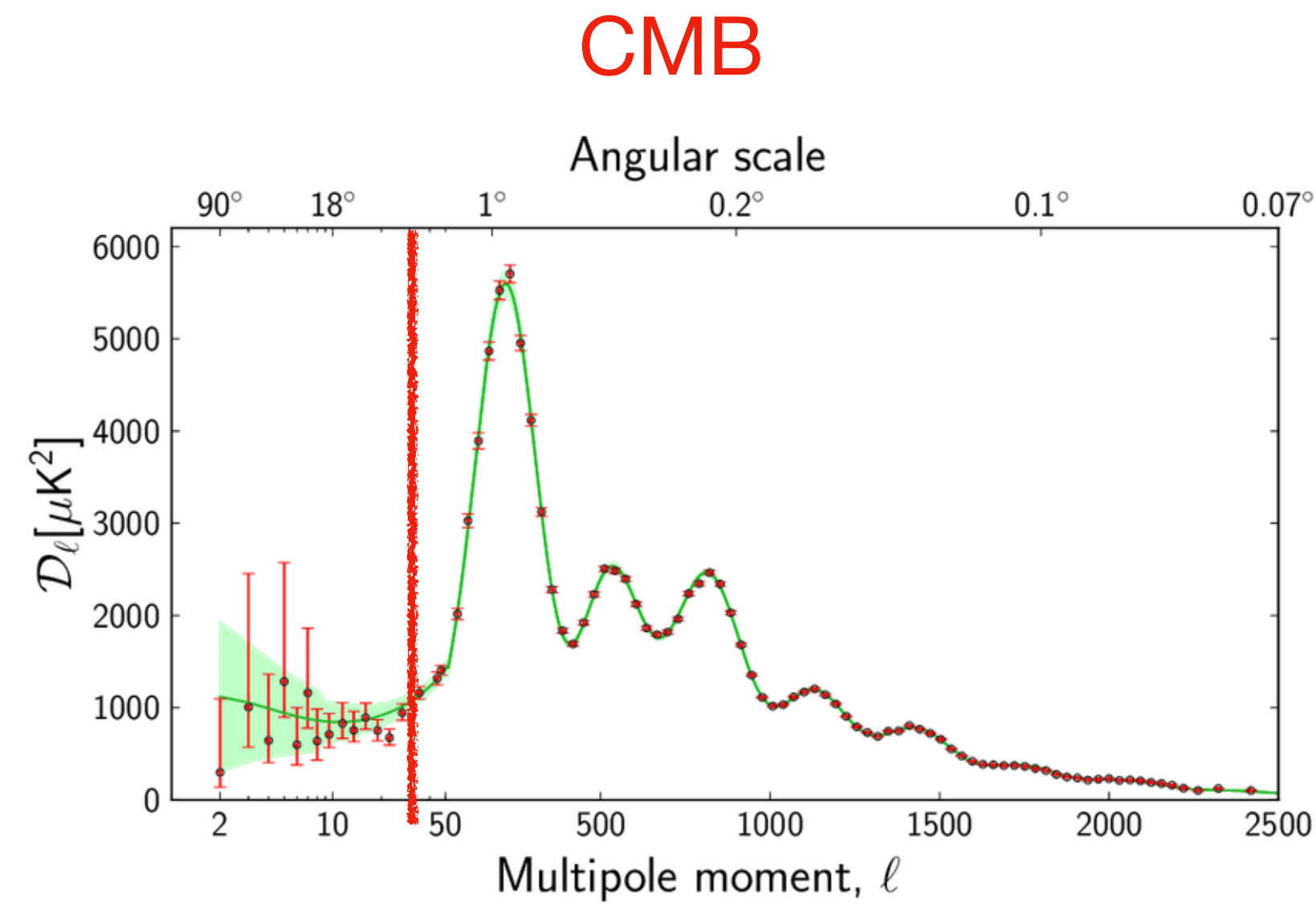


Caution: Zoomed in!

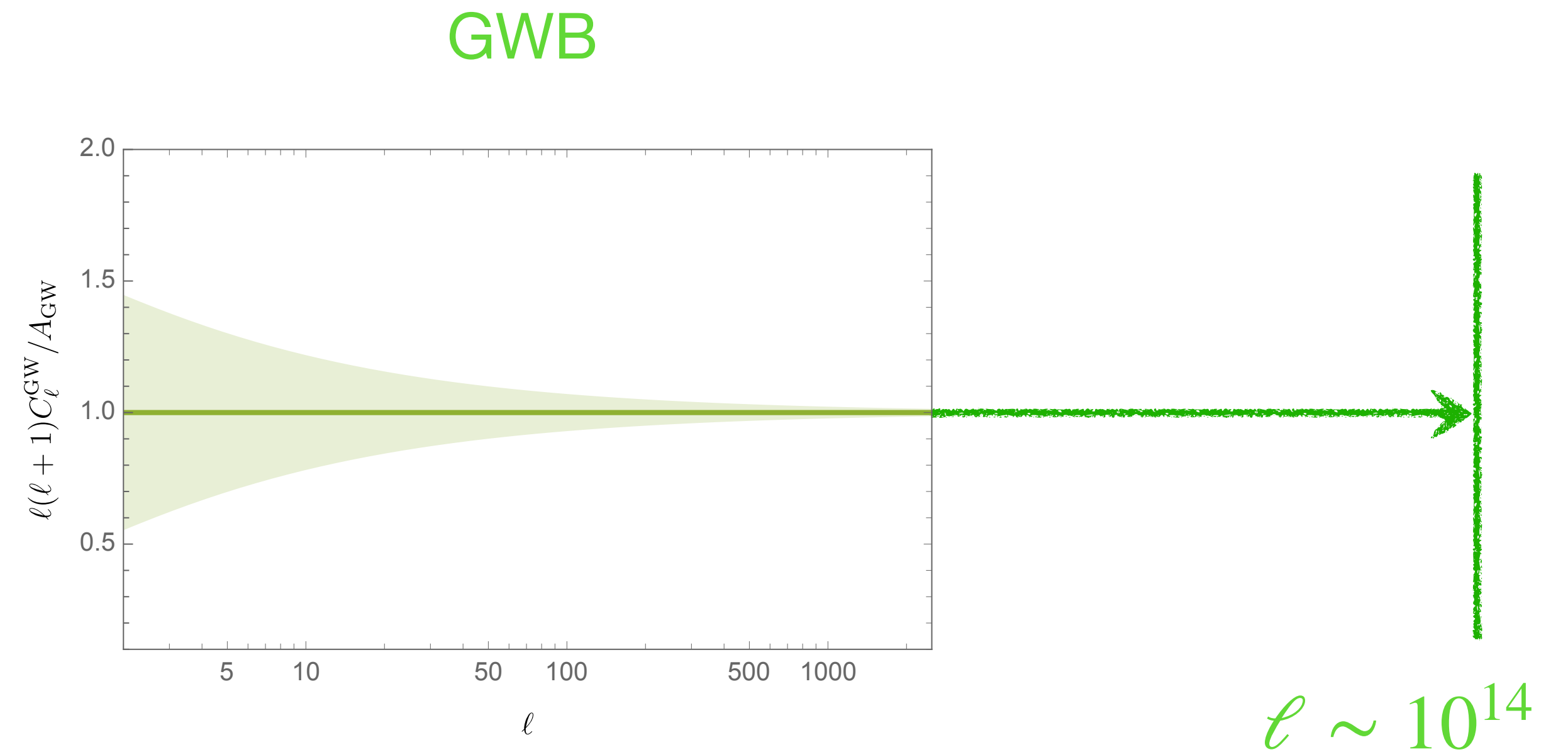
$> 10^{18}$ bubble collisions/arcsec²
⇒ Bubbles are unresolvable
sources : getting a course
grained picture of GWB

GWB is a “pristine” map

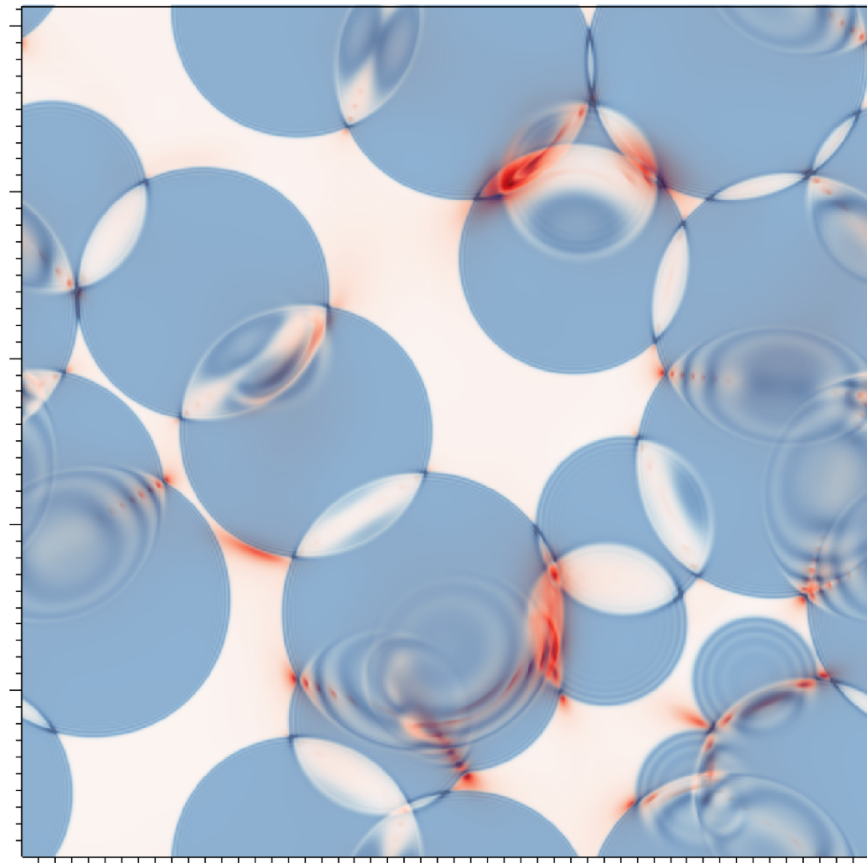
Earlier production + free-streaming of GW \rightarrow large range of scales is unaltered by sub-horizon physics



Planck



Energy density in GWB from PT (2)



Power released in bubble collision

$$\frac{dE_{\text{GW}}}{dt} \sim G_N \left(\frac{d^3 Q}{dt^3} \right)^2$$

Quadrupole moment

$$Q \sim \rho_{\text{lat}} r^5 \rightarrow \frac{d^3 Q}{dt^3} \sim \rho_{\text{lat}} \frac{r^5}{(\Delta t_{\text{PT}})^3}$$

Typical time scale /length scale
(Duration of the PT)

$$r \sim \Delta t_{\text{PT}} \equiv \beta^{-1}$$

GW energy density released in bubble collision

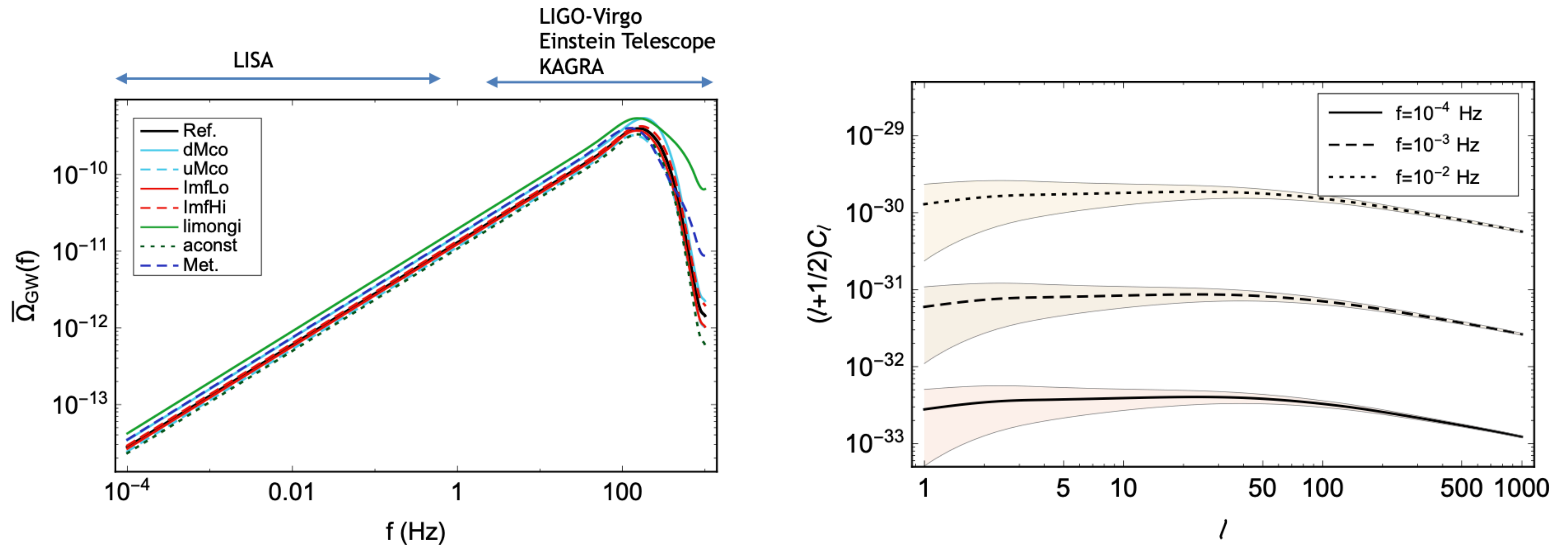
$$\rho_{\text{GW}} \sim \frac{dE_{\text{GW}}}{dt} \frac{\Delta t_{\text{PT}}}{r^3} \sim \frac{G_N \rho_{\text{lat}}^2}{\beta^2}$$

$$G_N \rho_{\text{total}} \sim H^2$$

$$\rho_{\text{GW}} \sim \left(\frac{H}{\beta} \right)^2 \left(\frac{\rho_{\text{lat}}}{\rho_{\text{total}}} \right)^2 \rho_{\text{total}}$$

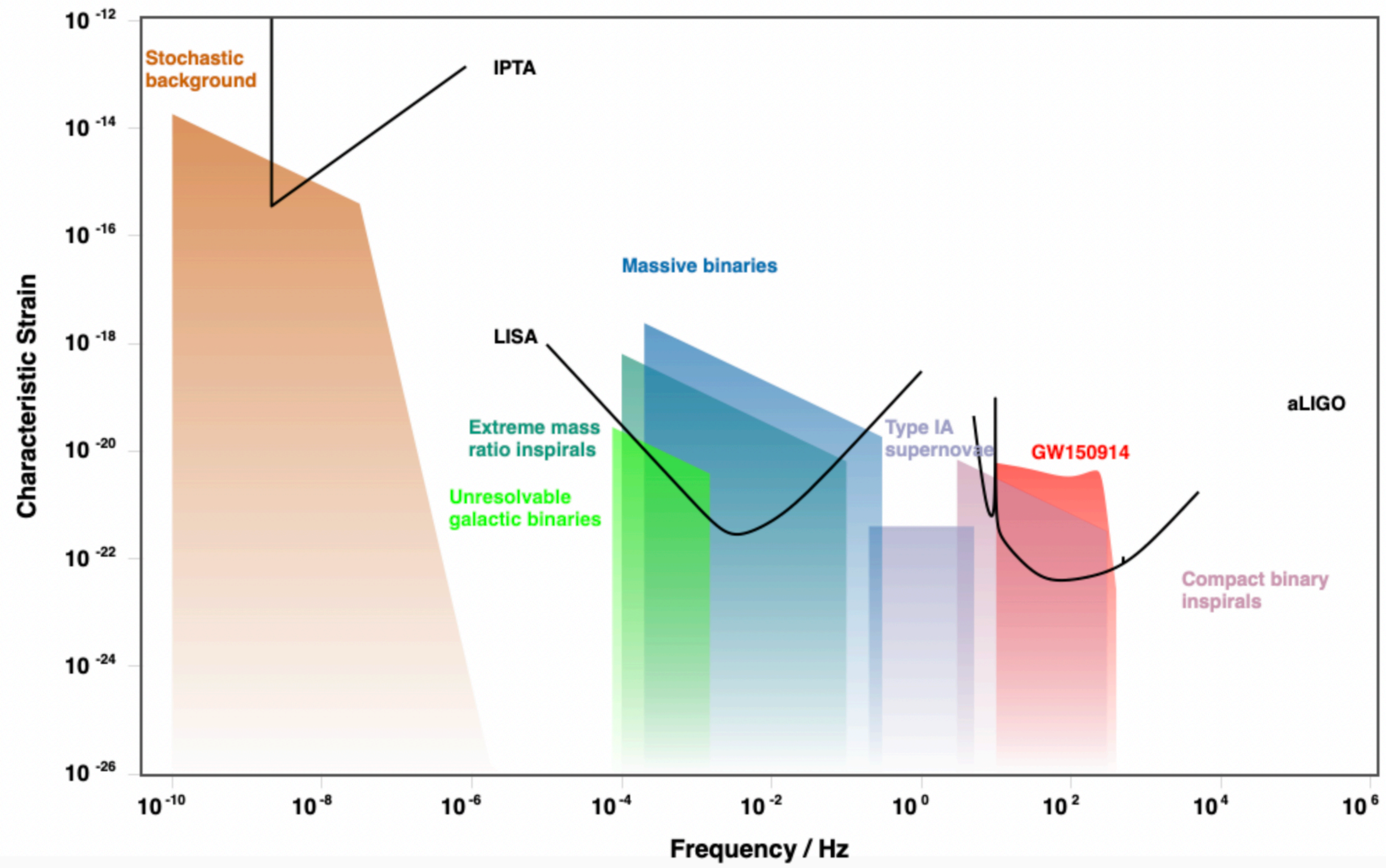
Astrophysical foregrounds in mHz range

Inspiring stellar-mass BH



Giulia Cusin, Irina Dvorkin, Cyril Pitrou, Jean-Philippe Uzan: 1904.07757v2

Also see: 2201.08782v from LISA working group



rhcole.com/apps/GWplotter/