

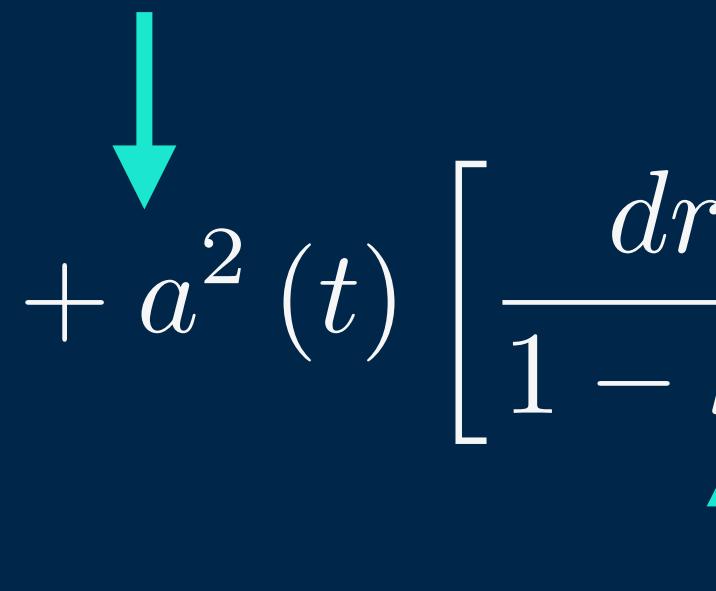
# Cosmology

## Lecture 2: Dynamics of the Universe and Thermal History

# Friedman-Lemaître-Robertson-Walker (FLRW) spacetime: the geometry of a **homogeneous-isotropic** universe

Conventionally  $a = 1$  Today

**Scale factor** – determined by matter content

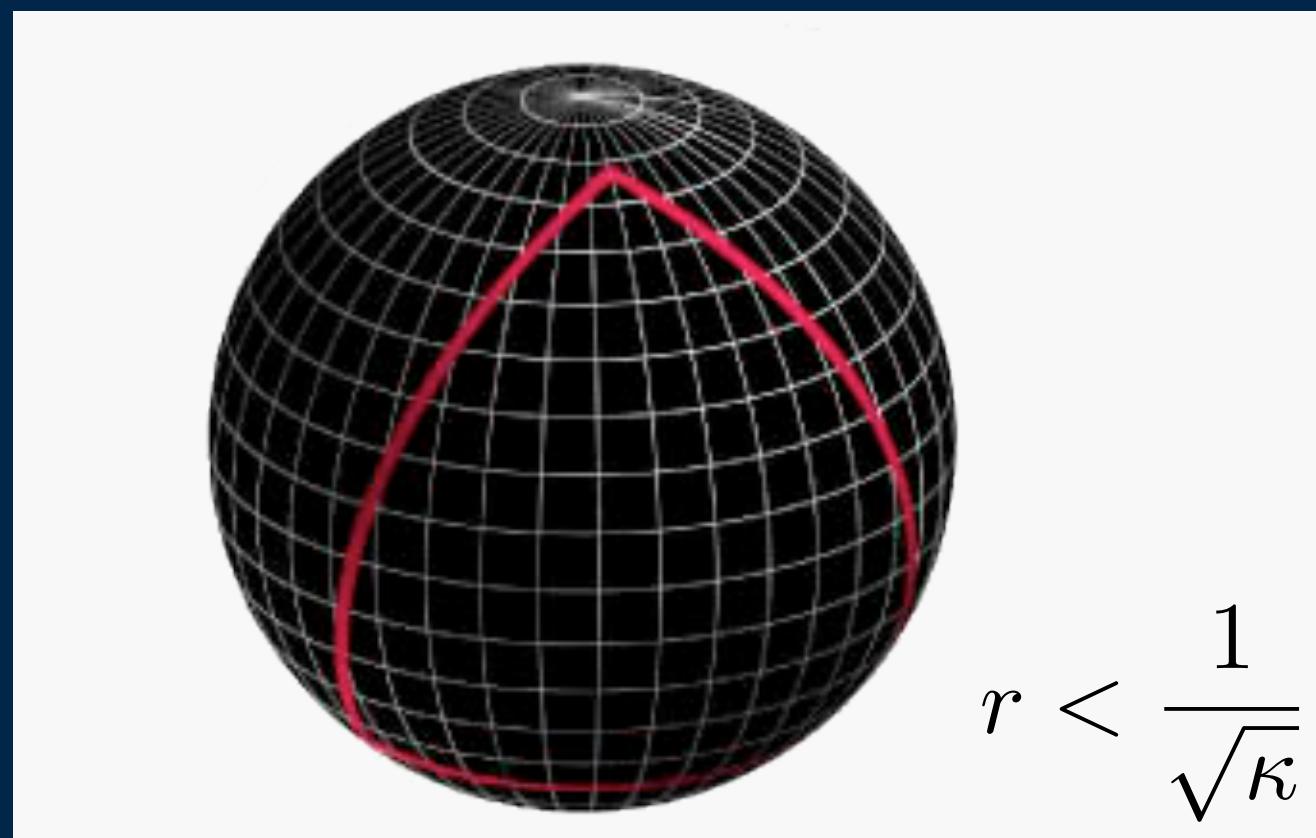


$$ds^2 = -dt^2 + a^2(t) \left[ \frac{dr^2}{1 - \kappa r^2} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

(Spherical coordinates)

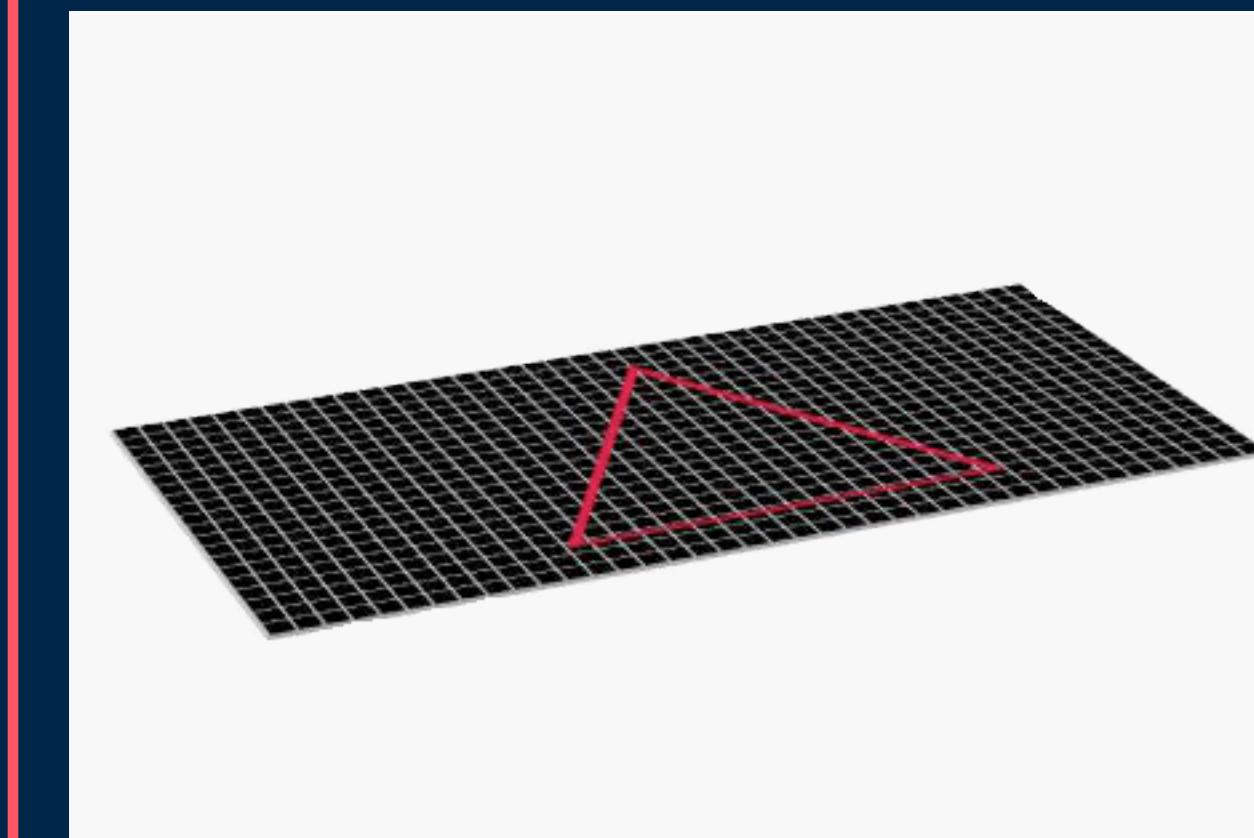
$$\kappa > 0$$

Positively curved (e.g. 3-sphere)



$$\kappa = 0$$

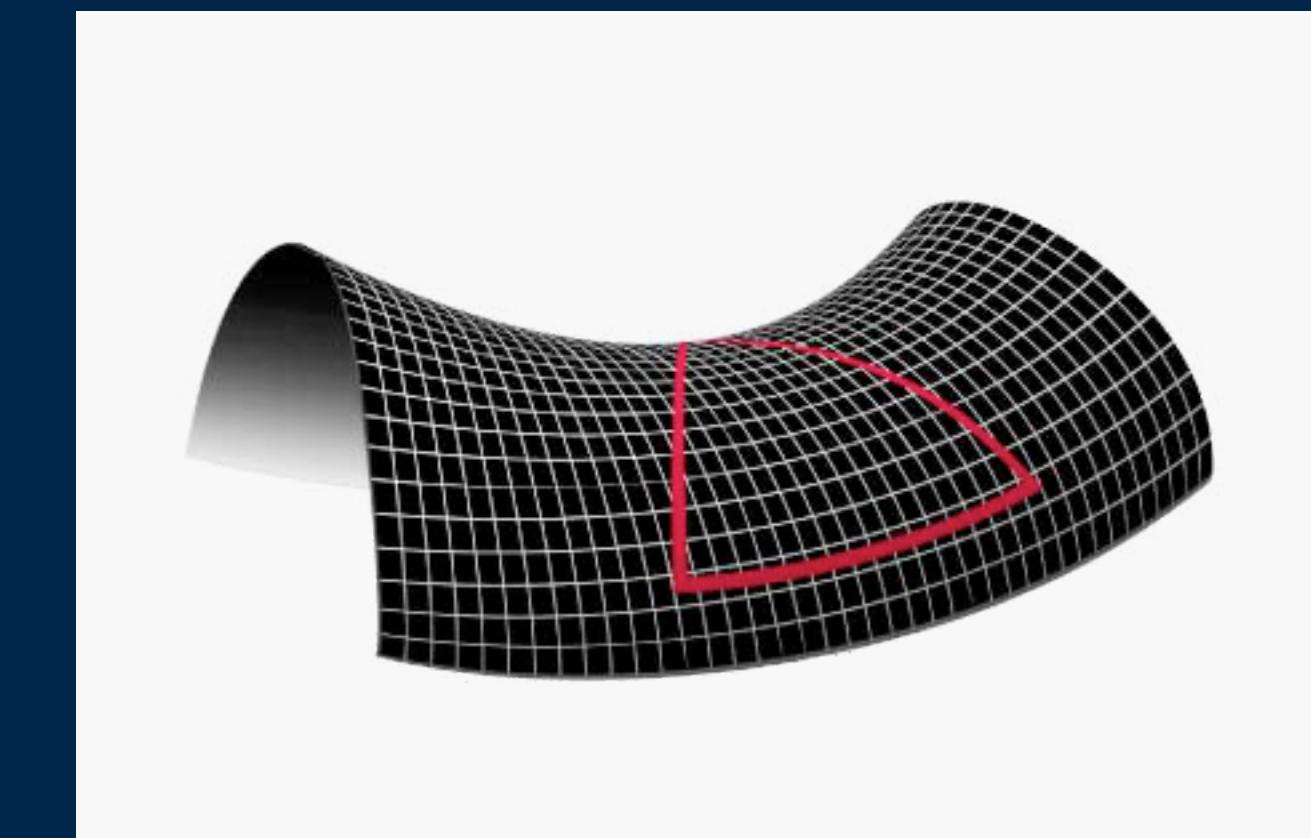
Spatially flat (Euclidean)



Preferred by the data

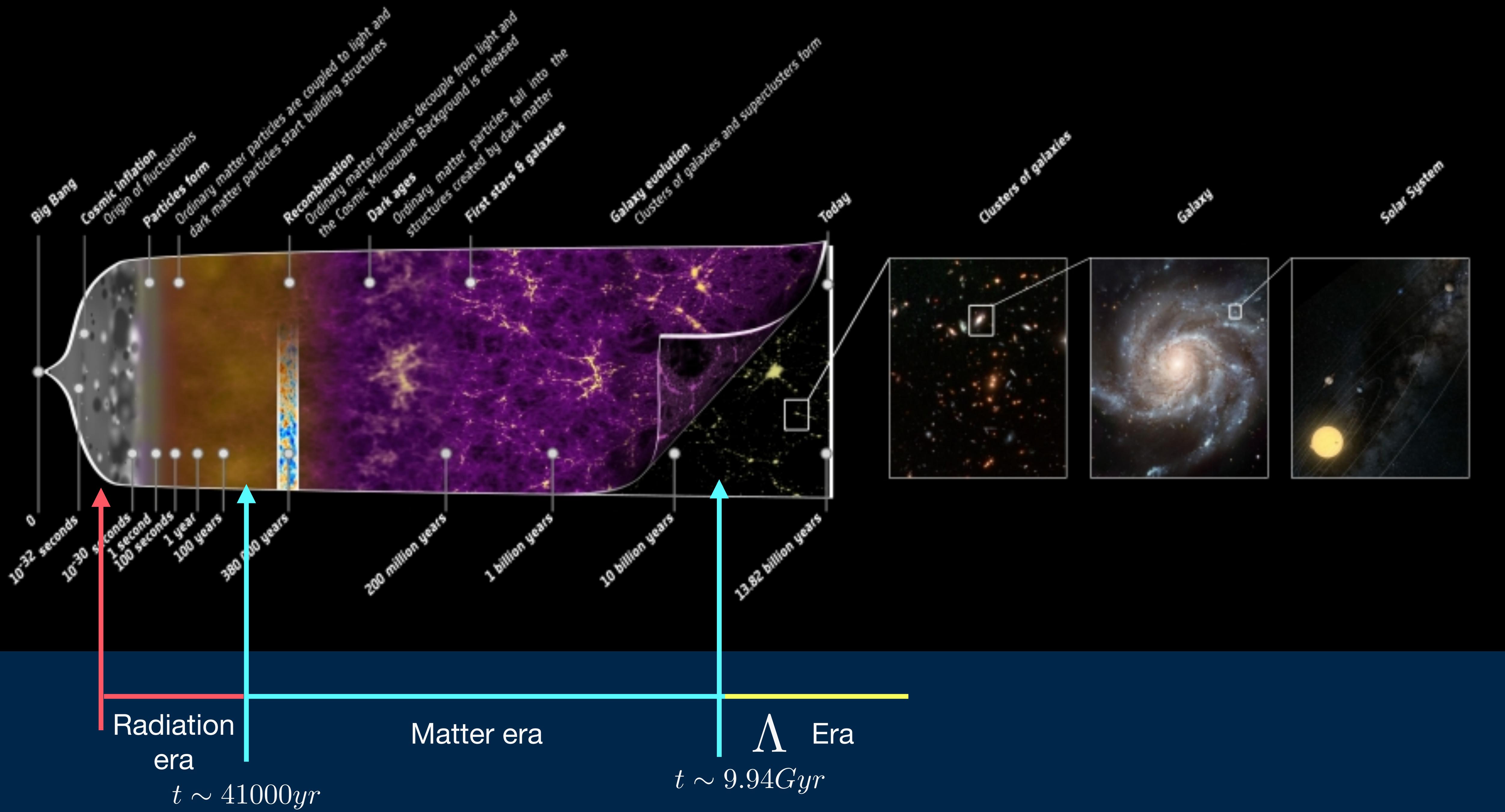
$$\kappa < 0$$

Negatively curved (e.g. hyperbolic 3-space)

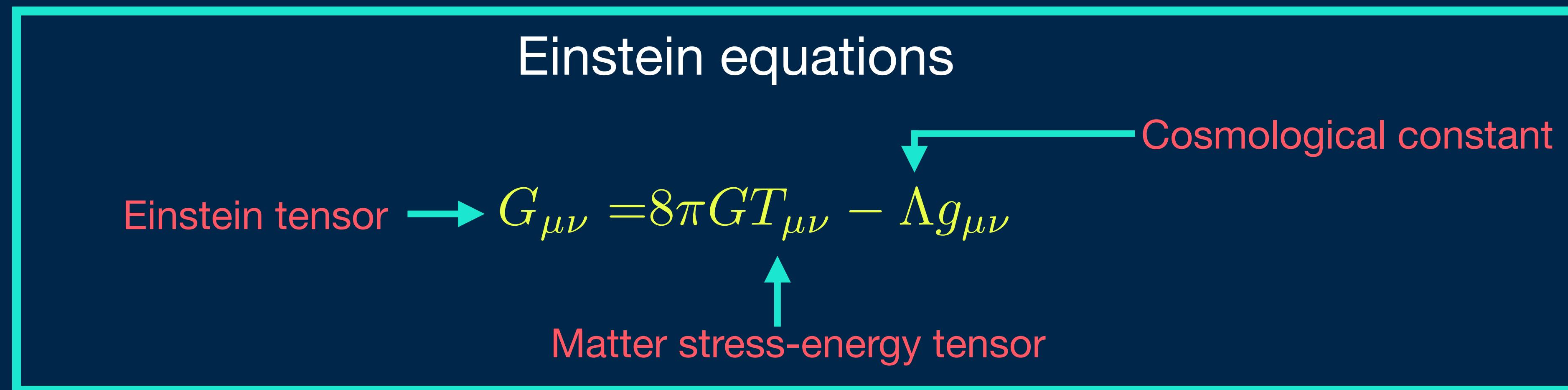


# Dynamics of homogeneous-isotropic Universe

How to determine  $a(t)$



# General Relativity – Gravity as geometry of spacetime



Schematically:

$$g_{\mu\nu} \xrightarrow{\partial} \Gamma_{\alpha\beta}^\mu \xrightarrow{\partial} R_{\mu\nu} \xrightarrow{\quad} G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu}$$
$$R$$

For any metric

$$\nabla_\mu G^\mu_\nu = 0 \quad \xrightarrow{\quad} \quad \text{Compatible with} \quad \nabla_\mu T^\mu_\nu = 0 \quad \xrightarrow{\quad} \quad \text{Stress-energy conservation}$$
$$\nabla_\rho g_{\mu\nu} = 0$$

# Spatially flat FLRW: $\kappa = 0$

(Cartesian coordinates)

Metric:

$$g_{\mu\nu}$$

Christoffel symbols

$$\Gamma_{\mu\nu}^{\alpha} \equiv \frac{1}{2} g^{\alpha\beta} (\partial_{\mu}g_{\beta\nu} + \partial_{\nu}g_{\beta\mu} - \partial_{\beta}g_{\mu\nu})$$

Ricci curvature

$$R_{\mu\nu} = \partial_{\alpha}\Gamma_{\mu\nu}^{\alpha} - \partial_{\nu}\Gamma_{\alpha\mu}^{\alpha} + \Gamma_{\alpha\lambda}^{\alpha}\Gamma_{\mu\nu}^{\lambda} - \Gamma_{\beta\nu}^{\alpha}\Gamma_{\mu\alpha}^{\beta}$$

Scalar curvature

$$R \equiv g^{\mu\nu} R_{\mu\nu}$$



$$ds^2 = -dt^2 + a(t)^2 \delta_{ij} dx^i dx^j$$

$$g_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & a^2 & 0 & 0 \\ 0 & 0 & a^2 & 0 \\ 0 & 0 & 0 & a^2 \end{pmatrix}$$

$$\Gamma_{0j}^i = H \delta_j^i$$

$$\Gamma_{ij}^0 = H g_{ij}$$

Hubble parameter

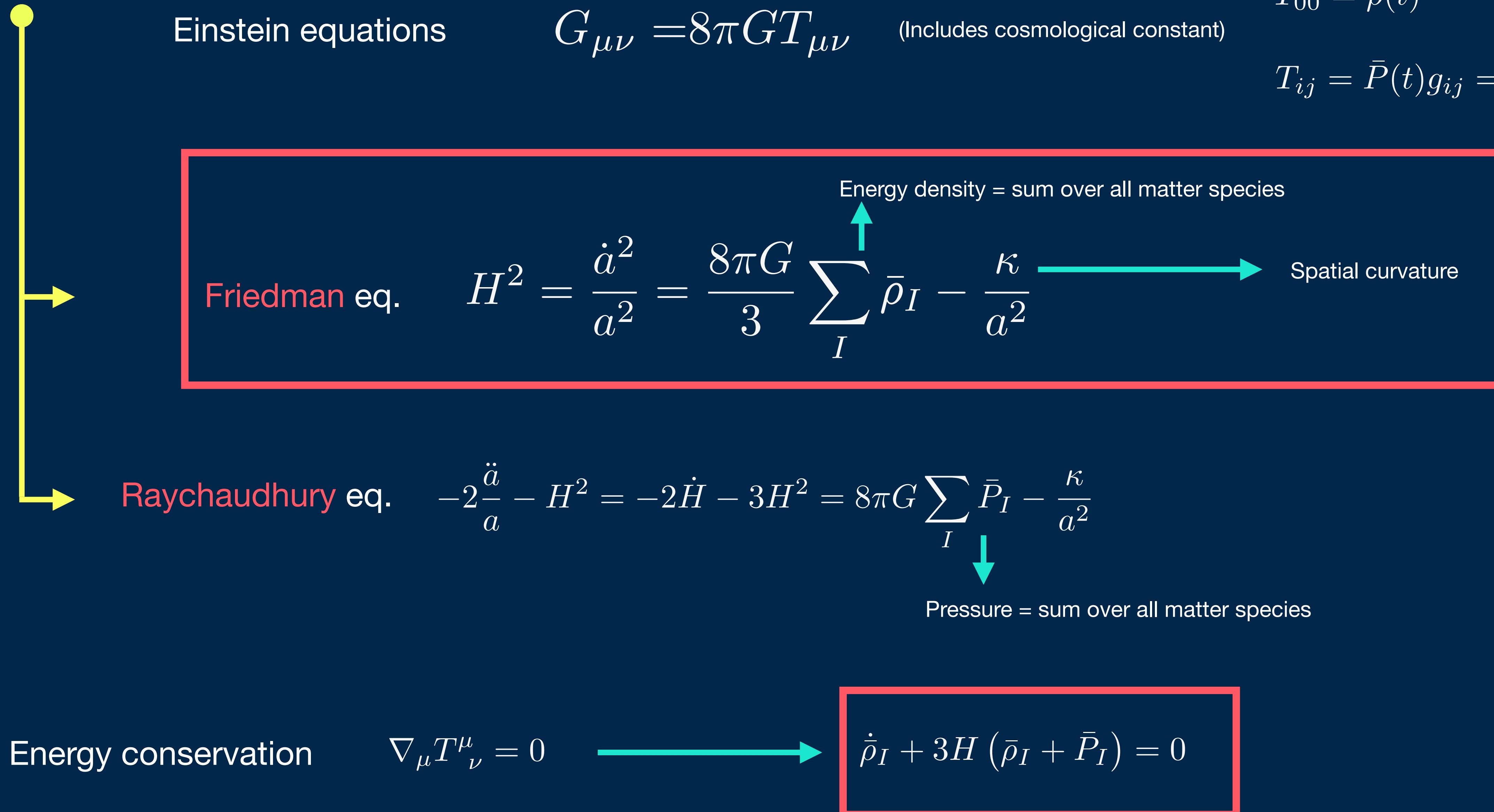
$$H(t) \equiv \frac{\dot{a}}{a}$$

Non-zero components

$$R_{00} = -3\dot{H} - 3H^2$$

$$R_{ij} = [\dot{H} + 3H^2] g_{ij}$$

$$R = 6 [\dot{H} + 2H^2]$$



$I$  Labels species: e.g. baryons, dark matter, photons, neutrinos, cosmological constant, etc

$$H^2 = \frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3} \sum_I \bar{\rho}_I - \frac{\kappa}{a^2}$$

Conventionally  $a = 1$  Today  
 $\bar{\rho}_I(a = 1) = \bar{\rho}_{0,I}$

**Critical density:**  $\bar{\rho}_{\text{crit}} \equiv \frac{3H^2}{8\pi G}$

**Relative density:**  $\Omega_I \equiv \frac{\bar{\rho}_I}{\bar{\rho}_{\text{crit}}} = \frac{8\pi G \bar{\rho}_I}{3H^2}$

Today  $\Omega_0 \equiv \sum_I \Omega_{0,I} \longrightarrow \kappa = (\Omega_0 - 1) H_0^2 = \Omega_K H_0^2$

Negligible curvature  $\Omega_K = -0.0106 \pm 0.0065$  CMB (Planck 2018)

Spatially flat

$$\sum_I \Omega_I = 1$$

# Equation of state: a relation between density and pressure

$$\bar{P} = w \bar{\rho}$$



$$\dot{\bar{\rho}}_I + 3H(\bar{\rho}_I + \bar{P}_I) = 0$$

↓ Solve

$$w = 1$$

Stiff fluid, massless scalar field

$$\rho_{\text{stiff}} \propto a^{-6}$$

$$w = \frac{1}{3}$$

Radiation (photons, massless neutrinos, gravitational waves, ...)

$$\rho_r \propto a^{-4}$$

$$w = 0$$

Pressureless matter, also called dust  
(cool baryons, cold dark matter, galaxies, ...)

$$\rho_m \propto a^{-3}$$

Conventionally  
 $a = 1$   
Today

$$w = -\frac{1}{3}$$

Spatial Curvature, or cosmic strings

$$\rho_{\text{strings}} \propto a^{-2}$$

$$w = -\frac{2}{3}$$

Domain walls

$$\rho_{\text{walls}} \propto a^{-1}$$

$$w = -1$$

Cosmological constant

$$\rho_\Lambda = \frac{\Lambda}{8\pi G} = \text{constant}$$

# Simple solutions

Friedman eq.  $3\frac{\dot{a}^2}{a^2} = 8\pi G \bar{\rho}$    $a(t)$

Today:  $3H_0^2 = 8\pi G \bar{\rho}_0$

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Stiff fluid  $w = 1$   $\bar{\rho}_{\text{stiff}} = \bar{\rho}_0 a^{-6}$   $a = (3H_0 t)^{1/3}$

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Radiation  $w = \frac{1}{3}$   $\bar{\rho}_r = \bar{\rho}_0 a^{-4}$   $a = \sqrt{2H_0 t}$

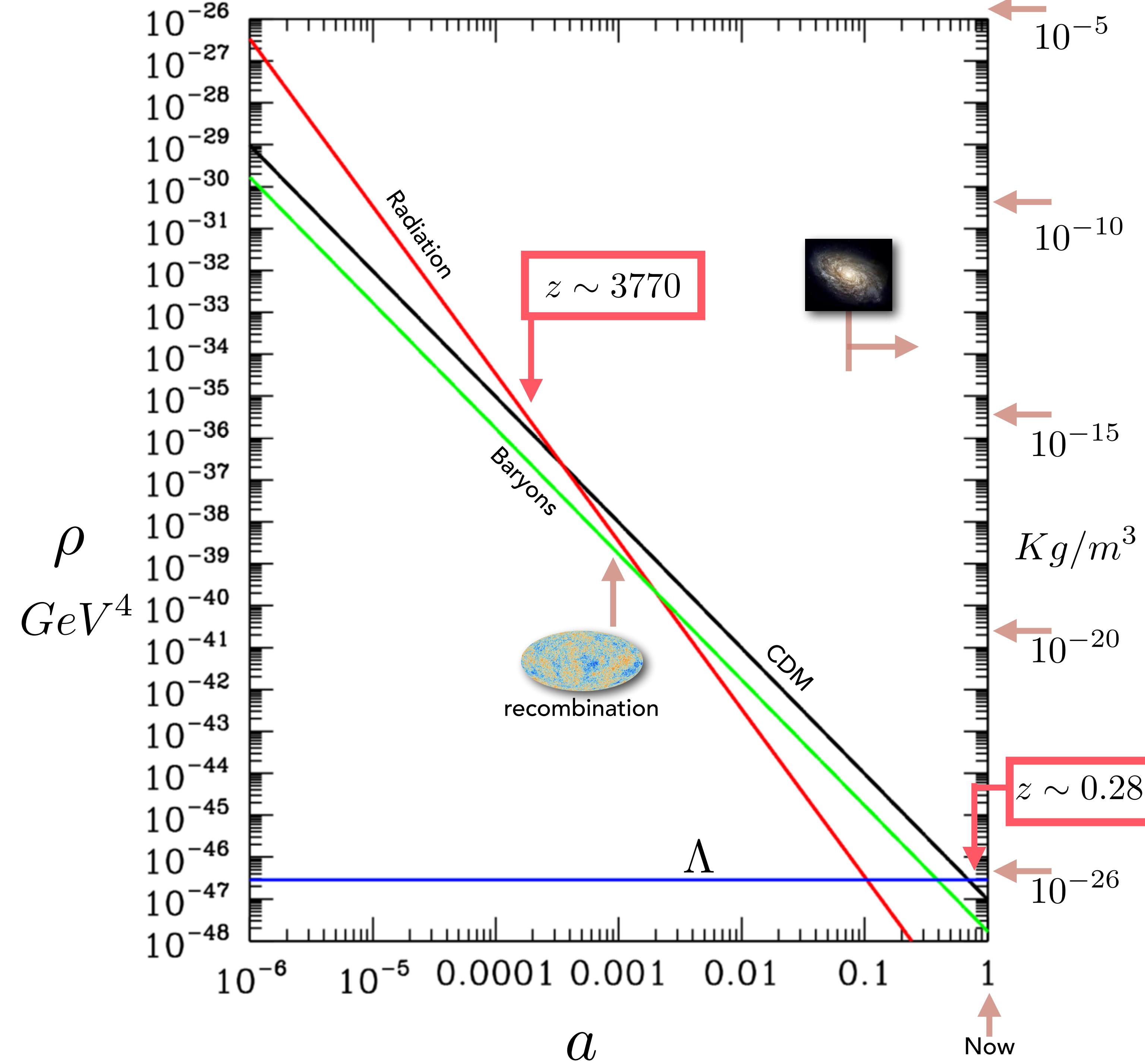
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Pressureless matter  $w = 0$   $\bar{\rho}_m = \bar{\rho}_0 a^{-3}$   $a = \left(\frac{3}{2}H_0 t\right)^{2/3}$

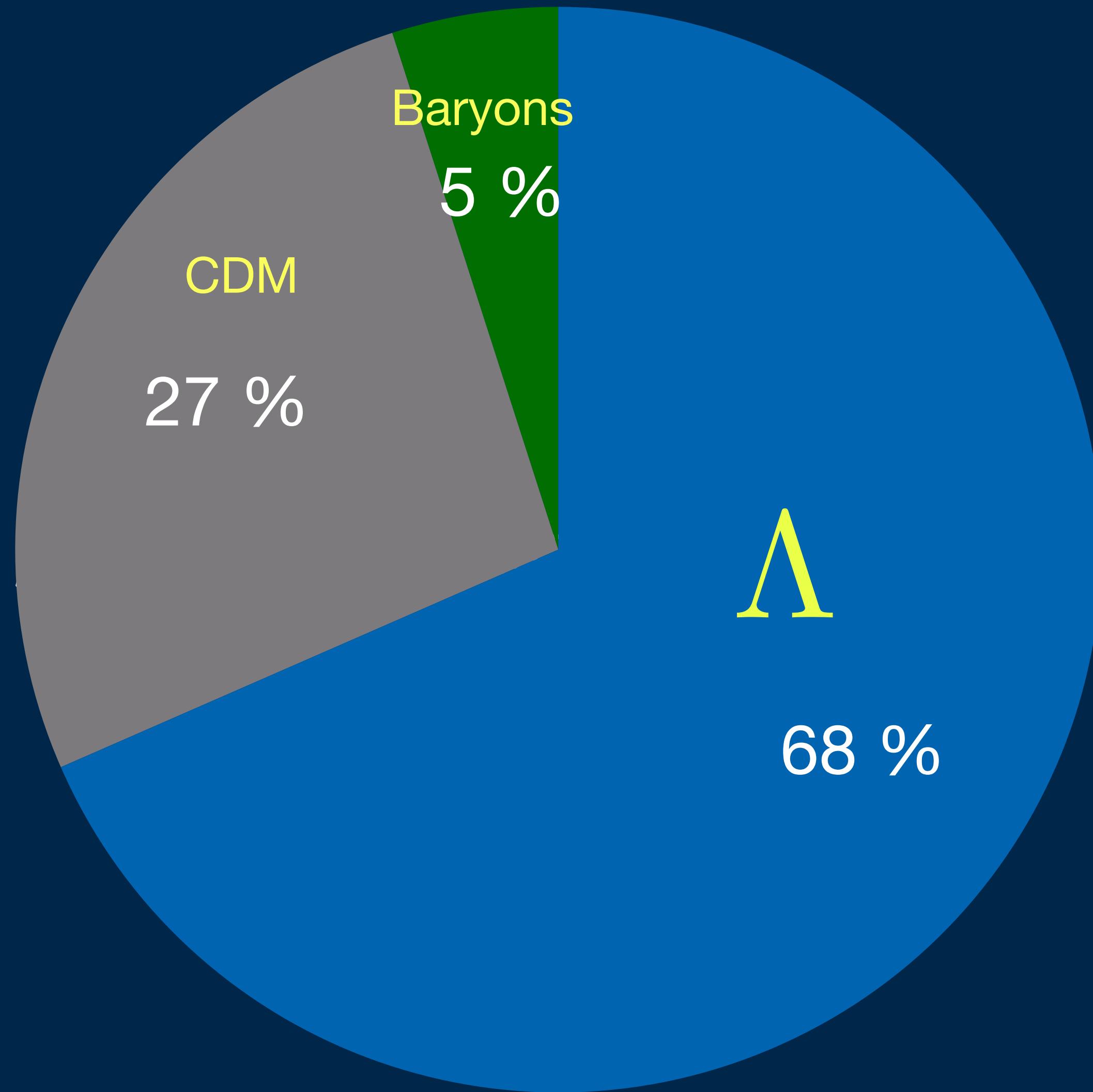
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Cosmological constant  $w = -1$   $\bar{\rho}_\Lambda = \bar{\rho}_0$   $a = e^{H_0 t}$

# Expansion eras

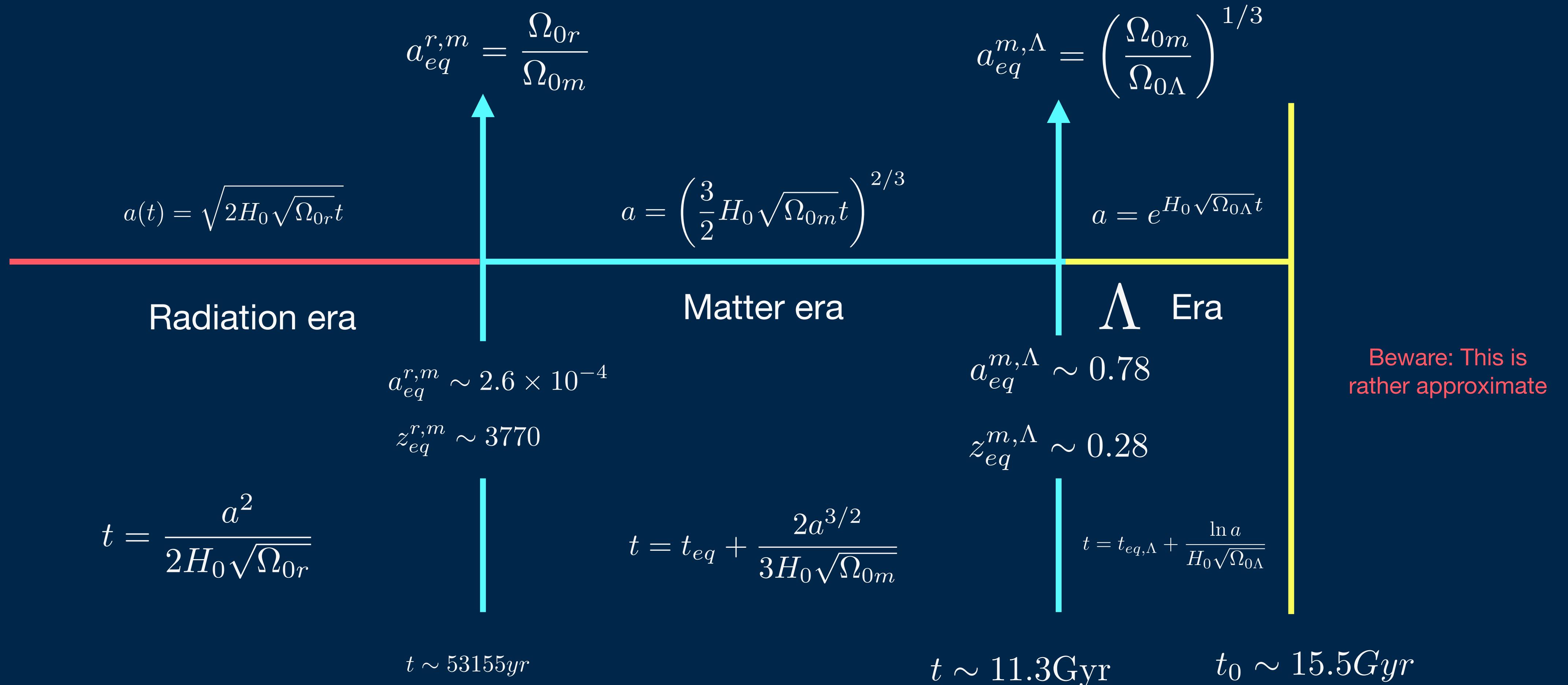


# Material content of the Universe



Hubble constant	$\left\{ \begin{array}{ll} H_0 = 67.36 \pm 0.54 & \text{CMB (Planck 2018)} \\ H_0 = 73.04 \pm 1.04 & \text{Local (SH0ES 2021)} \end{array} \right\}$	“Hubble tension”
Critical density	$\bar{\rho}_{crit,0} = \frac{3H_0^2}{8\pi G} = 1.878 \times 10^{-26} h^2 \text{ kg m}^{-3}$	
Negligible curvature	$\Omega_K = -0.0106 \pm 0.0065$	CMB (Planck 2018)
<b>Assuming flatness, CMB (Planck 2018)</b>		
Matter density	$\Omega_{0,m} = 0.3153 \pm 0.0073$	
Cosmological constant	$\Omega_{0,\Lambda} = 0.6847 \pm 0.0073$	
Cold dark matter	$\Omega_{0,c} h^2 = 0.1202 \pm 0.0014$	
Baryons	$\Omega_{0,b} \sim 0.2658$	
Photons	$\Omega_{0,b} h^2 = 0.02236 \pm 0.00015$	
Neutrinos	$\Omega_{0,b} \sim 0.0494$	
	$\Omega_{0\gamma} \sim 5.04 \times 10^{-5}$	
	$3.43 \times 10^{-5} \lesssim \Omega_{0\nu} \lesssim 0.01$	

# Expansion age of the Universe



Accurate way:

$$t(a_{ref}) = \frac{1}{H_0} \int_0^{a_{ref}} \frac{da}{a \sqrt{\frac{\Omega_{0r}}{a^4} + \frac{\Omega_{0m}}{a^3} + \Omega_{0\Lambda}}} \rightarrow t \sim 9.94 \text{ Gyr} \quad t_0 \sim 13.205 \text{ Gyr}$$

# Thermal history of the Universe

## Cosmic Microwave Background

$$\bar{T}_{CMB} = 2.7255K$$

Best Black body spectrum in the Universe

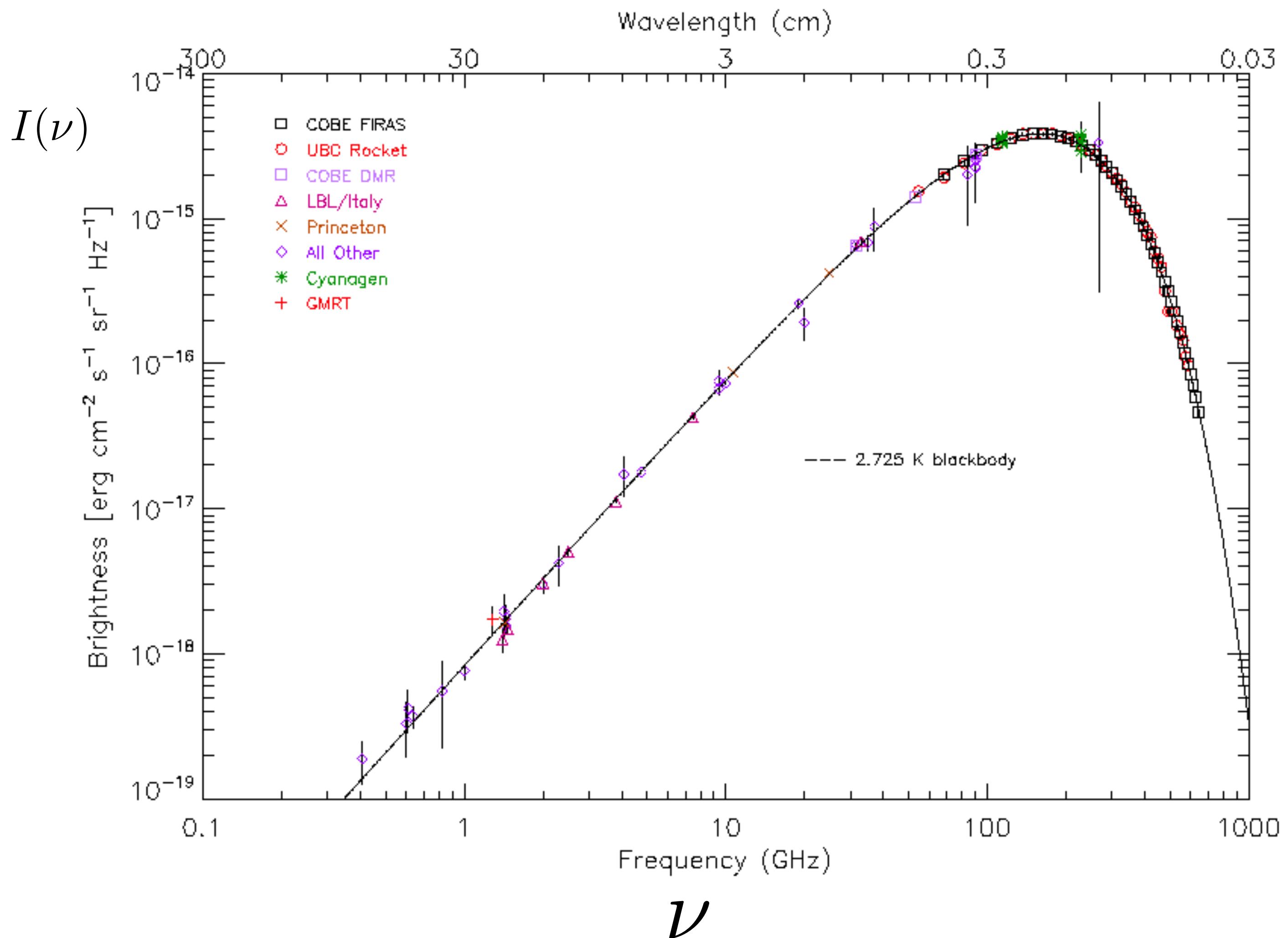
$$I(\nu) = \frac{4\pi\nu^3}{e^{2\pi\nu/T} - 1}$$

CMB at thermal equilibrium today

Expansions preserves Black Body spectrum

CMB at thermal equilibrium in the past

+ anything interacting with it.



Today	$z = 0$	$a = 1$	$t_0 \sim 13.7Gyr$	$T_{cmb} \sim 2.7255K$	$\Lambda$ domination
Distant galaxies	$z = 5$	$a \sim 0.17$	$t \sim 1.15Gyr$	$T_{cmb} \sim 16K$	
Photon decoupling	$z \sim 1100$	$a \sim 9 \times 10^{-4}$	$t \sim 350000yr$	$T_{cmb} \sim 3000K$	Matter domination
Matter-radiation equality	$z \sim 3700$	$a \sim 2.6 \times 10^{-4}$	$t \sim 41000yr$	$T_{cmb} \sim 10000K \sim eV$	
Nucleosynthesis ends	$z \sim 3 \times 10^8$	$a \sim 3 \times 10^{-9}$	$t \sim 200 - 300s$	$T \sim 0.05MeV$	
<b>CMB anisotropies</b>					
$e^+e^-$ annihilation	$z \sim 5 \times 10^9$	$a \sim 2 \times 10^{-10}$	$t \sim 1s$	$T \sim 0.5MeV$	Radiation domination
Neutrino decoupling n/p freeze-out	$z \sim 10^{10}$	$a \sim 10^{-10}$	$t \sim 0.2s$	$T \sim 1 - 2MeV$	
Quark-gluon phase transition	$z \sim 5 \times 10^{12}$	$a \sim 2 \times 10^{-13}$	$t \sim 10^{-5}s$	$T \sim 200MeV$	
Electroweak unification		$a \sim 2 \times 10^{-16}$	$t \sim 10^{-12}s$	$T \sim 100GeV$	?
GUT??		$a \sim 10^{-28}$	$t \sim 10^{-36}s$	$T \sim 10^{15}GeV$	Inflation: Scalar field domination
Quantum gravity??			$t \sim 10^{-43}s$		$V(\phi) \sim \Lambda_{inf}$
			Graviton decoupling		

# Particles at equilibrium

Distribution function:  $f(t, \vec{x}, \vec{p}) \rightarrow f(t, p)$

Describes collection of particles of given energy at time t



Physical momentum

$$E(p) = \sqrt{p^2 + m^2}$$

$$n = n(T)$$

Number density

$$\rho = \rho(T)$$

Energy density

$$P = P(T)$$

Pressure

Degrees of freedom

$$n = \frac{g}{2\pi^2} \int dp \frac{p^2}{e^{E(p)/T} \pm 1}$$



+ Fermions  
- Bosons

$$\rho = \frac{g}{2\pi^2} \int \frac{dp p^2}{e^{E(p)/T} \pm 1} E(p)$$

$$P = \frac{g}{2\pi^2} \int \frac{dp p^2}{e^{E(p)/T} \pm 1} \frac{p^2}{3E(p)}$$

Remember

$$T = \frac{T_0}{a}$$

# 1st law of thermodynamics

$$TdS = dE + PdV$$

↑                      ↑                      ←  
 $S = S(T, V)$  Entropy      Pressure  $P = P(T)$       Volume  $V = a^3$

Energy  $E = \rho V$  = energy density x volume  
 $\rho = \rho(T)$

$\frac{\partial S}{\partial T} = \frac{V}{T} \frac{\partial \rho}{\partial T}$   
 $\frac{\partial S}{\partial V} = \frac{\rho + P}{T}$

Consistency:  $\frac{\partial}{\partial V} \frac{\partial S}{\partial T} = \frac{\partial}{\partial T} \frac{\partial S}{\partial V}$  →  $dP = \frac{\rho + P}{T} dT$

$$dS = d \left[ \frac{(\rho + P)V}{T} \right] = 0$$

Entropy is conserved for equilibrium processes

→ Entropy density  $s = \frac{\rho + P}{T}$

# Relativistic species

$$m \ll T$$

$$E(p) \approx p$$

Number density

$$\rightarrow n = \frac{gT^3}{2\pi^2} \int_0^\infty \frac{x^2}{e^x \pm 1}$$

Energy density

$$\rightarrow \rho = \frac{gT^4}{2\pi^2} \int_0^\infty \frac{x^3 dx}{e^x \pm 1}$$

Pressure



$$P = \frac{1}{3}\rho$$

Entropy density



$$s = \frac{4}{3} \frac{\rho}{T} \propto n$$

$$n_B = \frac{g\zeta(3)}{\pi^2} T^3$$

Bosons

$$n_F = \frac{3}{4}n_B$$

Fermions

$$\rho_B = \frac{\pi^2 g}{30} T^4$$

Bosons

$$\rho_F = \frac{7}{8}\rho_B$$

Fermions

# Relativistic degrees of freedom

- Several relativistic species add together to the radiation energy density
- Relevant only during the radiation era
  - Neutrinos are then approximately massless

Effective massless degrees of freedom

$$\rho_r = \frac{\pi^2 g_*}{30} T_\gamma^4$$
$$g_* = \sum_{bosons,I} g_I \left( \frac{T_I}{T_\gamma} \right)^4 + \frac{7}{8} \sum_{fermions,I} g_I \left( \frac{T_I}{T_\gamma} \right)^4$$

$T \gtrsim 300\text{GeV}$

All standard model particles are relativistic:

$$g_* \approx 106.75$$

$1\text{MeV} \leq T \leq 100\text{MeV}$

Relativistic species are:  $\gamma, 3 \times \nu, e^\pm$   
( $m_{e^\pm} \sim 0.5\text{MeV}$ )

$$g_* = 2 + \frac{7}{8} \times 3 \times 2 + \frac{7}{8} \times 2 \times 2 = 10.75$$

$T \ll 1\text{MeV}$

Relativistic species are:  $\gamma, 3 \times \nu$

$$g_* = 2 + \frac{7}{8} \times 3 \times 2 \times \left( \frac{4}{11} \right)^{4/3} \approx 3.36$$

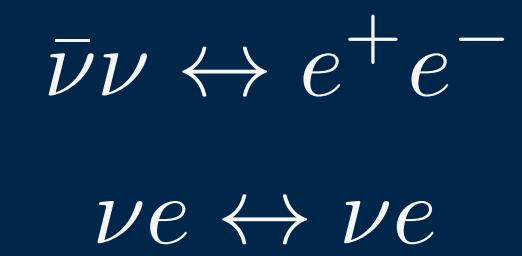
(Neutrinos become non-relativistic inside the matter era)

$N_{eff,\nu} \approx 3.0395$  Due to non-instantaneous neutrino decoupling + QED effects at finite T,  
see Magnano et al., Phys.Lett. B534, 8 (2002)

$$g_* \sim 3.38$$

# Neutrino decoupling

Neutrinos at equilibrium with other species



Interaction rate  $\Gamma = n\langle\sigma v\rangle \sim T^3(G_FT^2) \sim G_F^2 T^5$

Neutrino decoupling when  $\Gamma < H$

(This means that the expansion is faster than the interaction rate and prevents reactions above from occurring efficiently)

$$\frac{\Gamma}{H} \sim \frac{G_F^2 T^5}{\sqrt{8\pi G g_* T^4 / 3}} \sim \frac{G_F^2}{\sqrt{8\pi G g_*/3}} T^3 \sim \left(\frac{T}{1\text{MeV}}\right)^3$$

Neutrino decoupling when  $T < 1\text{MeV}$

Neutrinos maintain the same temperature as the other species but otherwise mind their own business

# $e^+e^-$ annihilation – Neutrino temperature

Below  $T \sim 0.5\text{MeV}$

$$e^+e^- \rightarrow \gamma \quad \text{Heat the photons}$$

Use entropy conservation

$$S_{before} = S_{after}$$

$$\rightarrow s_b a_b^3 = s_a a_a^3$$

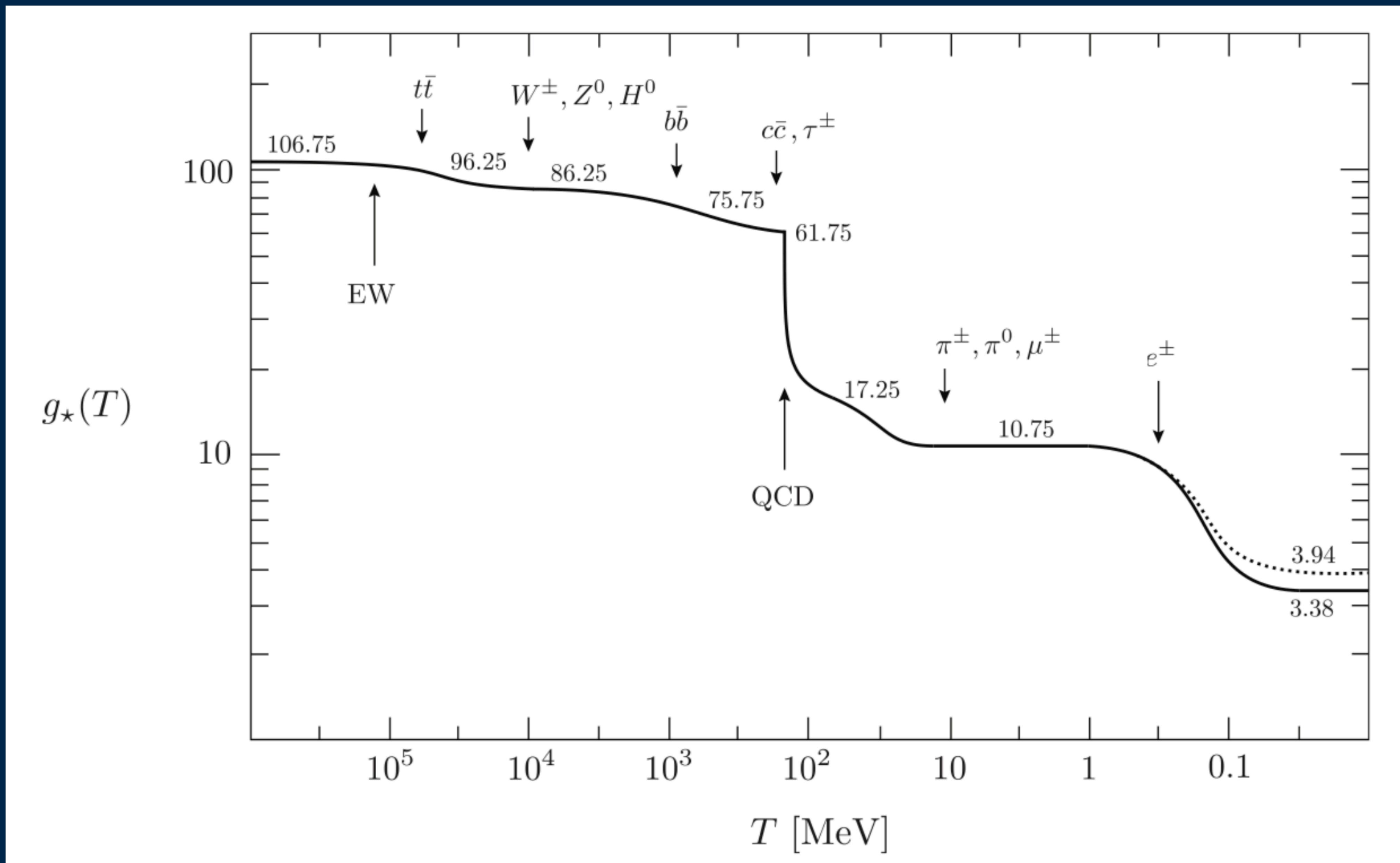
Before						After	
$s \rightarrow$		$\gamma$	$\nu$	$\bar{\nu}$	$e^+$	$e^-$	
2	3	3	2	2			

$$T_\nu = \left(\frac{4}{11}\right)^{1/3} T_\gamma$$

$$T_{\nu, \text{today}} \sim 1.94537K$$

# Relativistic degrees of freedom

From Julien Baur, PhD Thesis, see <https://inspirehep.net/literature/1765226>



# Non-relativistic species

$$m \gg T$$

$$E(p) \approx m$$

Number density →

$$n = g \left( \frac{mT}{2\pi} \right)^{3/2} e^{-m/T}$$

Energy density →

$$\rho = nm$$

Pressure →

$$P = nT \ll \rho \quad (P \approx 0)$$

e.g. massive neutrinos

$$\Omega_\nu = \frac{m_\nu}{94h^2eV}$$

Gershtein, Zel'dovich (1966)

Marx & Szalay (1972)

Cowsik & McClelland (1972)

For  $\Omega_\nu < 1$      $m_\nu \lesssim 45eV$

# Summary

$$\text{Friedman equation} \quad H^2 = \frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3} \sum_I \bar{\rho}_I - \frac{\kappa}{a^2}$$

$$\text{Energy conservation} \quad \dot{\bar{\rho}}_I + 3H(\bar{\rho}_I + \bar{P}_I) = 0$$

**Relative density:**  $\Omega_I \equiv \frac{\bar{\rho}_I}{\bar{\rho}_{\text{crit}}} = \frac{8\pi G \bar{\rho}_I}{3H^2}$

Today:  $\Omega_\Lambda \sim 0.68$      $\Omega_c \sim 0.27$      $\Omega_b \sim 0.05$

**Expansion age**  $t(a_{ref}) = \frac{1}{H_0} \int_0^{a_{ref}} \frac{da}{a \sqrt{\frac{\Omega_{0r}}{a^4} + \frac{\Omega_{0m}}{a^3} + \Omega_{0\Lambda}}}$

# Thermodynamics: entropy conserved

$$m \ll T \quad \text{Relativistic species:} \quad n \sim T^3 \quad \rho \sim T^4 \quad P = \frac{1}{3}\rho \quad s \sim T^3$$

$$\text{Fermions} \quad \times \frac{3}{4} \quad \times \frac{7}{8} \quad \text{relative to bosons}$$