

# Cosmology

## Lecture 3: Large scale structure and Cosmic Microwave Background



**FZU**

Institute of Physics of the  
Czech Academy of Sciences

**ceico**

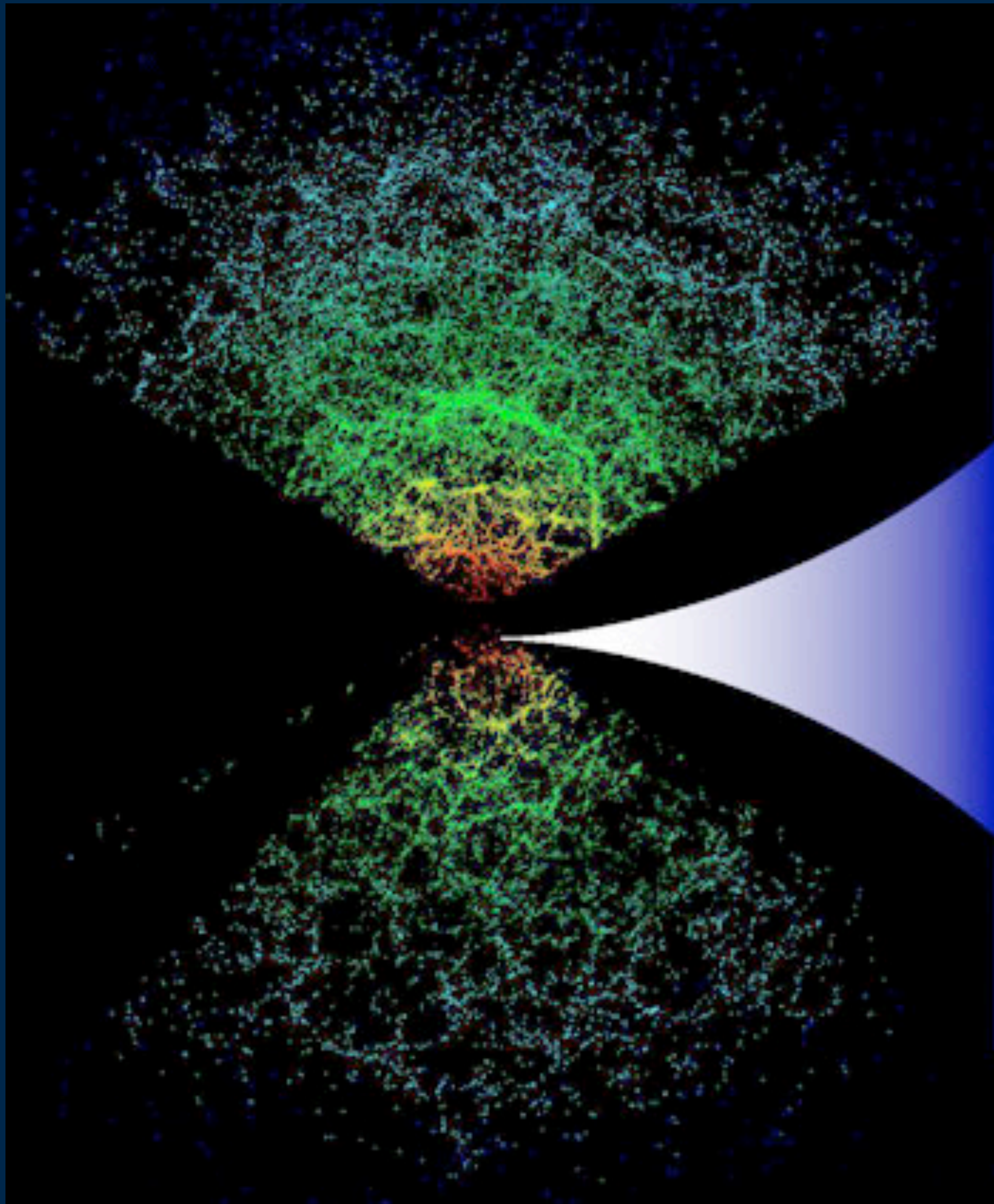
CENTRAL EUROPEAN INSTITUTE FOR  
COSMOLOGY AND FUNDAMENTAL PHYSICS



EUROPEAN UNION  
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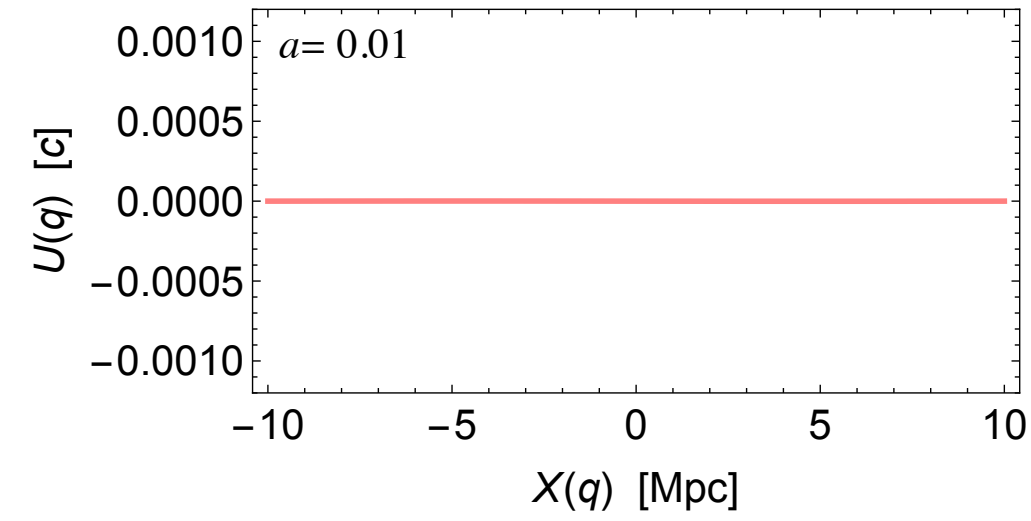


MINISTRY OF EDUCATION,  
YOUTH AND SPORTS

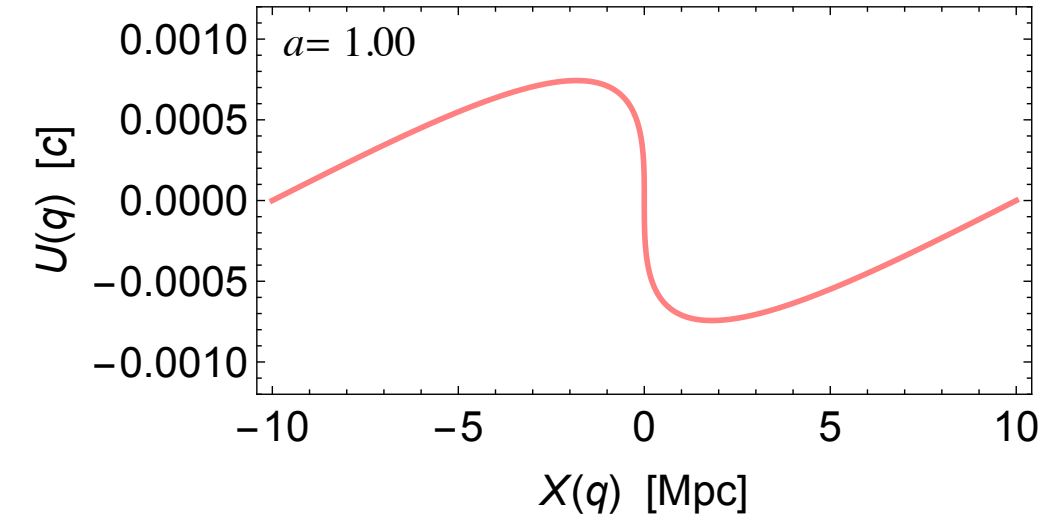


(from Sloan Digital Sky Survey)

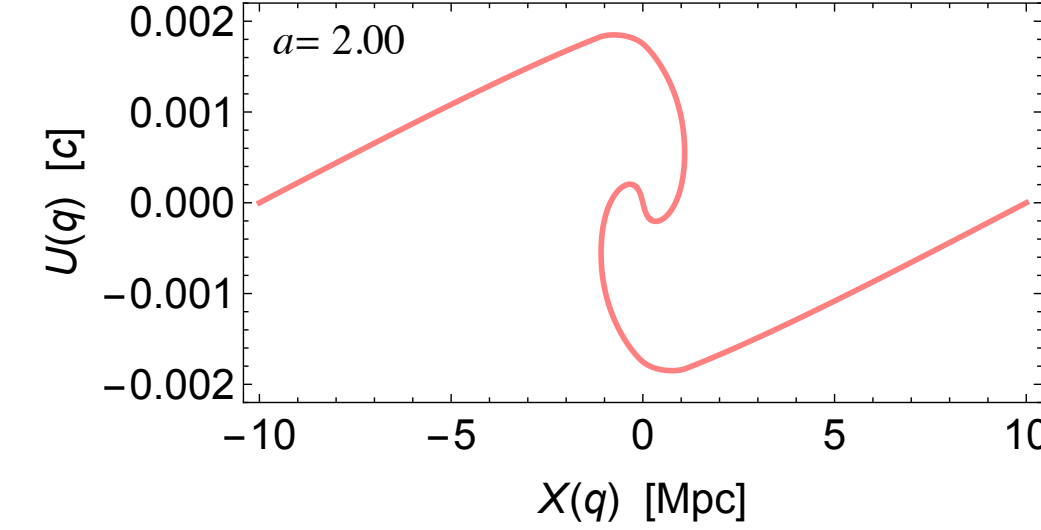
Initial condition



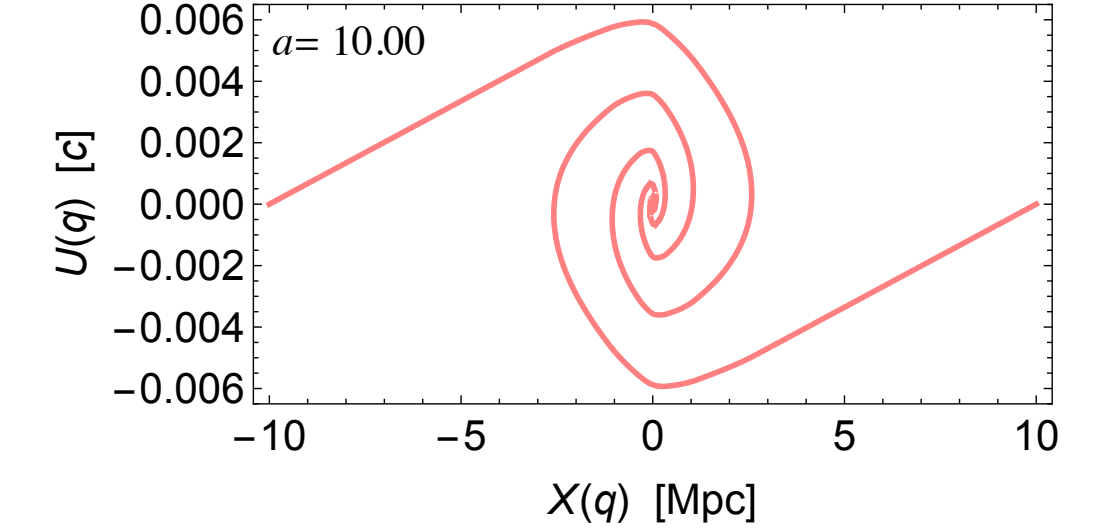
Shell crossing



Phase mixing



Further phase mixing

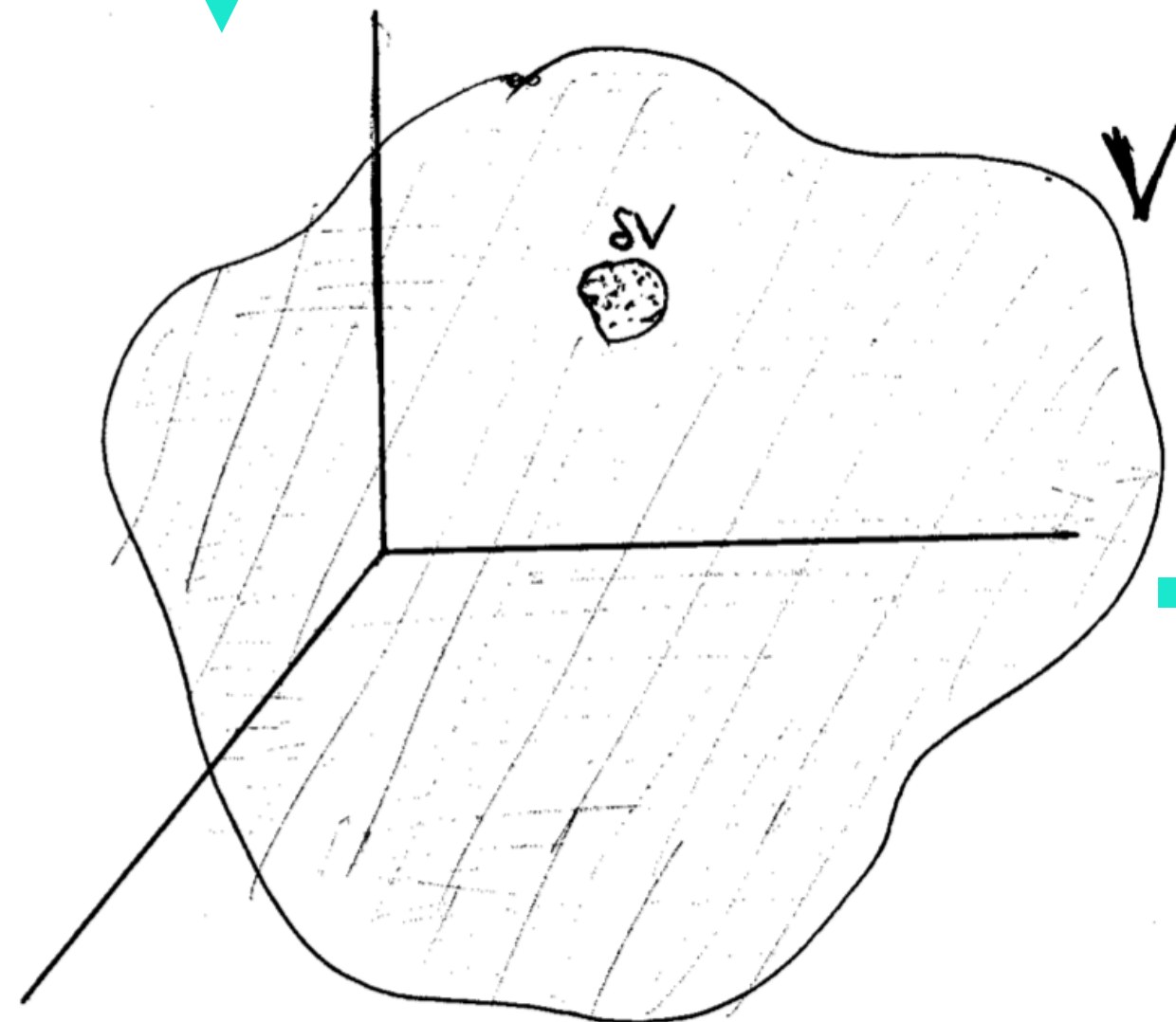
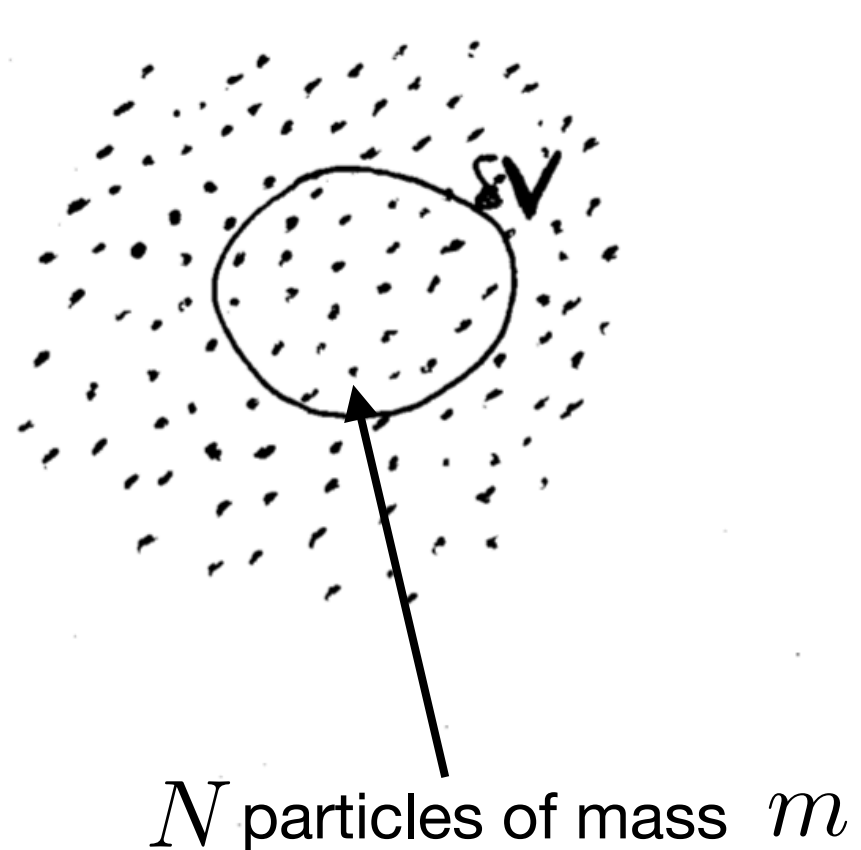


### Fluid approximation:

- Particle trajectories do not cross (no collisions)
- No phase mixing
- Any small  $\delta V$  contains large number of particles

Use N-body simulations, Vlasov solvers, Schoedinger-Poisson Method

- Gadget: <https://wwwmpa.mpa-garching.mpg.de/gadget4/>
- Enzo: <https://enzo-project.org/>
- Arepo: <https://arepo-code.org/>
- RAMSES: <https://www.ics.uzh.ch/~teyssier/ramses/RAMSES.html>



- Energy density field
- Velocity field

$$\rho(t, \vec{x}) = \lim_{\delta V \rightarrow 0} \frac{Nm}{\delta V}$$

$$\vec{v}(t, \vec{x})$$

# Newtonian Gravitational Collapse

$$\frac{dM}{dt} = \int dV \frac{\partial \rho}{\partial t}$$

**Continuity equation**  
(Mass change through closed surface)

$$\frac{\partial \rho}{\partial t} + \vec{v} \cdot \vec{\nabla} \rho = -\rho \vec{\nabla} \cdot \vec{v}$$

**Euler eq.**

(Newton-II: Acceleration = sum forces)

$$\rho \left[ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right] = -\vec{\nabla} P - \rho \vec{\nabla} \Phi$$

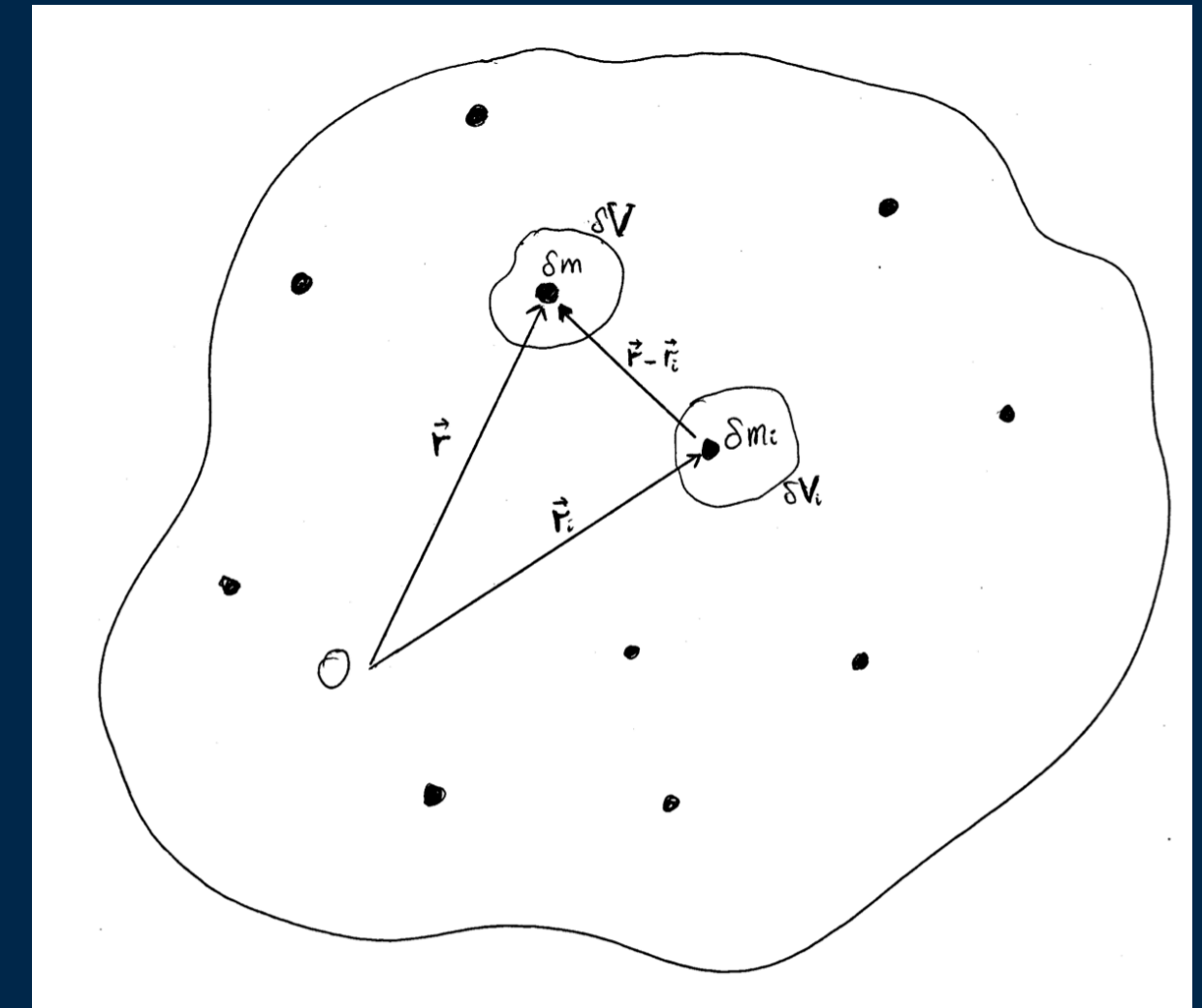
**Poisson eq.**  
(Gravity)

$$\nabla^2 \Phi = 4\pi G \rho$$

**Equation of state**

$$\vec{\nabla} P = C_s^2 \vec{\nabla} \rho$$

$$\frac{\delta F}{\delta V}(r) \sim -G \frac{\delta M}{\delta V}(r) \sum_i \frac{\delta m_i}{|r - r_i|^2} \sim \frac{\delta \Phi}{\delta V}$$



Use a trick — Taylor expansion

$$\rho = \bar{\rho} + \epsilon \delta \rho$$

$$\epsilon \ll 1$$

Set  $\epsilon^2 \rightarrow 0$

Equations become linear

$$\vec{v} = \bar{\vec{v}} + \epsilon \delta \vec{v}$$

background is homogeneous and static:

$$\frac{\partial \bar{\rho}}{\partial t} = \vec{\nabla} \bar{\rho} = 0 \quad \bar{\vec{v}} = 0$$

# Jeans instability

Define density contrast:  $\delta \equiv \frac{\delta\rho}{\rho}$

Use Fourier space  $\vec{\nabla}^2 \rightarrow -k^2$

$$\ddot{\delta} + (k^2 C_s^2 - 4\pi G \bar{\rho}) \delta = 0$$

$$k^2 C_s^2 > 4\pi G \bar{\rho}$$

Oscillations (stable)

$$k^2 C_s^2 < 4\pi G \bar{\rho}$$

Exponential collapse

Dividing line:  $k = \frac{\sqrt{4\pi G \bar{\rho}}}{C_s}$

Jeans length:  $\lambda_J = \frac{2\pi}{k} = C_s \sqrt{\frac{\pi}{G \bar{\rho}}}$

Only **wavelengths larger** than the **Jeans length** can **collapse** to form bound objects

# Collapse in an expanding Universe

Use Newtonian theory  
suffices for sub-horizon fluctuations  
for non-relativistic matter

Continuity equation	$\frac{\partial \rho}{\partial t} + \vec{v} \cdot \vec{\nabla} \rho = -\rho \vec{\nabla} \cdot \vec{v}$
Euler eq.	$\rho \left[ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right] = -\vec{\nabla} P - \rho \vec{\nabla} \Phi$
Poisson eq. (Gravity)	$\nabla^2 \Phi = 4\pi G \rho$
Equation of state	$\vec{\nabla} P = C_s^2 \vec{\nabla} \rho$

(Same eps as before)

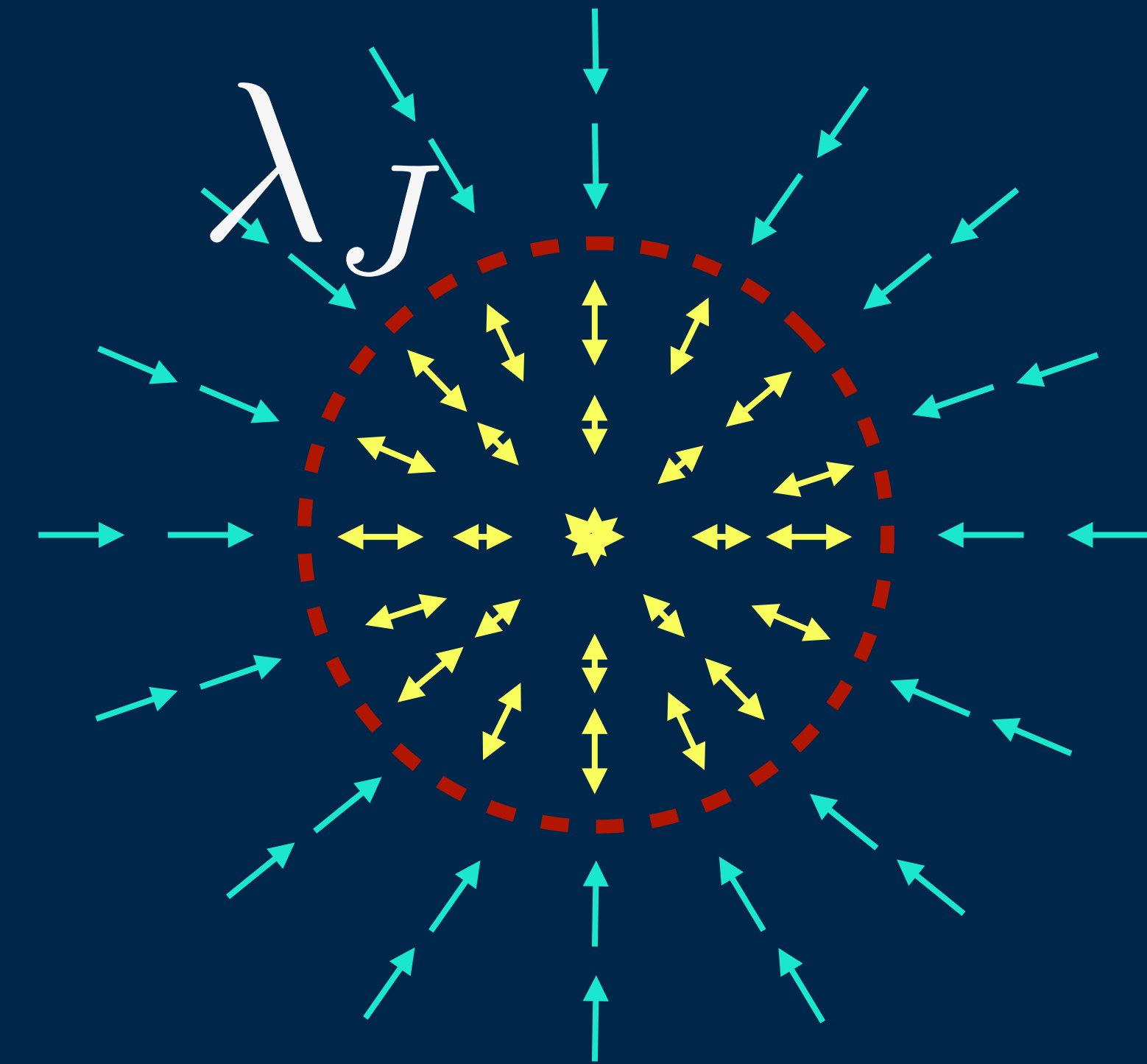
Taylor-expand  $\rho = \bar{\rho} + \epsilon \delta \rho$   
 $\vec{v} = \bar{\vec{v}} + \epsilon \delta \vec{v}$   $\epsilon \ll 1$   $\xrightarrow{\text{Set } \epsilon^2 \rightarrow 0}$  Equations become linear

Background is homogeneous and expanding:  $\bar{\vec{v}} = H \bar{\vec{x}}$

FRW energy conservation!  $\dot{\bar{\rho}} + 3H \bar{\rho} = 0$

# Collapse in an expanding Universe

→ 
$$\ddot{\delta} + 2H\dot{\delta} + \left( \frac{k^2 C_s^2}{a^2} - 4\pi G\bar{\rho} \right) \delta = 0$$



- Expansion introduces a damping term:  $2H\dot{\delta}$

- Scales are stretched by  $a(t)$

- Background density is evolving  $\bar{\rho} \propto a^{-3}$

- Jeans length is time dependent

$$\lambda_J = \frac{C_s}{a} \sqrt{\frac{\pi}{G\bar{\rho}}}$$

- A given perturbation may switch between growth and stasis

Differences arise

# Matter fluctuations: Growth or decay

**Within** Jeans length: **oscillations** (Bessel functions with decreasing amplitude)

**Outside** Jeans length:  $k \ll k_J$  ignore  $k^2 C_s^2$

1. **Matter domination:**  $a \sim t^{2/3}$        $8\pi G\bar{\rho} \sim \frac{1}{t^2}$

→  $\ddot{\delta} + \frac{4}{3t}\dot{\delta} - \frac{2}{3t^2}\delta = 0$       →  $\delta = \delta_0 t^{2/3} + \frac{\delta_1}{t}$       Power-law growth

2. **Radiation domination:**  $a \sim \sqrt{t}$        $8\pi G\bar{\rho} = 3H^2\Omega_m \rightarrow 0$

→  $\ddot{\delta} + \frac{1}{t}\dot{\delta} = 0$        $\delta = \delta_0 + \delta_1 \ln t$

Mészáros effect: During radiation era matter fluctuations grow only logarithmically at best



# Can baryons form structure?

It all depends on two things:

- the Baryon Jeans mass
- the age of the Universe

- Simplified assumptions:**
- only species are baryons+photons
  - recombination occurs at equality

**Before recombination:**  $a \approx \Omega_{0\gamma}^{1/4} \sqrt{2H_0 t}$  and  $c_s^2 = \frac{1}{3}$

Baryon Jeans length:  $\lambda_J = \sqrt{\frac{\pi}{3G\rho_b}} = \frac{2\pi}{3H_0} \sqrt{\frac{2}{\Omega_{0b}}} \Omega_{0\gamma}^{3/8} (2H_0 t)^{3/4}$

Horizon diameter  $\lambda_H = 4t$

$$\frac{\lambda_J}{\lambda_H} \propto \frac{1}{(H_0 t)^{1/4}}$$

As  $t \rightarrow 0 \rightarrow \lambda_J \gg \lambda_H$   
 $t \rightarrow t_* \rightarrow \frac{\lambda_J}{\lambda_H} \rightarrow \frac{\pi\sqrt{2}}{3} > 1$

Therefore before recombination  $\lambda_J > \lambda_H \rightarrow$  baryons cannot collapse to form structures

**After recombination**  $C_s^2 \rightarrow 0 \rightarrow \lambda_J \ll \lambda_H$  and baryons can collapse into structures

# Structure formation

## Sub-horizon fluctuations of uncoupled non-relativistic species

- Power-law growth during matter domination as  $t^{2/3}$
- Very small logarithmic growth during radiation domination

## Can baryons form structure by themselves?

YES, but... with a delay.

The universe is not old enough to produce the structure we see. A baryon-only universe would have looked very different.

## One way out: **Cold Dark Matter**\*

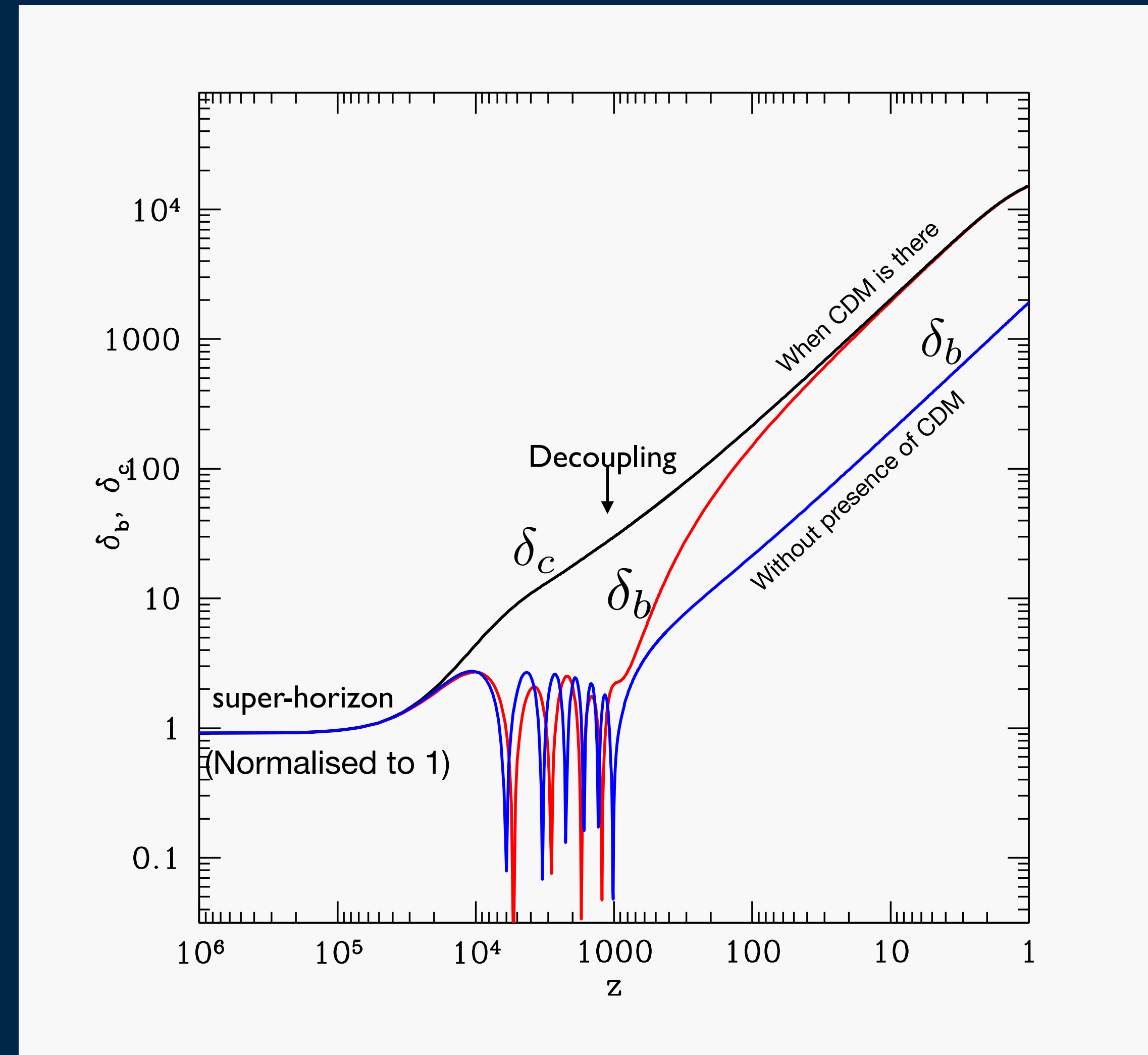
CDM does not couple to photons and gets a head start.

Once baryons decouple they follow the potential wells already created by CDM

$$\vec{\nabla}^2 \Phi = 4\pi G a^2 \bar{\rho}_m (\Omega_b \delta_b + \Omega_c \delta_c)$$

(\* another possibility is warm dark matter)

(\* yet another, possibly, extending General Relativity)



# Probes of Large Scale Structure

Correlation function

$$\xi(\vec{r}) = \int d^3 r' \langle \delta(\vec{r}) \delta(\vec{r}' + \vec{r}) \rangle$$

Gives the correlation of fluctuations separated by  $\vec{r}$

Matter power spectrum

$$\langle \delta_{\vec{k}} \delta_{\vec{k}'} \rangle = (2\pi)^3 P(k) \delta^{(3)}(\vec{k} - \vec{k}')$$

Correlation function is the Fourier transform of the Power spectrum

$$\xi(\vec{r}) = \int d^3 k e^{-i\vec{k}\cdot\vec{r}} P(k)$$

How do we estimate.  $P(k)$  ?

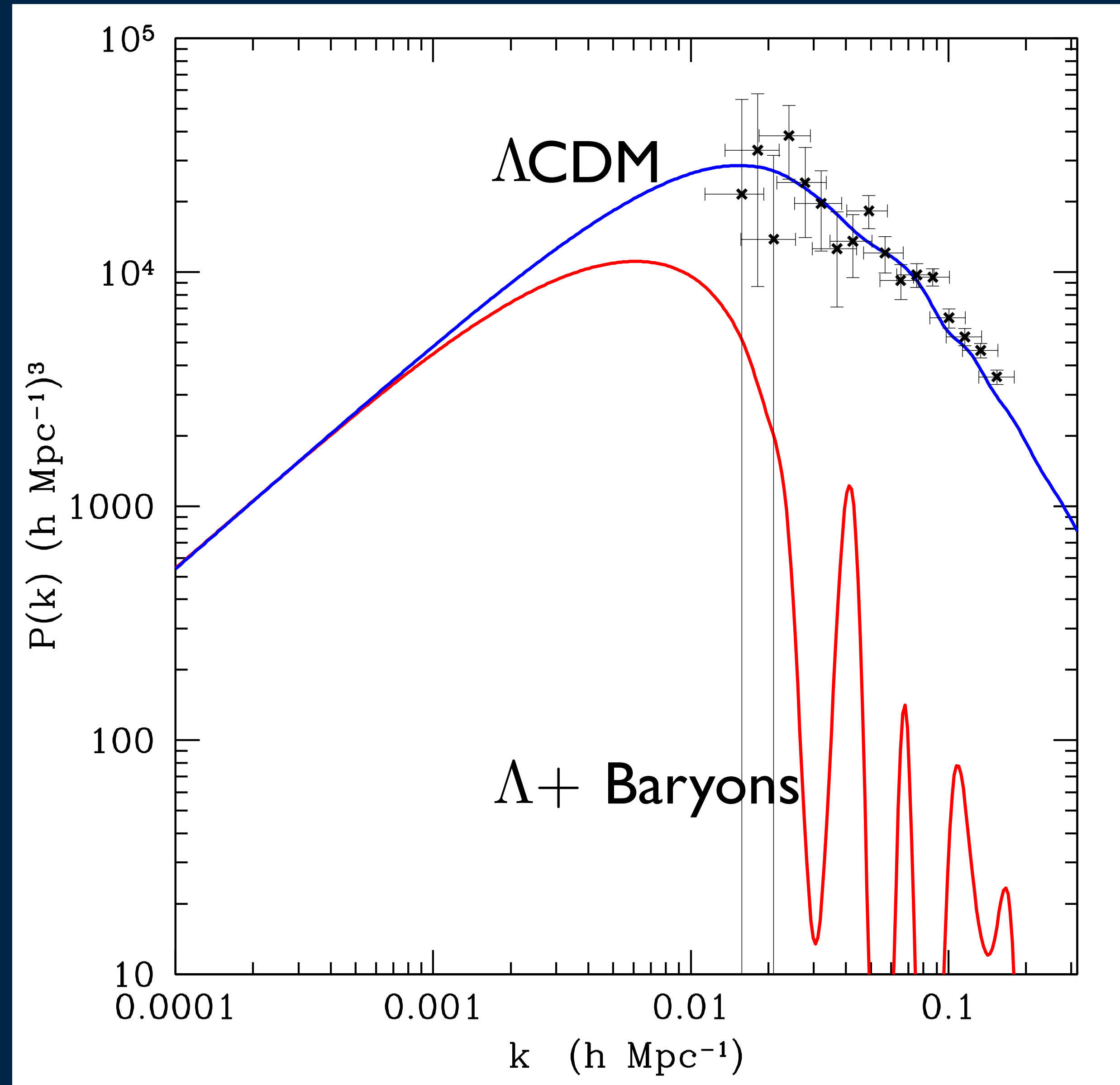
Can use the distribution of galaxies and their velocities

- Set of Power spectra
- Galaxy-Galaxy  $P_{gg}(k)$
  - Galaxy-Velocity  $P_{gv}(k)$
  - Velocity-Velocity  $P_{vv}(k)$

But galaxies may not trace the underlying matter field exactly (light does not follow mass)

$$\text{Bias } b \quad \text{s.t.} \quad \delta_g = b\delta \quad P_{gg}(k) = b^2 P(k)$$

# The matter power spectrum



# Neutrinos and Large Scale Structure

Collision-less particles stream out of over-dense regions and into under-dense regions (**free-streaming**)

Massive neutrinos undergo collision-less damping

$$\lambda_{FS} \approx 20 Mpc \times \frac{30 eV}{m_\nu}$$



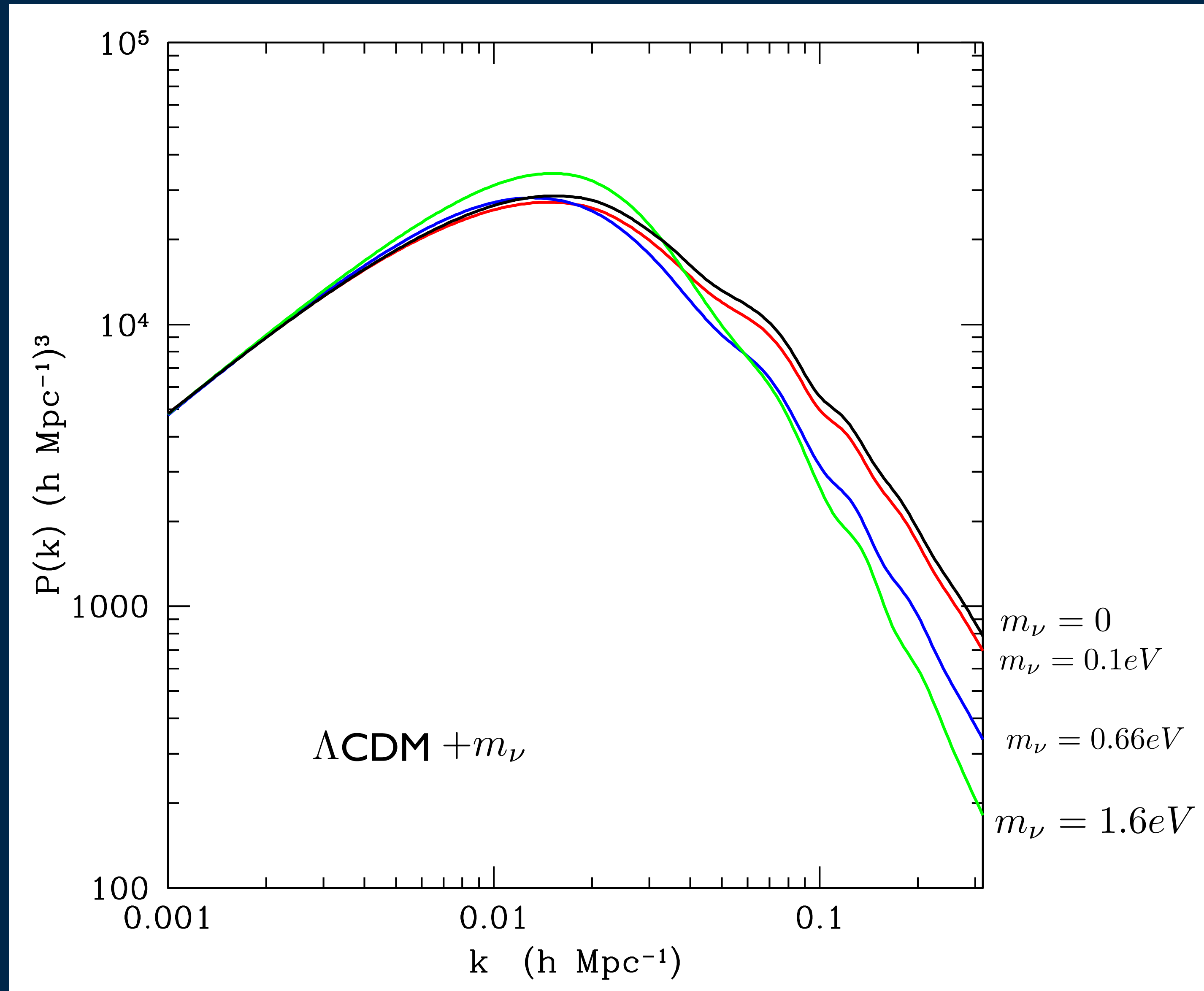
$$k_{FS} \approx 0.28 \frac{m_\nu}{1 eV} Mpc^{-1}$$

Structure formation is suppressed

Current limits:

$$\sum m_\nu < 0.24 eV \quad (\text{Planck 2018})$$

$$\sum m_\nu < 0.12 eV \quad \text{CMB} + \text{BAO} \quad (\text{Planck 2018})$$

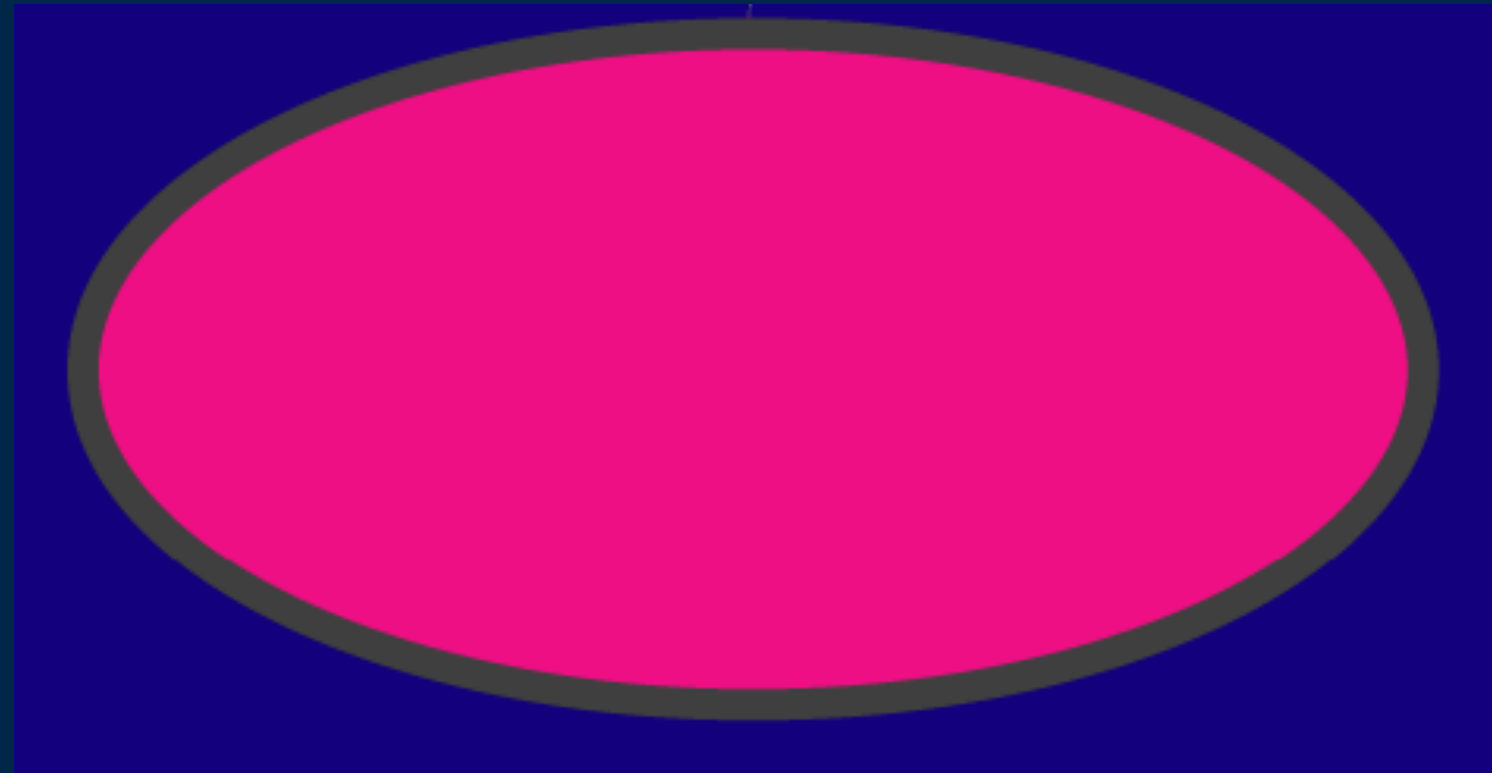


# CMB anisotropies

# CMB anisotropies (boring parts)

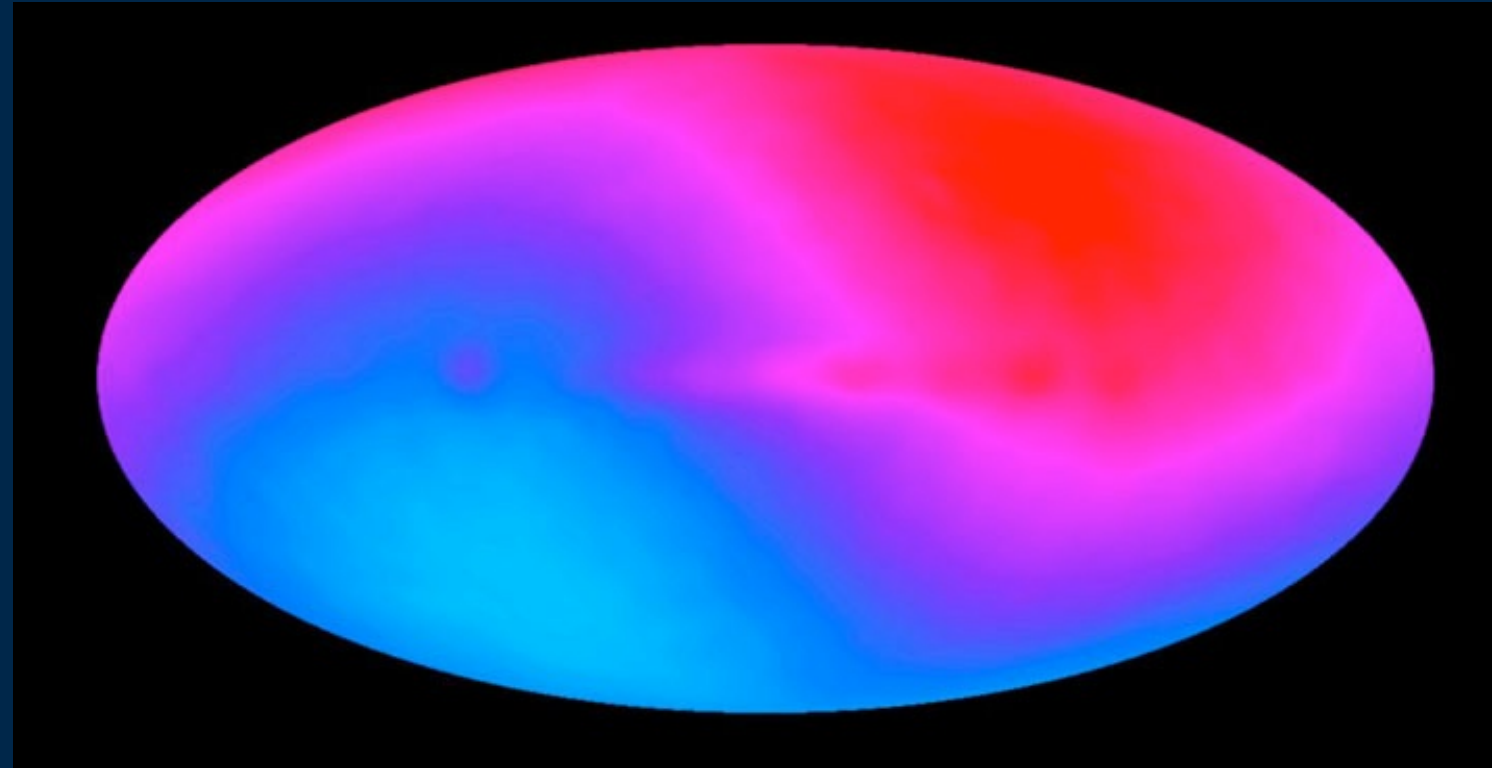
Monopole

$$\bar{T}_{CMB} = 2.7255K$$



Isotropic part,  
same everywhere

Dipole

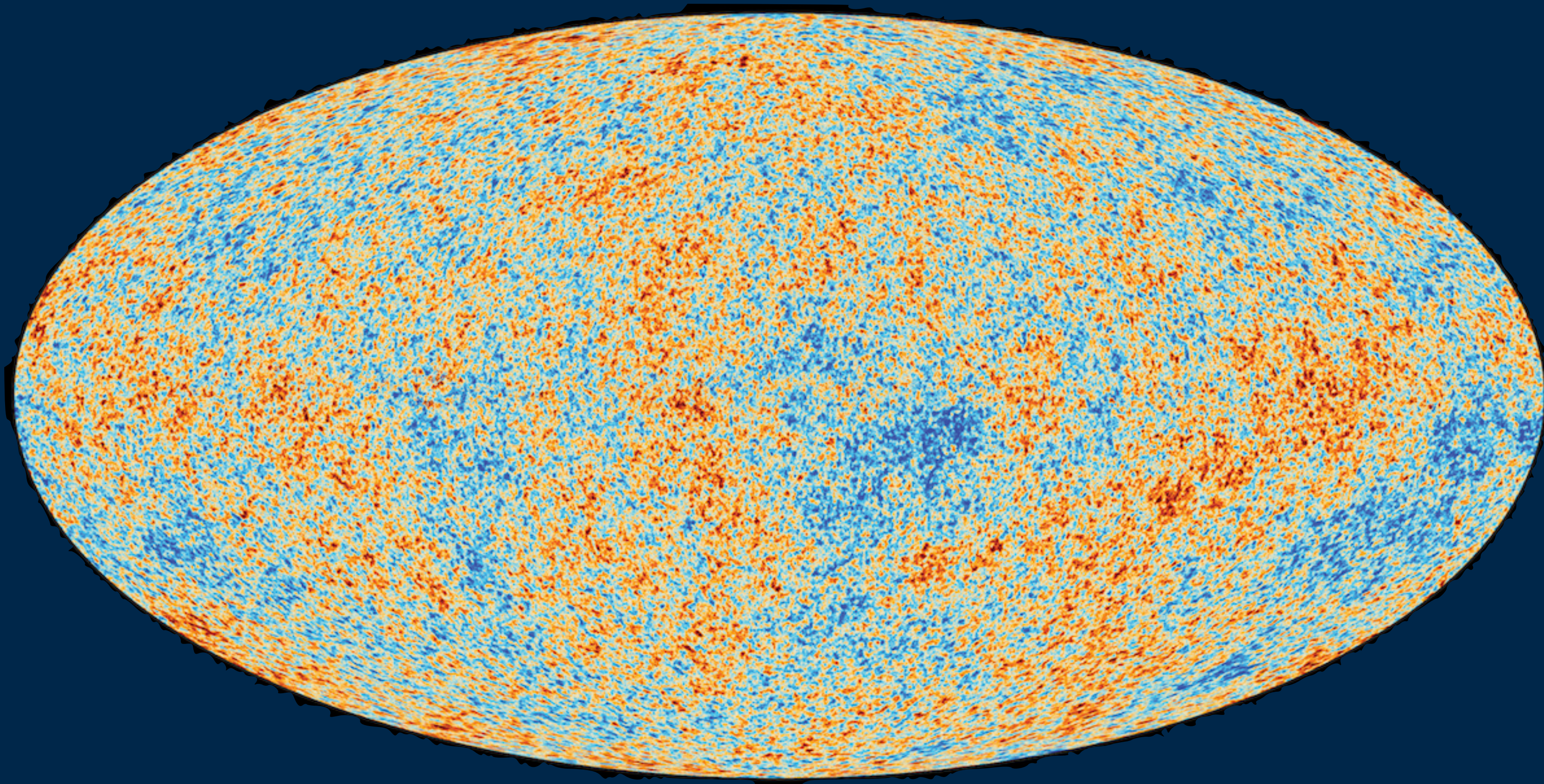


- 1000 times smaller than the monopole
- Doppler effect due to our motion wrt to rest frame of CMB
- Corresponding to a velocity of  $v = (627 \pm 22) \text{ km/s}$
- In the direction of galactic longitude  $l = (276 \pm 3)^\circ$   $b = (30 \pm 3)^\circ$

# CMB anisotropies (monopole + dipole subtracted)

Temperature slightly different on different patches on the sky.

1 part in 100 000



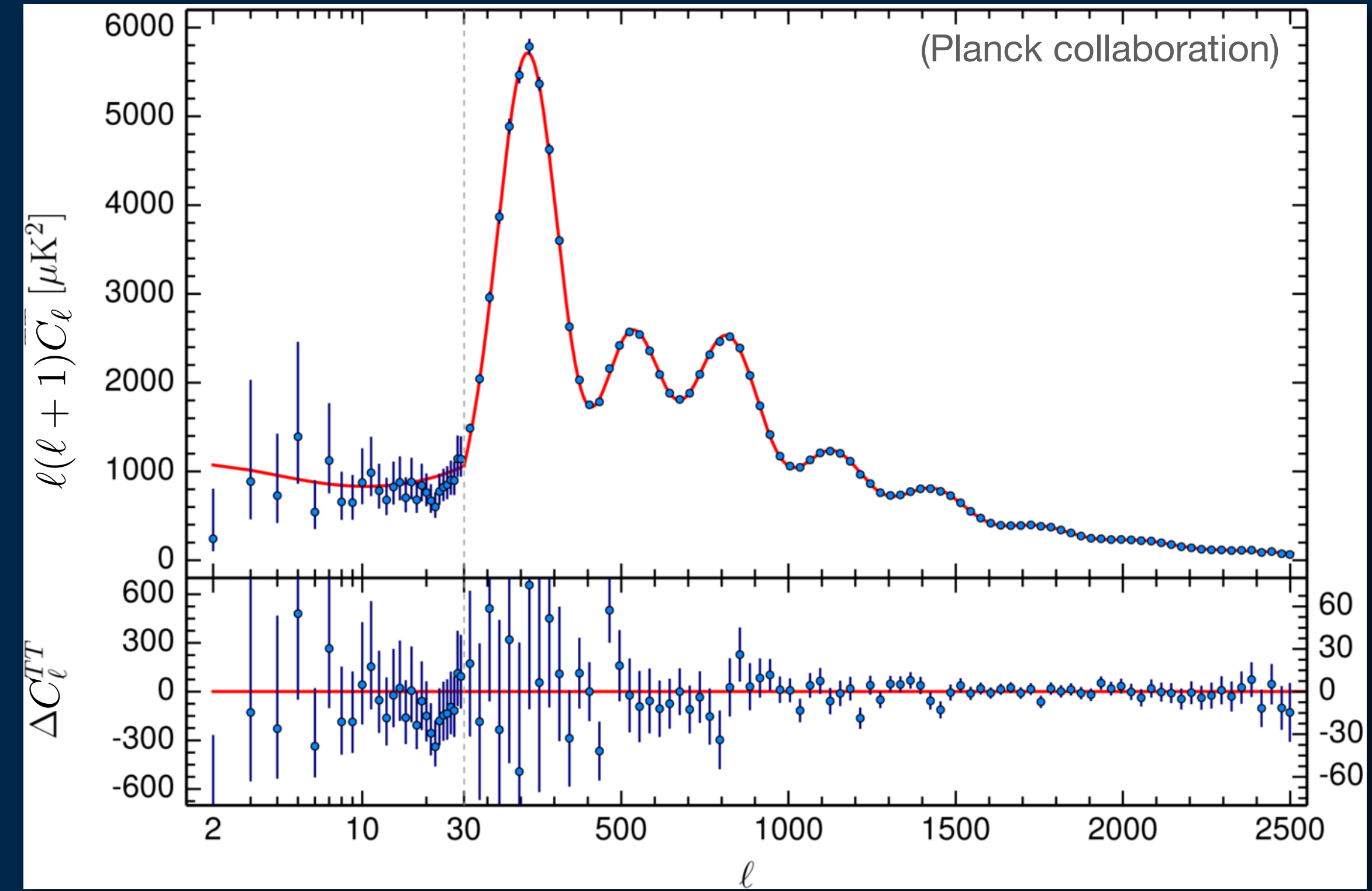
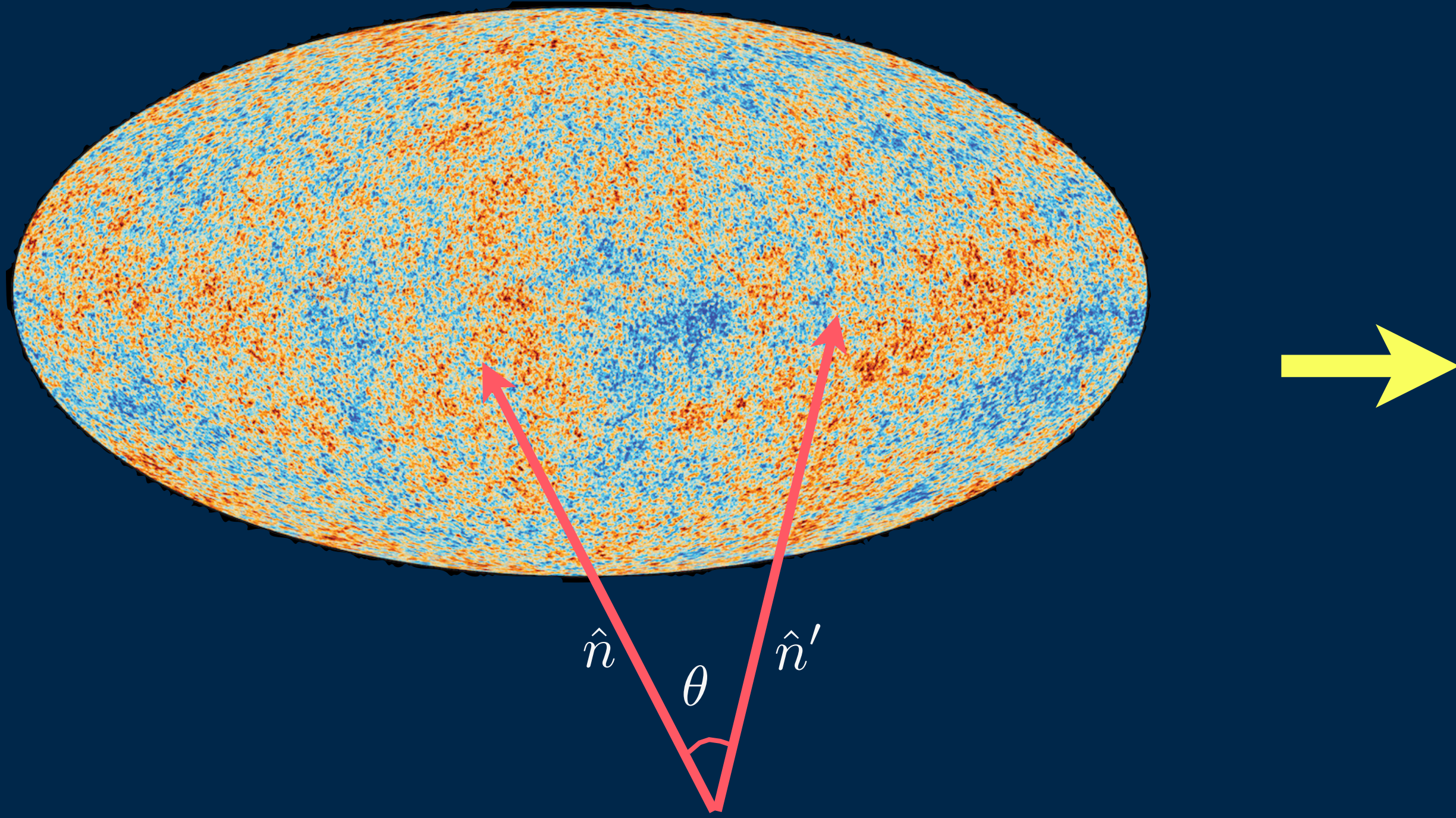
Planck Satellite

Temperature anisotropy  
in direction  $\hat{n}$

$$\Theta(\hat{n}) = \frac{T(\hat{n}) - \bar{T}_{CMB}}{\bar{T}_{CMB}} \lesssim 10^{-5}$$



# CMB angular power spectrum



$$\langle \Theta(\hat{n}) \Theta(\hat{n}') \rangle = \sum_{\ell} (2\ell + 1) C_{\ell} P_{\ell}(\cos \theta)$$

- $C_0$  Monopole
- $C_1$  Dipole
- $C_2$  Quadrupole
- ...

$C_{\ell}$  : Angular power spectrum

$P_{\ell}(\mu)$  : Legendre polynomials  
 $\mu = [-1, 1]$

Functions defined in  $[-1, 1]$   
 Can be expanded in terms of  $P_{\ell}(\mu)$

$$P_0 = 1 \quad P_1 = \mu \quad P_2 = \frac{3\mu^2 - 1}{2} \quad \dots$$

# Fluctuations in the Universe

CMB anisotropies: need to go beyond the homogeneous and isotropic Universe

Simplified assumption: *ALMOST* homogeneous and isotropic

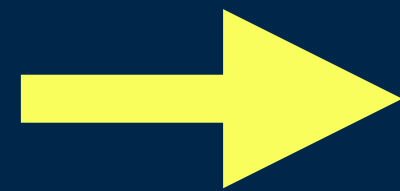
Split all variables into background (FLRW) + small fluctuation

e.g. metric  $ds^2 = -(1 + 2\Psi)dt^2 + a^2(1 - 2\Phi)\delta_{ij}dx^i dx^j$

Two gravitational potentials:  $\Psi(t, \vec{x})$   $\Phi(t, \vec{x})$

Photons:  $f(t, \vec{x}, \vec{p}) = \frac{1}{e^{p/T(t, \vec{x}, \vec{p})} - 1}$

Space-dependent temperature



$$\Delta(t, \vec{x}, \vec{p}) = \frac{T(t, \vec{x}, \vec{p}) - \bar{T}_\gamma(t)}{\bar{T}_\gamma(t)} = \sum_\ell (2\ell + 1) \Delta_\ell(t, k) P_\ell(\hat{x} \cdot \hat{p})$$

$$C_\ell = \frac{2}{\pi} \int dk k^2 P_0(k) |\Delta_\ell(t_0, k)|^2$$

Initial power spectrum (inflation)

Do the same for baryons, CDM and neutrinos (density fluctuation, velocity, ...)

Fourier expansion of every fluctuation variable

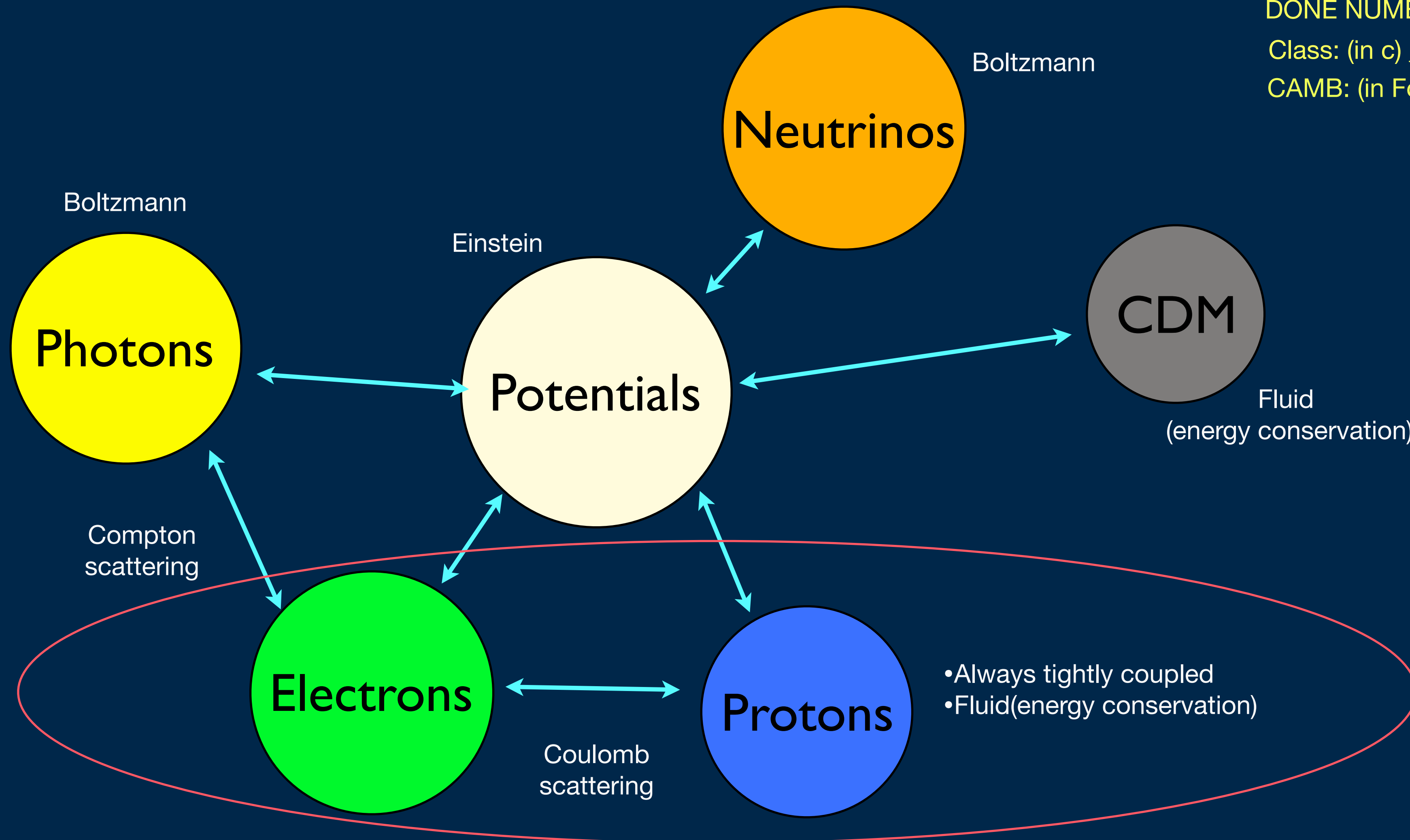
# Einstein-Fluid-Boltzmann system

The evolution of the metric potentials, baryon and dark matter fluid variables, and photon and neutrino temperature fluctuations

DONE NUMERICALLY with codes:

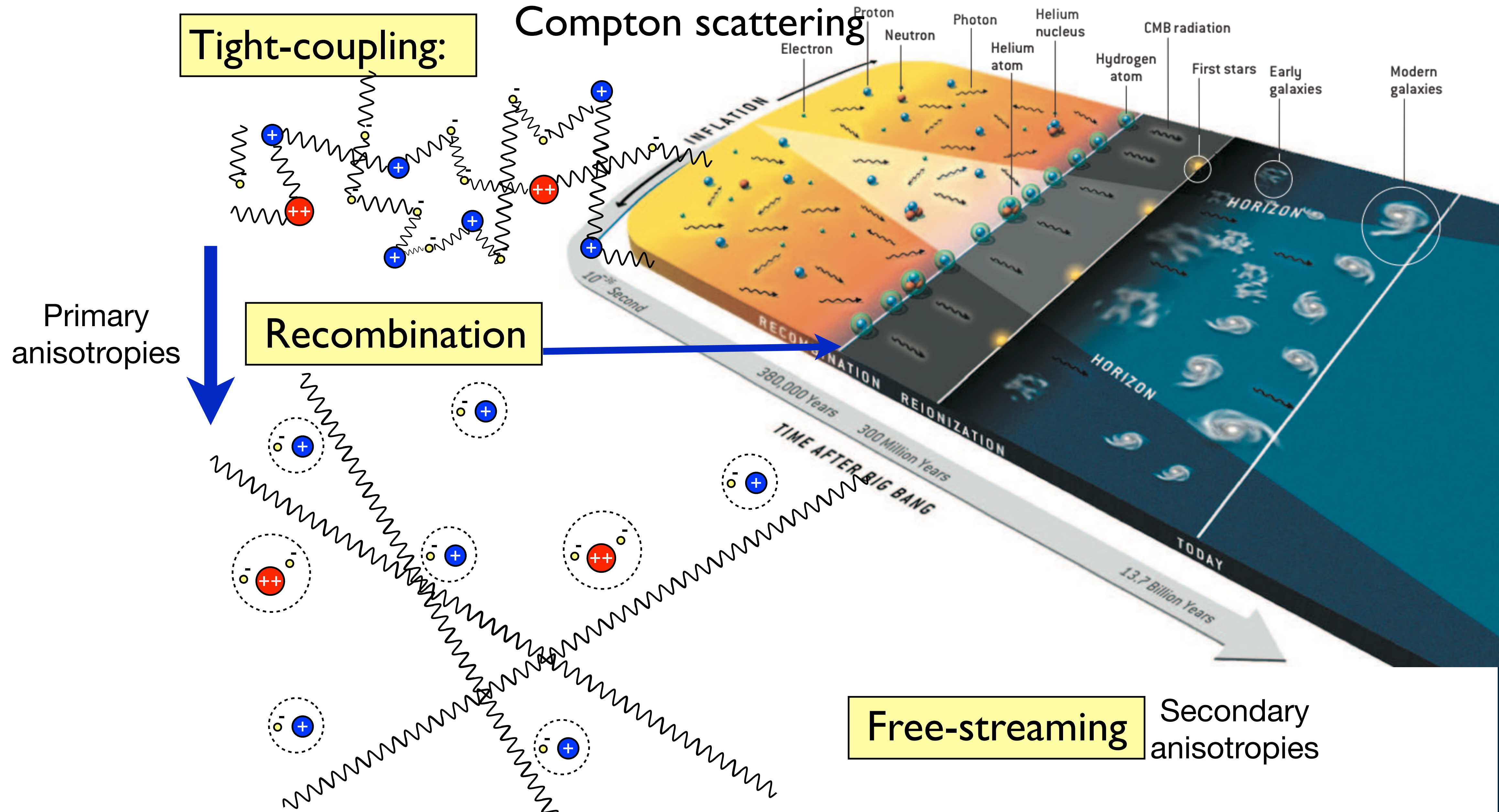
Class: (in c) <https://lesgourg.github.io>

CAMB: (in Fortran & Python) <https://camb.readthedocs.io/>

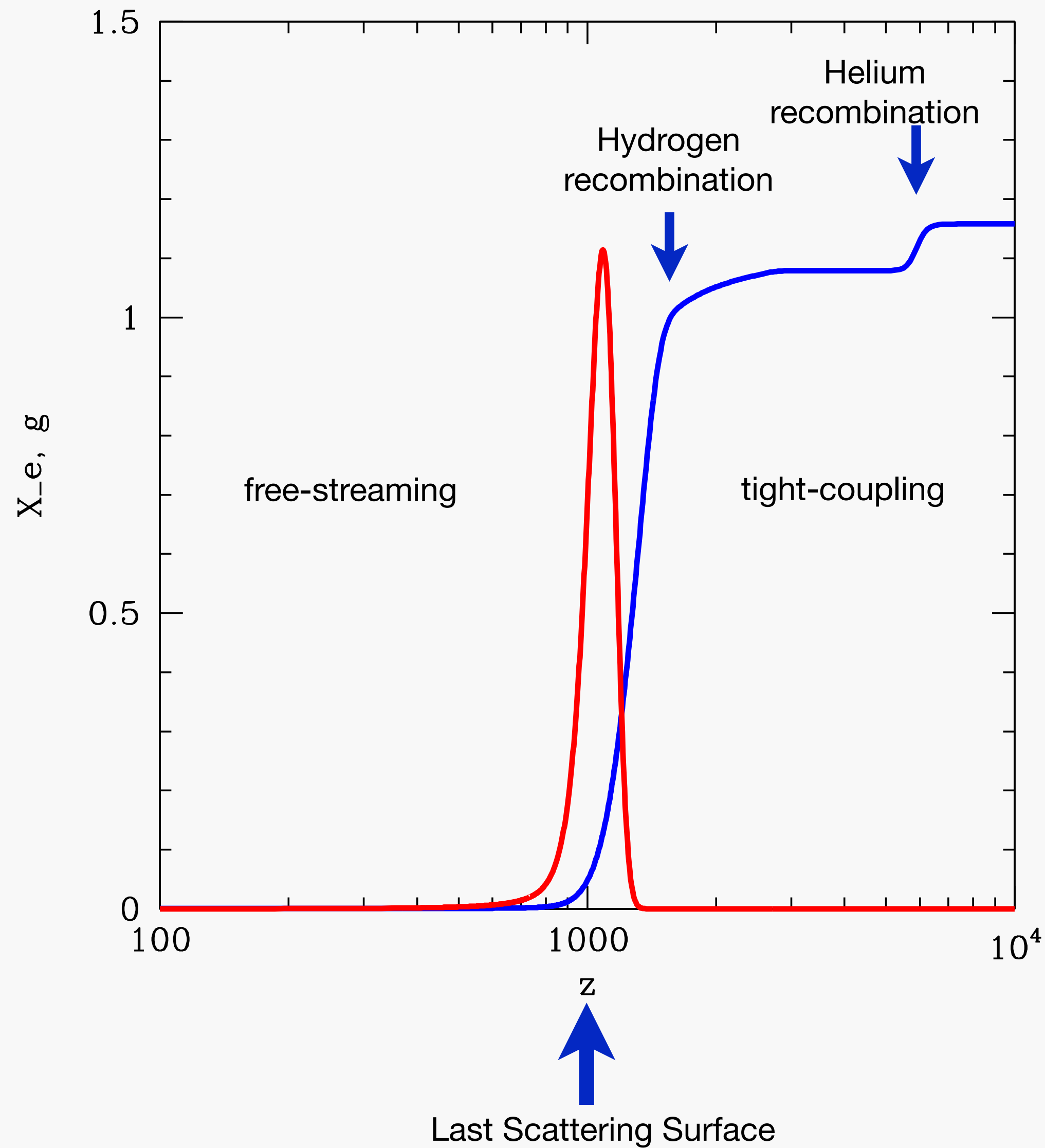


(from S.Dodelson)

# Tight-coupling and Free-streaming



# Visibility function

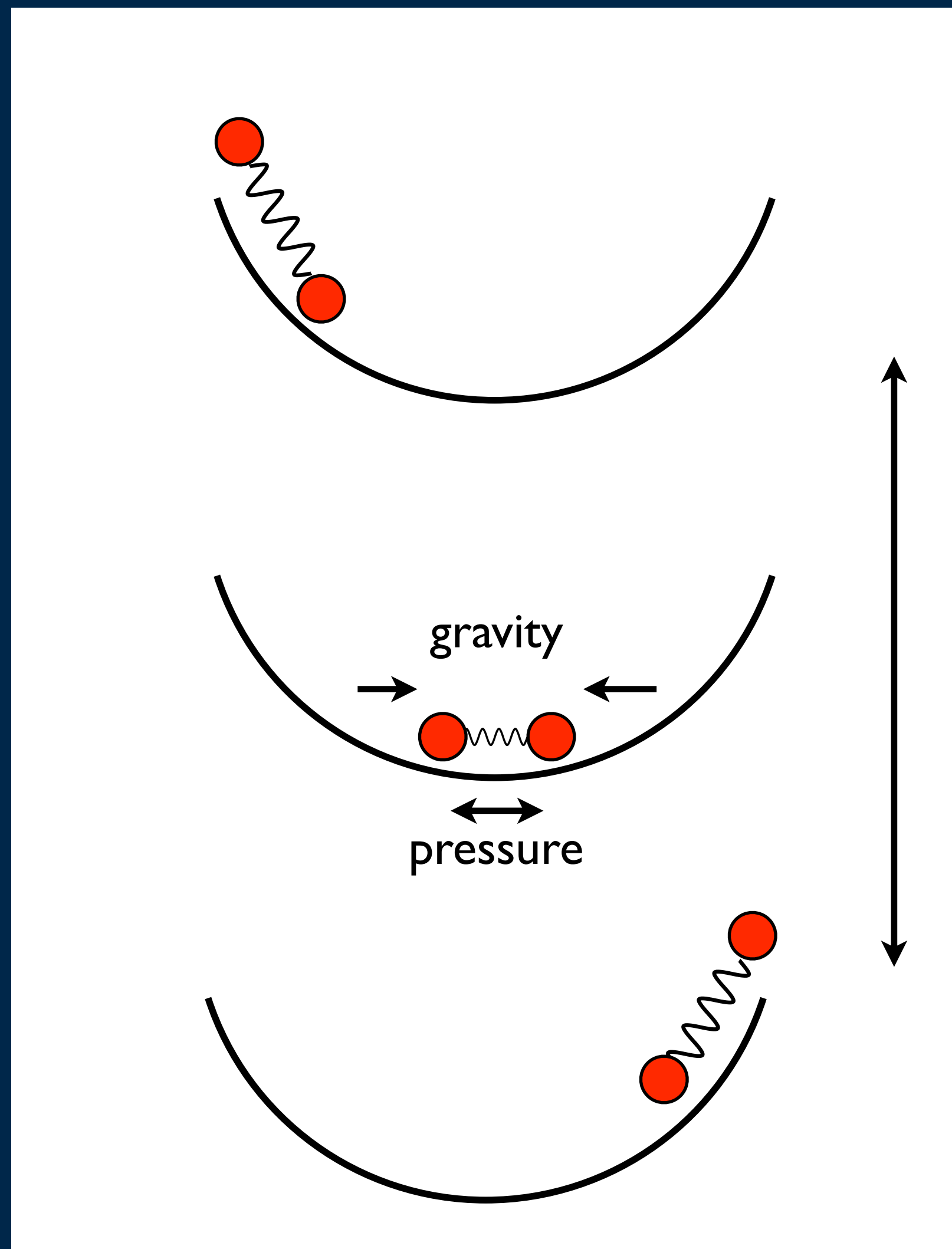


$X_e(z)$  Electron ionization fraction  
number of free electrons

$g(z)$  Visibility function

Proportional to the probability that an observed photon today has last scattered at redshift  $z$

# Acoustic oscillations



Tight-coupling very efficient — destroys all  $\Delta_\ell(t, k)$   $\ell \geq 2$

Full system of equations reduces to

$$\ddot{\Delta}_0 + \frac{\dot{R}}{1+R} \dot{\Delta}_0 + \frac{k^2}{3(1+R)} \Delta_0 = -\frac{k^2}{3} \Psi - \frac{1}{1+R} \frac{d}{dt} \left[ (1+R) \dot{\Phi} \right]$$

Baryon-photon ratio  $R = \frac{3\bar{\rho}_b}{4\bar{\rho}_\gamma}$

Gravity

Gravity

$$C_{\gamma b}^2 = \frac{1}{3(1+R)}$$

Sound horizon:  $r_s = \int_0^t C_{\gamma b} dt$

Assume slowly varying potentials and sound horizon

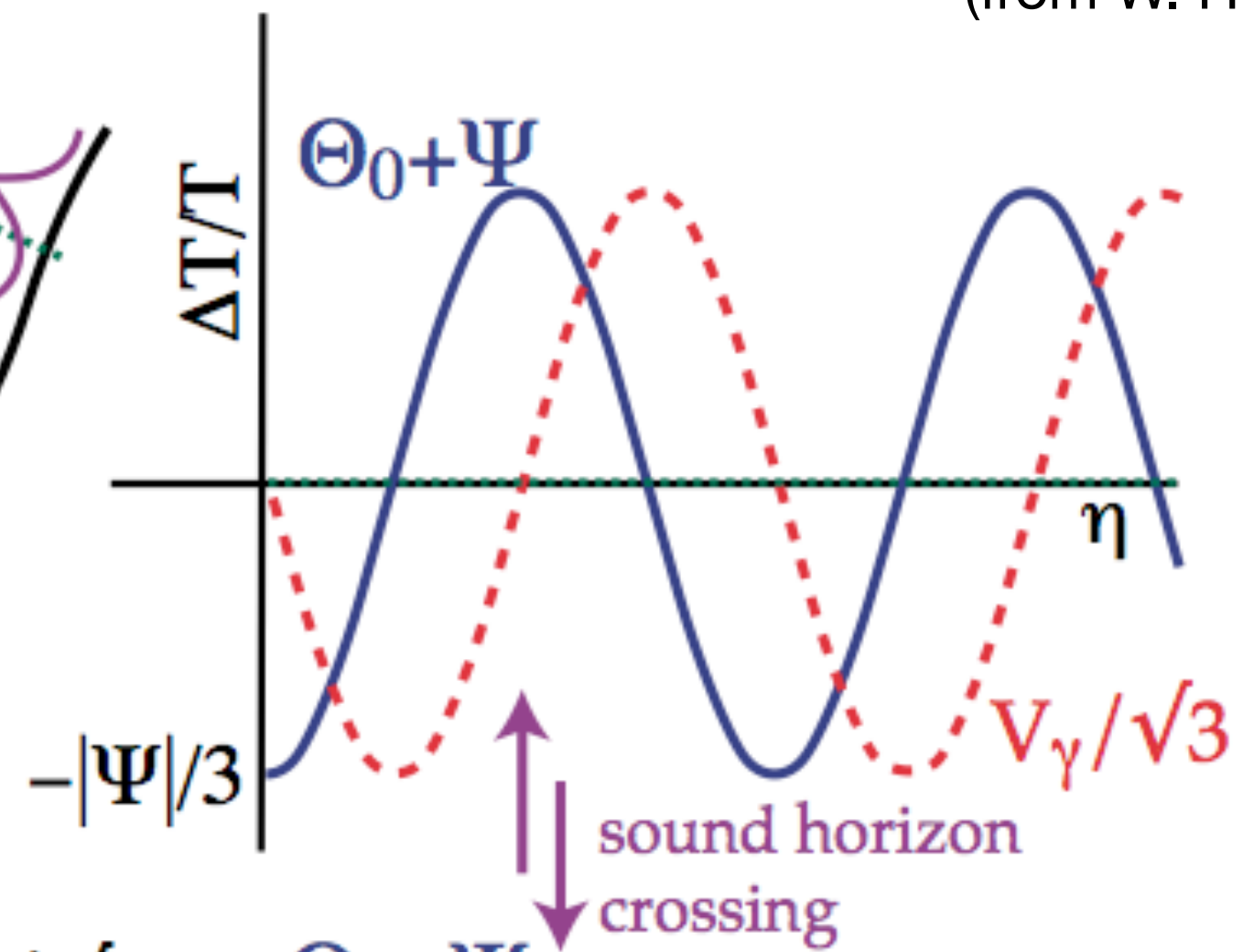
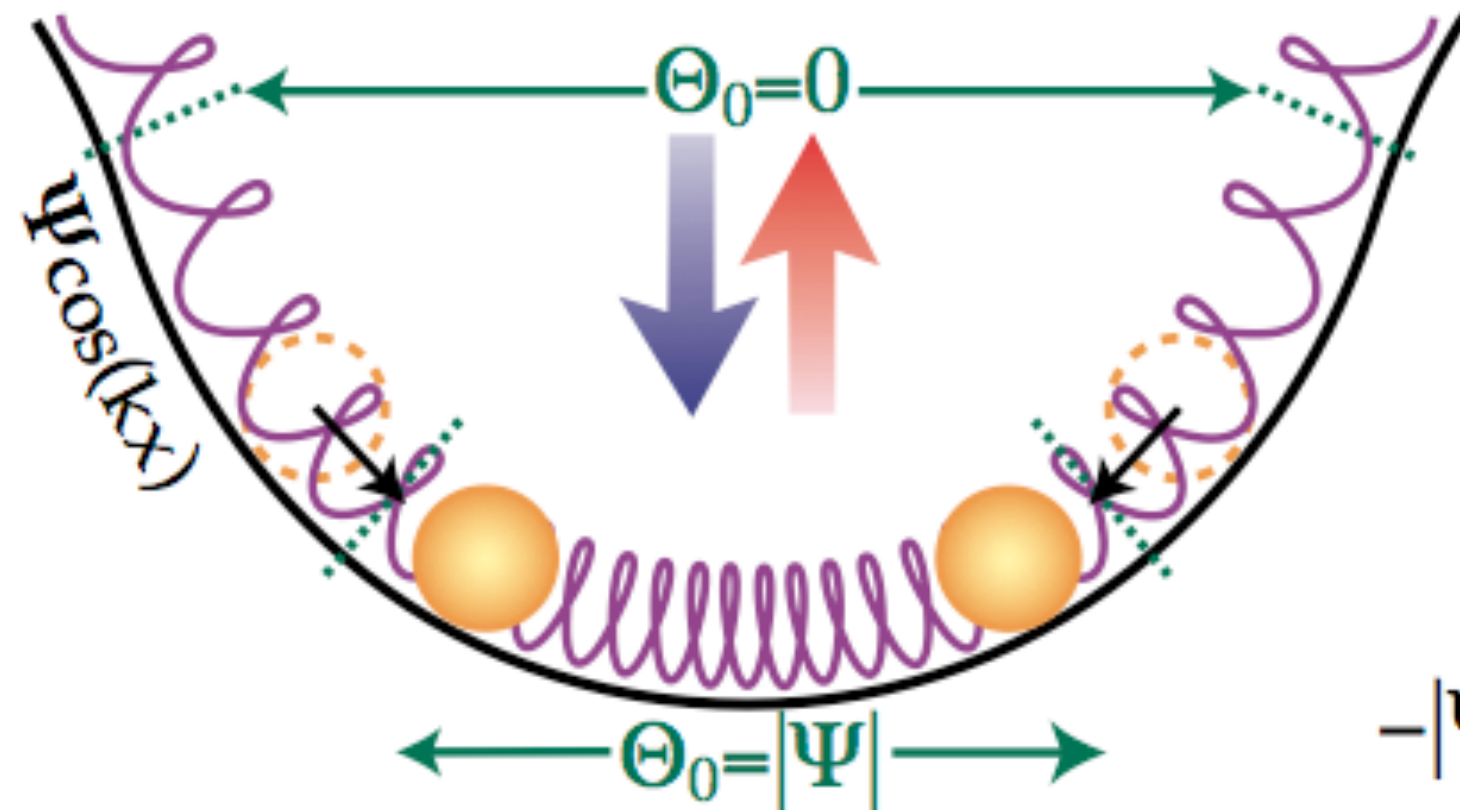
$$\ddot{\Delta}_0 + k^2 C_s^2 \Delta_0 \approx -\frac{k^2}{3} \Psi$$

$$(\Delta_0 + \Psi)(t_*) = A \cos(kr_s^* t_*) - R\Psi(t_*)$$

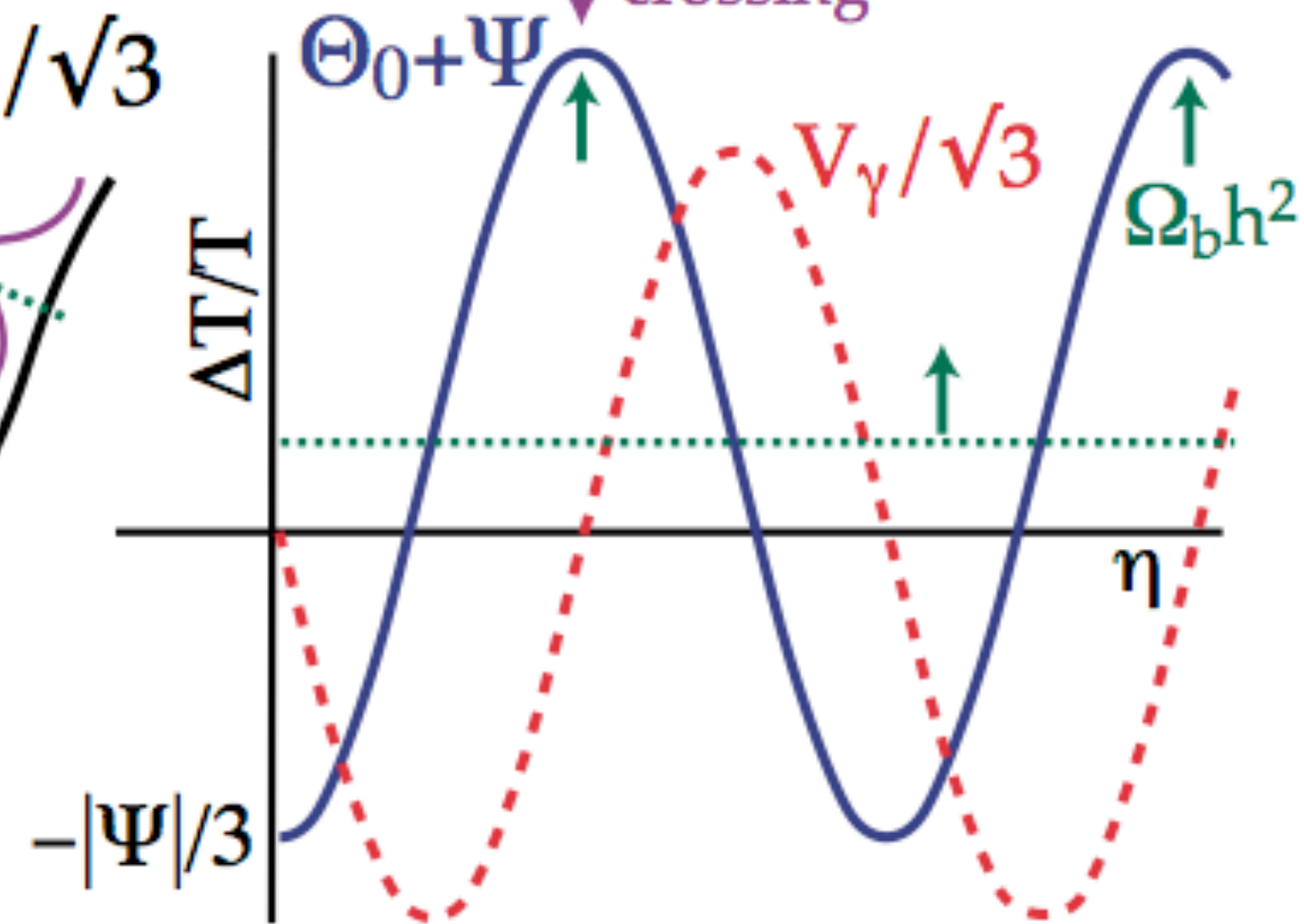
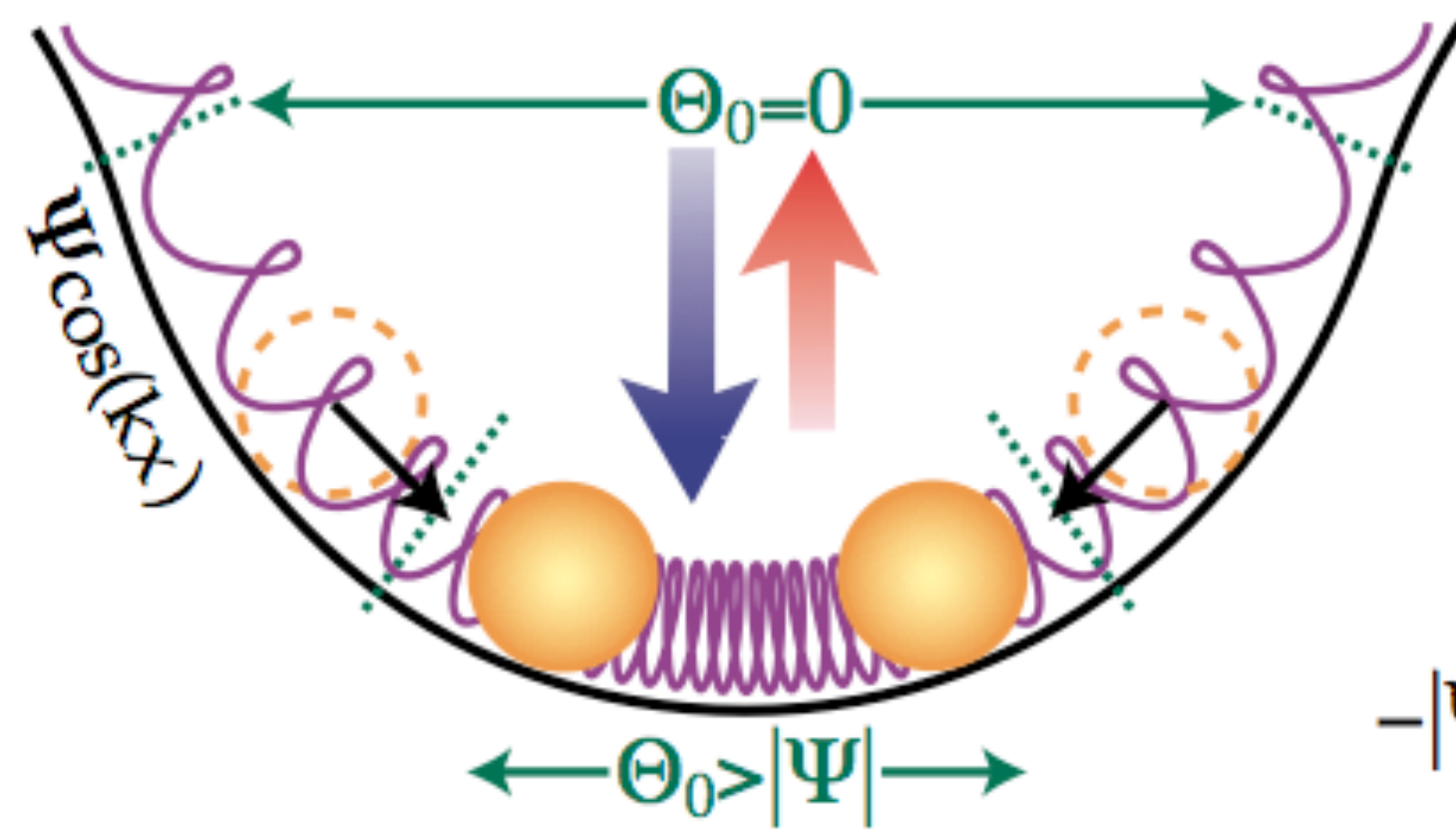
Oscillations displaced by gravity due to baryons  
Frequency set by sound horizon (baryons)

(from W. Hu)

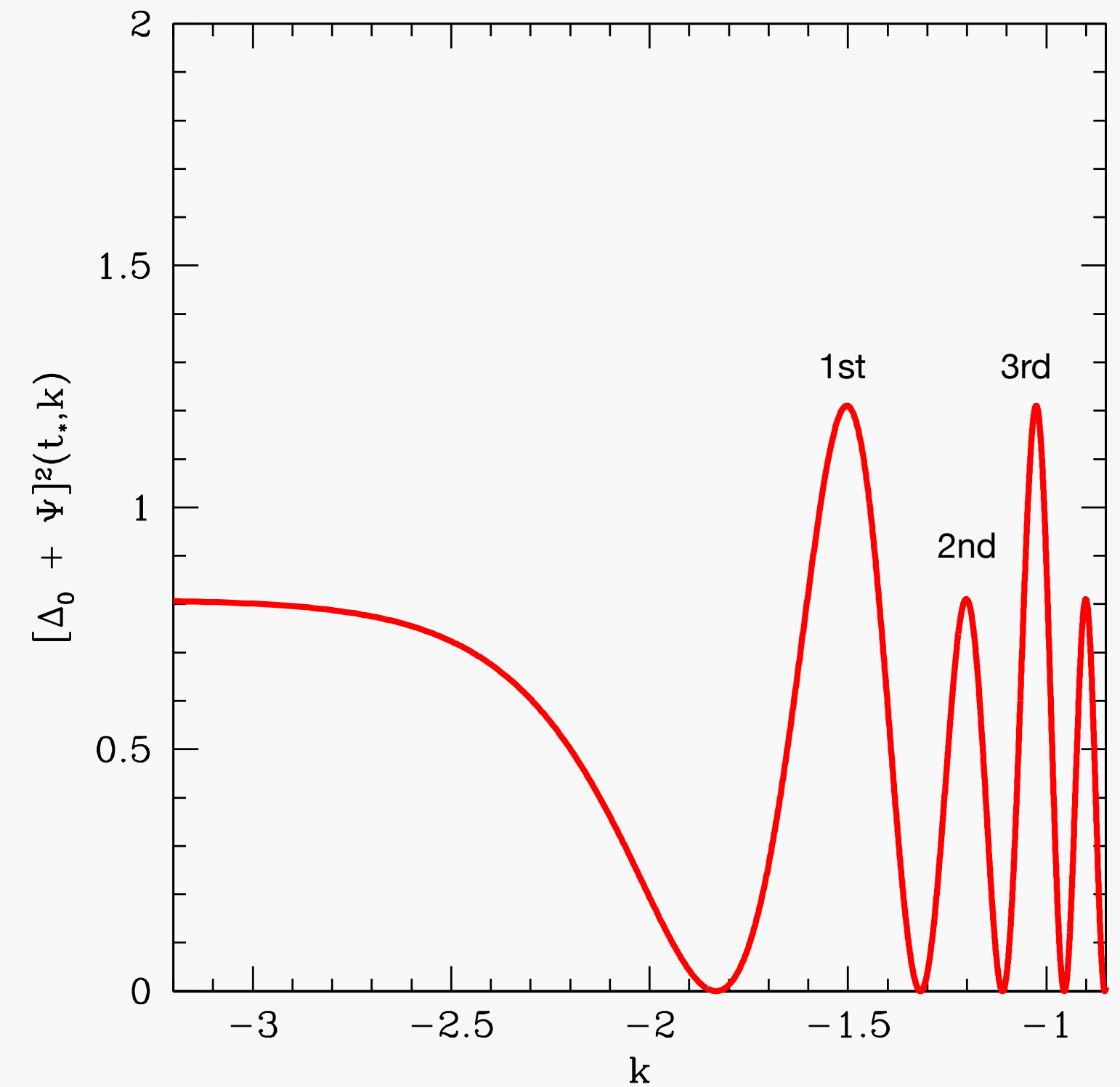
(a) Photons:  $c_s = 1/\sqrt{3}$



(b) Photons + Baryons:  $c_s < 1/\sqrt{3}$



At Last Scattering surface

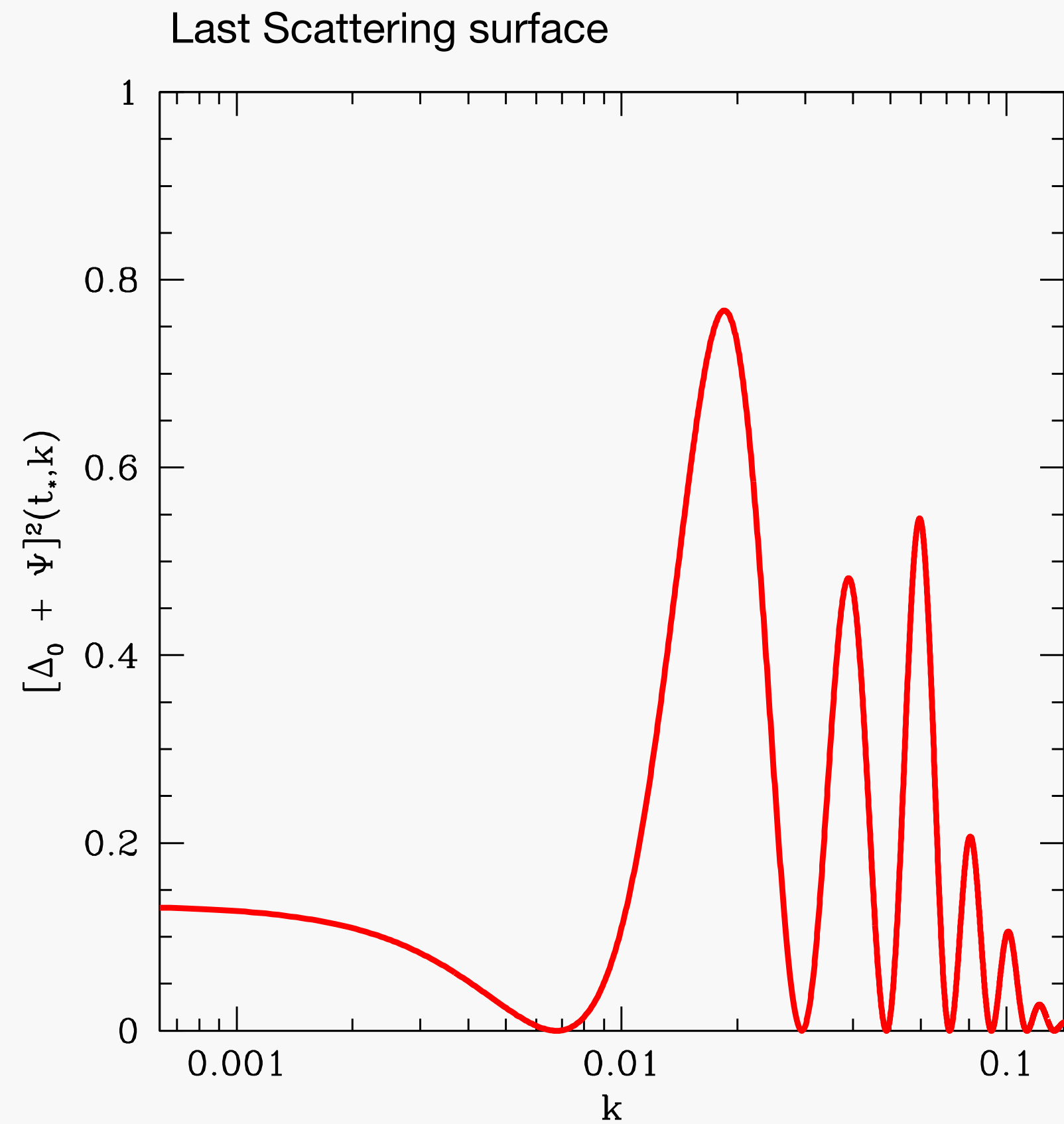


# Projection

After recombination  $\longrightarrow$  Free-streaming: photons follow null geodesics

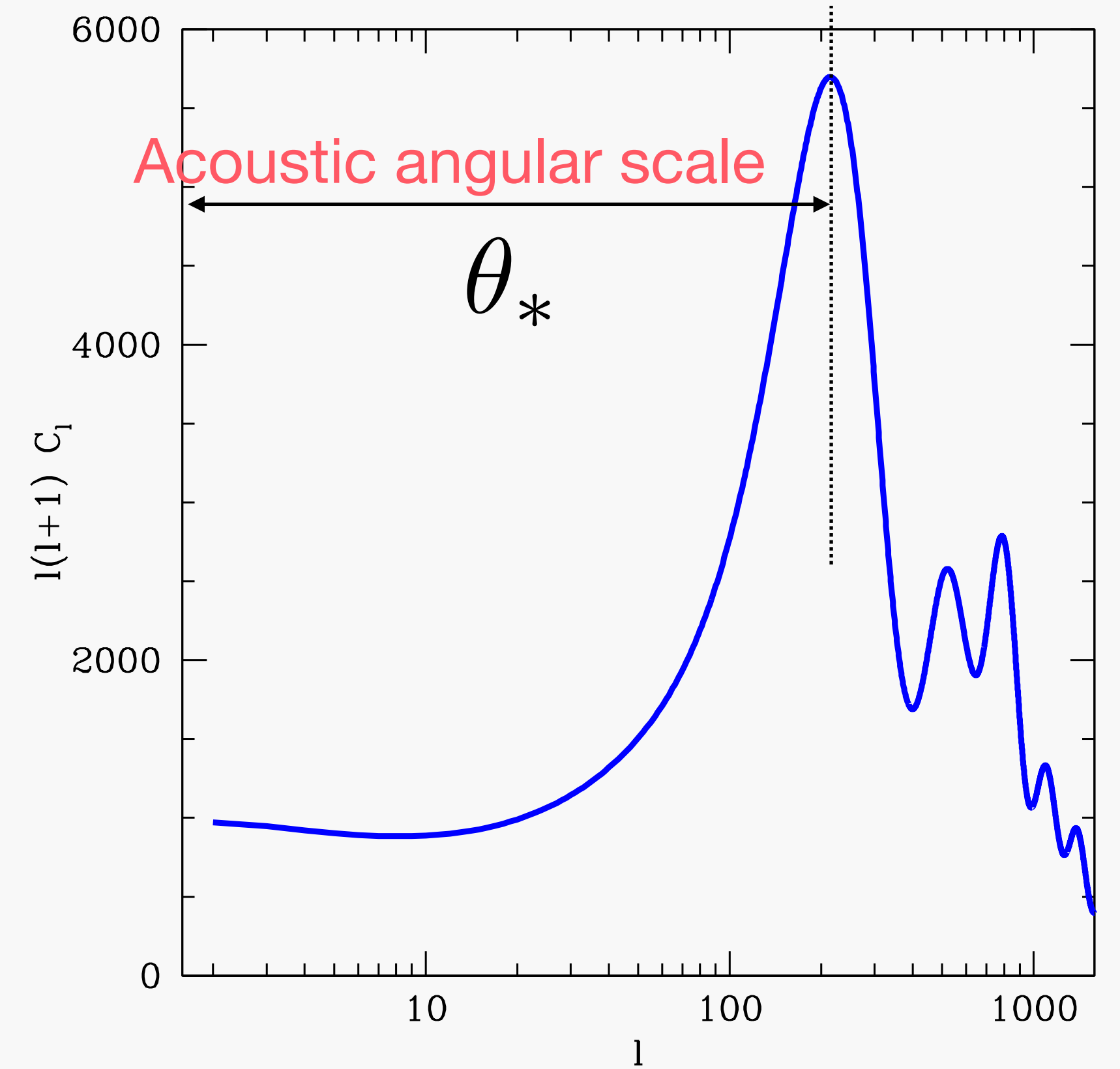
$$\Delta_\ell(t_0, k) \approx (\Delta_0 + \Psi)(t_*) j_\ell [k(t_0 - t_*)]$$

Strong dependence on: curvature,  $\Lambda$ , baryons



nth Peak location

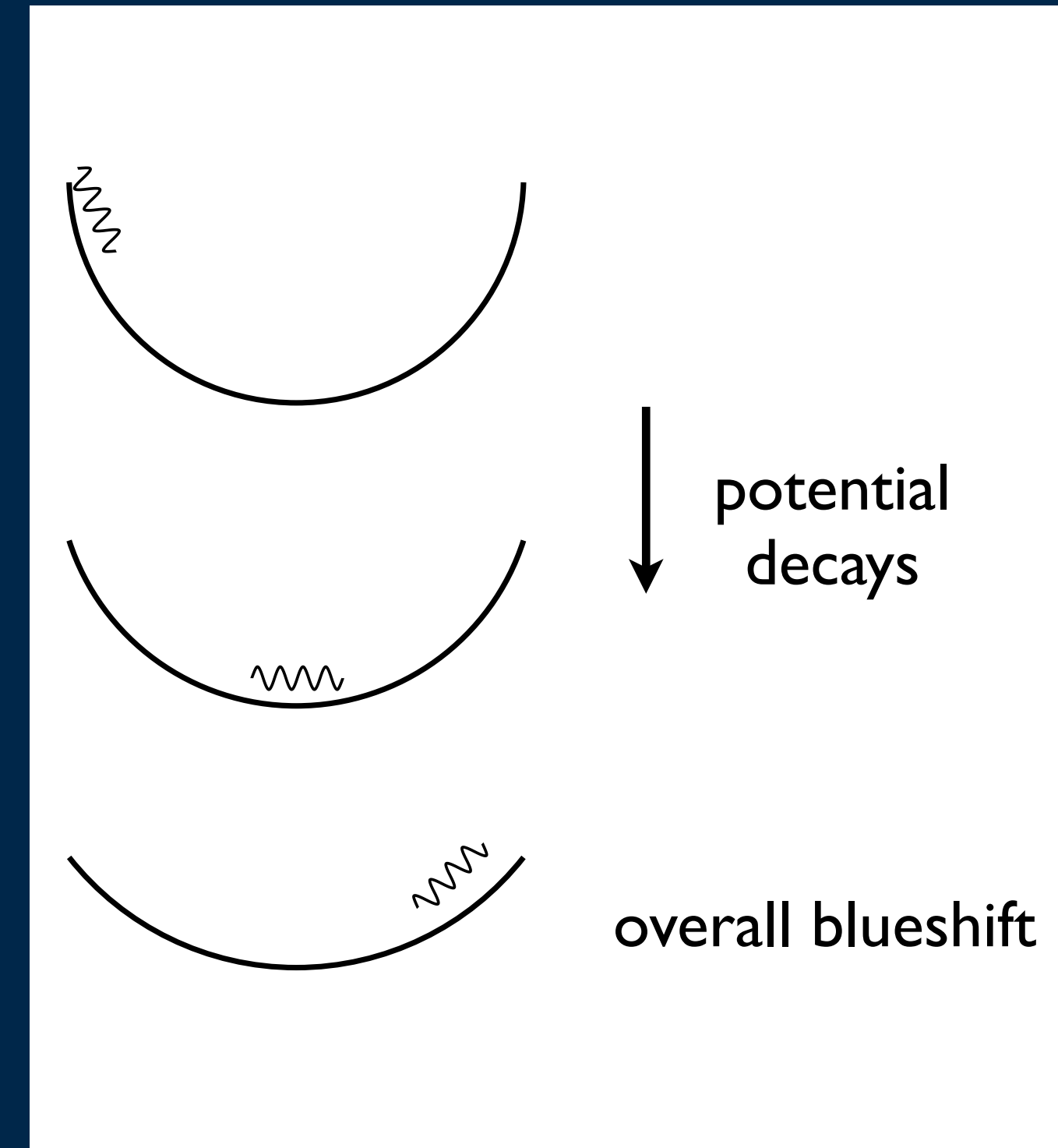
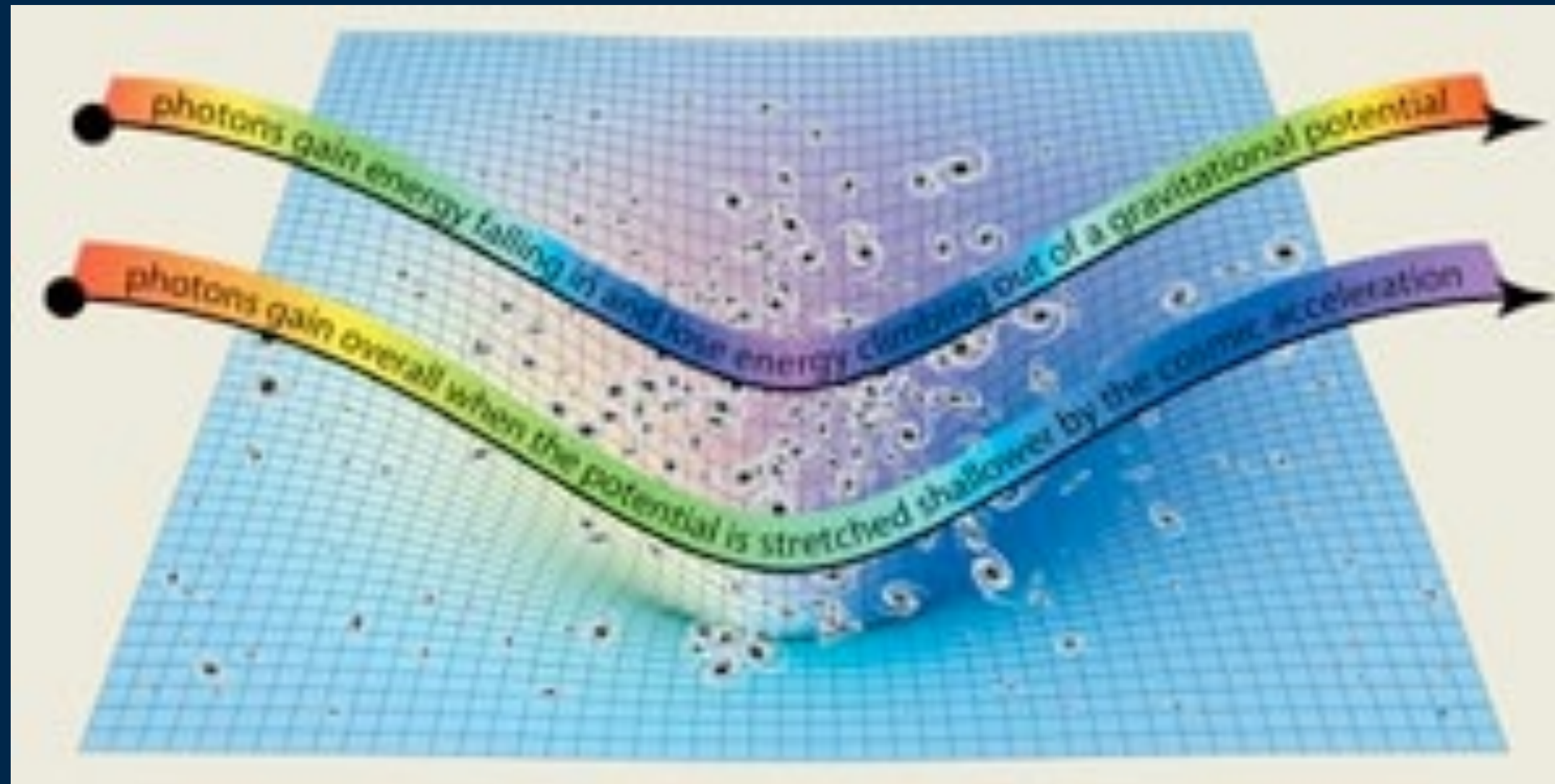
$$\ell_n = \frac{n\pi(t_0 - t_*)}{r_s}$$





# Integrated Sachs-Wolfe effect

Gravitational redshift of photons due to a potential well



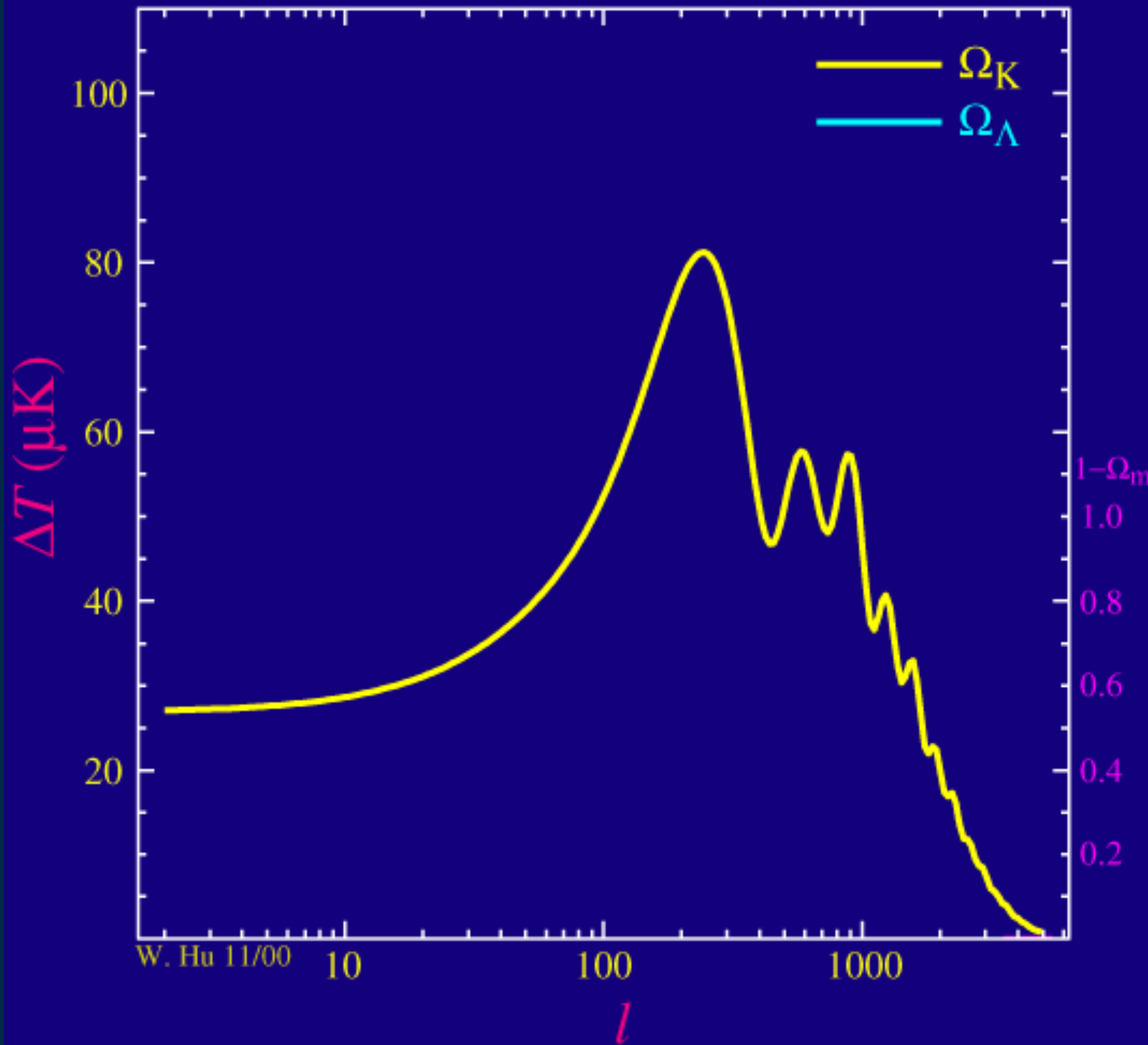
$$\dot{\Phi} \neq 0$$

potentials constant during **Matter domination** → **No ISW**

Radiation-to-matter transition: **early ISW** effect → **small** scales

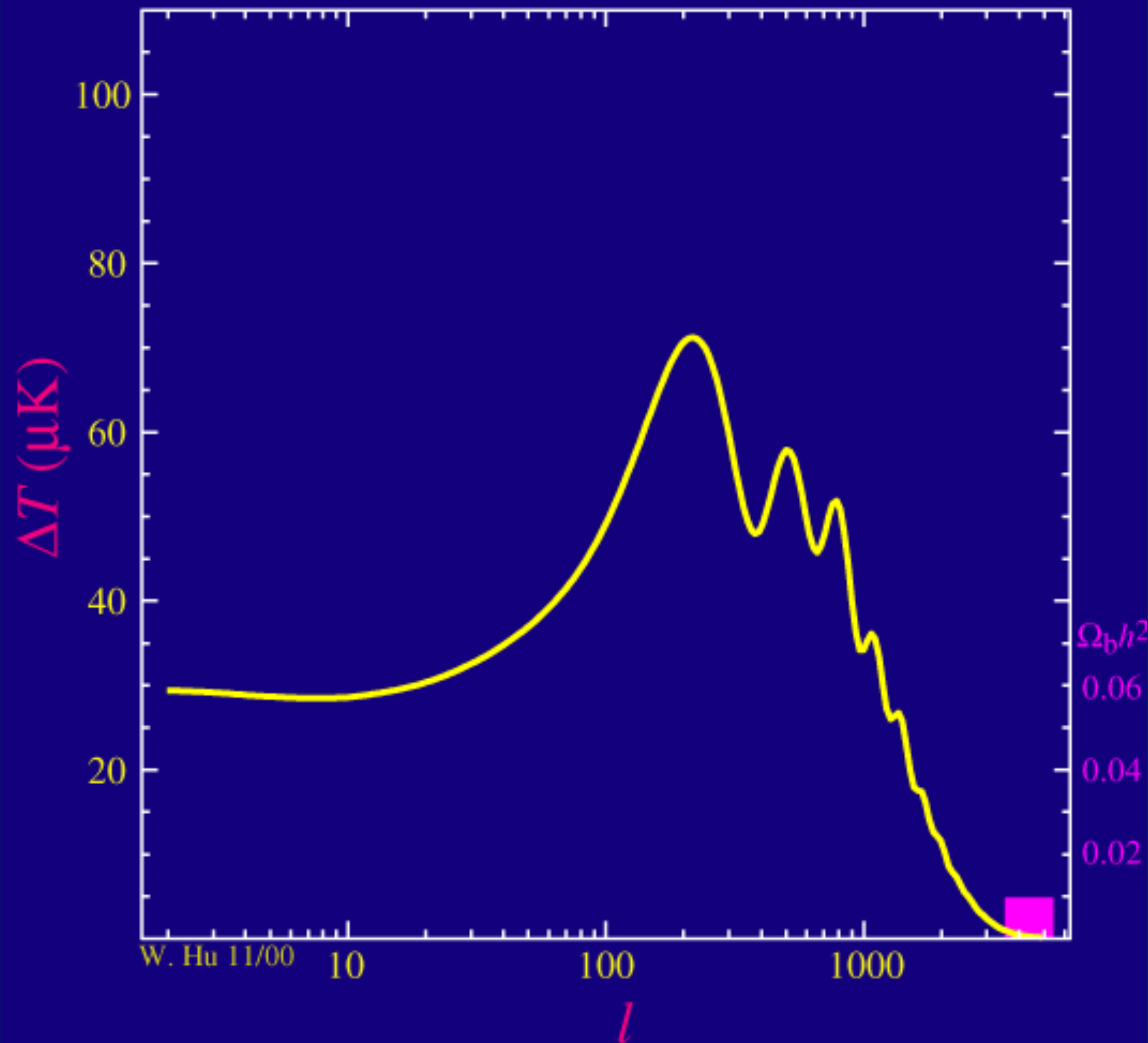
(Extra radiation or less matter)

$\Lambda$  Era: **late ISW** effect → **large** scales



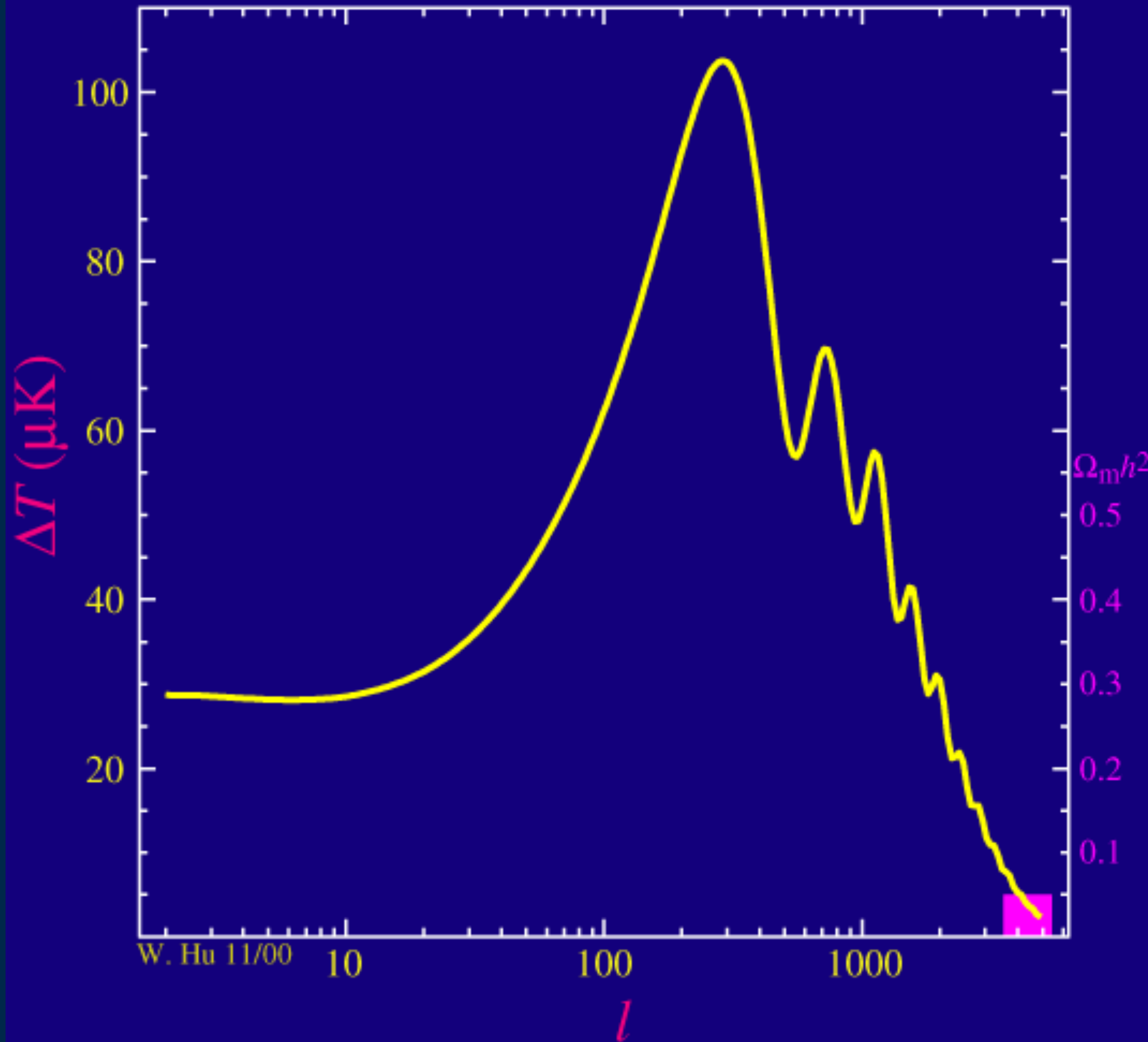
ANIMATION:

From Wayne Hu's webpage  
<http://background.uchicago.edu/~whu>



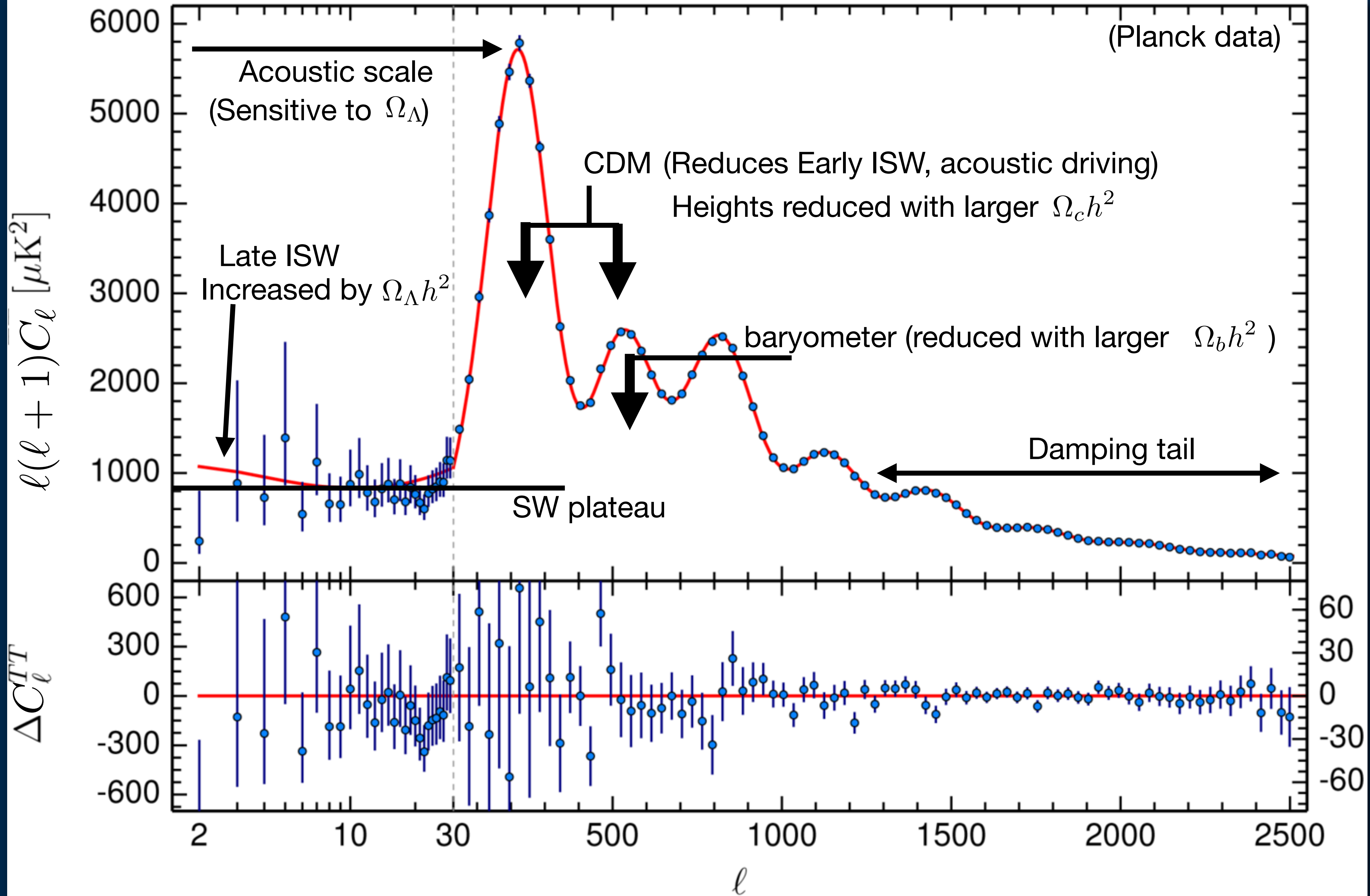
ANIMATION:

From Wayne Hu's webpage  
<http://background.uchicago.edu/~whu>

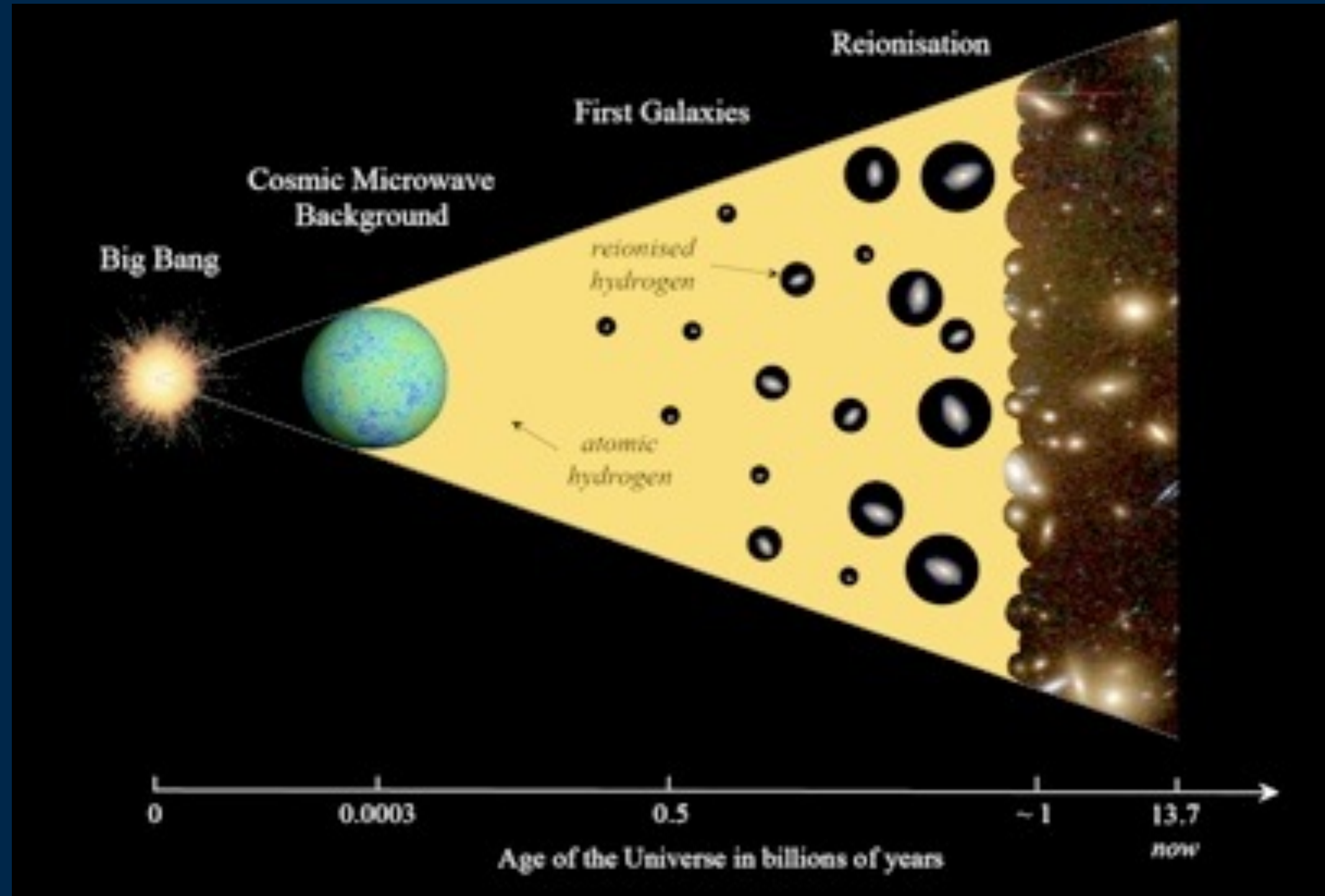


ANIMATION:

From Wayne Hu's webpage  
<http://background.uchicago.edu/~whu>



# Reionization

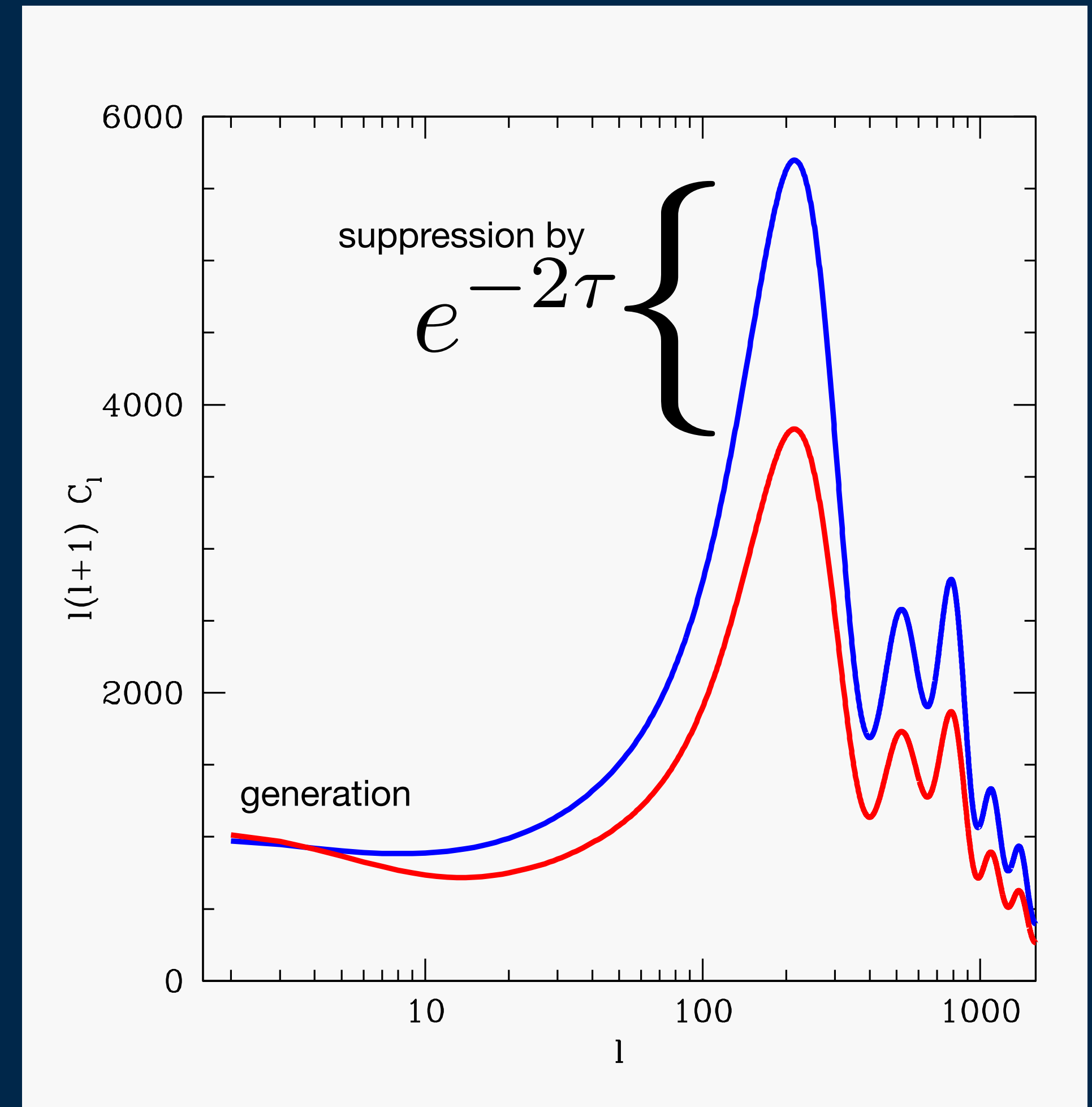


It's tight-coupling again!

optical depth to reionization  $\tau = \sigma_T \int_{t_0}^{t_r} \frac{X_e(t)}{a^2} dt$

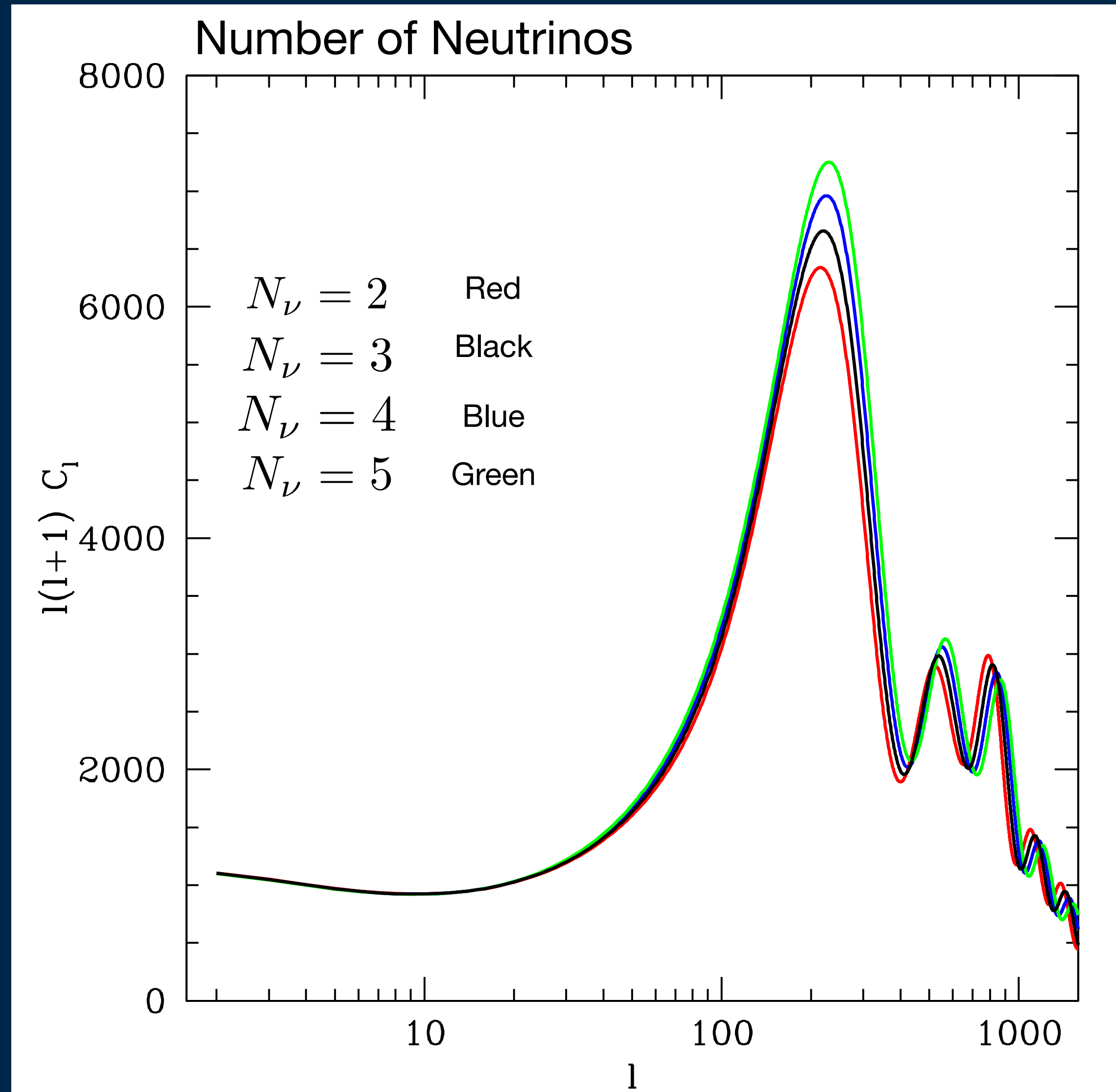
Large scales: generation of anisotropies

Small scales: suppression of anisotropies  
(remember diffusion damping)



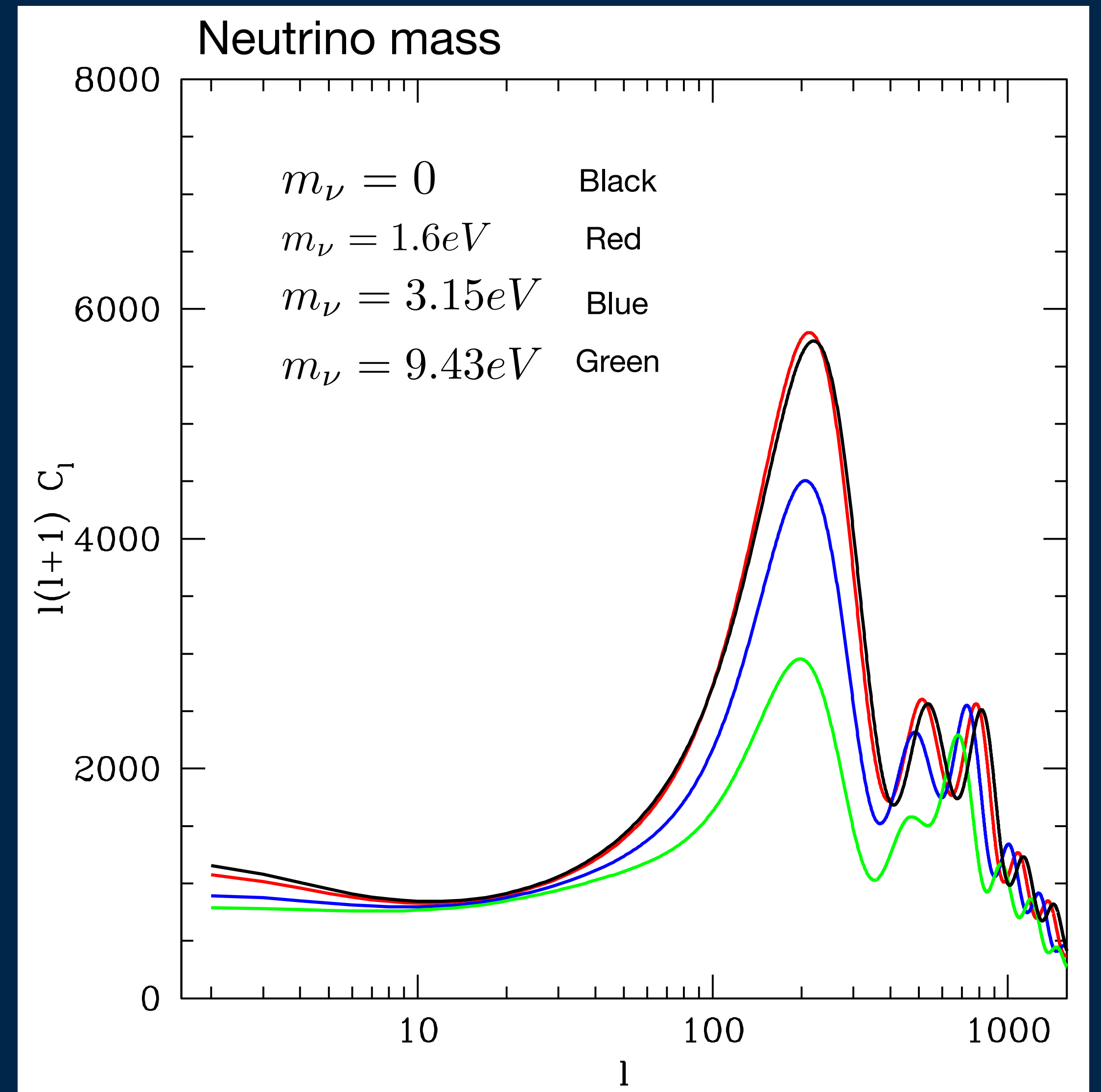
# Neutrinos and CMB

For CMB primary anisotropies neutrinos are still relativistic



More neutrinos shift matter-radiation equality later

- Shifts peaks to smaller scales
- Increases power around 1st peak



Largely insensitive to small masses

Large masses: neutrinos contribute like CDM

# Planck Collaboration (2018)



Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits	TT,TE,EE+lowE+lensing+ 68% limits
$\Omega_b h^2$	$0.02212 \pm 0.00022$	$0.02249 \pm 0.00025$	$0.0240 \pm 0.0012$	$0.02236 \pm 0.00015$	$0.02237 \pm 0.00015$	$0.02242 \pm 0.00014$
$\Omega_c h^2$	$0.1206 \pm 0.0021$	$0.1177 \pm 0.0020$	$0.1158 \pm 0.0046$	$0.1202 \pm 0.0014$	$0.1200 \pm 0.0012$	$0.11933 \pm 0.00091$
$100\theta_{MC}$	$1.04077 \pm 0.00047$	$1.04139 \pm 0.00049$	$1.03999 \pm 0.00089$	$1.04090 \pm 0.00031$	$1.04092 \pm 0.00031$	$1.04101 \pm 0.00029$
$\tau$	$0.0522 \pm 0.0080$	$0.0496 \pm 0.0085$	$0.0527 \pm 0.0090$	$0.0544^{+0.0070}_{-0.0081}$	$0.0544 \pm 0.0073$	$0.0561 \pm 0.0071$
$\ln(10^{10} A_s)$	$3.040 \pm 0.016$	$3.018^{+0.020}_{-0.018}$	$3.052 \pm 0.022$	$3.045 \pm 0.016$	$3.044 \pm 0.014$	$3.047 \pm 0.014$
$n_s$	$0.9626 \pm 0.0057$	$0.967 \pm 0.011$	$0.980 \pm 0.015$	$0.9649 \pm 0.0044$	$0.9649 \pm 0.0042$	$0.9665 \pm 0.0038$
$H_0$ [km s <sup>-1</sup> Mpc <sup>-1</sup> ]	$66.88 \pm 0.92$	$68.44 \pm 0.91$	$69.9 \pm 2.7$	$67.27 \pm 0.60$	$67.36 \pm 0.54$	$67.66 \pm 0.42$
$\Omega_\Lambda$	$0.679 \pm 0.013$	$0.699 \pm 0.012$	$0.711^{+0.033}_{-0.026}$	$0.6834 \pm 0.0084$	$0.6847 \pm 0.0073$	$0.6889 \pm 0.0056$
$\Omega_m$	$0.321 \pm 0.013$	$0.301 \pm 0.012$	$0.289^{+0.026}_{-0.033}$	$0.3166 \pm 0.0084$	$0.3153 \pm 0.0073$	$0.3111 \pm 0.0056$
$\Omega_m h^2$	$0.1434 \pm 0.0020$	$0.1408 \pm 0.0019$	$0.1404^{+0.0034}_{-0.0039}$	$0.1432 \pm 0.0013$	$0.1430 \pm 0.0011$	$0.14240 \pm 0.00087$
$\Omega_m h^3$	$0.09589 \pm 0.00046$	$0.09635 \pm 0.00051$	$0.0981^{+0.0016}_{-0.0018}$	$0.09633 \pm 0.00029$	$0.09633 \pm 0.00030$	$0.09635 \pm 0.00030$
$\sigma_8$	$0.8118 \pm 0.0089$	$0.793 \pm 0.011$	$0.796 \pm 0.018$	$0.8120 \pm 0.0073$	$0.8111 \pm 0.0060$	$0.8102 \pm 0.0060$
$S_8 \equiv \sigma_8(\Omega_m/0.3)^{0.5}$	$0.840 \pm 0.024$	$0.794 \pm 0.024$	$0.781^{+0.052}_{-0.060}$	$0.834 \pm 0.016$	$0.832 \pm 0.013$	$0.825 \pm 0.011$
$\sigma_8 \Omega_m^{0.25}$	$0.611 \pm 0.012$	$0.587 \pm 0.012$	$0.583 \pm 0.027$	$0.6090 \pm 0.0081$	$0.6078 \pm 0.0064$	$0.6051 \pm 0.0058$
$z_{re}$	$7.50 \pm 0.82$	$7.11^{+0.91}_{-0.75}$	$7.10^{+0.87}_{-0.73}$	$7.68 \pm 0.79$	$7.67 \pm 0.73$	$7.82 \pm 0.71$
$10^9 A_s$	$2.092 \pm 0.034$	$2.045 \pm 0.041$	$2.116 \pm 0.047$	$2.101^{+0.031}_{-0.034}$	$2.100 \pm 0.030$	$2.105 \pm 0.030$
$10^9 A_s e^{-2\tau}$	$1.884 \pm 0.014$	$1.851 \pm 0.018$	$1.904 \pm 0.024$	$1.884 \pm 0.012$	$1.883 \pm 0.011$	$1.881 \pm 0.010$
Age [Gyr]	$13.830 \pm 0.037$	$13.761 \pm 0.038$	$13.64^{+0.16}_{-0.14}$	$13.800 \pm 0.024$	$13.797 \pm 0.023$	$13.787 \pm 0.020$
$z_*$	$1090.30 \pm 0.41$	$1089.57 \pm 0.42$	$1087.8^{+1.6}_{-1.7}$	$1089.95 \pm 0.27$	$1089.92 \pm 0.25$	$1089.80 \pm 0.21$
$r_*$ [Mpc]	$144.46 \pm 0.48$	$144.95 \pm 0.48$	$144.29 \pm 0.64$	$144.39 \pm 0.30$	$144.43 \pm 0.26$	$144.57 \pm 0.22$
$100\theta_*$	$1.04097 \pm 0.00046$	$1.04156 \pm 0.00049$	$1.04001 \pm 0.00086$	$1.04109 \pm 0.00030$	$1.04110 \pm 0.00031$	$1.04119 \pm 0.00029$
$z_{drag}$	$1059.39 \pm 0.46$	$1060.03 \pm 0.54$	$1063.2 \pm 2.4$	$1059.93 \pm 0.30$	$1059.94 \pm 0.30$	$1060.01 \pm 0.29$
$r_{drag}$ [Mpc]	$147.21 \pm 0.48$	$147.59 \pm 0.49$	$146.46 \pm 0.70$	$147.05 \pm 0.30$	$147.09 \pm 0.26$	$147.21 \pm 0.23$
$k_D$ [Mpc <sup>-1</sup> ]	$0.14054 \pm 0.00052$	$0.14043 \pm 0.00057$	$0.1426 \pm 0.0012$	$0.14090 \pm 0.00032$	$0.14087 \pm 0.00030$	$0.14078 \pm 0.00028$
$z_{eq}$	$3411 \pm 48$	$3349 \pm 46$	$3340^{+81}_{-92}$	$3407 \pm 31$	$3402 \pm 26$	$3387 \pm 21$
$k_{eq}$ [Mpc <sup>-1</sup> ]	$0.01041 \pm 0.00014$	$0.01022 \pm 0.00014$	$0.01019^{+0.00025}_{-0.00028}$	$0.010398 \pm 0.000094$	$0.010384 \pm 0.000081$	$0.010339 \pm 0.000063$
$100\theta_{s,eq}$	$0.4483 \pm 0.0046$	$0.4547 \pm 0.0045$	$0.4562 \pm 0.0092$	$0.4490 \pm 0.0030$	$0.4494 \pm 0.0026$	$0.4509 \pm 0.0020$



The End