



Hot QCD Matter

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Lecture 1: Tools

Lecture 2: Initial conditions: partonic structure and global observables

Lecture 3: Collective flow and hydrodynamics

Lecture 4: Jets and other hard probes

My approach to these lectures

The field of hot QCD matter spans the boundaries of nuclear, particle and condensed matter physics, and string theory

The field is also relatively young, with many phenomena not yet understood on a fundamental level

Think about it more like Condensed Matter Physics than Particle Physics: the Lagrangian is known precisely (QED, QCD), but many interesting phenomena cannot be calculated from first principles

→ Extensive use of effective theories and modeling

This is an opportunity for new ideas and concepts, but also a barrier to the outsider to sort out what is really known and what is conjectured

I will make no attempt to be comprehensive

Rather, I will discuss a limited number of topics that are well-established experimentally and have a connection to well-founded theory

Outline: Lecture 1

Theory Tools

- Basics of QCD
- Finite Temperature QCD

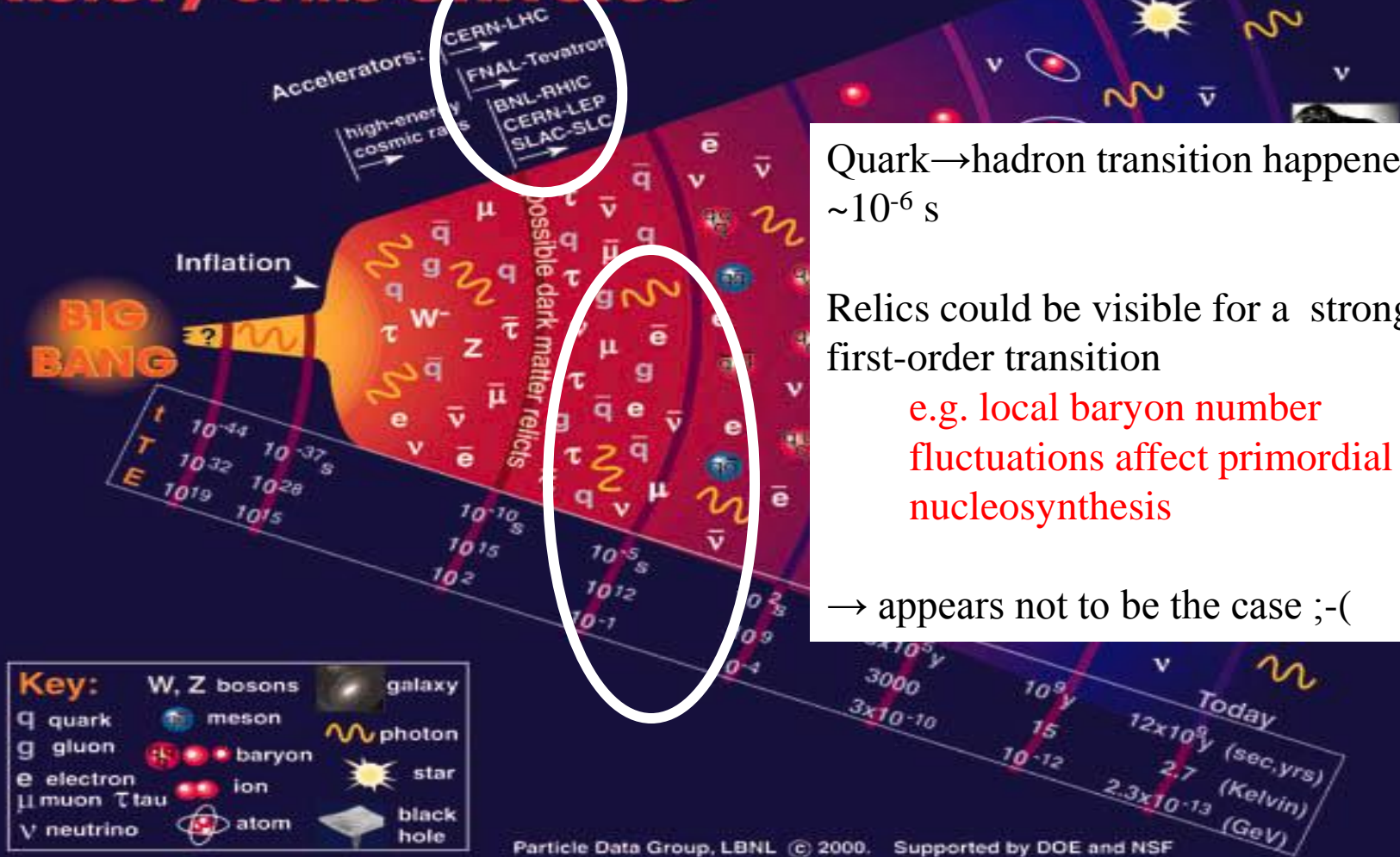
Experimental Tools

- Colliders
- Detectors

Analysis Tools

- Relativistic Kinematics
- Characterization of nuclear collisions

History of the Universe



Quark → hadron transition happened at $\sim 10^{-6}$ s

Relics could be visible for a strongly first-order transition

e.g. local baryon number fluctuations affect primordial nucleosynthesis

→ appears not to be the case ;-(

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Quantum Chromo-dynamics: the field theory of the strong (nuclear) force

Same basic structure as QED (electromagnetism)....

....except that gluons ("photons" of strong force) carry (color) charge...

Gluons

$$A_{\mu}^a \left\{ \begin{array}{l} \text{color } a = 1, \dots, 8 \\ \text{spin } \epsilon_{\mu}^{\pm} \end{array} \right.$$

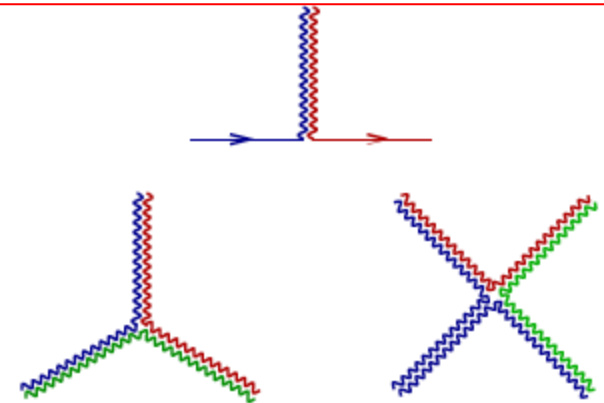
Dynamics: Generalized Maxwell (Yang-Mills)

....so they interact among themselves, generating much more complex structures...

$$\mathcal{L} = \bar{q}_f (i\not{D} - m_f) q_f - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$

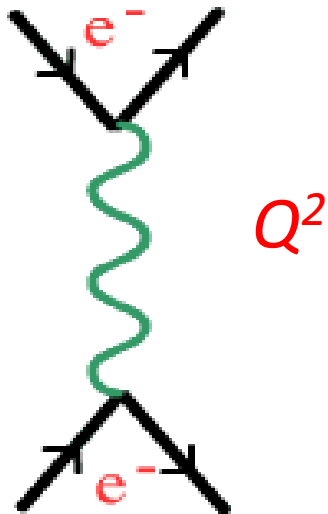
$$G_{\mu\nu}^a = \partial_{\mu} A_{\nu}^a - \partial_{\nu} A_{\mu}^a + gf^{abc} A_{\mu}^b A_{\nu}^c$$

$$i\not{D}q = \gamma^{\mu} (i\partial_{\mu} + gA_{\mu}^a \not{t}^a) q$$



Field theory: “running” of the coupling

Consider the interaction of two elementary particles:



Momentum transfer Q^2

small $Q^2 \Rightarrow$ large distance scales

large $Q^2 \Rightarrow$ small distance scales

Quantum mechanics:

Virtual pairs (loops) screen bare interaction

\Rightarrow momentum-dependent interaction strength

Running of the coupling: QED vs QCD

$$a \circ \frac{g^2}{4p}$$

negative

QED:

$$\frac{1}{\alpha(Q^2)} \approx \frac{1}{\alpha(\mu^2)} - \frac{1}{3\pi} \log\left(\frac{|Q^2|}{\mu^2}\right)$$

Smaller $|Q^2|$ (larger distance) \Rightarrow weaker coupling

- similar to screening of charge in di-electric material

QCD:

$$\frac{1}{\alpha_s(Q^2)} \approx \frac{1}{\alpha_s(\mu^2)} + \frac{11N_{color} - 2n_{flavor}}{12\pi} \log\left(\frac{|Q^2|}{\mu^2}\right)$$

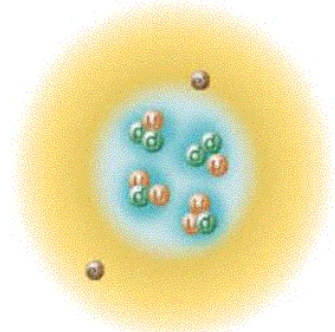
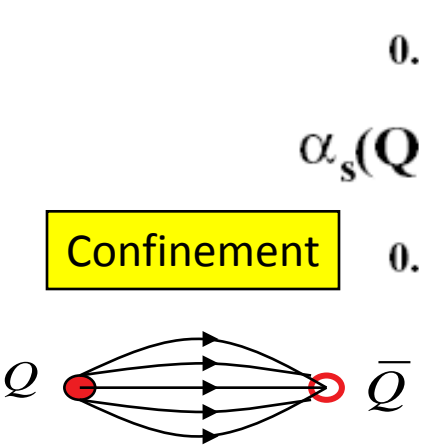
$=+(33-12)/12\pi = \text{positive!}$

Smaller $|Q^2|$ (larger distance) \Rightarrow larger coupling

And that makes a huge difference!

QCD: running of α_s

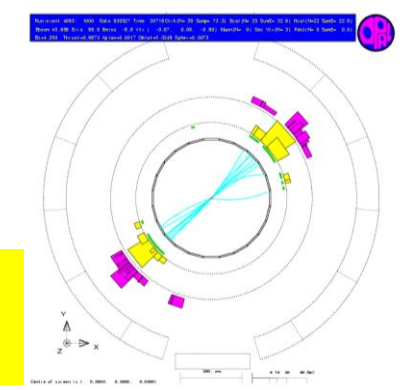
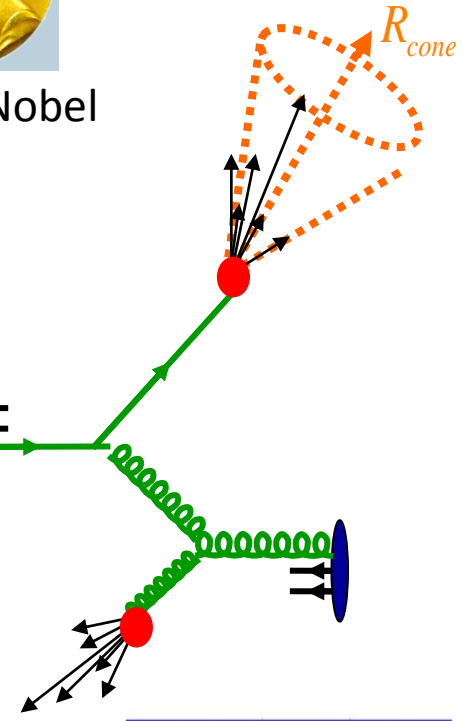
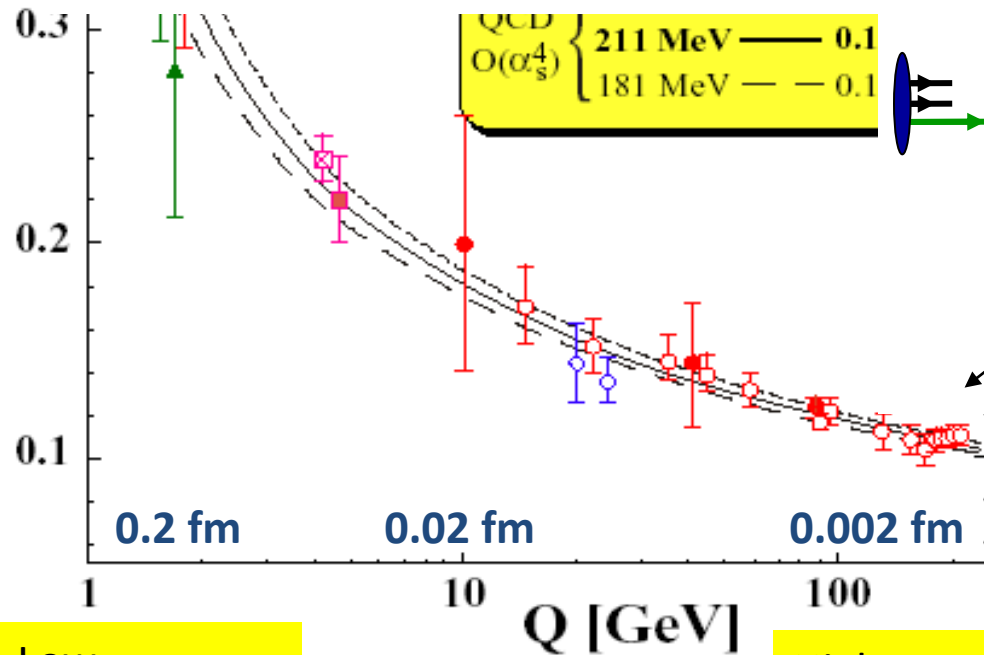
Asymptotic Freedom



Gross Politzer Wilczek

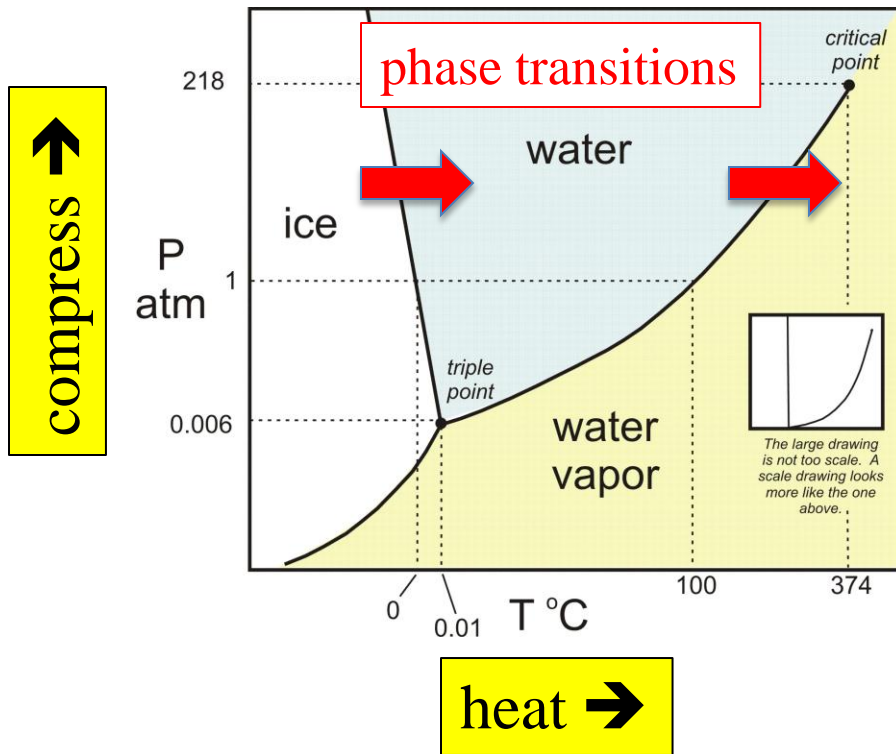


2004 Nobel Prize

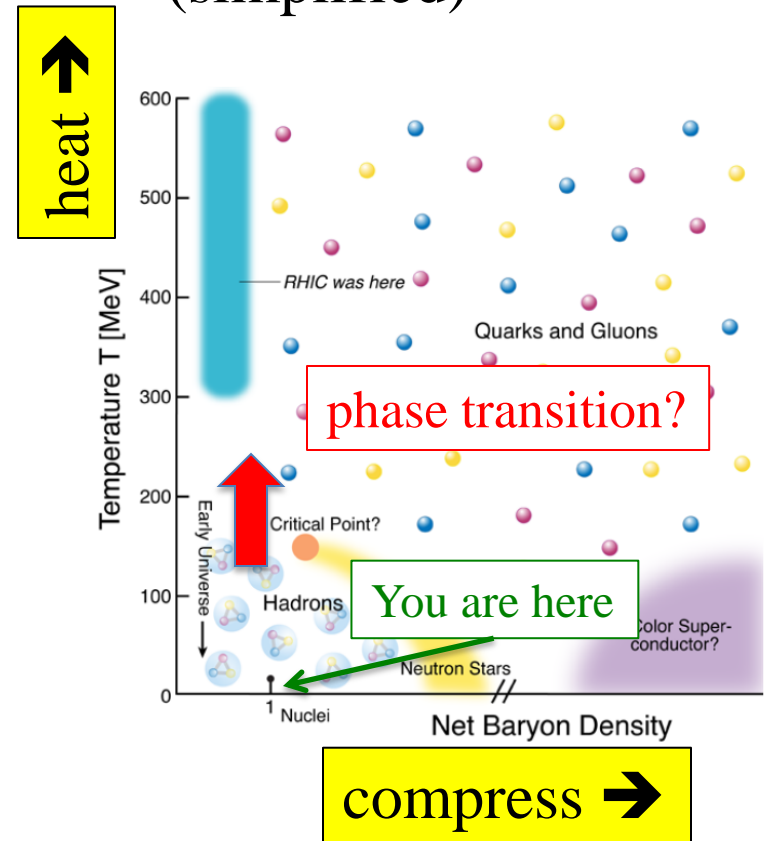


Now let's think about "matter"

Phase diagram of water (simplified)



Phase diagram of QCD (simplified)



Quantitative QCD thermodynamics

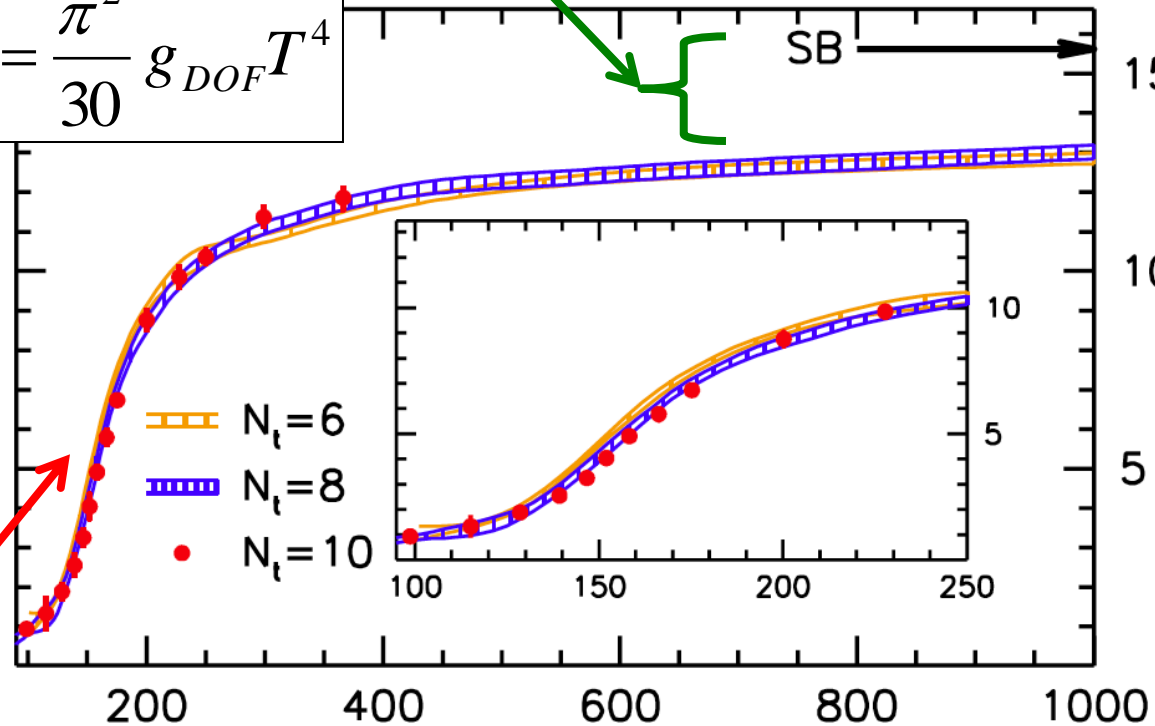
Finite temperature QCD calculated numerically on the lattice ($\mu_B=0$)

Slow convergence to non-interacting Steffan-Boltzmann limit
 What carries energy - complex bound states of q+g? “strongly-coupled” plasma?

Energy density

$$\varepsilon = \frac{\pi^2}{30} g_{DOF} T^4$$

$$\frac{\varepsilon}{T^4}$$



Cross-over, not sharp phase transition
 (like ionization of atomic plasma)

Temperature [MeV]

S. Borsanyi et al., JHEP 1011, 077 (2010)

Exploration of hot QCD Matter: what are the questions? (partial list)

What is the nature of QCD Matter at finite temperature?

- What is its phase structure?
- What is its equation of state?
- What are its effective degrees of freedom?
 - Is it a (trivial) gas of non-interacting quarks and gluons, or a fluid of interacting quasi-particles?
- What are its symmetries?
- Is it correctly described by Lattice QCD or does it require new approaches, and why?

What are the dynamics of QCD matter at finite temperature?

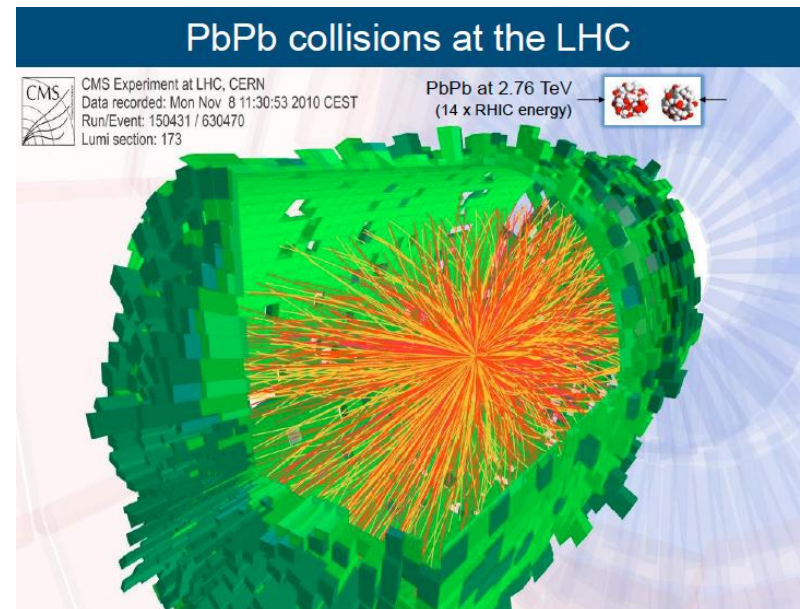
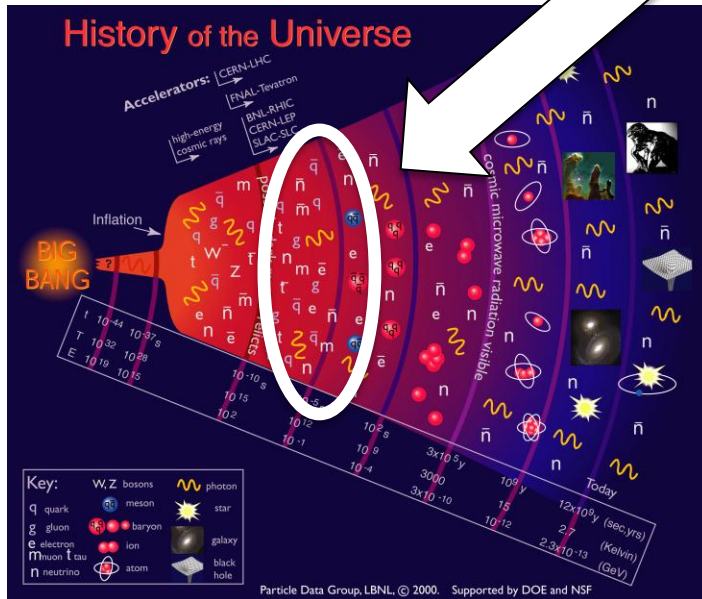
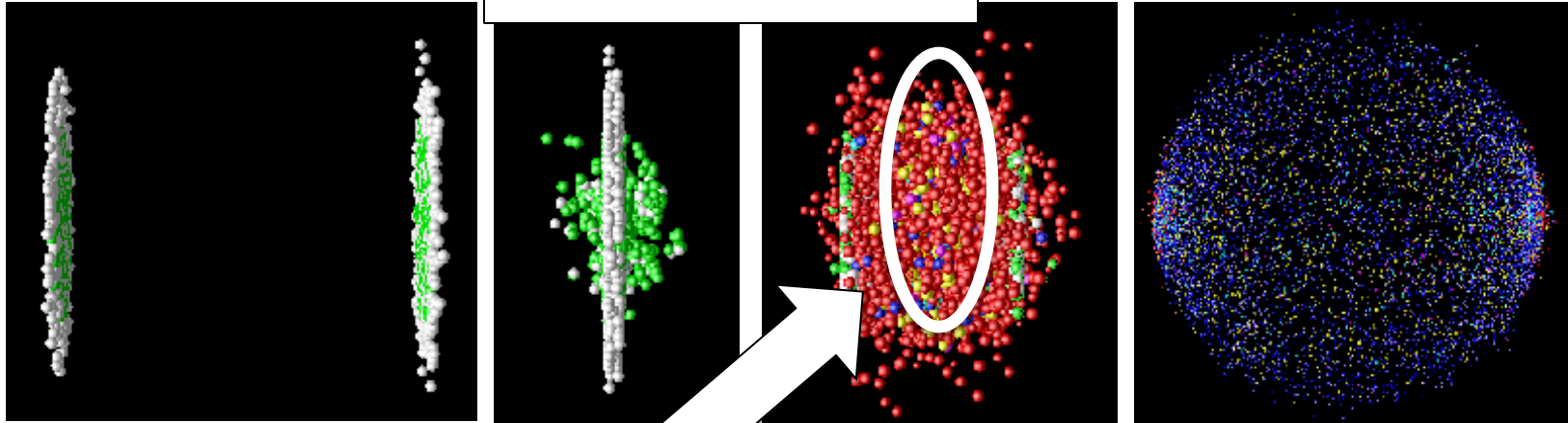
- What is the order of the (de-)confinement transition?
- How is chiral symmetry restored at high T, and how?
- Is there a QCD critical point?
- What are its transport properties?

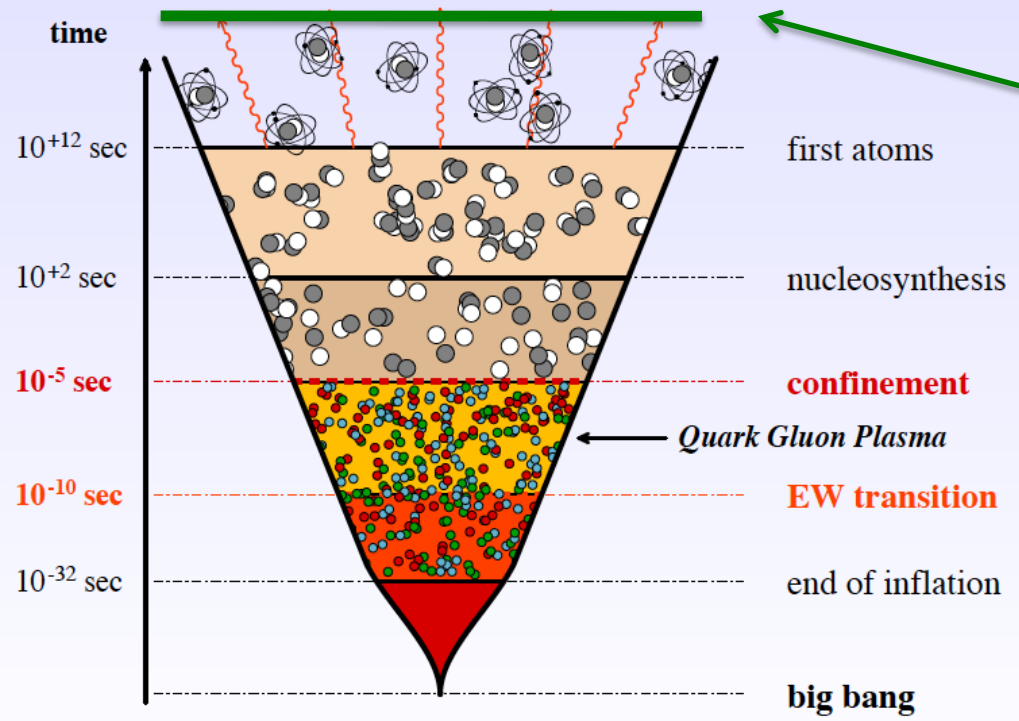
Can hot QCD matter be related to other physical systems?

Can we study hot QCD matter experimentally?

Studying hot QCD in the Laboratory: high energy collisions of heavy nuclei

Model calculation



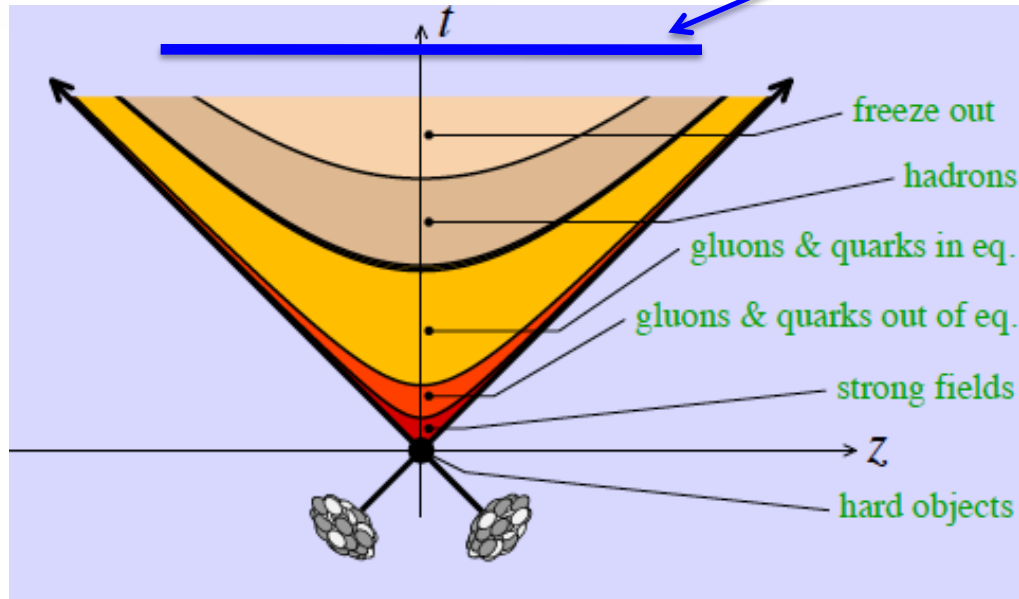


Measurements with telescopes

Evolution of the Early Universe

Measurements with colliders

Evolution of a Heavy Ion Collision



Experimental exploration of hot QCD Matter: what are the issues?

Intensive thermodynamic quantities ($T, P, \varepsilon, \mu, \dots$) are only defined for systems in (quasi- or local-) equilibrium

- QCD Lattice calculates equilibrated matter (e.g. at fixed T)

But nuclear collisions are highly dynamic:

- “Fireball” starts blowing apart the instant it is generated
- Fireball lifetime \sim few fm/c
- no *a priori* reason that quasi-equilibration should be achieved on this time-scale

No *ab initio* theory to describe full dynamical evolution of the fireball

Experimental study of hot QCD Matter: Strategy

Experiment:

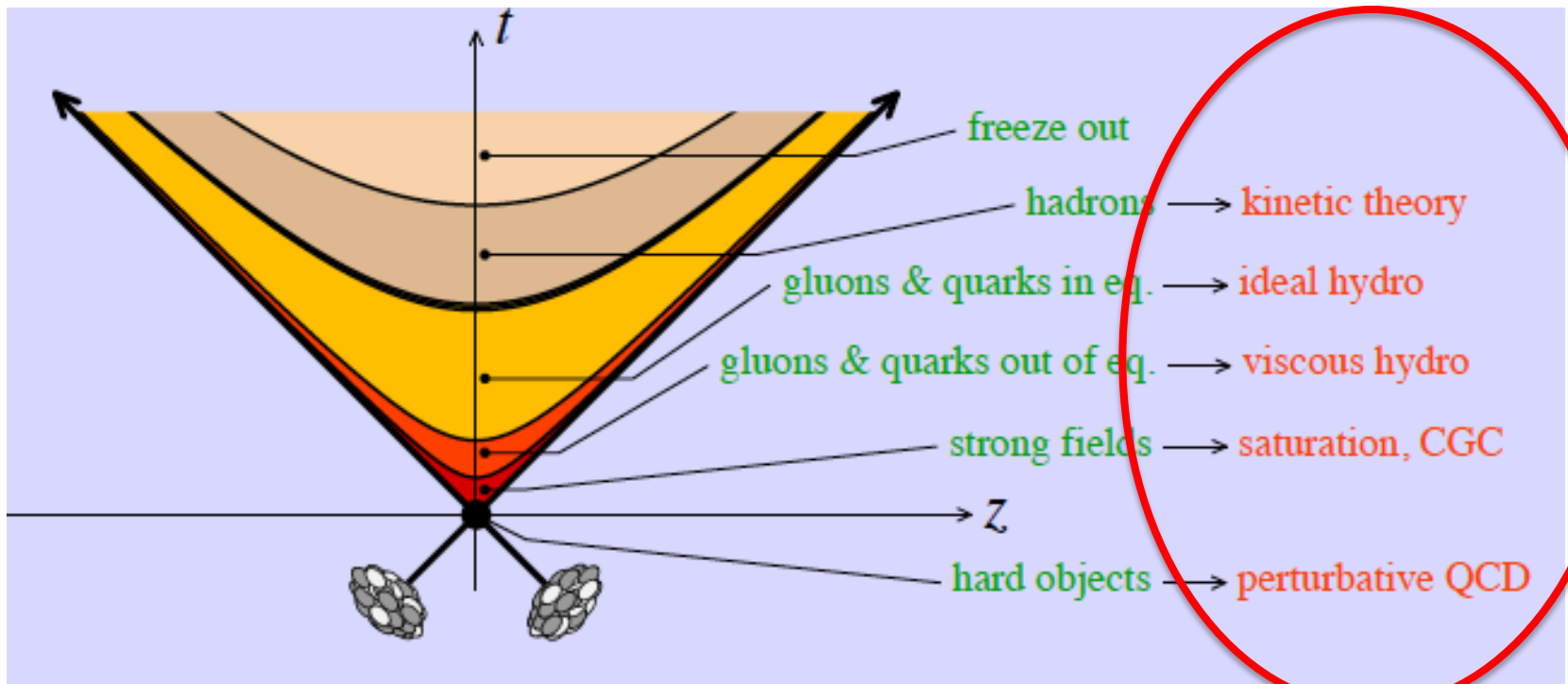
- No *ab initio* theory → interpretation via comparison to reference systems:
p+p, p/d+A, light ion collisions,...
- Vary system size: quantitative control over collision geometry
- Choose observables with close connection to theory and controlled modeling
- Over-determined measurements: multiple, systematically ~independent observables sensitive to the same underlying physics

Theory: models and effective theories for different stages of fireball evolution

- initial state: modified pdfs, saturation models,...
- hard probes: pQCD-based modeling
- collective expansion: viscous relativistic hydrodynamics
- hadronic phase: detailed Monte Carlos

Experiment+Theory:

- detailed comparison and mutual calibration
- evolution with \sqrt{s} : RHIC vs LHC



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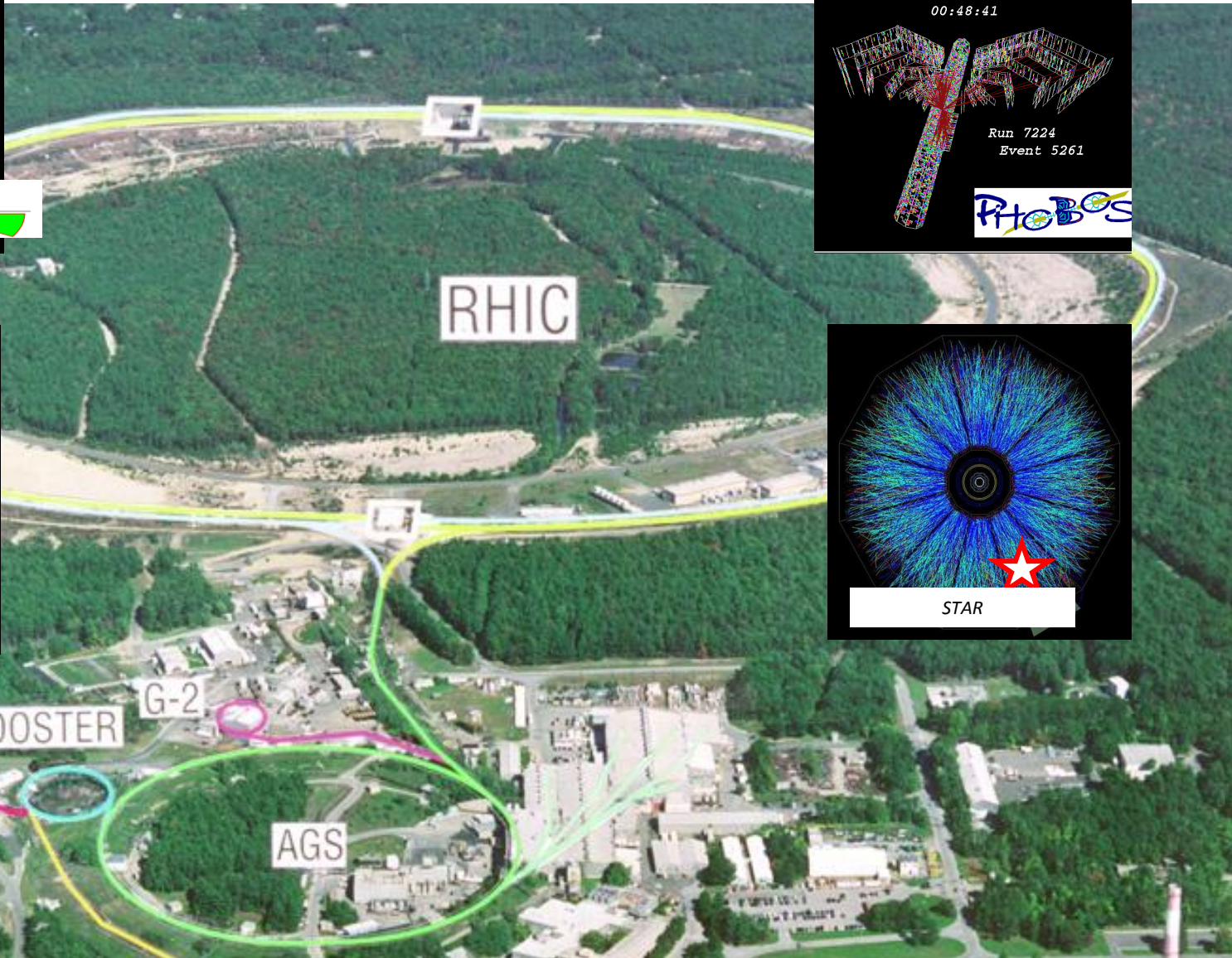
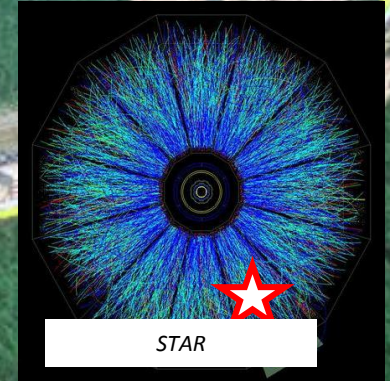
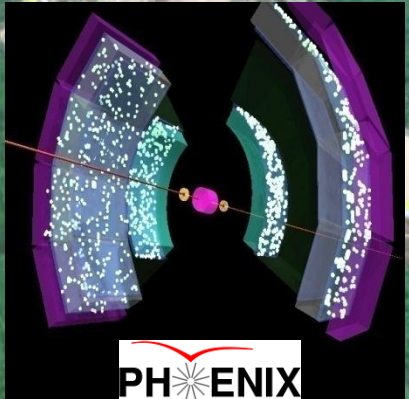
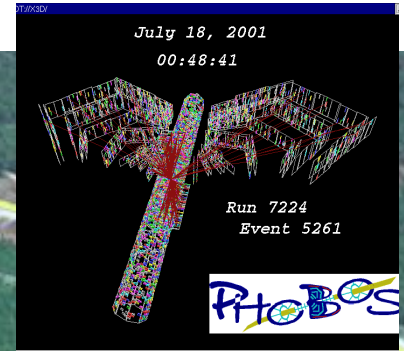
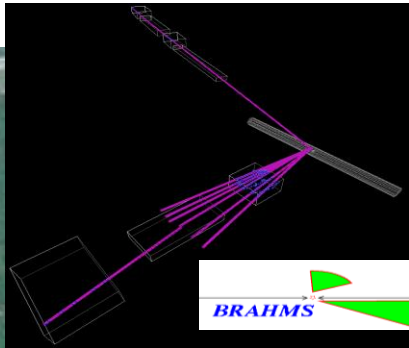
Analysis Tools

- Relativistic Kinematics
- Characterization of nuclear collisions

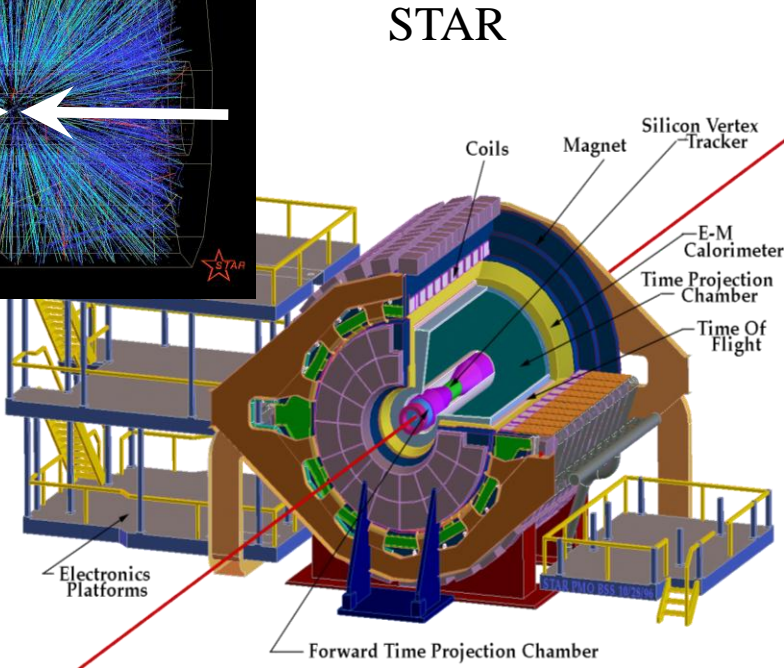
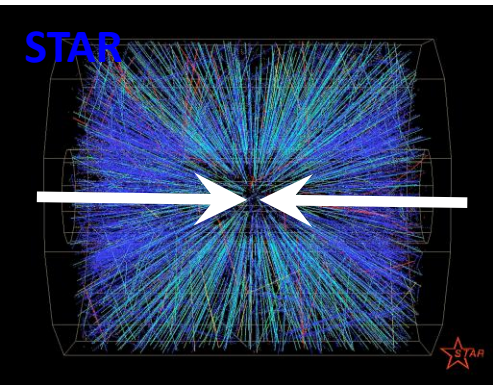
The Relativistic Heavy Ion Collider Brookhaven National Laboratory



The Relativistic Heavy Ion Collider (BNL)



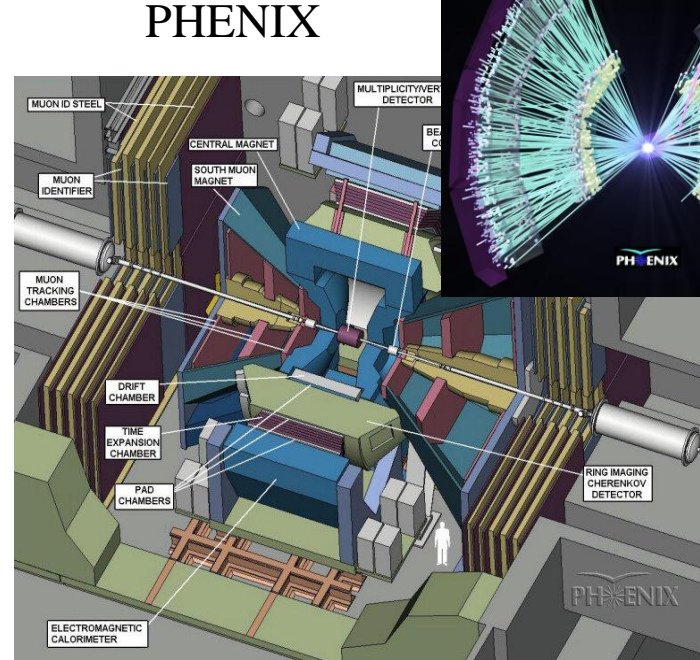
STAR and PHENIX at RHIC



2π coverage, $-1 < \eta < 1$
for tracking + (coarse) EMCal

PID by TOF, dE/dx (STAR), RICH (PHENIX)

Optimised for acceptance
(correlations, jet-finding)



Partial coverage $2 \times 0.5\pi$, $-0.35 < h < 0.35$
Finely segmented calorimeter
+ forward muon arm

Optimised for high-pt π^0 , γ , e , J/ψ
(EMCal, high trigger rates)

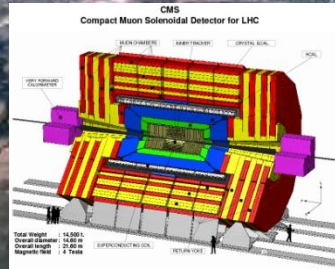
(PHOBOS, BRAHMS, more specialised)

Large Hadron Collider at CERN

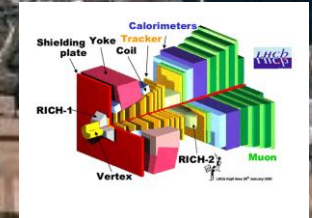
p+p at $\sqrt{s}=7$ (14) TeV

Pb+Pb at $\sqrt{s}=2.76$ (5.5) TeV

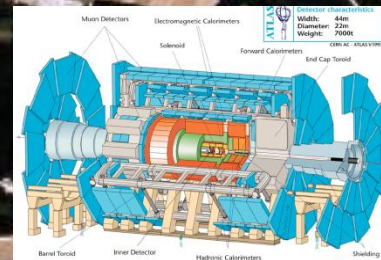
heavy ion running: 4 physics weeks/year



CMS



LHCb

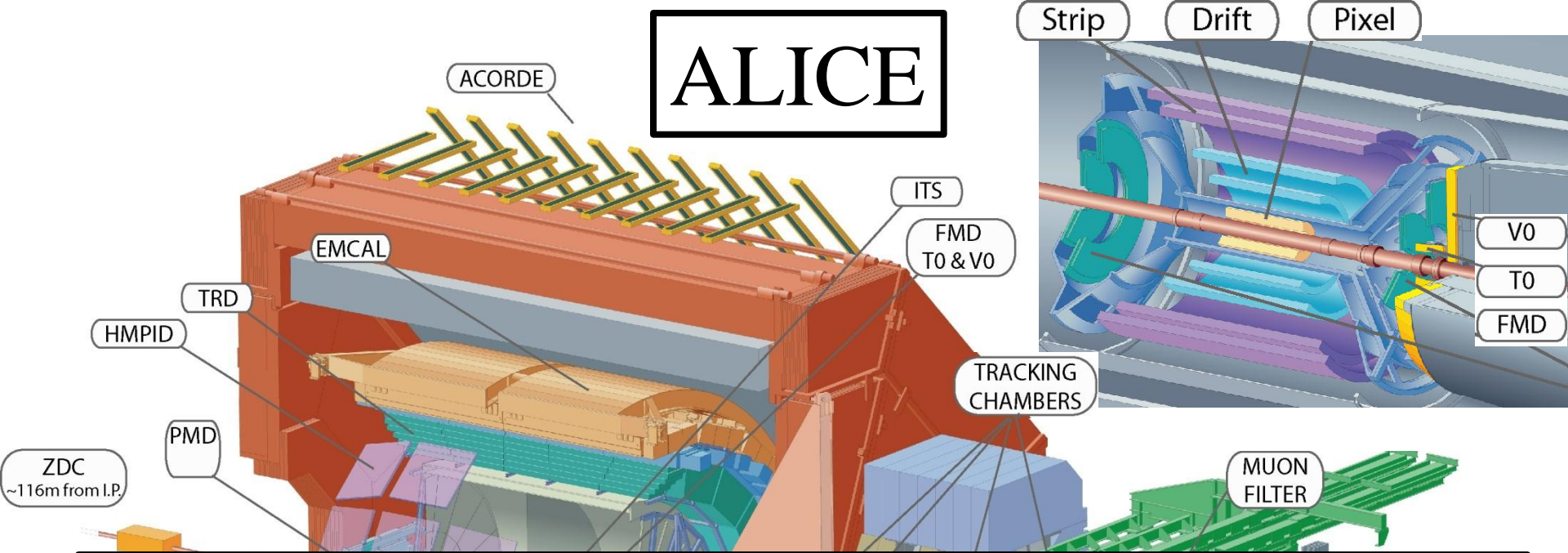


ATLAS



ALICE

ALICE



ALICE is the comprehensive heavy ion experiment at the LHC

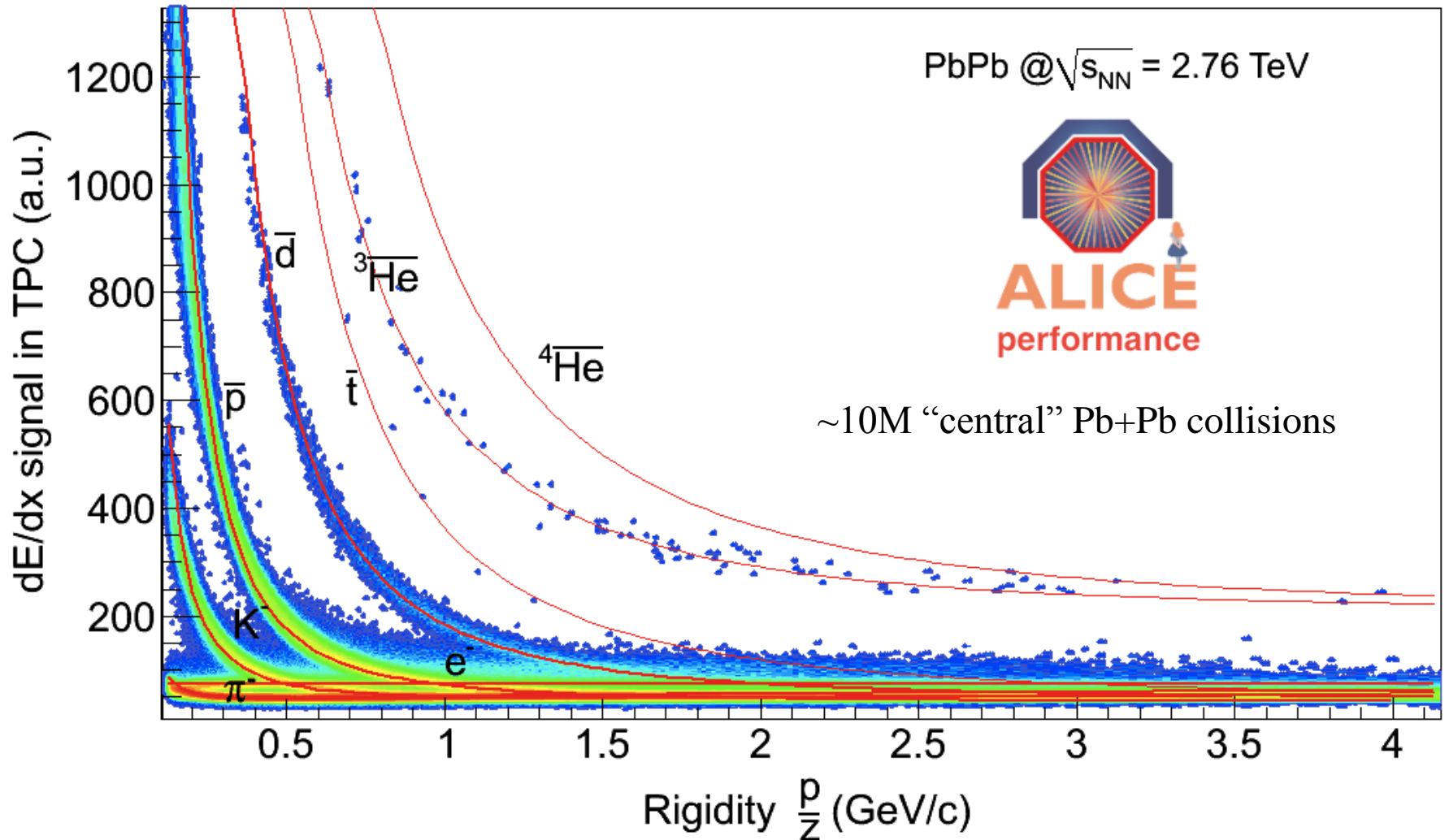
Design optimized for

- huge particle multiplicities of nuclear collisions
- efficient tracking over wide momentum range
- extensive particle identification
- low mass around vertex → low p_T measurements

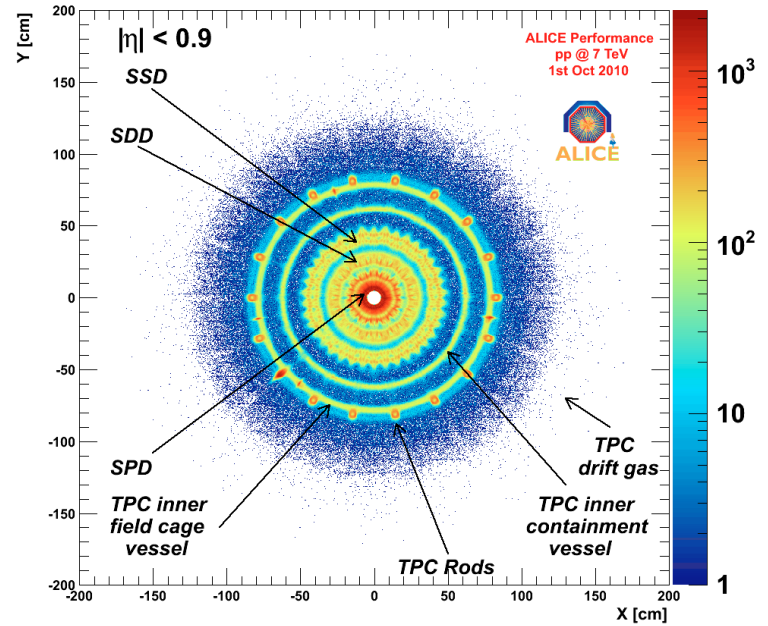
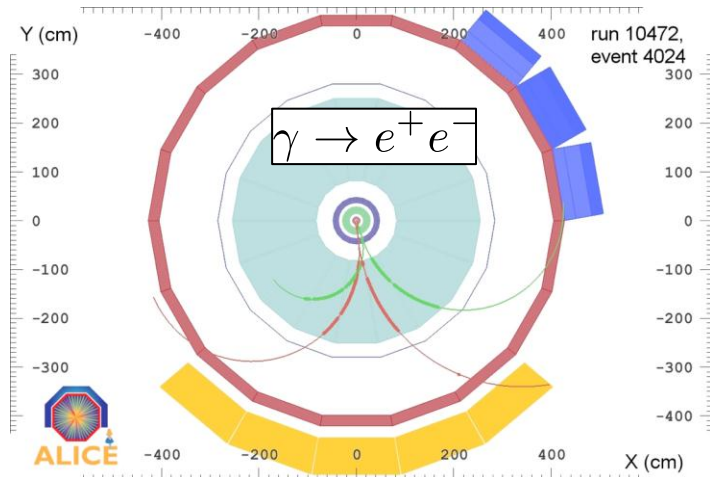
QGP “temperature” $\sim \Lambda_{\text{QCD}} \sim$ few hundred MeV

Detector Performance

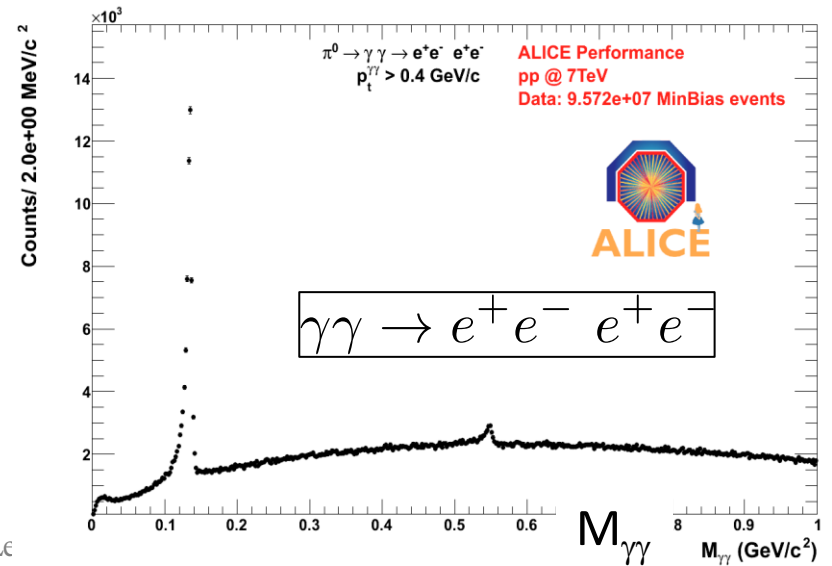
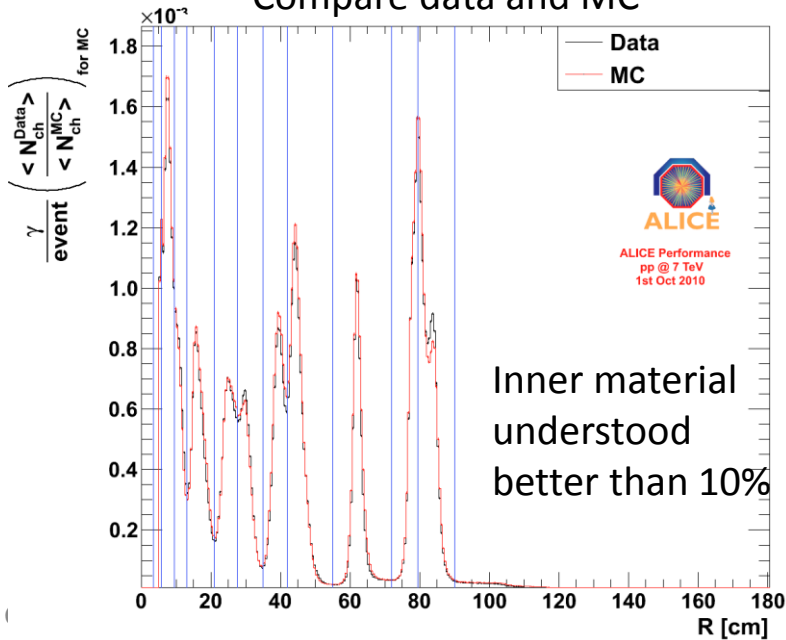
ALICE Particle ID (TPC dE/dx)



ALICE: Tomography via γ -conversions



Compare data and MC



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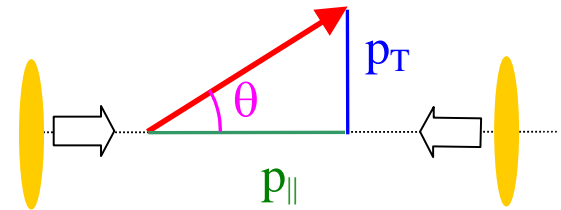
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Kinematics for Inclusive Reactions

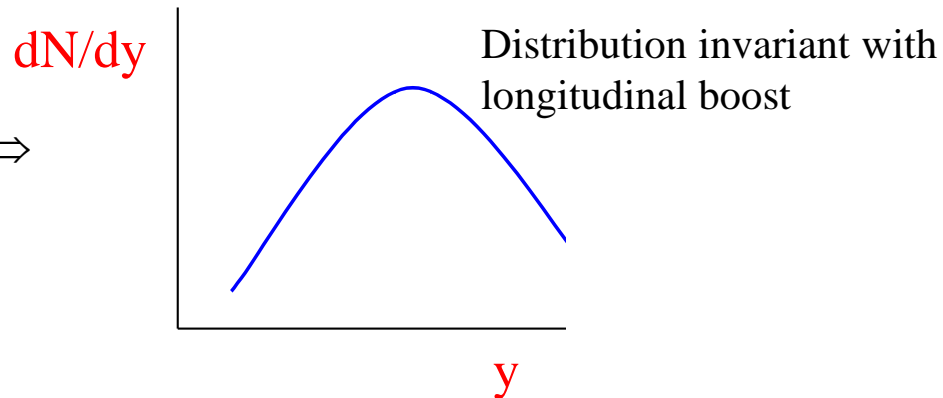
Rapidity

$$y = \frac{1}{2} \ln \left(\frac{E + p_{\parallel}}{E - p_{\parallel}} \right)$$



Rapidity is differentially boost-invariant

$$\delta y \sim \frac{\delta p_{\parallel}}{E} \Rightarrow$$



Pseudo-rapidity

$$y \rightarrow \eta = -\ln \left[\tan(\theta/2) \right]$$

for $m/p \ll 1$

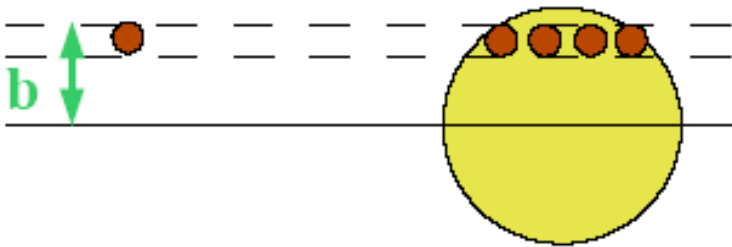
Invariant production cross section

$$E \frac{d^3 \sigma}{d^3 p} = \frac{d^2 \sigma}{2\pi p_T dy dp_T}$$

Nuclear geometry and hard processes: Glauber theory

Glauber scaling for hard processes with large momentum transfer

- short coherence length \Rightarrow successive NN collisions independent
- p+A is incoherent superposition of N+N collisions



Normalized nuclear density $r(b, z)$:

$$\int dz db \rho(b, z) = 1$$

Nuclear thickness function

$$T_A(b) = \int_{-\infty}^{\infty} dz \rho(b, z)$$

Inelastic cross section for
p+A collisions:

$$\sigma_{pA}^{inel} = \int d\vec{b} \left(1 - [1 - T_A(b) \sigma_{NN}^{inel}]^A \right)$$

$$\sigma_{pA}^{hard} \simeq A \cdot \sigma_{NN}^{hard} \int d\vec{b} T_A(b) = A \sigma_{NN}^{hard}$$

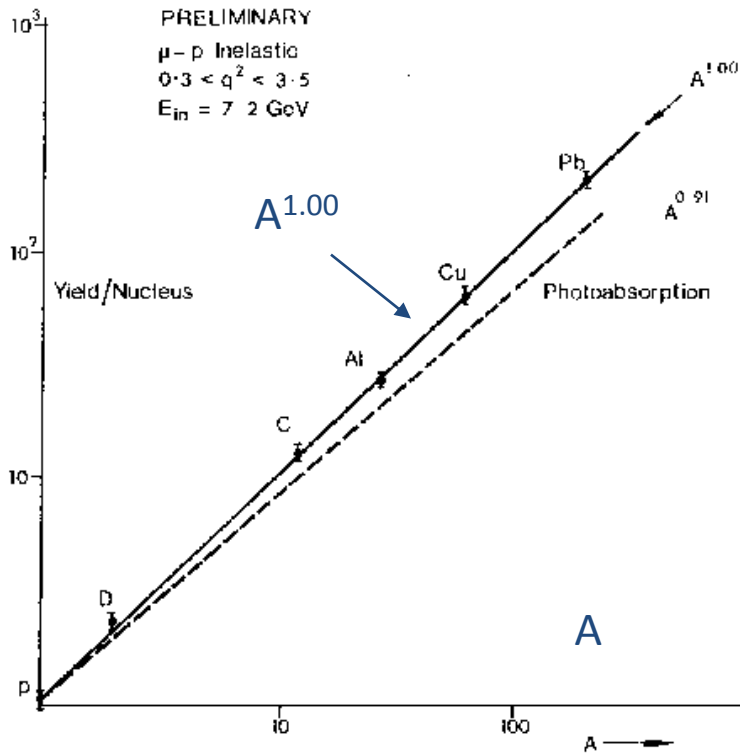
Experimental tests of Glauber scaling: hard cross sections in p(μ)+A collisions

Glauber scaling expectation:

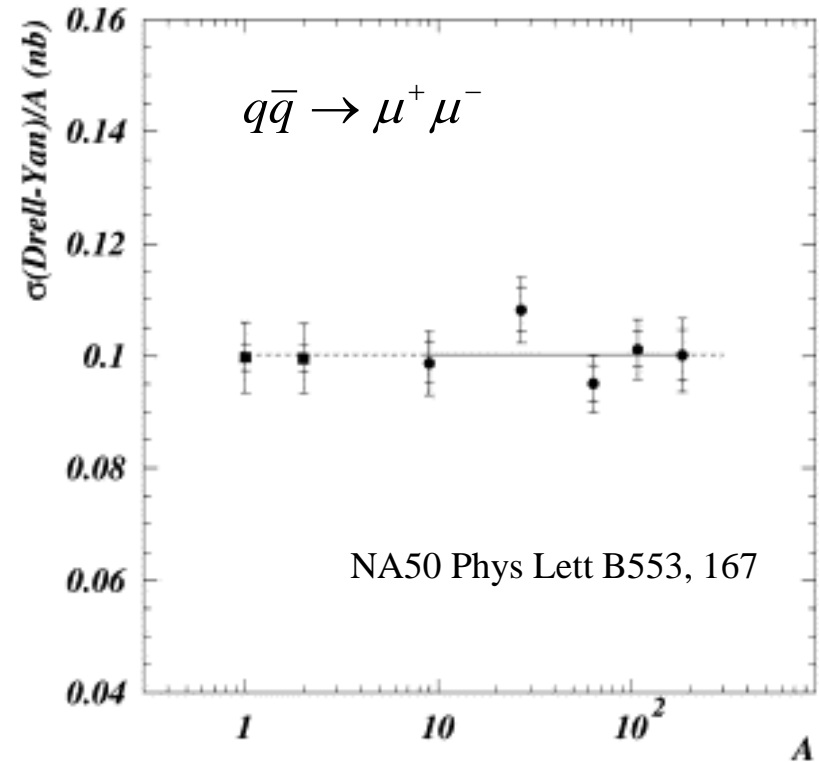
$$\sigma_{pA}^{hard} = A \sigma_{NN}^{hard}$$

σ_{inel} for 7 GeV muons on nuclei

M. May et al, Phys Rev Lett 35, 407 (1975)



$\sigma_{Drell-Yan}/A$ in p+A at SPS



Hard cross sections in p+A scale as $A^{1.0}$

Glauber Theory for A+B Collisions

Nuclear overlap function:

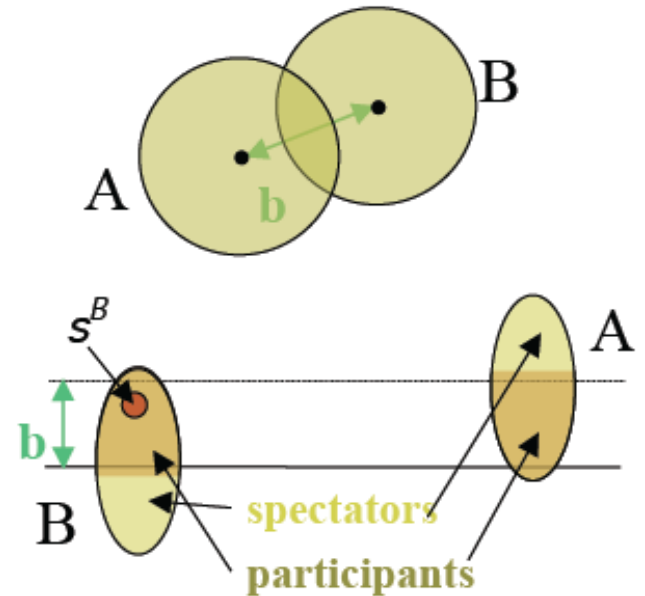
$$T_{AB}(\vec{b}) = \int d\vec{s} T_A(\vec{s}) T_B(\vec{s} - \vec{b})$$

Average number of binary NN collisions for B nucleon at coordinate \vec{s}_B :

$$N_{bin}^{nA}(\vec{b} - \vec{s}_B) = A \cdot T_A(\vec{b} - \vec{s}_B) \cdot \sigma_{nn}^{inel}$$

Average number of binary NN collisions for A+B collision with impact parameter b :

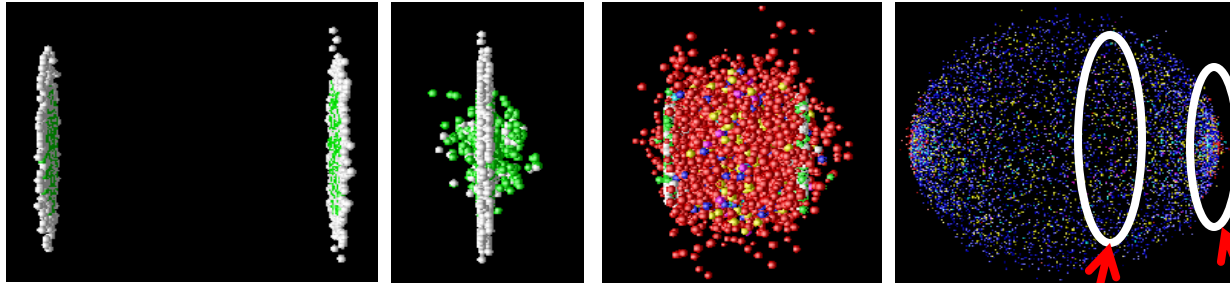
$$\begin{aligned} N_{bin}^{AB}(b) &= B \int d\vec{s}_B T_B(\vec{s}_B) \cdot N_{bin}^{nA}(\vec{b} - \vec{s}_B) \\ &= AB \cdot T_{AB}(b) \cdot \sigma_{nn}^{inel} \end{aligned}$$



Measuring collision geometry I

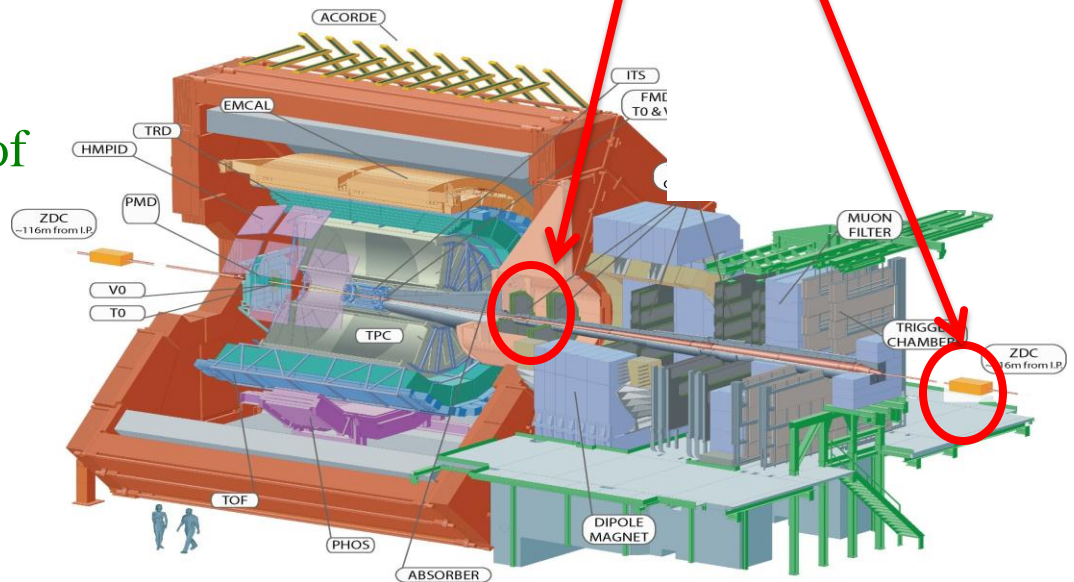
Nuclei are “macroscopic”

→ characterize collisions by impact parameter



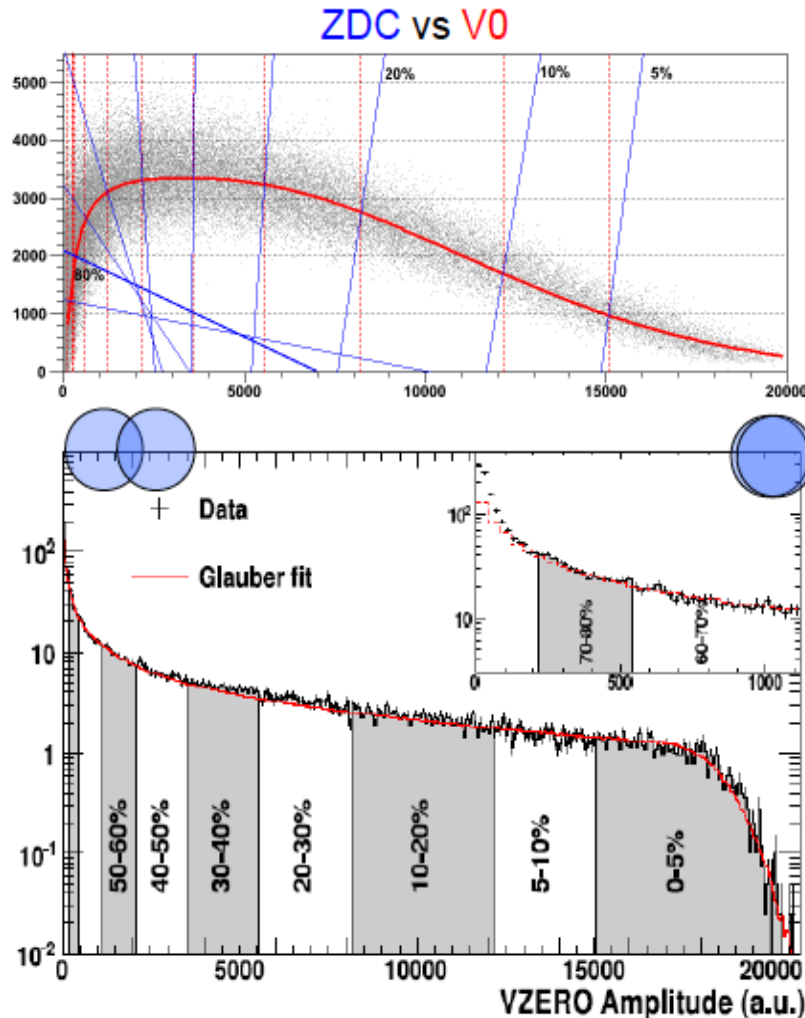
Correlate particle yields from
~causally disconnected parts of
phase space

→ correlation arises from
common dependence on
collision impact parameter



Measuring collision geometry II

Forward neutrons



Charged hadrons $\eta \sim 3$

- Order events by centrality metric
- Classify into percentile bins of “centrality”

HI jargon: “0-5% central”

Connect to Glauber theory via particle production model:

- N_{bin} : effective number of binary nucleon collisions (~5-10% precision)
- N_{part} : number of (inelastically scattered) “participating” nucleons

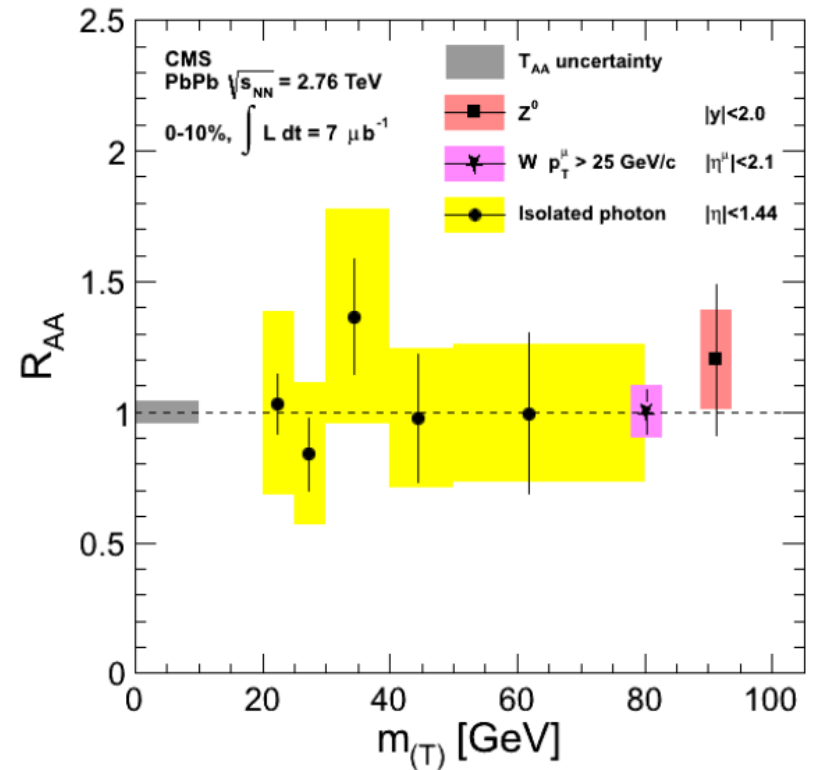
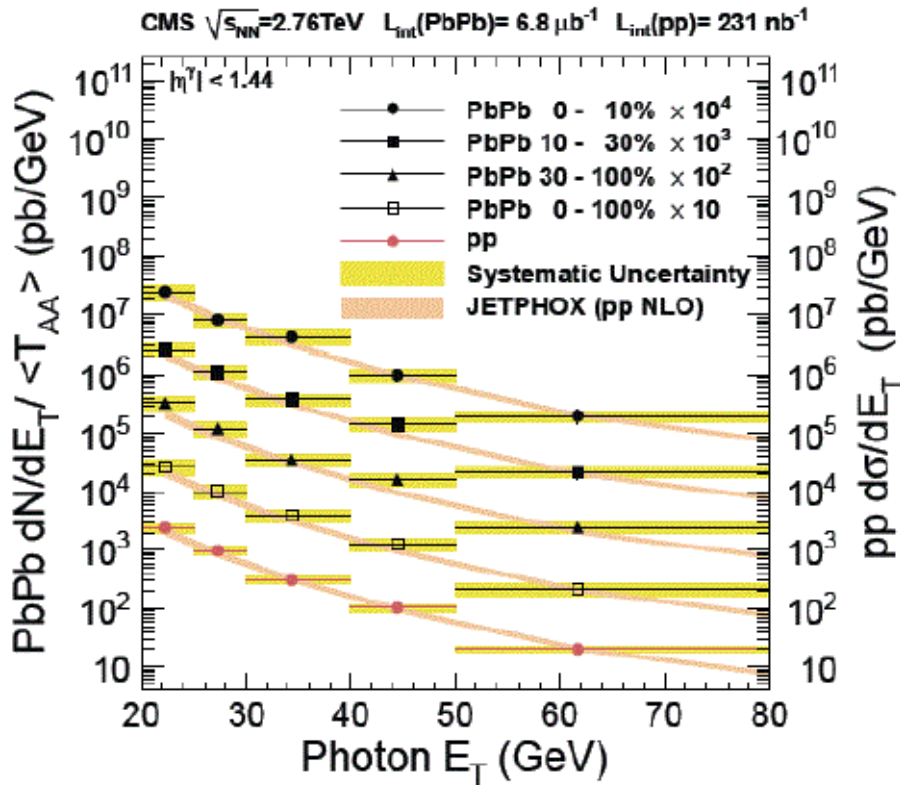
Scaling of cross sections using Glauber theory plays a central role in quantitative analysis of experimental measurements and connection to theory.

Let's test it experimentally in A+A collisions...

Glauber test at LHC:

Scaling of direct photon, Z, W yields in Pb+Pb vs p+p

$$R_{AA} = \frac{d\sigma_{AA}^{hard}/dp_T}{\langle T_{AA} \rangle \cdot d\sigma_{pp}^{hard}/dp_T}$$



EW boson yields all scale with N_{bin} : Glauber OK for hard processes

Summary of Lecture 1: what are the questions? (partial list)

What is the nature of QCD Matter at finite temperature?

- What is its phase structure?
- What is its equation of state?
- What are its effective degrees of freedom?
 - Is it a (trivial) gas of non-interacting quarks and gluons, or a fluid of interacting quasi-particles?
- What are its symmetries?
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What are the dynamics of QCD matter at finite temperature?

- What is the order of the (de-)confinement transition?
- How is chiral symmetry restored at high T, and how?
- Is there a QCD critical point?
- What are its transport properties?

Can QCD matter be related to other physical systems?

Can we study hot QCD matter experimentally?

References

QCD

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http://pdg.lbl.gov/2004/reviews/contents_sports.html
- QCD and jets: CTEQ web page and summer school lectures
<http://www.phys.psu.edu/~cteq/>
- Handbook of Perturbative QCD, Rev. Mod. Phys. 67, 157–248 (1995)
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Heavy Ion Physics

- Results from the Relativistic Heavy Ion Collider, B. Mueller and J. Nagle; Ann. Rev. Nucl. Part. Sci. 56, 93 (2006), nucl-th/0602029
- Heavy Ion Collisions at the LHC – Last Call for Predictions, N. Armesto et al. (ed.); J. Phys. G35 054001 (2008), arXiv:0711.0974
- New Developments in Relativistic Viscous Hydrodynamics, P. Romatschke; Int. J. Mod. Phys. E19, 1-53 (2010), arXiv:0902.3663
- The theory and phenomenology of perturbative QCD-based jet quenching, A. Majumder and M. van Leeuwen; arXiv:1002.2206
- Gauge/String Duality, Hot QCD and Heavy Ion Collisions, J. Casalderrey-Solana et al.; arXiv:1101.0618

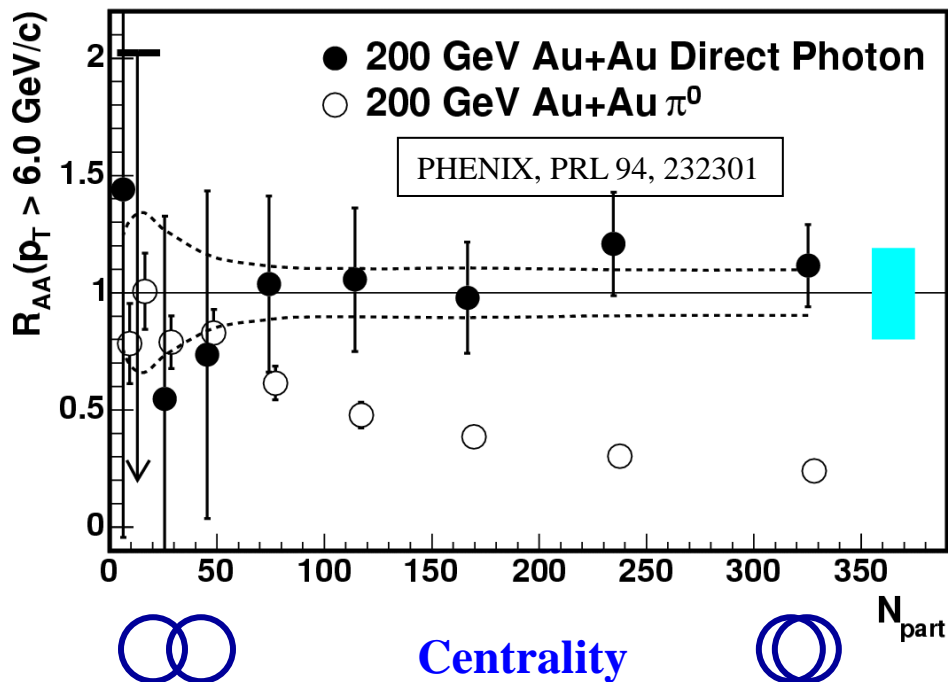
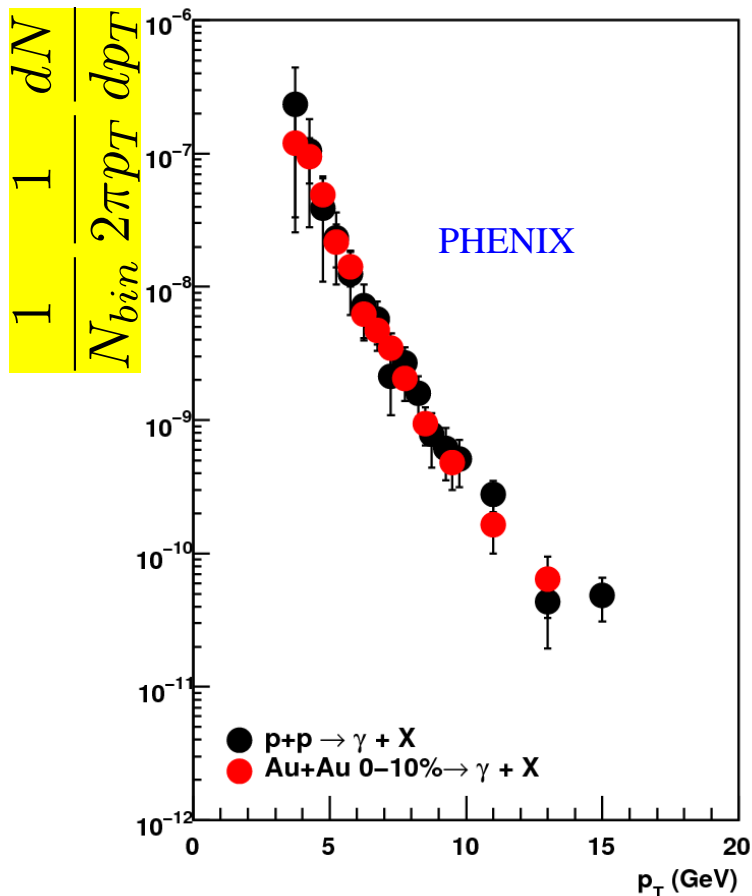
Backup

Glauber test at RHIC:

Scaling of direct photon yield in p+p vs. Au+Au

Direct γ : N_{bin} -scaled
inclusive yield

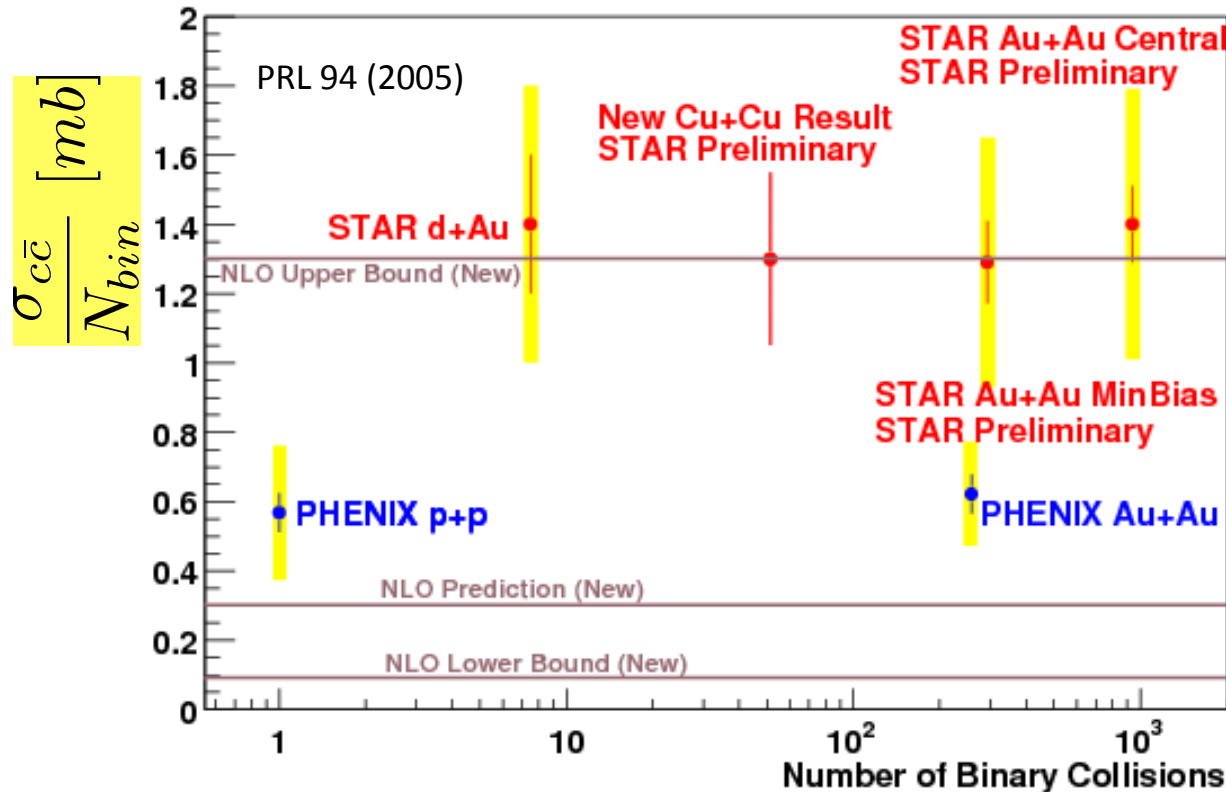
$$R_{AA} = \frac{dN_{Au+Au}/dp_T}{N_{bin} \cdot dN_{p+p}/dp_T}$$



Direct γ yield scales with N_{bin}

Glauber test at RHIC:

Scaling of charm total production cross section



NLO prediction:
 $m \approx 1.3 \text{ GeV}$, reasonably
 hard scale at $p_T=0$

Total charm cross section scales with N_{bin} in A+A

(Sizable disagreement between STAR and PHENIX?)