

# SMEFT STUDIES WITH EIC AND LHEC DIS PSEUDO DATA

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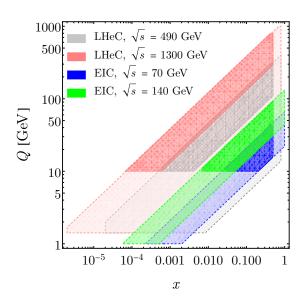




# **OUR WORK IN A NUTSHELL**

ightharpoonup We study the m BSM potential of the LHeC and the EIC

detailed accounting of anticipated uncertainties on pseudo data



- Multidimensional fits
  - of NC DIS cross section and asymmetries performed in the **SMEFT** framework

★ We show that both the EIC and LHeC can improve upon the existing bound on the Z-boson couplings





# **BEFORE STARTING ...**



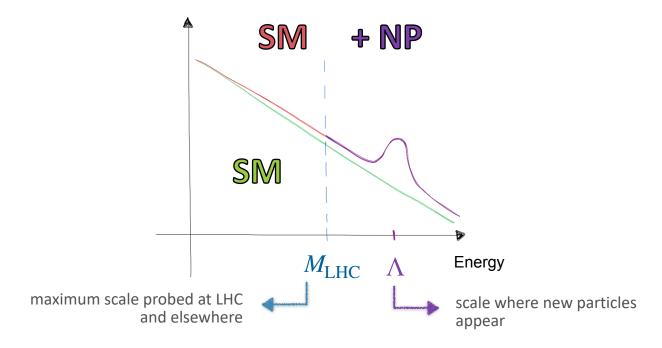
one possibile idea could be to use xFitter to **benchmark** our findings





# **BEYOND STANDARD MODEL**





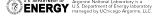
All **new physics** is assumed to be heavier than all SM states and accessible collider energy



SMEFT

Standard Model Effective Field Theory





# **SMEFT**

#### **Standard Model Effective Field Theory**

$$\mathscr{L}_{\mathrm{SMEFT}} = \mathscr{L}_{\mathrm{SM}} + \sum_{n>4} \frac{1}{\Lambda^{n-4}} \sum_k \mathcal{C}_k^{(n)} O_k^{(n)}$$
 in this work  $n=6$ 

- → Patterns and correlations among operators and observables are key
- ✓ What might we find?

Best case: a non-zero value for  $\mathcal{C}_k$ 

indicating a mass scale slightly above probed values

Otherwise: stringent constraints on the  $\mathcal{C}_k$ 

suggest where to focus future searches

# **WILSON COEFFICIENTS**

#### **Dimension 6 operators**

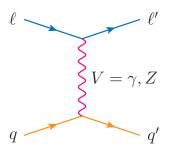
Warsaw basis

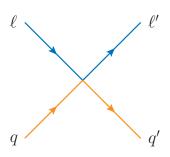
# 17 Wilson coefficients affect NC DIS matrix elements at LO

ffV			semi-leptonic four-fermion	
$C_{\varphi WB}$	$O_{\varphi WB} = (\varphi^{\dagger} \tau^I \varphi) W^I_{\mu\nu} B^{\mu\nu}$	$C_{\ell q}^{(1)}$	$O_{\ell q}^{(1)} = (ar{\ell} \gamma_{\mu} \ell) (ar{q} \gamma^{\mu} q)$	
$C_{arphi D}$	$O_{\varphi D} = (\varphi^{\dagger} D_{\mu} \varphi)^* (\varphi^{\dagger} D^{\mu} \varphi)$	$C_{\ell q}^{(3)}$	$(ar{\ell}\gamma_{\mu} au^I\ell)(ar{q}\gamma^{\mu} au^Iq)$	
$C_{arphi\ell}^{(1)}$	$O_{arphi\ell}^{(1)} = (arphi^\dagger i\stackrel{\leftrightarrow}{D}_\mu arphi)(ar{\ell}\gamma^\mu\ell)$	$C_{eu}$	$O_{eu}=(ar{e}\gamma_{\mu}e)(ar{u}\gamma^{\mu}u)$	
$C_{arphi\ell}^{(3)}$	$O_{arphi\ell}^{(3)} = (arphi^\dagger i \stackrel{\leftrightarrow}{D}_\mu   au^I arphi) (ar{\ell} \gamma^\mu  au^I \ell)$	$C_{ed}$	$O_{ed}=(ar{e}\gamma_{\mu}e)(ar{d}\gamma^{\mu}d)$	
$C_{arphi e}$	$O_{arphi e} = (arphi^\dagger i\stackrel{\leftrightarrow}{D}_\mu arphi)(ar{e}\gamma^\mu e)$	$C_{\ell u}$	$O_{\ell u} = (ar{\ell} \gamma_{\mu} \ell) (ar{u} \gamma^{\mu} u)$	
$C_{arphi q}^{(1)}$	$O^{(1)}_{arphi q} = (arphi^\dagger i\stackrel{\leftrightarrow}{D}_\mu arphi) (ar{q} \gamma^\mu q)$	$C_{\ell d}$	$O_{\ell d} = (ar{\ell} \gamma_\mu \ell) (ar{d} \gamma^\mu d)$	
$C_{arphi q}^{(3)}$	$O_{\varphi q}^{(3)} = (\varphi^{\dagger} i \stackrel{\leftrightarrow}{D}_{\mu} \tau^{I} \varphi)(\bar{q} \gamma^{\mu} \tau^{I} q)$	$C_{qe}$	$O_{qe}=(ar{q}\gamma_{\mu}q)(ar{e}\gamma^{\mu}e)$	
$C_{arphi u}$	$O_{arphi u} = (arphi^\dagger i \stackrel{\leftrightarrow}{D}_\mu arphi) (ar{u} \gamma^\mu u)$	£.0	q: left handed doublets	
$C_{arphi d}$	$O_{arphi d} = (arphi^\dagger i\stackrel{\leftrightarrow}{D}_\mu arphi)(ar{d}\gamma^\mu d)$	e, u, d: right handed singlets		
$C_{\ell\ell}$	$O_{\ell\ell} = (\bar{\ell}\gamma_{\mu}\ell)(\bar{\ell}\gamma^{\mu}\ell)$	arphi : SU(2) Higgs doublet		



# **DEEP INELASTIC SCATTERING**





# 

Feynman diagrams for the partonic process mediating

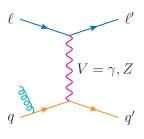
$$\ell + H \to \ell' + X$$

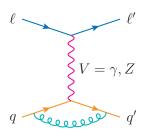


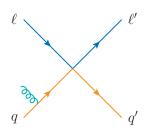
corrections modify only the quark lines

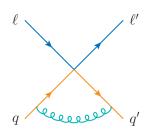


they are identical for both SM and SMEFT cross sections









# **SMEFT THEORY PREDICTIONS**

#### **Structure of observables**

We linearize our SMEFT expressions  $\,\mathcal{T}\,$ 

$$\mathcal{T} = \mathcal{T}^{\mathrm{SM}} + \sum_{k} \mathcal{C}_{k} \, \, \delta \mathcal{T}_{k} + \mathcal{O}(\mathcal{C}_{k}^{2})$$

k runs over the active Wilson coefficients

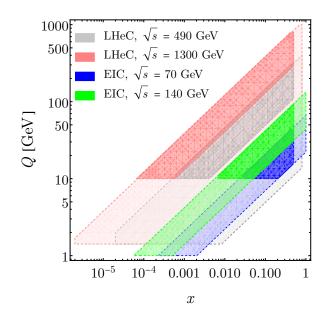
SMEFT shift associated with the Wilson coefficient  $C_k$ 

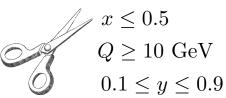
# SMEFT cross section

$$\sigma = \sigma^{\text{SM}} + \sum_{i}^{N_{\text{d6}}} \frac{C_i}{\Lambda^2} k_i + \sum_{ij}^{N_{\text{d6}}} \frac{C_i C_j}{\Lambda^4} \tilde{k}_{ij} + \dots$$

# KINEMATIC COVERAGE

#### of pseudo data





# ★ complementarity of LHeC and EIC

Data set label	Data set configuration	Observable		
LHeC1	$60 \text{ GeV} \times 1000 \text{ GeV } e^- p, \ P_{\ell} = 0, \ \mathcal{L} = 100 \text{ fb}^{-1}$			
m LHeC2	60 GeV × 7000 GeV $e^-p$ , $P_\ell = -80\%$ , $\mathcal{L} = 100 \text{ fb}^{-1}$			
m LHeC3	60 GeV × 7000 GeV $e^-p$ , $P_\ell = +80\%$ , $\mathcal{L} = 30 \text{ fb}^{-1}$	<b></b>		
${ m LHeC4}$	60 GeV × 7000 GeV $e^+p$ , $P_\ell = +80\%$ , $\mathcal{L} = 10 \text{ fb}^{-1}$	$\sigma_{ m NC}$		
m LHeC5	60 GeV × 7000 GeV $e^-p$ , $P_{\ell} = -80\%$ , $\mathcal{L} = 1000 \text{ fb}^{-1}$			
${ m LHeC6}$	60 GeV × 7000 GeV $e^-p$ , $P_\ell = +80\%$ , $\mathcal{L} = 300 \text{ fb}^{-1}$			
LHeC7	60 GeV × 7000 GeV $e^+p$ , $P_\ell = 0\%$ , $\mathcal{L} = 100 \text{ fb}^{-1}$			
D4	$10 \text{ GeV} \times 137 \text{ GeV } e^-D, \ P_{\ell} = 80\%, \ \mathcal{L} = 100 \text{ fb}^{-1}$			
D5	18 GeV × 137 GeV $e^-D$ , $P_{\ell} = 80\%$ , $\mathcal{L} = 15.4 \text{ fb}^{-1}$	$igg] A_{ m PV} = igg $		
P4	$10 \text{ GeV} \times 275 \text{ GeV } e^- p, \ P_{\ell} = 80\%, \ \mathcal{L} = 100 \text{ fb}^{-1}$	]		
P5	18 GeV × 275 GeV $e^-p$ , $P_{\ell} = 80\%$ , $\mathcal{L} = 15.4 \text{ fb}^{-1}$			
$\Delta \mathrm{D4}$	The same as D4 but with $P_{\ell} = 0$ and $P_{H} = 70\%$			
$\Delta \mathrm{D5}$	The same as D5 but with $P_{\ell} = 0$ and $P_{H} = 70\%$	$igg  \Delta A_{ m PV}$		
$\Delta \mathrm{P4}$				
$\Delta  ext{P5}$	The same as P5 but with $P_\ell=0$ and $P_H=70\%$			

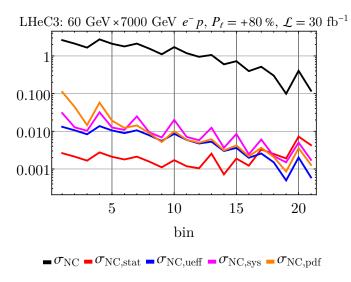




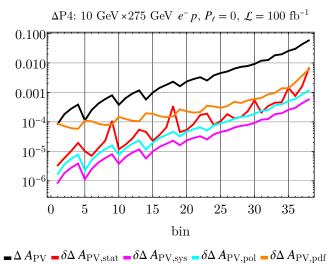
# **UNCERTAINTIES**

bins are ordered first in Q, then in x

# LHeC



# EIC





NNPDF 3.1 NLO

NNPDF POL 1.1

# PSEUDODATA AND $\chi^2$

 $oldsymbol{e}$  experiment b bin



# uncertainties

$$\mathcal{O}_{m{e},m{b}} = \mathcal{O}_{m{b}}^{ ext{SM}} + r_{m{e},m{b}} \; \delta\mathcal{O}_{ ext{unc},m{b}} + \sum_{j} r'_{j,m{e}} \; \delta\mathcal{O}_{ ext{cor}_j,m{b}}$$
 data

# pseudodata

each correlated error has its own random coefficient  $r'_{i,e}$ 

$$\chi_e^2 = \sum_{i=1}^{N_{\text{bin}}} (\mathcal{T}_i - \mathcal{O}_{e,i}) V_{ij}^{-1} (\mathcal{T}_j - \mathcal{O}_{e,j})$$

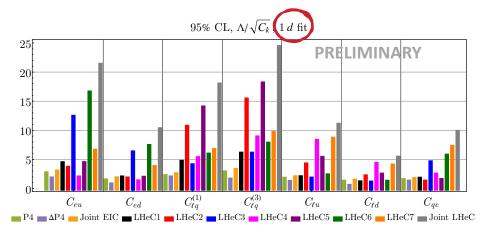


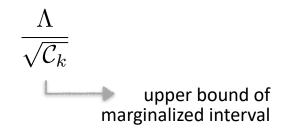


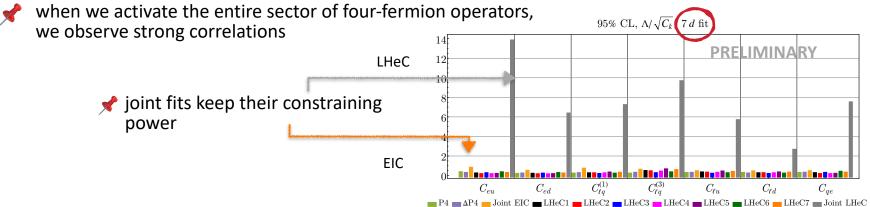




# **EFFECTIVE UV SCALES**

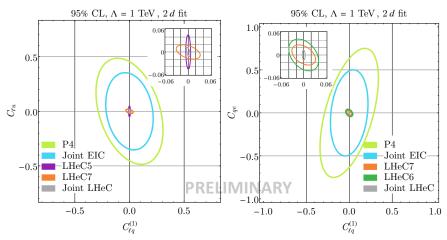






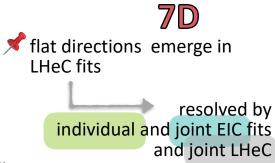
# **CONFIDENCE ELLIPSES**

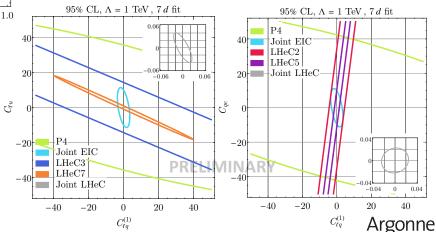
# Four-fermion $\mathcal{C}_k$



2D

- no flat directions here
- LHeC more constraining than EIC



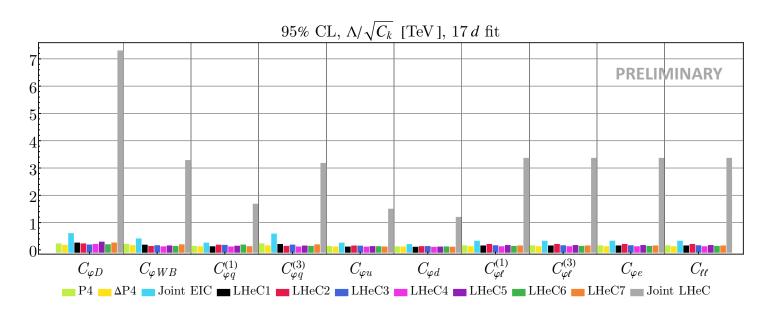


# **17D FIT**

#### **UV** effective scales



17-parameter fit on the Wilson coefficients that induce the semi-leptonic four-fermion contact interaction





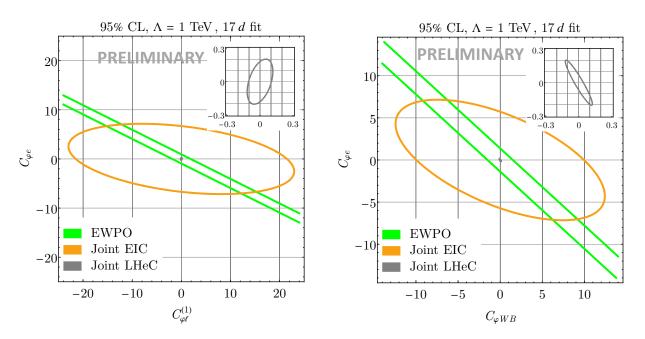
very similar situation to 7D fit, joint fits have the most constraining power





# **17D FIT**

#### **Confidence ellipses**

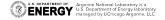


#### 34D EWPO fits

J. Ellis, M. Madigan, K. Mimasu V. Sanz, T. You JHEP04 (2021) 279

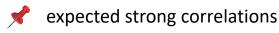


resolving blind spots observed in the fits of Wilson coefficients in the f f V sector using Z and W pole observables (EWPO) data



#### **PRELIMINARY**

### joint LHeC fit

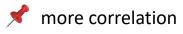


$$(\mathcal{C}_{\ell q}^{(1)}\,,\mathcal{C}_{\ell q}^{(3)})$$

unexpected strong correlations

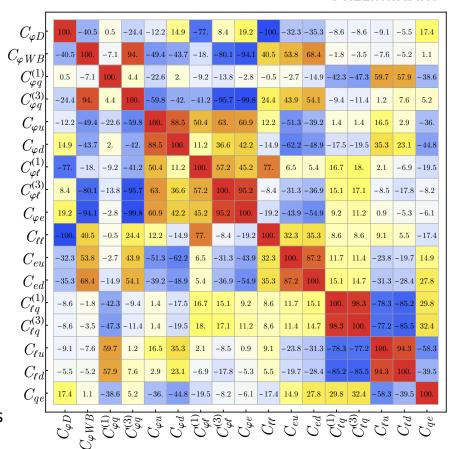
$$(\mathcal{C}_{\ell\ell}, \mathcal{C}_{\varphi D})$$
  $(\mathcal{C}_{\ell u}, \mathcal{C}_{\ell d})$ 

non trivial relation with experimental uncertainties



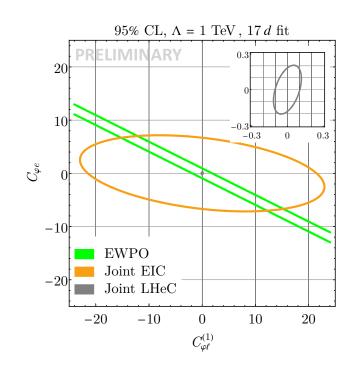


weaker bounds



# **CONCLUSIONS**

- ✓ We performed 1D, 2D, 7D and 17D fits
  of EIC and LHeC pseudo data
- LHeC can probe scales up to 7 TeV
- (Some of) the **blind spots** observed by Ellis et al. are **resolved** by ElC and LHeC



**EIC and LHeC have both great potential for BSM studies** 

# **OUTLOOK**



# perform impact studies



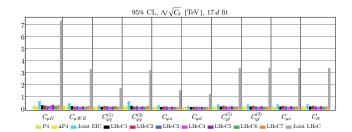






growing interest in the community

see Tim's talk









# **OBSERVABLES**

#### **Reduced Neutral Current (NC) DIS cross section**

$$\frac{\mathrm{d}^2 \sigma_{r, \text{NC}}^{\ell}}{\mathrm{d}x \, \mathrm{d}Q^2} = \left\{ \frac{2\pi\alpha^2}{xQ^4} [1 + (1-y)^2] \right\}^{-1} \frac{\mathrm{d}^2 \sigma_{\text{NC}}^{\ell}}{\mathrm{d}x \, \mathrm{d}Q^2}$$

$$\frac{d^2 \Delta \sigma_{r,NC}^{\ell}}{dx \, dQ^2} = \left\{ \frac{4\pi\alpha^2}{xQ^4} [1 + (1-y)^2] \right\}^{-1} \frac{d^2 \Delta \sigma_{NC}^{\ell}}{dx \, dQ^2}$$



Unpolarized hadron

Polarized hadron

#### Parity-violating (PV) DIS asymmetries

unpolarized

$$A_{\mathrm{PV}} = rac{\sigma_{\mathrm{NC}}^{+} - \sigma_{\mathrm{NC}}^{-}}{\sigma_{\mathrm{NC}}^{+} + \sigma_{\mathrm{NC}}^{-}}$$



polarized

$$\Delta A_{\rm PV} = \frac{\Delta \sigma_{\rm NC}^0}{\sigma_{\rm NC}^0}$$

# WILSON COEFFICIENTS BOUNDS

# Marginalized bound of $\mathcal{C}_k$

$$\Delta C_k = \sqrt{\Delta \chi^2(1,c)(F^{-1})_{kk}}$$
 quai

F is the Fisher matrix of the fitted parameters

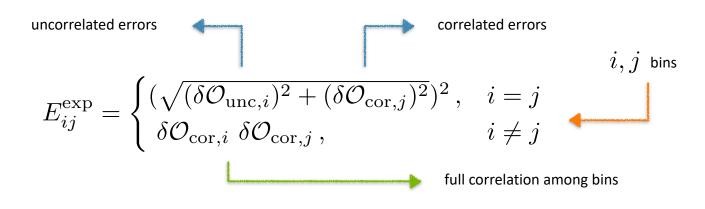
quantile of the  $\chi^2$  distribution for d=1,2 marginalized fitted parameters at confidence level c

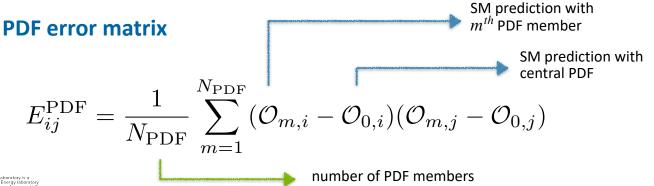
# Confidence ellipses

$$(C_k \quad C_{k'}) \begin{pmatrix} (F^{-1})_{kk} & (F^{-1})_{kk'} \\ (F^{-1})_{k'k} & (F^{-1})_{k'k'} \end{pmatrix}^{-1} \begin{pmatrix} C_k \\ C_{k'} \end{pmatrix} = \Delta \chi^2(2, c)$$

# **COVARIANCE MATRIX**

#### **Experimental error matrix**





# **COVARIANCE MATRIX**

#### **Total covariance matrix**

$$V_{ij} = E_{ij}^{\text{exp}} + E_{ij}^{\text{PDF}}$$

for joint fits, with more than one experiment

# joint covariance matrix

LHeC case

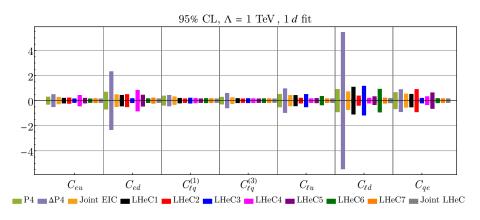
$$V = egin{pmatrix} V_1 & J_{12} & \cdots & J_{17} \ & V_2 & \cdots & J_{27} \ & & \ddots & dots \ & & V_7 \end{pmatrix} egin{pmatrix} J_{nn'} = J_{nn'}^{ ext{exp}} + J_{nn'}^{ ext{PDF}} \ & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & &$$

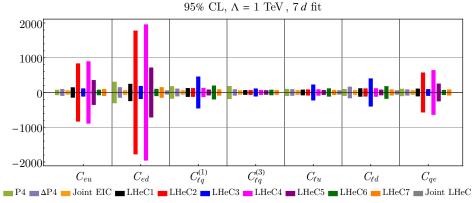
 $V_n$  is the covariance matrix of the  $n^{th}$  LHeC set



# **MARGINALIZED BOUNDS**

#### **Four-fermion Wilson coefficients**

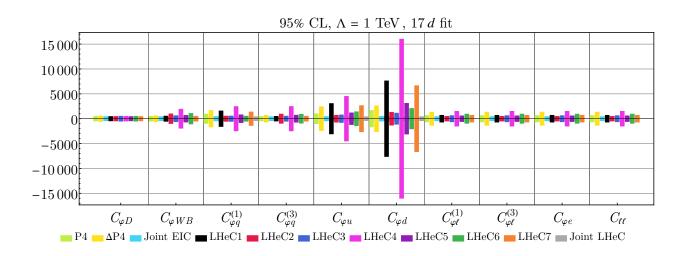






# **17D FIT**

# **Marginalized bounds**



# **17D FIT**

# **Marginalized bounds**

	Joint EIC	Joint LHeC	EW diboson, Higgs, and top data
C -			
$C_{\varphi D}$	[-3.8, 3.8]	[-0.019, 0.019]	[-1.6, 0.81]
$\frac{\Lambda}{\sqrt{C_{\varphi D}}}$	0.51	7.2	0.91
$C_{\varphi WB}$	[-9.9, 9.9]	[-0.098, 0.098]	[-0.36, 0.73]
$\frac{\Lambda}{\sqrt{C_{\varphi WB}}}$	0.32	3.2	1.4
$C_{\varphi q}^{(1)}$	[-38., 38.]	[-0.40, 0.40]	[-0.27, 0.18]
$\frac{\Lambda}{\sqrt{C_{\varphi q}^{(1)}}}$	0.16	1.6	2.1
$C_{\varphi q}^{(3)}$	[-4.1, 4.1]	[-0.11, 0.11]	[-0.11, 0.012]
$\frac{\Lambda}{\sqrt{C_{\varphi q}^{(3)}}}$	0.49	3.1	4.1
$C_{\varphi u}$	[-38., 38.]	[-0.51, 0.51]	[-0.63, 0.25]
$\frac{\Lambda}{\sqrt{C_{\varphi u}}}$	0.16	1.4	1.5
$C_{\varphi d}$	[-84., 84.]	[-0.82, 0.82]	[-0.91, 0.13]
$\frac{\Lambda}{\sqrt{C_{\varphi d}}}$	0.11	1.1	1.4
$C_{\varphi\ell}^{(1)}$	[-18., 18.]	[-0.094, 0.094]	[-0.19, 0.41]
$\frac{\Lambda}{\sqrt{C_{\varphi\ell}^{(1)}}}$	0.23	3.3	1.8
$C_{\varphi\ell}^{(3)}$	[-4.1, 4.1]	[-0.060, 0.060]	[-0.13, 0.055]
$\frac{\Lambda}{\sqrt{C_{\varphi\ell}^{(3)}}}$	0.49	4.1	3.3
$C_{\varphi e}$	[-5.7, 5.7]	[-0.16, 0.16]	[-0.41, 0.79]
$\frac{\Lambda}{\sqrt{C_{\varphi e}}}$	0.42	2.5	1.3
$C_{\ell\ell}$	[-7.7, 7.7]	[-0.039, 0.039]	[-0.084, 0.02]
$\frac{\Lambda}{\sqrt{C_{tt}}}$	0.36	5.1	4.4