Recent Light Hypernuclei Measurements from STAR Experiment

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- Introduction \bigcirc
- Review of previous hypernuclei measurements from STAR
- Recent results of hypernuclei from STAR \bigcirc
	- Internal structure
		- Branching ratios, lifetimes, $Λ$ binding energies
	- Production mechanism in heavy-ion collisions
		- Yields, collectivity
	- Discovery of 4 $\frac{4}{\Lambda}H$
- Summary \bigcirc
- **Outlook** \bigcirc

Outline

Introduction: what and why

• What are hypernuclei?

- Why hypernuclei?
	- Probe hyperon-nucleon (Y-N) interaction
	- Strangeness in high density nuclear matter EoS of neutron stars.
		- Equation-of-State (EoS) of neutron star

^ΛHe

Marian Danysz (right) and Jerzy Pniewski (left) **Marian Danysz (right) and Jerzy Pniewski (left)** discovered hypernuclei in 1952 **discovered hypernuclei in 1952** Fridth Dattyse (right) and octey I mewsix (icty

Xiujun Li/ USTC&UT relation (panel b) satisfying the 2*M*! constraint. The horizontal lines and bands in panel (b) show the observational data of

Introduction: how

- Experimentally, we can make measurements related to:
- 1. Internal structure Internal structure of hypernuclei
	- Lifetime, binding energy, branching ratios etc. Understanding hypernuclei structure can provide insights to the Y-N interaction Loosely bounded, binding energy, branching ratios ... Uniong hypernuclei structure suri provide morghte to the T-N interaction
- 2. Production mechanism
	- Spectra, collectivity etc. production yields/mechanisms, collectivity …

The process of hypernuclei formation in violent heavy-ion collisions is not well *understood* \mathcal{F}

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• During the BES-II program, STAR utilized the fixed-target (FXT) setup, which

Introduction: RHIC BES program

extends the energy reach below $\sqrt{s_{NN}}$ = 7.7 GeV, down to 3.0 GeV

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Hypernuclei and STAR BES-**List of BES-II datasets:**

• Hypernuclei measurements are scarce in heavy-ion collision experiments

A. Andronic et al. PLB (2011) 697:203–207

- At low beam energies, hypernuclei production is expected to be enhanced due to high baryon density
	- Datasets with large statistics taken during BES-II
	- → A great opportunity to study hypernuclei production

Previous hypernuclei measurements from STARSTAR

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Measurement of mass difference and binding energies of ${}_{\wedge}^{3}$ H and **Nature Phys. 16 (2020) 409 (STAR)** $\frac{3}{\Lambda}$ H and $\frac{3}{\Lambda}$ H

STAR collaboration **made the discovery of the anti-hyper triton. Science 328, 58 (2010) (STAR)**

Lifetime measurement of **Science 328, 58 (2010) (STAR) PRC 97, 054909 (2018) (STAR)** 3 $\rm \frac{3}{\Lambda}H$

Particle in the internal fication and hypernuclei reconstruction

- Particle identification from energy loss measurement using TPC
- KF particle package $^{[1]}$ is used for signal reconstruction
- Hypernuclei reconstructed via their weak decay channels: ${}_{\Lambda}^{3}H\rightarrow{}^{3}He + \pi^{-}$ ${}_{\Lambda}^{3}H\rightarrow d + p + \pi^{-}$ ${}_{\Lambda}^{4}H\rightarrow{}^{4}He + \pi^{-}$ ${}_{\Lambda}^{4}$

$$
\pi^- \quad {}^{4}_{\Lambda}\text{He} \rightarrow {}^{3}\text{He} + \text{p} + \pi^-
$$

[1]Zyzak M, Kisel I, Senger P. Online selection of short-lived particles on many-core computer architectures in the CBM experiment at FAIR[R]. Collaboration FAIR: CBM, 2016.

Vertex

^p ^¼ ⁷.² GeV are not presented here due

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	-

Hypernuclei signal reconstruction

$\frac{3}{4}$ H branching ratio $\frac{3}{4}H$ branching ratio R_3

Relative branching ratio: $R_3 =$

• $B_\Lambda \to$ provide constraints to Y-N interaction

- Using $\sqrt{s_{NN}}$ = 3.0 GeV data:
	- $R_3 = 0.272 \pm 0.030(stat.) \pm 0.042(syst.)$
	- Model comparison suggesting a weakly-bounded state for $\frac{3}{\Lambda}$ H

• Stronger constraints on absolute B.R.s and $^{3}_{\Lambda}H$ internal structure models 3 $^{\text{5}}_{\Lambda} \text{H}$

PRC 102,

064002

 (2020) B_A (MeV)

 0.2

 0.1

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[1]A. Gal and H. Garcilazo, PLB 791, 48 (2019) [2]J.G. Congleton, J. Phys. G 18, 339 (1992)

$\frac{3}{\Lambda}$ H, $\frac{4}{\Lambda}$ H and $\frac{4}{\Lambda}$ He lifetimes

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$^{4}_{\Lambda}$ H

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$^{4}_{\Lambda}$ He

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Using $\sqrt{s_{NN}}$ = 3.0 GeV and 7.2 GeV datasets:

- ${}_{\Lambda}^{3}$ H: $\tau = 221 \pm 15$ (stat.) ± 19 (syst.)[ps]
- $^{4}_{\Lambda}$ H: $\tau = 218 \pm 6$ (stat.) ± 13 (syst.)[ps]
- $^{4}_{\Lambda}$ He: $\tau = 229 \pm 23$ (stat.) ± 20 (syst.)[ps]
- Indication of shorter lifetimes for ${}_{\Lambda}^{3}H$, ${}_{\Lambda}^{4}H$ and ${}_{\Lambda}^{4}He$ than that of free Λ (with 1.8 σ , 3.0 σ , 1.1 σ respectively)
- Consistent with former measurements and world average values
- $\tau_{3,H}$: consistent with calculation including pion FSI^{LI} and calculation with Λ d 2-body picture $^{[2]}$ within 1 σ $^{3}_{\Lambda}\text{H}$ [1]
- $\tau_{A\dot{H}}$ and $\tau_{A\dot{H}e}$: consistent with expectations from isospin rule

Precision ${}_{\lambda}^{3}H$ and ${}_{\lambda}^{4}H$ measurements provide tight constraints on models. ${}^{3}_{\Lambda}\rm{H}$ and ${}^{4}_{\Lambda}\rm{H}$

$\rm\,B_{\Lambda}$ and $\rm\Delta B_{\Lambda}$ of $^{4}_{\Lambda}\rm H$ and B_Λ and ΔB_Λ of $^4_\Lambda$ H and $^4_\Lambda$ He EN ANA AR BDT A CUIU AND A \blacksquare

STAR, PLB 834, 137449 (2022) *STAR Collaboration Physics Letters B 834 (2022) 137449*

• A binding energy $B_\Lambda = (M_\Lambda + M_{\text{core}} - M_{\text{hypernucleus}})c^2$ • \blacksquare dashed vertical lines are drawn at "B4 • In binding energy $D_A - (M_A + M_{\text{core}} - M_{\text{hypernucleus}})$

• The ground state \mathbf{B}_{Λ} are directly measured: $\Delta B_{\Lambda}^{4}(0^{+}) = B_{\Lambda}({}^{4}_{\Lambda}\text{He},0^{+}) - B_{\Lambda}({}^{4}_{\Lambda}\text{H},0^{+})$ ground-state binding energies are from this analysis. The values for excited states for e • The ground state \mathbf{B}_{Λ} are directly measured

!H and

$$
\Delta B_{\Lambda}^-(0^+) = B_{\Lambda}(\Lambda H, 0^+) - B_{\Lambda}(\Lambda H, 0^+)
$$
\n• For excited states, the results are obtained by combining with the γ -ray transition energies E_{γ} ^{J-PARC E13, PRL 115, 222501}
\n
$$
D_{\Lambda}^A(A \times A \times A) = D_{\Lambda}^A(A \times A) = D_{\Lambda}^A(A
$$

$$
B_{\Lambda}^{4}(^{4}_{\Lambda}\text{He/H},1^{+}) = B_{\Lambda}^{4}(\Lambda\text{He/H},0^{+}) - E_{\gamma}^{4}(\Lambda\text{He/H})
$$
 consist
\n
$$
\Delta B_{\Lambda}^{4}(1^{+}) = B_{\Lambda}^{4}(\Lambda\text{He},1^{+}) - B_{\Lambda}^{4}(\Lambda\text{H},1^{+})
$$
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And the 1900 km in the 1900 km i Eq. (1). The mass of \mathbb{R}^n and \mathbb{R}^n and \mathbb{R}^n is taken from \mathbb{R}^n the mass of \mathbb{R}^n is taken from \mathbb{R}^n the mass of \mathbb{R}^n and \mathbb{R}^n is taken from \mathbb{R}^n the mass of \mathbb{R}^n a

• Mirror hypernuclei ${}^{4}_{\Lambda}\text{H}$ and ${}^{4}_{\Lambda}\text{He}$: opportunity to study charge symmetry breaking (CSB) effect in $A = 4$ hypernuclei and a solid solid black line solid solid solid solid solid solid solid solid solid error bars show statistical u \mathbf{b} and \mathbf{b} in except states was calculated to be \mathbf{b} $\mathsf{I} \cap \mathsf{A} = 4$ in woernuclei $\mathsf{I} \cap \mathsf{A} = 2$ \mathcal{A} discussed in the introduction and shown as solid black circle black c

> (2015) CERN-Lyon-Warsaw, PLB. 62, 467 (1976)

- d by combining \bullet CSB in 0^+ and 1^+ states are comparable and have opposite signs Lations predictions in both ground i states and excited states and excited states \blacksquare (denoted as \blacksquare Lyon-Warsaw, PLB. 62, 467 (1976) \blacksquare \blacksquare \blacksquare \blacksquare \blacksquare \blacksquare \blacksquare \blacksquare $\frac{1}{2}$. **1.** $\frac{1}{2}$ $\sinh(\theta)$ signals and the mirror of the JOSILE SIGNS
- Consistent with theoretical calculations within large uncertainties \sim prediction matches our measurements. This matches our measurements. This may indicate that \sim [√]*s*NN ⁼ 3 GeV. By using the ^γ -ray transition energies of the ext with theoretical calculations \blacksquare energies of the measurement states are also extracted. The CSB effects are als \bullet binds between the ground states between the ground states of 4

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Coalescence parameters (rc, pc) are tuned to fit the data, see backup* **Transport model (JAM) with coalescence afterburner* qualitatively ately reproduces trends of $^*_\Lambda \mathrm{H}$ rapidity distributions seen in data, but

Hypernuclei production at 3 GeV 4 **ТУ**

- Different trends in the $\overline{\Lambda}$ H rapidity disti-
	- Transport model (JAM) with coalescence approximately reproduces trends of $^{4}_{\Lambda}H$ rapidity distributions seen in data, but fails to reproduce the trend of $^3_\Lambda \text{H}$ in 10-50% • Transport model (JAM) with coalescence approximately reproduces trends of $^{4}_{\Lambda}H$ rapidity distributions

• Different trends in the $^4_\Lambda \rm H$ rapidity distribution in central (0-10%) and mid-central (10-50%) collisions at $\sqrt{s_{NN}}$ = 3.0 GeV • Different trends in the $^4_\Lambda \rm H$ rapidity distribution in central (0-10%) and mid-central (10-50%) collisions at $\sqrt{s_{NN}}$ = 3.0 GeV (10-10) and mid-central (10-50%) collisions in the central (10-50%) collisions in the central (10-50%) collisi
The central (10-50%) collisions in the central (10-50%) collisions in the central (10-50%) collisions in the c

***Uncertainty in R3 (19%) not shown*

- First observation of ${}^{3}_{\Lambda}H$ and ${}^{4}_{\Lambda}H$ directed flow (v₁) in mid-central 5-40% Au+Au collisions at 3 GeV
- Mid-rapidity v_1 slopes of ${}^3_\Lambda H$ and ${}^4_\Lambda H$ follow baryon mass scaling.

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[arXiv:2211.16981](https://arxiv.org/abs/2211.16981) accepted by PRL

 \rightarrow Imply coalescence process to be the dominant formation mechanism for $^3_\Lambda \text{H}$ and $^4_\Lambda \text{H}$ **production in 3 GeV Au+Au collisions**

Energy dependence of hypernuclei production in heavy-ion collisions

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- T. Reichert, J. Steinheimer et al, arXiv:2210.11876(2022) (UrQMD, Thermal-FIST)

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- Suppression of ${}_{\Lambda}^{3}H/{}^{3}He$ yield ratios compared to that of Λ/p
	- Observed at both 0-10% and 10-40% centrality in Au+Au collisions at 3 GeV.
- The $^{4}_{\Lambda}\text{H}/^{4}\text{He}$ yield ratios are comparable to that of Λ/p
- **• UrQMD model with coalescence describes the tendency of the distributions reasonably well, suggesting coalescence mechanism for hypernuclei formation. Suggest coalescence mechanism and creation of excited A = 4 hypernuclei**

• Non-monotonic behavior in light-to-hyper-nuclei ratio vs A

- Thermal model calculations including excited $^{4}_{\Lambda}\text{H}^{*}$ feed-down **show a similar trend** exponential dependence of yields/(2J+1) vs A \mathbf{r} is a factor of \mathbf{r} above exponential fit to \mathbf{r} $\frac{1}{4}$ II^k food decess. Λ ll iccu-down ι-4ς
1 $\frac{1}{2}$ IUW d SI
	- Feed-down from excited state enhances ${}^{4}_{\Lambda}H$ production

 \to Coalescence parameters $\rm B_A$ of $^{\rm A}_{\Lambda}H$ and $^{\rm A}$ He follow similar tendency versus $\rm p_T$, rapidity and centrality, indicating that N-N and Y-N interactions that drive coalescence dynamics in these collisions are similar

S.Zhang, PLB 684(2010)224

- B_A : Coalescence parameters

 S_3 < 1: relative suppression of ${}^{3}_{\Lambda}H/{}^{3}He$ compared to Λ/p S_4 ~1, $S_4 > S_3$: ${}^{4}_{\Lambda}H/{}^{4}He$ is comparable to Λ/p ■ Expect ~1 if no suppression

No obvious kinematic and centrality dependence of $\mathrm{S}_{3,4}$ observed at 3 GeV.

 10^{-1}

at 3 GeV

- Strangeness population factor $\mathbf{S}_{\mathbf{A}}$
	- **Relative suppression of hypernuclei** production compared to light nuclei production

$$
S_A = \frac{AH}{AHe \times \frac{\Lambda}{p}} = \frac{B_A(AH)(p_T)}{B_A(AHe)(p_T)} \text{ of}
$$

STAR, Science 328 (2010) 58 ALICE, PLB 754 (2016) 360 E864, PRC 70 (2004) 024902 NA49, J.Phys.Conf.Ser.110(2008)032010

A. Andronic et al, PLB 697 (2011) 203 (Thermal (GSI)) S. Zhang, PLB 684(2010)224 (Coal.+AMPT) T. Reichert, J. Steinheimer et al, arXiv:2210.11876(2022) (UrQMD, Thermal-FIST)

STAR preliminary \sqrt{m} in Thermal Model (GSI) **•** Data show a hint of an increasing trend from $\overline{s_{NN}}$ = 3.0 GeV to 2.76 TeV

Energy dependence of S_3

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• For coalescence models, the energy dependence is sensitive to the source radius (Δr)

• Thermal-FIST, which includes feed-down to p and ³He from unstable nuclei, describes the data reasonably well 3 $\rm He$ from unstable nuclei, describes the $\rm S_3$

Provide constraints for hypernuclei production Xiujun Li@ATHIC2023 **Water of Service CEVI models in the high-baryon-density region**

See Tan Lu's poster@QM2022

First observation of $\frac{4}{4}$ **First observation of**

- The **heaviest antimatter hypernucleus** observed up to now
	- New opportunity for the study of matterantimatter asymmetry

Summary

Enhanced hypernuclei yields at low energies allow precision measurement.

- Precision ${}^3_\Lambda \text{H}$, ${}^4_\Lambda \text{H}$ lifetimes and relative branching ratio R_3 of ${}^3_\Lambda \text{H}$ measured $\;\rightarrow$ strong constraints on hypernuclei models
-

Internal structure

First discovery of 4 $\frac{4}{\Lambda}H$

• A binding-energy difference between ${}^{4}_{\Lambda}\text{H}$ and ${}^{4}_{\Lambda}\text{He}$ shows hint of CSB effect for A=4 hypernuclei

• Relative suppression of $^{3}_{\Lambda}H/^{3}He$ compared to Λ/p and $^{4}_{\Lambda}H/^{4}He \to$ support creation of $^{4}_{\Lambda}H^{*}$,

Production

- v_1 slopes of Λ , $^3_\Lambda \text{H}$ and $^4_\Lambda \text{H}$ follow baryon mass scaling $\;\rightarrow$ support coalescence picture • First ${}^3_\Lambda \text{H}$ and ${}^4_\Lambda \text{H}$ dN/dy vs y and energy dependence of dN/dy @ high $\mu_{\text{B}} \to$ constraints to
- hypernuclei production models
- coalescence picture; weak centrality/kinematic dependence for S_3 and S_4

Outlook

Coalesc. (JAM)

PHQMD

14.6

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-

• e.g.
$$
{}_{\Lambda\Lambda}^4
$$
He $\rightarrow {}_{\Lambda}^4$ He π , ${}_{\Lambda\Lambda}^5$ He $\rightarrow {}_{\Lambda}^5$ He π

Outlook

-
-
-

• e.g.
$$
{}_{\Lambda\Lambda}^4
$$
He $\rightarrow {}_{\Lambda}^4$ He π , ${}_{\Lambda\Lambda}^5$ He $\rightarrow {}_{\Lambda}^5$ He π

Backups

• UrQMD cascade + coalescence in slide 14:

• Coalescence takes place if the spatial coordinates and relative momenta

• UrQMD+ coalescence in slide 16:

- Assuming two parameter sets (a) and (b) for ${}^{3}_{\Lambda}H$.
	- (a) $\Delta r = 9.5$ fm, similar to ${}_{\Lambda}^{3}H$ size.
	- (b) $\Delta r = 4.3$ fm, similar to triton size.
- $\sqrt{s_{NN}} \leq 20$ GeV, UrQMD cascade + coalescence; $\sqrt{s_{NN}}\geq$ 20 GeV, UrQMD hybrid + coalescence; Δp djusted to match each other at 20 GeV.

Model parameters

- of constituents are within a sphere of radius (Δr , Δp)
	- JAM + coalescence:

 ${}^{3}_{\Lambda}$ H 3-body signal ⁵
ΔH

If rotate proton or $pi-(d+p)$, it can not well describe the data mass.

3

 3.005 3.01 3.015 3.02

Mass($p\pi d$) (GeV/ c^2)

$\frac{3}{4}H \Lambda$ -binding energy²⁸⁰ $\frac{3}{\Lambda}H\Lambda$

- Recent results:
	- STAR 2020(Nature Phys. 16 (2020) 409):
		- $B_A = 410 \pm 120$ (stat.) ± 110 (syst.) keV
		- World average: 181 ± 48 keV
	- ALICE 2021(arXiv:2209.07360, 2022):
		- $B_A = 72 \pm 63(stat.) \pm 36(syst.)$ keV
		- World average(PLB 837 (2023)137639): 148 ± 40 keV

 $\bigcap_{n=1}^{\infty} 2^{n} 9^n 6^n + 0^2 8(\text{stat }) \pm 0 11(\text{syst }) \text{ MeV } c^{-2}$

NATURE PHYSICS | VOL 16 | APRIL 2020 | 409–412 | www.nature.com/naturephysics **411**Chaudhari-69 0 Juric-73 ALICE-21 Keyes-70 Prakash-61 -200

Lifetime

- Lifetime τ extracted via $N(t) = N_0 e^{-L/\beta \gamma c \tau}$
- Λ lifetime cross check : 267 \pm 4 ps, consistent with PDG value (263 \pm 2 ps)
- ${}^{3}_{\Lambda}$ H and ${}^{4}_{\Lambda}$ H lifetimes from 3.0 GeV consistent with 7.2 GeV results

