

# Combined Traps

Francis Robicheaux  
Auburn University

Single Particle Properties  
Plasma Properties

Support from DOE and NSF

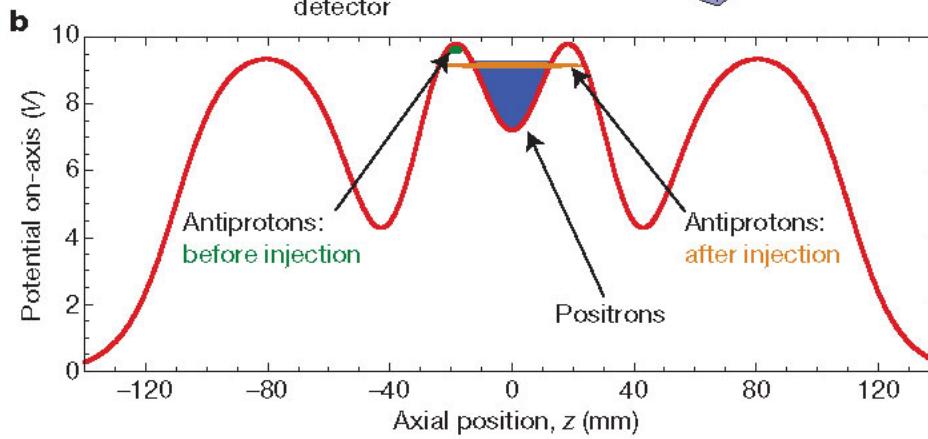
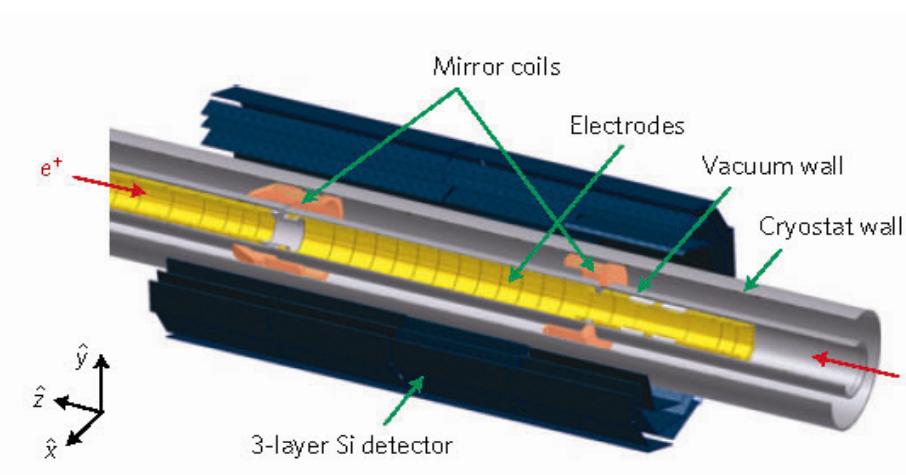
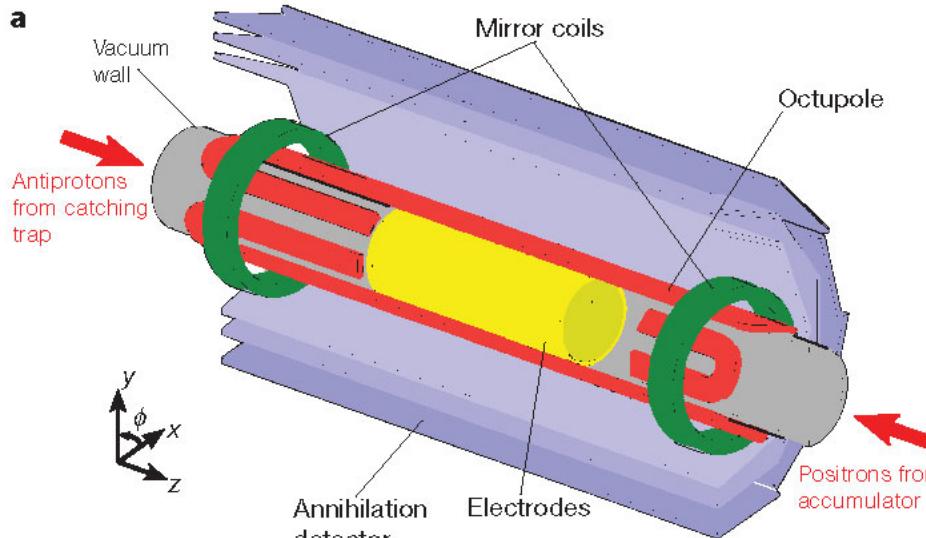
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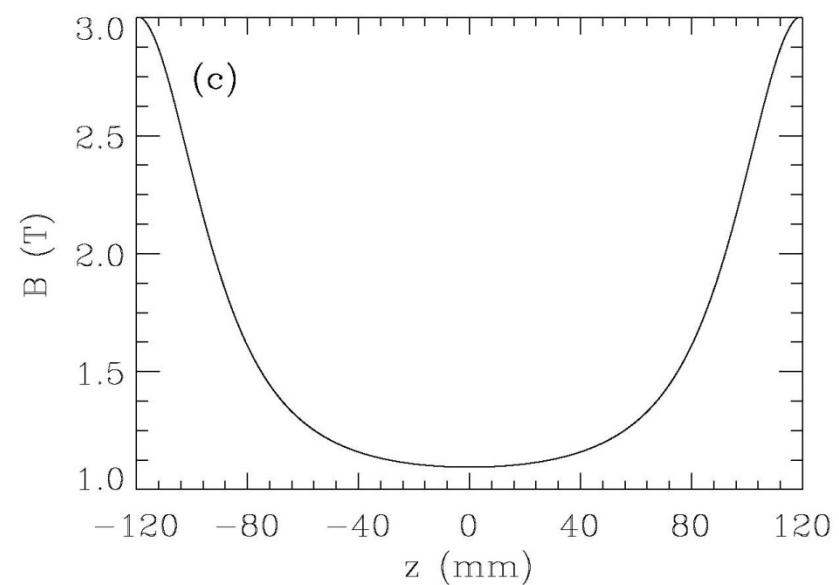
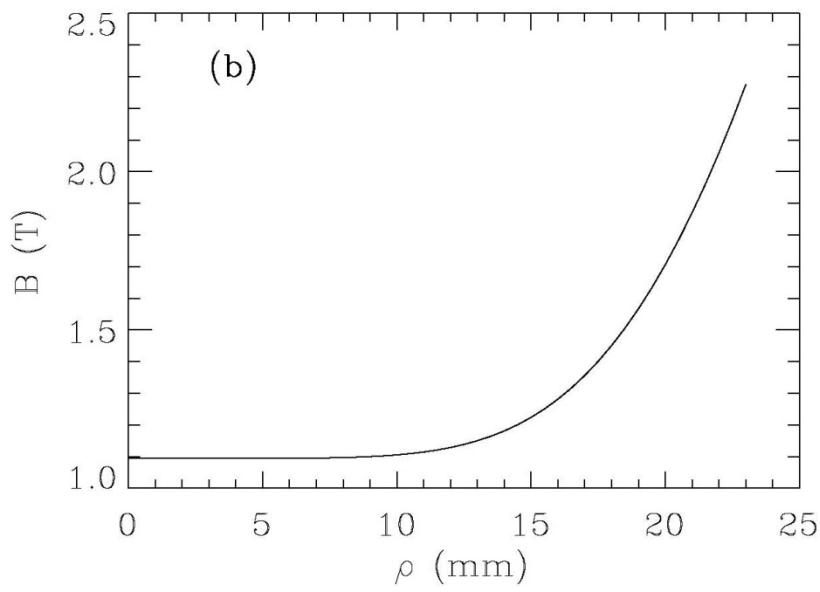
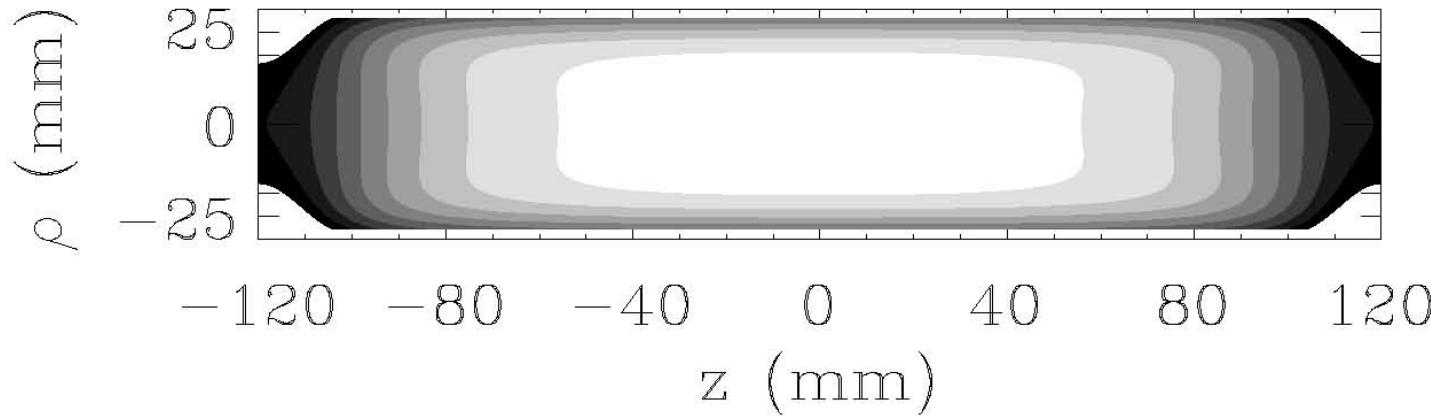
# Generic Combined Trap



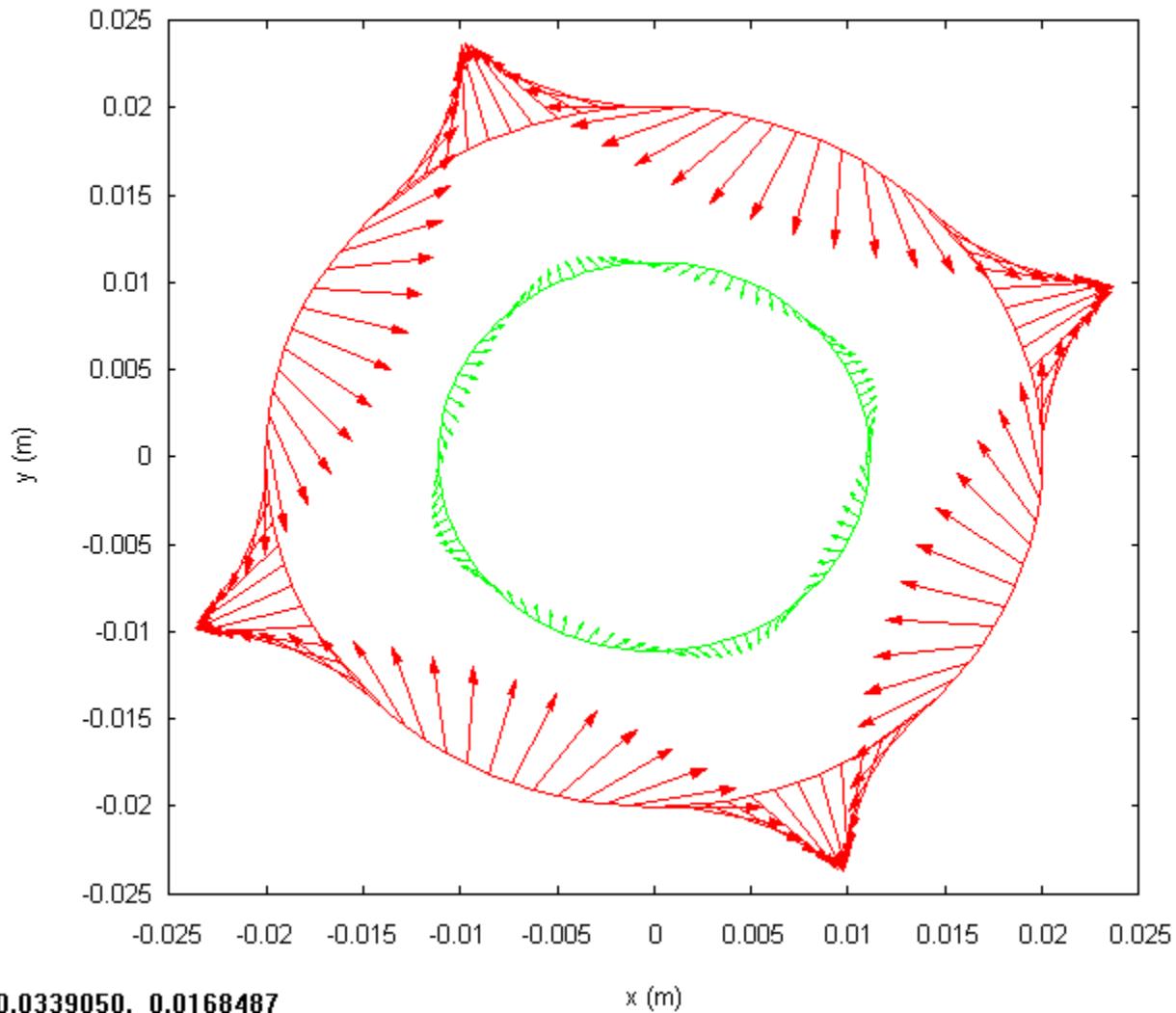
Trapping of charged particles

Axial electric potential well + axial B-field

# Trap For Neutral Atom



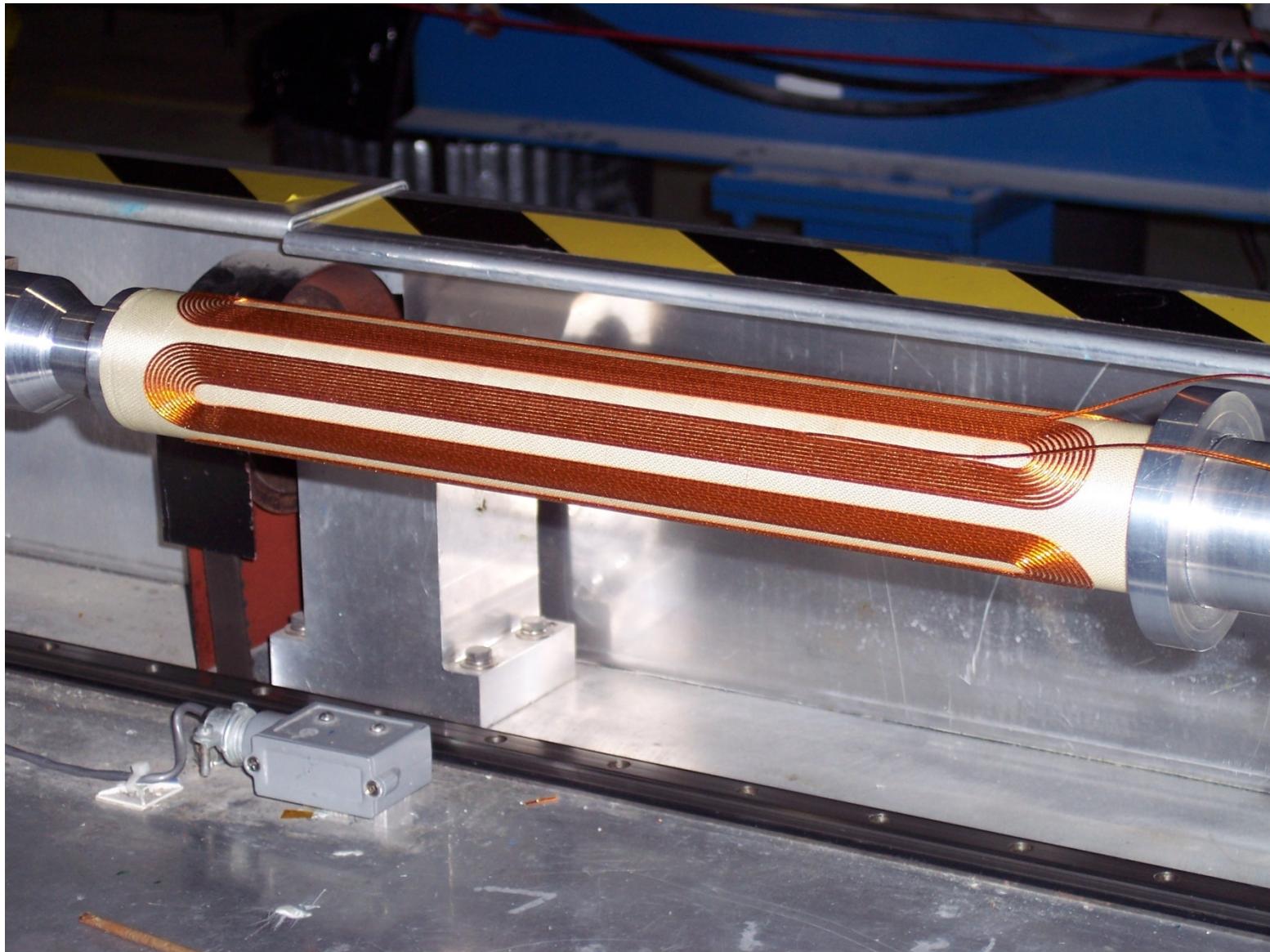
# Octupole B-field



Decreases  
very rapidly  
as  $r$  goes to 0

red =  $0.9 r_w$   
green =  $0.5 r_w$

# Actual Magnets: Octupole

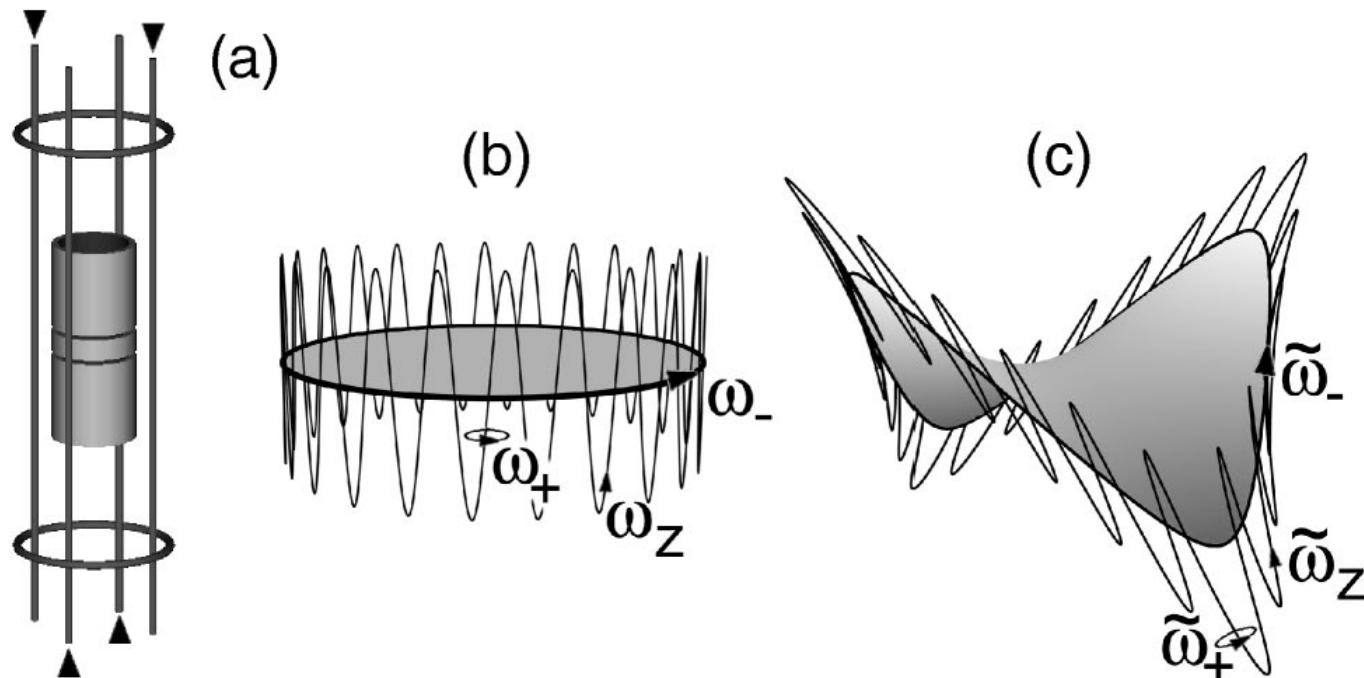


## Stability of a Charged Particle in a Combined Penning-Ioffe Trap

T. M. Squires, P. Yesley, and G. Gabrielse

*Department of Physics, Harvard University, Cambridge, Massachusetts 02138*

(Received 19 January 2001)



Simplified quadrupole  
case

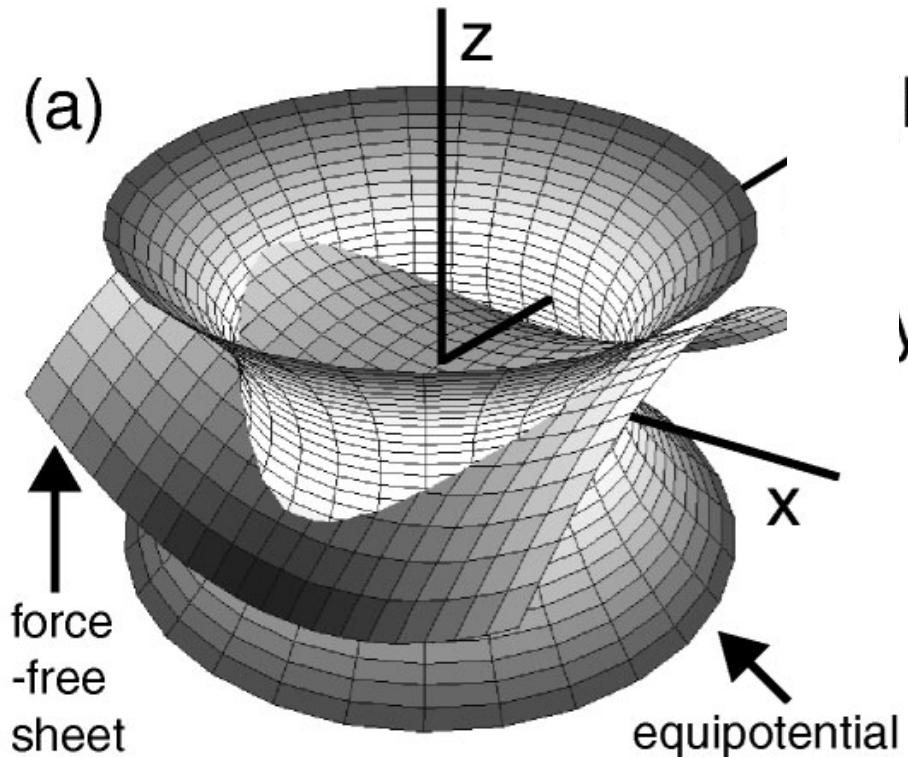
# Restoring region

$$PE = \frac{1}{2} m \omega_z (z^2 - r^2/2)$$

$$\mathbf{B} = (x/R_0, -y/R_0, 1) B_0$$

Edge of restoring region  
defined by  $\mathbf{F} \cdot \mathbf{B} = 0$

$$z = \frac{1}{2} (x^2 - y^2)/R_0$$



# Adiabatic Invariants

Fastest motion is cyclotron motion: not in xy plane because  $B$  is tilted

Next fastest motion is axial motion (not purely  $z$ )

Last is magnetron motion

Avoid resonances then trapped for exponentially long time

# Critical loss radius in a Penning trap subject to multipole fields

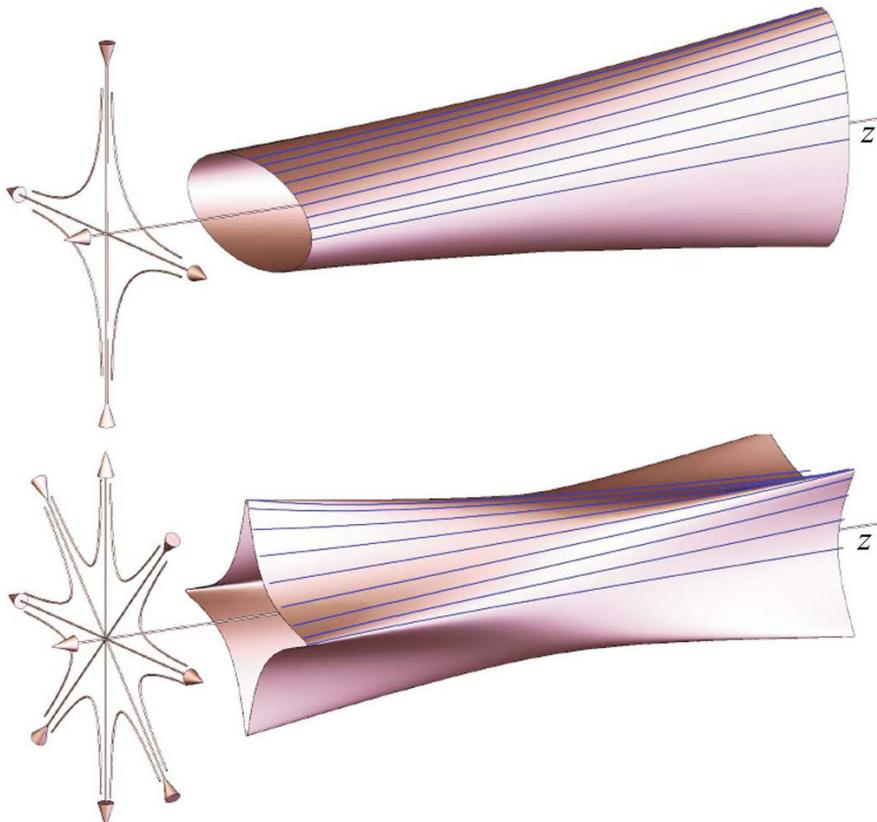
J. Fajans,<sup>1,a)</sup> N. Madsen,<sup>2</sup> and F. Robicheaux<sup>3</sup>

<sup>1</sup>*Department of Physics, University of California at Berkeley and the Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

<sup>2</sup>*Department of Physics, Swansea University, Swansea SA2 8PP, United Kingdom*

<sup>3</sup>*Department of Physics, Auburn University, Auburn, Alabama 36849-5311, USA*

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Initial circle following  
field lines to left and right  
for quadrupole and  
octupole

Single particle calculations

Distance  $z \gg r$

# Radius vs. Initial Angle

Define scaled length

$$L = B_w z / (B_z R_w)$$

$B_w$  = B at wall

$B_z$  = axial B

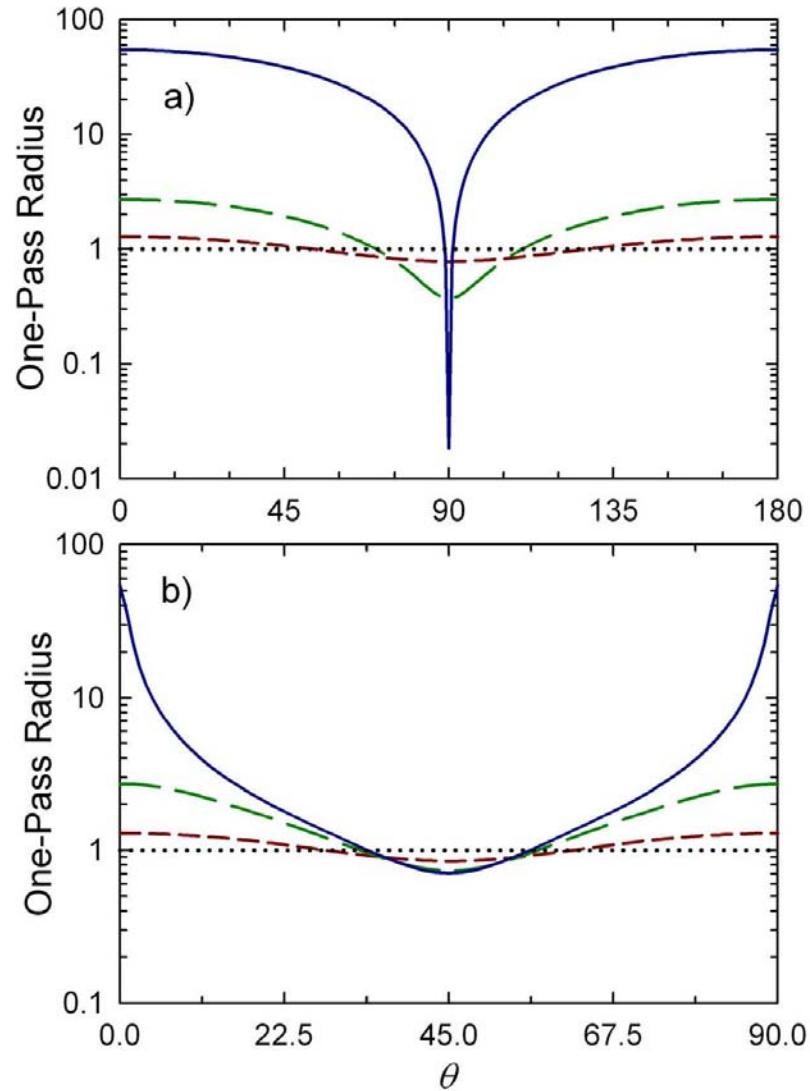
$R_w$  = radius of wall

**Quadrupole**

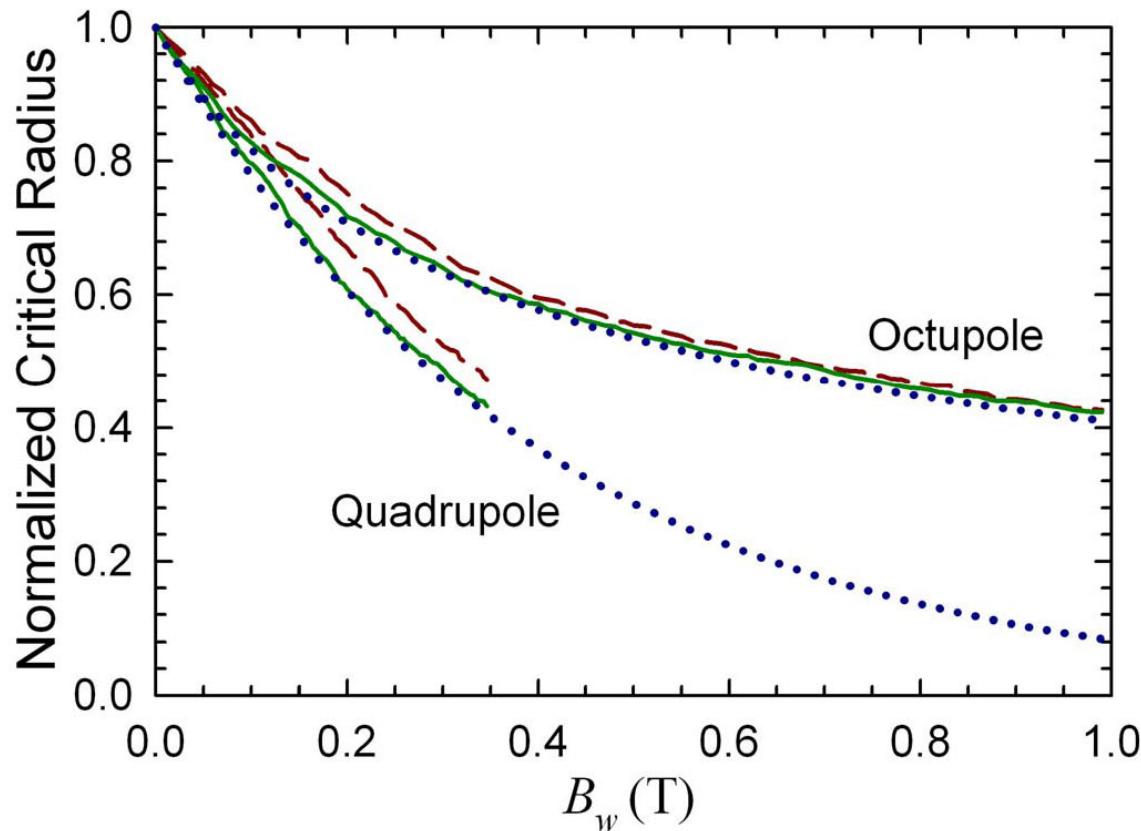
$$r = r_0 \exp(L/2)$$

**Octupole**

$$r = r_0 / \sqrt{1 - 2 L (r_0 / R_w)^2}$$

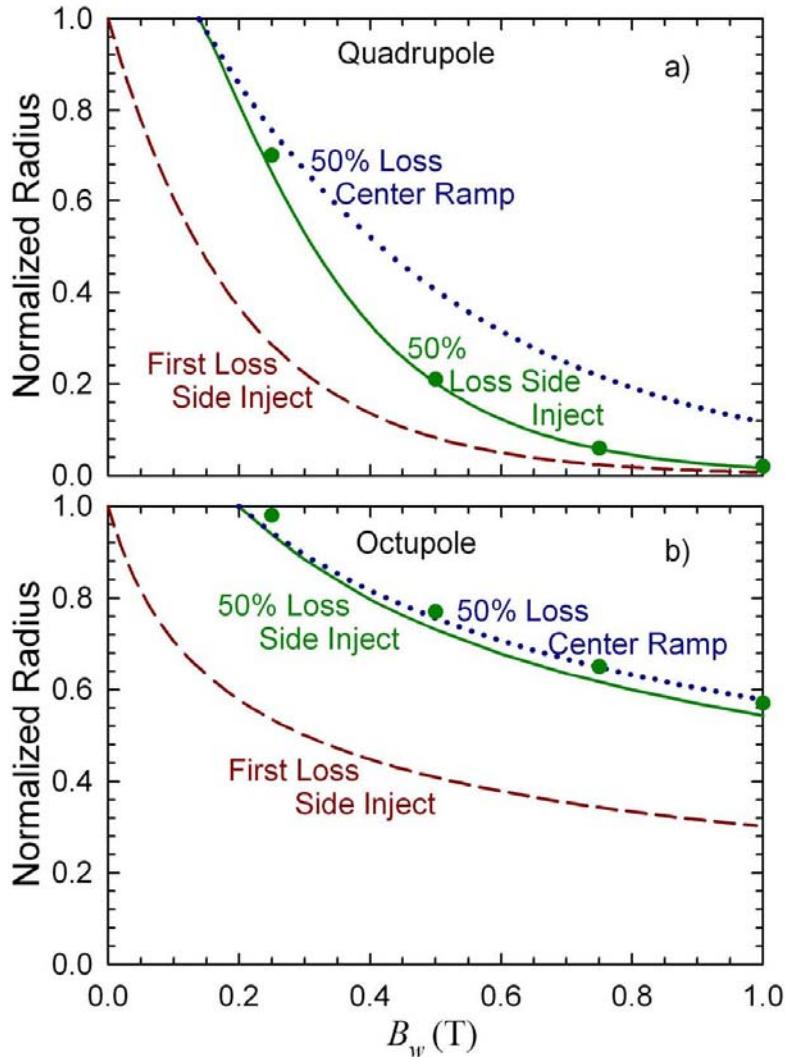


# Critical Radius vs. $B_w$



Dotted blue from inverting previous equation  
Green line (numerical)  
Red dashed (numerical no rotation)

# Comparison of Loss vs. $B_w$

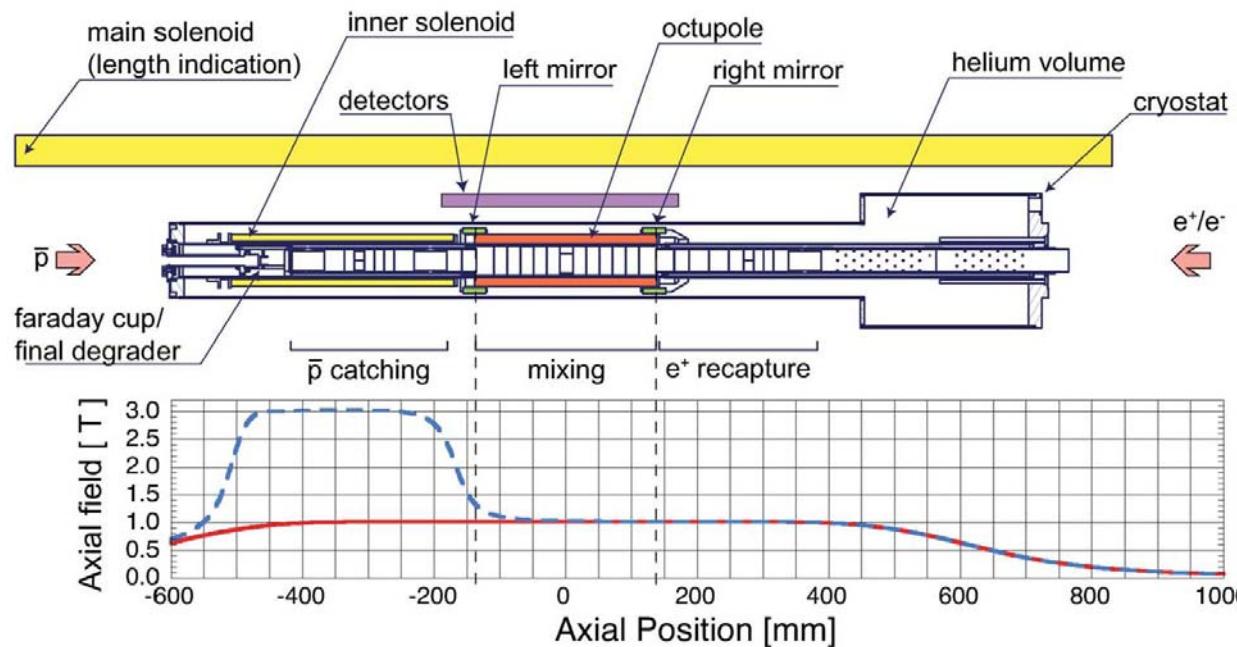


No collisions

What fraction of  
trajectories hit walls  
for quadrupole and  
octupole fields

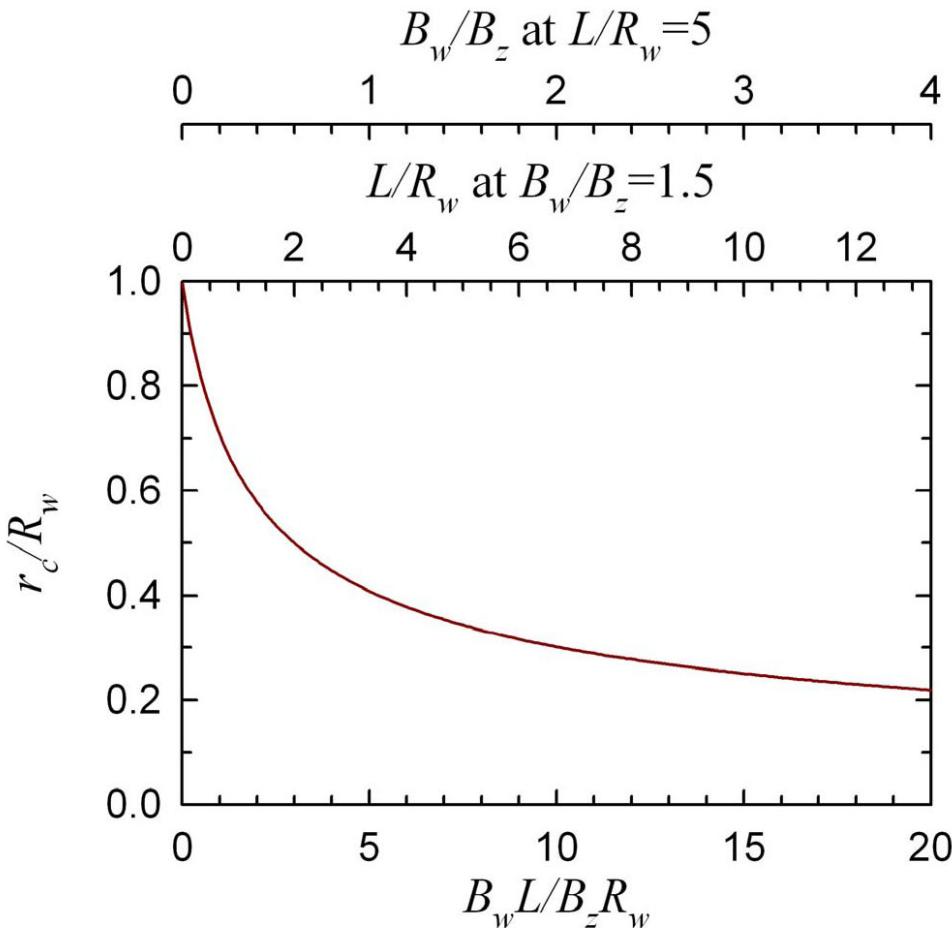
# A novel antiproton radial diagnostic based on octupole induced ballistic loss

G. B. Andresen,<sup>1</sup> W. Bertsche,<sup>2</sup> P. D. Bowe,<sup>1</sup> C. C. Bray,<sup>3</sup> E. Butler,<sup>2</sup> C. L. Cesar,<sup>4</sup> S. Chapman,<sup>3</sup> M. Charlton,<sup>2</sup> J. Fajans,<sup>3</sup> M. C. Fujiwara,<sup>5</sup> R. Funakoshi,<sup>6</sup> D. R. Gill,<sup>5</sup> J. S. Hangst,<sup>1</sup> W. N. Hardy,<sup>7</sup> R. S. Hayano,<sup>6</sup> M. E. Hayden,<sup>8</sup> A. J. Humphries,<sup>2</sup> R. Hydomako,<sup>9</sup> M. J. Jenkins,<sup>2</sup> L. V. Jørgensen,<sup>2</sup> L. Kurchaninov,<sup>5</sup> R. Lambo,<sup>4</sup> N. Madsen,<sup>2</sup> P. Nolan,<sup>10</sup> K. Olchanski,<sup>5</sup> A. Olin,<sup>5</sup> R. D. Page,<sup>10</sup> A. Povilus,<sup>3</sup> P. Pusa,<sup>10</sup> F. Robicheaux,<sup>11</sup> E. Sarid,<sup>12</sup> S. Seif El Nasr,<sup>7</sup> D. M. Silveira,<sup>4</sup> J. W. Storey,<sup>5</sup> R. I. Thompson,<sup>9</sup> D. P. van der Werf,<sup>2</sup> J. S. Wurtele,<sup>3</sup> and Y. Yamazaki<sup>13</sup>



Increase octupole strength to shave off antiprotons!

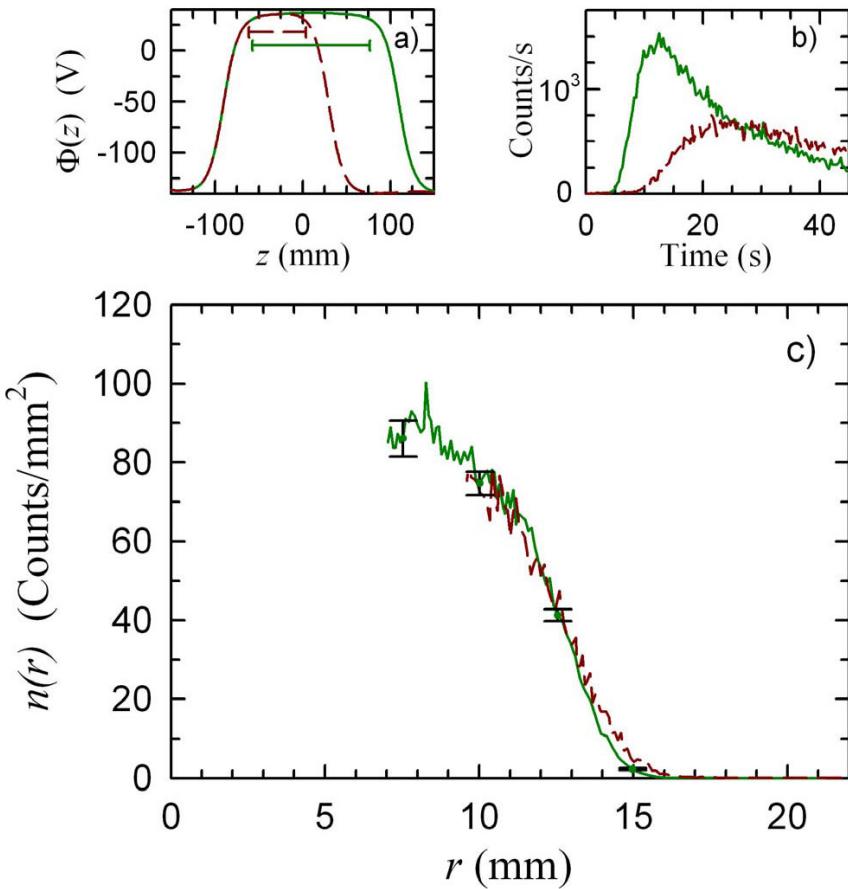
# Critical Radius



Critical radius depends on length of plasma, axial B, octupole B at wall, and wall radius

Increase octupole B with time and use wall hits vs time to extract radial distribution

# Application to Radial Density

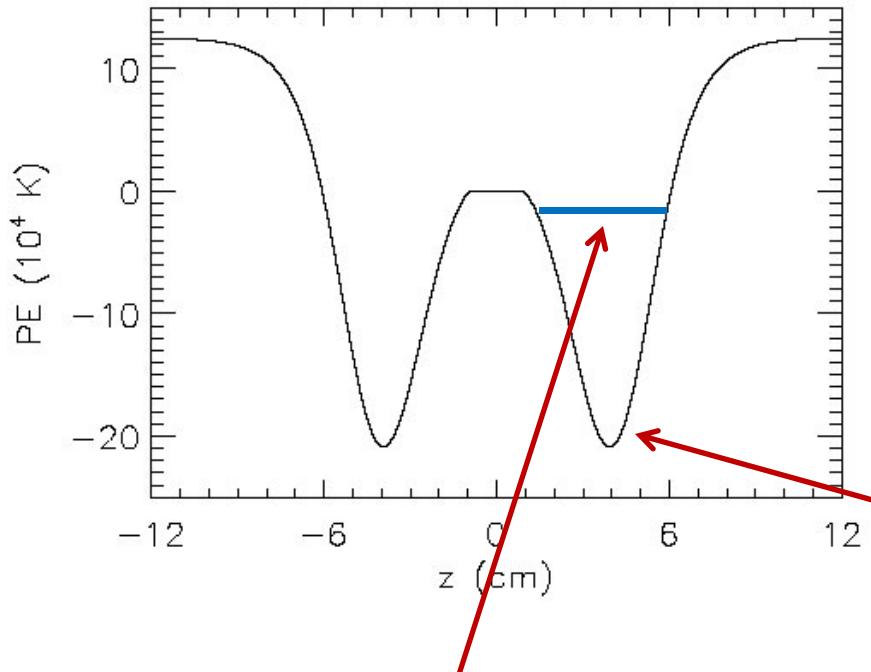


- a) potentials
- b) counts vs time
- c) convert to density vs radius

Although different time dependence of hits get same reconstructed density

Limitation: there is minimum r can extract

# What Happens When Mix



antiprotons

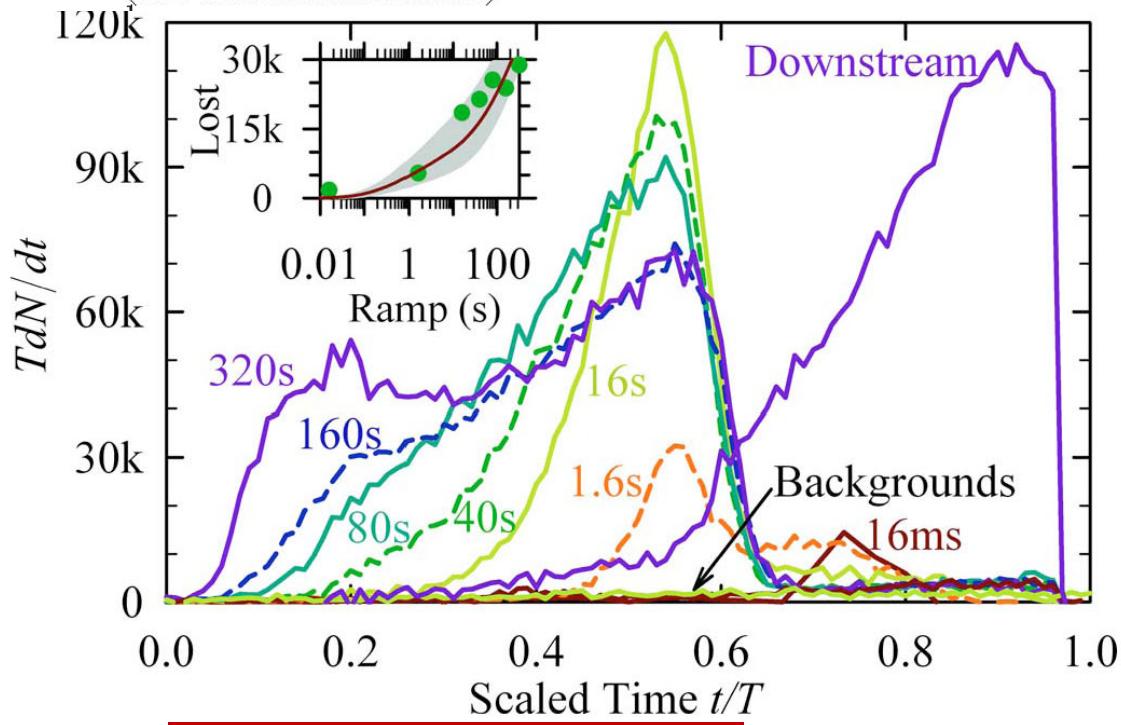
Mixing antiprotons  
that start out below the  
positron plasma

Slowly raise the  
bottom of the PE

Small radius antiprotons should  
be stable

# Magnetic multipole induced zero-rotation frequency bounce-resonant loss in a Penning–Malmberg trap used for antihydrogen trapping

G. B. Andresen,<sup>1</sup> W. Bertsche,<sup>2</sup> C. C. Bray,<sup>3</sup> E. Butler,<sup>2</sup> C. L. Cesar,<sup>4</sup> S. Chapman,<sup>3</sup> M. Charlton,<sup>2</sup> J. Fajans,<sup>3</sup> M. C. Fujiwara,<sup>5</sup> D. R. Gill,<sup>5</sup> W. N. Hardy,<sup>6</sup> R. S. Hayano,<sup>7</sup> M. E. Hayden,<sup>8</sup> A. J. Humphries,<sup>2</sup> R. Hydomako,<sup>9</sup> L. V. Jørgensen,<sup>2</sup> S. J. Kerrigan,<sup>2</sup> J. Keller,<sup>3</sup> L. Kurchaninov,<sup>5</sup> R. Lambo,<sup>4</sup> N. Madsen,<sup>2</sup> P. Nolan,<sup>10</sup> K. Olchanski,<sup>5</sup> A. Olin,<sup>5</sup> A. Povilus,<sup>3</sup> P. Pusa,<sup>10</sup> F. Robicheaux,<sup>11</sup> E. Sarid,<sup>12</sup> S. Seif El Nasr,<sup>6</sup> D. M. Silveira,<sup>7,13</sup> J. W. Storey,<sup>5</sup> R. I. Thompson,<sup>9</sup> D. P. van der Werf,<sup>2</sup> J. S. Wurtele,<sup>3</sup> and Y. Yamazaki<sup>13</sup>  
(ALPHA Collaboration)



Loss to walls

“Mysterious” loss  
of antiprotons  
when mixing the  
antiprotons with e<sup>+</sup>

More loss when  
ramp slower

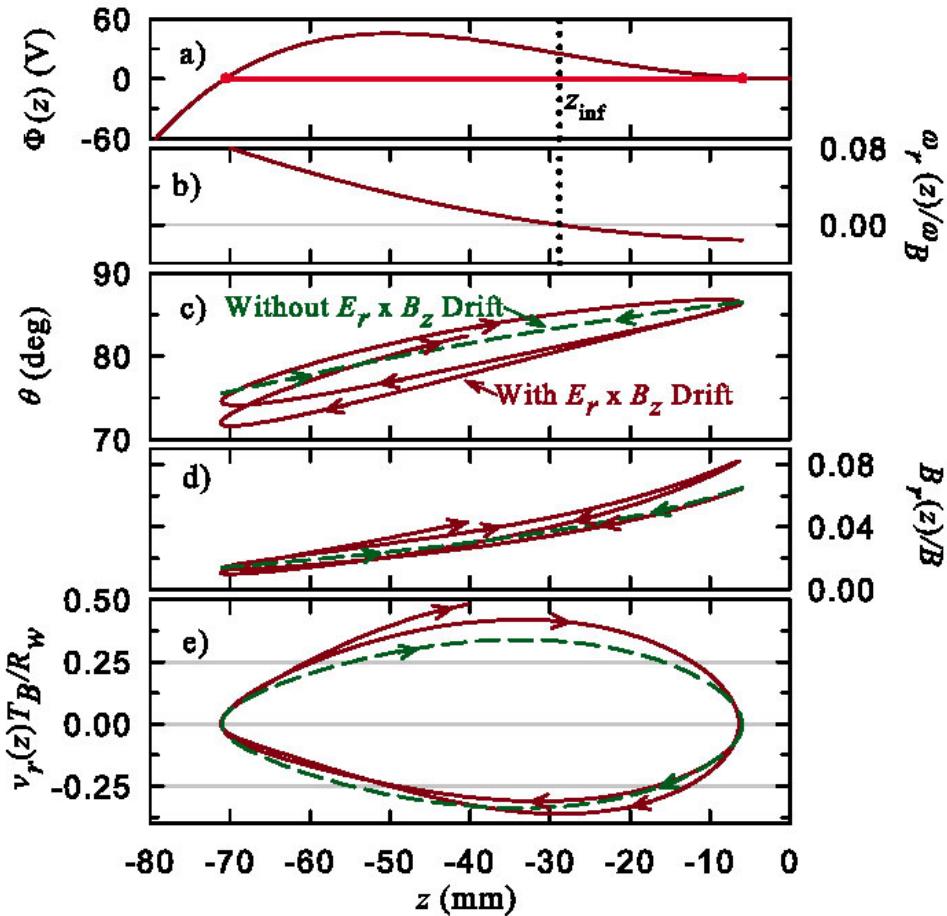
# Why a Mystery?

Radius of antiprotons should have been small enough to have almost all with less than critical radius.

Time scale too fast for collisions.

Time scale too slow for immediate transport.

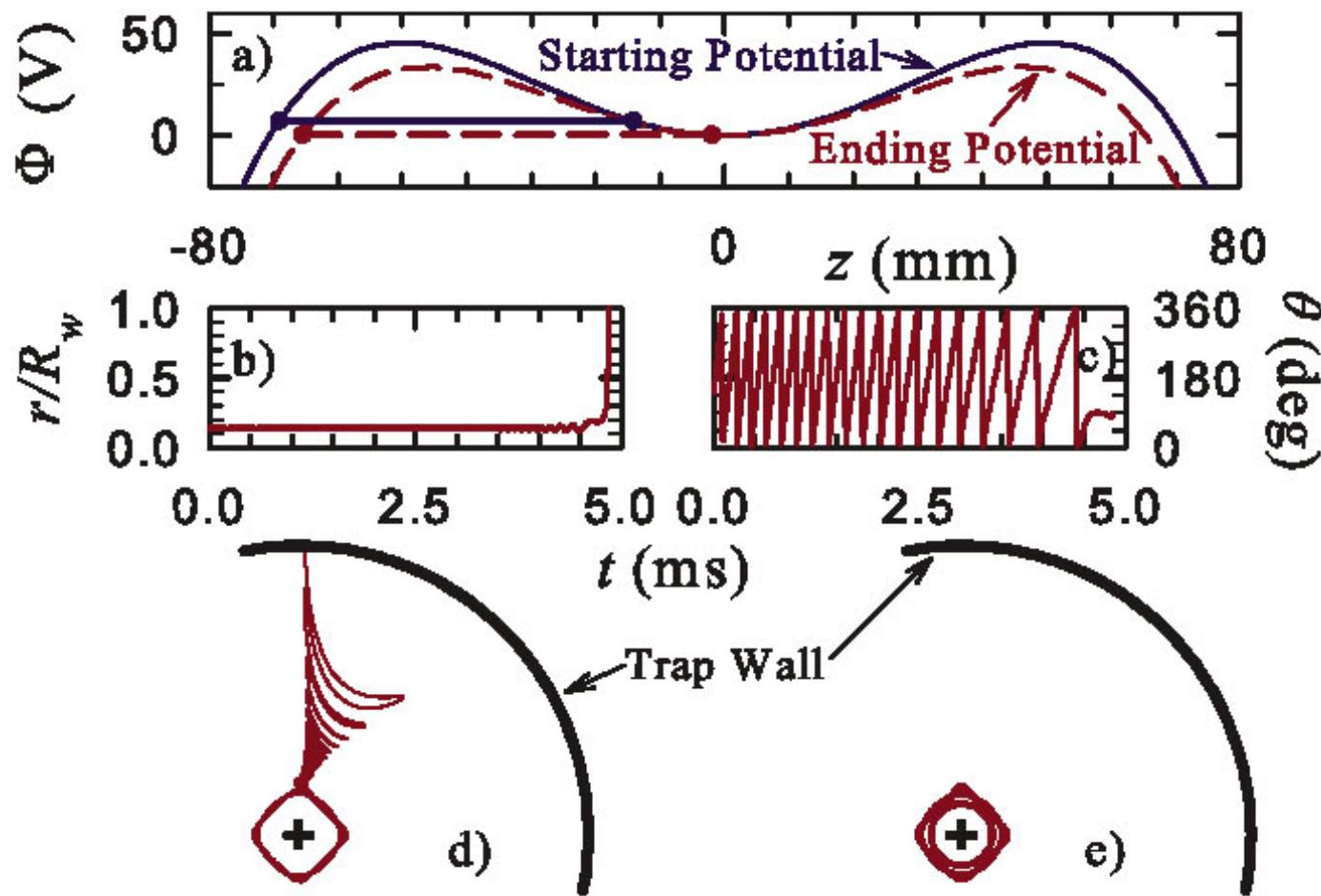
# Rotation Around Axis Stalls



As the electric potential is elongated, the rotation direction around axis changes sign.

As  $t$  changes, reach condition where 0 rotation  $\Rightarrow$  leads to resonance loss.

# Example Trajectory



Plots b-d are for same trajectory. Plot e has slightly different initial conditions and isn't lost.

# Regular and chaotic motion of anti-protons through a nested Penning trap with multipole magnetic fields

Jingjing Zhang, C L Taylor, J D Hanson and F Robicheaux

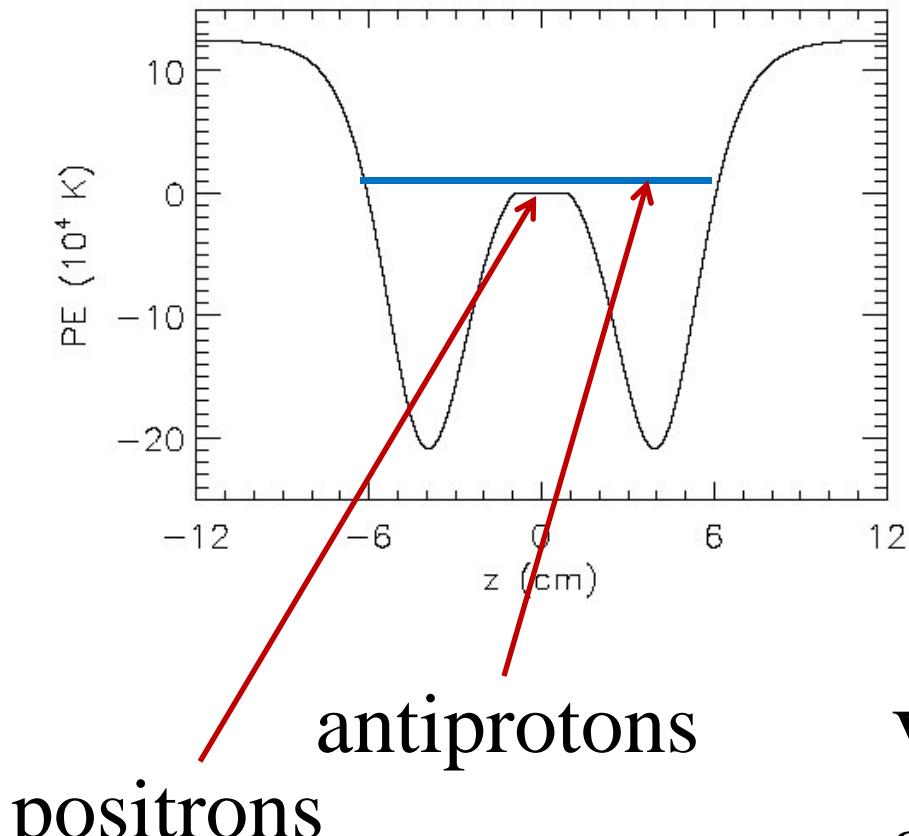
Department of Physics, Auburn University, AL 36849-5311, USA

Received 13 November 2006, in final form 29 December 2006

Published 19 February 2007

What does the antiproton trajectory look like when they are just above the level of the positrons.

# Antiproton Motion When Above

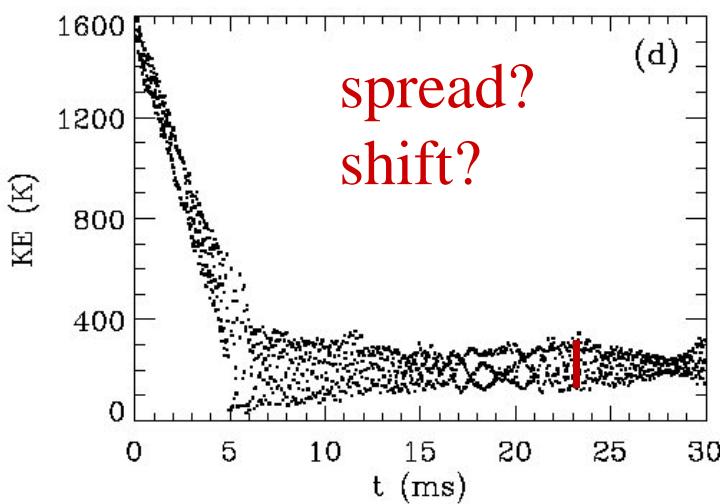
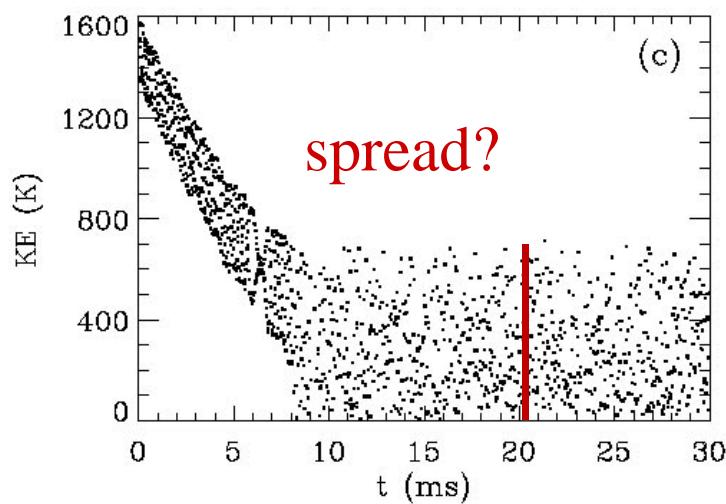
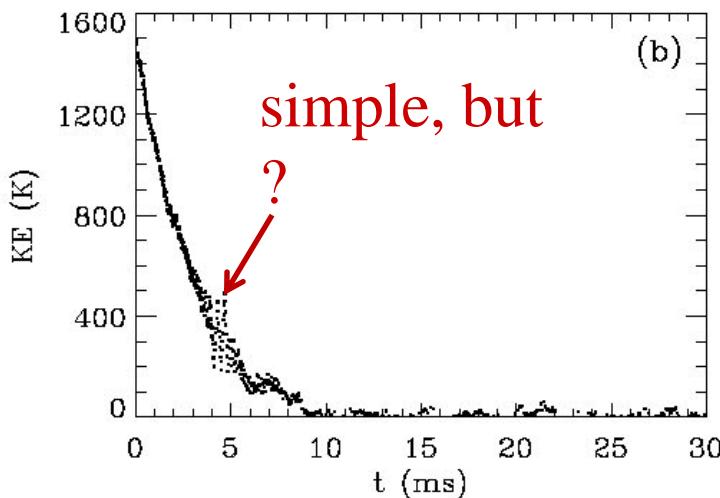
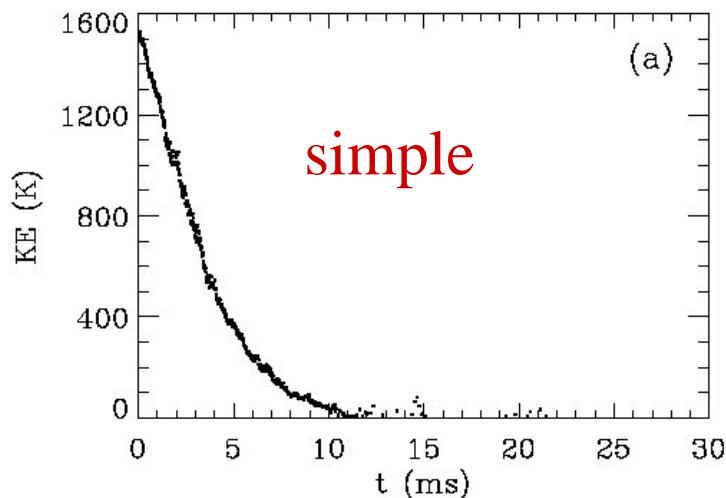


Antiprotons lose energy to the positrons by collisions/waves

Anything strange happens?

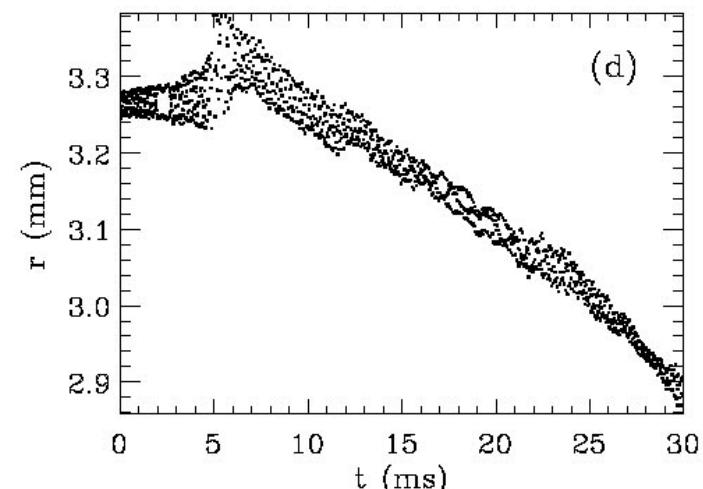
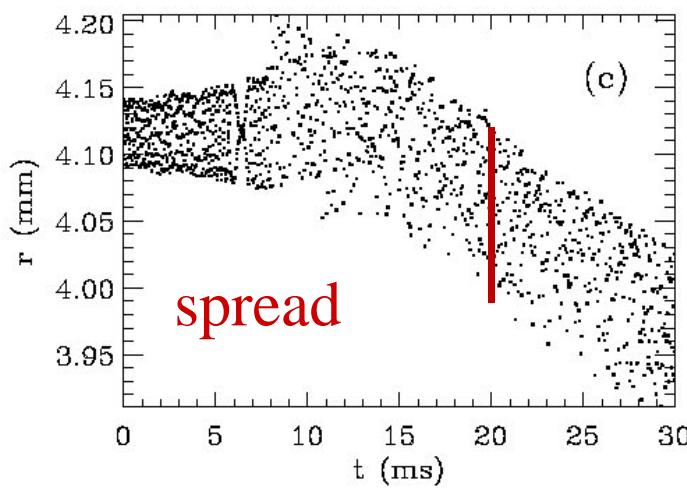
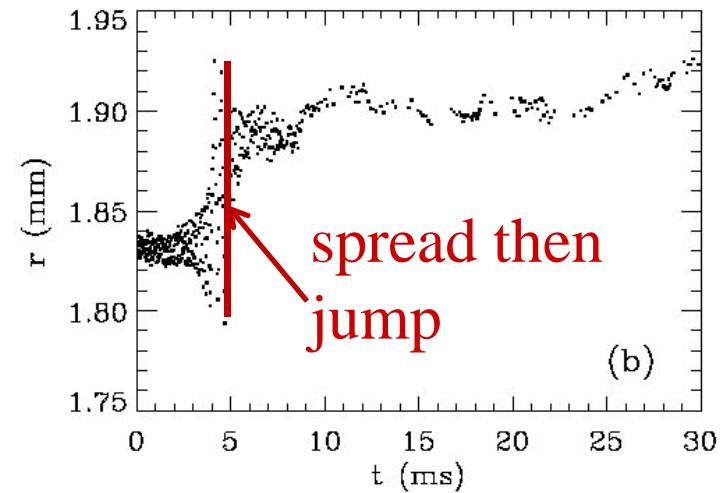
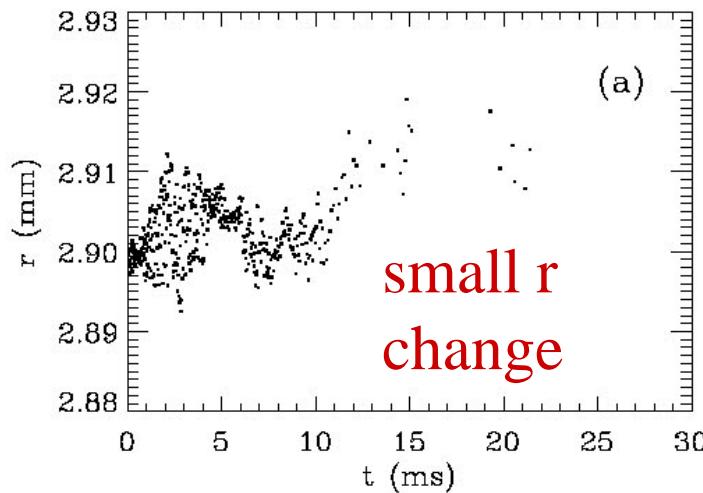
Very small radius so antiprotons motion should be simple

# KE at $z=0$ vs. $t$



4 different trajectories

# $r$ at $z=0$ vs. $t$

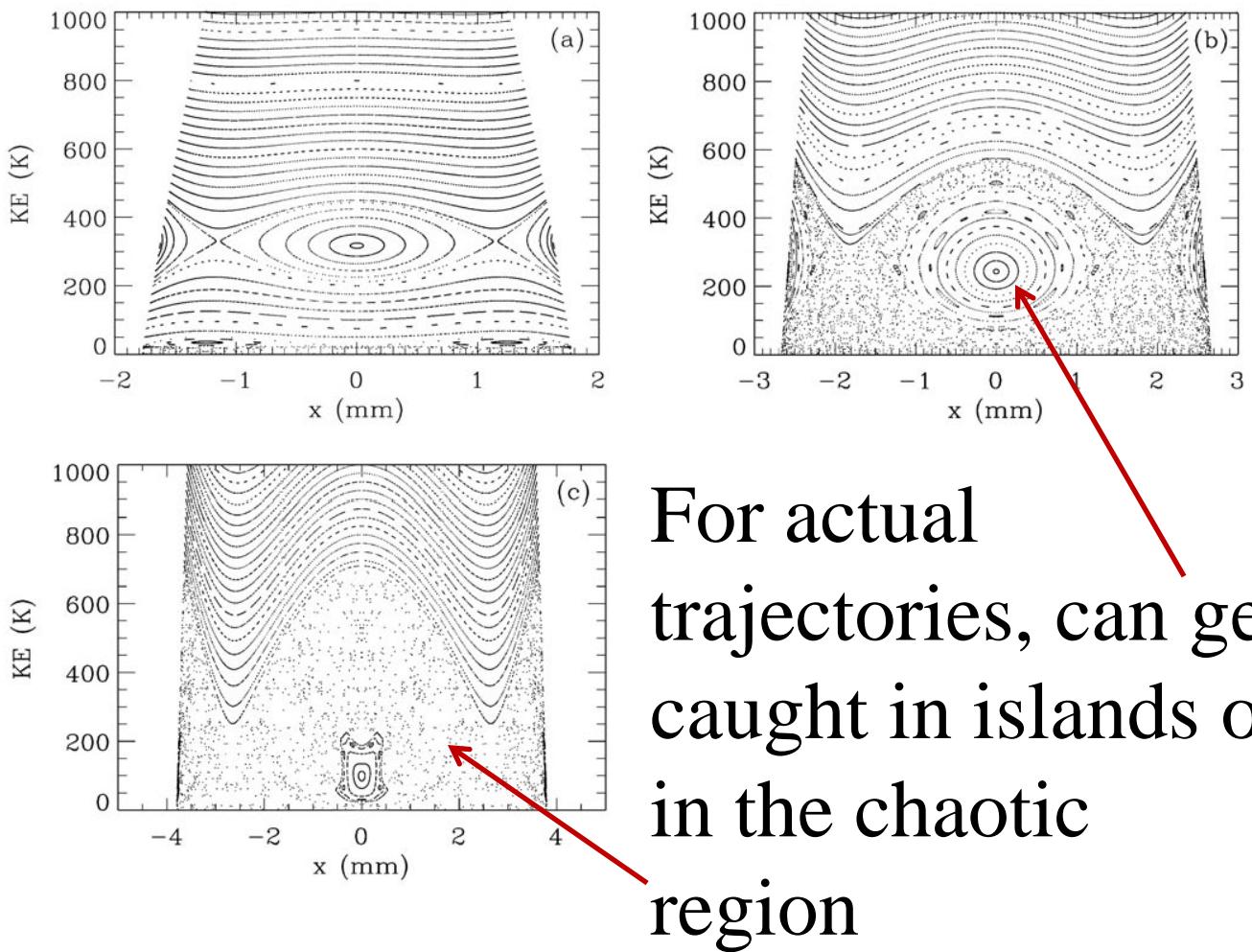


same 4 trajectories

# Poincare Surface of Section

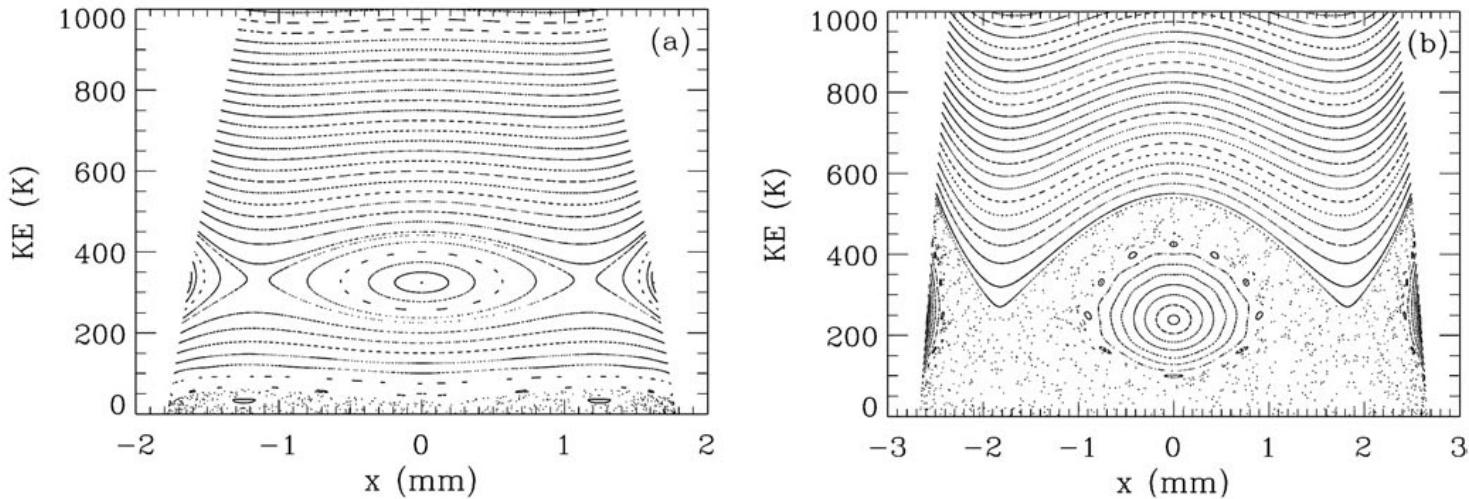
No KE  
loss to  
positrons

a to c is  
increasing  
radius



For actual  
trajectories, can get  
caught in islands or  
in the chaotic  
region

# Simple Map



Envelope is due to conservation of E (PE changes with r)

Large KE only small rotation in each axial bounce

Main island from rotation in positron  $\sim 1/4$

Size of chaotic region increases with r because variation of PE larger as r increases

# End of 1 Particle Motion

# Plasmas Change Things

Mirror magnets don't introduce qualitative changes because cylindrical symmetry

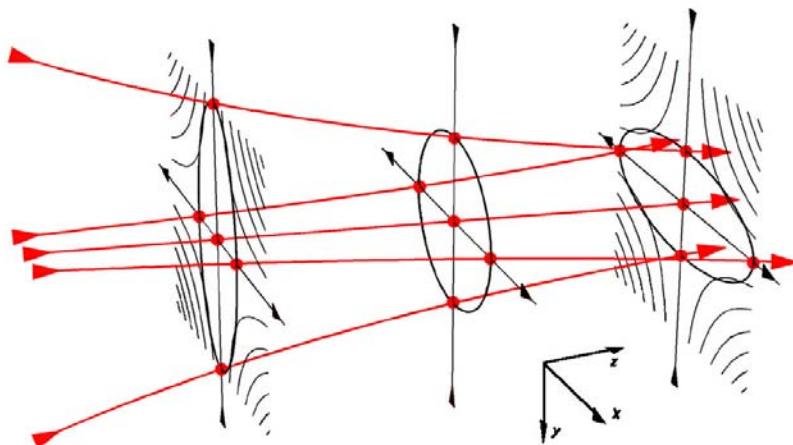
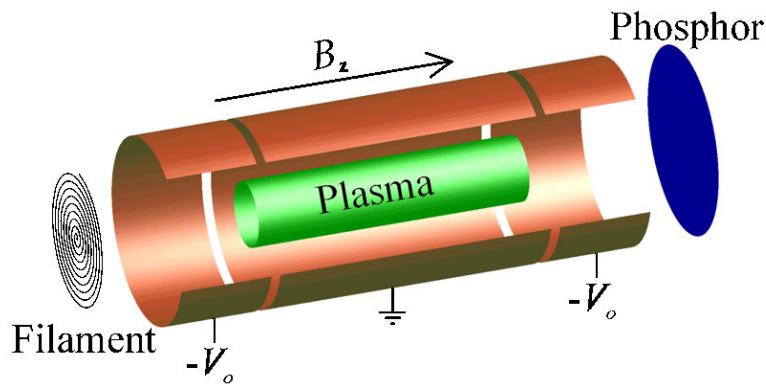
Collisions and waves give problems for holding plasmas with multipole fields because break cylindrical symmetry

# Quadrupole-Induced Resonant-Particle Transport in a Pure Electron Plasma

E. P. Gilson\* and J. Fajans†

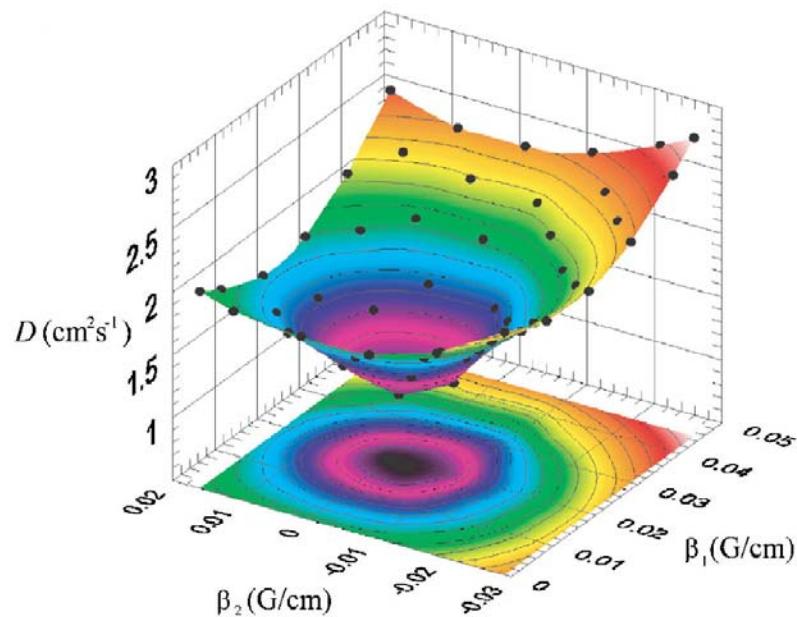
*Department of Physics, University of California, Berkeley, California 94720*

(Received 12 August 2002; revised manuscript received 28 October 2002; published 8 January 2003)

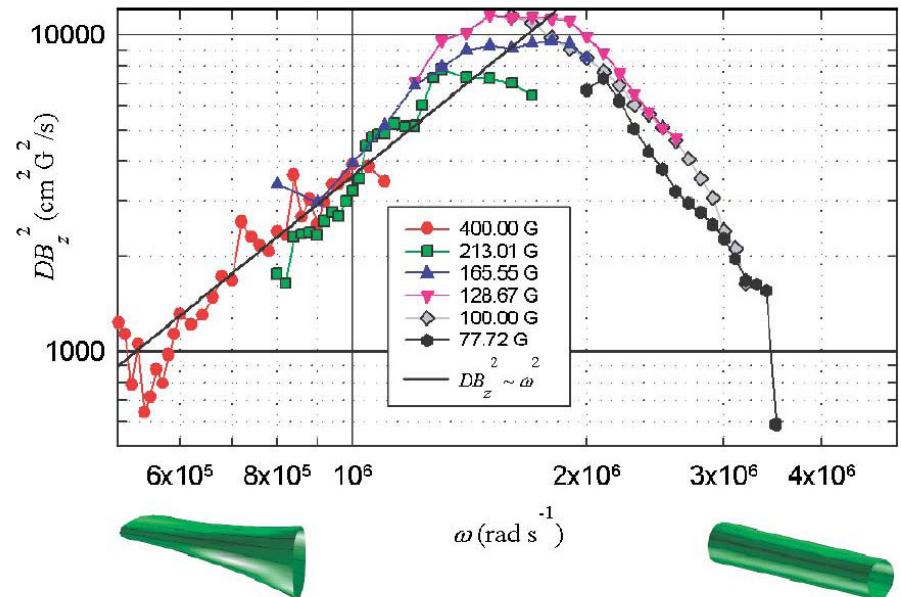


Look at transport in  
pure electron plasma vs  
quadrupole strength

# Diffusion With Quadrupole

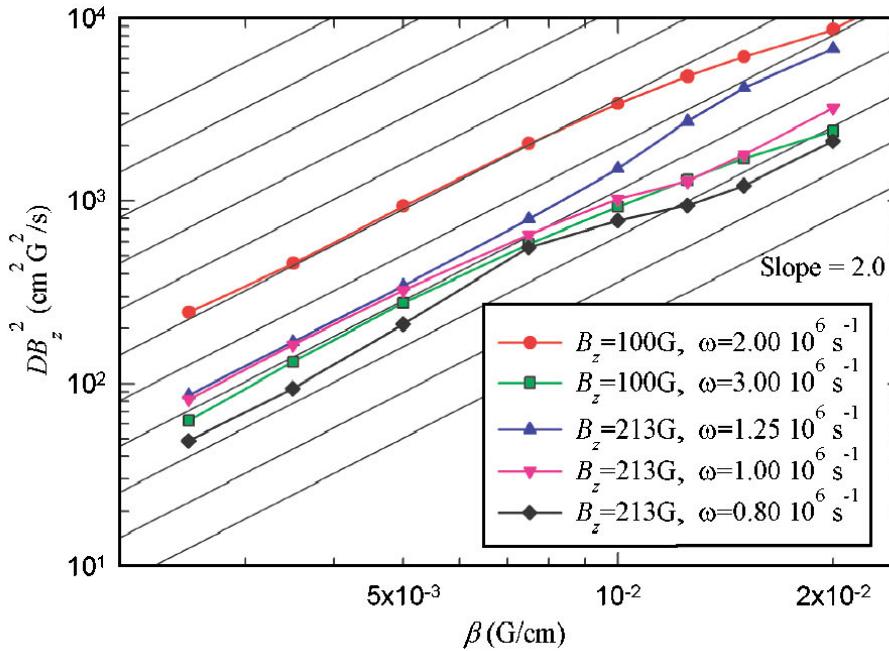


Diffusion as the  
quadrupole strength  
increases



Scaled diffusion vs  
plasma rotation  
frequency

# Scaling of Diffusion



Diffusion scales as square of quadrupole strength

# Theory of asymmetry-induced transport in a non-neutral plasma

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*Occidental College, Physics Department, Los Angeles, California 90041*

T. M. O'Neil

*University of California at San Diego, Department of Physics, La Jolla, California 92093*

(Received 19 January 1999; accepted 29 March 1999)

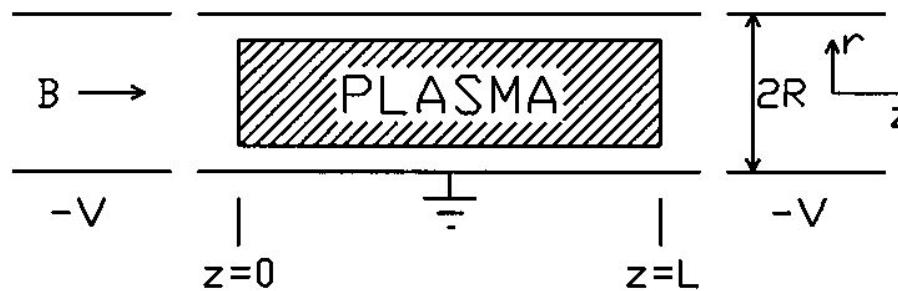


FIG. 1. Schematic of the plasma model used for this theory. The plasma is assumed to have flat ends and be of length  $L$ .

Transport dominated by particles that move in resonance with asymmetry

Agreement with experiment???

D. L. Eggleston, Phys. Plasmas 17, 042304 (2010).

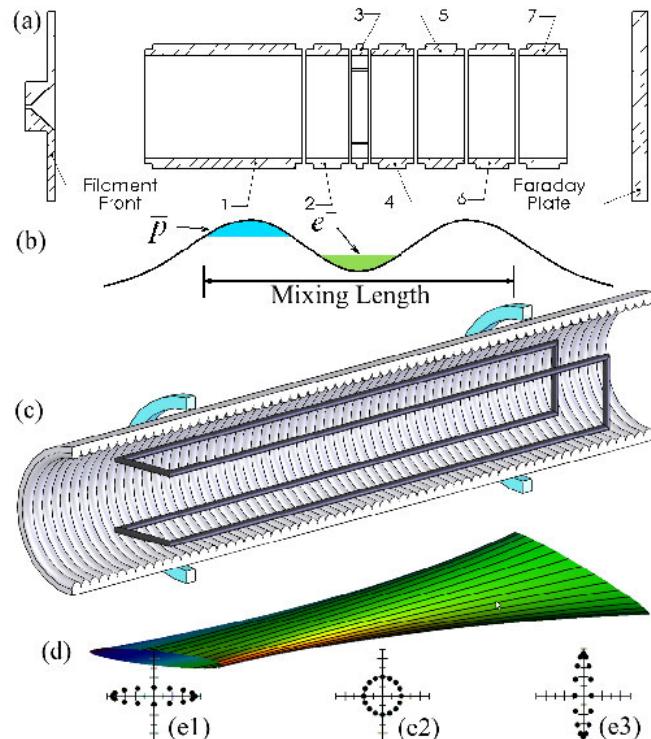
# Effects of Extreme Magnetic Quadrupole Fields on Penning Traps and the Consequences for Antihydrogen Trapping

J. Fajans,<sup>1,\*</sup> W. Bertsche,<sup>1</sup> K. Burke,<sup>1</sup> S. F. Chapman,<sup>1</sup> and D. P. van der Werf<sup>2</sup>

<sup>1</sup>*Department of Physics, University of California at Berkeley, and the Lawrence Berkeley National Laboratory, Berkeley California 94720, USA*

<sup>2</sup>*Department of Physics, University of Wales Swansea, Singleton Park SA2 8PP, United Kingdom*

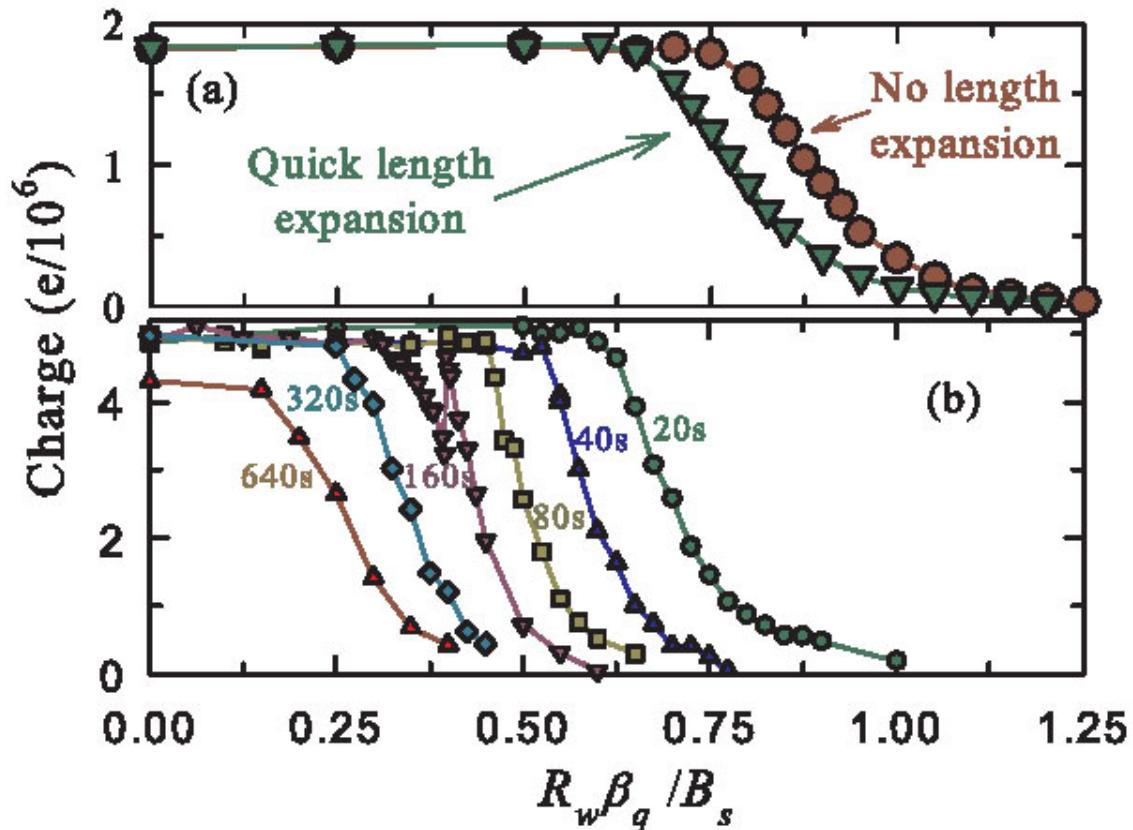
(Received 16 November 2004; published 4 October 2005)



## Immediate and diffusive transport to walls

“...suggest quadrupoles cannot be used to trap antihydrogen.”

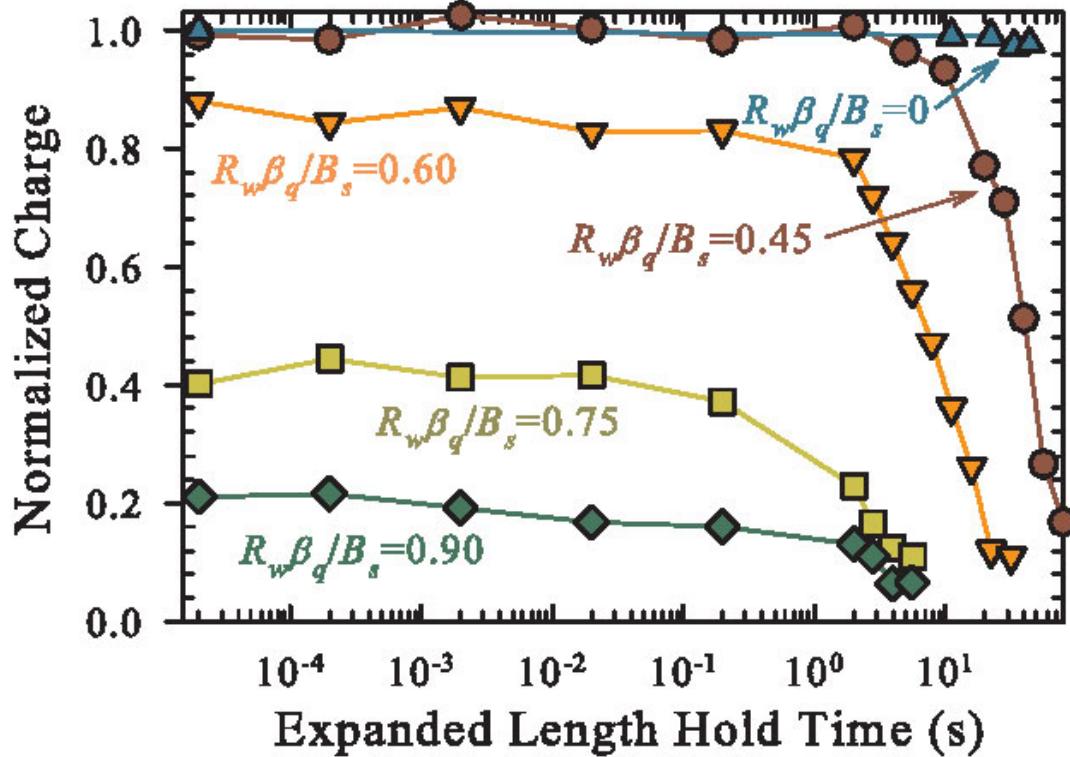
# Charge Remaining



Expand from 1 to 4 cm

Remaining as a function of quadrupole strength: different hold times (2 cm)

# Charge Remaining



Remaining as a function of hold time: different quadrupole strengths (4 cm)

# Antimatter Plasmas in a Multipole Trap for Antihydrogen

G. Andresen,<sup>1</sup> W. Bertsche,<sup>2</sup> A. Boston,<sup>3</sup> P. D. Bowe,<sup>1</sup> C. L. Cesar,<sup>4</sup> S. Chapman,<sup>2</sup> M. Charlton,<sup>5</sup> M. Chartier,<sup>3</sup> A. Deutsch,<sup>2</sup> J. Fajans,<sup>2</sup> M. C. Fujiwara,<sup>6</sup> R. Funakoshi,<sup>7</sup> D. R. Gill,<sup>6</sup> K. Gomberoff,<sup>2,8</sup> J. S. Hangst,<sup>1</sup> R. S. Hayano,<sup>7</sup> R. Hydomako,<sup>9</sup> M. J. Jenkins,<sup>5</sup> L. V. Jørgensen,<sup>5</sup> L. Kurchaninov,<sup>6</sup> N. Madsen,<sup>5</sup> P. Nolan,<sup>3</sup> K. Olchanski,<sup>6</sup> A. Olin,<sup>6</sup> A. Povilus,<sup>2</sup> F. Robicheaux,<sup>10</sup> E. Sarid,<sup>11</sup> D. M. Silveira,<sup>4</sup> J. W. Storey,<sup>6</sup> H. H. Telle,<sup>5</sup> R. I. Thompson,<sup>9</sup> D. P. van der Werf,<sup>5</sup> J. S. Wurtele,<sup>2</sup> and Y. Yamazaki<sup>12</sup>

(ALPHA Collaboration)

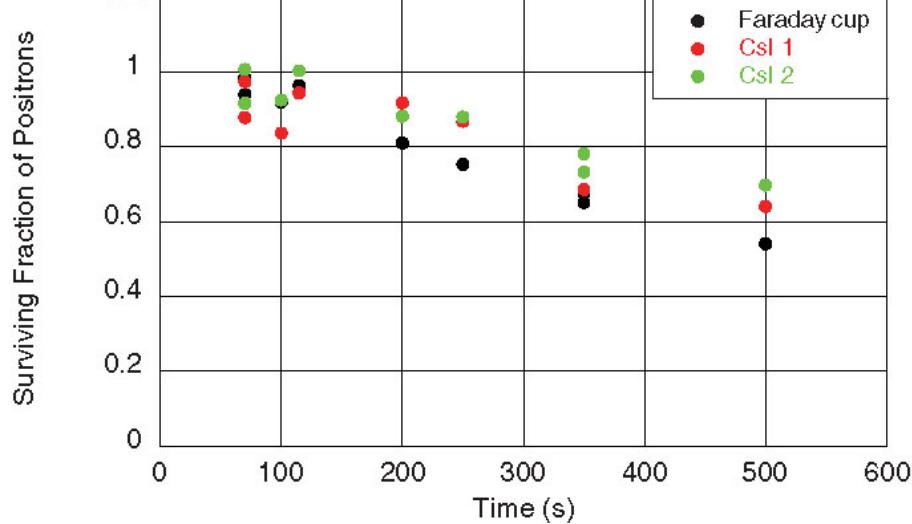
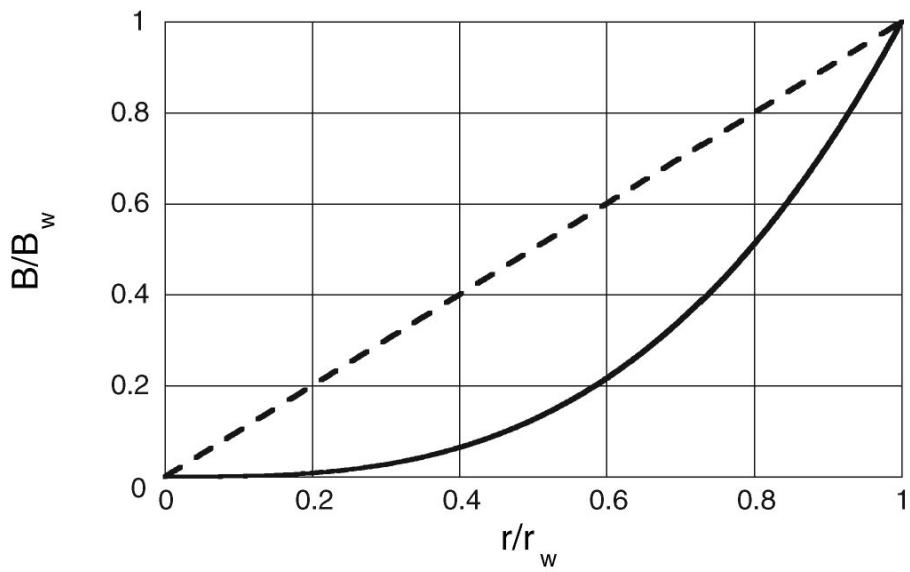
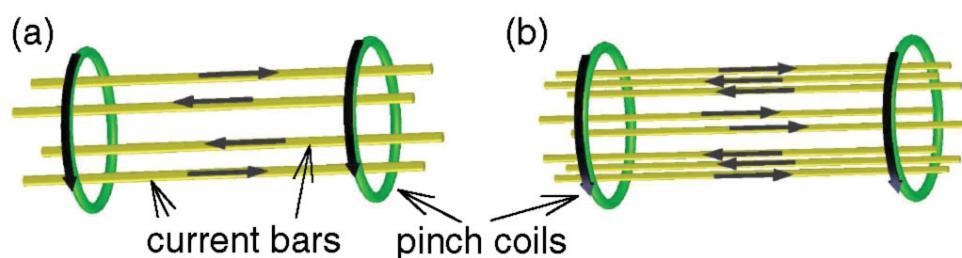
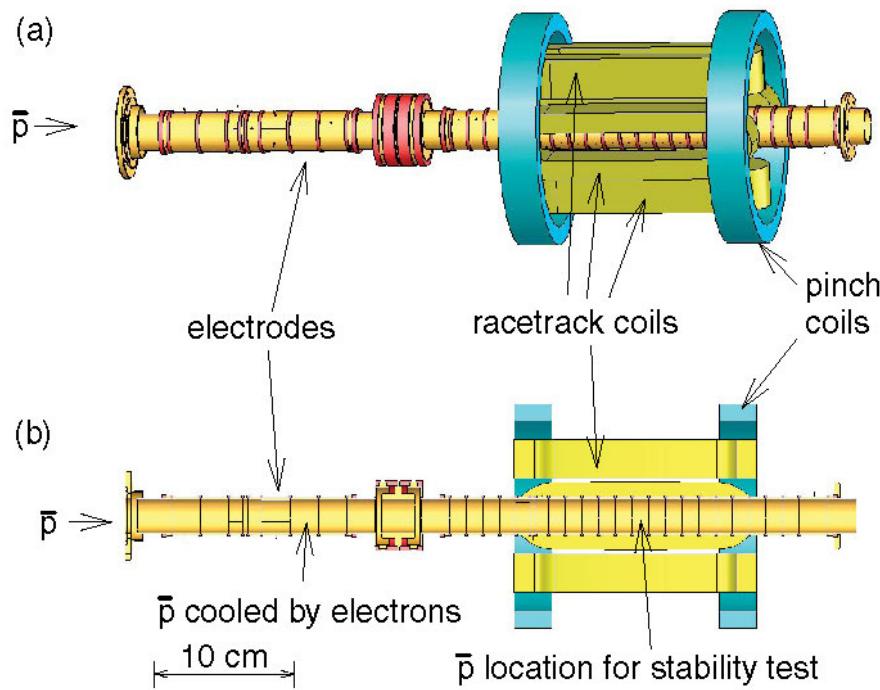


FIG. 2. Magnetic field strength versus radius for an ideal quadrupole (dashed line) and an ideal octupole (solid line).  $B_w$  is the field at the inner wall (radius  $r_w$ ) of the Penning trap.

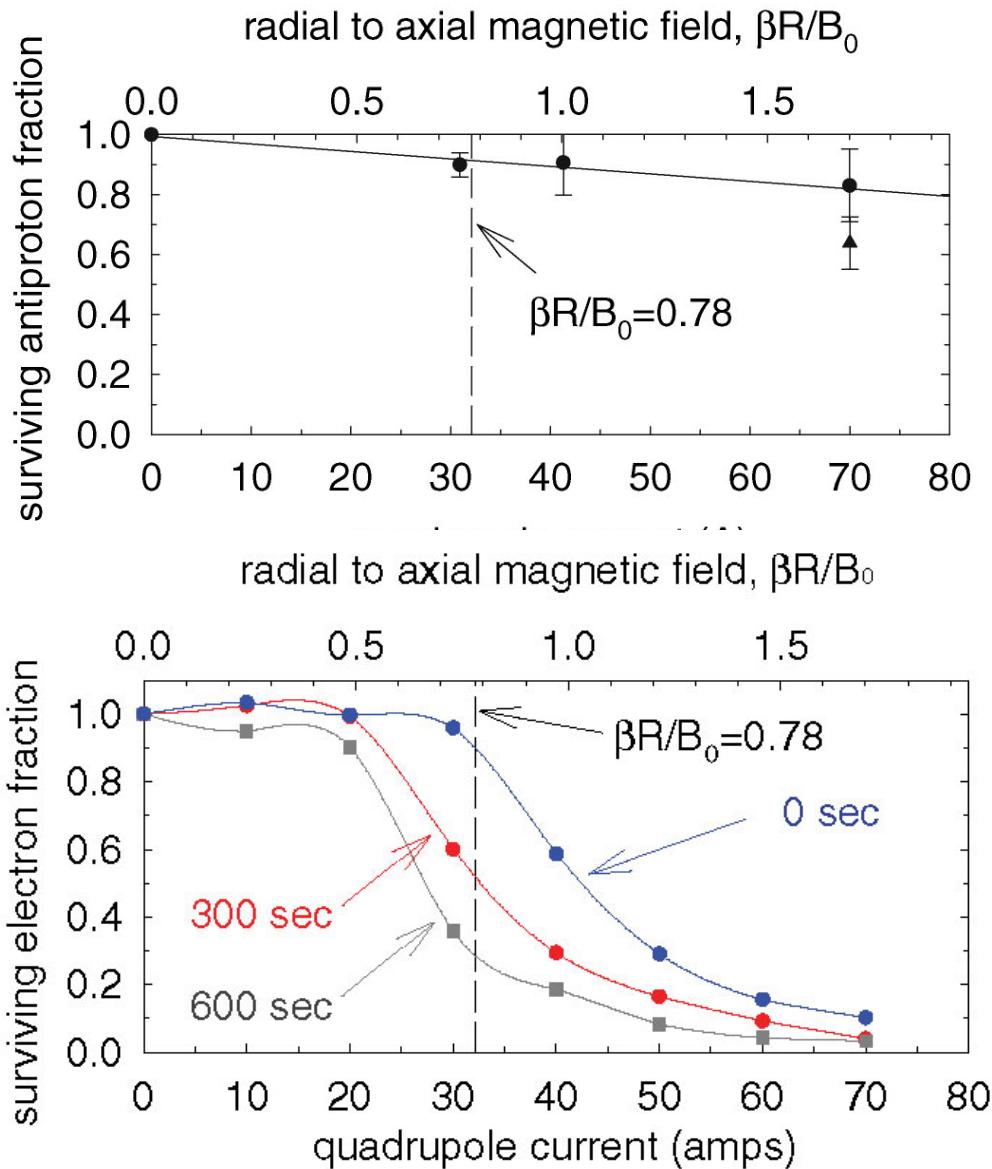
# Antiproton Confinement in a Penning-Ioffe Trap for Antihydrogen

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(ATRAP Collaboration)



# Results



Trap depth is 375 mK for ground state H when  $\beta R/B_0 = 0.78$

Fraction of trapped antiH scales like trap depth $^{3/2}$

# Abstract & Conclusion

“This first demonstration with antiprotons suggests that quadrupole Ioffe traps can be superimposed upon antiproton and positron traps to attempt to capture antihydrogen atoms as they form, *contrary to conclusions of previous analyses.*”

“Quadrupole Ioffe traps thus seem to have a role to play for antihydrogen formed by charge exchange, *in spite of earlier claims to the contrary.*”

(emphasis added)

# Simulations of plasma confinement in an antihydrogen trap

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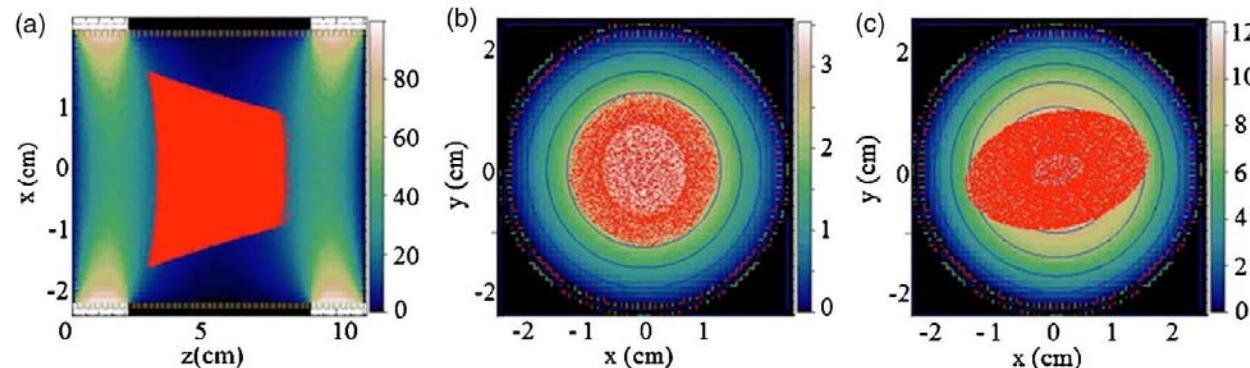
A. Friedman, D. Grote, and J.-L. Vay

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California 94720, USA*

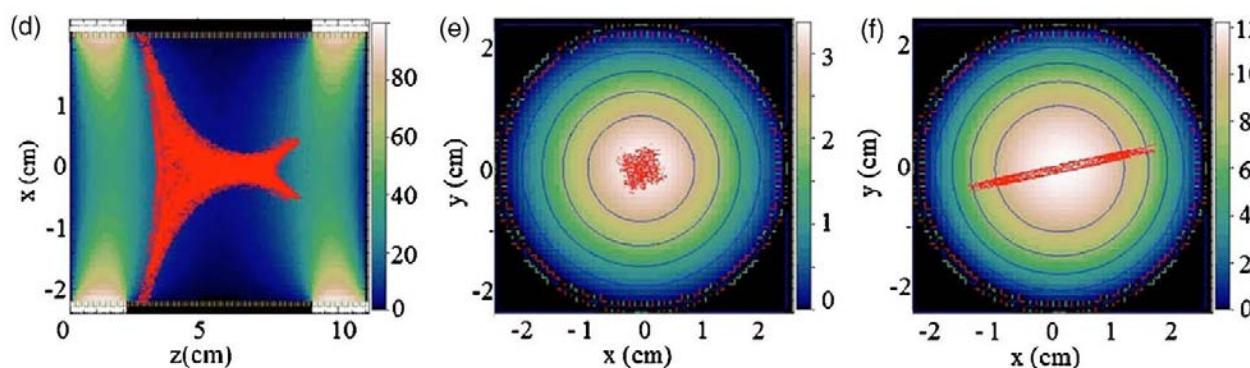
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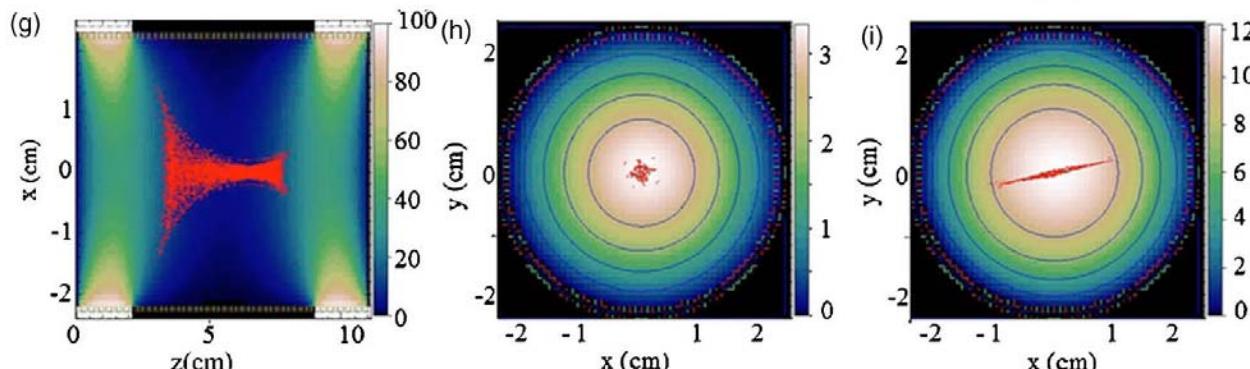
# Quadrupole



$0.6 \mu\text{s}$



$1.0 \mu\text{s}$



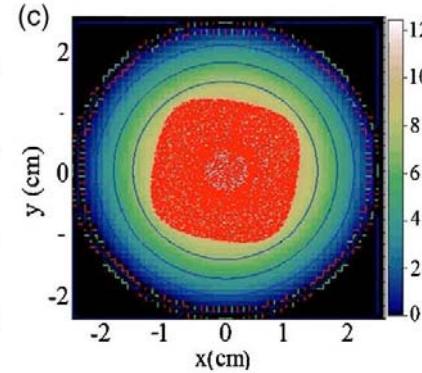
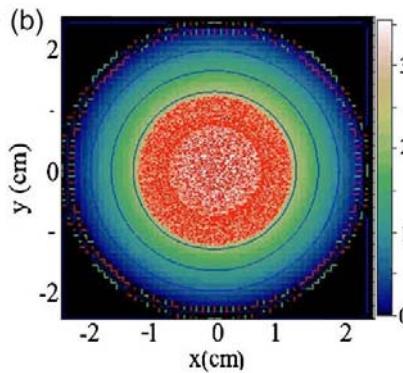
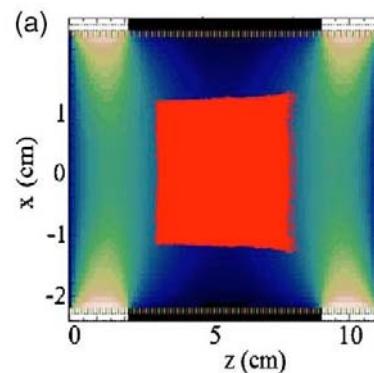
$3.0 \mu\text{s}$

project y

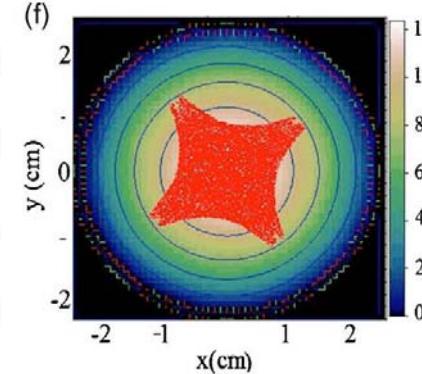
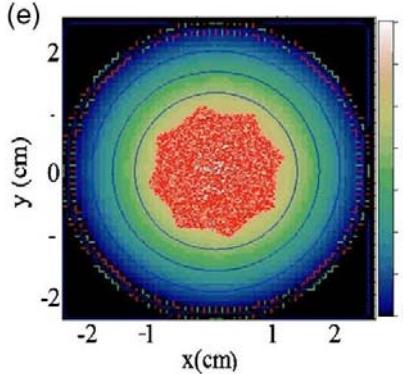
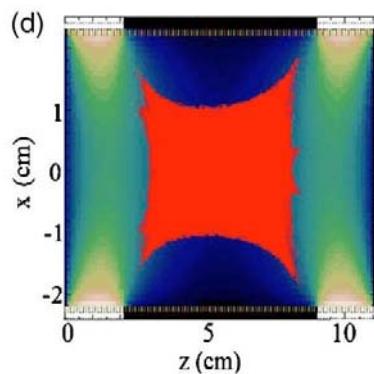
center

edge

# Octupole



0.6  $\mu$ s



3.0  $\mu$ s

project y

center

edge

# Survival vs. time

