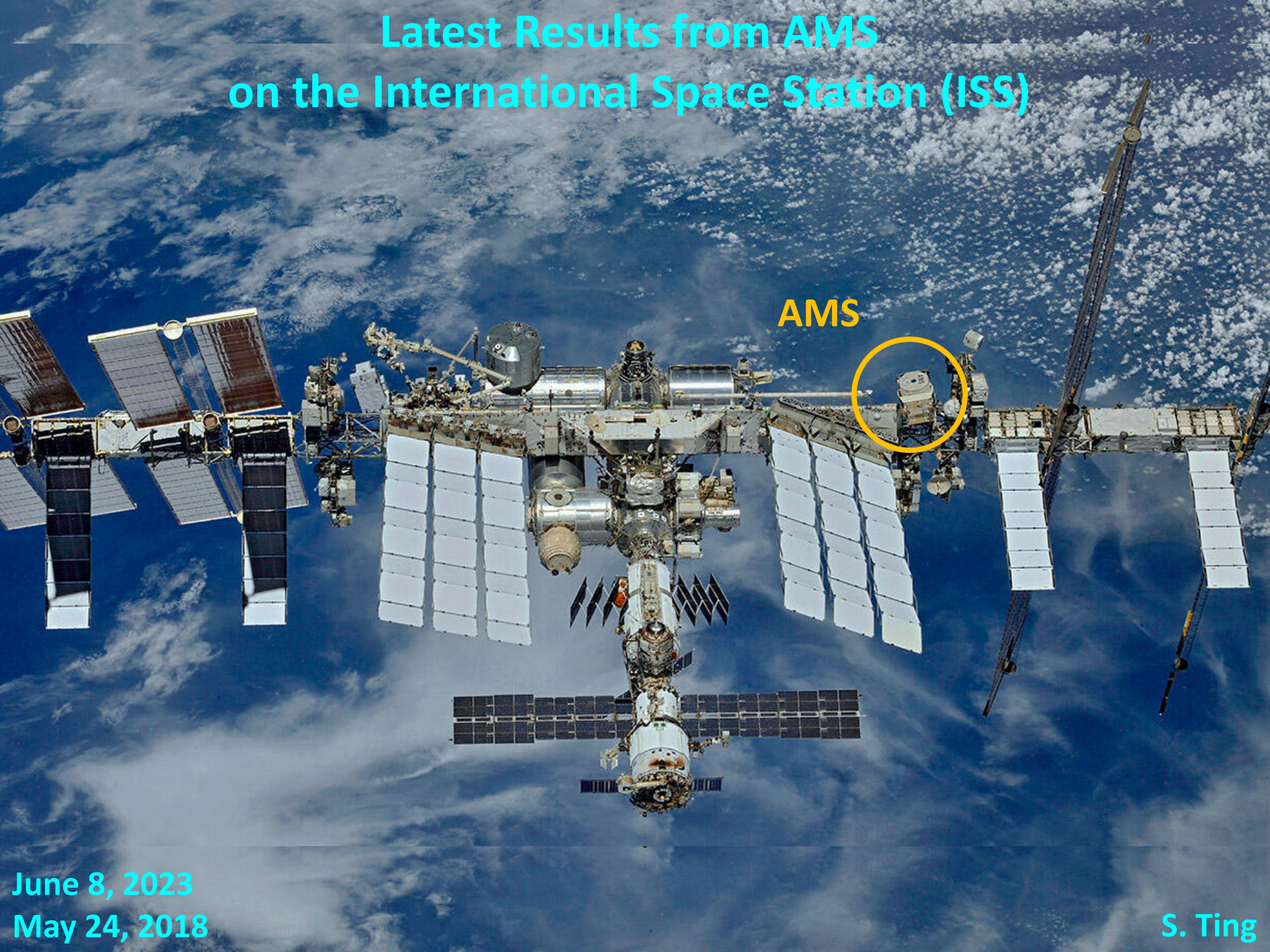


# Latest Results from AMS on the International Space Station (ISS)



AMS

June 8, 2023  
May 24, 2018

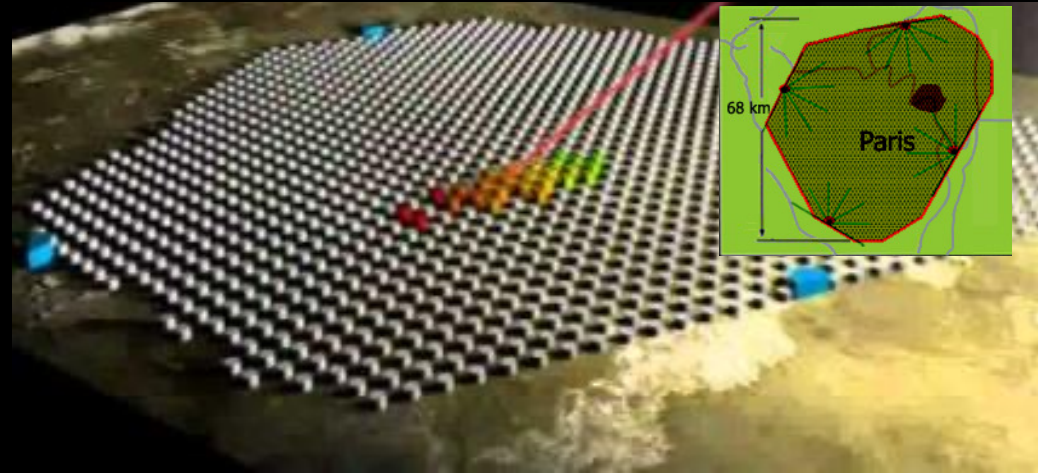
S. Ting

# Examples of Current Ground-Based Cosmic Ray Experiments

## The Pierre Auger Observatory



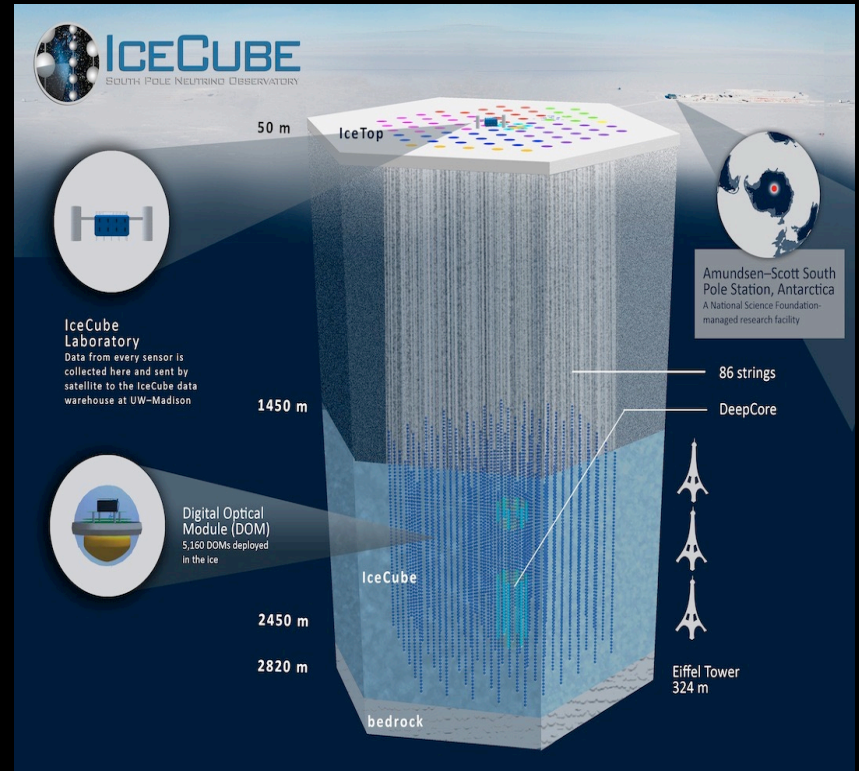
HAWC



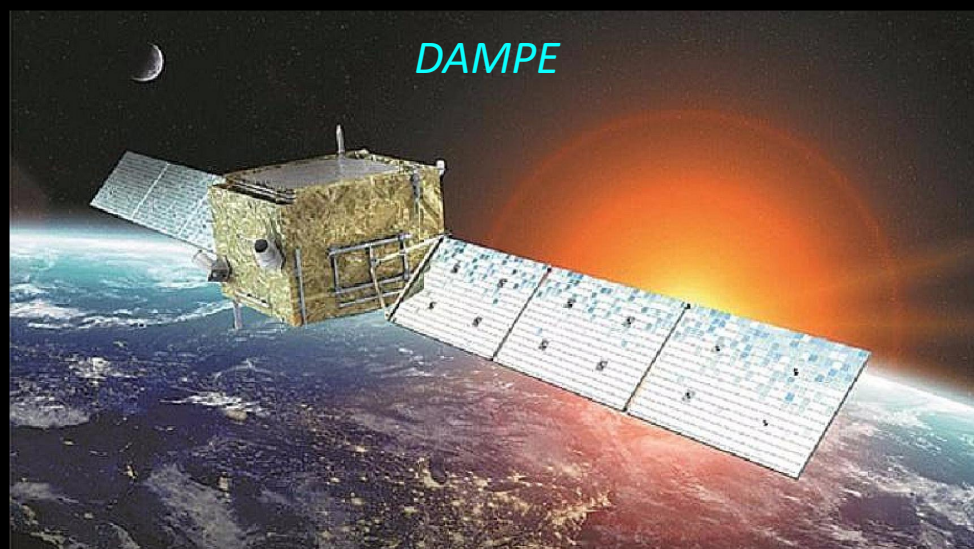
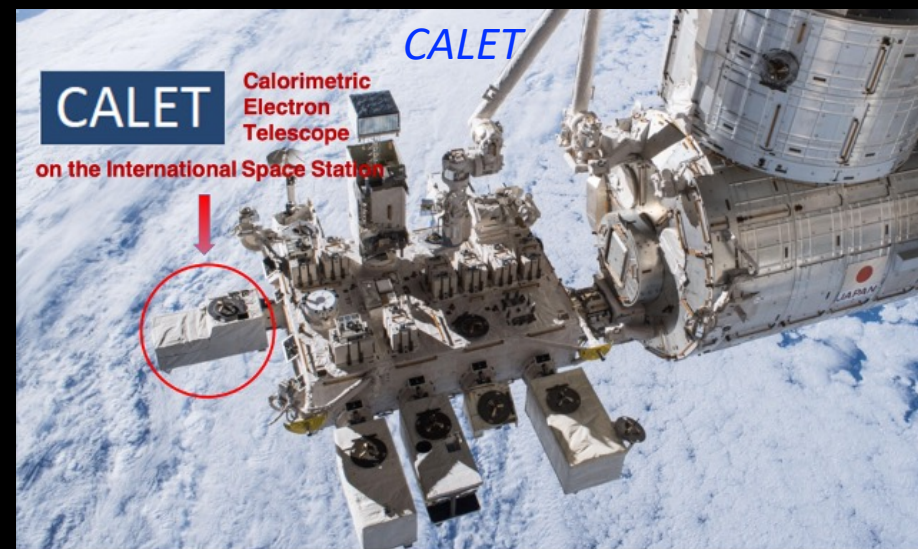
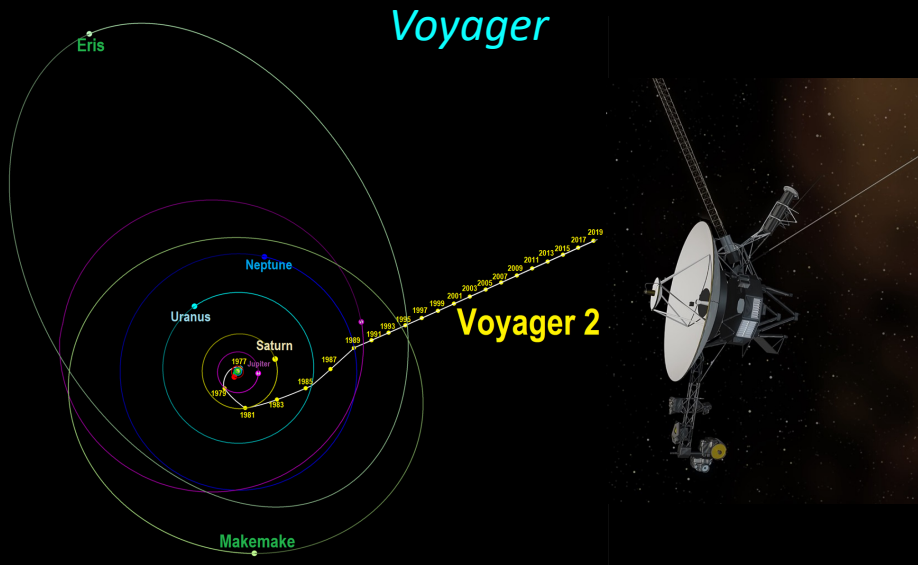
LHAASO



HESS



# Examples of Current Non-magnetic, Calorimeter Experiments in Space



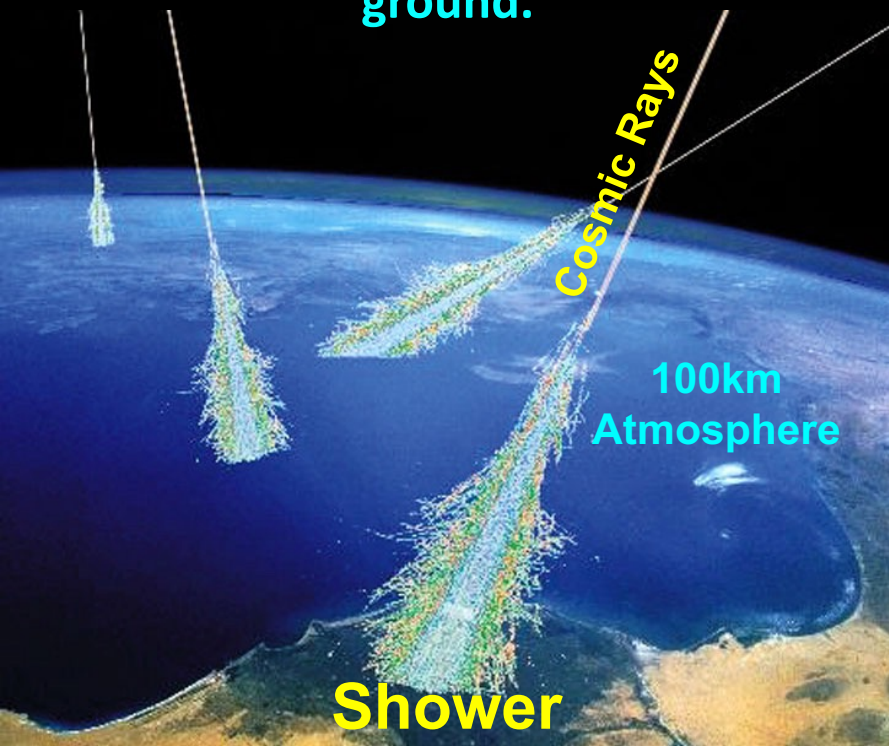
# AMS on the Space Station:

Physics of Dark Matter, Antimatter, the Origin of the Cosmos, and new phenomena through the precision, long-duration measurement of charged cosmic rays

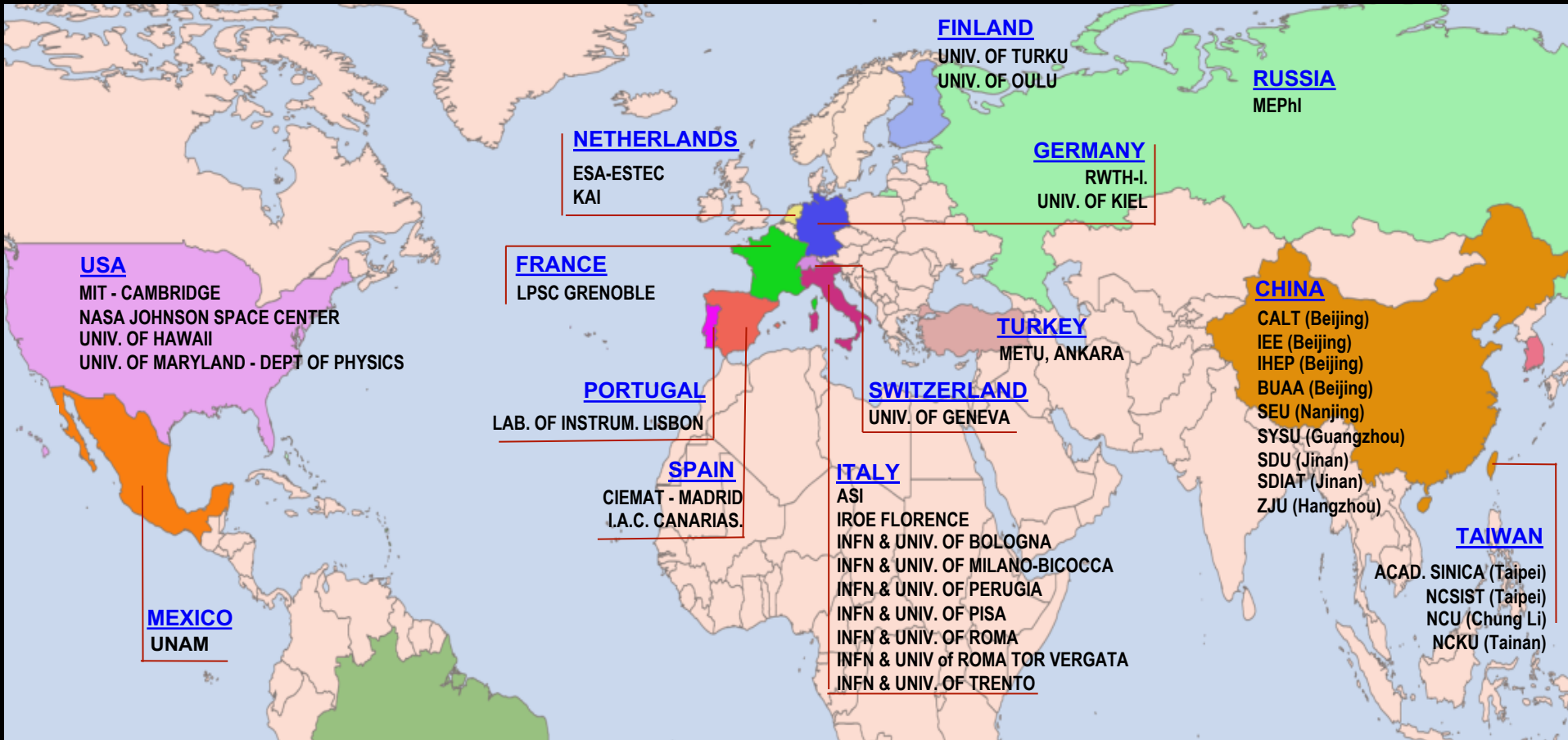
Charged cosmic rays have mass.  
They are absorbed by the  
100 km of Earth's atmosphere  
(10m of water).

The properties ( $\pm Z$ ,  $P$ ) of charged cosmic rays cannot be studied on the ground.

To measure cosmic ray charge and momentum requires a magnetic spectrometer in space



# Alpha Magnetic Spectrometer experiment (AMS) on the Space Station



**We thank Professor F. Gianotti, CERN;  
NASA and DOE, U.S.; ESA; ASI and INFN, Italy; DLR, RWTH-Aachen, and Julich, Germany;  
CNES and CNRS/IN2P3, France; CDTI, Spain; Academia Sinica and NCKU, Taiwan;  
IHEP, SDU, SDIAT, and SEU, China; UniGe, Switzerland;  
for their strong support.**

# The strong support from NASA and CERN has been essential to the success of AMS



**Ken Bollweg, Manager, AMS Project Office, leads the interface between AMS and NASA**



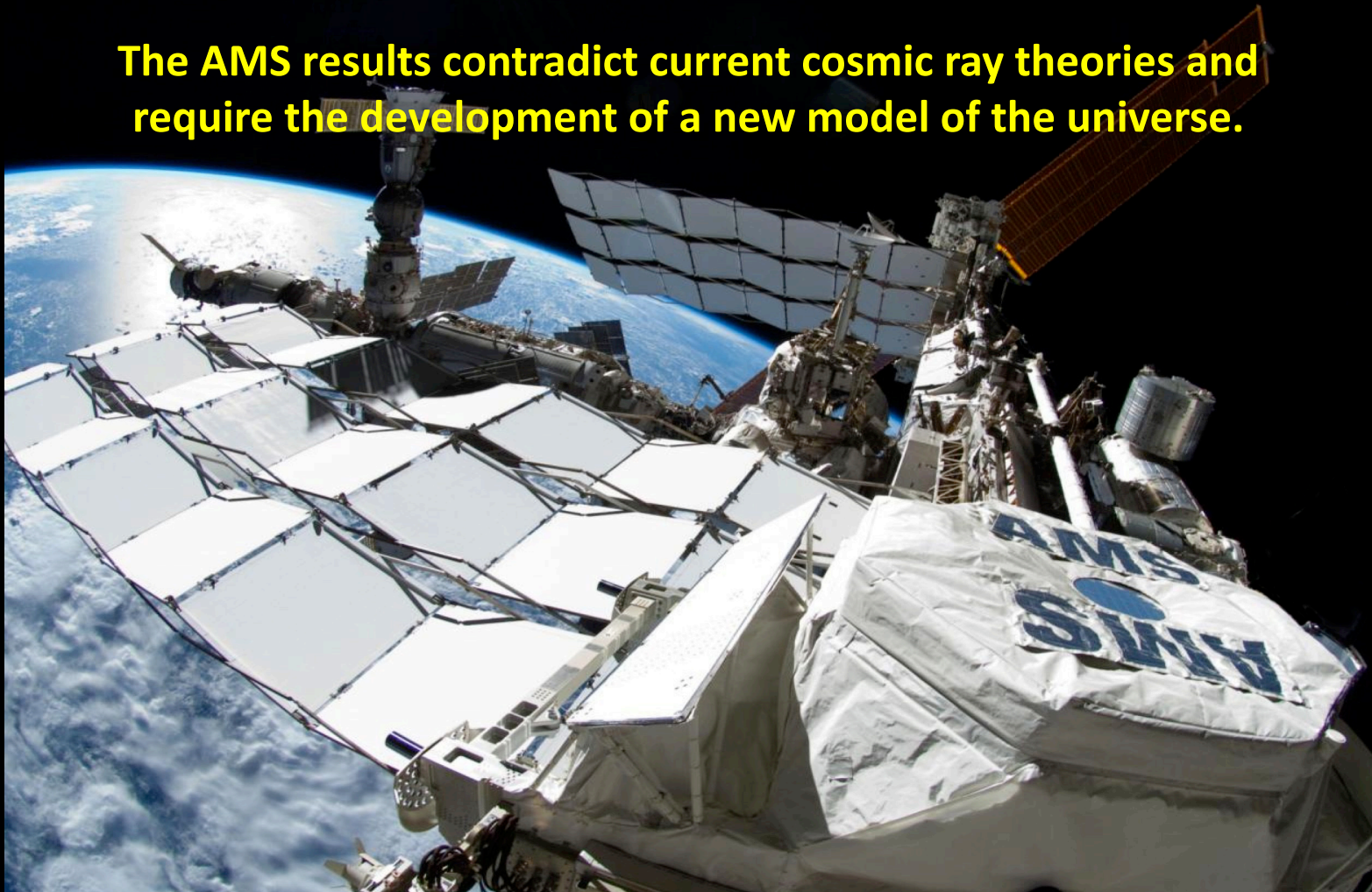
**Dr. Corrado Gargiulo led the assembly of AMS at CERN and is leading the current AMS upgrade to increase the acceptance by 300%**

**We also thank Dieter Schinzel, Andrzej Siemko, and Bernd Panzer-Steindel for their strong support**

**In the past hundred years, measurements of charged cosmic rays by balloons and satellites have typically had  $\sim(30-50)\%$  accuracy.**

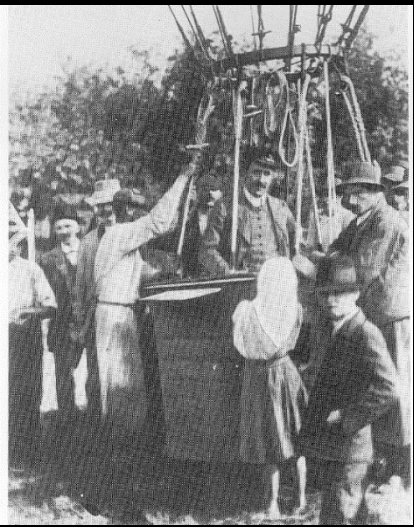
**AMS is providing cosmic ray information with  $\sim 1\%$  accuracy. The improvement in accuracy and extension of the energy range is providing new insights.**

**The AMS results contradict current cosmic ray theories and require the development of a new model of the universe.**



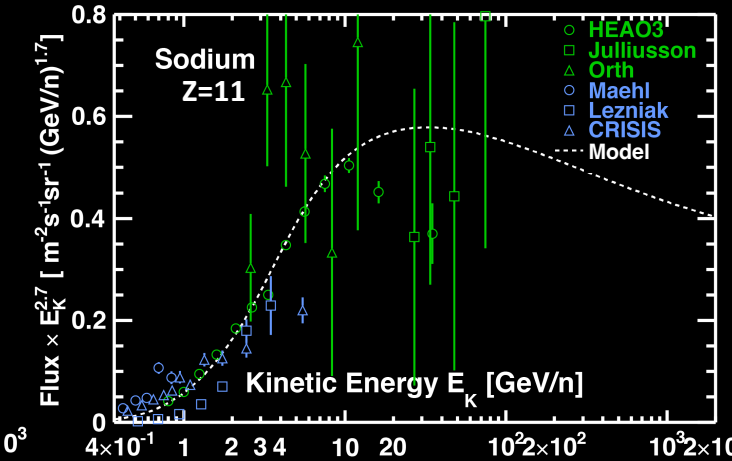
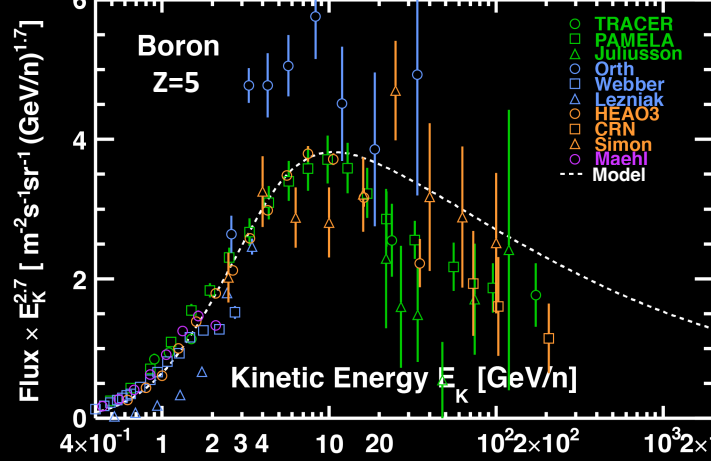
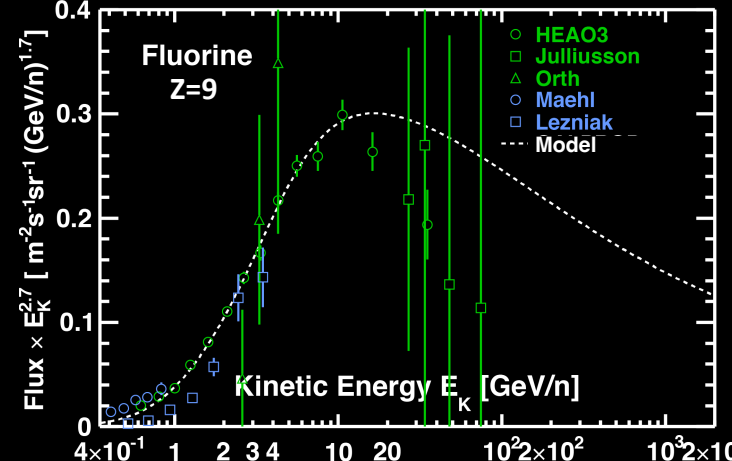
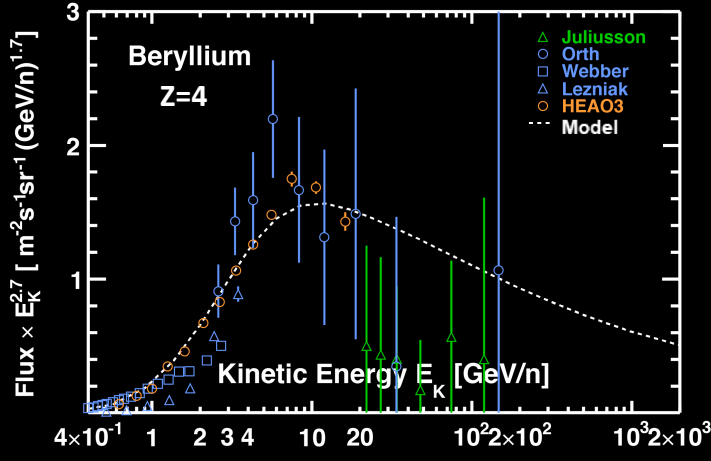
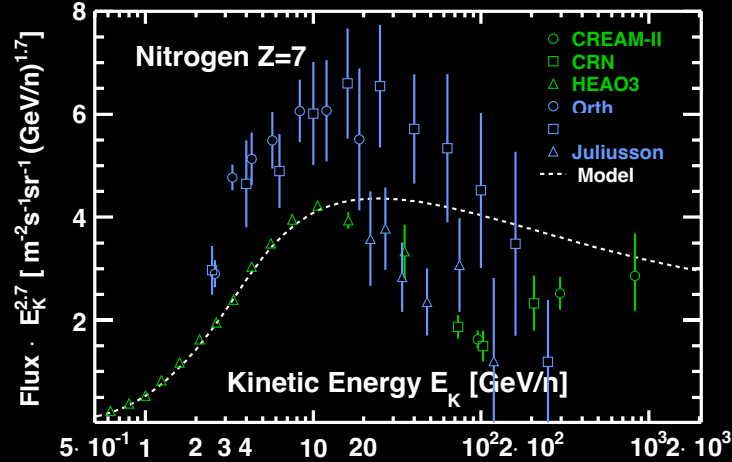
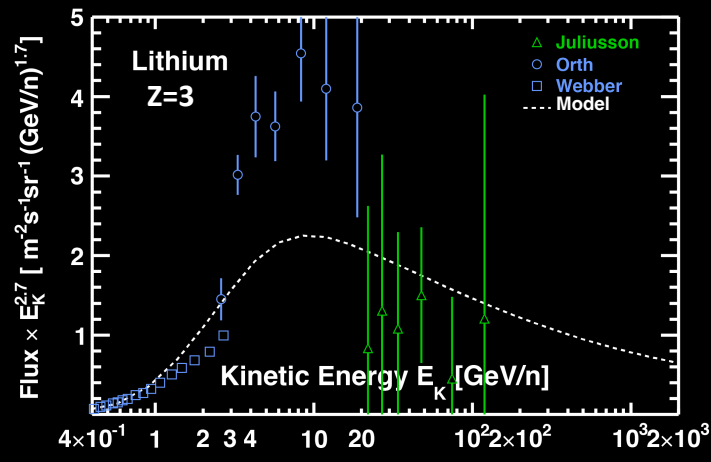
# Cosmic Rays in the last 100 years

1912:  
Discovery of Cosmic  
Rays



Victor Hess  
Nobel Prize (1936)

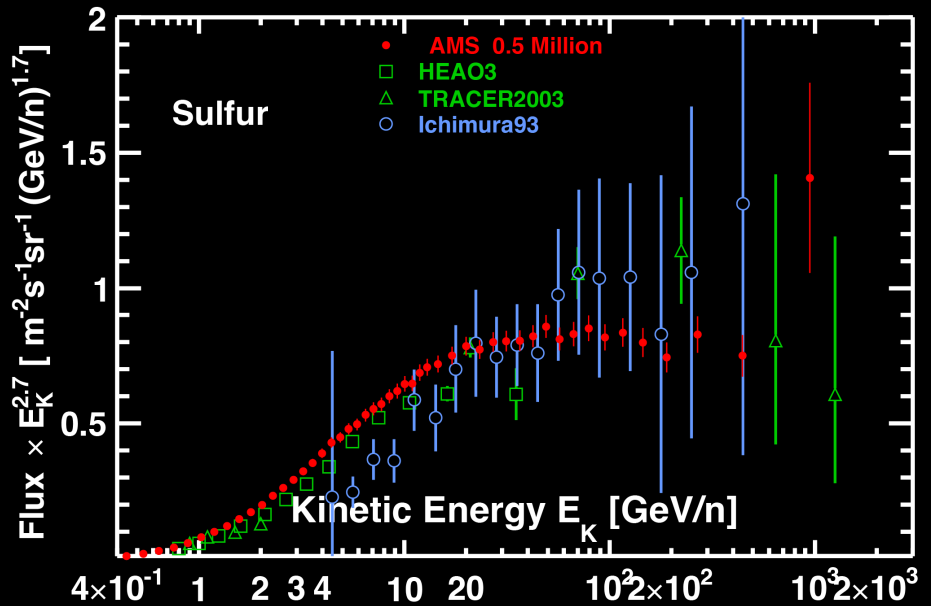
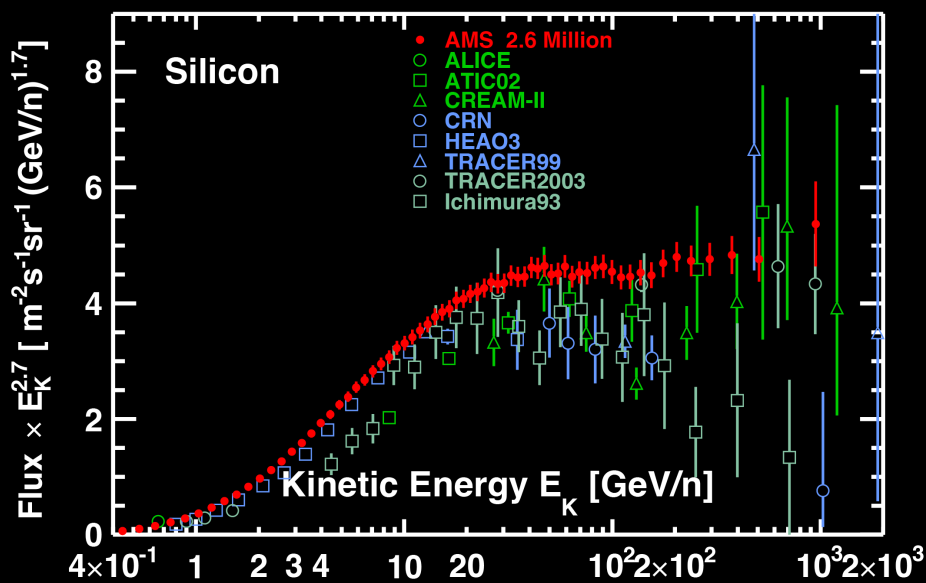
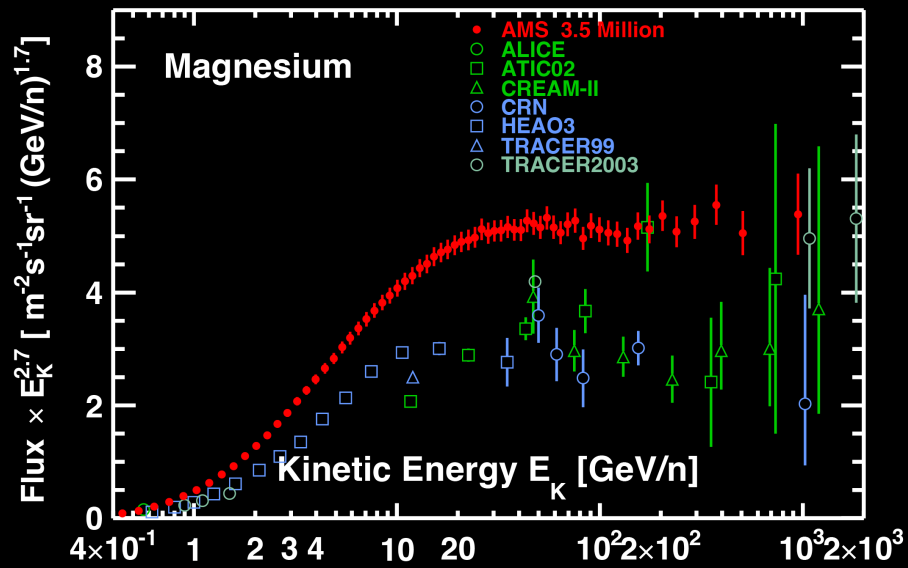
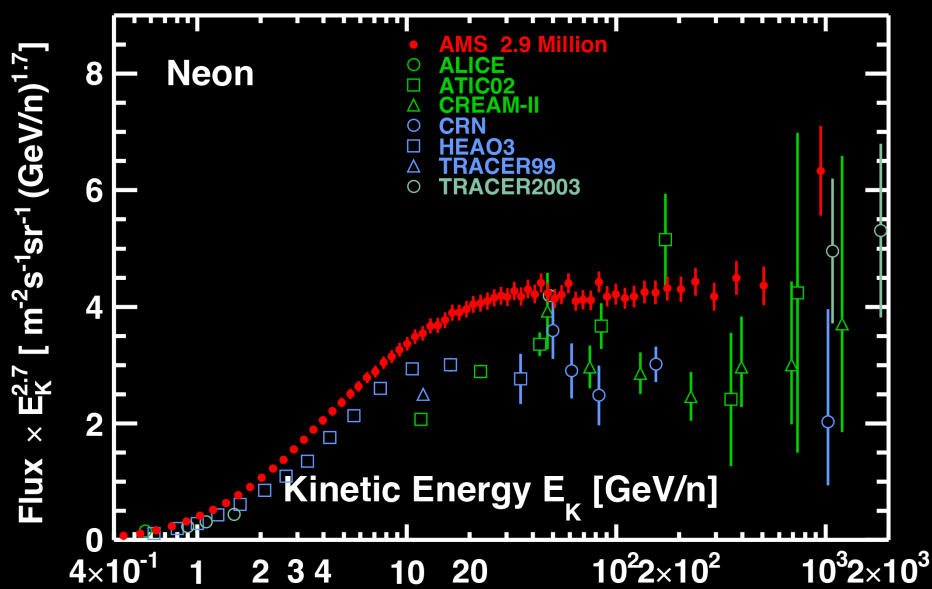
Before AMS:  
Many theoretical  
models agree with  
experimental data





# Examples of AMS Results compared with earlier measurements

## The precision AMS results cannot be explained by current models.



# AMS is a space version of a precision detector used at CERN accelerators

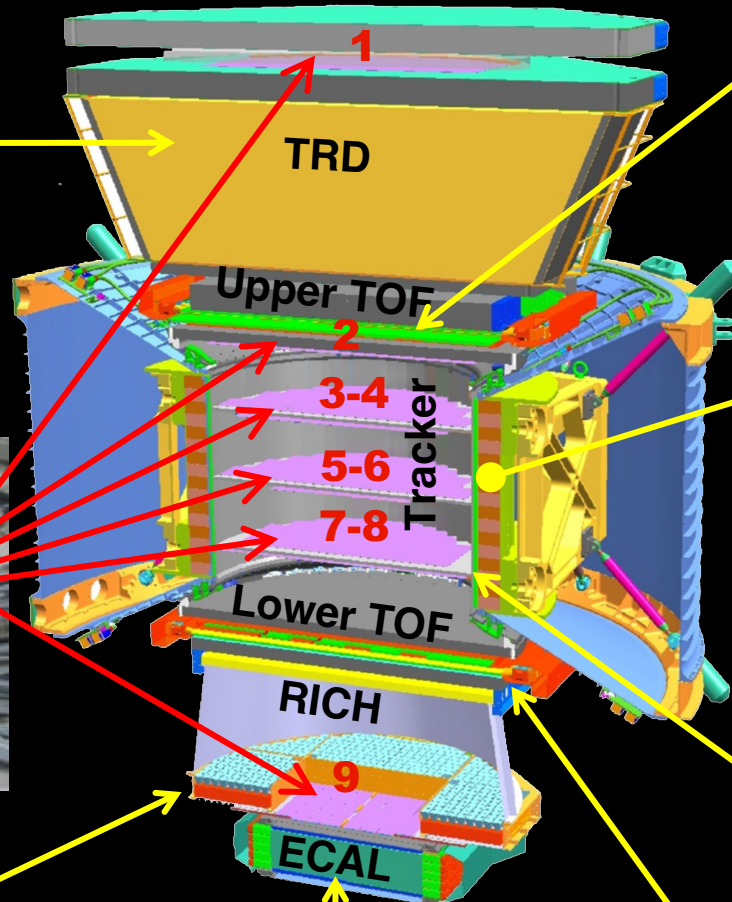
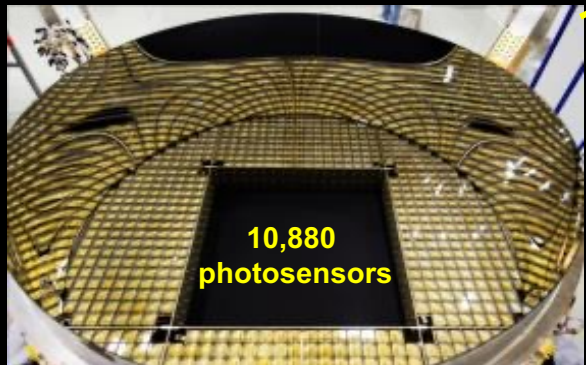
**Transition Radiation Detector (TRD)**  
identify  $e^+$ ,  $e^-$



**Silicon Tracker**  
measure Z, P



**Ring Imaging Cerenkov (RICH)**  
measure Z, E



**Upper TOF** measure Z, E



**Magnet** identify  $\pm Z, P$



**Anticoincidence Counters (ACC)**  
reject particles from the side



**Lower TOF** measure Z, E



**Electromagnetic Calorimeter (ECAL)**  
measure E of  $e^+$ ,  $e^-$



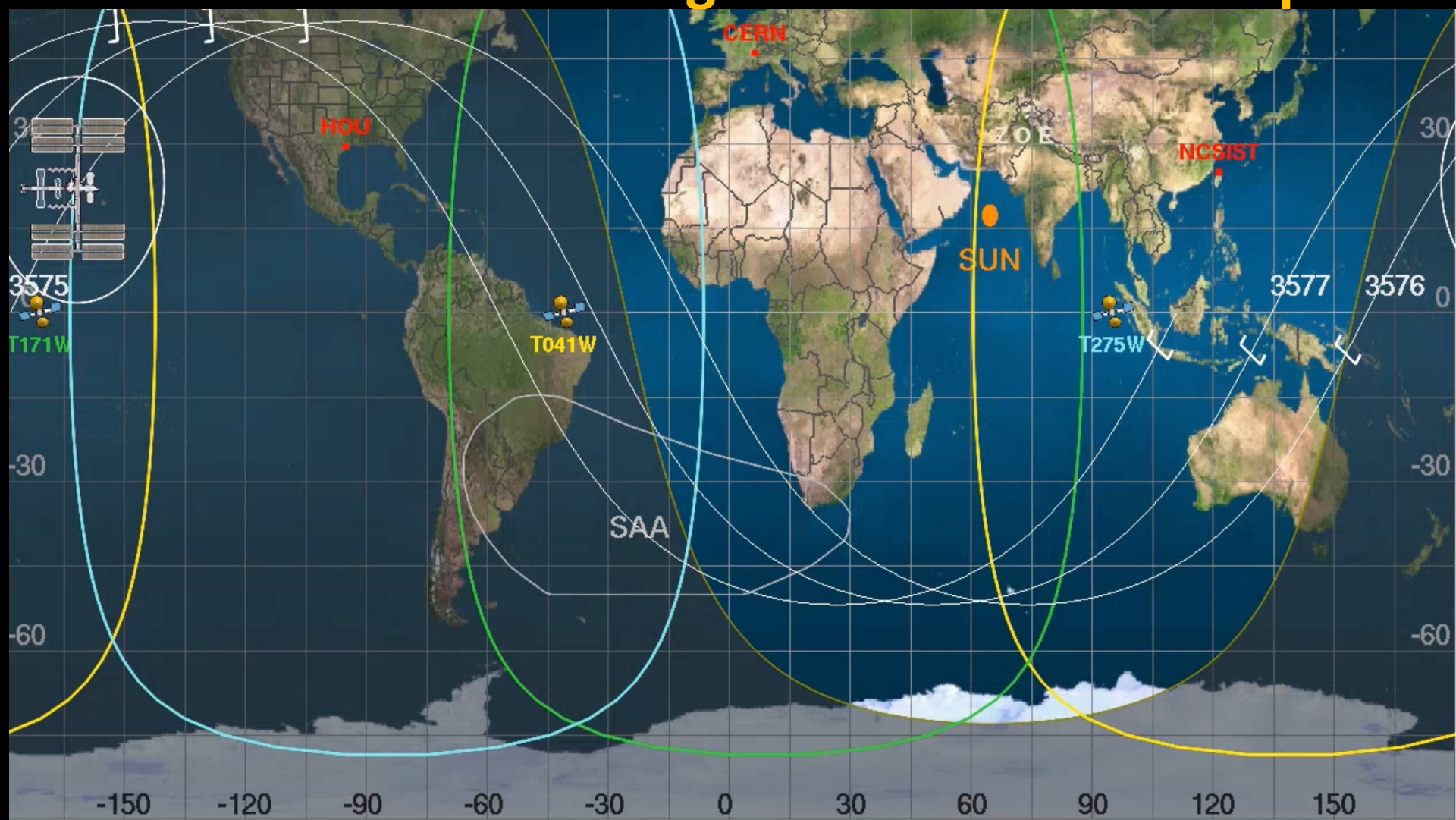


5m x 4m x 3m  
7.5 tons



**Endeavour approaching the Space Station, May 18, 2011**

# 2011-2023: AMS is taking data without interruption



220,182,212,893

# AMS Payload Operations Control Center at CERN



**Flight and Ground Operations**

**Detector Operations**

**Thermal Operations**

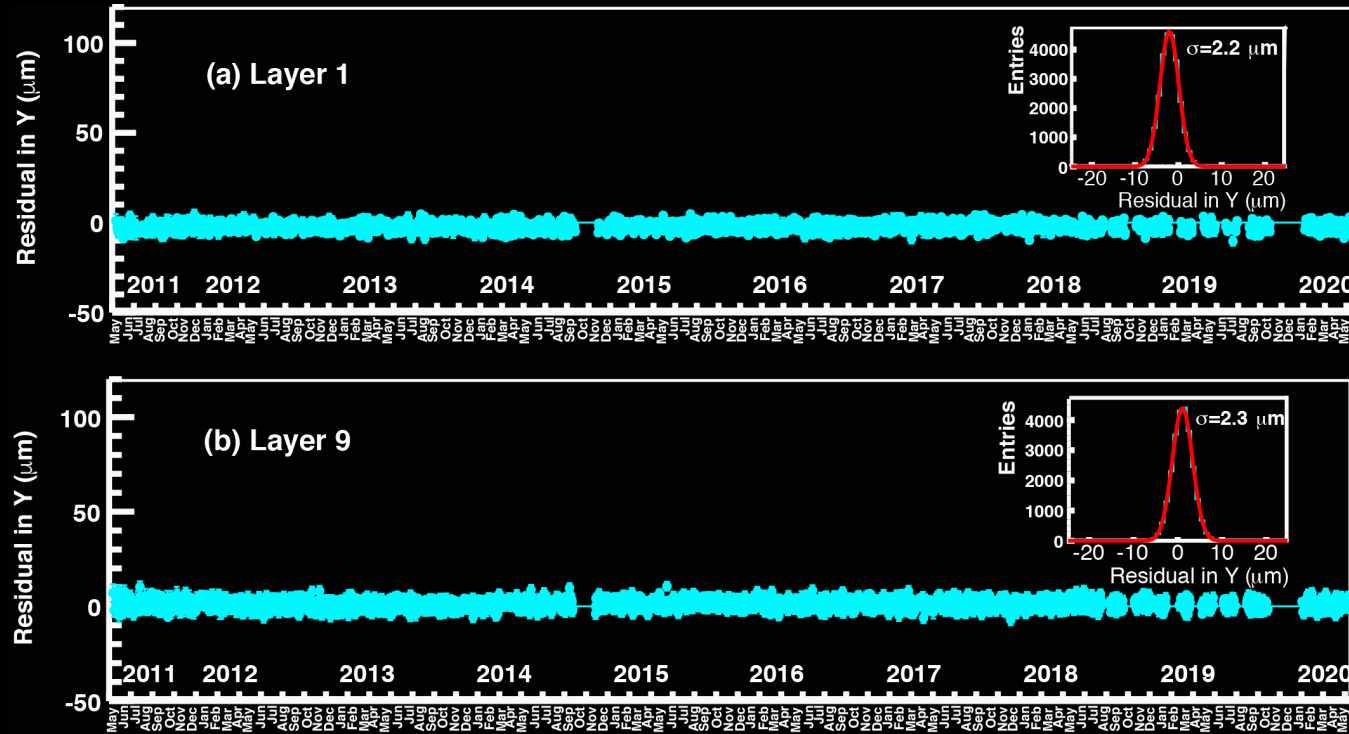
**Fully staffed 24 hours per day, 365 days per year**

# Continuous Verification of Detector Performance on Orbit

## Example: Tracker Alignment

Inner tracker alignment  
( $< 1$  micron)  
monitored with IR lasers

Outer layers monitored every 2 minutes by cosmic rays

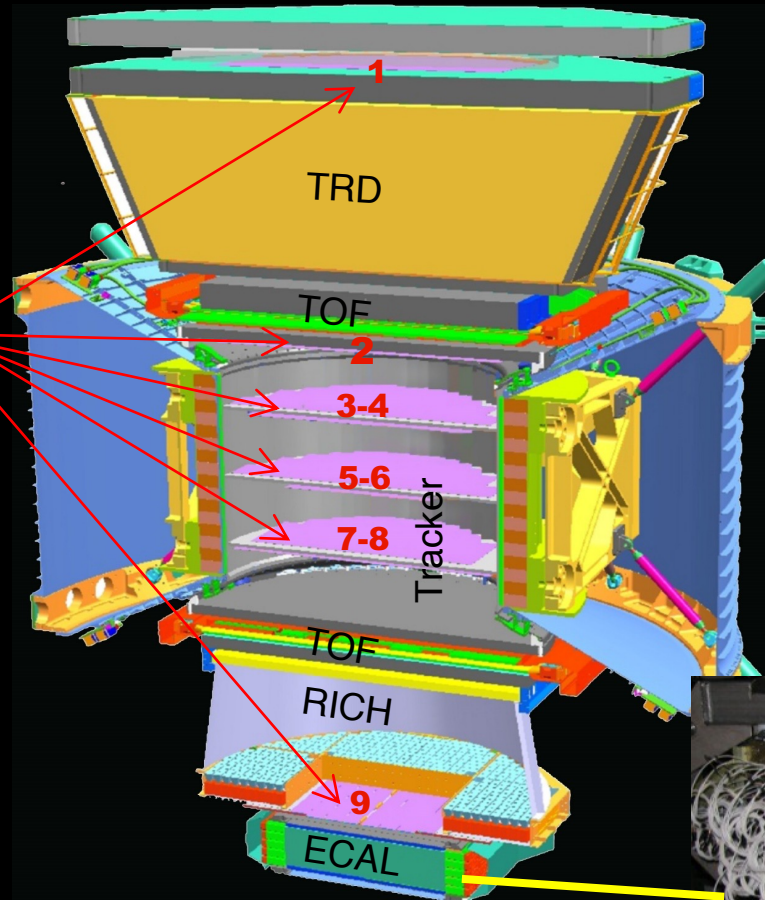


stable to 2 microns over 10 years

## Detector Improvements in space

With 12 years of operation and analysis of more than 220 billion events, we now know the details of the detector response and this has enabled us to make significant improvements in the measurements.

**Silicon Tracker**



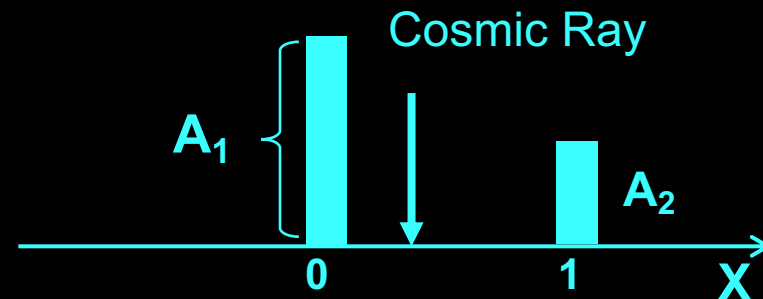
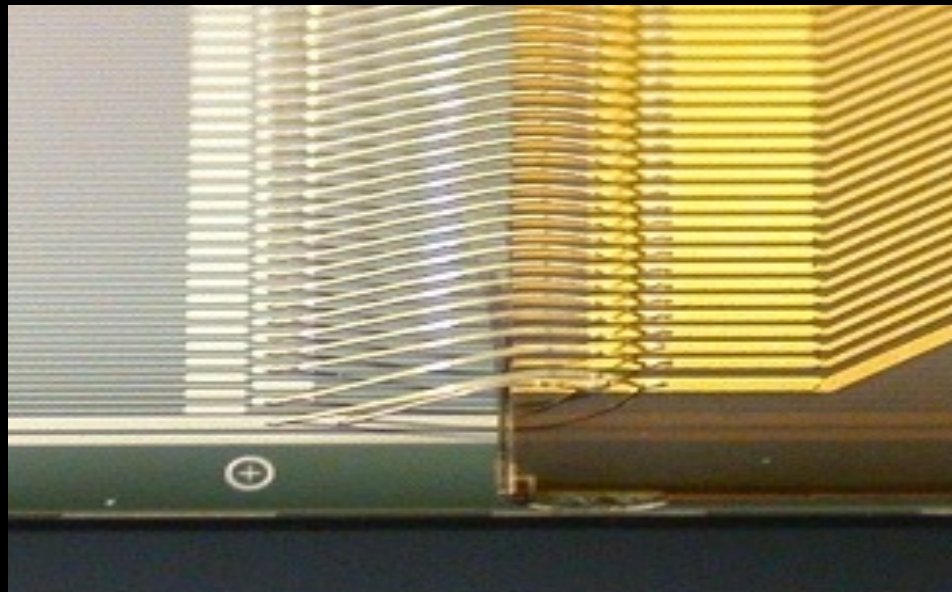
**ECAL**





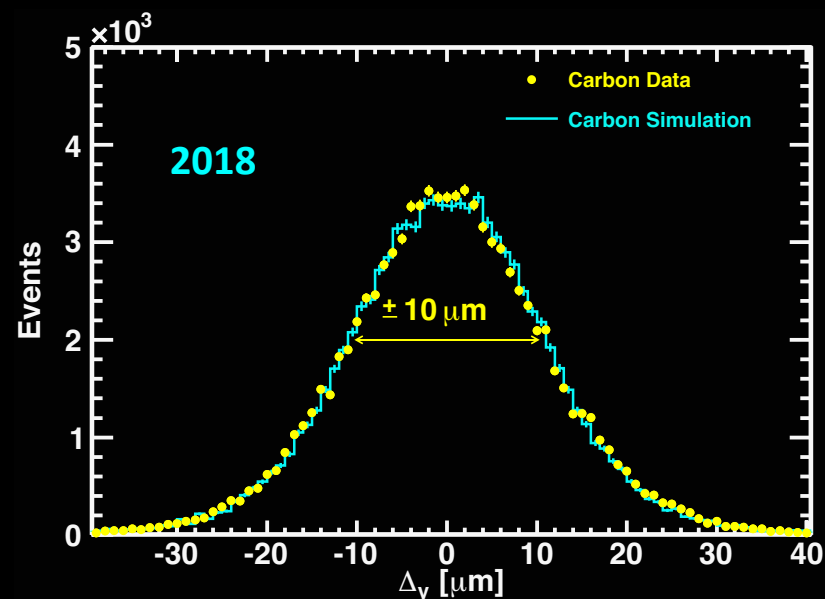
# Coordinate resolution improvement

The maximum amplitude is defined as  $A_1$  ( $x=0$ ) and the next largest adjacent strip as  $A_2$  ( $x=1$ ).



Traditional method: The  $A_1/A_2$  ratio provides the particle coordinate

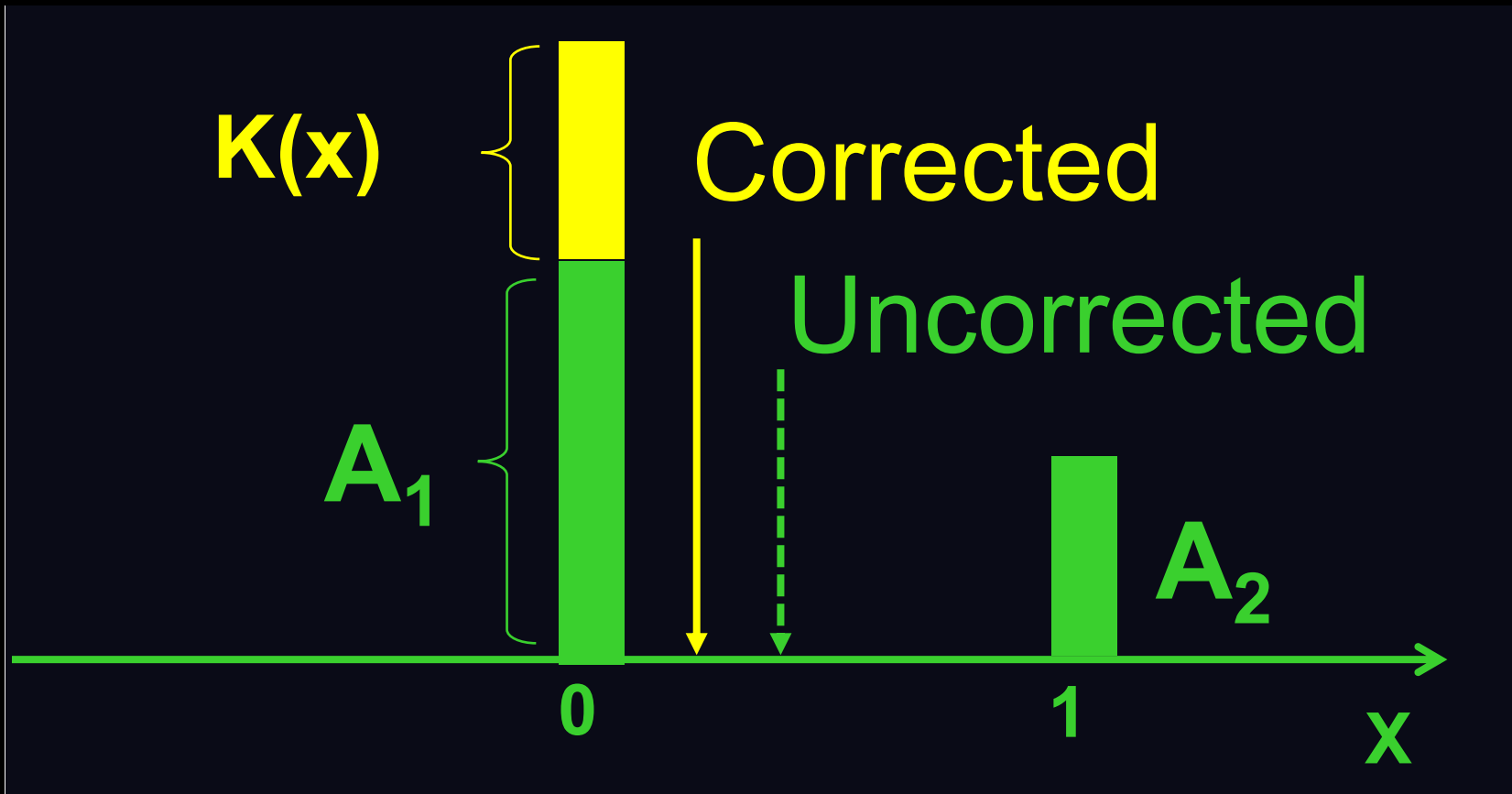
The signals should be proportional to  $Z^2$ .  
For  $Z > 3$ ,  $A_1$  starts to be non-linear,  
causing resolution degradation.



# Coordinate resolution improvement

To correct for the non-linear effect, we find a function  $K(x)$

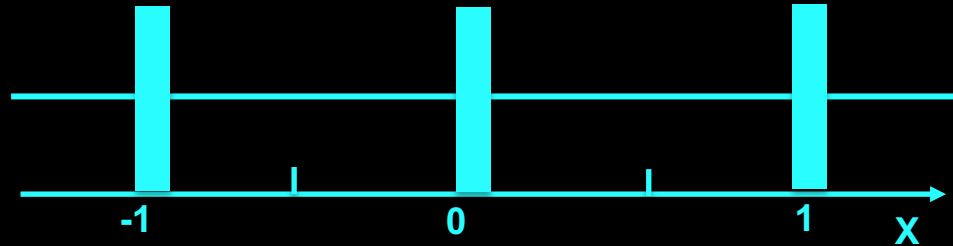
which will restore the non-linear amplitudes  $A_1 \rightarrow A_1 + K(x)$



# Coordinate resolution improvement

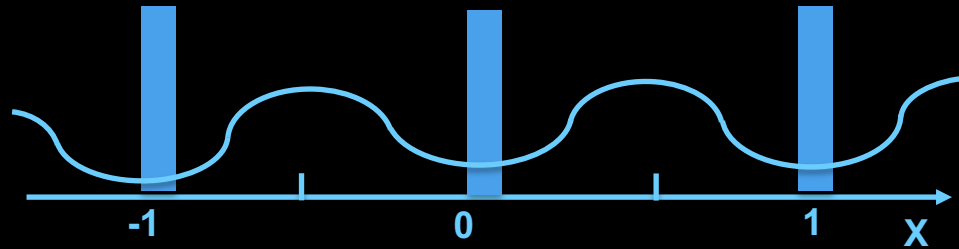
*Cosmic rays come uniformly from all directions*

Tracker linear response



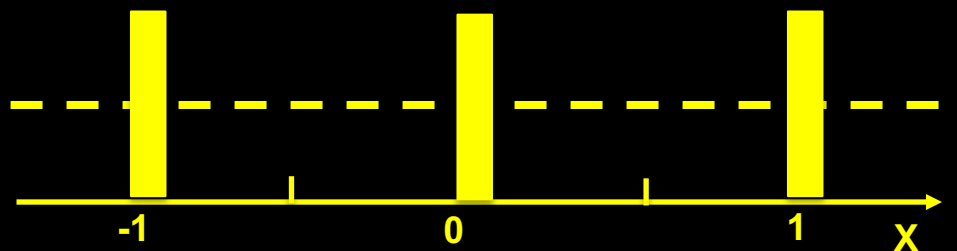
Uniform event position density

Tracker non-linear response

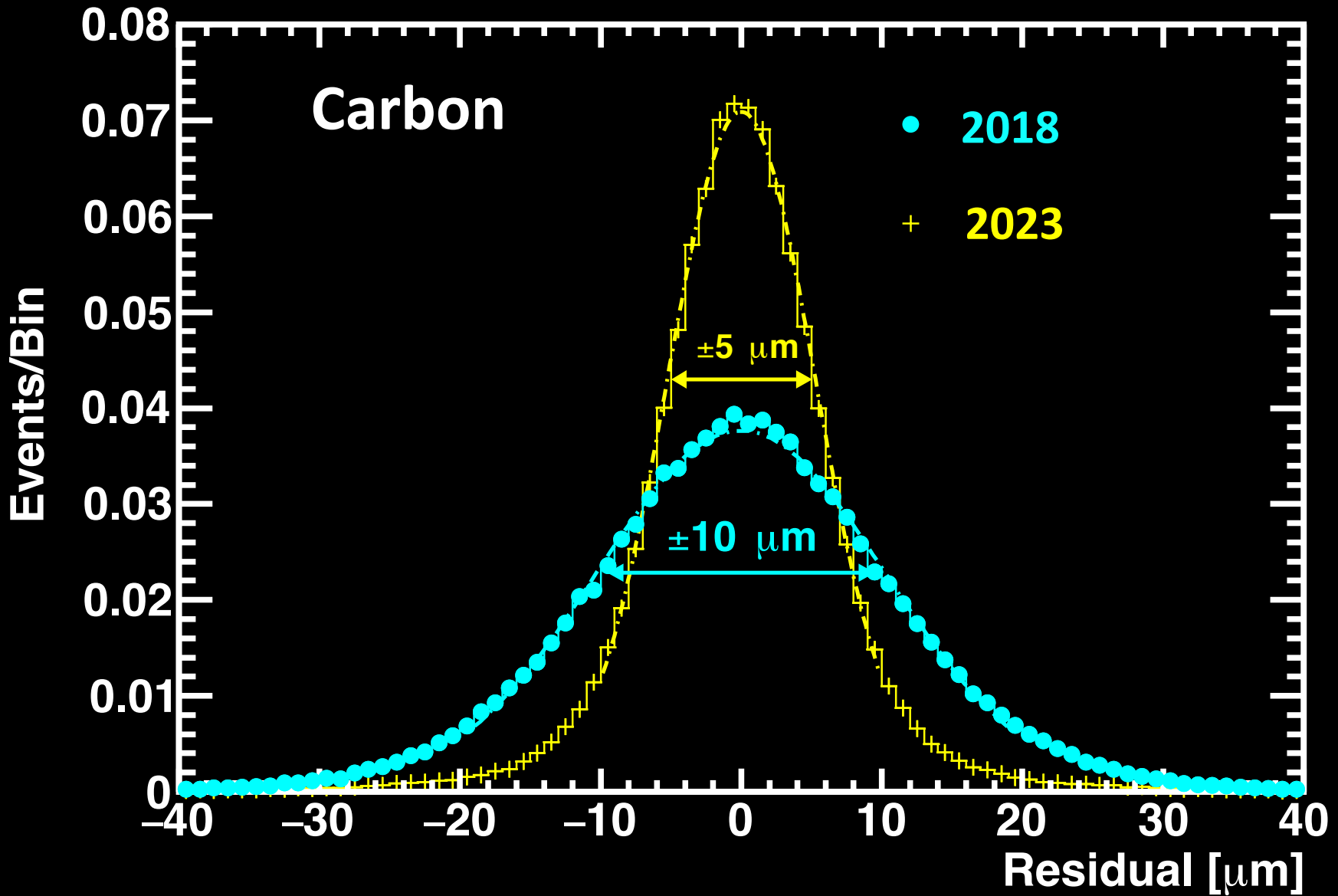


Event density not uniform

**We determine  $K(x)$  by requiring a uniform event density:**

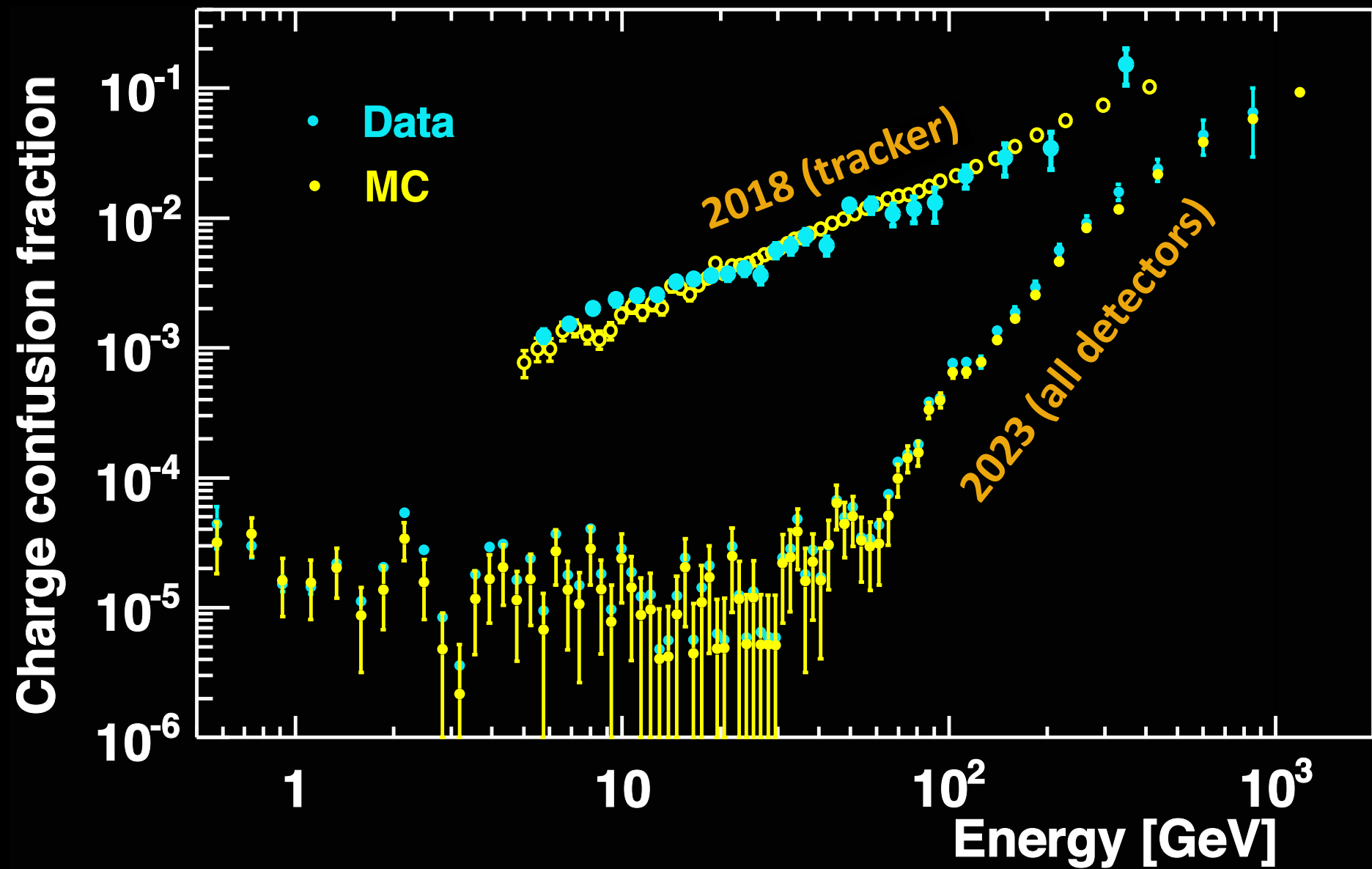


# Improvements in the coordinate resolution



# Improvements in the $e^\pm$ charge-sign determination from all detectors

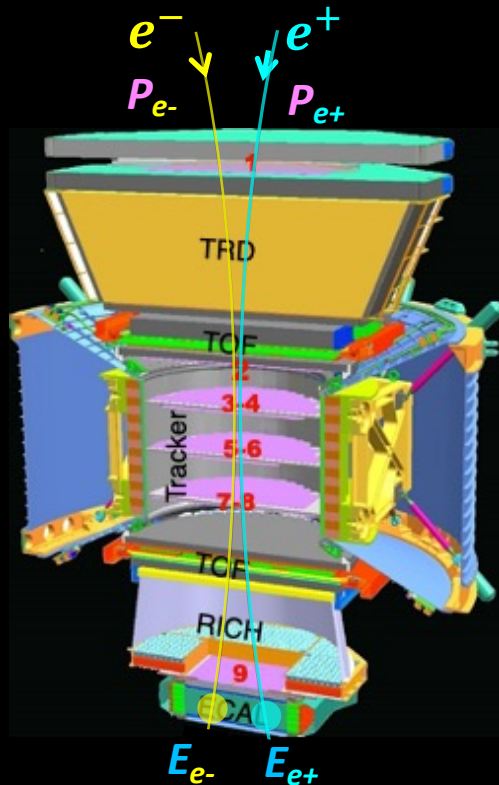
The charge confusion is reduced by a factor of 10 at 500 GeV



# Continuous Momentum Scale Verification (Unique Advantage of a Magnetic Spectrometer)

In AMS, the largest systematic error in the determination of the fluxes at the highest energies is due to the uncertainty in the absolute momentum scale.

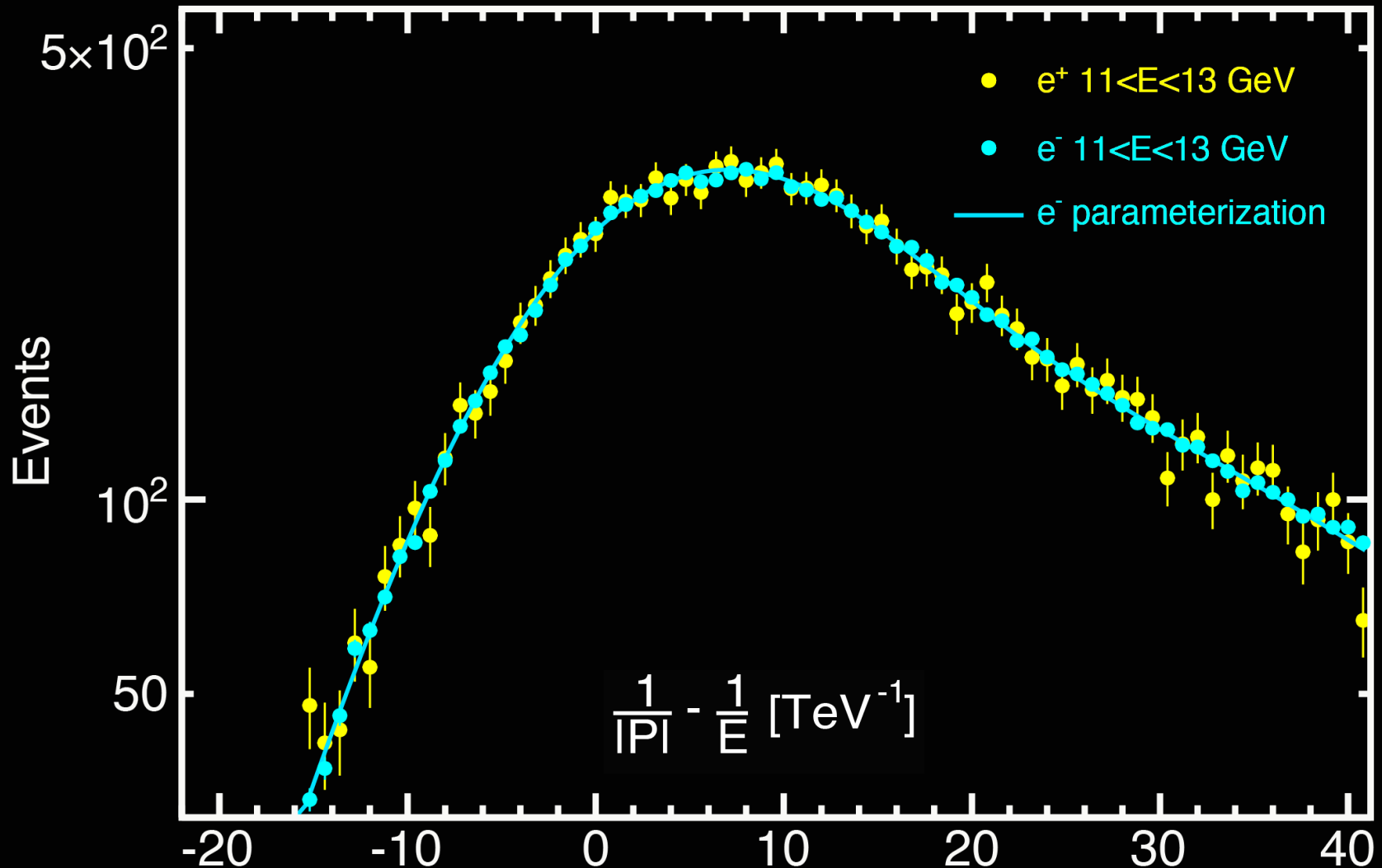
In space continuous outgassing of the supporting structure can affect the position of the tracker sensors at the sub-micron level.



A shift in the central tracker planes  
of 0.5 microns  
is sufficient to create  
a momentum shift of 10% at 1 TeV  
and bias flux measurements.

# Momentum Scale Verification

with 4 independent measurements of  $P_{e^+}$ ,  $E_{e^+}$ ,  $P_{e^-}$ ,  $E_{e^-}$



The accuracy of the momentum scale is determined to be  $1/(34,000 \text{ GeV})$ ;  
i.e., at 1 TeV the uncertainty is less than 3%

# Examples on the importance of the determination of energy scale:

In calorimeters (like CALET on the ISS) there is no way to know the absolute energy scale

CALET, PRL [126](#), 241101 (2021)

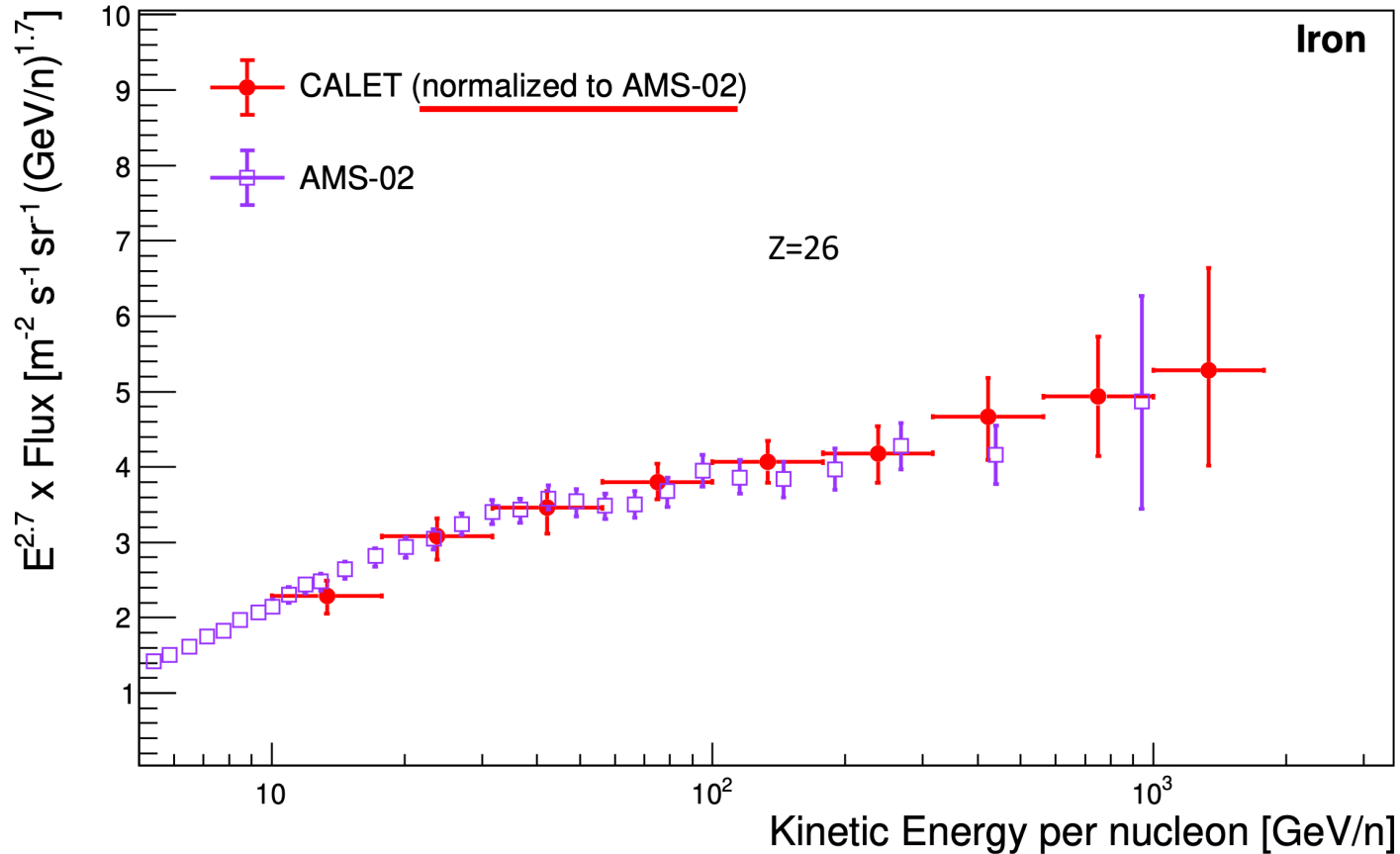


FIG. S12. Iron flux (with multiplicative factor  $E^{2.7}$ ) measured by CALET (red points) with 4 bins/decade, multiplied by 1.20 for comparison with the AMS-02 results [S3]. The error bars of the CALET data are the quadrature sum of statistical and systematic uncertainties.



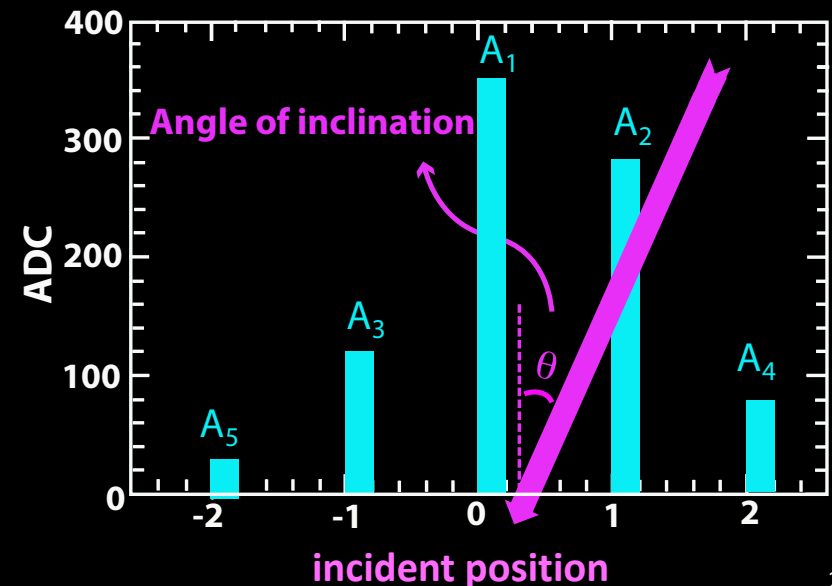
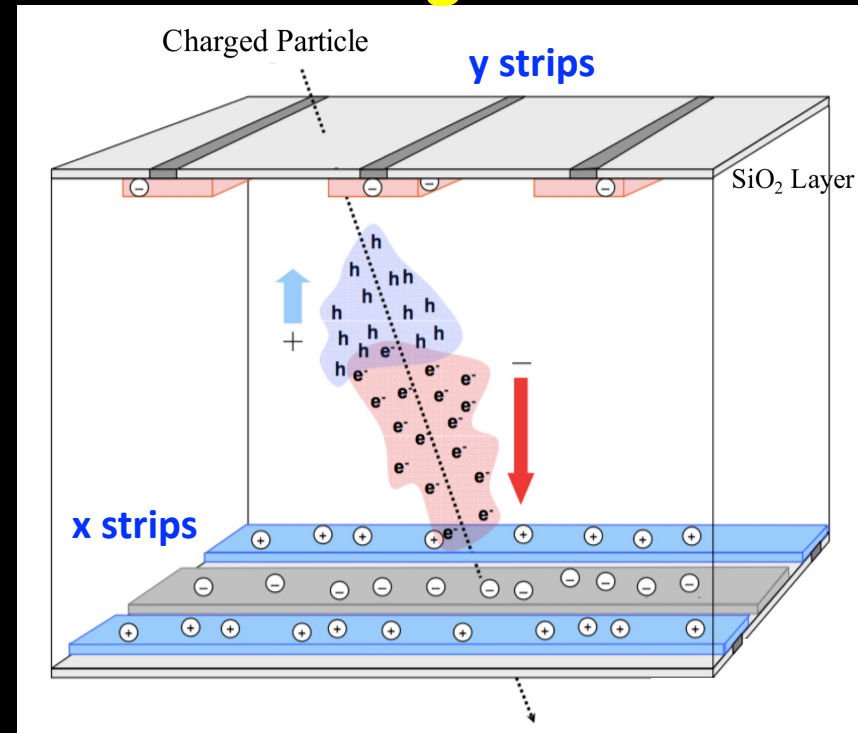
# Charge Resolution improvement for high-Z

The Tracker measures the charge  $|Z|$  with both the x- and y-strips.

The energy loss deposition is collected by several strips on each side.

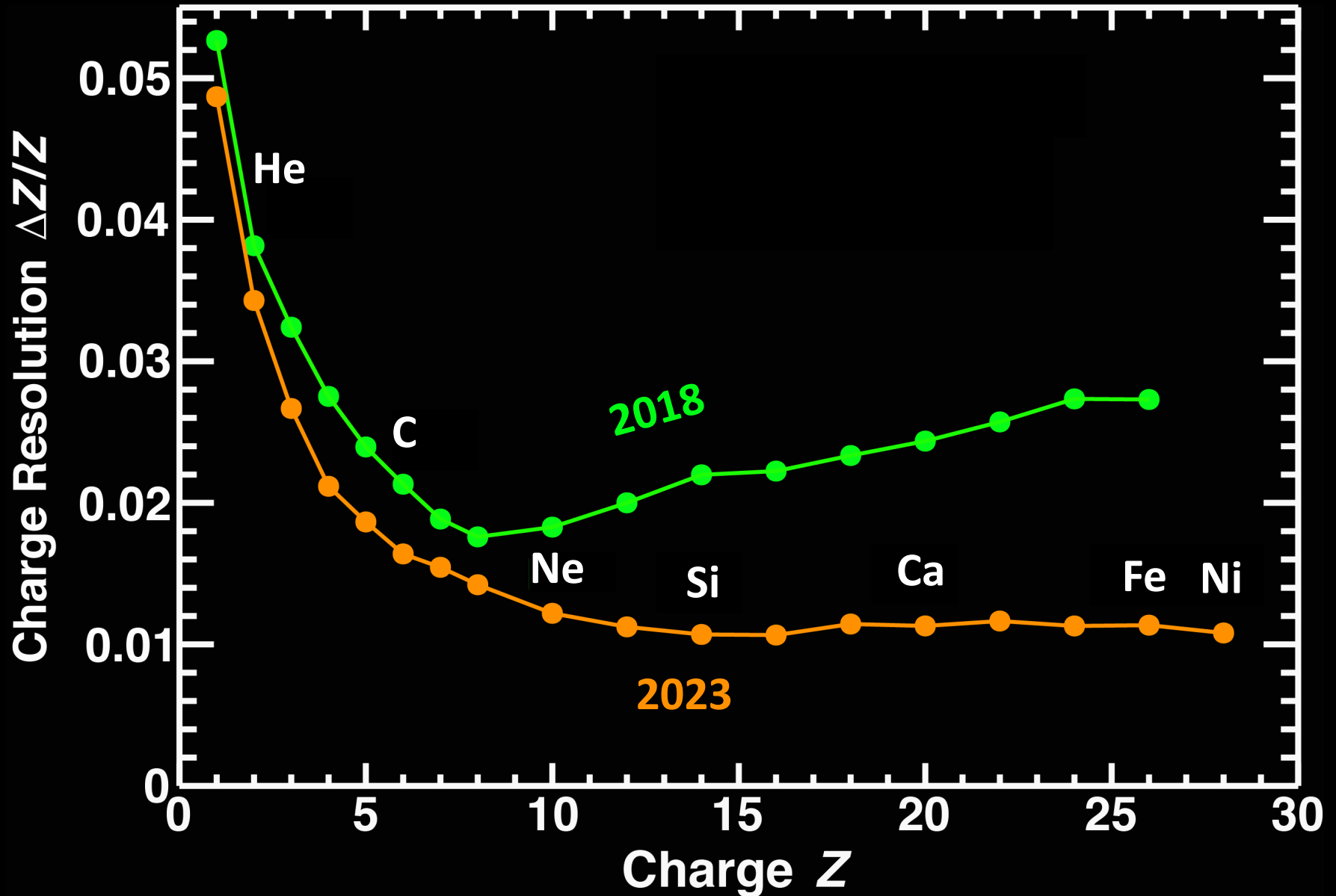
The amplitudes are related to the incident position and angle  $\theta$ .

The improvement corrects for the saturation strip-by-strip.

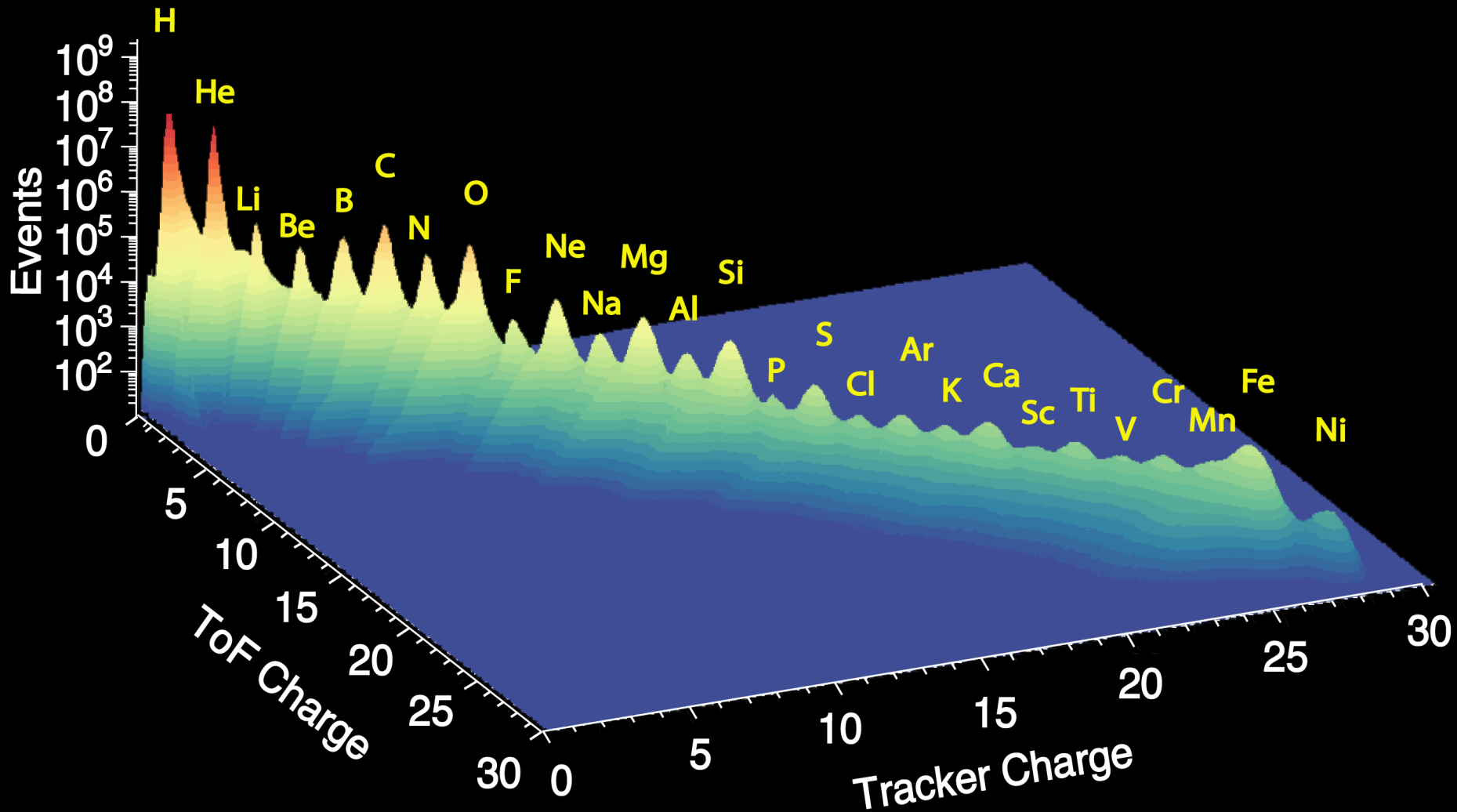


# Improvement in the charge resolution

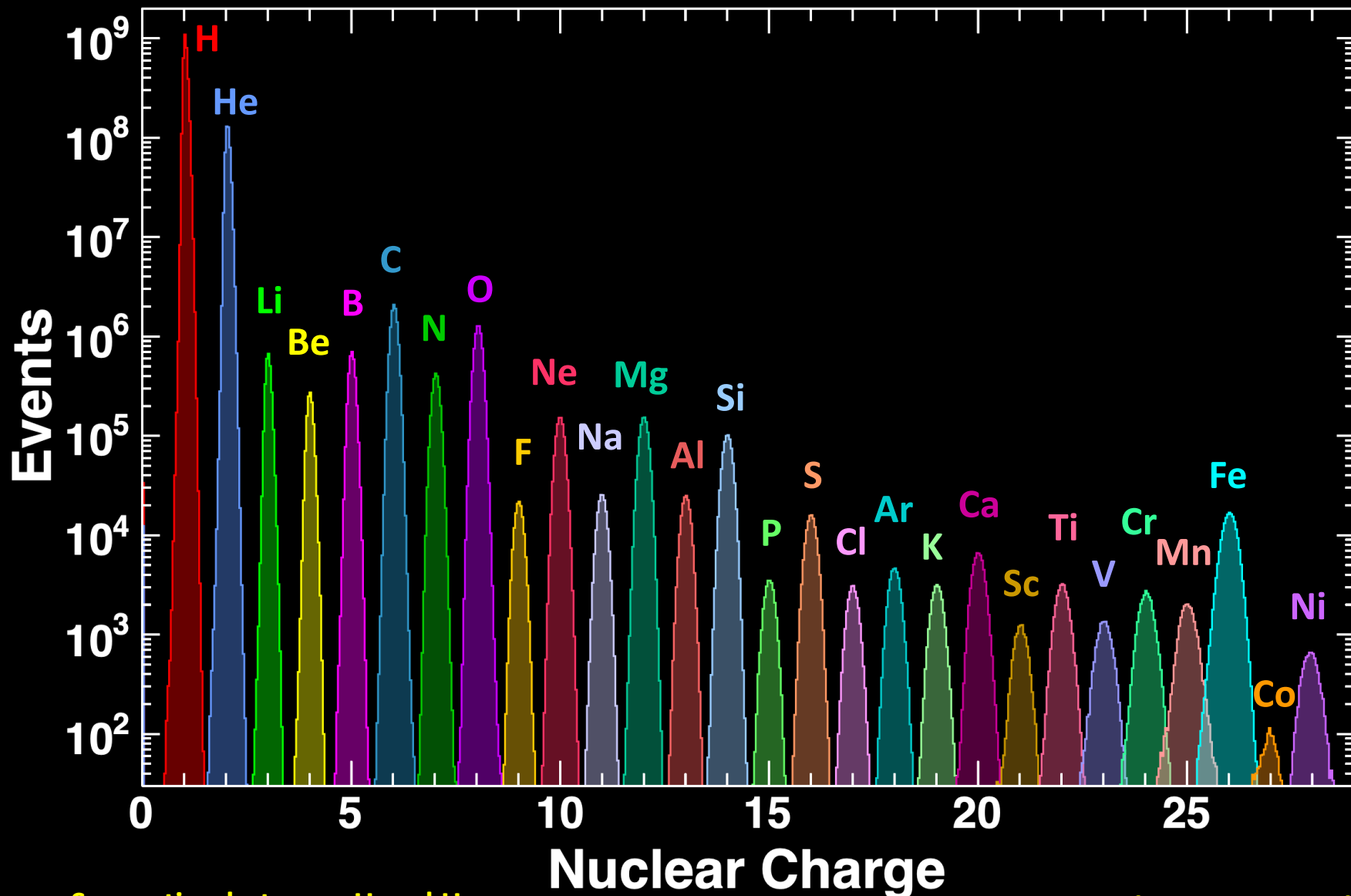
Improved by 200% and more



# Charge Measurement in 2018



# Charge Measurement in 2023

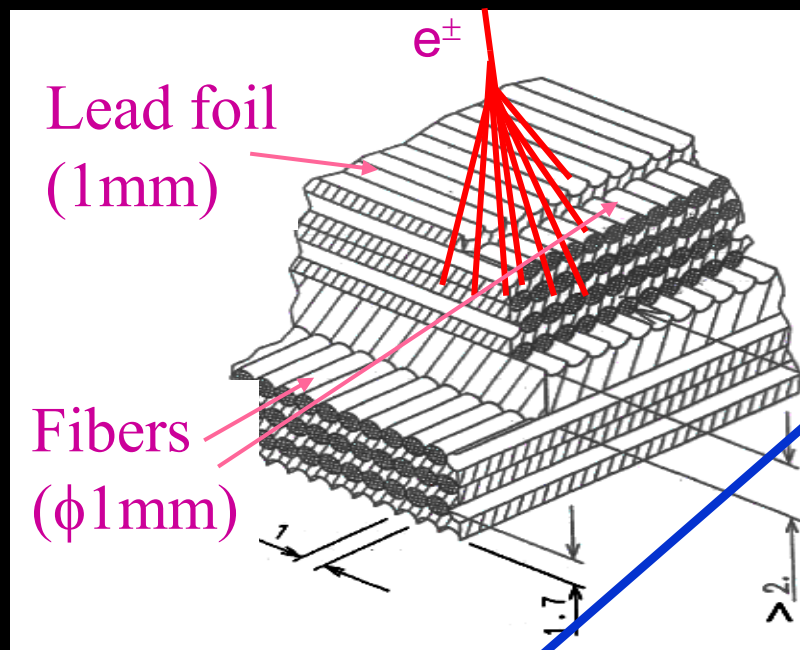


Separation between H and He  
is better than 1 in  $10^9$

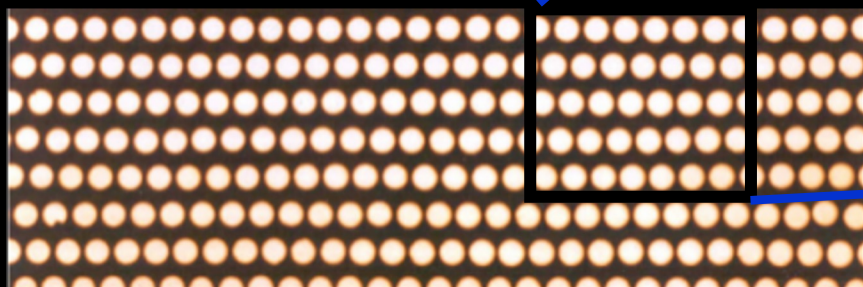
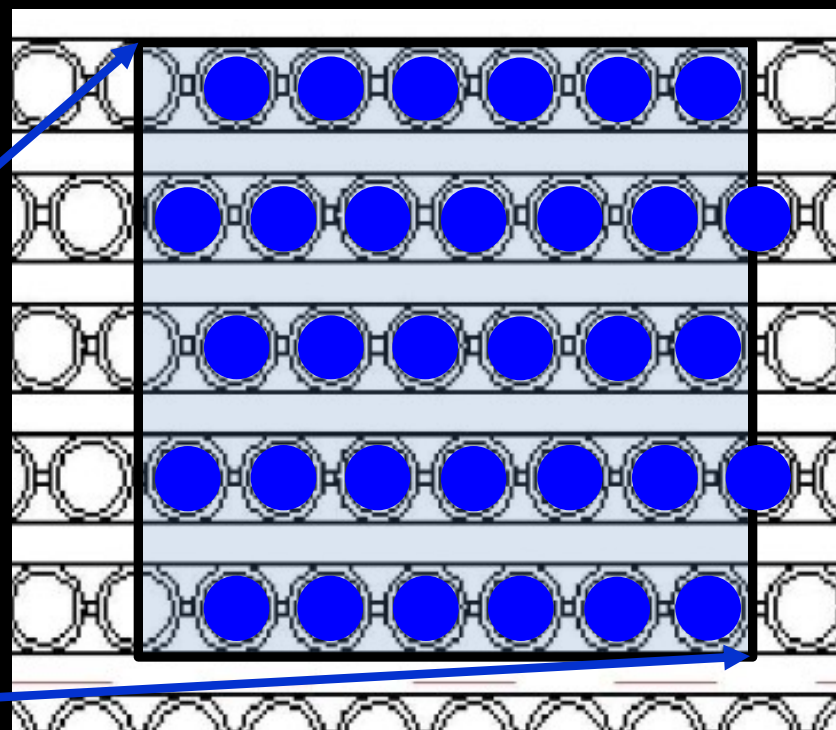
Separation between Fe and Co  
is better than 1 in  $10^2$

# Improvement in Electromagnetic Calorimeter (ECAL)

**3-D measurement** over 17 radiation lengths  
of the directions and energies of electrons and positrons

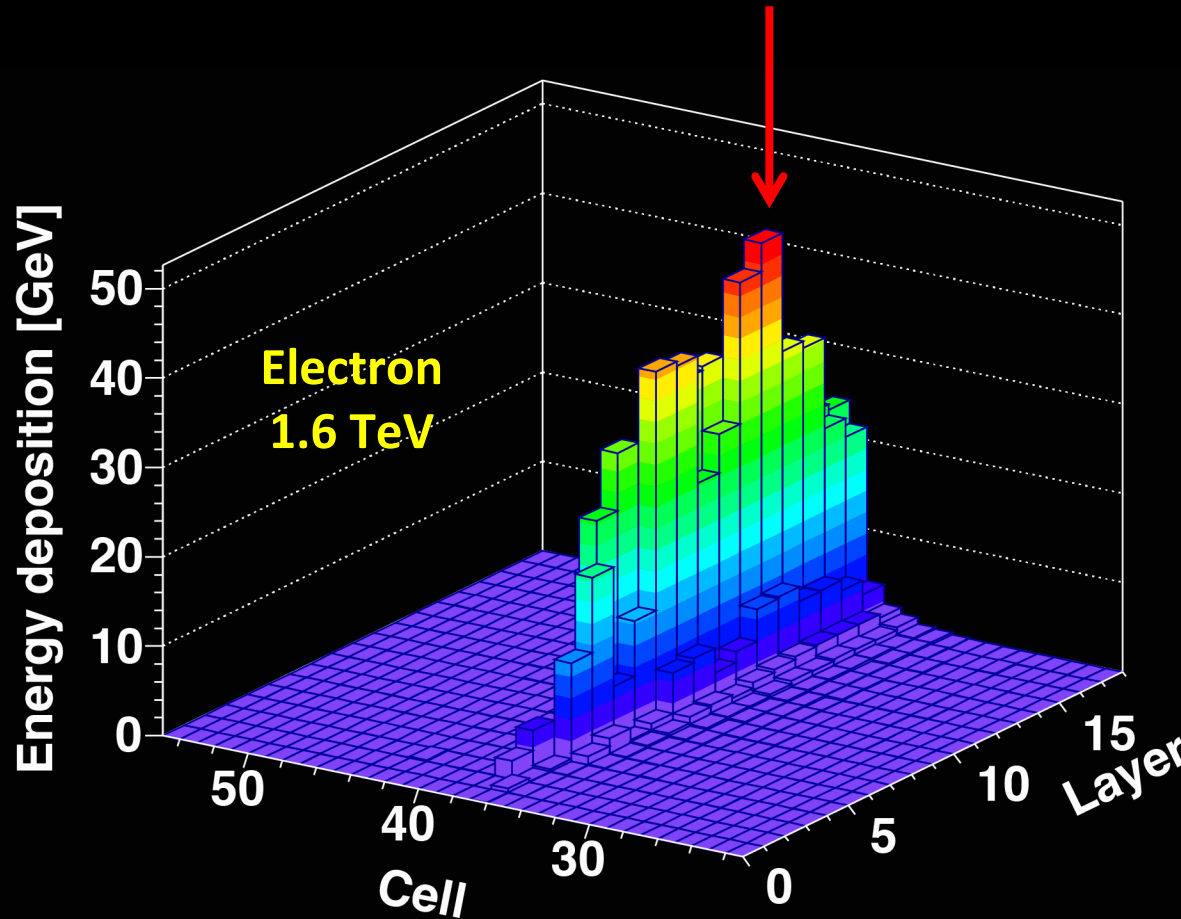


One of 1296 cells ( $9 \times 9 \text{ mm}^2$ )



# High energy measurements for positrons and electrons

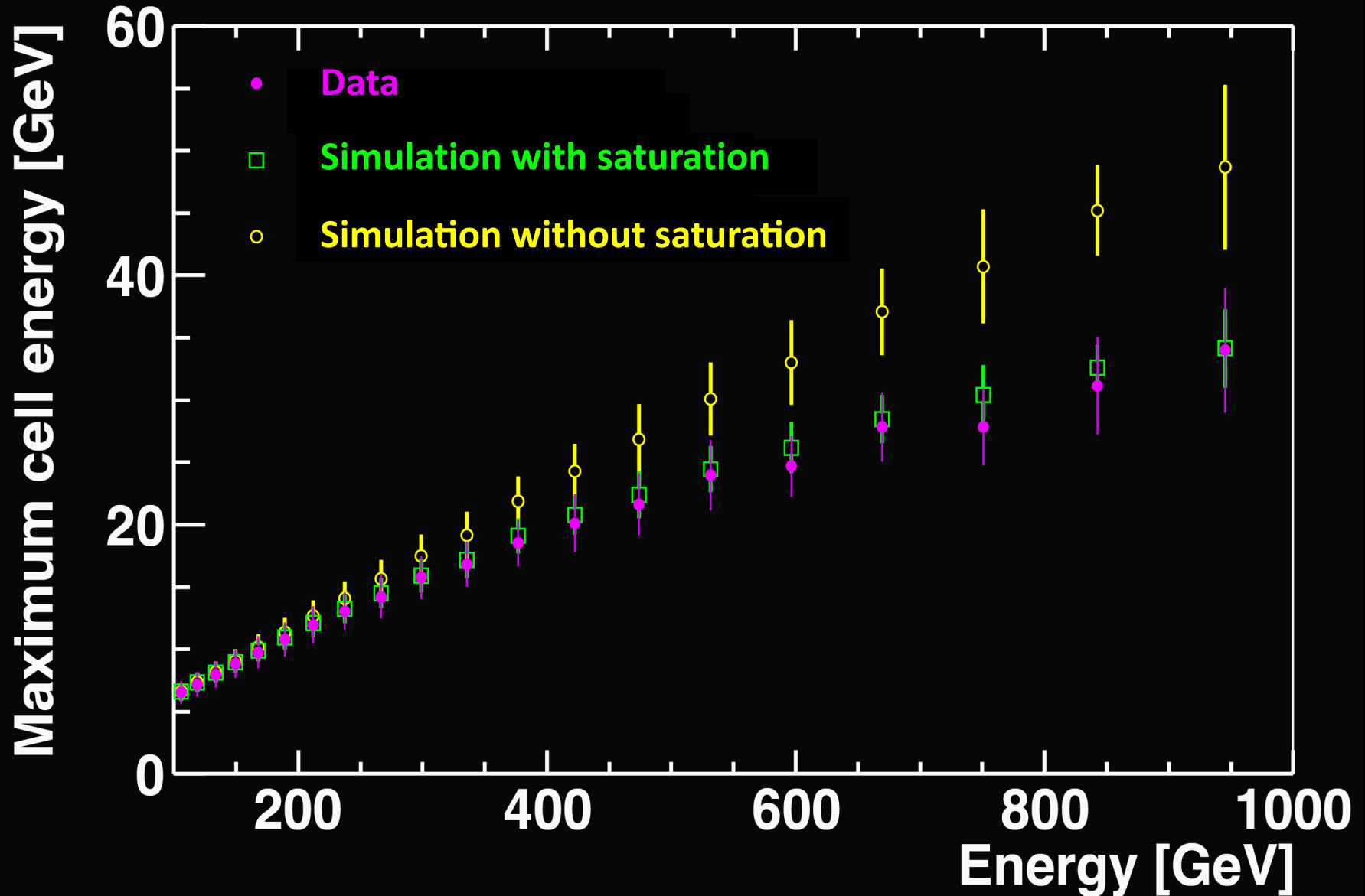
For the highest energy  $e^\pm$ , the response in cells in the center of the shower is saturated in the fibers.



The rest of the cells carry all the information about the shower

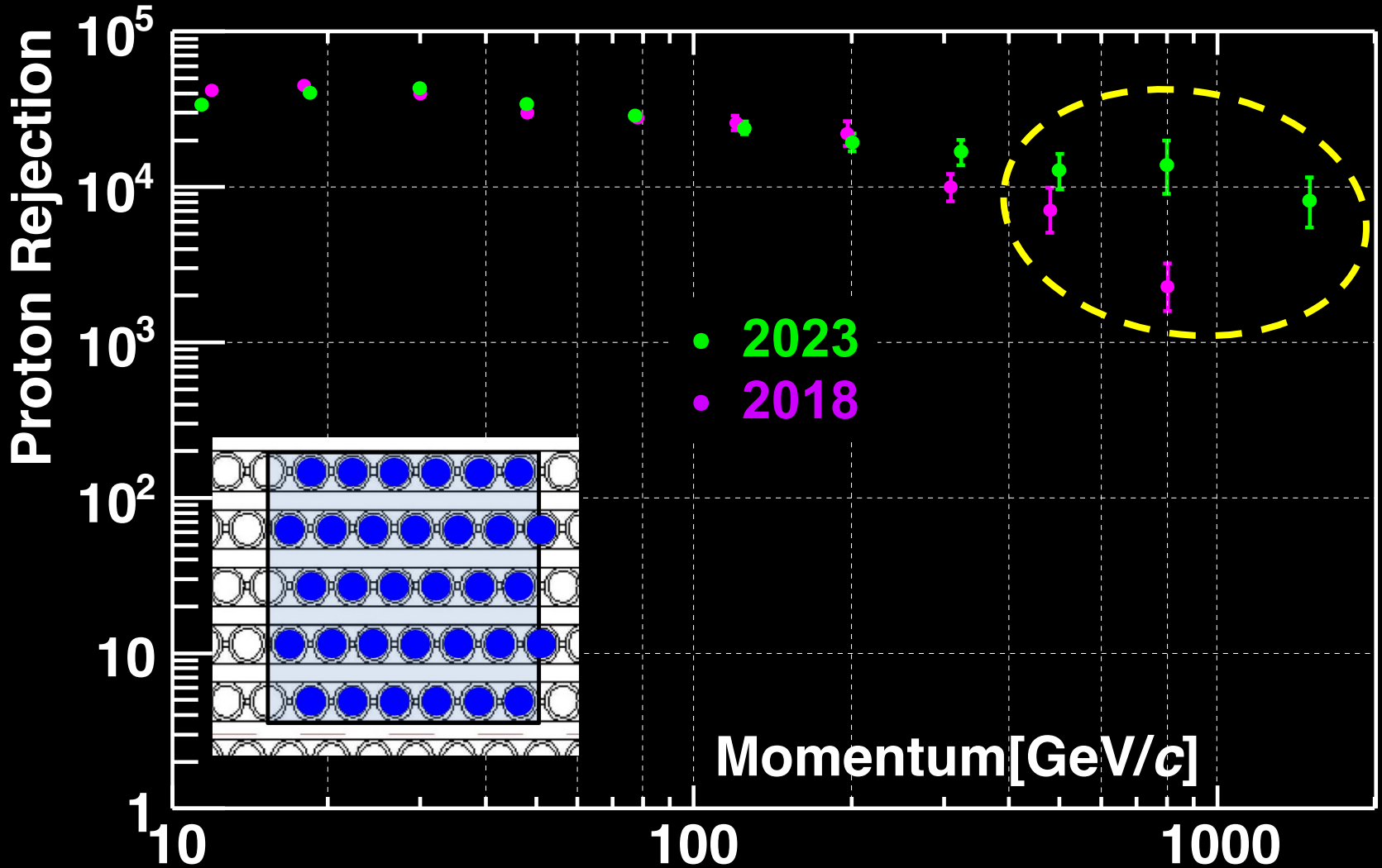
In space the rate is low. We have developed a technique that uses only these surrounding cells to reconstruct the overall shower, its energy and direction.

# Improvement in high-energy measurements for $e^\pm$



The total energy correction for 1 TeV Electromagnetic showers is 8%

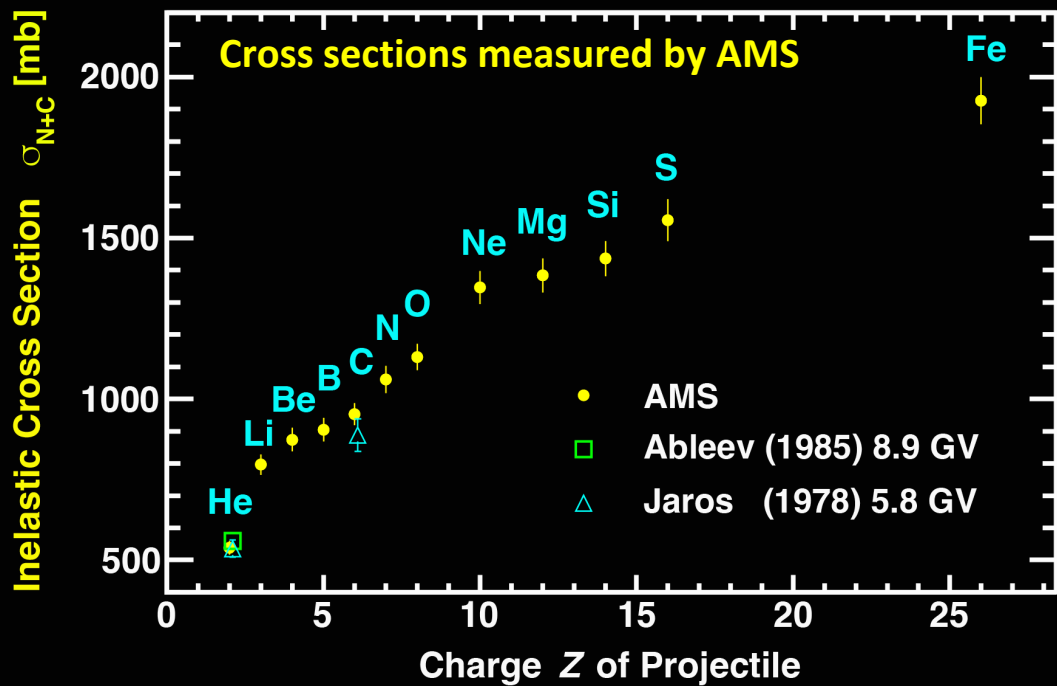
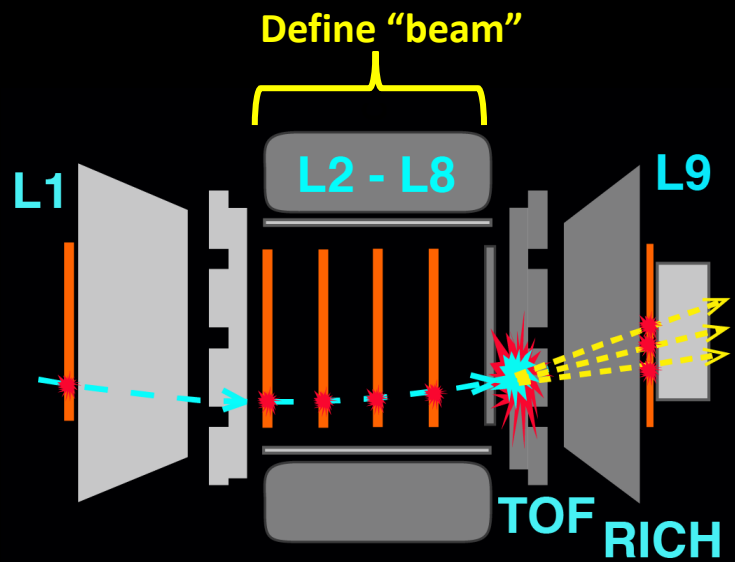
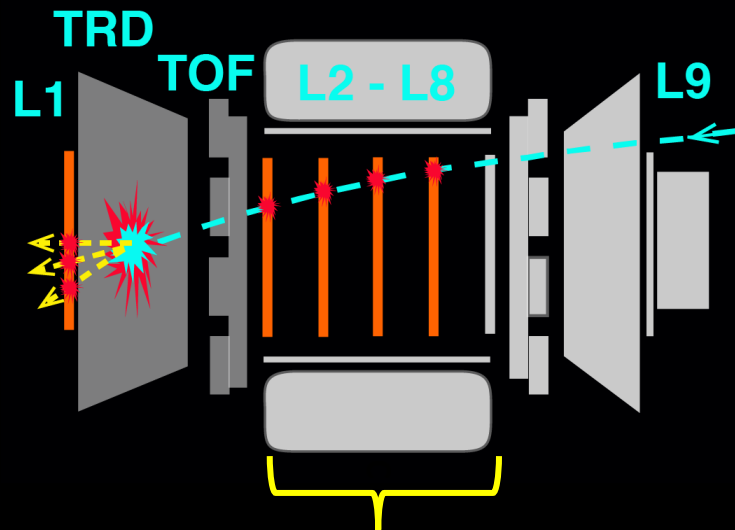
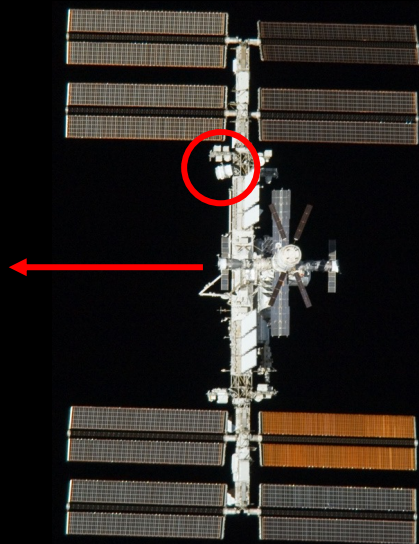
# Improvement in the high-energy proton rejection by modeling energy deposition in each fiber of the ECAL



At 1 TeV, proton rejection power improved by a factor of 5.  
This enables AMS to measure positrons and electrons to 2 TeV.



# Precision measurements of interactions in the detector for accurate flux determination



# AMS on ISS

AMS 2011-2025

Continuous data-taking

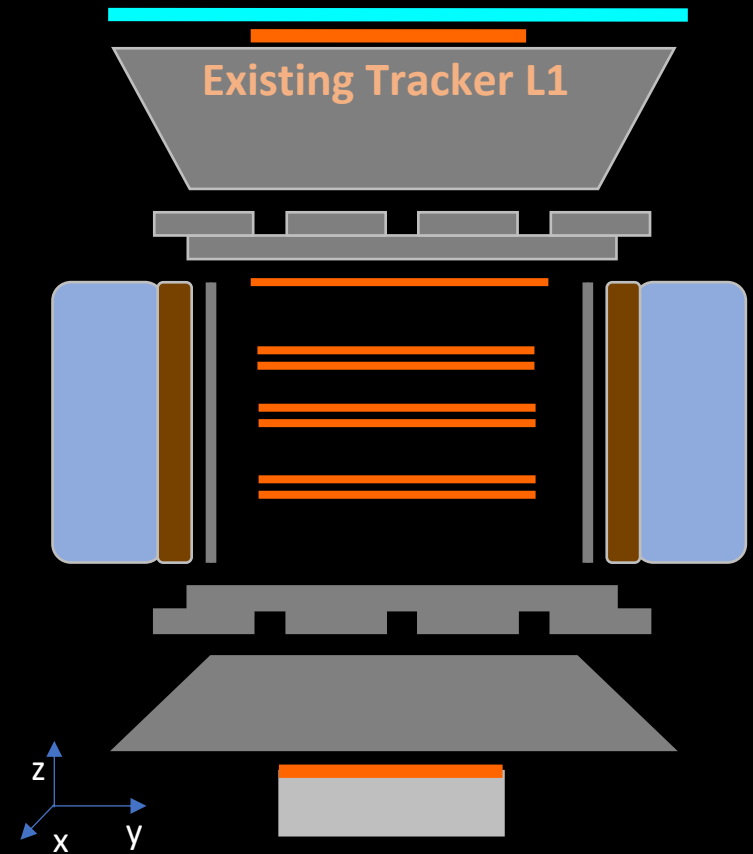


Latest Results: 2011-2022

AMS 2025-2030

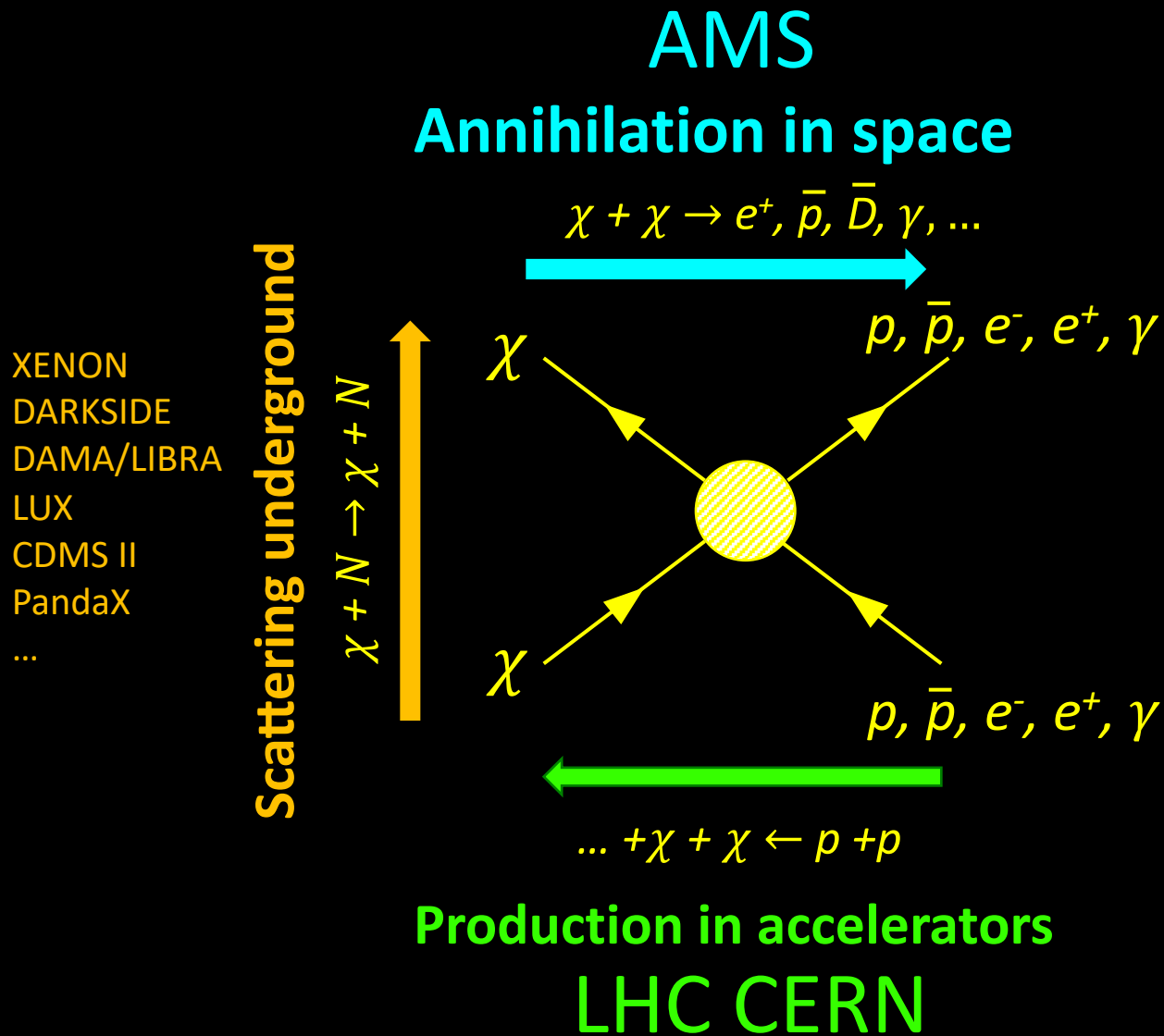
New 8m<sup>2</sup> Silicon Tracker Layer

Acceptance increased to 300%



and Projections

# Three independent methods to search for Dark Matter $\chi$



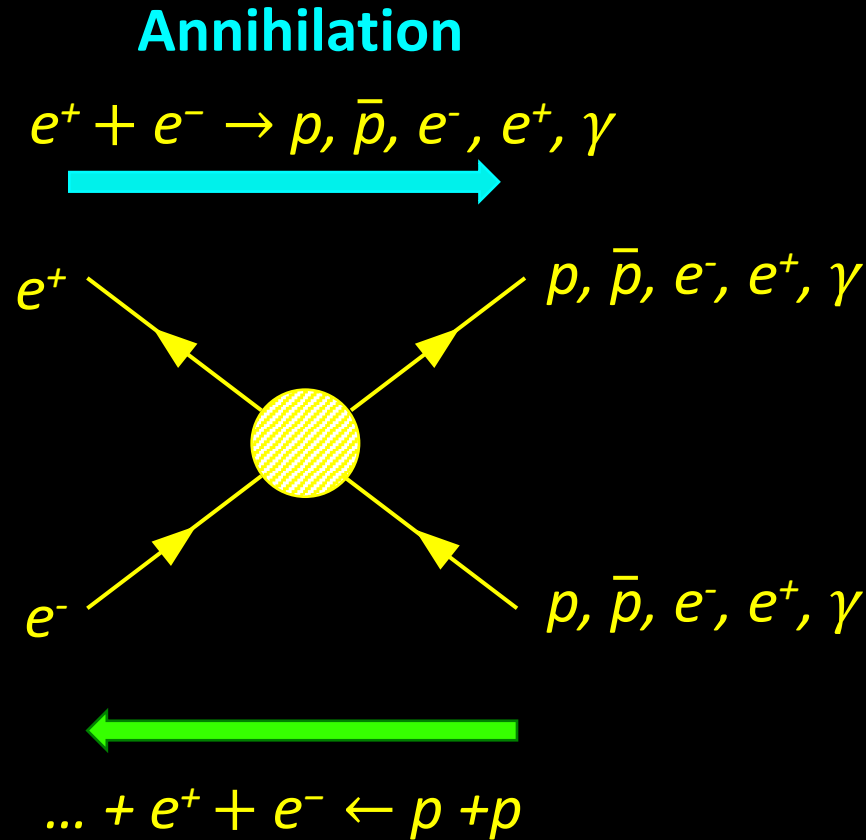
# Physics of electrons and protons

SPEAR, DORIS, PEP, PETRA, LEP, ...  $\psi, \tau$

SLAC ... partons, electroweak

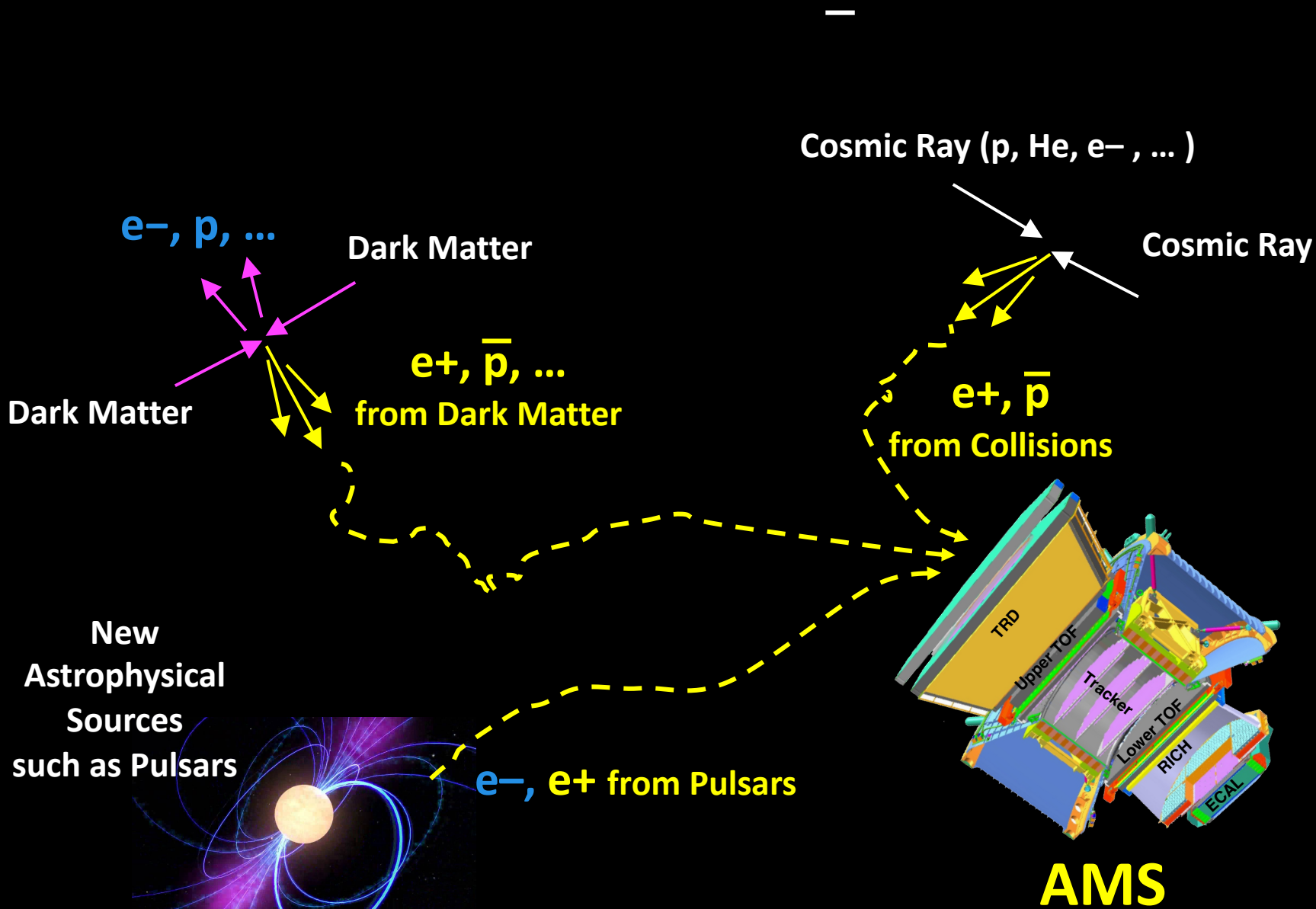
Scattering

$e + p \rightarrow p, \bar{p}, e^-, e^+, \gamma$



BNL, FNAL, LHC ...  $J, Y, T, Z, W, h^0$

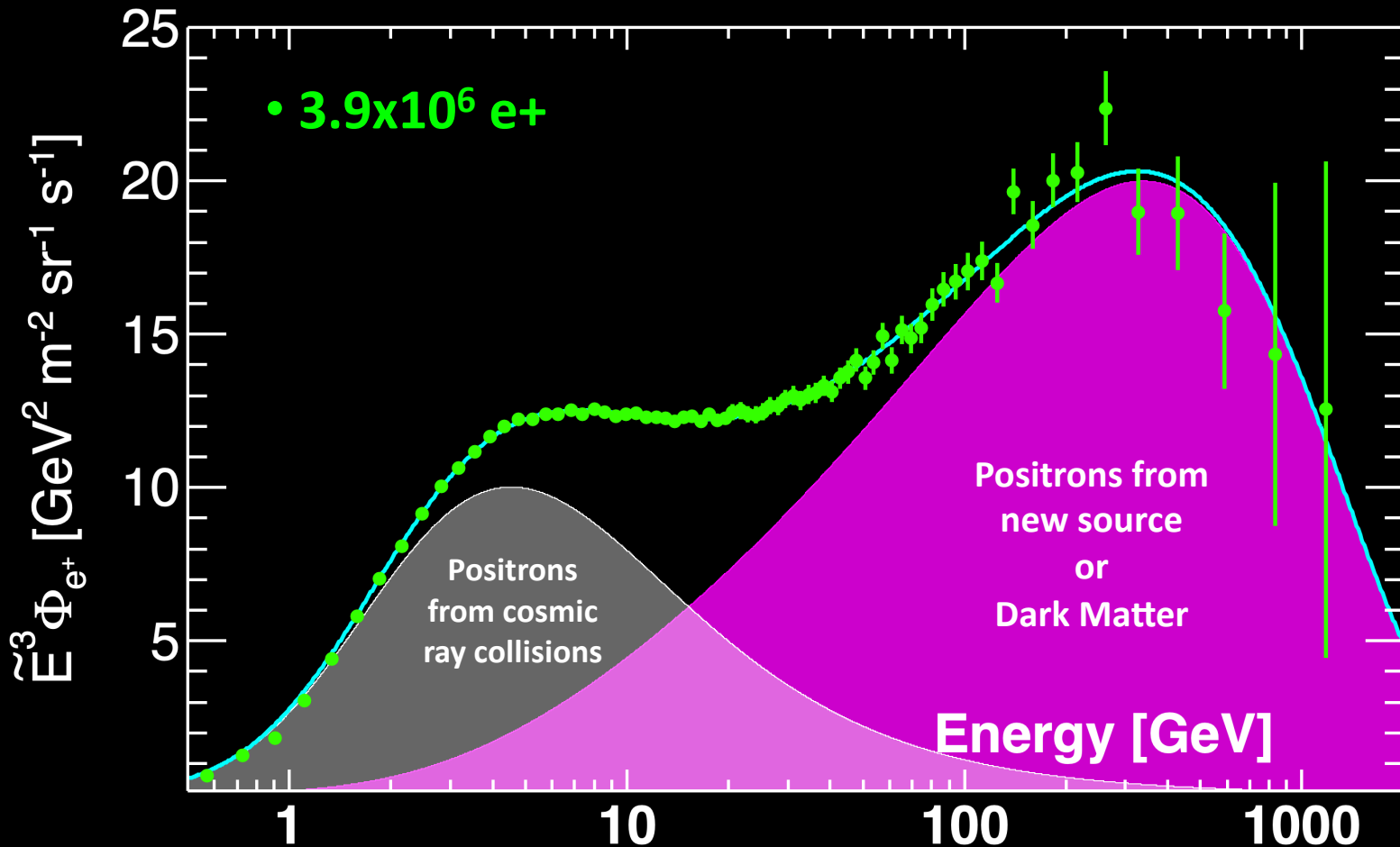
# Latest AMS Results on $e^+$ , $e^-$ , and $\bar{p}$



The positron flux is the sum of low-energy part from cosmic ray collisions plus a high-energy part from pulsars or dark matter both with a cutoff energy  $E_s$ .

$$\Phi_{e^+}(E) = \frac{E^2}{\hat{E}^2} \left[ C_d (\hat{E}/E_1)^{\gamma_d} + C_s (\hat{E}/E_2)^{\gamma_s} \exp(-\hat{E}/E_s) \right]$$

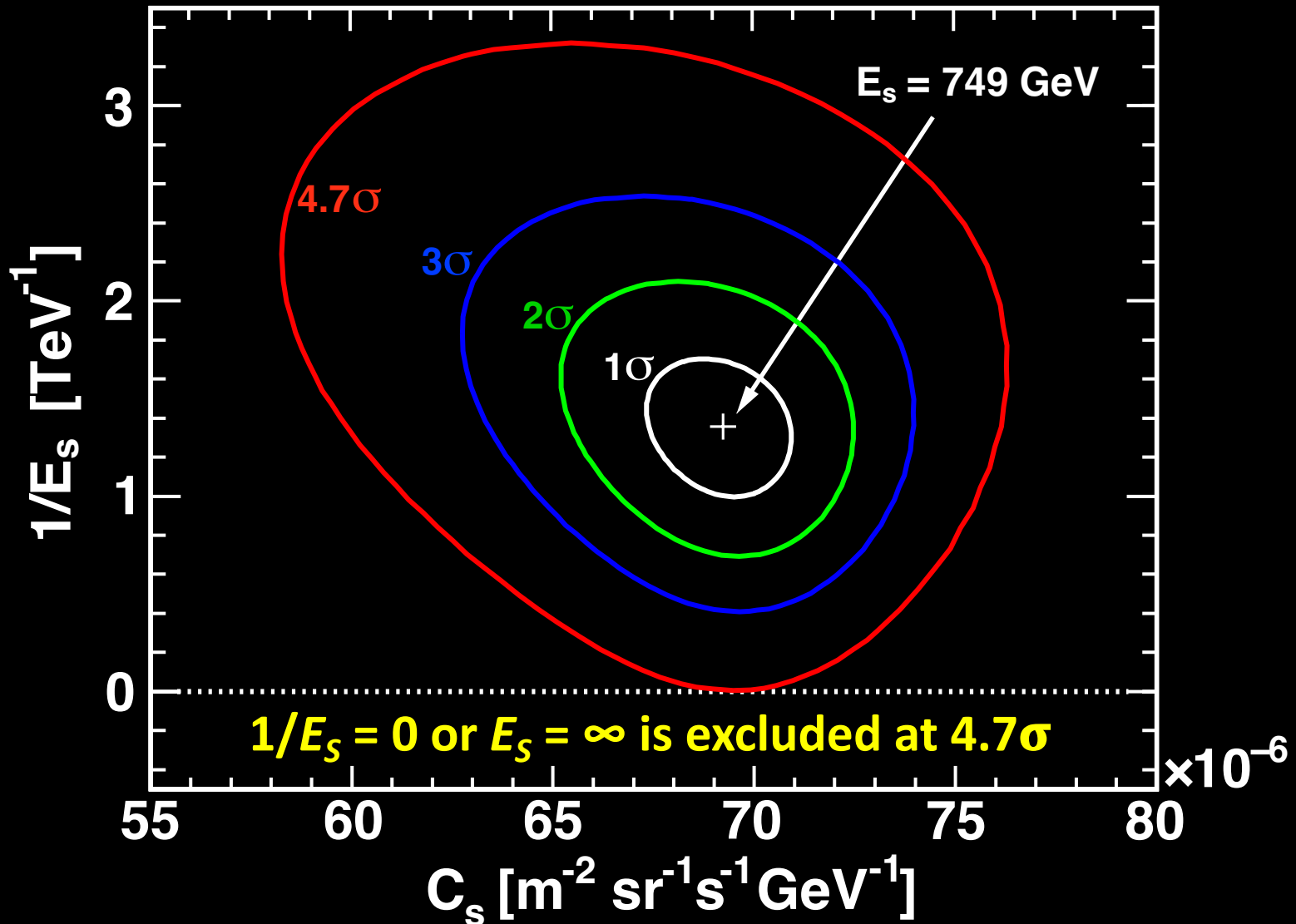
Solar Collisions
Pulsars or Dark Matter



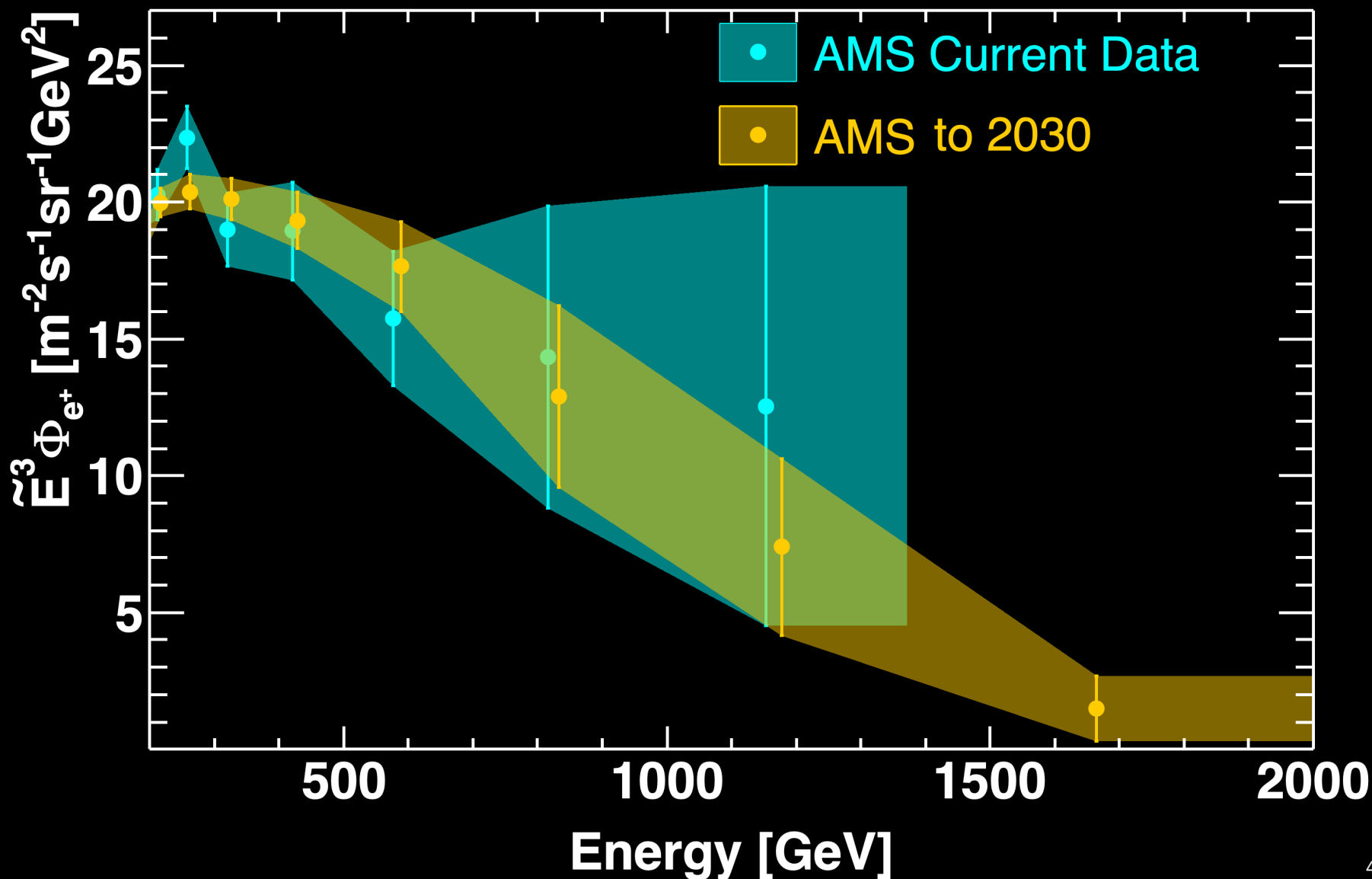
The existence of the finite cutoff energy ( $4.7\sigma$ ) is a new and unexpected observation

# Determination of the cutoff energy $E_s$

$$\Phi_{e^+}(E) = \frac{E^2}{\hat{E}^2} \left[ \overset{\text{Collisions}}{C_d (\hat{E}/E_1)^{\gamma_d}} + \overset{\text{New Source or Dark Matter}}{C_s (\hat{E}/E_2)^{\gamma_s} \exp(-\hat{E}/E_s)} \right]$$



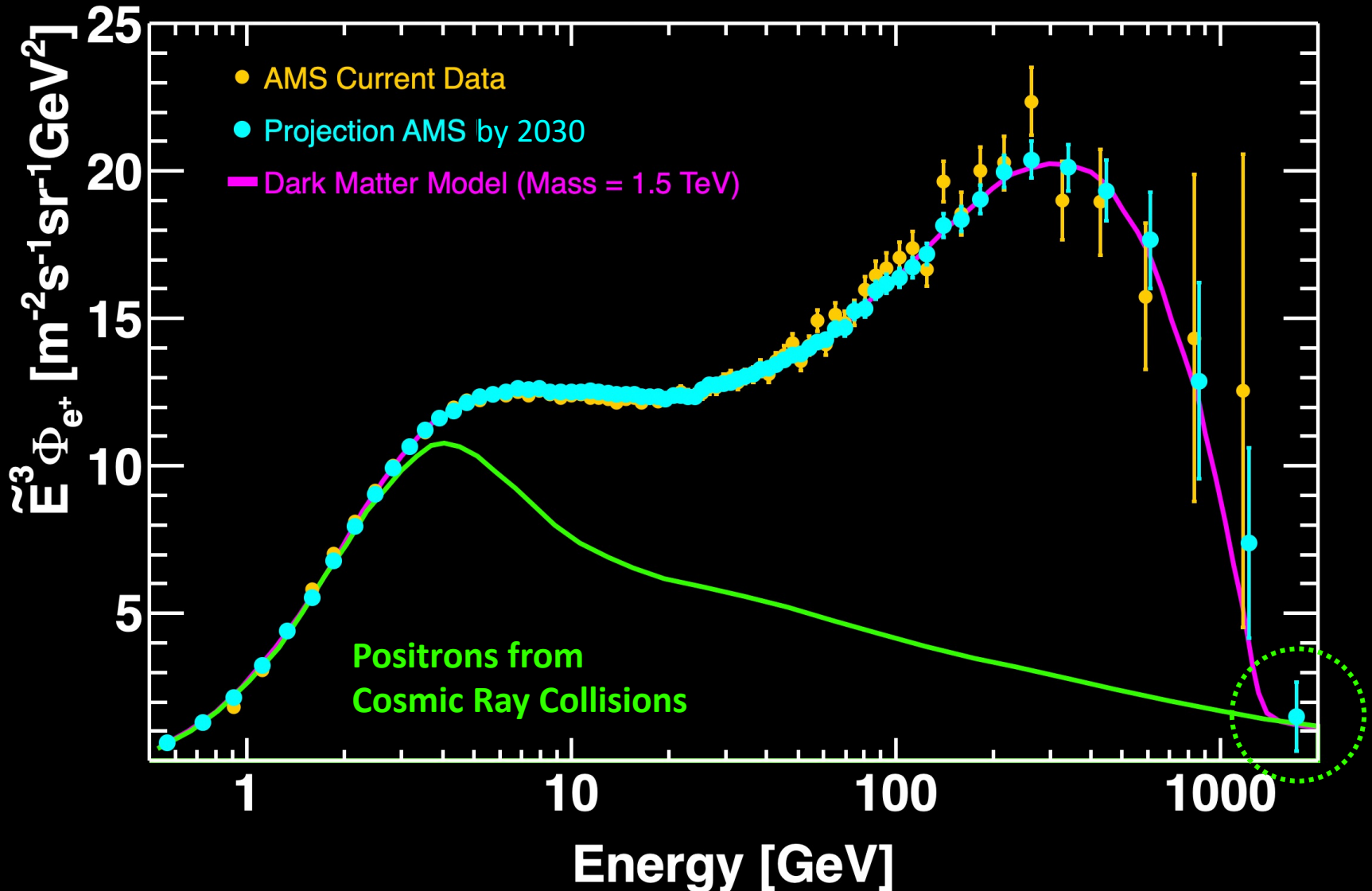
**By 2030, AMS will extend the energy range of the positron flux measurement from 1.4 to 2 TeV and reduce the error by a factor of two compared to current data**





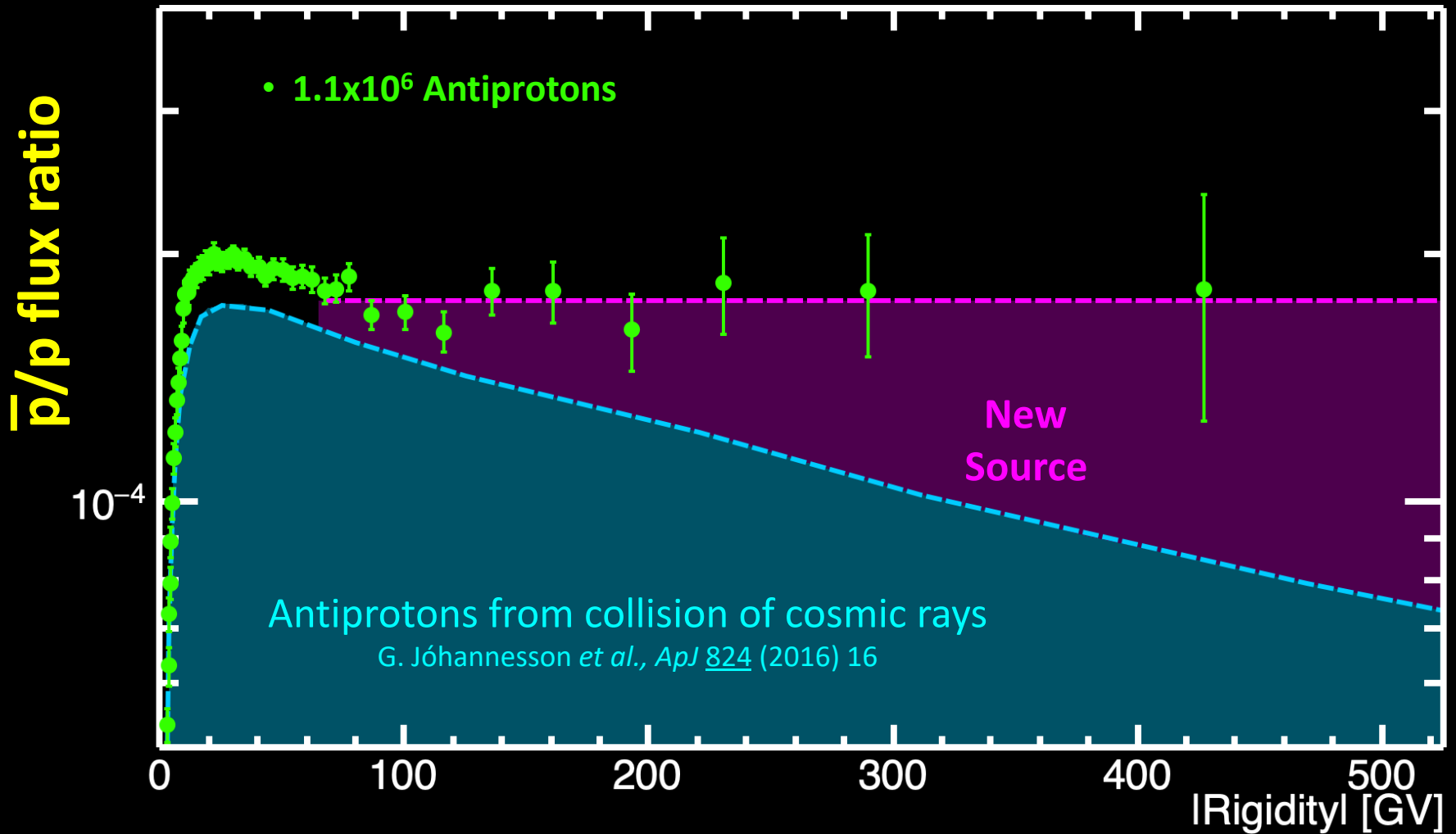
# Determination of the Origin of Cosmic Positrons by 2030

AMS will ensure that the measured high energy positron spectrum indeed drops off quickly and, at the highest energies, the positrons only come from cosmic ray collisions as predicted by dark matter models



# Properties of Cosmic Antiprotons

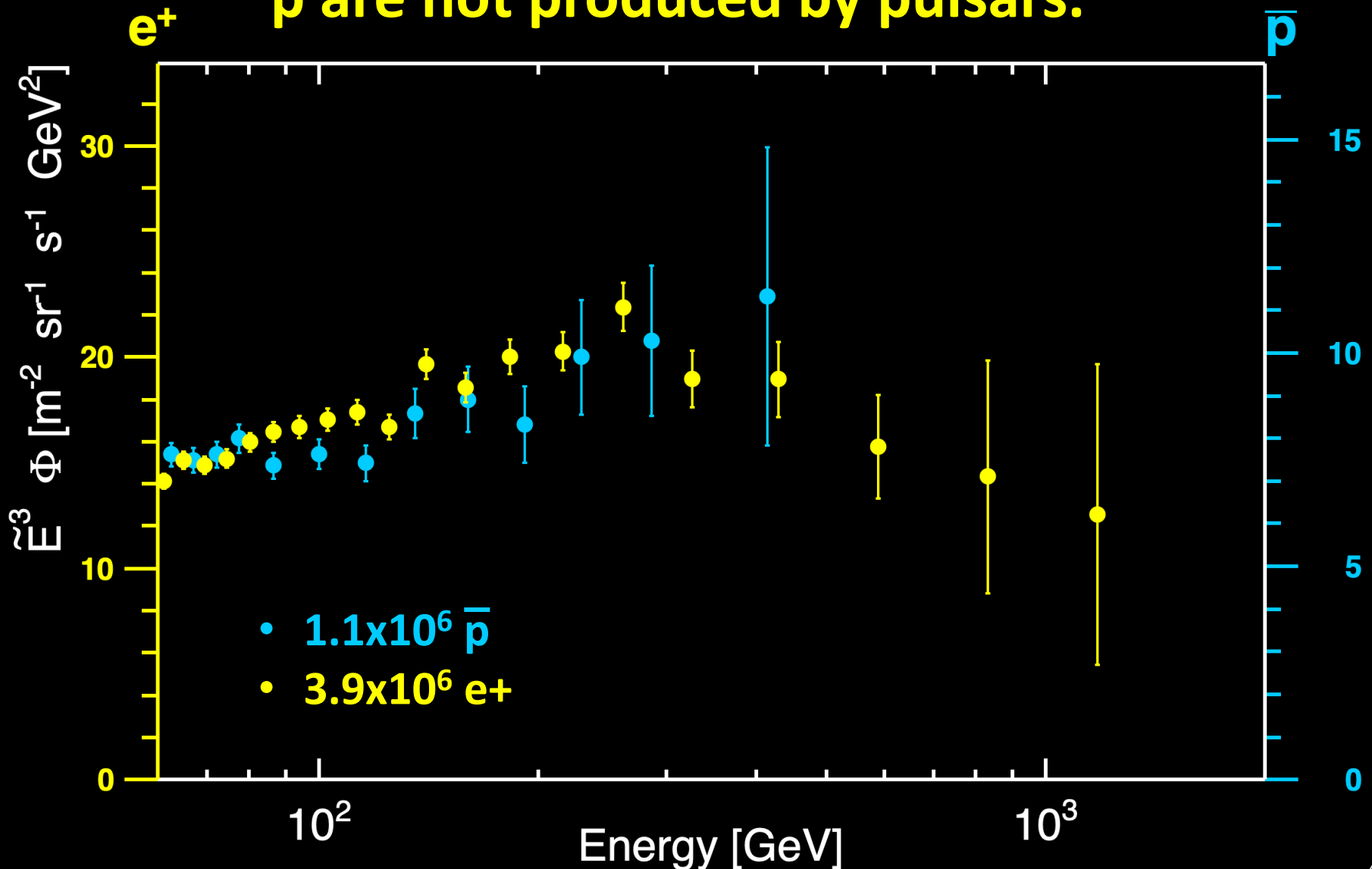
The antiproton-to-proton flux ratio shows that above 60 GV the ratio is energy independent.



# Properties of Cosmic Antiprotons

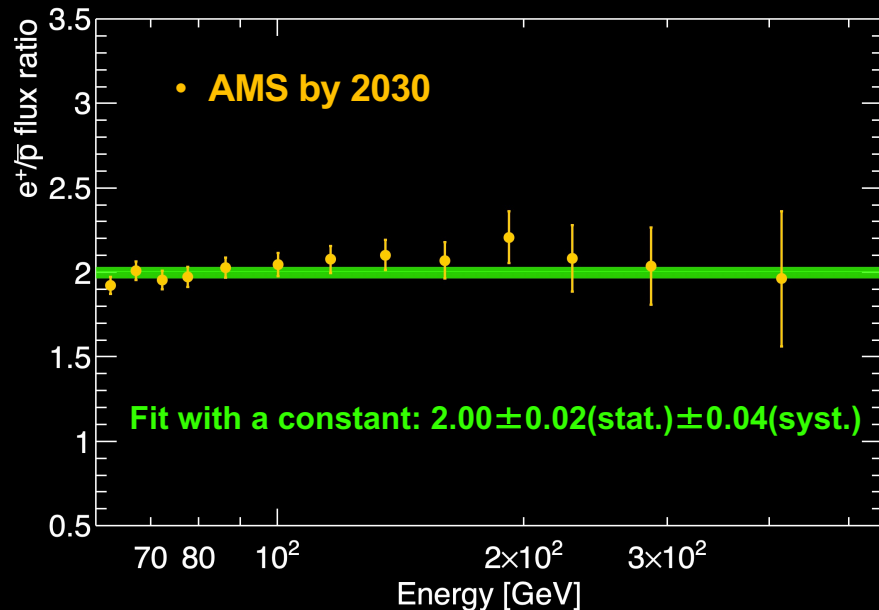
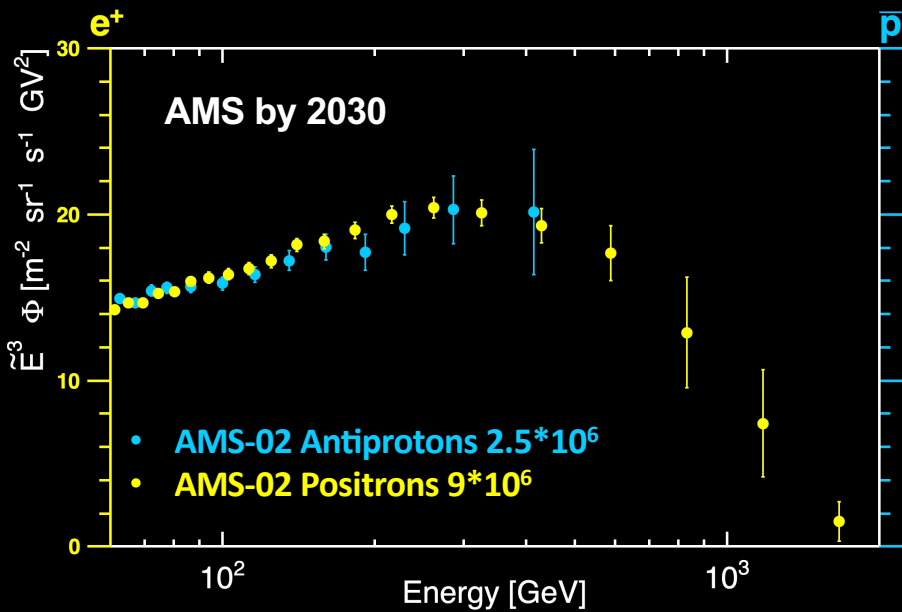
The  $\bar{p}$  and  $e^+$  fluxes have identical rigidity dependence.

$\bar{p}$  are not produced by pulsars.



# By 2030, AMS will greatly improve the accuracy of the antiproton spectra

The identical behaviour of positrons and antiprotons excludes the pulsar origin of positrons

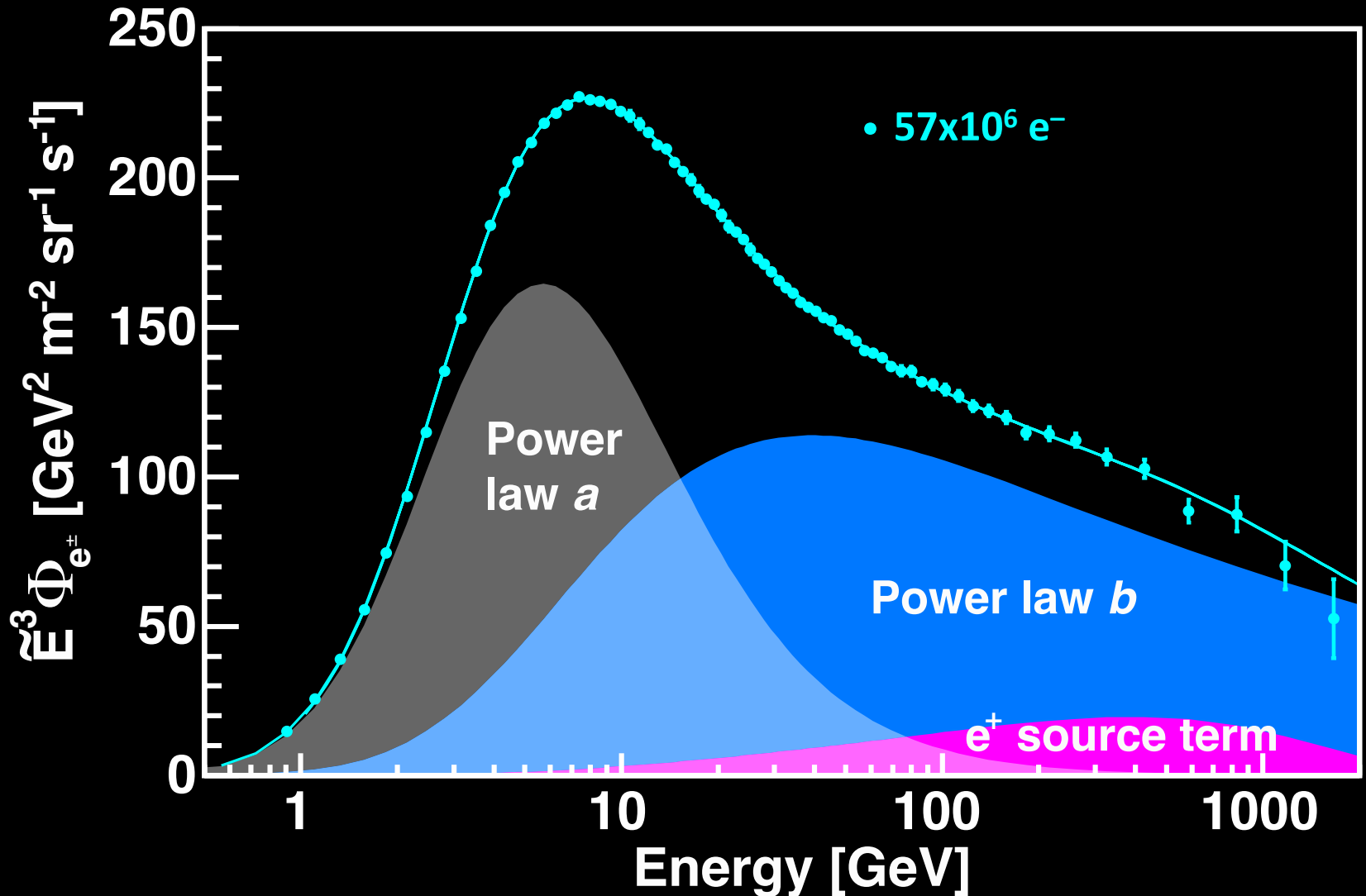


# AMS Result on the electron spectrum

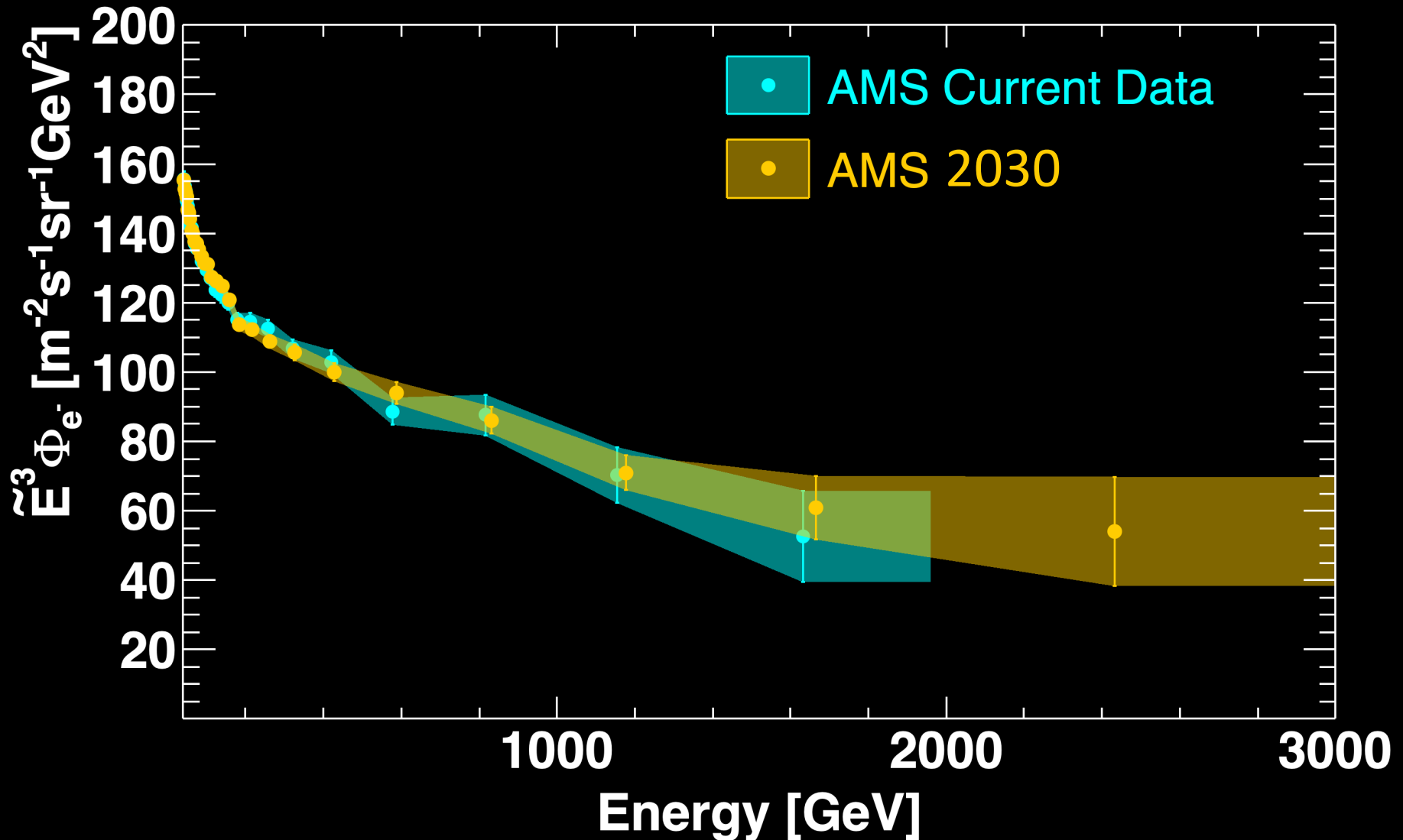
The spectrum fits well with two power laws ( $a$ ,  $b$ ) and a source term like positrons

$$\Phi_{e^-}(E) = \frac{E^2}{\widehat{E}^2} (C_a \widehat{E}^{\gamma_a} + C_b \widehat{E}^{\gamma_b} + \text{Positron Source Term})$$

Solar    Power law  $a$     Power law  $b$

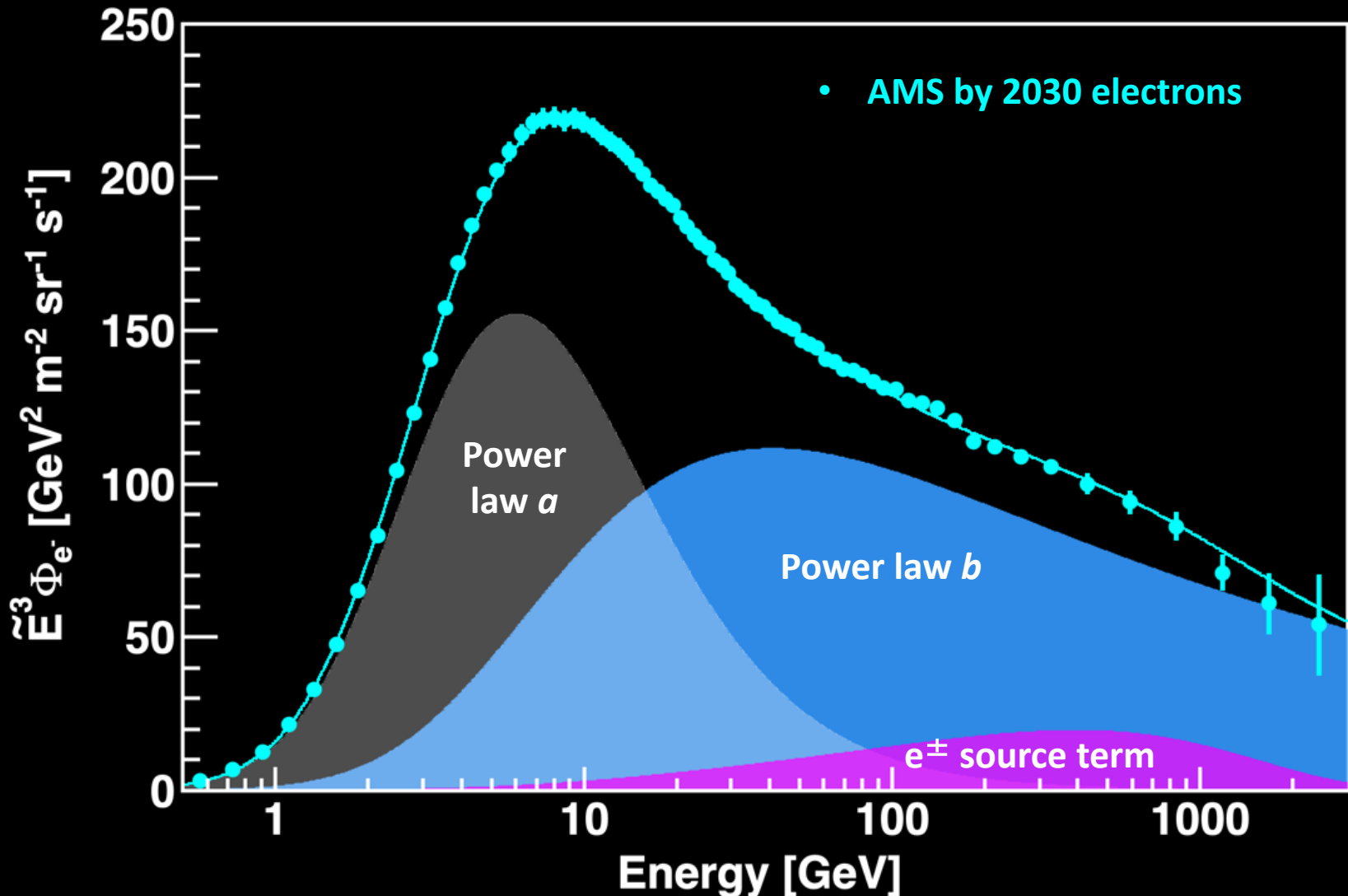


**By 2030, AMS will extend the energy range of the electron flux measurement from 2 to 3 TeV  
and reduce the error by a factor of two compared to current data**

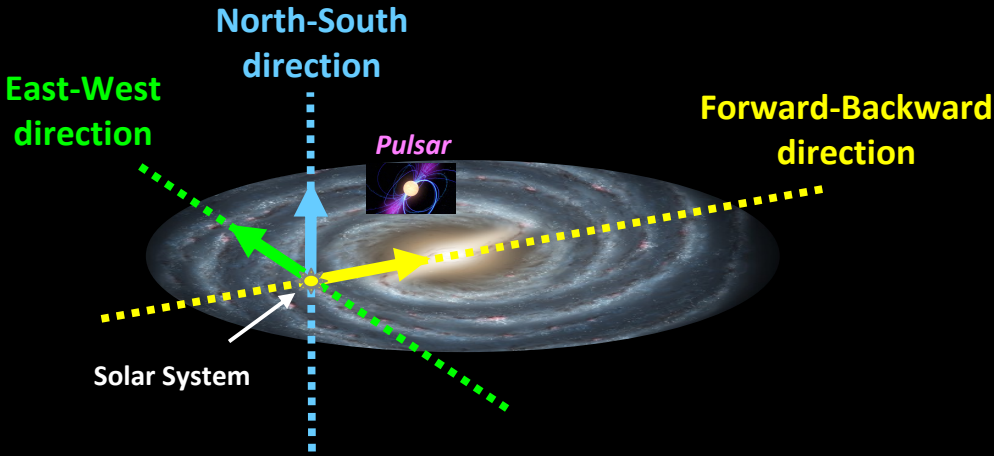


By 2030, the charge-symmetric nature of the high energy source will be established at the  $4\sigma$  level

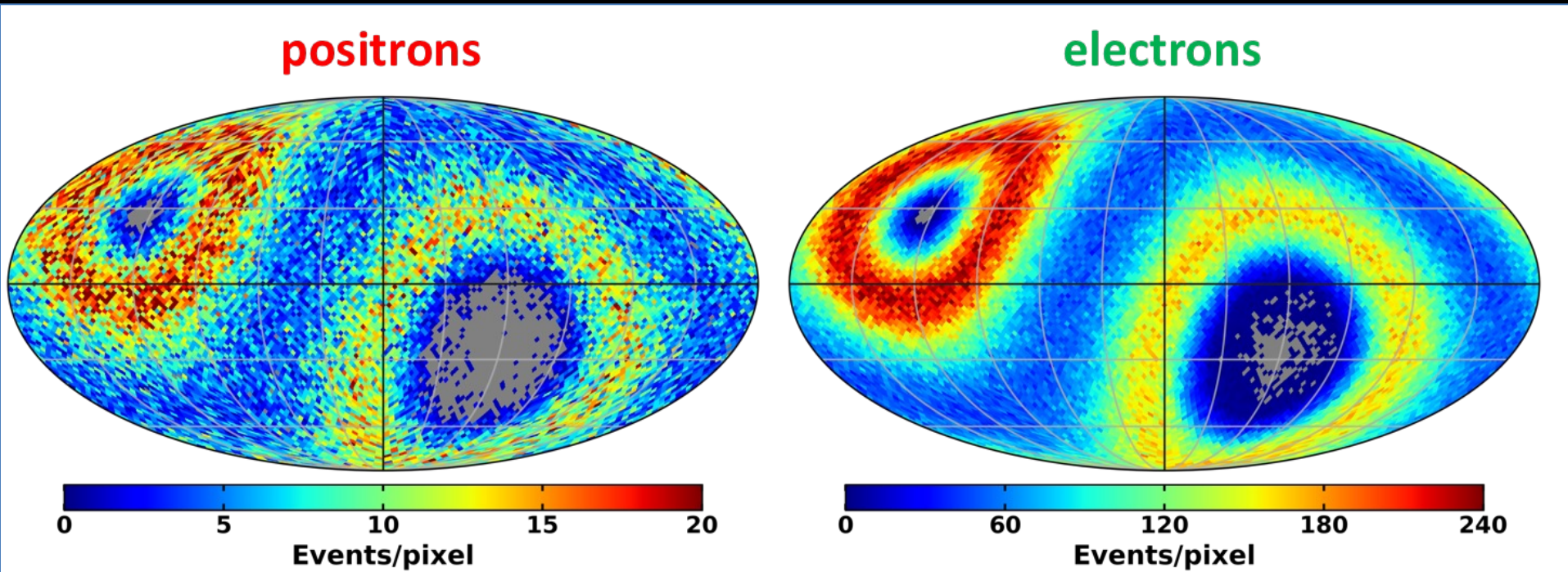
New sources, like Dark Matter or Pulsars, produce equal amounts of  $e^+$  and  $e^-$



# Positron Anisotropy and Dark Matter



Astrophysical point sources will imprint a higher anisotropy on the arrival directions of energetic positrons than a smooth dark matter halo.



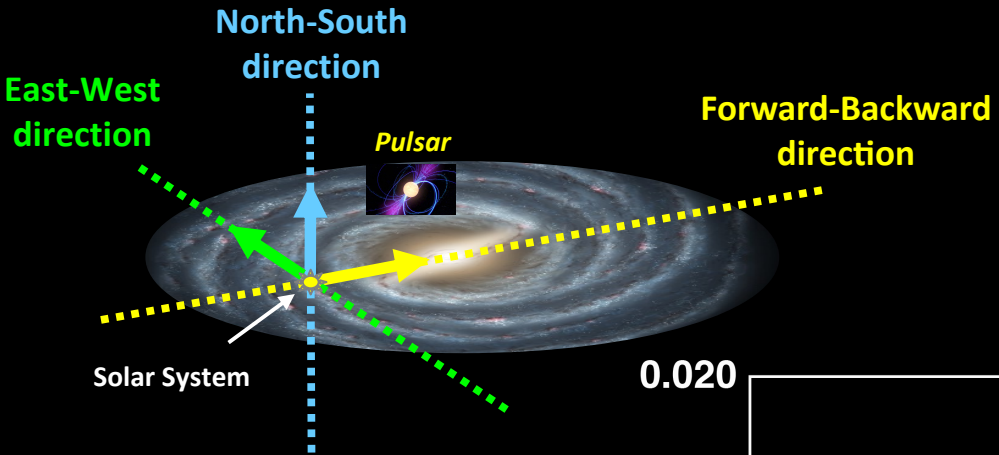
Dipole anisotropy:

$$\delta = 3\sqrt{C_1/4\pi}$$

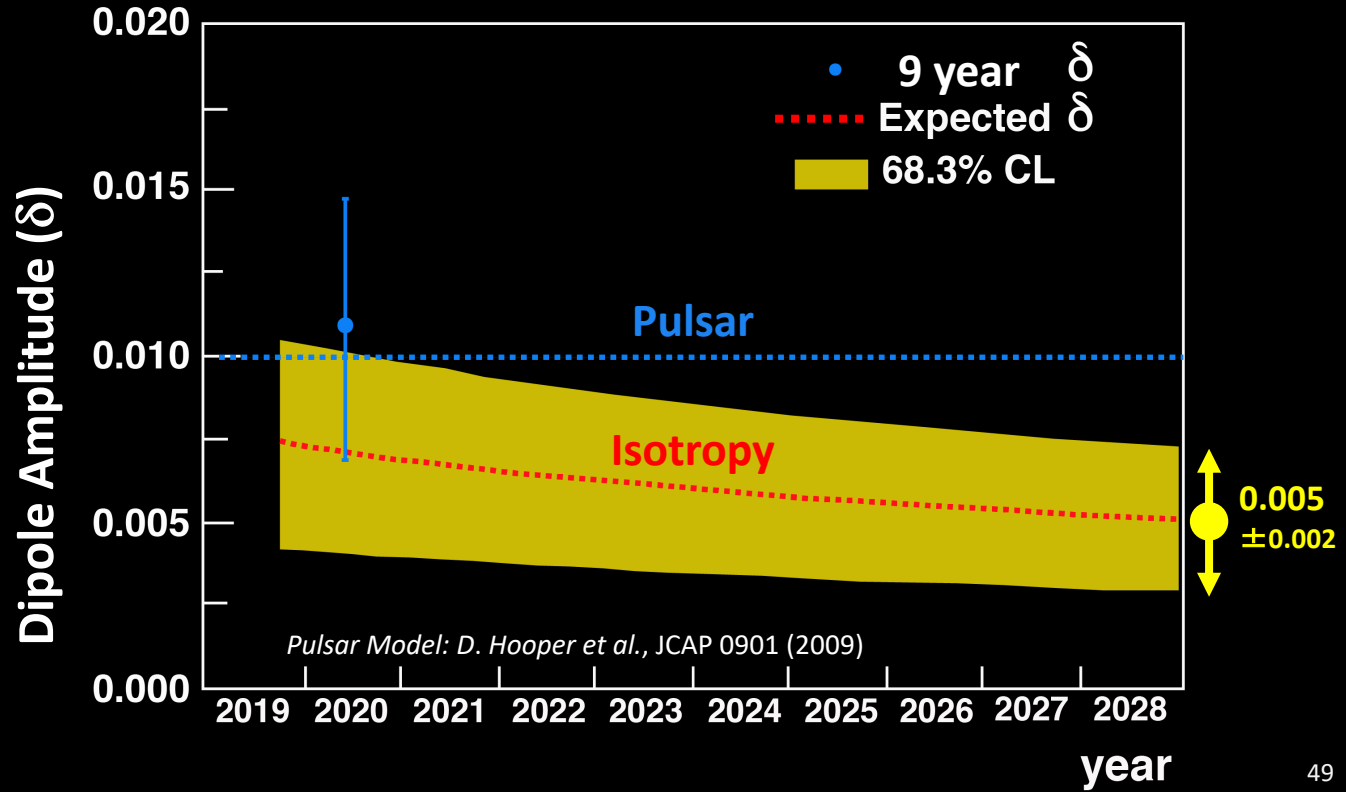
$C_1$  is the dipole moment



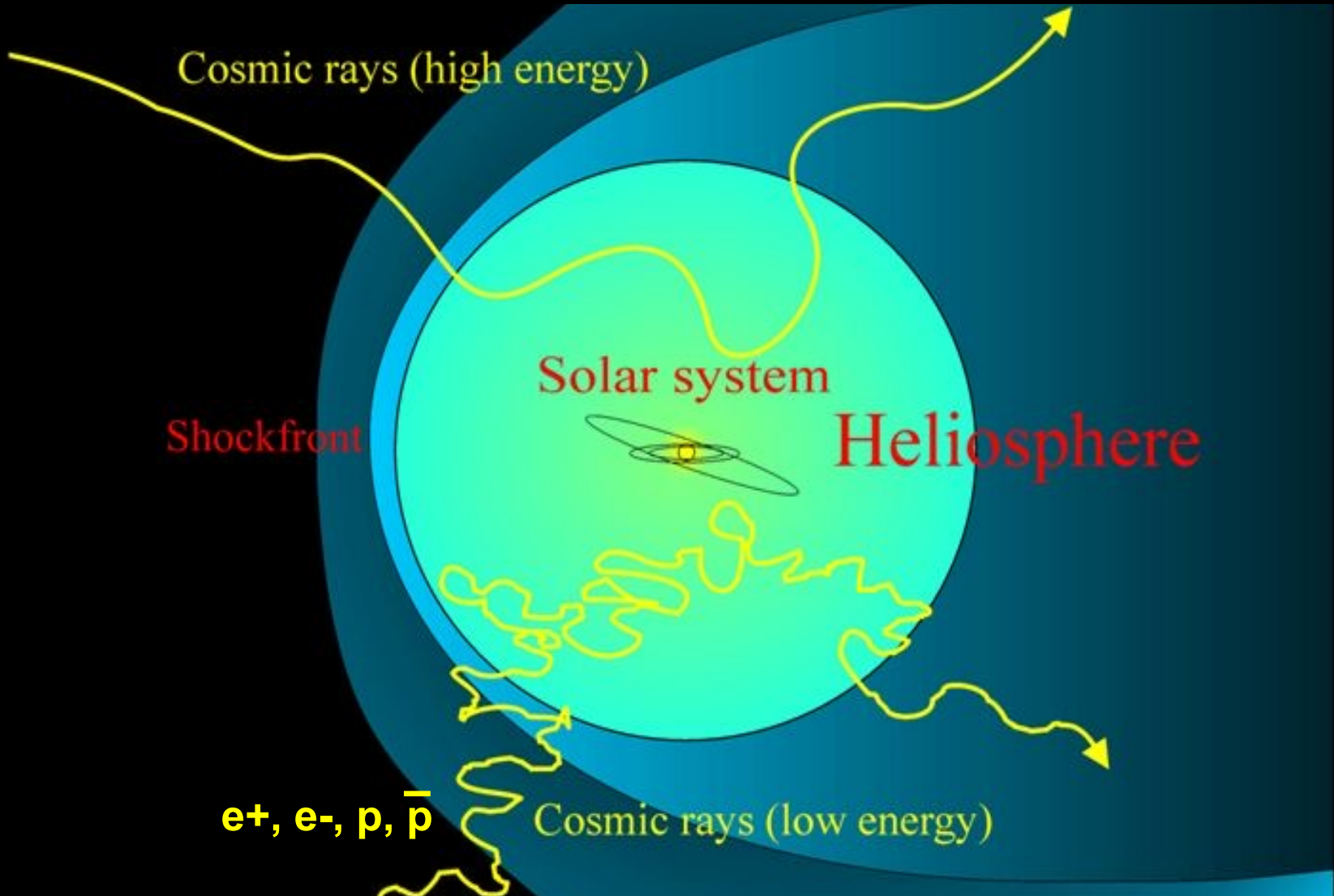
**By 2030, the positron statistics will allow us to measure the anisotropy accurately to permit a separation between dark matter and pulsars at the 99.93% C.L.**



**Positrons from a pulsar would have a preferential direction**



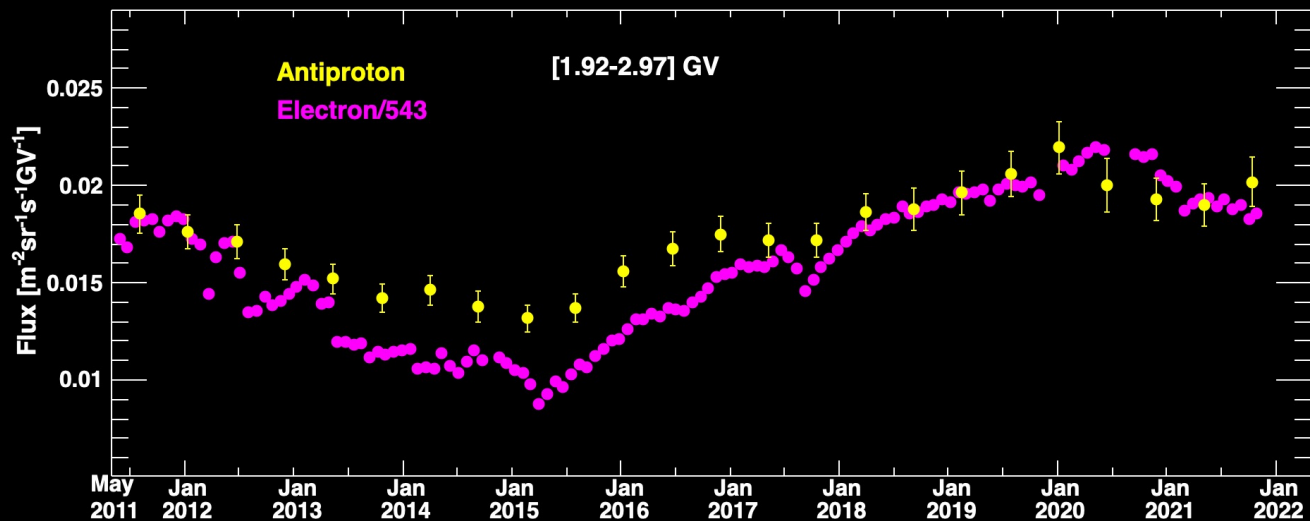
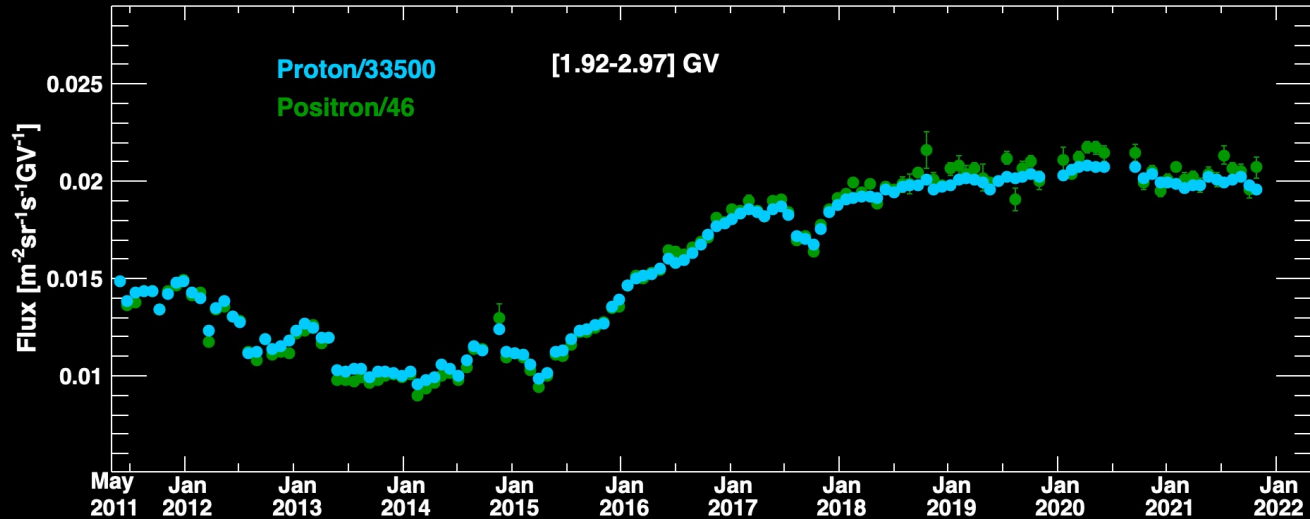
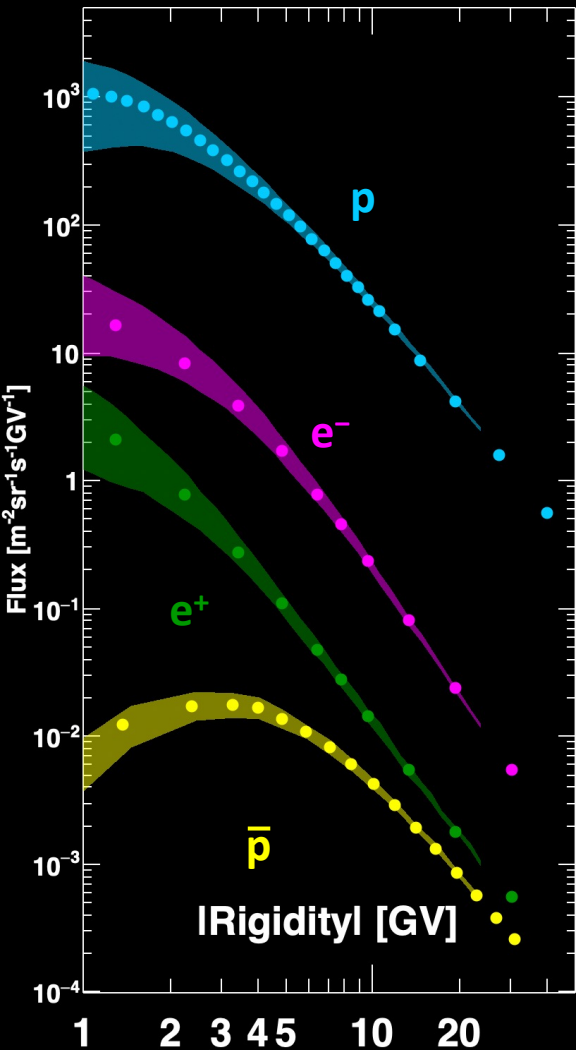
# Latest AMS Results on Elementary Particles ( $e^+$ , $e^-$ , $p$ , $\bar{p}$ , ...) in the Heliosphere over an **11-year Solar Cycle** (2011-2022)



**By 2030, AMS will study Heliosphere physics over 22-year Solar Cycle**

# Elementary Particles ( $e^+$ , $e^-$ , $p$ , $\bar{p}$ , ...)

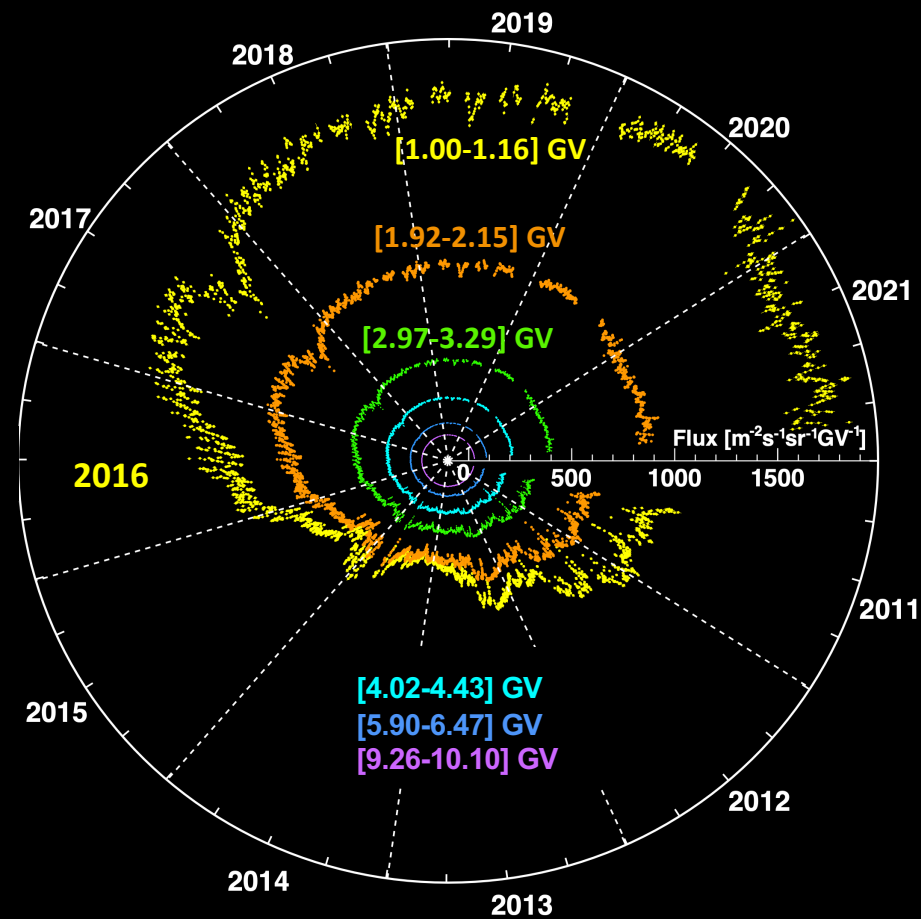
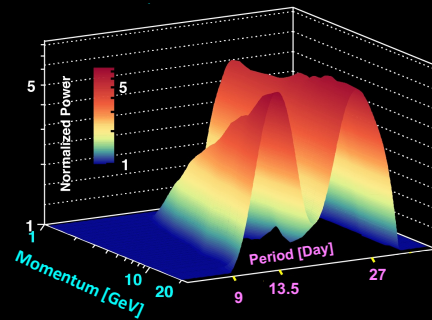
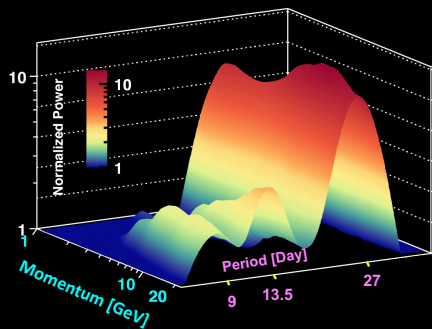
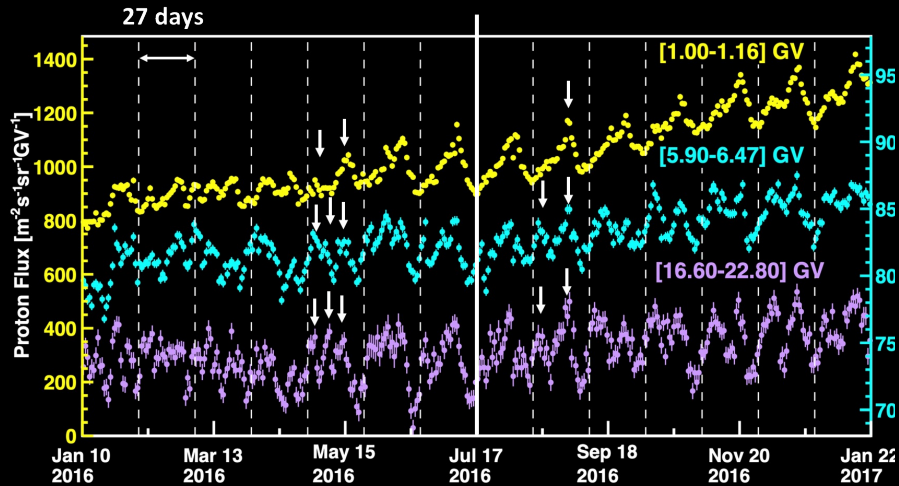
## in the Heliosphere over an 11-year Solar Cycle (2011-2022)



# Daily Protons in the Heliosphere

## Frequency Analysis of daily fluxes in 2016

Double-peak and triple-peak structures are visible



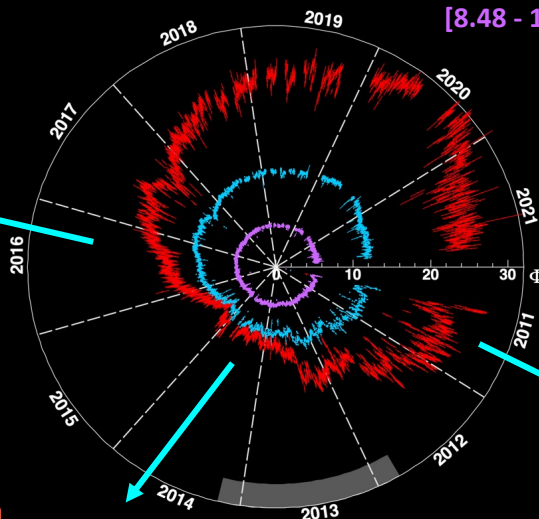
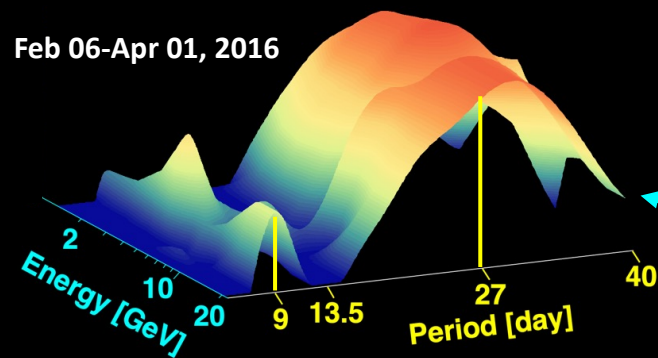
Recurrent variations with periods of 27, 13.5, and 9 days are observed. Unexpectedly, in 2016, the strength of the 9 and 13.5 day periodicity increases with increasing rigidity up to 10 and 20 GeV, respectively.

The model expects the strength of the periodicity to decrease with increasing rigidity.

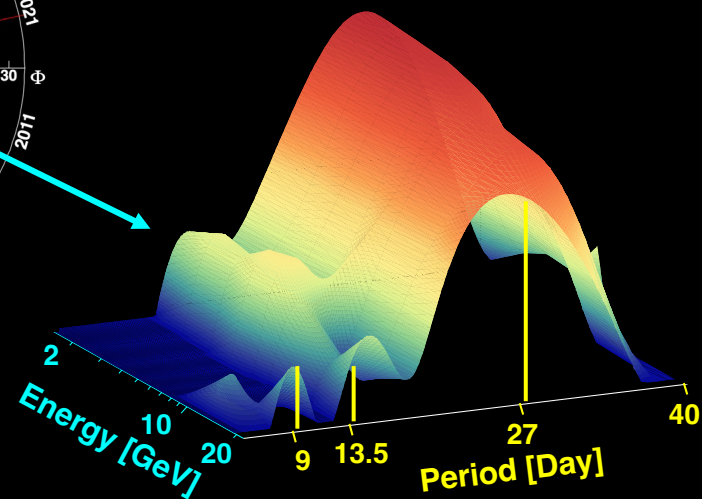
# First Observation of Periodicity in the Daily Electron Flux in the Heliosphere

[1.00 - 1.71] GV  $\Phi_{e^-}$   
 [2.97 - 4.02] GV  $\Phi_{e^-} \times 2.5$   
 [8.48 - 11.0] GV  $\Phi_{e^-} \times 20$

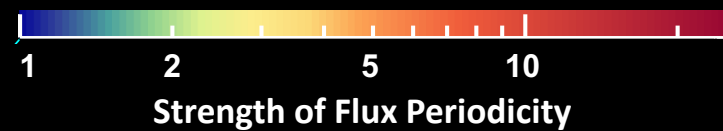
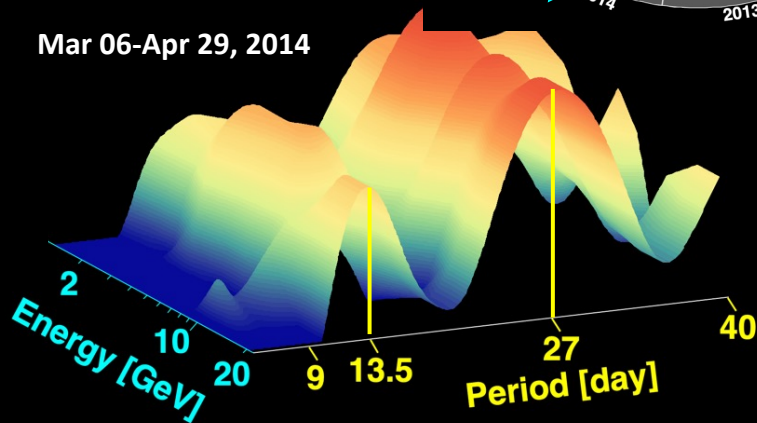
Feb 06-Apr 01, 2016



May 20-Dec. 16, 2011

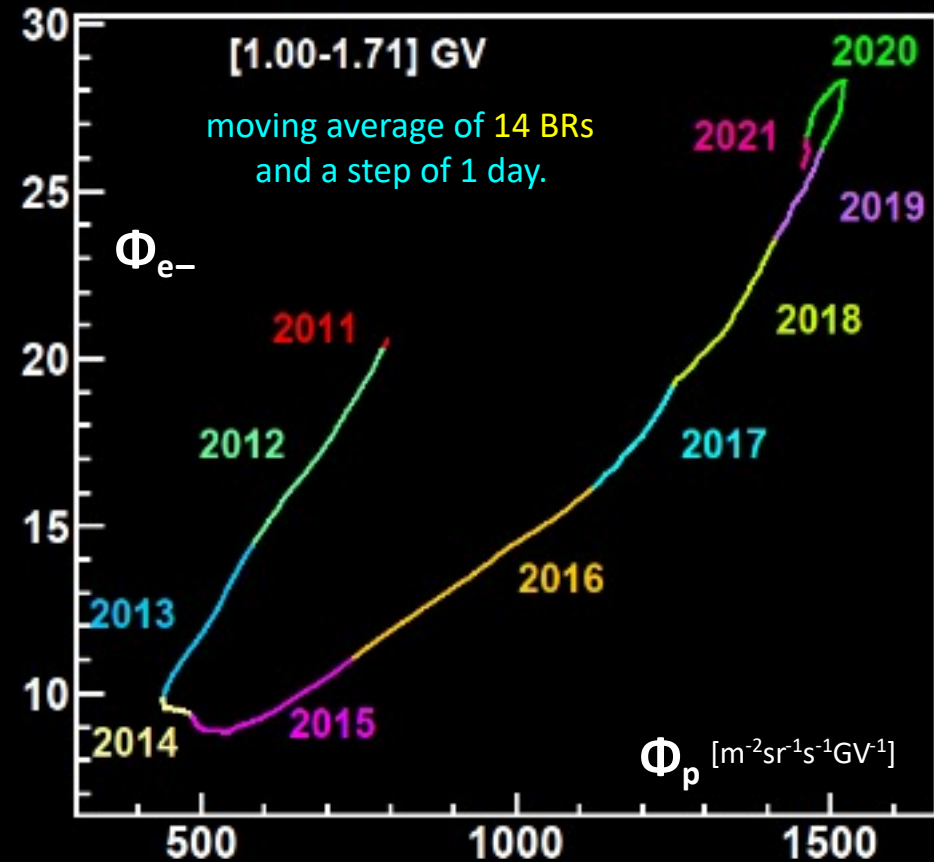
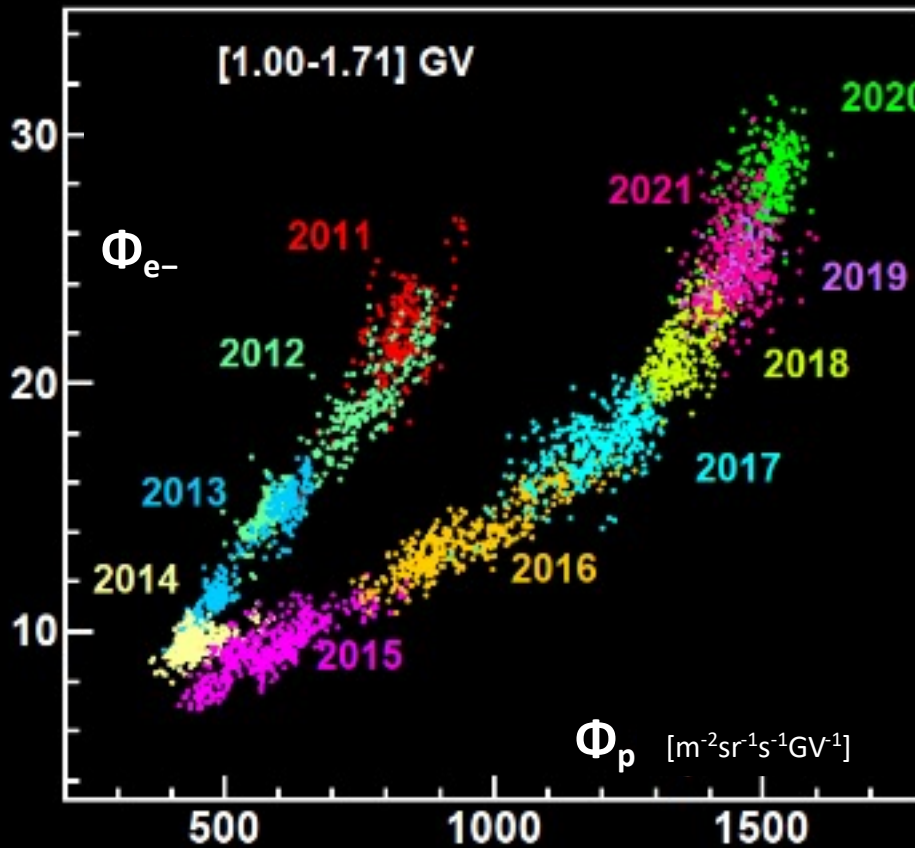


Mar 06-Apr 29, 2014



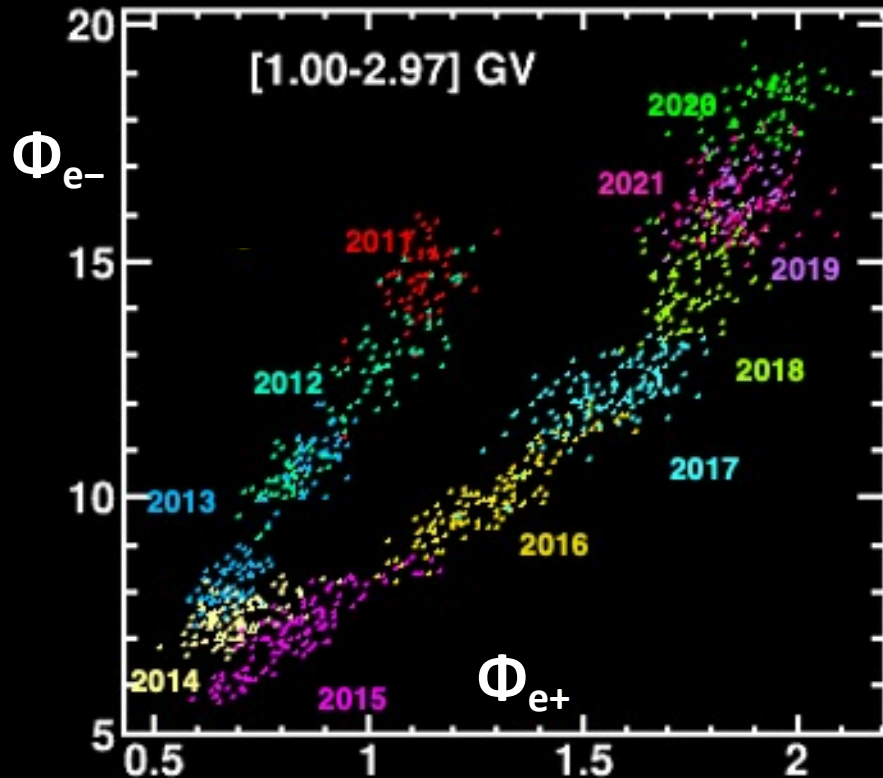
# Elementary Particles in the Heliosphere

Observation of a hysteresis between daily electron flux  $\Phi_{e^-}$  and daily proton flux  $\Phi_p$

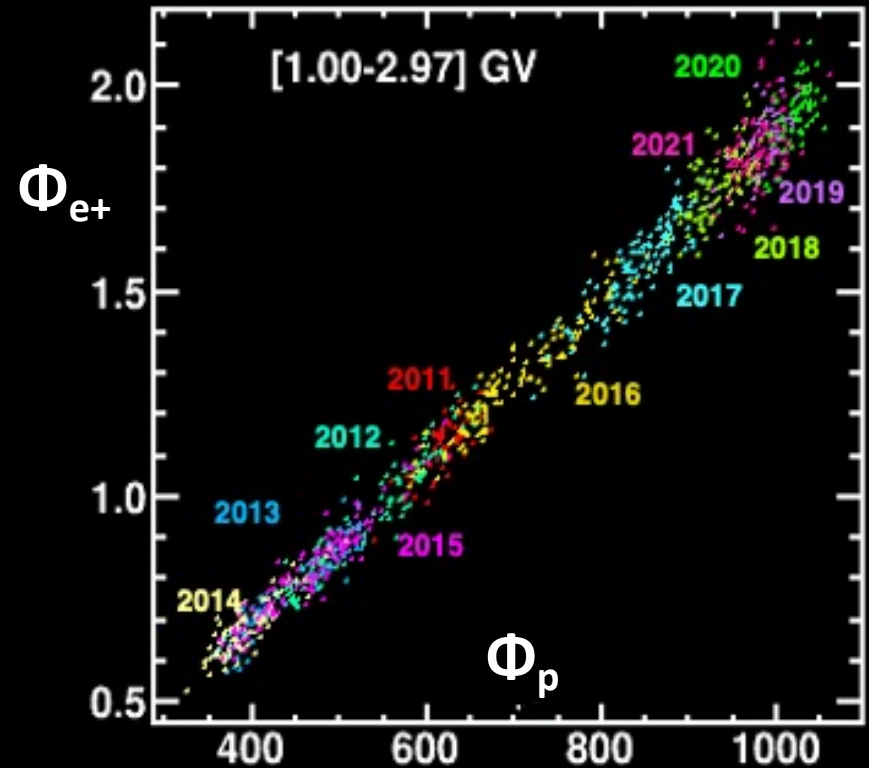


# Relationship between charge and mass

Equal mass, Opposite charge

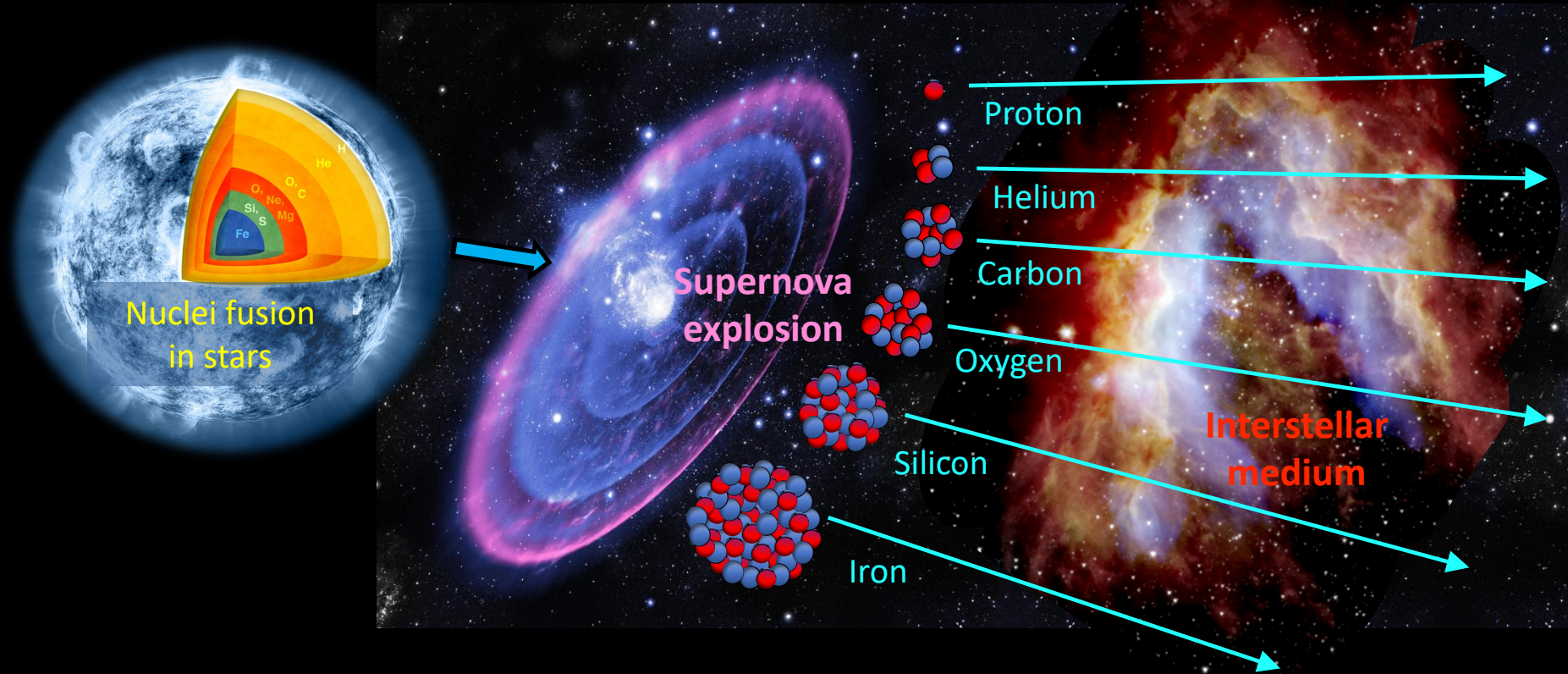


Equal charge, different mass



# Latest AMS Results on Primary Cosmic Rays

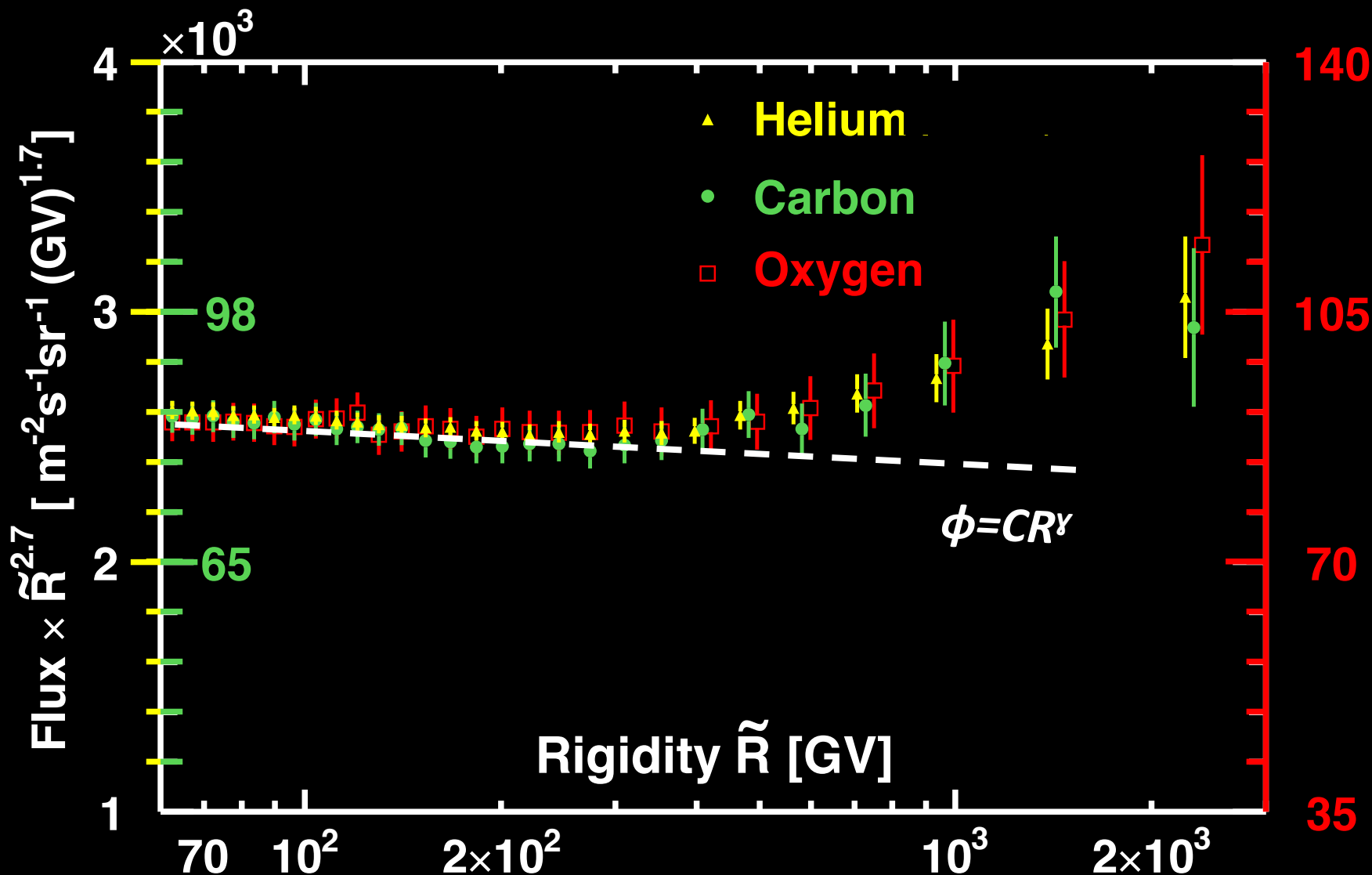
Primary cosmic rays p, He, C, O, ..., Si, ..., Fe are produced during the lifetime of stars and accelerated by supernovae. They propagate through interstellar medium before they reach Earth.



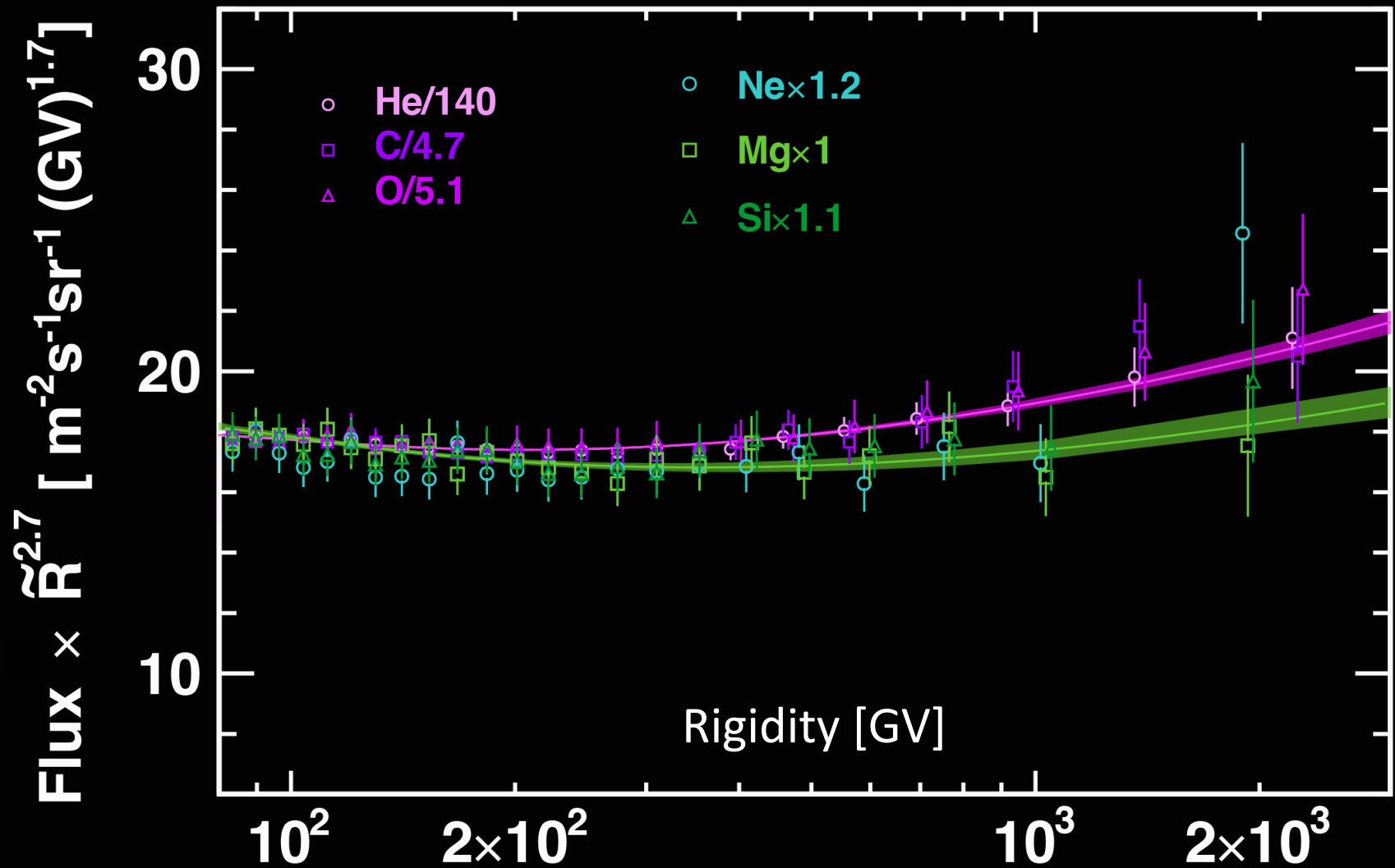
Measurements of primary cosmic ray fluxes are fundamental to understanding the origin, acceleration, and propagation processes of cosmic rays in the Galaxy.



Unexpectedly, above 60 GV, the light primary cosmic rays He-C-O have **identical** rigidity ( $R=P/Z$ ) dependence. In the traditional understanding  $\phi=CR^\gamma$ , AMS found that He-C-O **harden** in the same way above  $\sim 200$  GV.

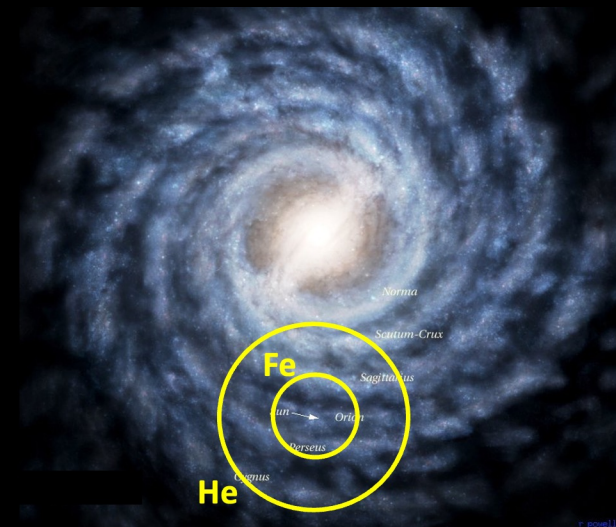
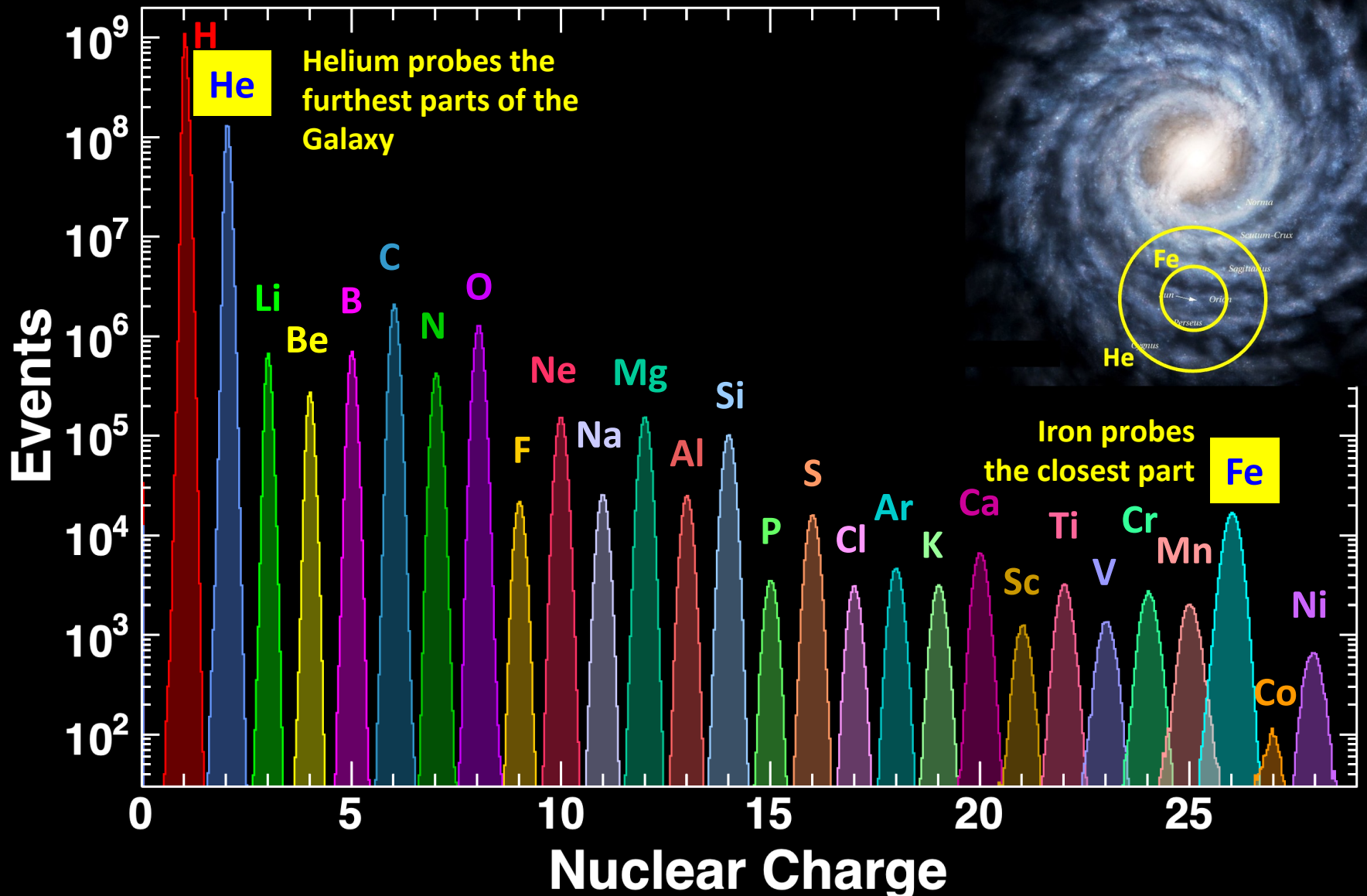


Heavier elements Ne-Mg-Si have their own rigidity dependence, different than the dependence of light elements He-C-O.

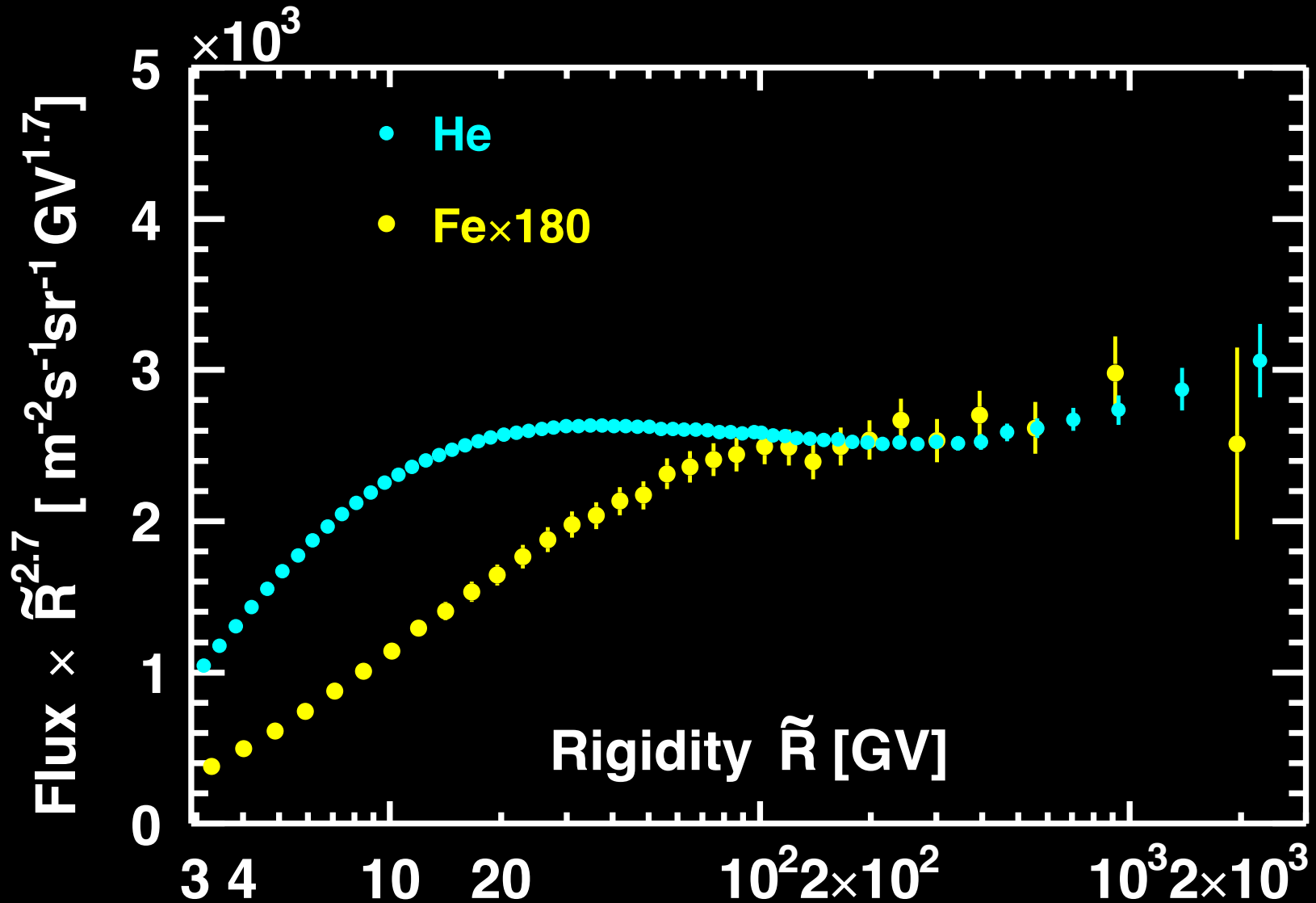


Primary cosmic rays have at least two classes.

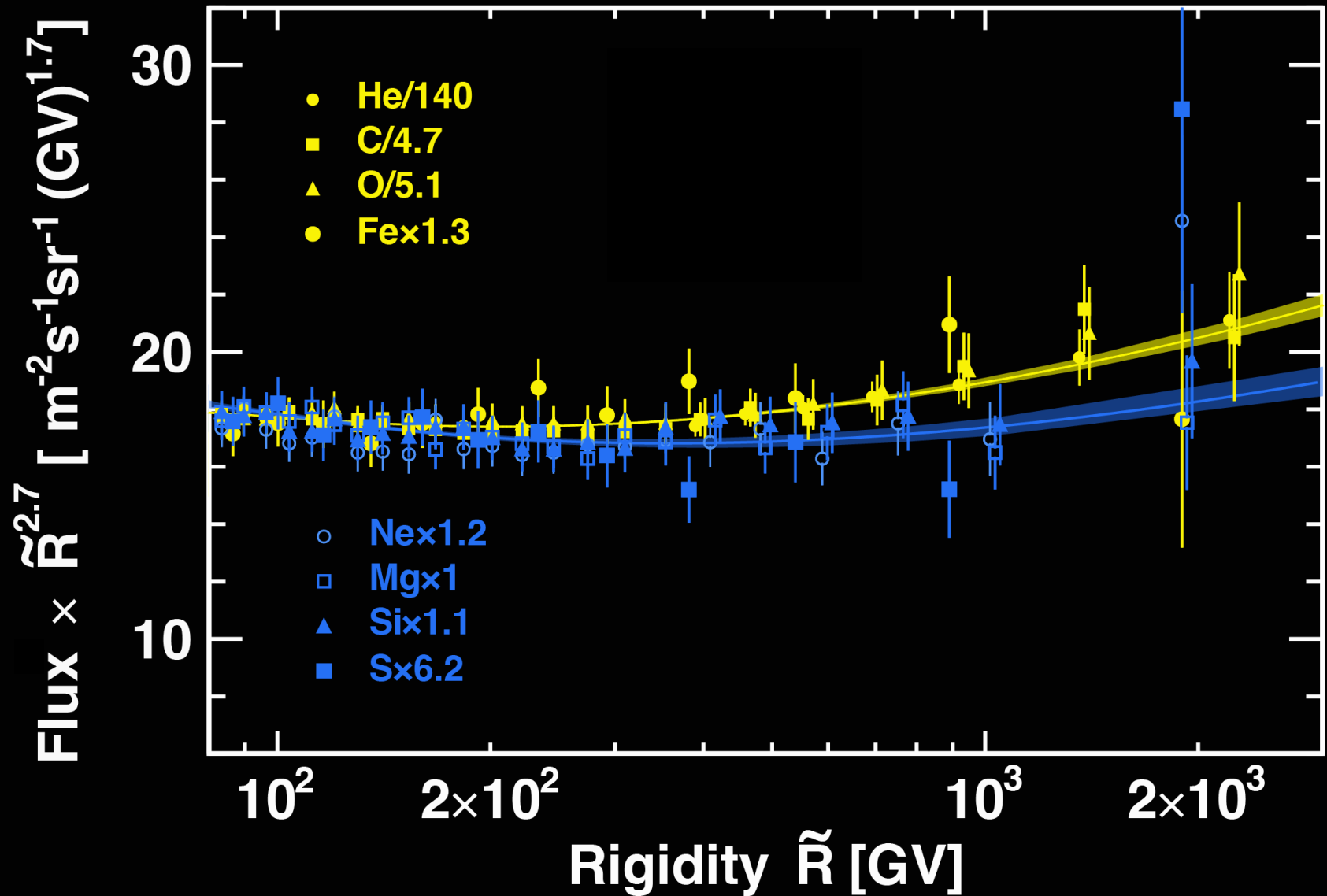
Iron is a very important element in cosmic ray theories because it is the heaviest element produced during stellar evolution. Iron has a large interaction rate with the interstellar medium and comes from the closest part of the Galaxy.



**Iron** is in the **He-C-O primary cosmic ray group** instead of the expected heavy Ne-Mg-Si group.



# AMS Results on all 8 primary elements: They are in two classes

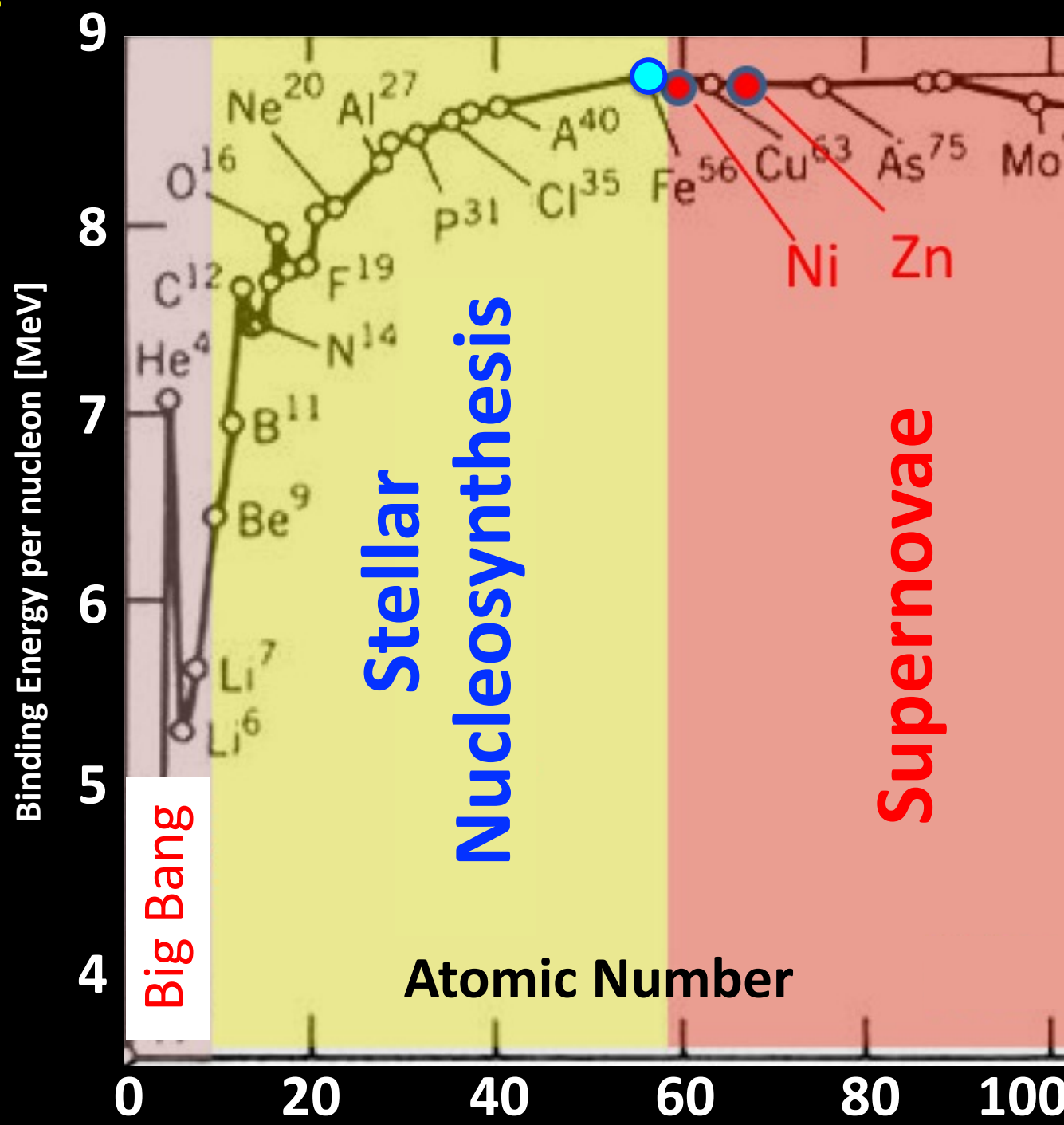


# On the origin of cosmic rays

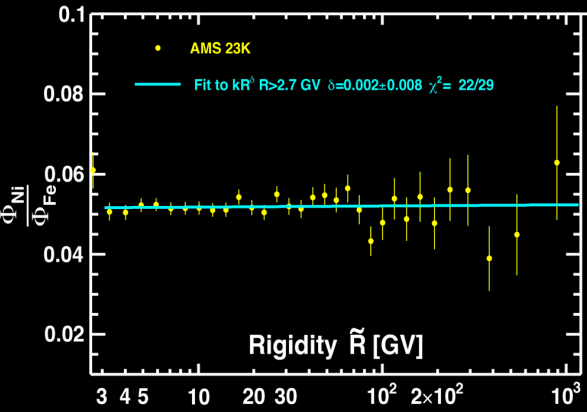
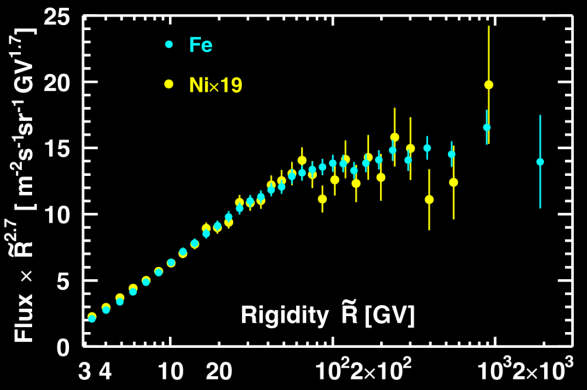
Iron is the heaviest element produced by Stellar Nucleosynthesis

Nickel is the lightest element created by supernovae.

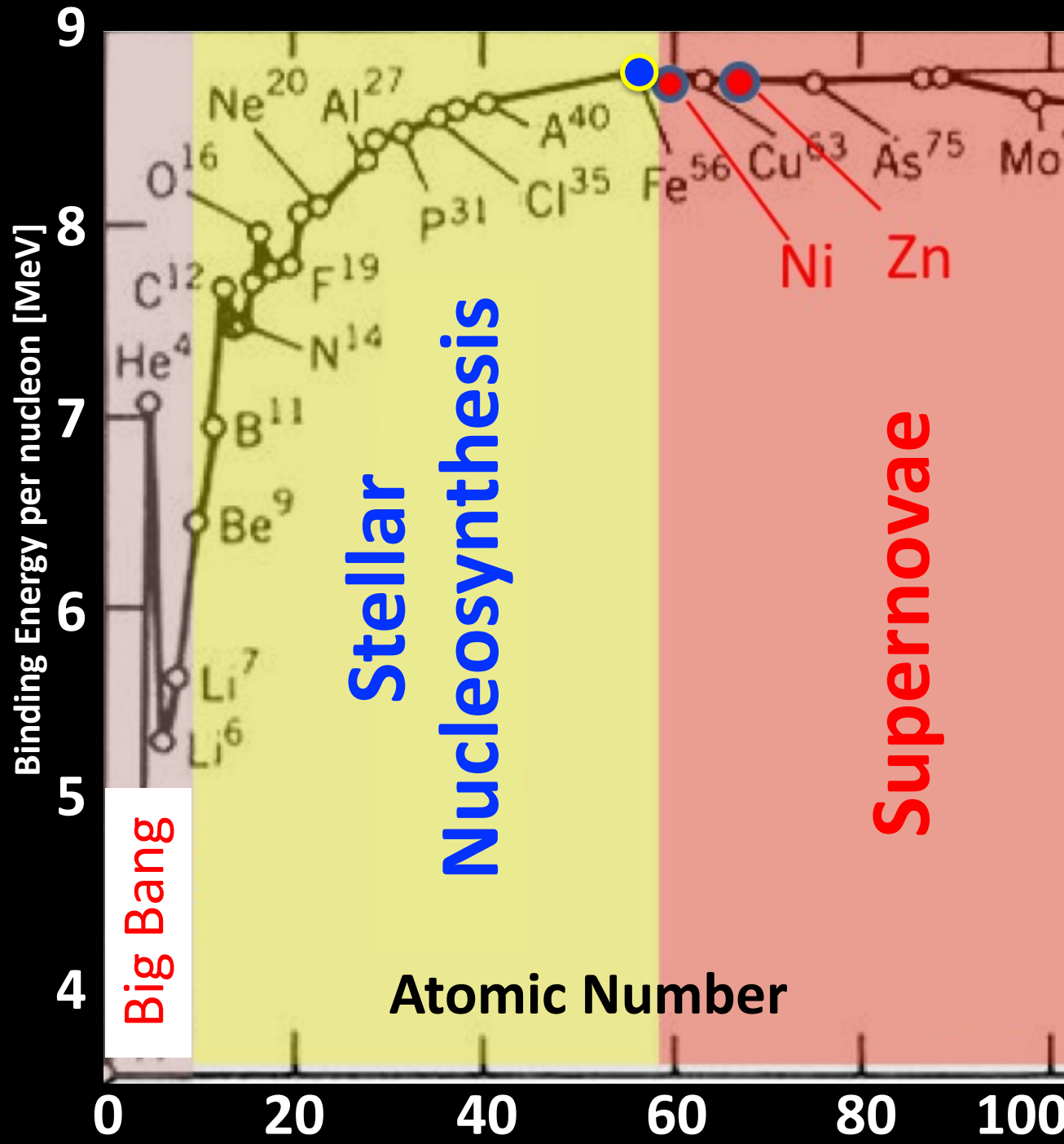
In what way are they different?



Within the given statistics, the rigidity dependence of Ni is similar to Fe

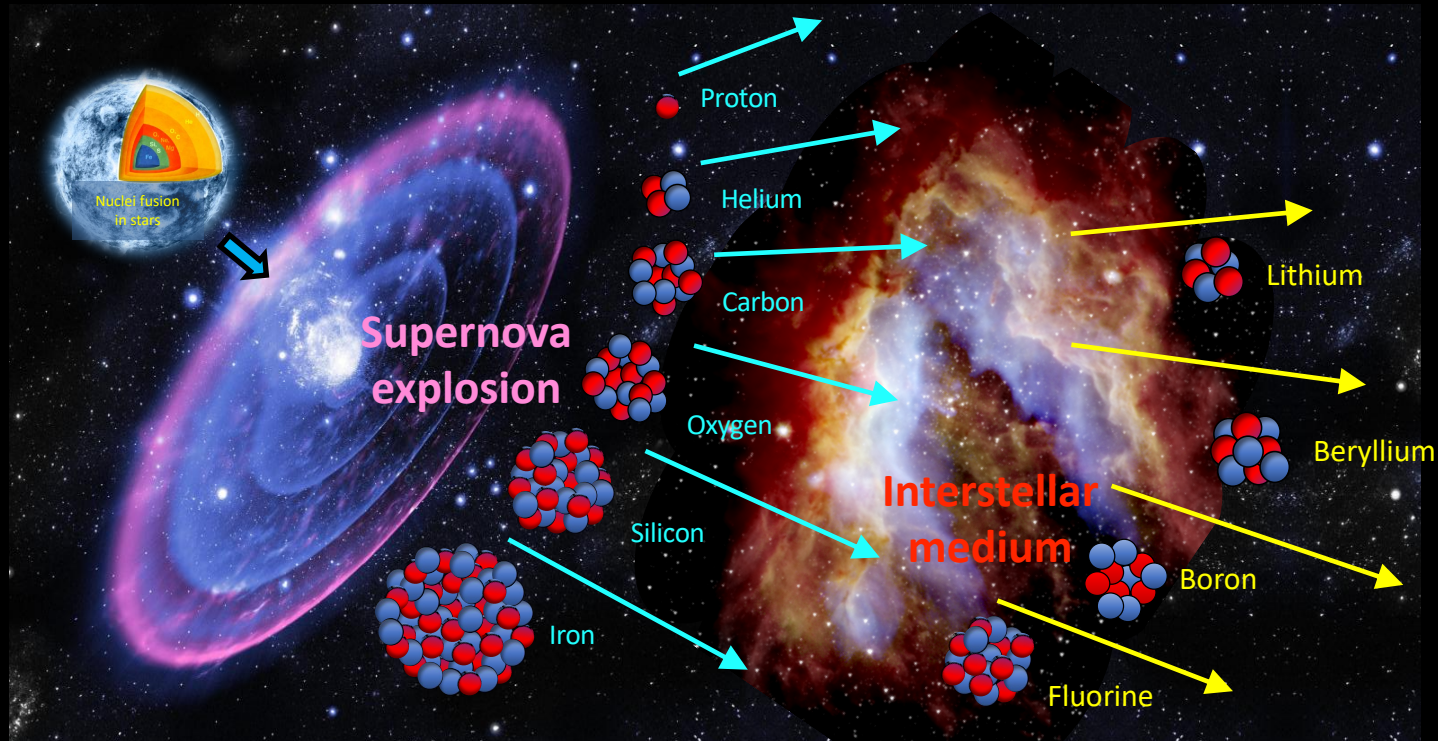


By 2030, we will provide a much more sensitive comparison among Fe, Ni, and Zn.



# Latest AMS Results on Secondary Cosmic Ray Nuclei

Secondary Li, Be, B, and F nuclei in cosmic rays are produced by the collision of primary cosmic ray C, O, Ne, Mg, Si, ..., Fe with the interstellar medium.

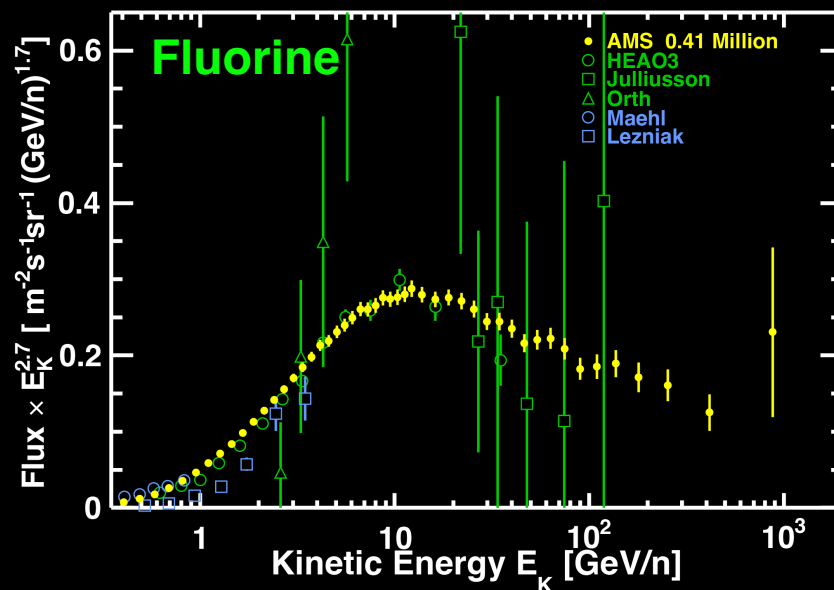
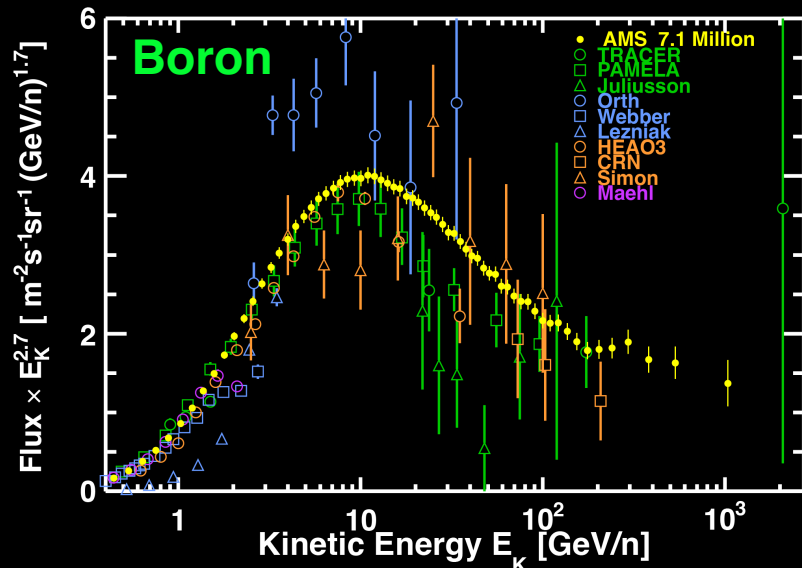
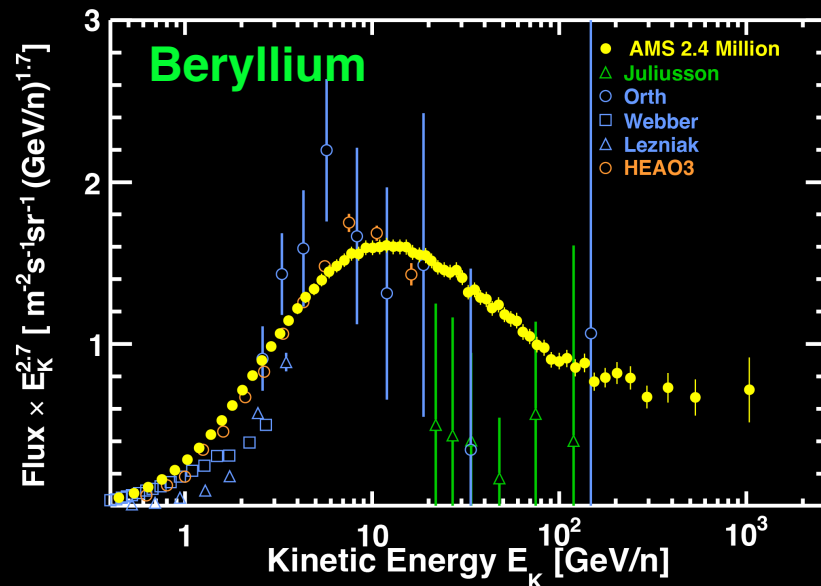
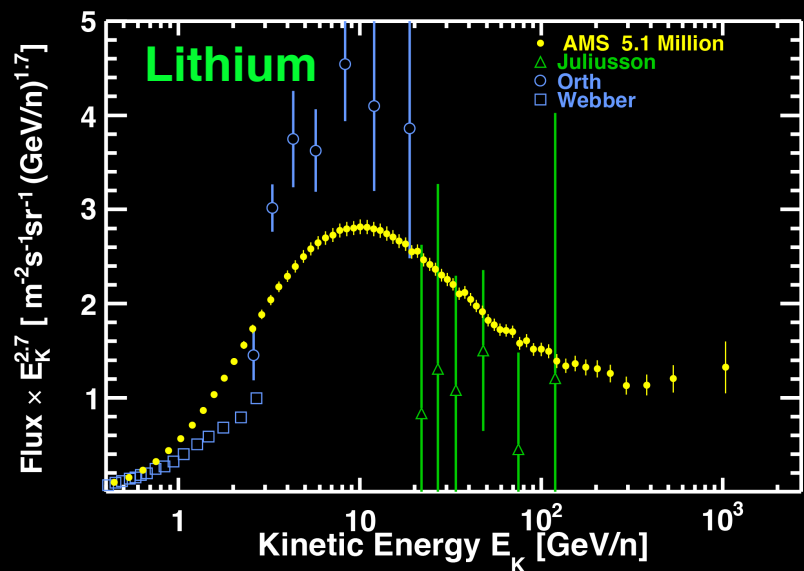


Measurements of the secondary cosmic ray nuclei fluxes are important in understanding the propagation of cosmic rays in the Galaxy.

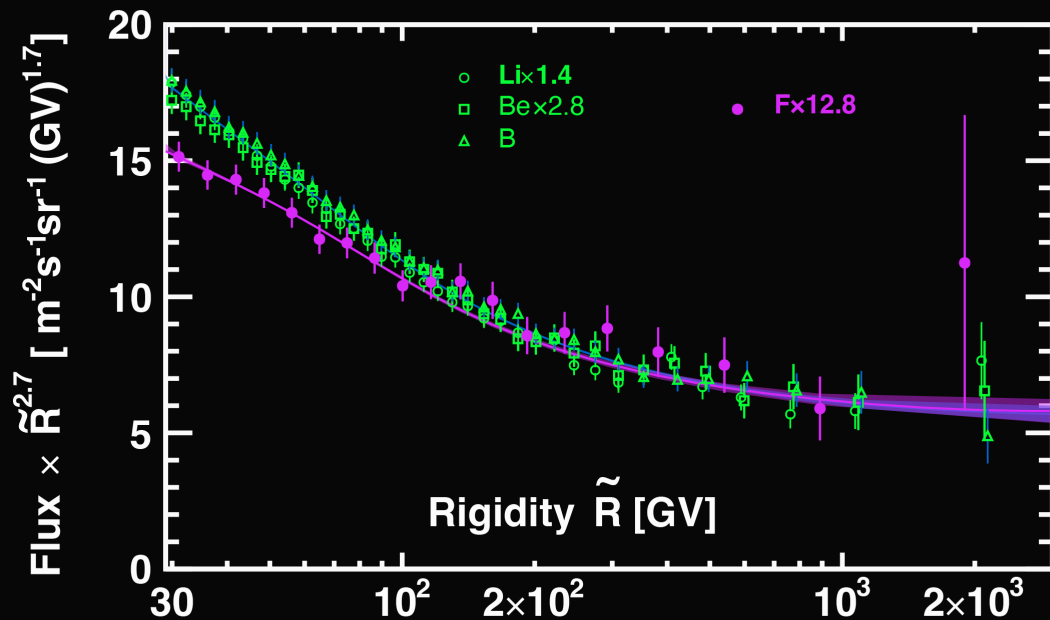
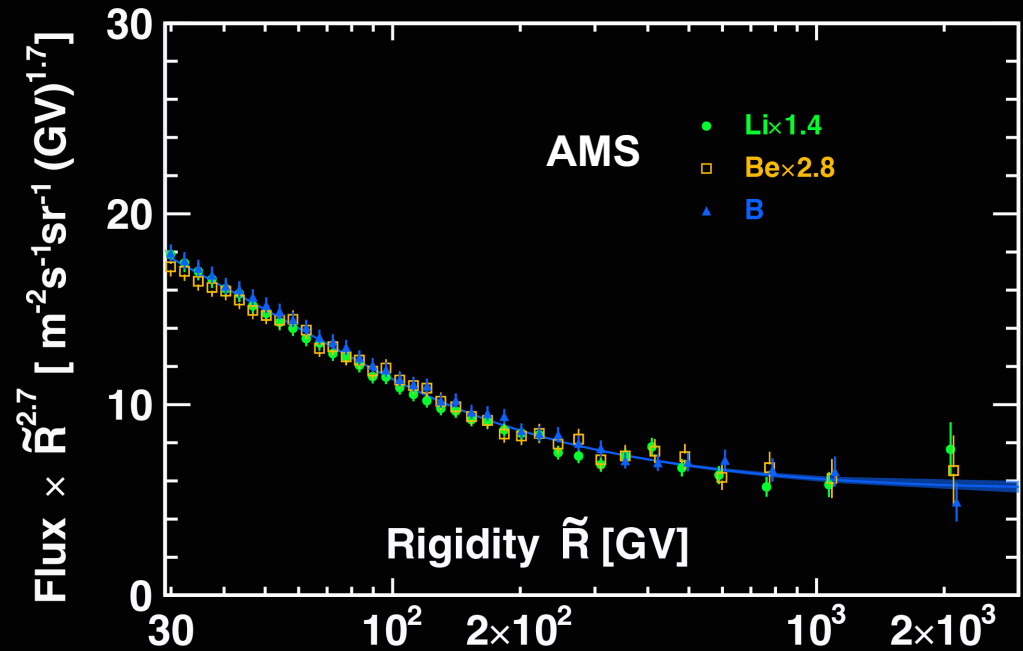
**Lithium, Beryllium, Boron, and Fluorine on Earth are produced by cosmic rays.**



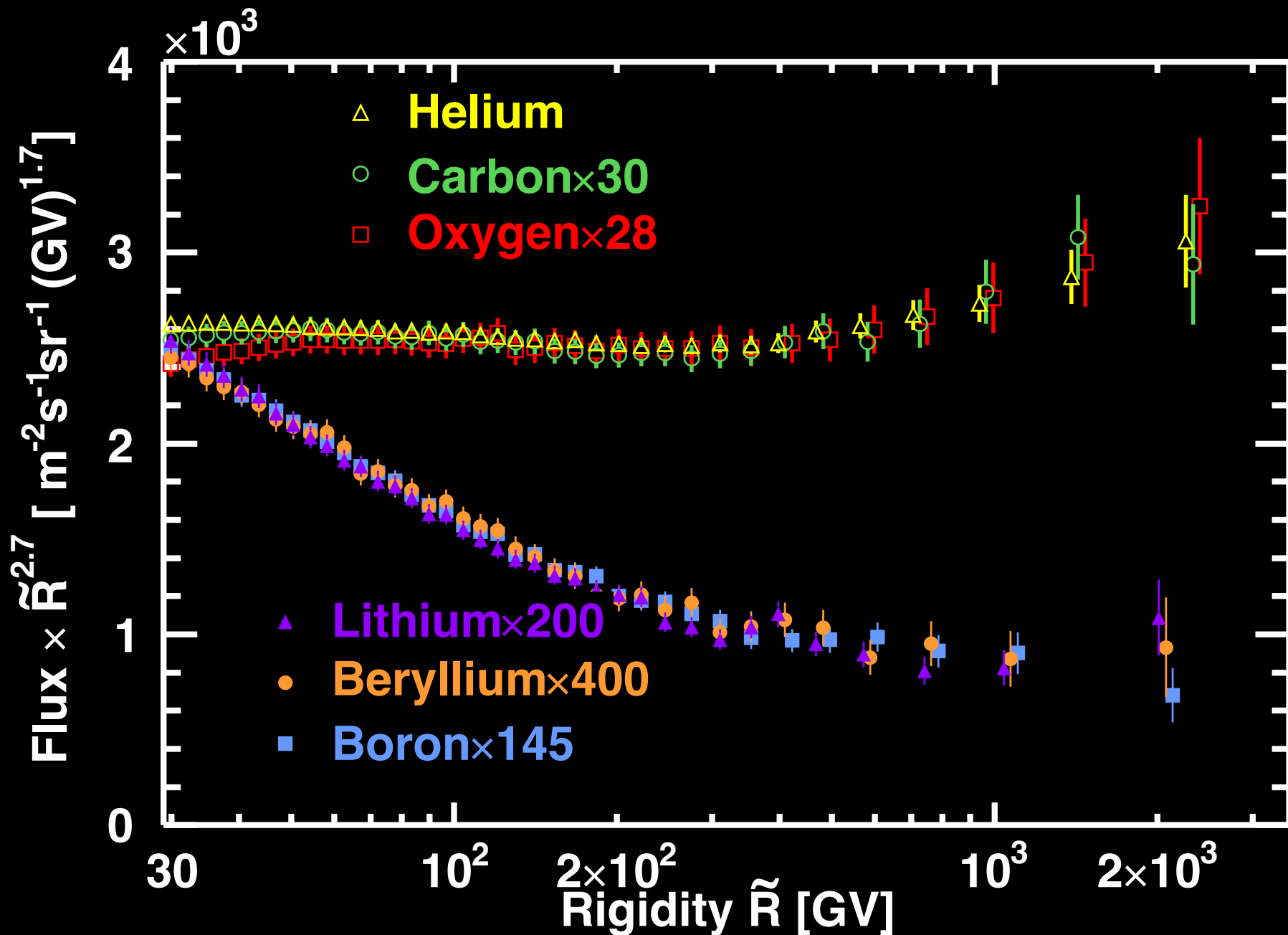
# AMS Secondary nuclei results compared with earlier measurements



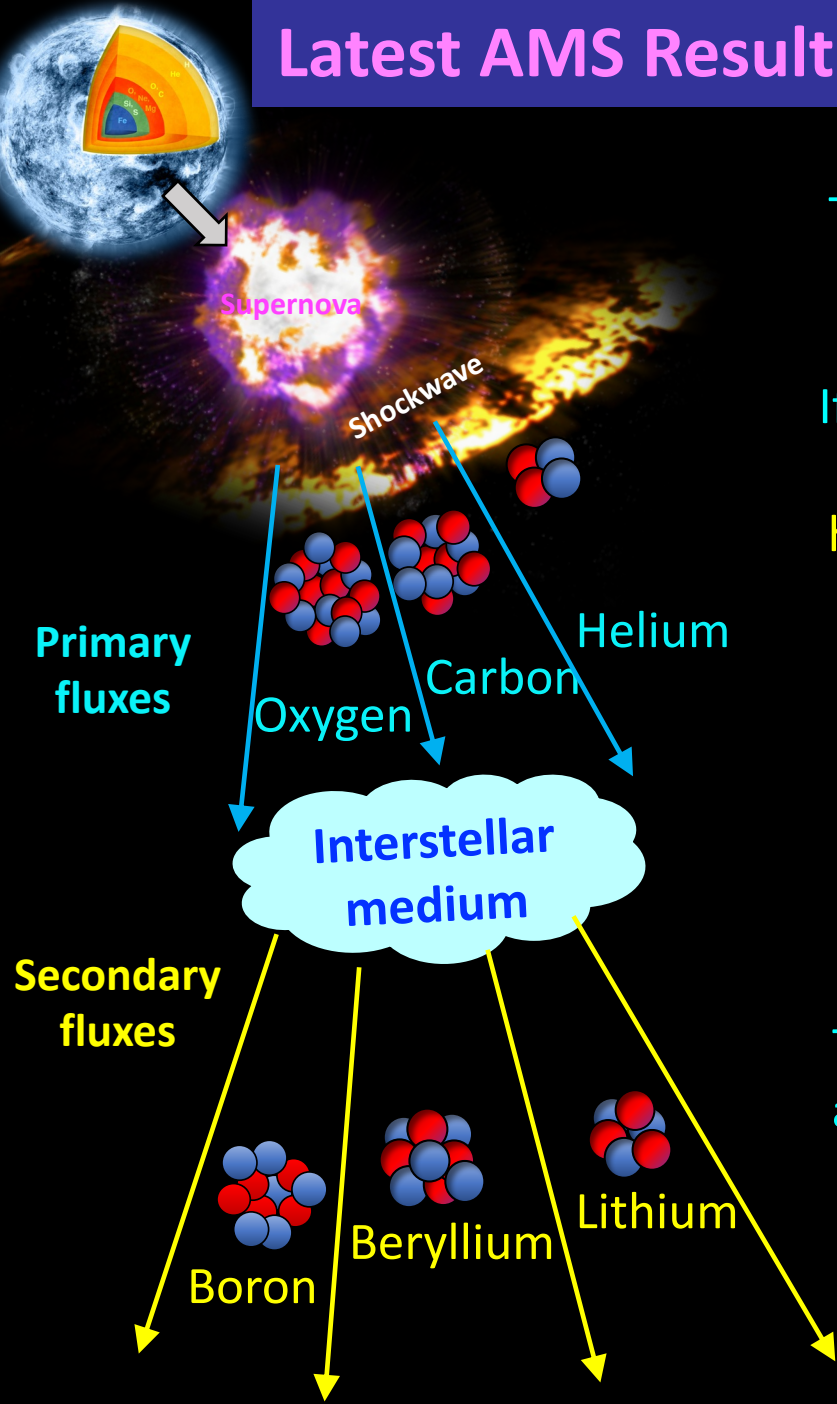
# Secondary cosmic rays also have two classes of rigidity dependence



# He-C-O primaries compared with Li-Be-B secondaries



# Latest AMS Results on Secondary-to-Primary Ratios



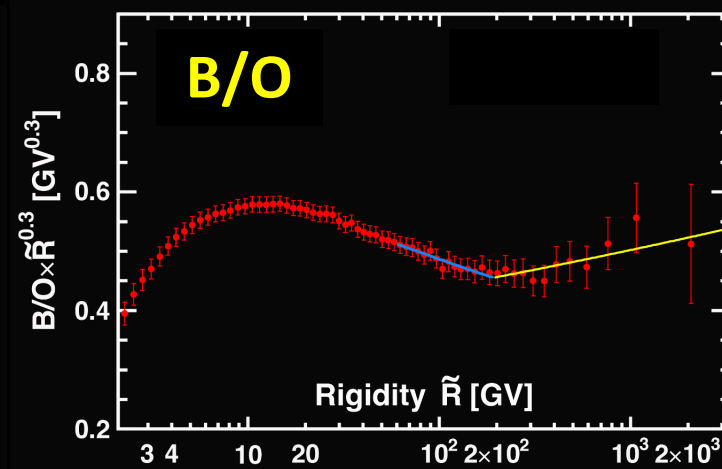
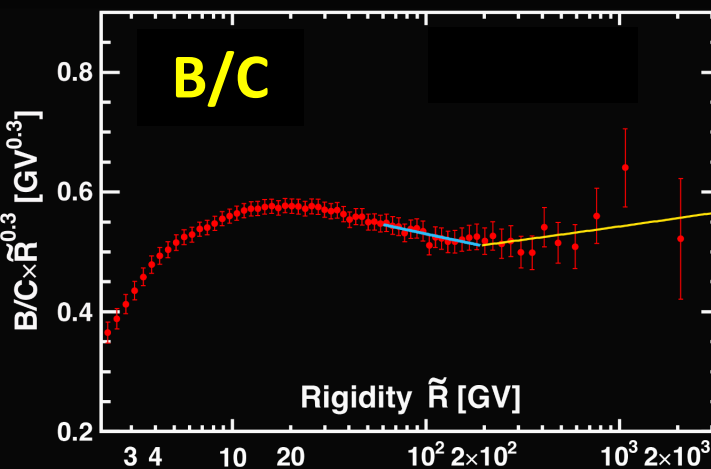
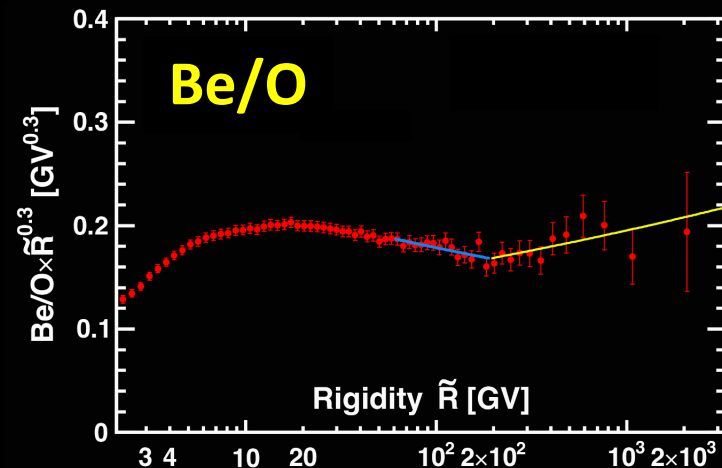
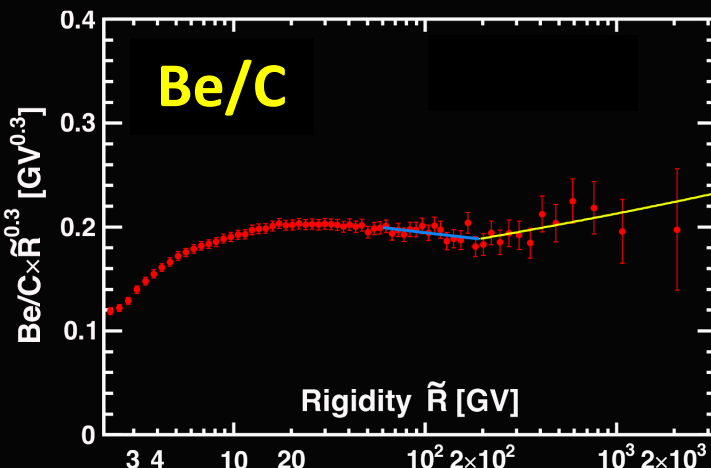
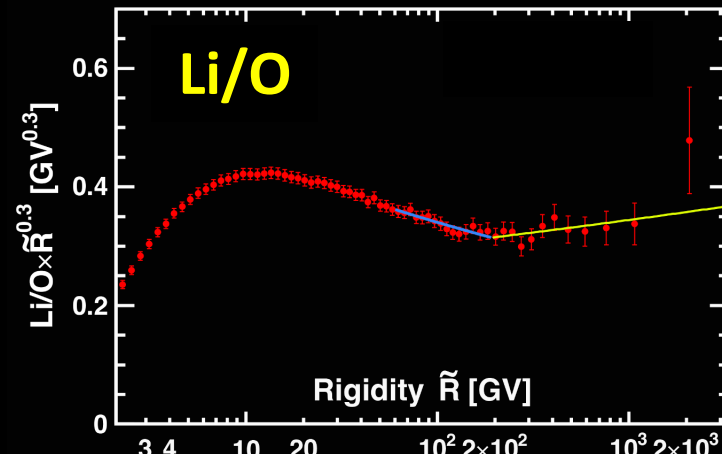
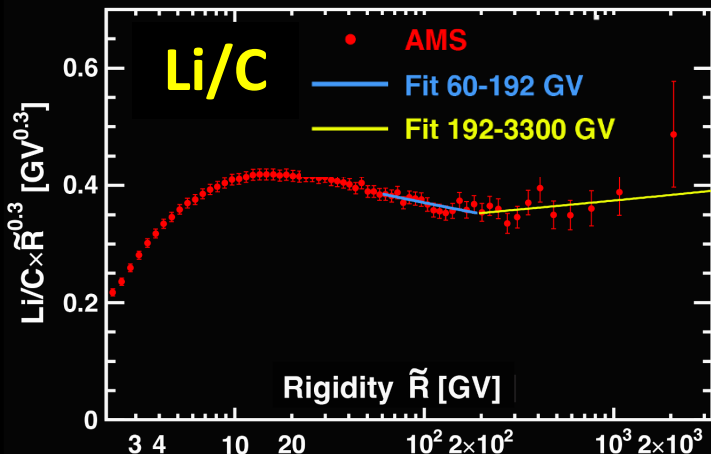
There are many theoretical models describing the behavior of cosmic rays.

If the **hardening** in cosmic rays is related to the injected spectra at their source, then similar **hardening** is expected for both secondary and primary cosmic rays.

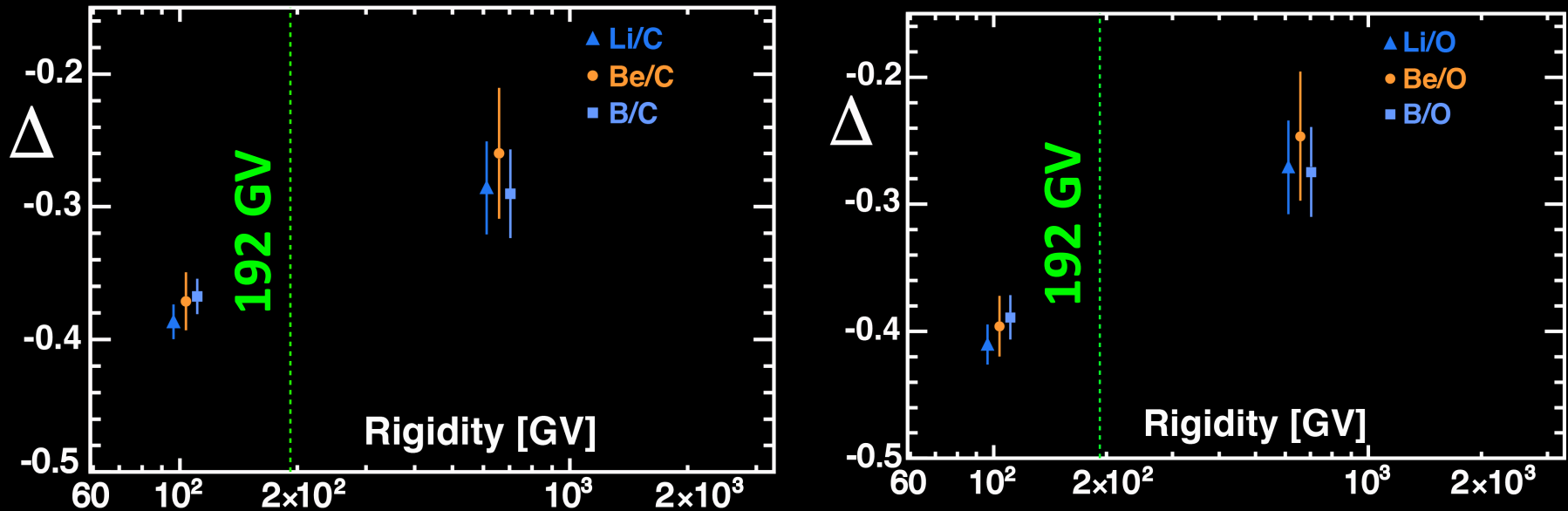
If the **hardening** is related to propagation properties in the Galaxy, then a stronger **hardening** is expected for the secondary with respect to the primary cosmic rays.

The theoretical models have their limitations, as none of them predicted the AMS observed spectral behavior of the primary cosmic rays He, C, and O nor the secondary cosmic rays Li, Be, and B.

# Secondary -to- Primary Ratios



# Secondary/Primary Ratios ( $kR^\Delta$ ) are rigidity dependent



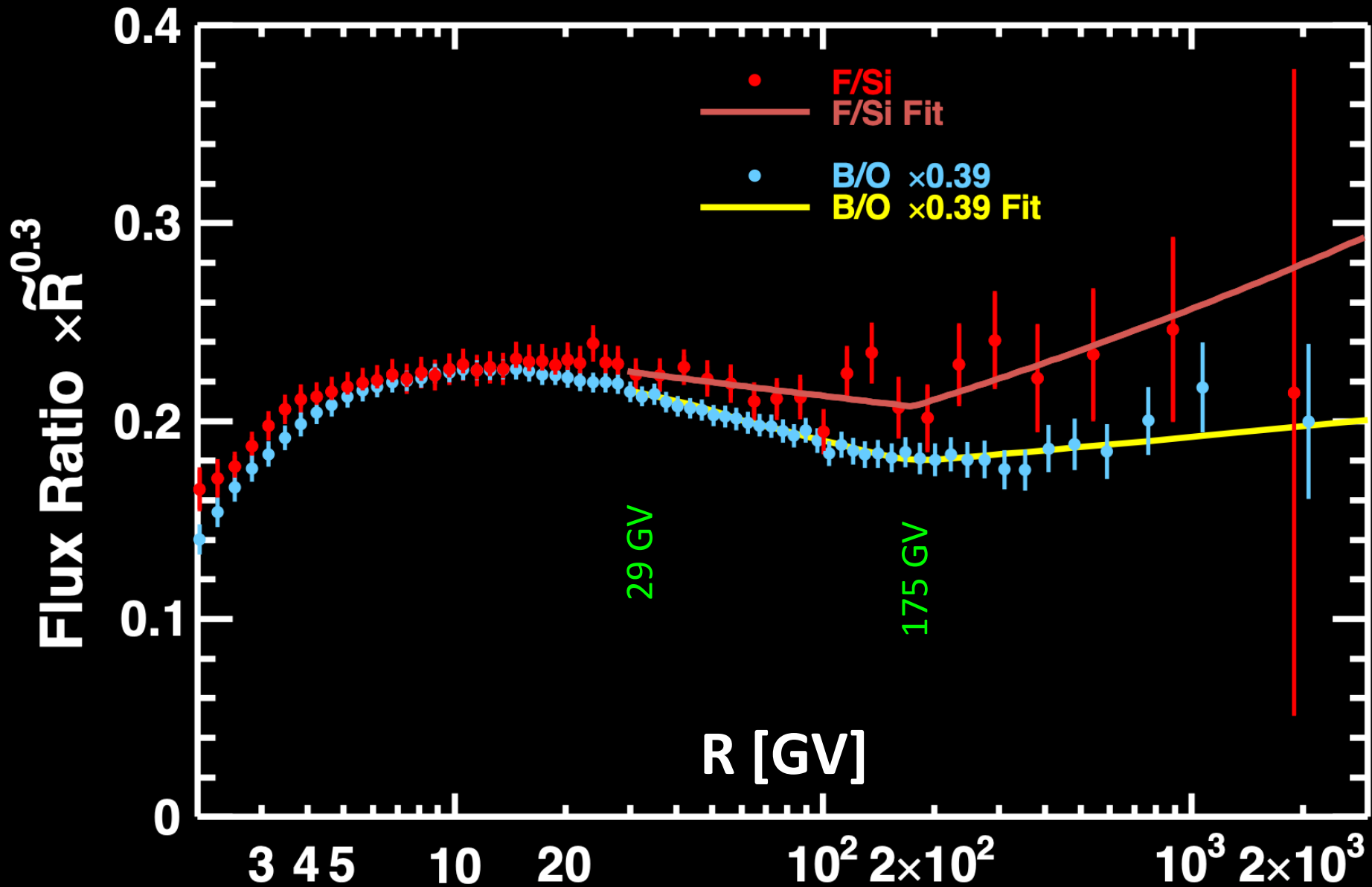
$\Delta$  in two rigidity intervals (60 – 192 GV and 192 – 3300 GV) exhibit an average hardening of  $0.11 \pm 0.02$ .

The significance of this change is  $5.5\sigma$ .

Above  $\sim 200$  GV secondary cosmic rays harden twice as much as primaries.

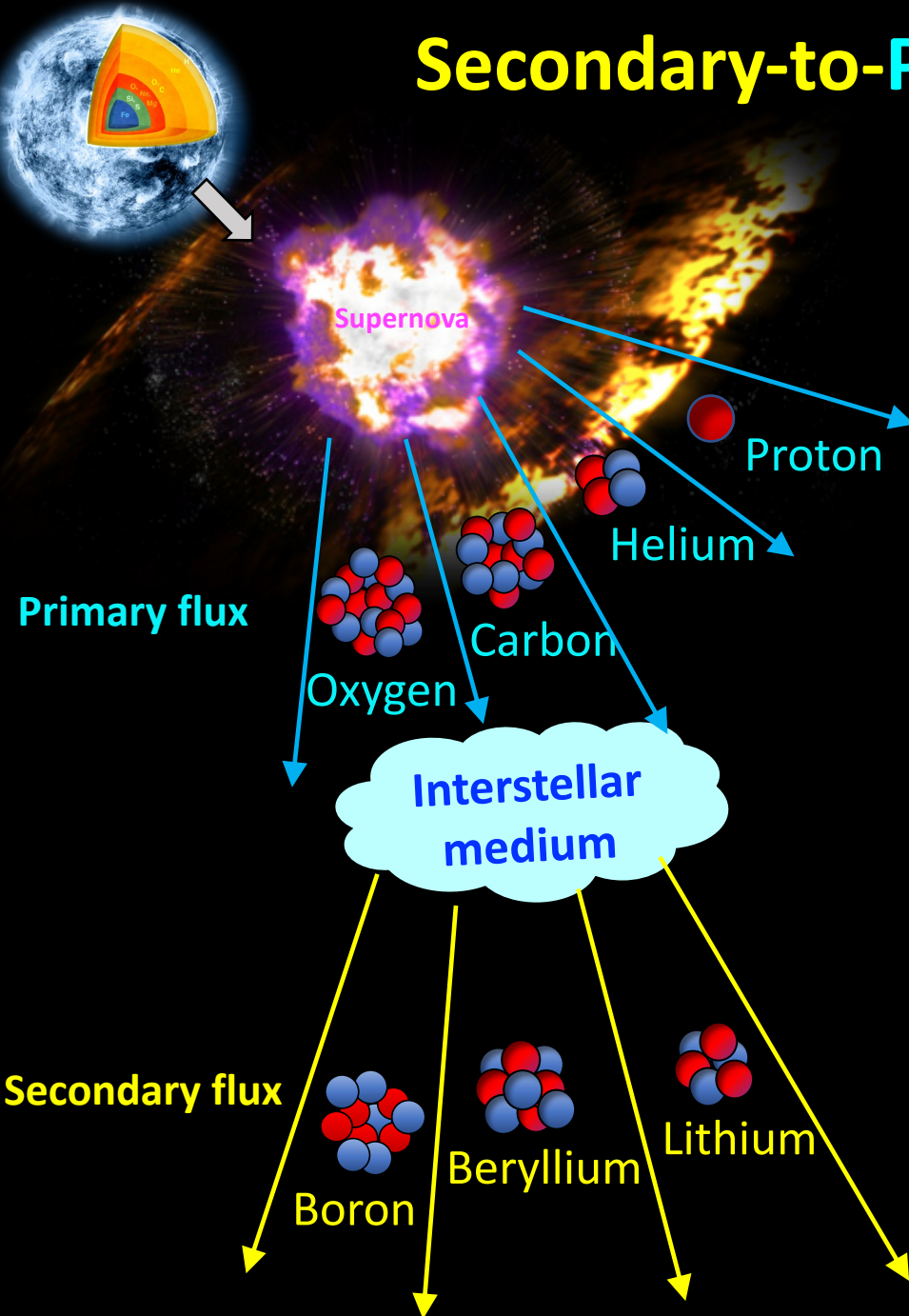
This strongly supports that the hardening is related to propagation properties in the Galaxy.

# F/Si (high-Z) compared to B/O (low-Z)



The Secondary-to-Primary ratios are different for high-Z and low-Z

# Secondary-to-Primary Ratios



Before AMS, the secondary-to-primary ratios (B/C ...) were assumed to be  $\propto R^\Delta$  with  $\Delta$  a constant (independent of  $R$  and  $Z$ ).

**AMS results:**

$\Delta$  depends on  $Z$

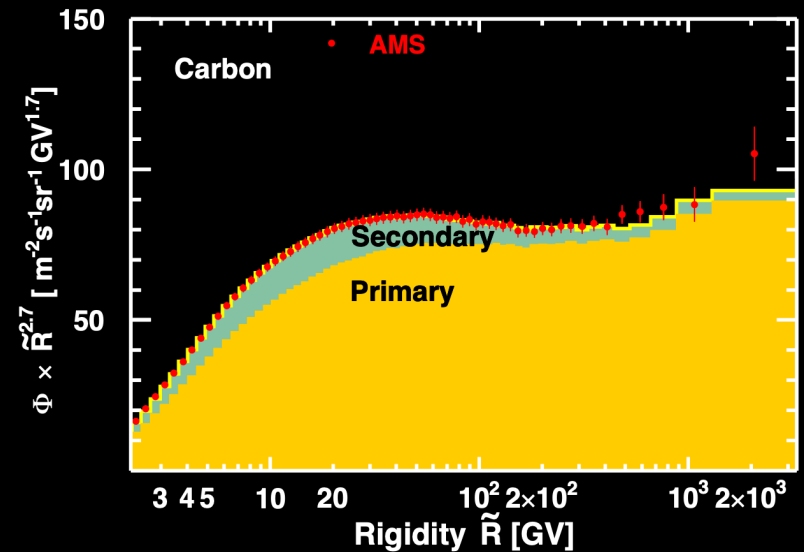
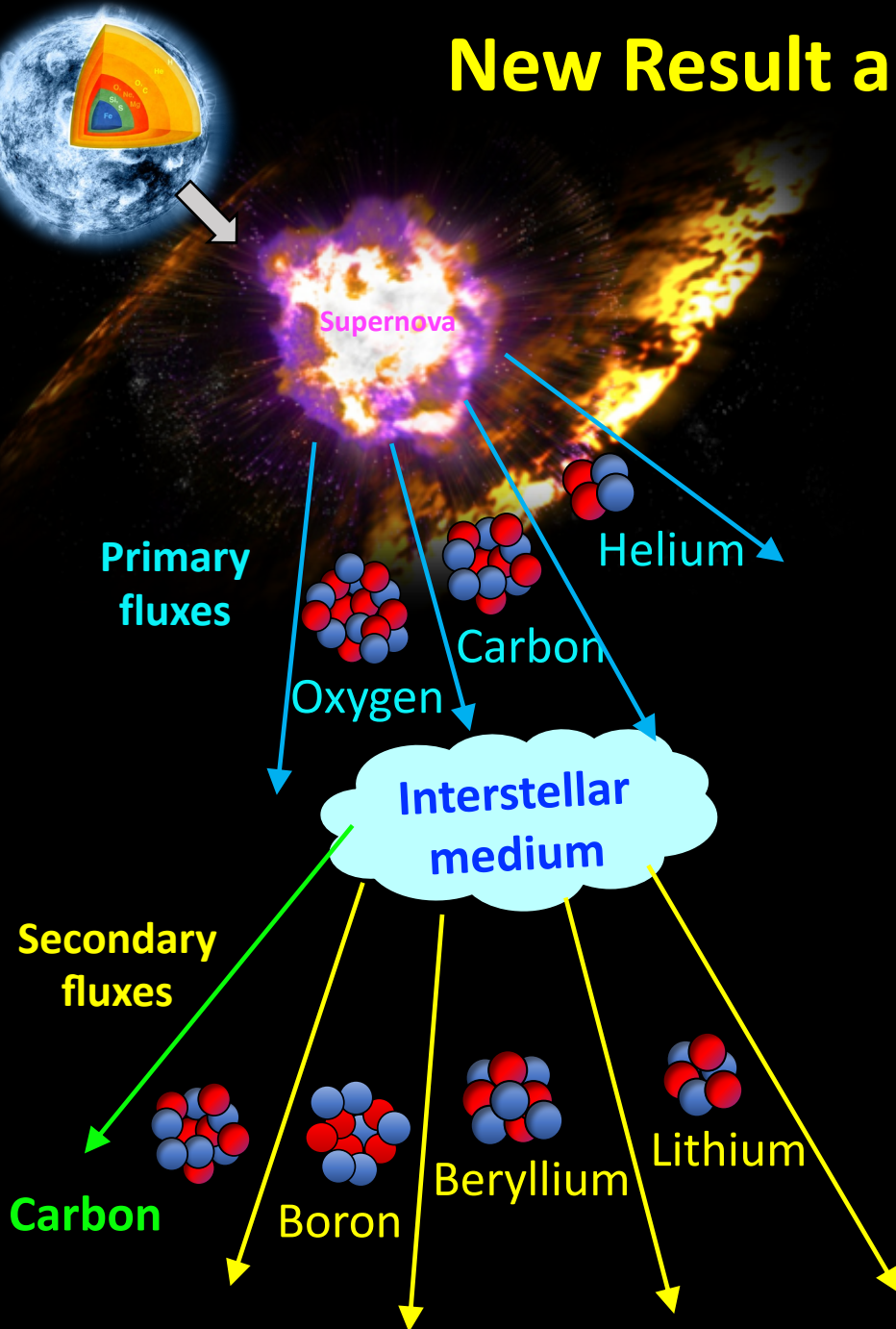
$\Delta$  depends on  $R$

The traditional B/C measurement does not describe the Interstellar Medium



# New Result and its Implications

Before AMS, taking into account the long-standing idea that **C is pure primary** and **B is pure secondary**, the **(B/C) ratio** has been used in models to describe cosmic ray propagation



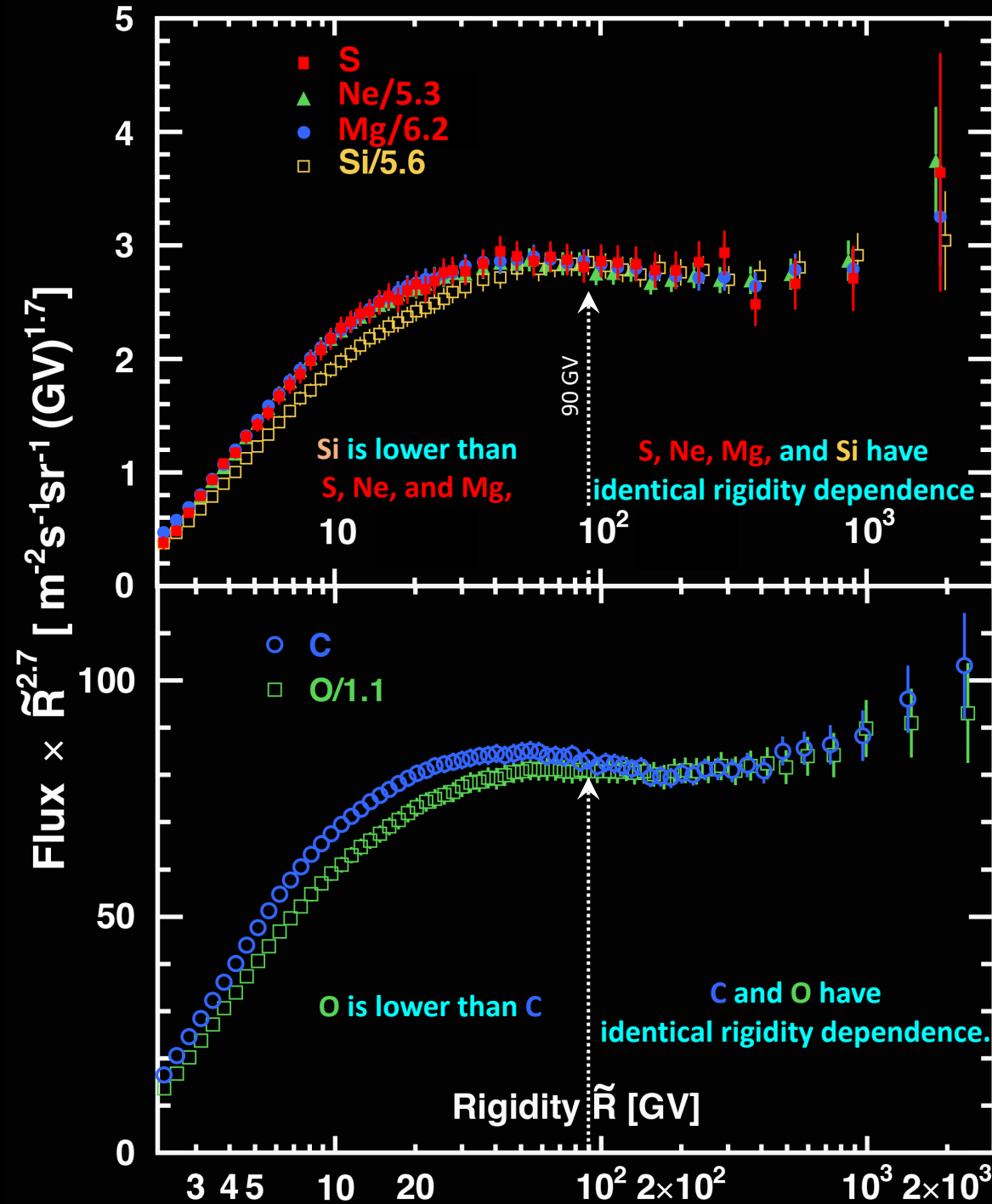
**But C is NOT pure primary.**  
Question: how to use **(B/C)** in cosmic ray models?

# Examples of Theoretical Papers based on Boron to Carbon Ratio

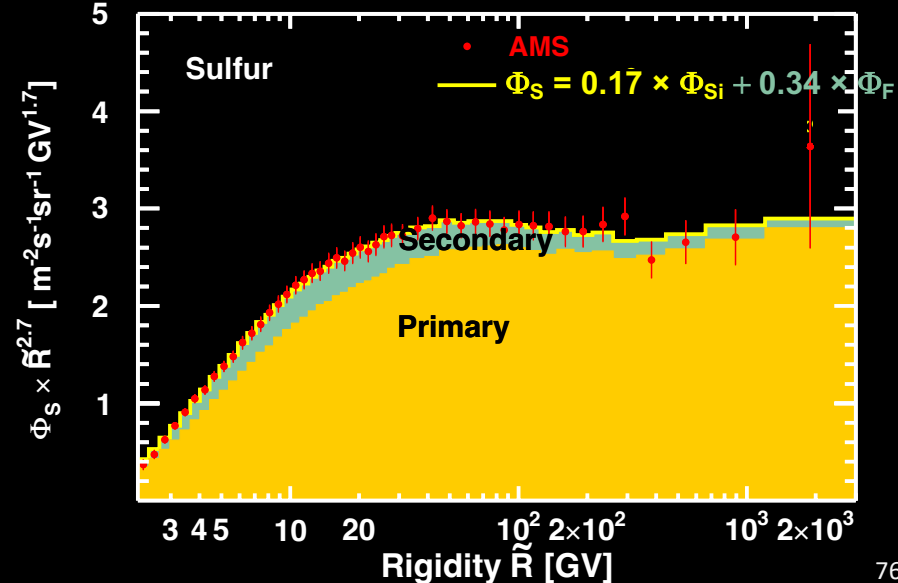
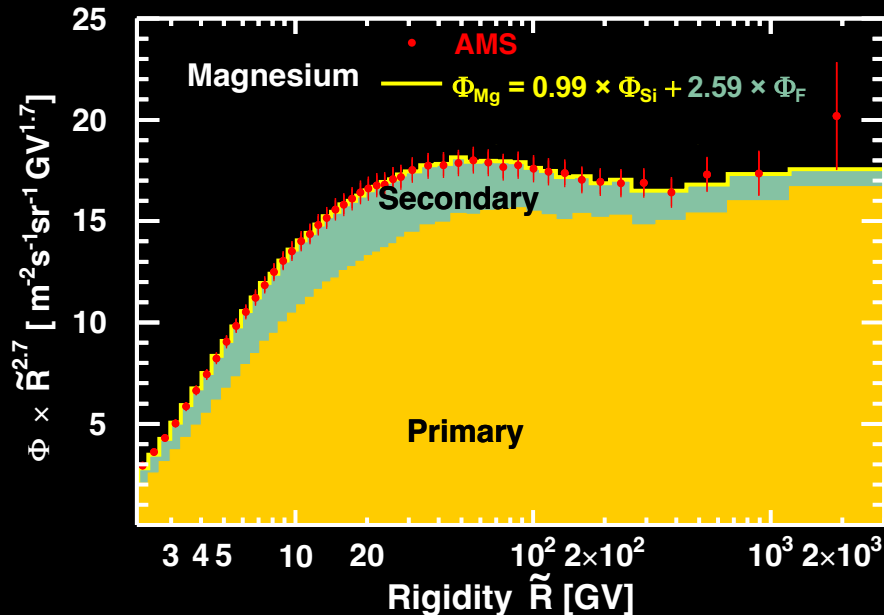
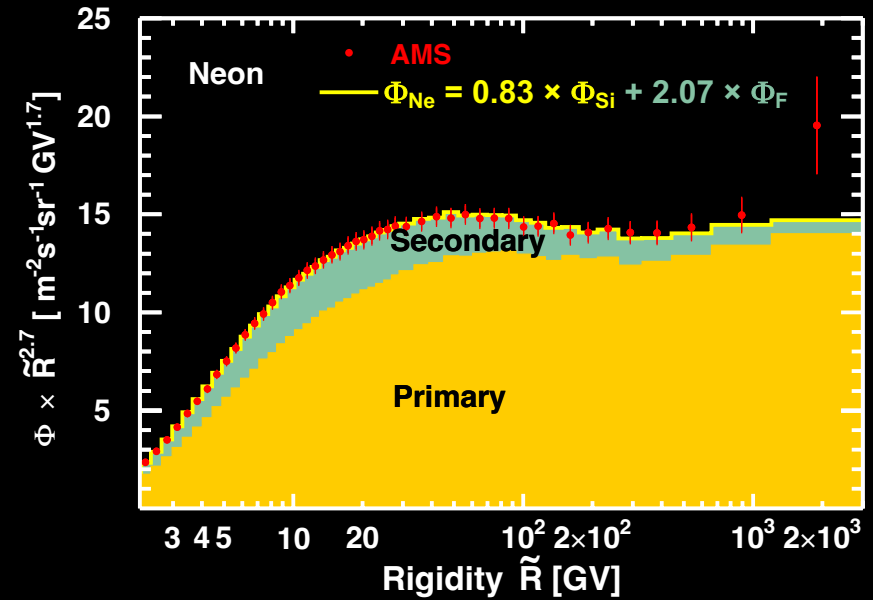
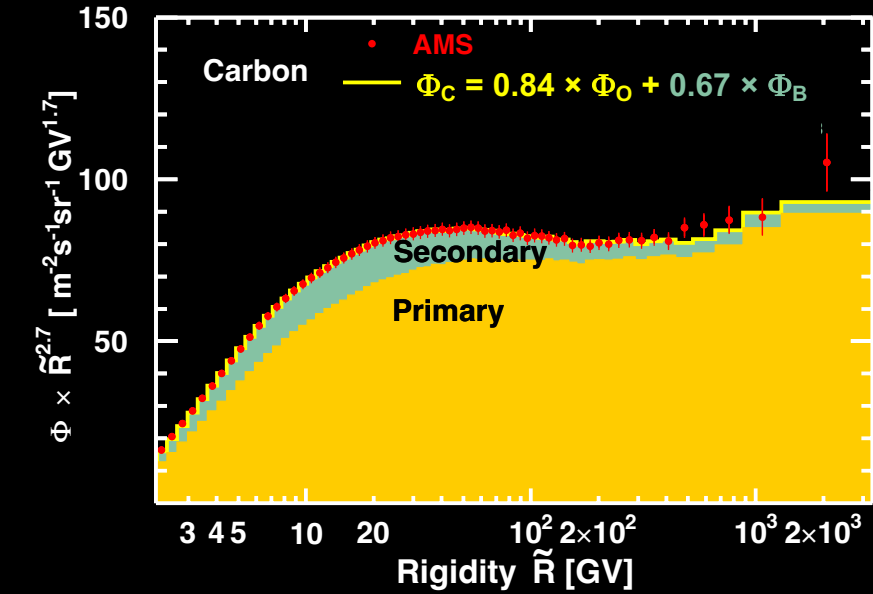
- 1) P.-p. Zhang, et al., "Ultra-high-energy diffuse gamma-ray emission from cosmic-ray interactions with the medium surrounding acceleration sources", *Phys. Rev. D* 105, 023002 (2022).
- 2) R. C. Anjos and F. Catalani, "Galactic Center as an efficient source of cosmic rays", *Phys. Rev. D* 101, 123015 (2020).
- 3) R. Diesing and D. Caprioli, "Nonsecondary origin of cosmic ray positrons", *Phys. Rev. D* 101, 103030 (2020).
- 4) Y. Génolini et al., "Cosmic-ray transport from **AMS-02 boron to carbon ratio** data: Benchmark models and interpretation", *Phys. Rev. D* 99, 123028 (2019).
- 5) L. Derome et al., "Fitting **B/C** cosmic-ray data in the **AMS-02** era: a cookbook", *Astron. Astrophys.* 627, A158 (2019).
- 6) I. Cholis, T. Linden, and D. Hooper, "A robust excess in the cosmic-ray antiproton spectrum: Implications for annihilating dark matter", *Phys. Rev. D* 99, 103026 (2019).
- 7) M. Kachelrieß, A. Neronov, and D. V. Semikoz, "Cosmic ray signatures of a 2-3 Myr old local supernova", *Phys. Rev. D* 97, 063011 (2018).
- 8) G. Giacinti, M. Kachelrieß and D.V. Semikoz, "Reconciling cosmic ray diffusion with Galactic magnetic field models", *J. Cosmol. Astropart. Phys.* 07 (2018) 051.
- 9) É. Jaupart, É. Parizot, and D. Allard, "Contribution of the Galactic centre to the local cosmic-ray flux", *Astron. Astrophys.* 619, A12 (2018).
- 10) Y. Génolini et al., "Indications for a High-Rigidity Break in the Cosmic-Ray Diffusion Coefficient", *Phys. Rev. Lett.* 119, 241101 (2017).
- 11) D. Eichler, "An Alternative Explanation of the Varying **Boron-to-carbon Ratio** in Galactic Cosmic Rays", *Astrophys. J.* 842, 50 (2017).
- 12) Q. Yuan et al., "Propagation of cosmic rays in the **AMS-02** era", *Phys. Rev. D* 95, 083007 (2017).
- 13) N. Tomassetti, "Solar and nuclear physics uncertainties in cosmic-ray propagation", *Phys. Rev. D* 96, 103005 (2017).
- 14) P. Lipari, "interpretation of the cosmic ray positron and antiproton fluxes", *Phys. Rev. D* 95, 063009 (2017).
- 15) P. Blasi, "On the spectrum of stable secondary nuclei in cosmic rays", *Mon. Not. R. Astron. Soc.* 471, 1662 (2017).
- 16) J. C. Joshi and S. Razzaque, "A self-consistent model of cosmic-ray fluxes and positron excess: roles of nearby pulsars and a sub-dominant source population", *JCAP* 09 (2017) 029.
- 17) D. Benyamin et al., "The **B/C** and Sub-Iron/Iron Cosmic Ray Ratios-Further Evidence in Favor of the Spiral-Arm Diffusion Model", *Astrophys. J.* 826, 1 (2016).
- 18) M. D'Angelo, P. Blasi, and E. Amato, "Grammage of cosmic rays around Galactic supernova remnants", *Phys. Rev. D* 94, 083003 (2016).
- 19) Y. Génolini et al., "Theoretical uncertainties in extracting cosmic-ray diffusion parameters: **the boron-to-carbon ratio**", *Astron. Astrophys.* 580, A9 (2015).
- 20) H. B. Jin, Y. L. Wu, and Y. F. Zhou, "Cosmic ray propagation and dark matter in light of the latest **AMS-02** data", *J. Cosmol. Astropart. Phys.* 09 (2015) 049.
- 21) N. Tomassetti and F. Donato, "The connection between the positron fraction anomaly and the spectral features in galactic cosmic-ray hadrons" *Astrophys. J. Lett.* 803, L15 (2015).
- 22) R. Aloisio, P. Blasi, and P. D. Serpico, "Nonlinear cosmic ray Galactic transport in the light of **AMS-02** and **Voyager** data", *Astron. Astrophys.* 583, A95 (2015).
- 23) I. Cholis and D. Hooper, "Constraining the origin of the rising cosmic ray positron fraction with **the boron-to-carbon ratio**", *Phys. Rev. D* 89, 043013 (2014).
- 24) D. Benyamin et al., "Recovering the observed **B/C ratio** in a dynamic spiral-armed cosmic ray model", *Astrophys. J.* 782, 34 (2014).
- 25) P. Mertsch and S. Sarkar, "**AMS-02** data confront acceleration of cosmic ray secondaries in nearby sources", *Phys. Rev. D* 90, 061301 (2014).
- 26) M. Kachelrieß and S. Ostapchenko, "**B/C ratio** and the PAMELA positron excess", *Phys. Rev. D* 87, 047301 (2013).
- 27) A. Obermeier et al., "The **boron-to-carbon abundance ratio** and galactic propagation of cosmic radiation", *Astrophys. J.* 752, 69 (2012).
- 28) N. Tomassetti and F. Donato, "Secondary cosmic-ray nuclei from supernova remnants and constraints on the propagation parameters", *Astron. Astrophys.* 544, A16 (2012).
- 29) D. Maurin, et al., "Systematic uncertainties on the cosmic-ray transport parameters Is it possible to reconcile **B/C** data with  $\delta = 1/3$  or  $\delta = 1/2$ ?", *Astron. Astrophys.* 516, A67 (2010).
- 30) R. Cowsik and B. Burch, "Positron fraction in cosmic rays and models of cosmic-ray propagation", *Phys. Rev. D* 82, 023009 (2010).
- 31) A. Shalchi and I. Büsching, "Influence of turbulence dissipation effects on the propagation of low-energy cosmic rays in the galaxy", *Astrophys. J.* 725, 2110 (2010).
- 32) P. Mertsch and S. Sarkar, "Testing Astrophysical Models for the PAMELA Positron Excess with Cosmic Ray Nuclei", *Phys. Rev. Lett.* 103, 081104 (2009).
- 33) C. Evoli, et al., "Cosmic ray nuclei, antiprotons and gamma rays in the galaxy: a new diffusion model", *J. Cosmol. Astropart. Phys.* 10 (2008) 018.
- 34) A. Castellina and F. Donato, "Diffusion coefficient and acceleration spectrum from direct measurements of charged cosmic ray nuclei", *Astropart. Phys.* 24, 146 (2005).
- 35) I. Büsching et al., "Cosmic-Ray Propagation Properties for an Origin in Supernova Remnants", *Astrophys. J.* 619, 314 (2005).
- 36) M. S. Potgieter and U. W. Langner, "The heliospheric modulation of cosmic ray **boron and carbon**", *Ann. Geophys.* 22, 3729 (2004).
- 37) I. V. Moskalenko et al., "Challenging Cosmic-Ray Propagation with Antiprotons: Evidence for a "Fresh" Nuclei Component?", *Astrophys. J.* 586, 1050 (2003).
- 38) M. Maurin, R. Taillet, and F. Donato, "New results on source and diffusion spectral features of Galactic cosmic rays: **I- B/C ratio**" *Astron. Astrophys.* 394, 1039 (2002).
- 39) I. V. Moskalenko et al., "Secondary Antiprotons and Propagation of Cosmic Rays in the Galaxy and Heliosphere", *Astrophys. J.* 565, 280 (2002).
- 40) -----

# Detailed Properties of Primary Cosmic Rays

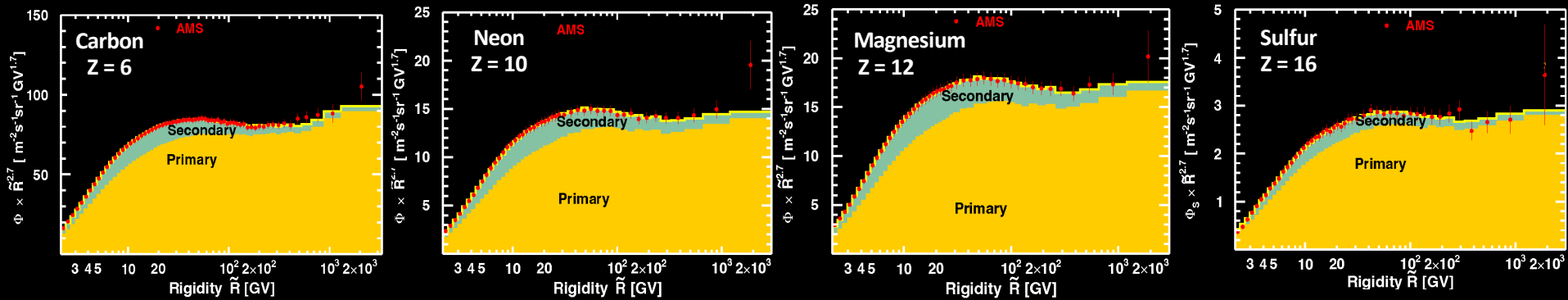
The discrepancies in functional behavior at low rigidity indicate that many primary cosmic rays have a significant secondary component



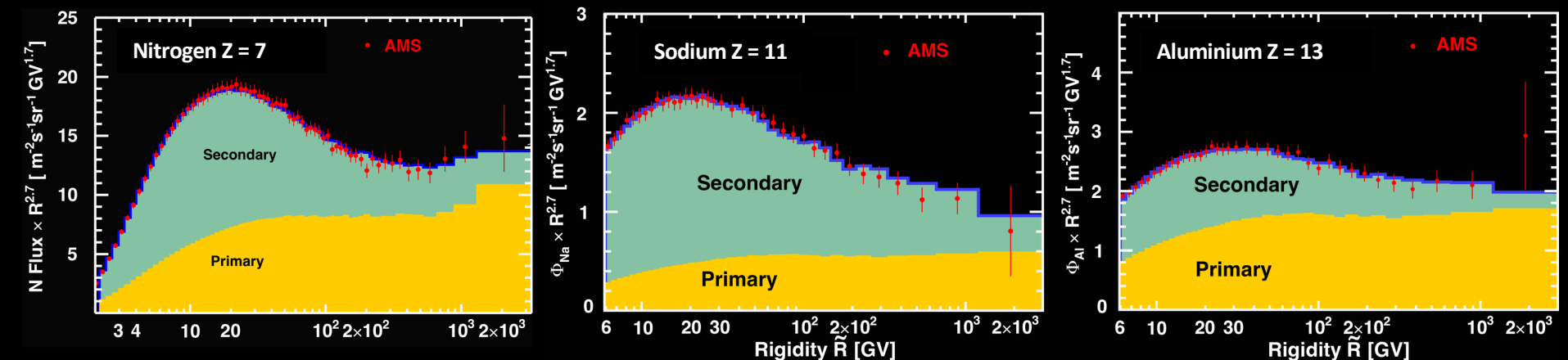
**New, unexpected observation:**  
**Traditional primary cosmic rays C, Ne, Mg, and S fluxes are not pure primary;**  
**they all have a significant secondary component**



# Even-Z nuclei and Odd-Z nuclei have distinctly different primary and secondary composition

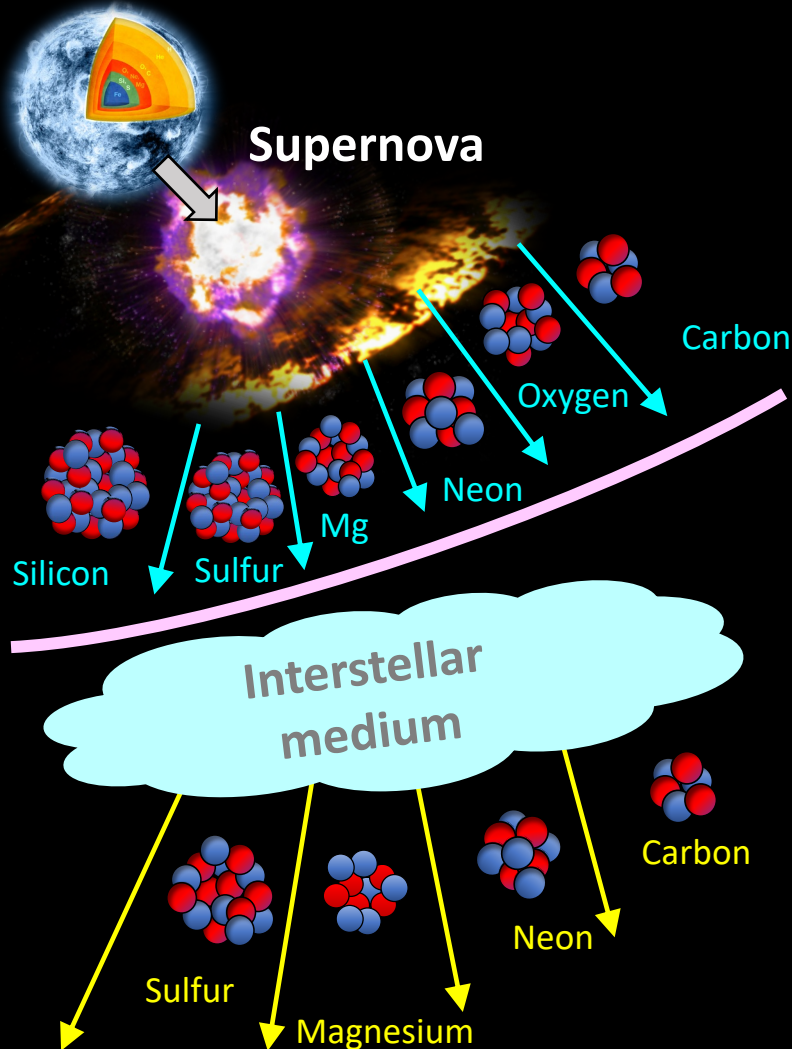


Even-Z nuclei are dominated by primaries



Odd-Z nuclei have more secondaries than even-Z

# Model-independent measurements of the relative abundances at the source (before cosmic ray propagation)



Abundance  
Ratio

Value at  
the Source

$$\Phi_{\text{C}} / \Phi_{\text{O}} \quad 0.836 \pm 0.025$$

$$\Phi_{\text{Ne}} / \Phi_{\text{Si}} \quad 0.833 \pm 0.025$$

$$\Phi_{\text{Mg}} / \Phi_{\text{Si}} \quad 0.994 \pm 0.029$$

$$\Phi_{\text{S}} / \Phi_{\text{Si}} \quad 0.167 \pm 0.006$$

$$\Phi_{\text{N}} / \Phi_{\text{O}} \quad 0.092 \pm 0.002$$

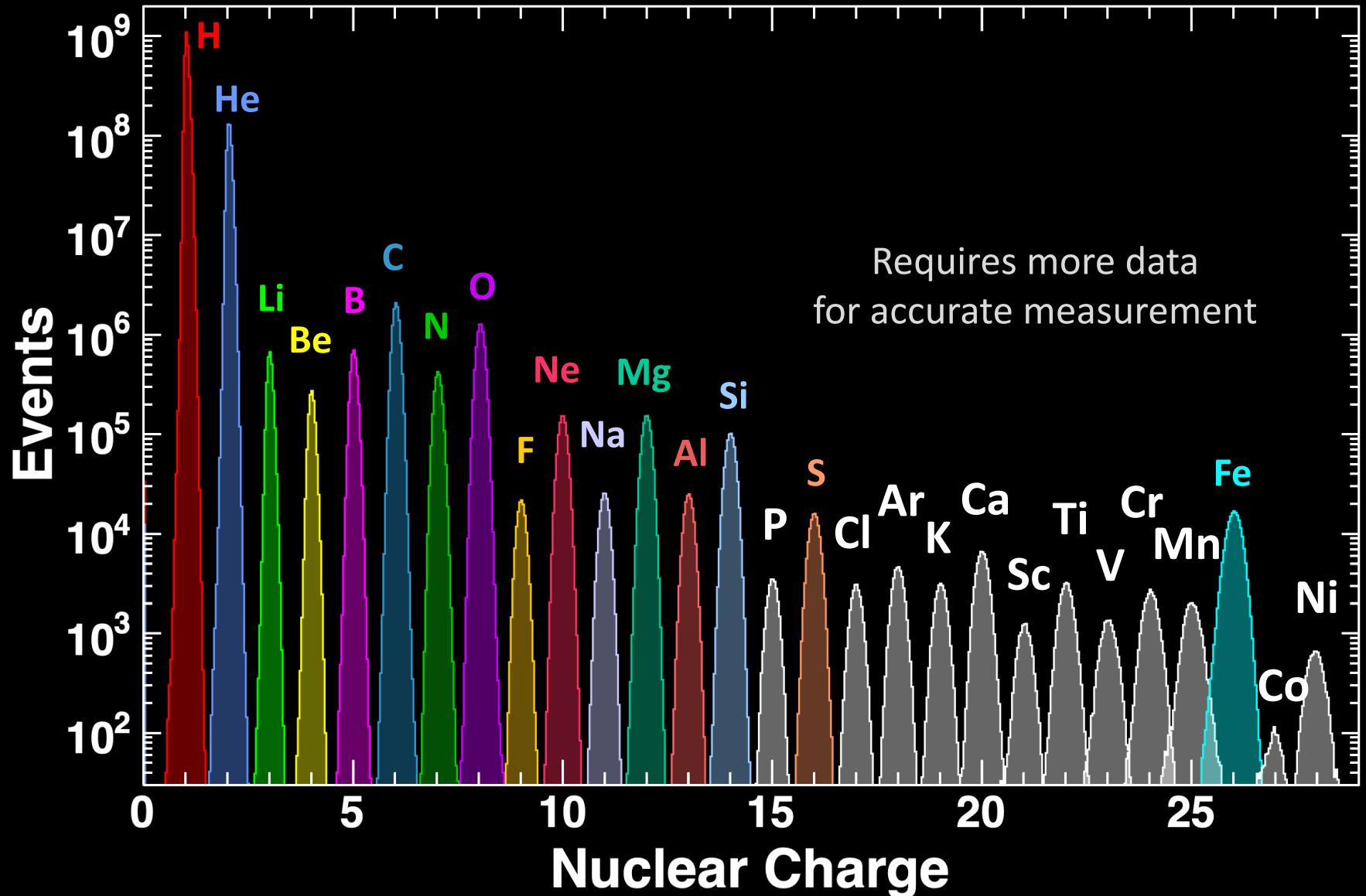
$$\Phi_{\text{Na}} / \Phi_{\text{Si}} \quad 0.036 \pm 0.003$$

$$\Phi_{\text{Al}} / \Phi_{\text{Si}} \quad 0.103 \pm 0.004$$

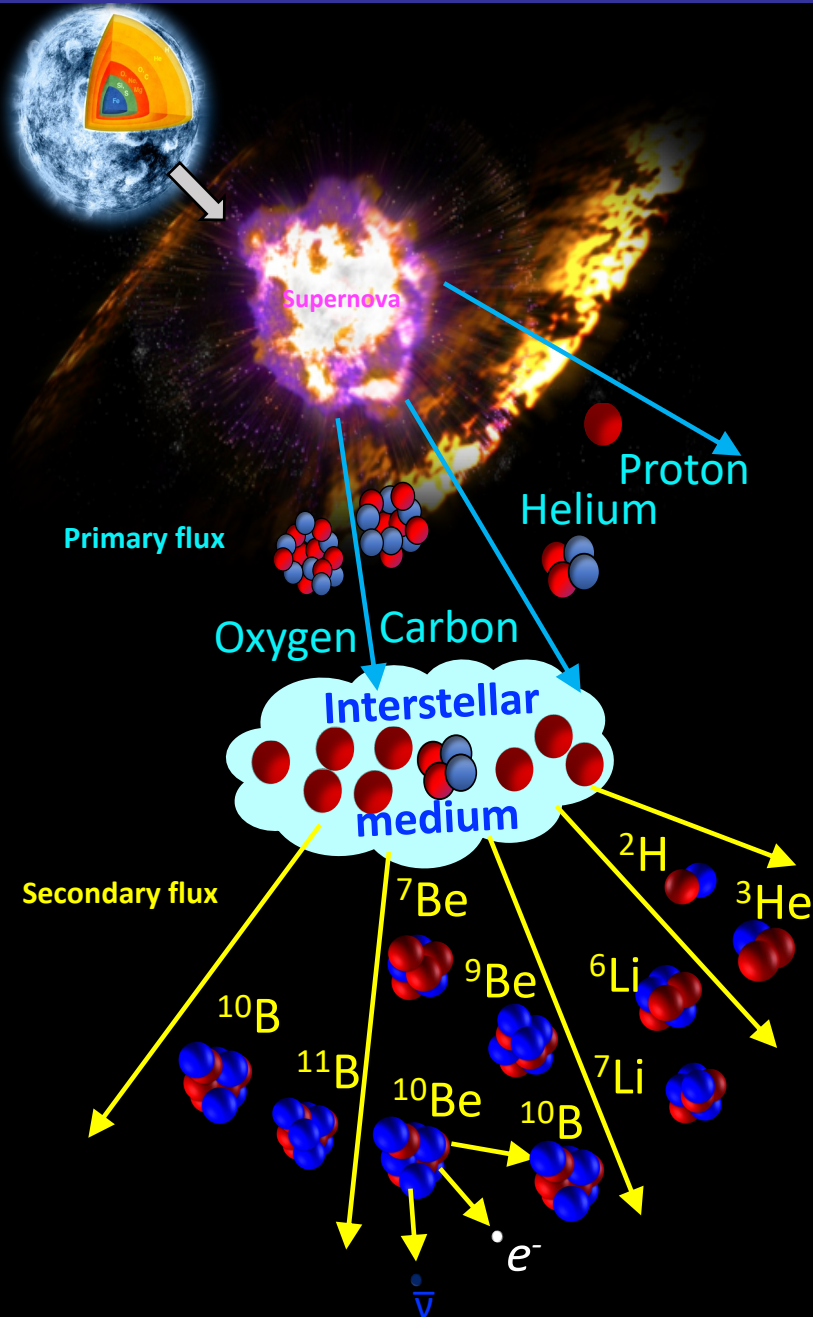
# Cosmic Ray Nuclei by 2030

AMS will provide complete and accurate spectra for the

29 elements and provide the foundation for a comprehensive theory of cosmic rays.

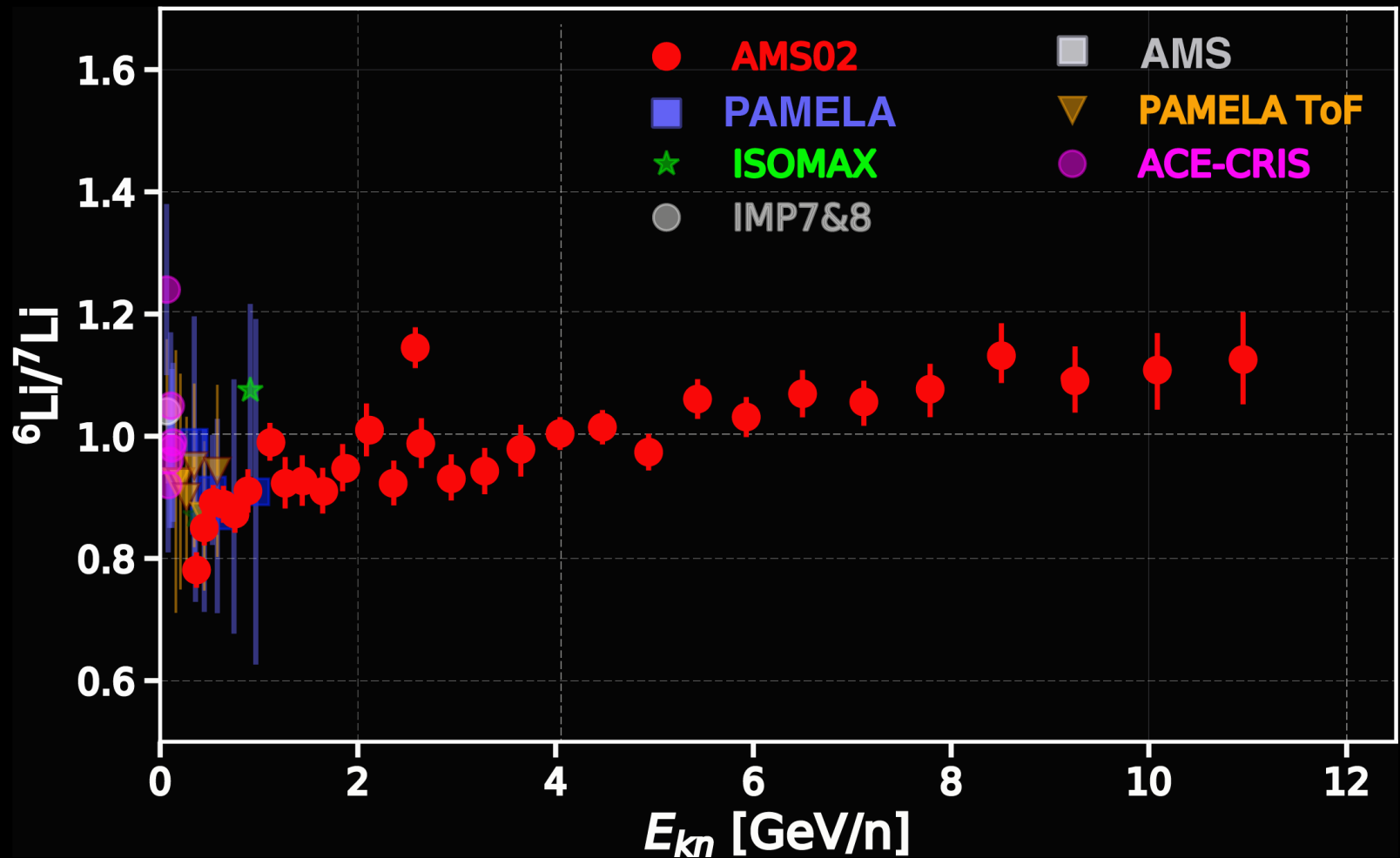


# Latest AMS Results on Cosmic Isotopes





# Lithium Isotopes



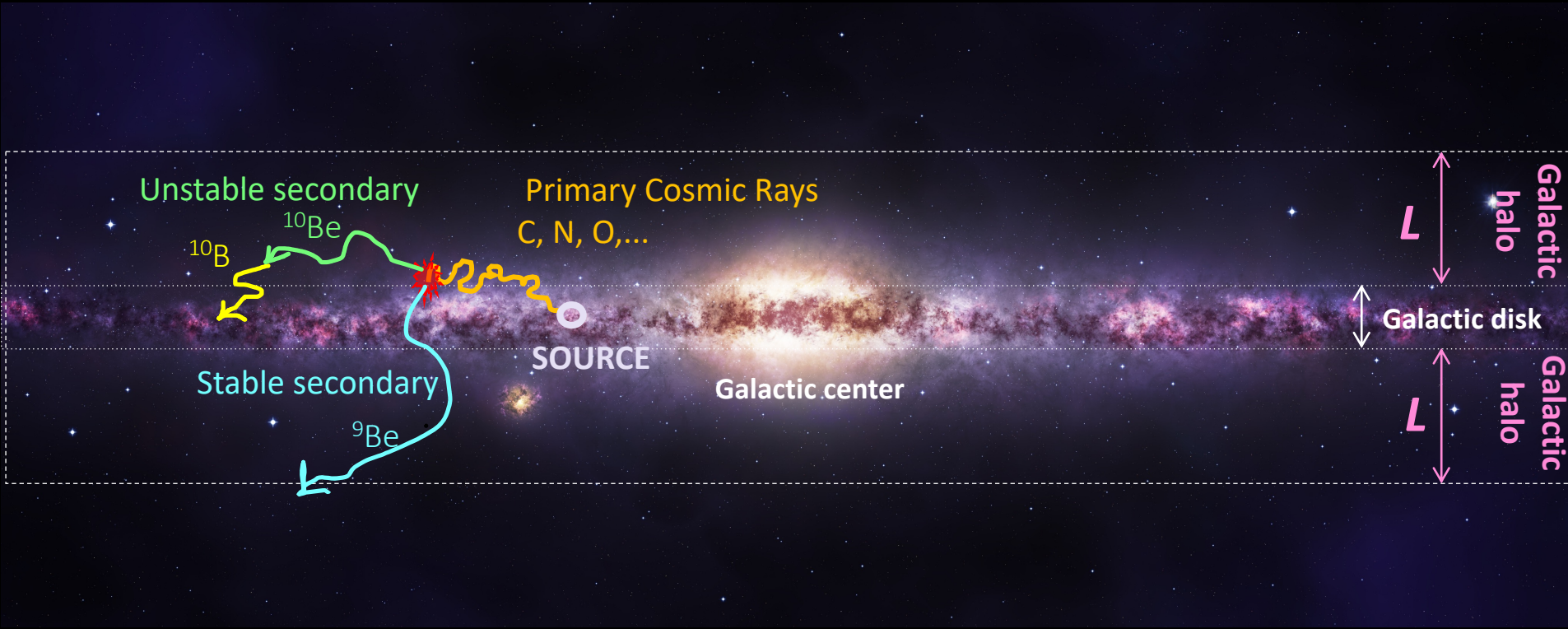
Boschini et al. (ApJ **889**, 167, 2020) predicted primary  $^7\text{Li}$  from a new mechanism so that, from 5 – 20 GeV/n,  $^6\text{Li}/^7\text{Li} = 0.6 \pm 0.1$ .

No detectable primary  $^7\text{Li}$  component.

# Be Isotopes

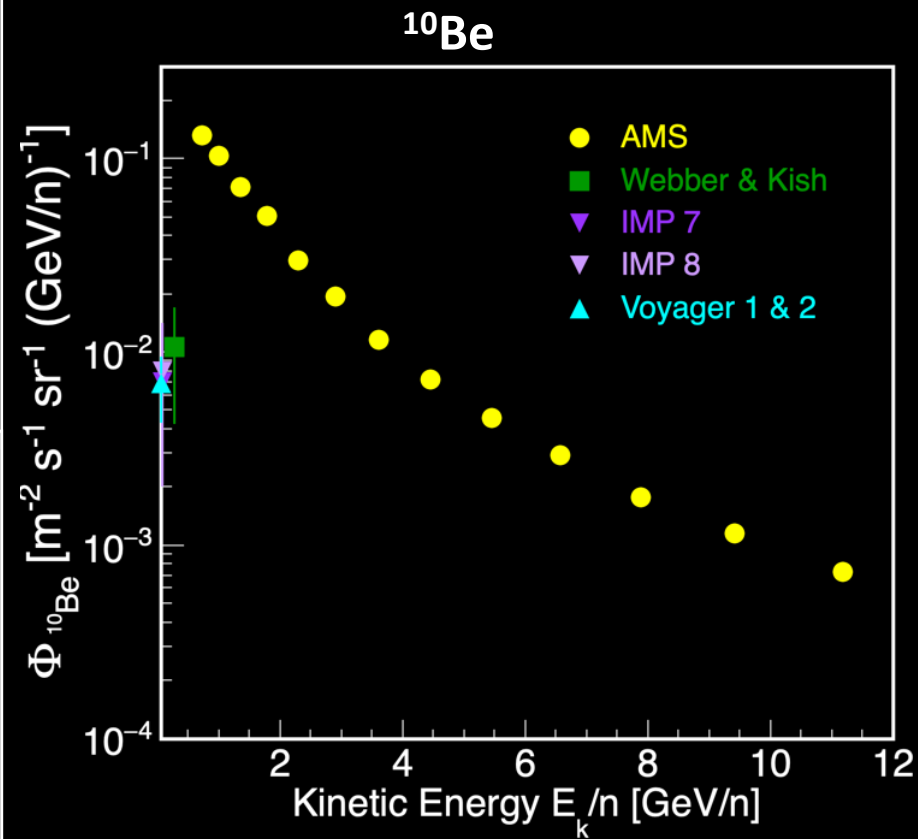
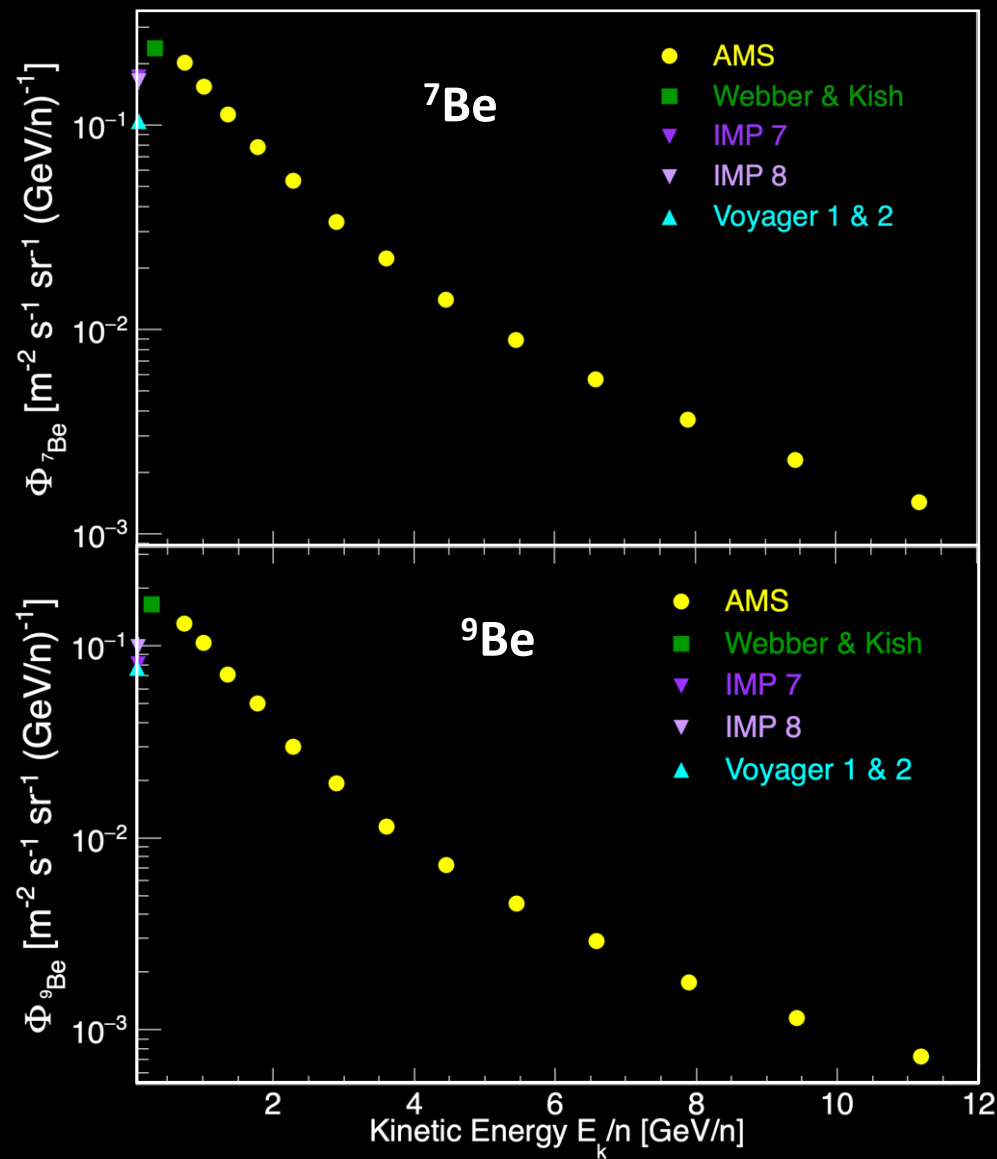
Beryllium nuclei have three isotopes,  ${}^7\text{Be}$ ,  ${}^9\text{Be}$ , and  ${}^{10}\text{Be}$ .

Stable  ${}^9\text{Be}$  propagate in the entire galactic halo while  ${}^{10}\text{Be}$  decay to  ${}^{10}\text{B}$  before reaching the boundary of the Galaxy.

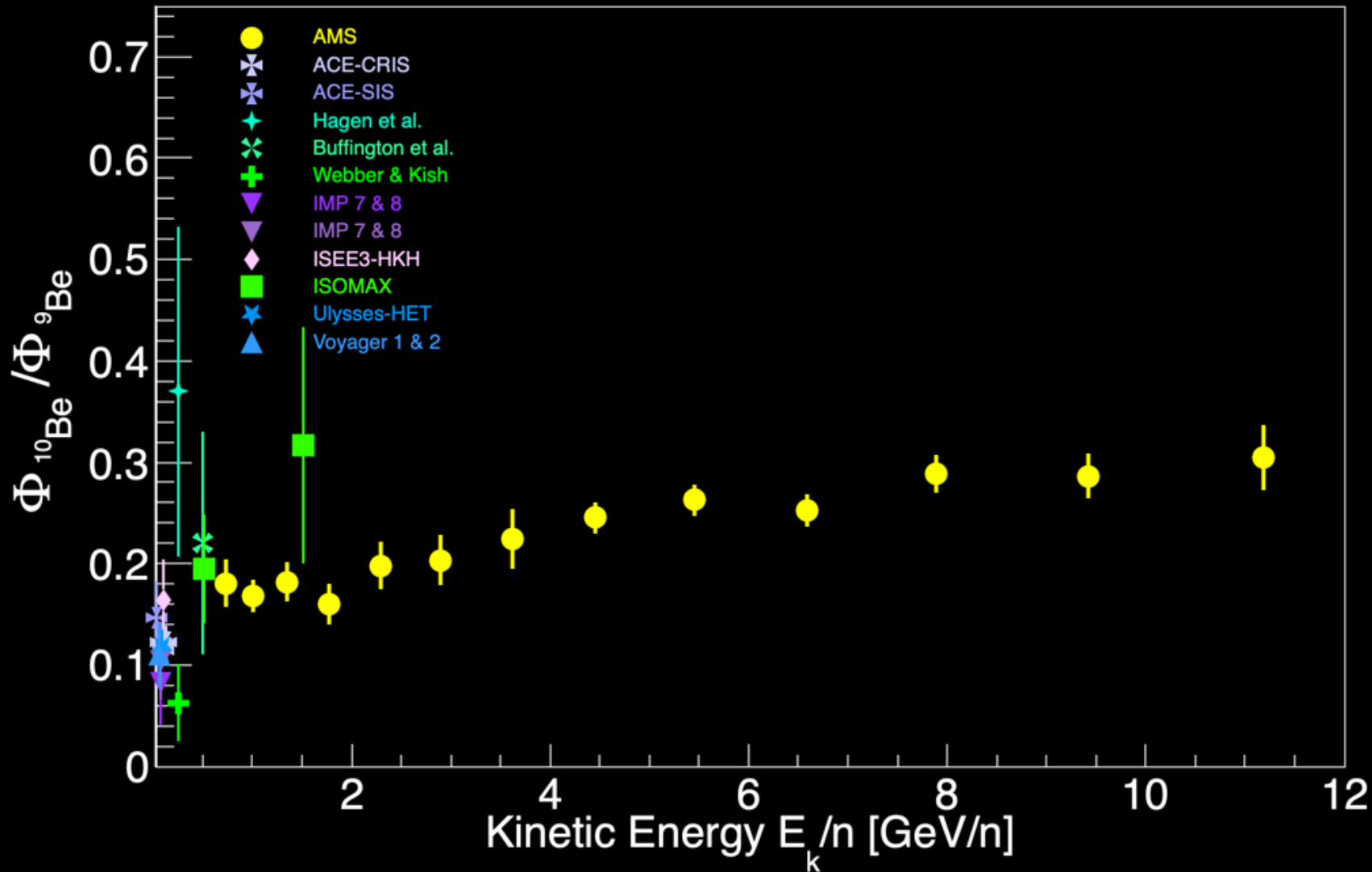


The ratio of unstable-to-stable,  ${}^{10}\text{Be}/{}^9\text{Be}$ , measures the Galactic halo size  $L$ .  
 $L$  determines the galactic cosmic ray propagation (or diffusion) volume.

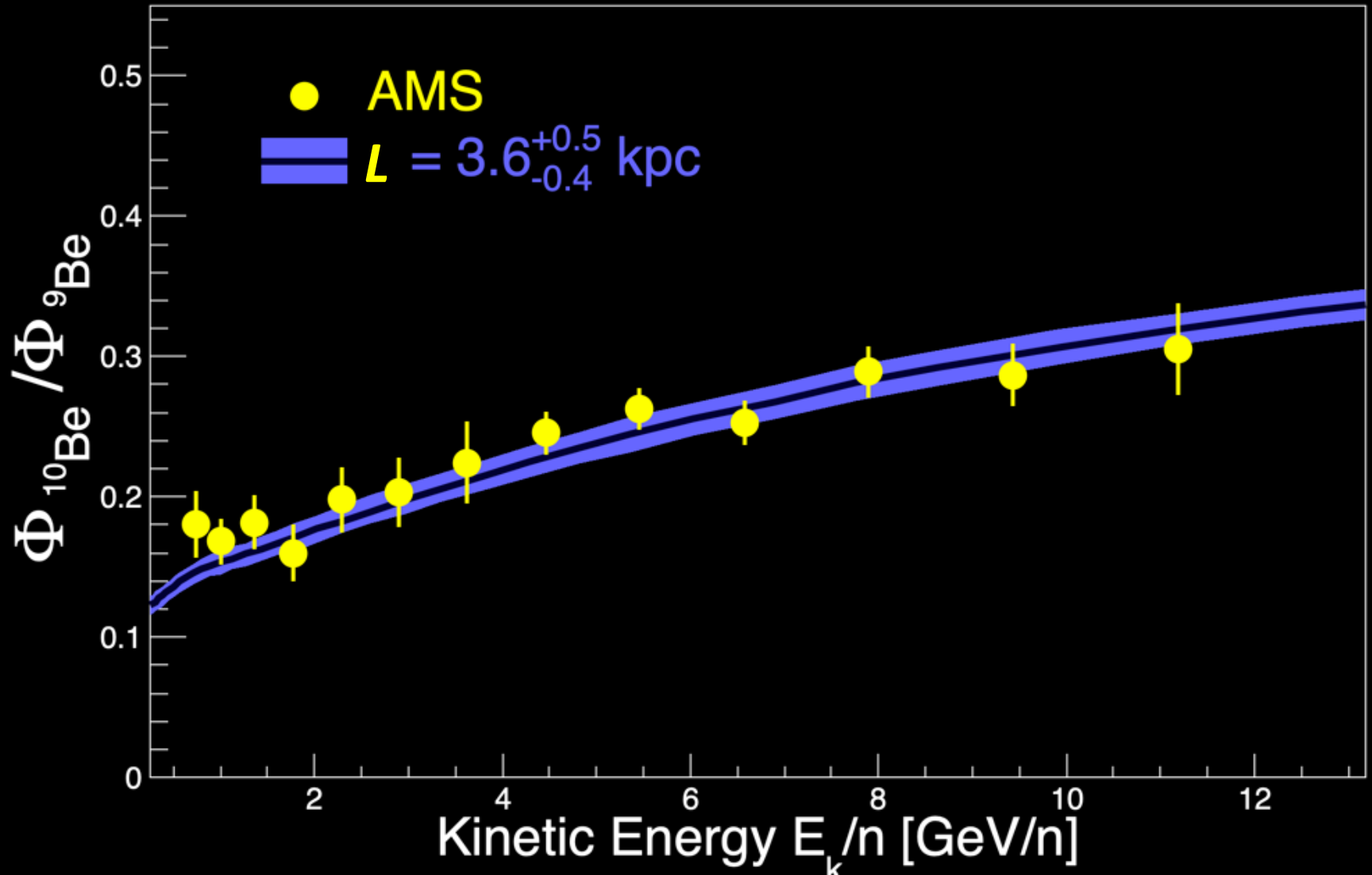
# AMS Results on Beryllium isotopes



# Latest AMS Results on the $^{10}\text{Be}/^9\text{Be}$ ratio



# Measurement of the galactic halo size $L$



The model that fits the AMS measurement has a halo size  $L$  of  $3.6^{+0.5}_{-0.4}$  kiloparsecs

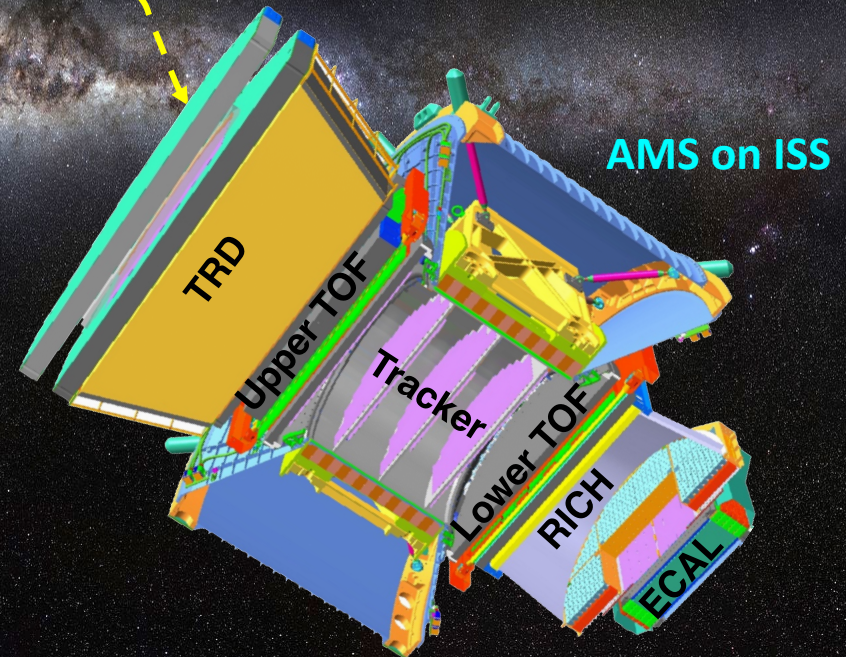
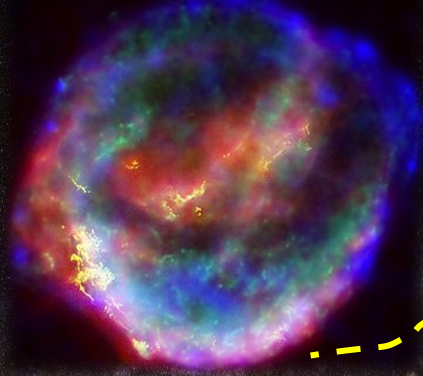
# Latest AMS Results on Heavy Antimatter

Matter is defined by its mass  $M$  and charge  $Z$ .

Antimatter has the same mass  $M$  but opposite charge  $-Z$ .

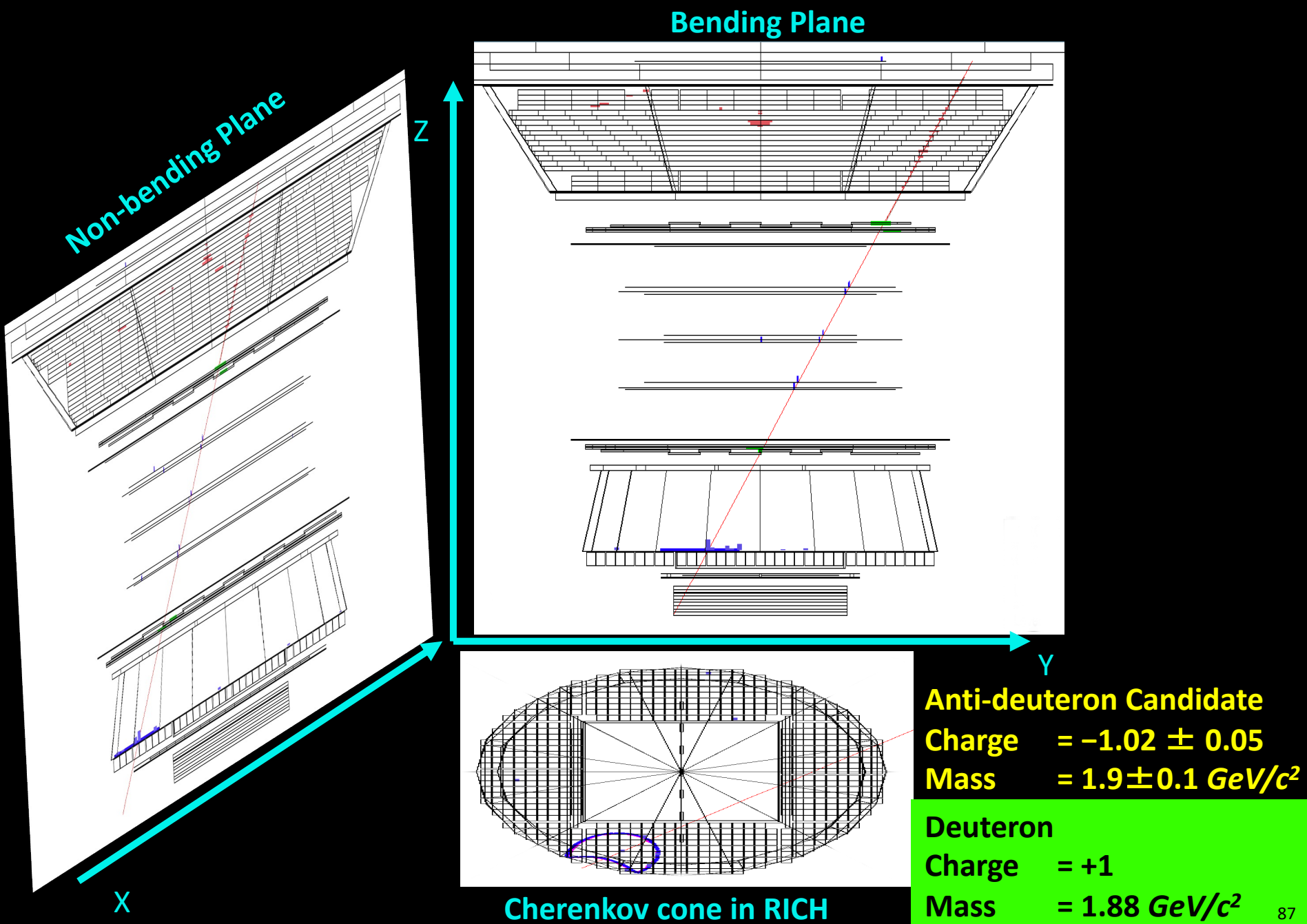
$\bar{D}$ ,  $\bar{He}$ ,  $\bar{C}$ ,  $\bar{O}$  ...

Antimatter Star

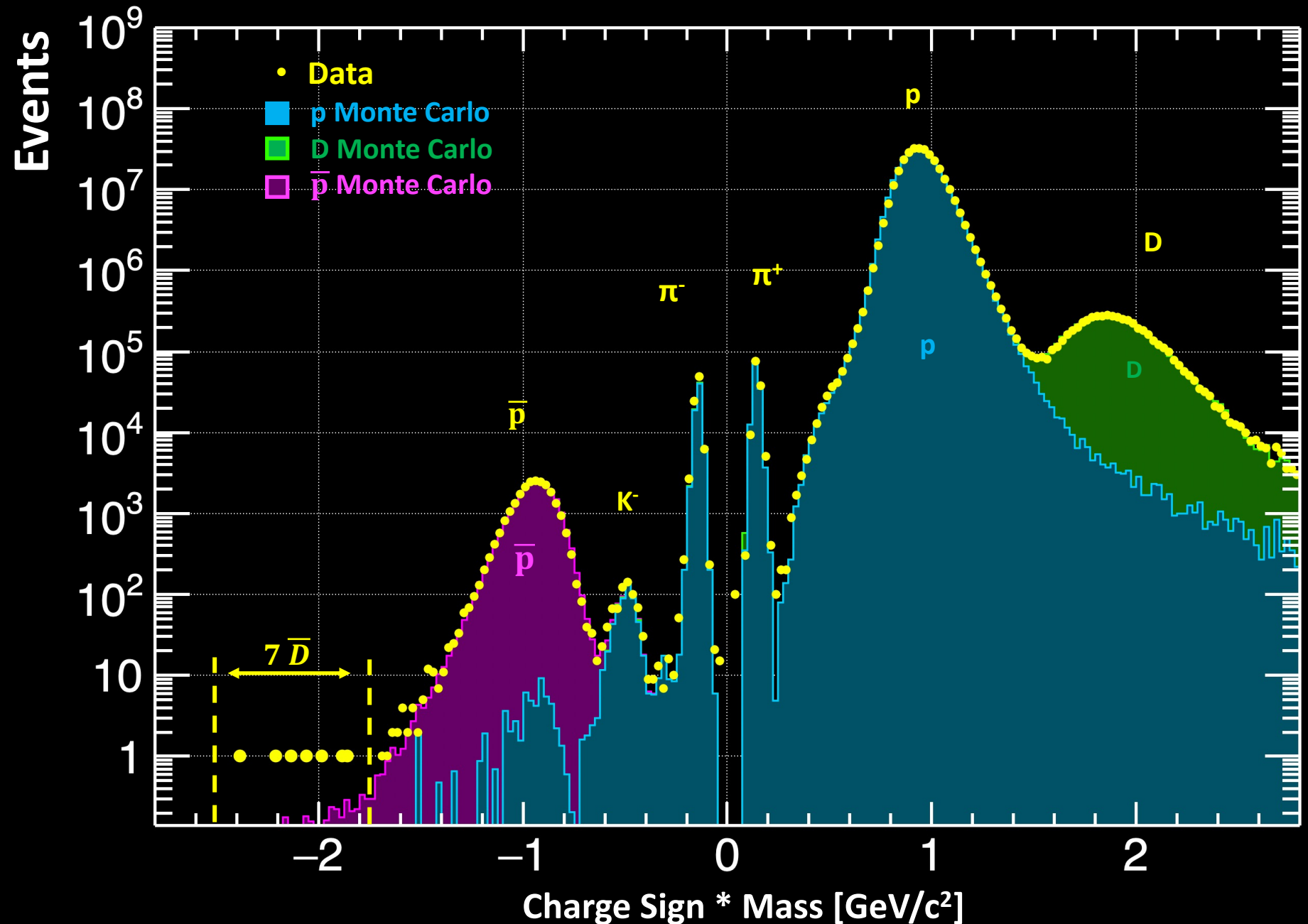


AMS is a unique antimatter spectrometer in space

# An Anti-Deuteron Candidate from ~100 million deuterons and ~10 billion protons



# Current AMS Anti-Deuteron Results





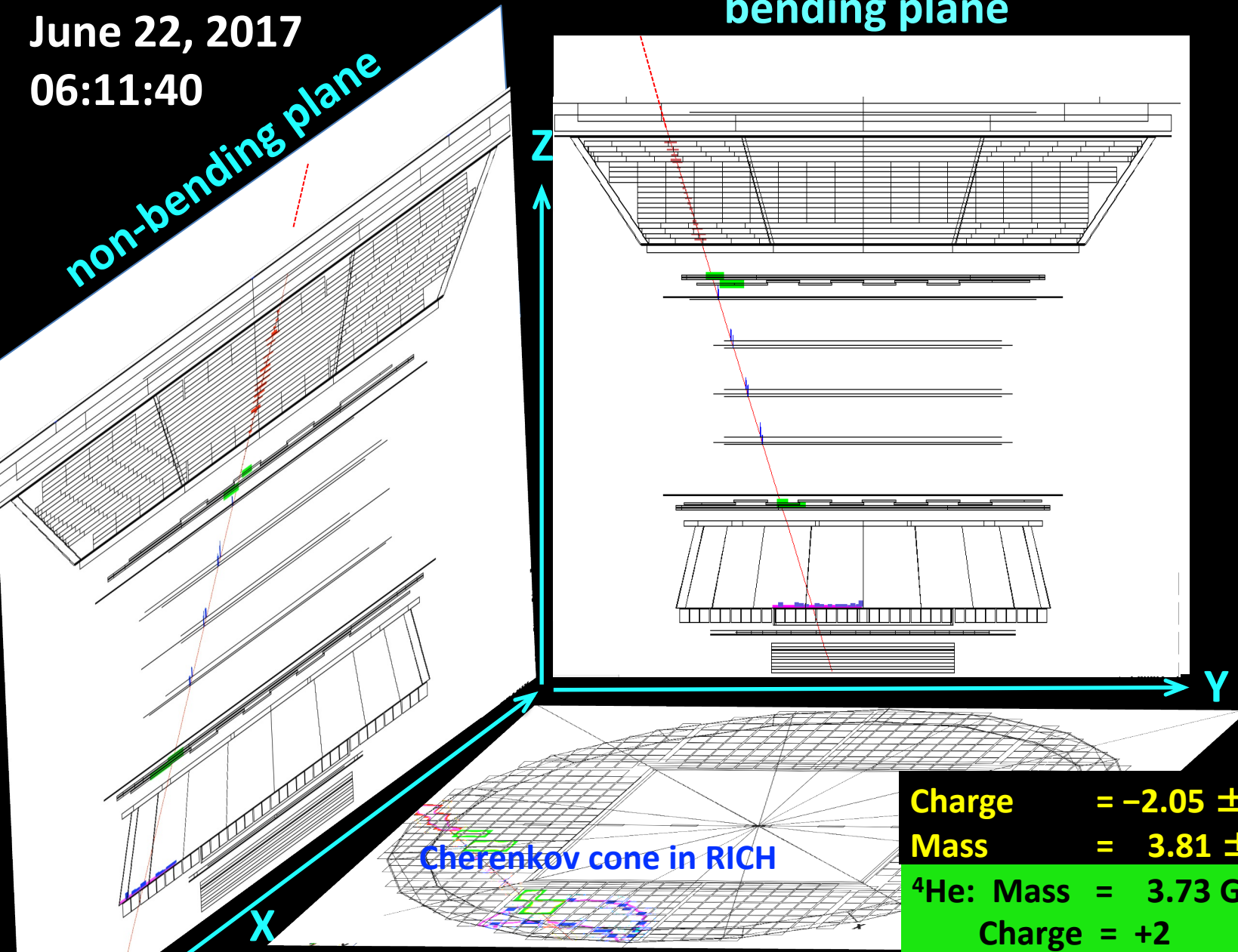
# Anti-<sup>4</sup>Helium Event

June 22, 2017

06:11:40

non-bending plane

bending plane



Charge =  $-2.05 \pm 0.05$

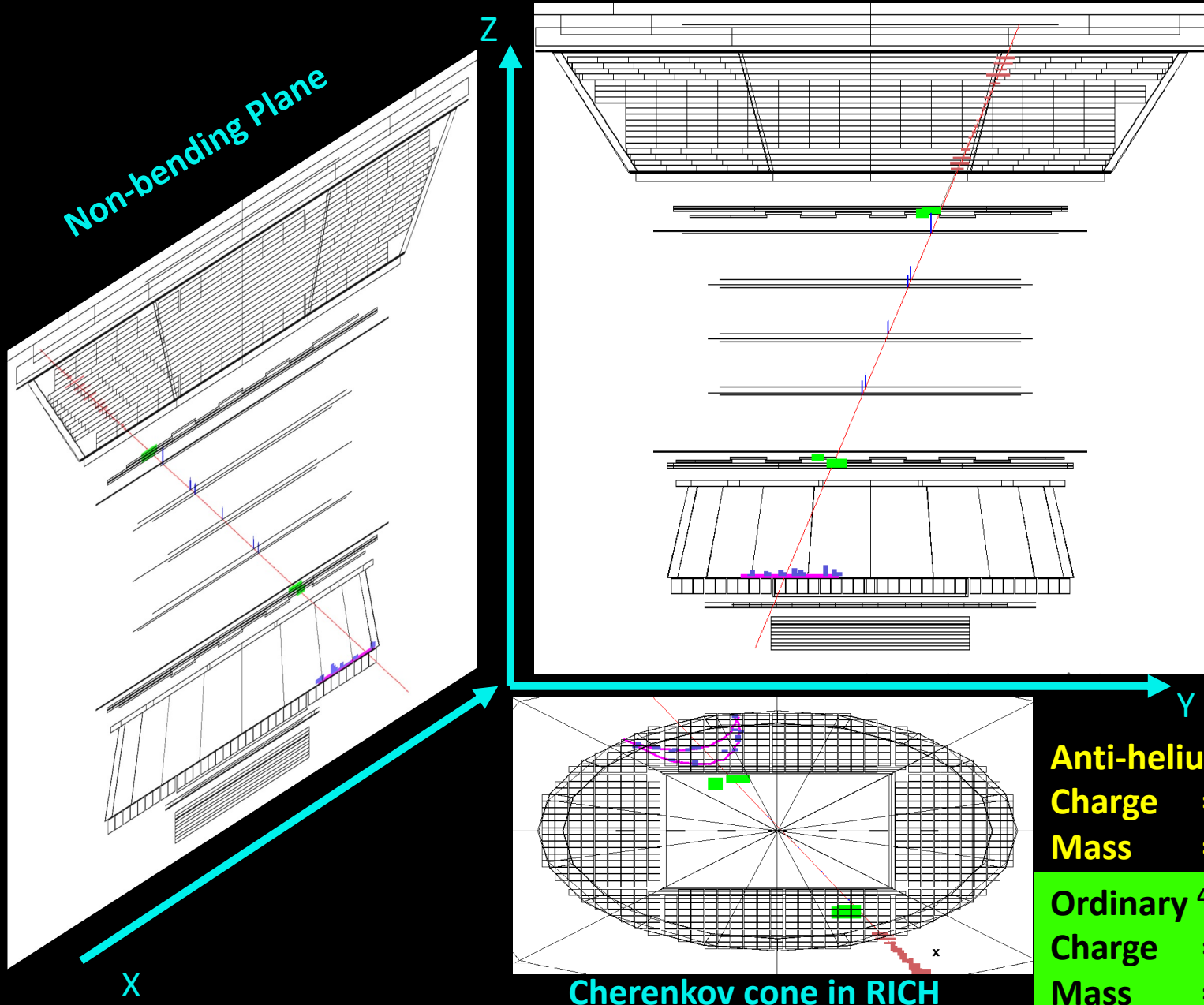
Mass =  $3.81 \pm 0.29 \text{ GeV}/c^2$

<sup>4</sup>He: Mass =  $3.73 \text{ GeV}/c^2$

Charge = +2

# Another Anti-Helium Event (Sept. 20, 2022)

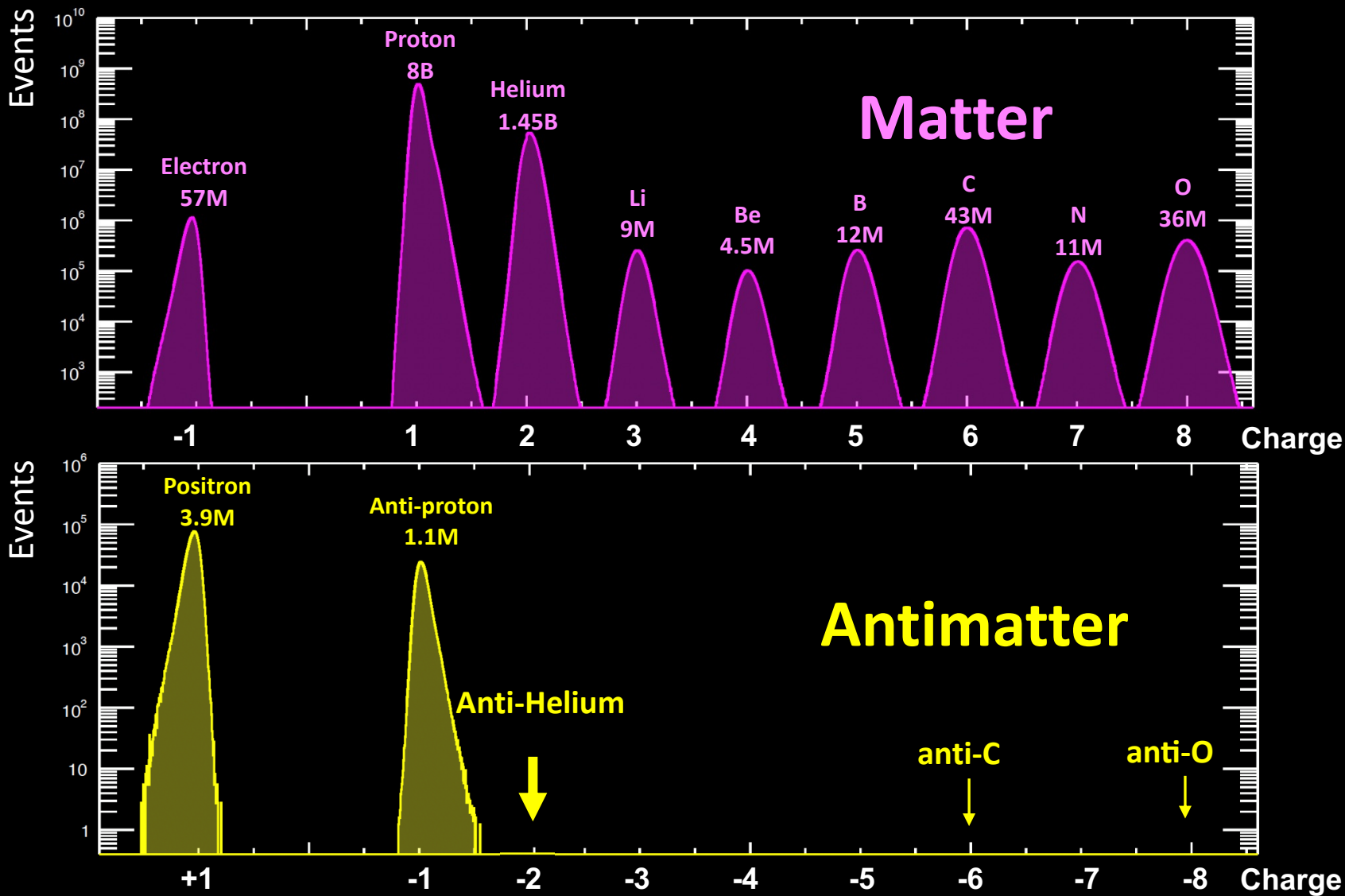
Bending Plane



**Anti-helium Candidate**  
Charge = -2  
Mass =  $3.15 \pm 0.53$

**Ordinary  $^4\text{He}$**   
Charge = +2  
Mass =  $3.73 \text{ GeV}/c^2$

# Current Matter and Antimatter Statistics



By 2030, AMS will have additional measurement points in the study of antimatter: anti-deuterons, anti-helium, anti-carbon and anti-oxygen.

PHYSICS

## The Space Station's Crown Jewel

A fancy cosmic-ray detector, the Alpha Magnetic Spectrometer, is about to scan the cosmos for dark matter, antimatter and more

By George Musser, staff editor

**T**HE WORLD'S MOST ADVANCED COSMIC-RAY DETECTOR TOOK 16 YEARS AND \$2 billion to build, and not long ago it looked as though it would wind up mothballed in some warehouse. NASA, directed to finish building the space station and retire the shuttle by the end of 2010, said it simply did not have room in its schedule to launch the instrument anymore. Saving it took a lobbying campaign by physicists and intervention by Congress to extend the shuttle program. And so the shuttle *Endeavour* is scheduled to take off on April 19 for the express purpose of delivering the Alpha Magnetic Spectrometer (AMS) to the International Space Station.

Cosmic rays are subatomic particles and atomic nuclei that zip and zap through space, coming from ordinary stars, supernovae explosions, neutron stars, black holes and who knows what—the last category naturally being of greatest interest and the main impetus for a brand-new instrument. Dark matter is one of those possible mystery sources. Clumps of the stuff out in space might occasionally release blazes of particles that would set the detectors alight. Some physicists also speculate that our planet might be peppered with the odd antistar coming from distant galaxies made not of matter but of its evil antitwin.

The spectrometer's claim to fame is that it can tell the ordinary from the extraordinary, which otherwise are easily conflated. No other instrument has the combination of detectors that can tease out all the properties of a particle: mass, velocity, type, electric charge. Its closest predecessor is the PAMELA instrument, launched by a European consortium in 2006. PAMELA has seen hints of dark matter and other exotica, but its findings remain ambiguous because it lacks the ability to distinguish a low-mass antiparticle, such as a positron, from a high-mass ordinary particle with the same electric charge, such as a proton.

The AMS instrument is a monster by the standards of the space program, with a mass of seven metric tons (more than 14 times heavier than PAMELA) and a power consumption of 2,400 watts. In a strange symbiotic way, it and the space station have come to justify each other's existence. The station satisfies the instrument's thirst for power and orbital boosts; the spectrometer, although it could never fully placate the station's many skeptics, at least means the outpost will do world-class research. As CERN's Large Hadron Collider plumbs the depths of nature on the ground, the Alpha Magnetic Spectrometer will do the same from orbit.

SCIENTIFIC AMERICAN ONLINE

For more information on how the Alpha Magnetic Spectrometer works, visit [ScientificAmerican.com/may2011/ams](http://ScientificAmerican.com/may2011/ams)

### Time of Flight System 1

**PURPOSE:** Measure particle velocity and charge.  
**DESIGN:** Sheets of transparent polymer that glow when a charged particle passes through.  
**OPERATION:** A pair of these detectors times how fast the particle takes to cover the length of the instrument.

### Magnet

**PURPOSE:** Bend paths of charged particles.  
**DESIGN:** Permanent magnet with a field strength of 0.45 tesla. This magnet replaces the cryogenic superconducting magnet used in the original design, giving the instrument a longer lifetime.  
**OPERATION:** When passing through, a positively charged particle is deflected to the left, a negatively charged one to the right.

### Silicon Tracker

**PURPOSE:** Measure particle charge and momentum.  
**DESIGN:** Nine planes of particle detectors.  
**OPERATION:** The detectors trace out the path of each particle through the magnetic field.

### Transition Radiation Detector

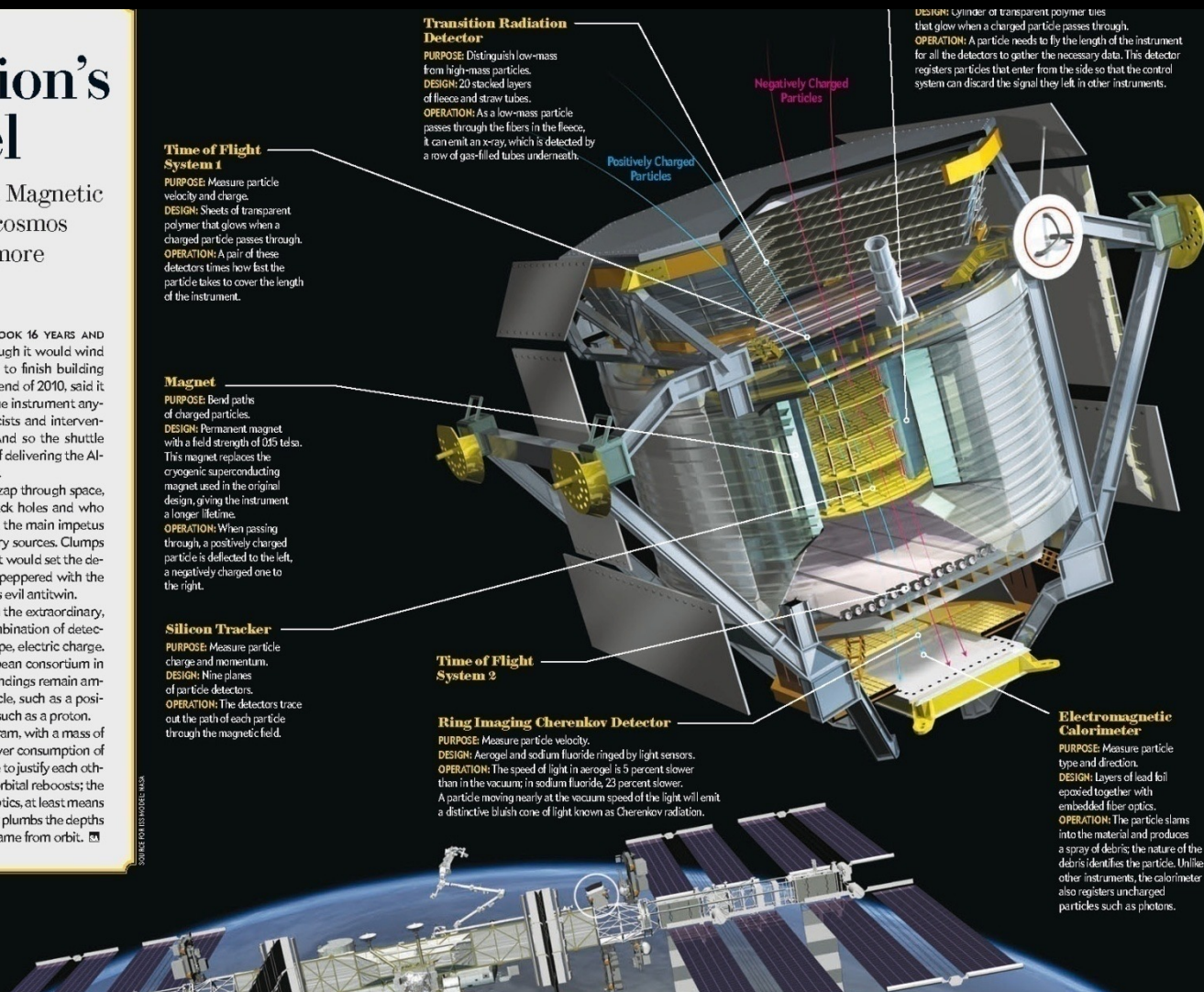
**PURPOSE:** Distinguish low-mass from high-mass particles.  
**DESIGN:** 20 stacked layers of fleece and straw tubes.  
**OPERATION:** As a low-mass particle passes through the fibers in the fleece, it can emit an x-ray, which is detected by a row of gas-filled tubes underneath.

### Time of Flight System 2

### Ring Imaging Cherenkov Detector

**PURPOSE:** Measure particle velocity.  
**DESIGN:** Aerogel and sodium fluoride ringed by light sensors.  
**OPERATION:** The speed of light in aerogel is 9 percent slower than in the vacuum; in sodium fluoride, 23 percent slower. A particle moving nearly at the vacuum speed of the light will emit a distinctive bluish cone of light known as Cherenkov radiation.

**DESIGN:** Cylinder of transparent polymer tubes that glow when a charged particle passes through.  
**OPERATION:** A particle needs to fly the length of the instrument for all the detectors to gather the necessary data. This detector registers particles that enter from the side so that the control system can discard the signal they left in other instruments.



**In twelve years on the ISS, AMS has recorded more than 220 billion cosmic rays. The accuracy and characteristics of the data simultaneously from many different types of cosmic rays require the development of a new comprehensive model of the universe. AMS will continue to collect data to 2030 with an upgraded detector.**