

QUANTUM INFORMATION TOOLS AT THE INTERFACE BETWEEN QUANTUM THEORY AND GRAVITY

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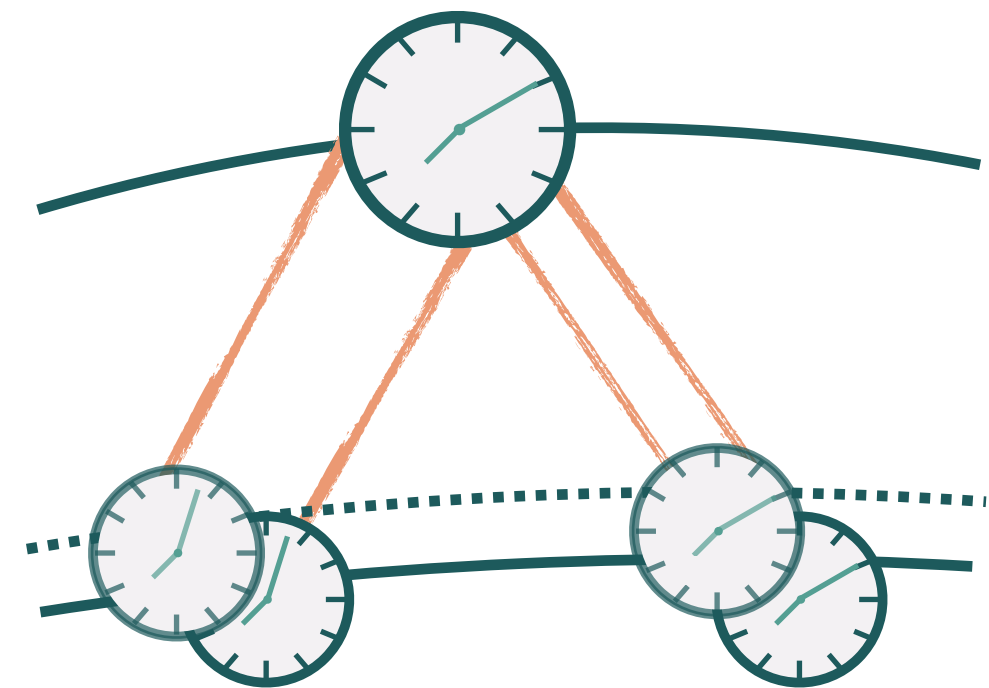
Image credits: J. Palomino

Basics of Quantum Gravity
22-25 May 2023

LECTURE 4: QUANTUM CLOCKS AS PROBES OF NONCLASSICAL SPACETIME

- Quantum clock: the basics
- Quantum clocks in classical spacetime
- Quantum clocks in nonclassical spacetime: limits to measurability and relative localization of events
- The gravitational quantum switch: quantum clocks and superposition of temporal order

QUANTUM CLOCKS IN CLASSICAL SPACETIME

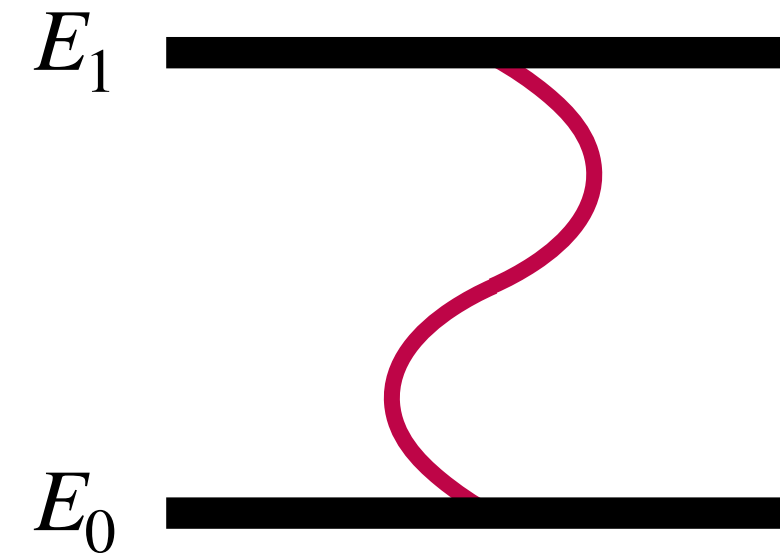


How do quantum clocks encode information about spacetime?

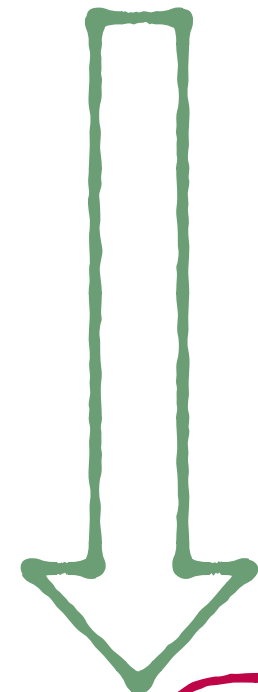
WHAT IS A QUANTUM CLOCK?



=



$$|\psi_0\rangle = \frac{1}{\sqrt{2}}(|E_0\rangle + |E_1\rangle)$$



$$|\psi_t\rangle = \frac{1}{\sqrt{2}}(|E_0\rangle + e^{\frac{i}{\hbar}(E_1-E_0)t} |E_1\rangle)$$

-1

$$H_C = E_0 |E_0\rangle\langle E_0| + E_1 |E_1\rangle\langle E_1|$$

$$\langle \psi_{t_\perp} | \psi_0 \rangle = 0$$

ORTHOGONALISATION TIME

$$t_\perp = \frac{\pi\hbar}{(E_1 - E_0)}$$

$$E = mc^2$$

$$m \rightarrow m + \frac{\hat{H}_I}{c^2}$$

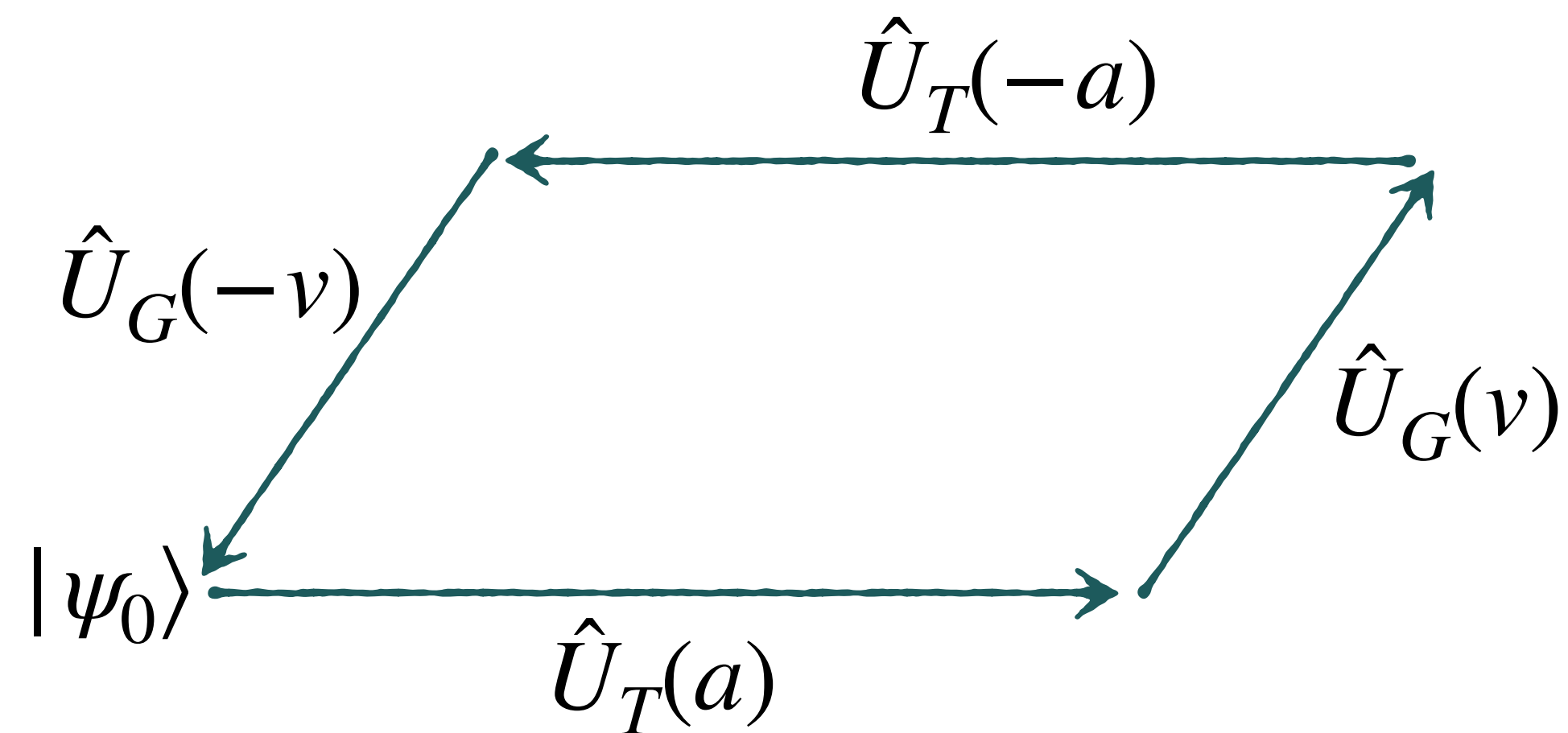
**THE MASS BECOMES
AN OPERATOR!**

MASS AND PROPER TIME

$$g_{\mu\nu} p^\mu p^\nu = c^{-2} E_{rest}$$

$$E_{rest} = mc^2 + H_C$$

Total rest energy
(Mass + internal degrees of freedom)



Translation operator

$$\hat{U}_T(a) = e^{\frac{i}{\hbar} a \hat{p}}$$

Galilean boost operator

$$\hat{U}_G(v) = e^{\frac{i}{\hbar} v \hat{G}}$$

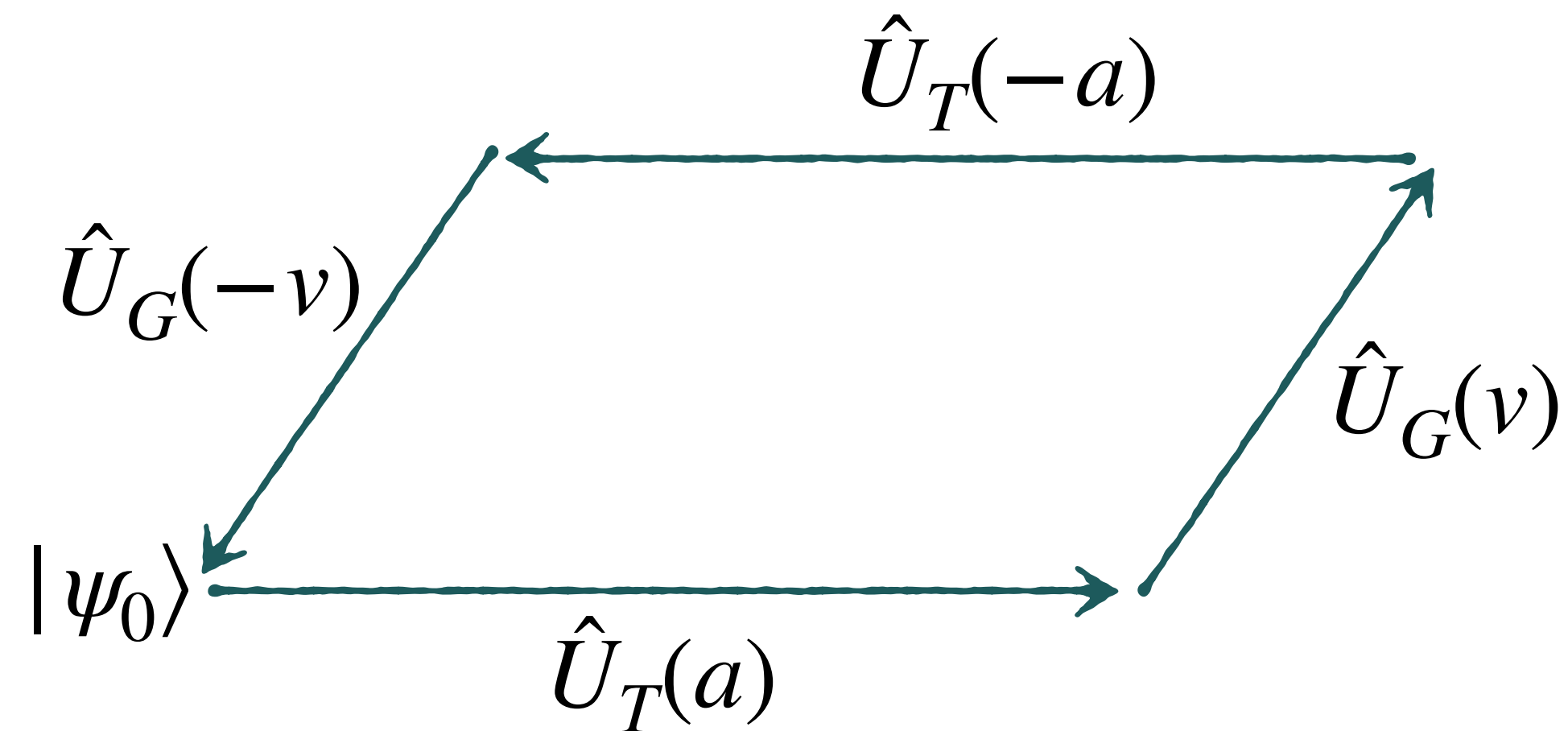
$$\hat{G} = \hat{p}t - m\hat{x}$$

MASS AND PROPER TIME

$$g_{\mu\nu} p^\mu p^\nu = c^{-2} E_{rest}$$

$$E_{rest} = mc^2 + H_C$$

Total rest energy
(Mass + internal degrees of freedom)



$$|\psi_0\rangle \rightarrow e^{\frac{i}{\hbar} m a v} |\psi_0\rangle$$

SUPERPOSITION OF MASSES?

$$|\psi_0^{(m_1)}\rangle + |\psi_0^{(m_2)}\rangle \rightarrow |\psi_0^{(m_1)}\rangle + e^{\frac{i}{\hbar} \Delta m \cdot a v} |\psi_0^{(m_2)}\rangle$$

Bargmann: mass superselection rule

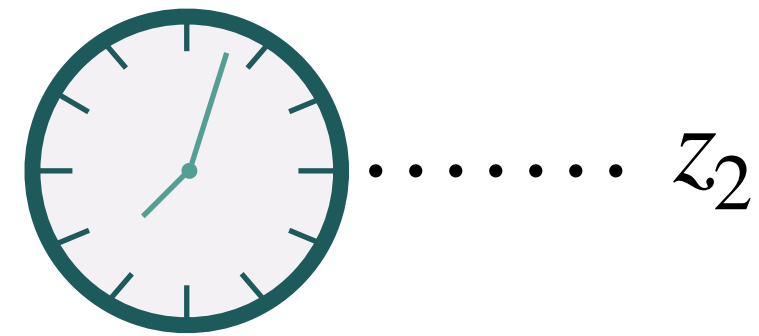
$$e^{\frac{i}{\hbar} m a v} = e^{\frac{i}{\hbar} \Delta \tau c^2 \Delta m}$$

Greenberger, PRL 87 (2001)

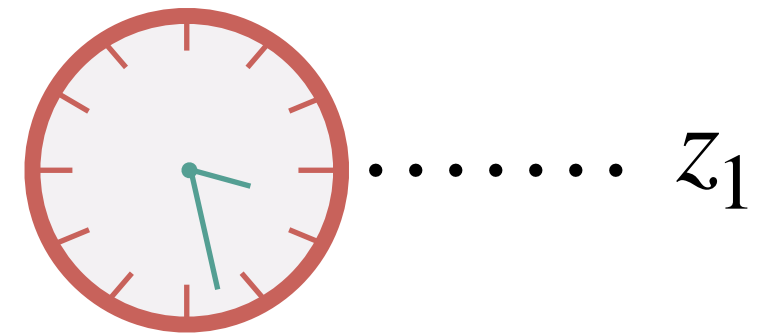
GRAVITY INFLUENCES HOW A CLOCK TELLS THE TIME

“Lower is slower”

$$\tau_2 = \Delta t(1 + V_N(z_2))$$





$$\tau_1 = \Delta t(1 + V_N(z_1))$$



Article | [Published: 16 February 2022](#)

Resolving the gravitational redshift across a millimetre-scale atomic sample

[Tobias Bothwell](#) , [Colin J. Kennedy](#), [Alexander Aepli](#), [Dhruv Kedar](#), [John M. Robinson](#), [Eric Oelker](#), [Alexander Staron](#) & [Jun Ye](#) 

[Nature](#) **602**, 420–424 (2022) | [Cite this article](#)

QUANTUM CLOCKS AS PROBES OF SPACETIME

REST FRAME $\frac{d}{d\tau} |\psi\rangle_C = \hat{H}_C |\psi\rangle_C$

LAB FRAME $\frac{d}{dt} |\psi\rangle_C = \dot{\tau} \hat{H}_C |\psi\rangle_C$

$$\dot{\tau} = \sqrt{g_{\mu\nu} \frac{dx^\mu}{dt} \frac{dx^\nu}{dt}}$$

$$c^2 d\tau^2 = g_{\mu\nu} dx^\mu dx^\nu$$

Line element in
classical spacetime

Internal and external degrees of freedom

$$|\Psi_t\rangle = e^{-\frac{i}{\hbar}(\hat{H}_{ext} + \dot{\tau} \hat{H}_C)t} |\phi_0\rangle |\psi_0\rangle_C = \int D[x(t)] \phi_0(x_0) e^{\frac{i}{\hbar} \int ds (S_{ext} + \dot{\tau} \hat{H}_C)} |x(t)\rangle |\psi_0\rangle_C$$

semiclassical evolution
for external d.o.f.

QUANTUM CLOCKS AS PROBES OF SPACETIME

THE NEWTONIAN LIMIT

$$g_{00}(x) = 1 + \frac{2\Phi(x)}{c^2}, \quad g_{0i}(x) = 0, \quad g_{ij}(x) = -1.$$

$$S = mc^2 \int d\tau = mc^2 \int dt \sqrt{g_{00} - \frac{|\dot{x}|^2}{c^2}}$$
$$\approx mc^2 \int dt \left\{ 1 - \frac{|\dot{x}|^2}{2c^2} + \frac{\Phi(x)}{c^2} + O(c^{-4}) \right\}$$

$$c^2 d\tau^2 = g_{\mu\nu} dx^\mu dx^\nu$$

Line element in classical spacetime

$$H = \frac{|p|^2}{2m} + m\Phi(x)$$

Internal energy contributes to the total mass

$$H = \frac{\hat{p}^2}{2} \left(m + \frac{\hat{H}_I}{c} \right)^{-1} + \left(m + \frac{\hat{H}_C}{c^2} \right) \Phi(\hat{x})$$

Coupling between internal and external degrees of freedom

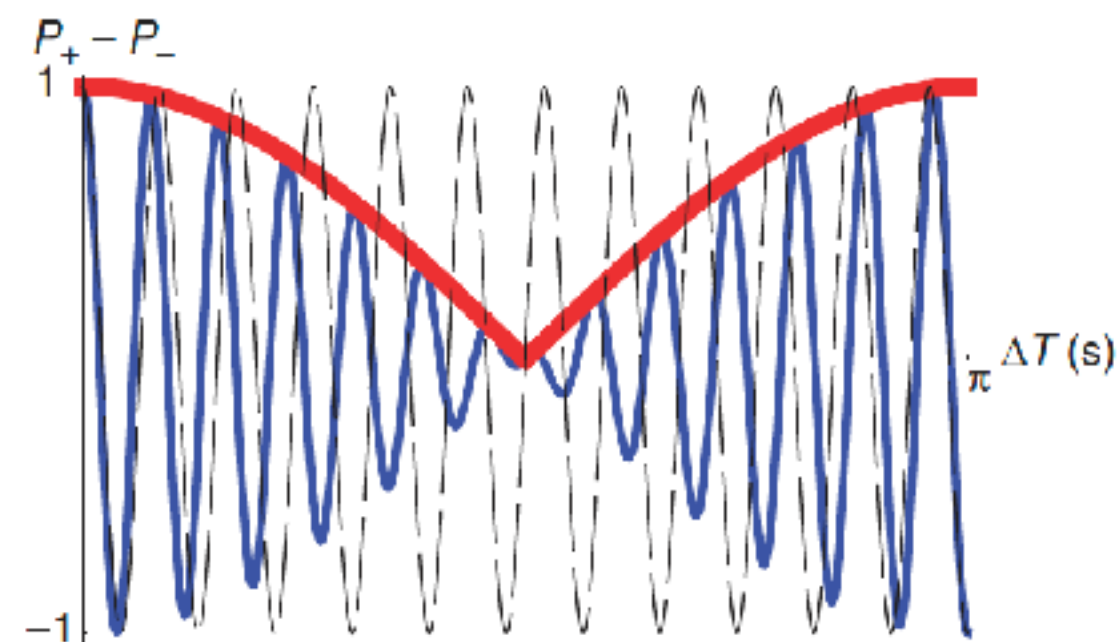
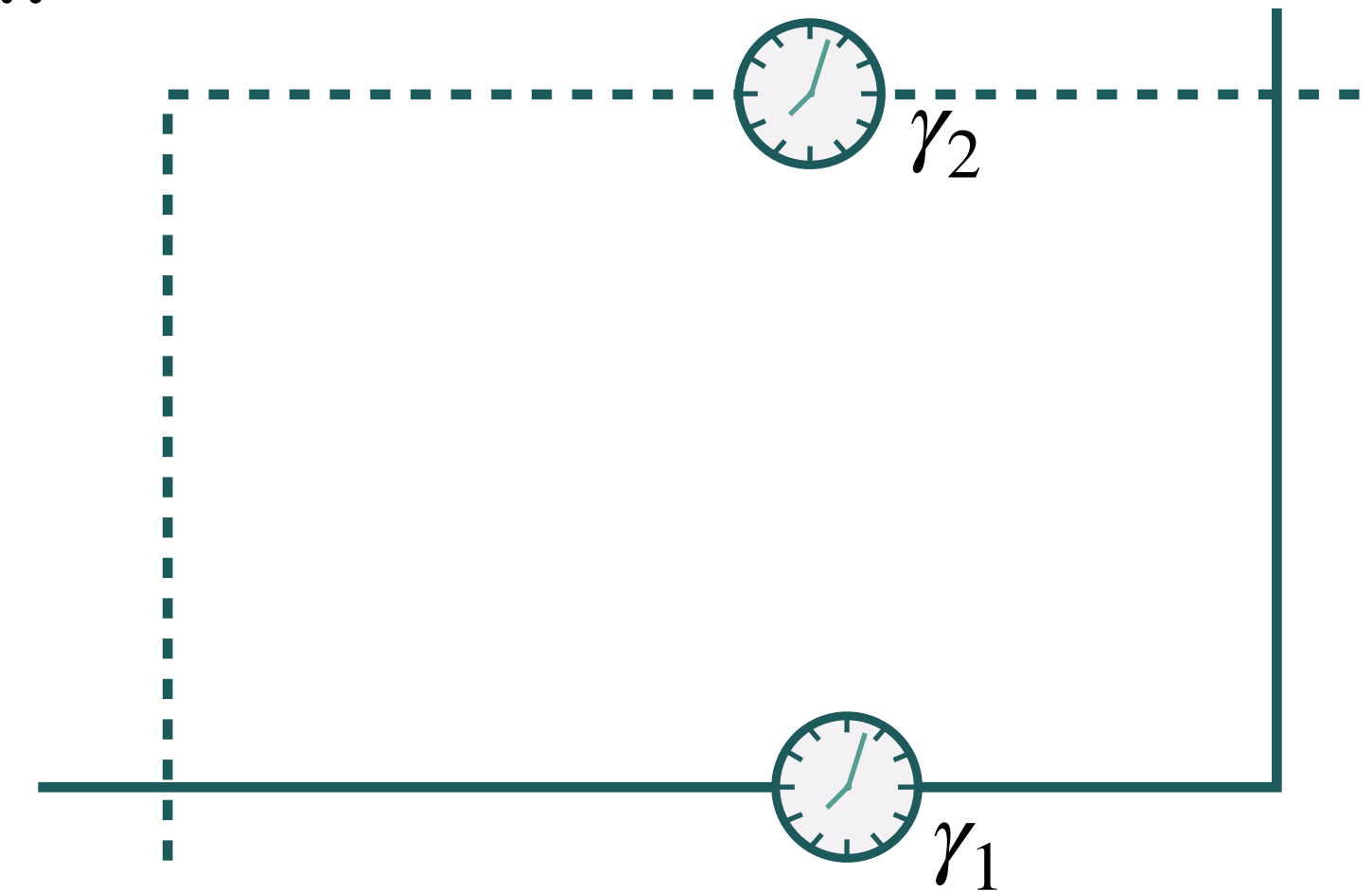
QUANTUM CLOCKS AS PROBES OF SPACETIME

Quantum clocks test the principle of linear superposition and gravitational time dilation together.

“LOWER IS SLOWER”

$$|\tau_0\rangle = \frac{1}{\sqrt{2}} (|E_0\rangle + |E_1\rangle)$$

$$\frac{1}{\sqrt{2}} (|\gamma_1\rangle |\tau_1\rangle + e^{i\Delta\phi} |\gamma_2\rangle |\tau_2\rangle)$$

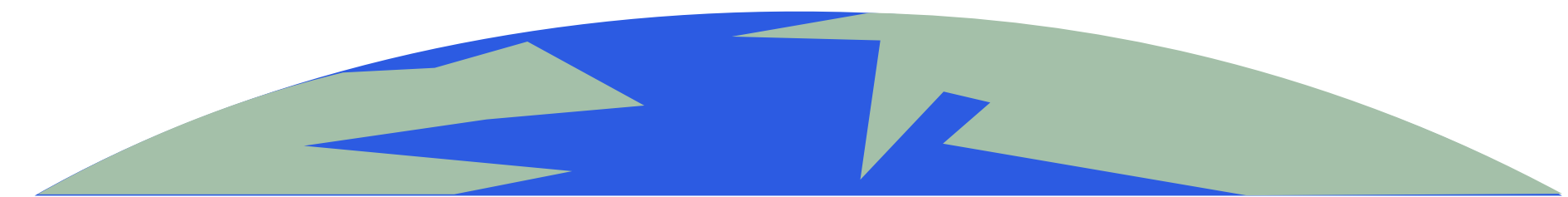


$$\mathcal{V} = |\langle \tau_1 | \tau_2 \rangle|^2$$

Visibility
(interference)

$$\mathcal{D} = 1 - |\langle \tau_1 | \tau_2 \rangle|^2$$

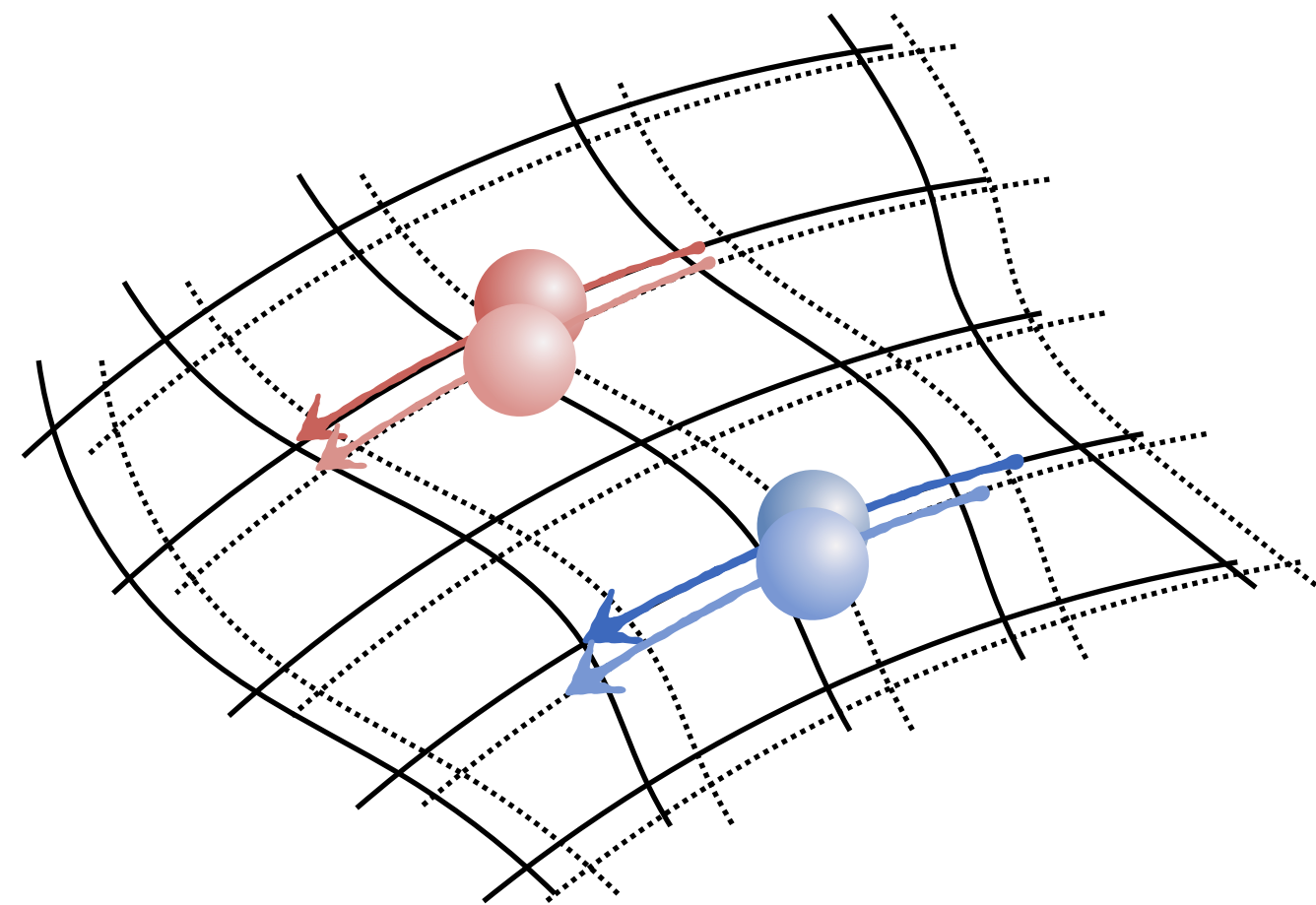
Distinguishability
(which-way information)



The clock carries “which-way” information in its proper time.

Zych, Costa, Pikovski, Brukner, Nat. Commun. (2011)

QUANTUM CLOCKS IN NONCLASSICAL SPACETIME

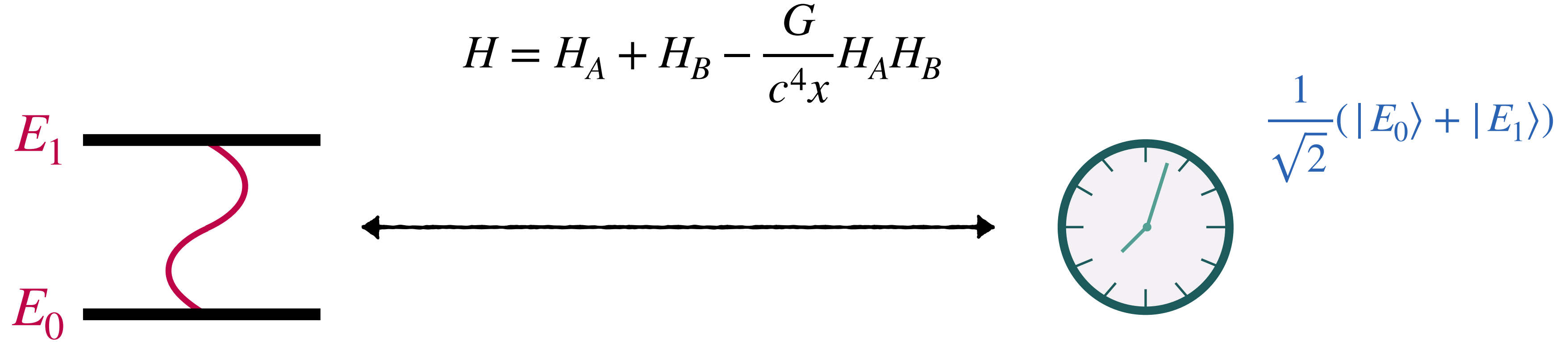


*Can quantum clocks probe
nonclassical spacetime?*

ENTANGLEMENT OF QUANTUM CLOCKS THROUGH GRAVITY

$$E_0 = 0$$

$$E_1 = \Delta E$$



$$|\psi_t\rangle = e^{-\frac{i}{\hbar} \hat{H} t} |\psi_0\rangle$$

$$\propto |E_0\rangle_A (|E_0\rangle_B + e^{-\frac{i}{\hbar} \Delta E t} |E_1\rangle_B) + |E_1\rangle_A (|E_0\rangle_B + e^{-\frac{i}{\hbar} \left(1 - \frac{G\Delta E}{c^4 x}\right) \Delta E t} |E_1\rangle_B)$$

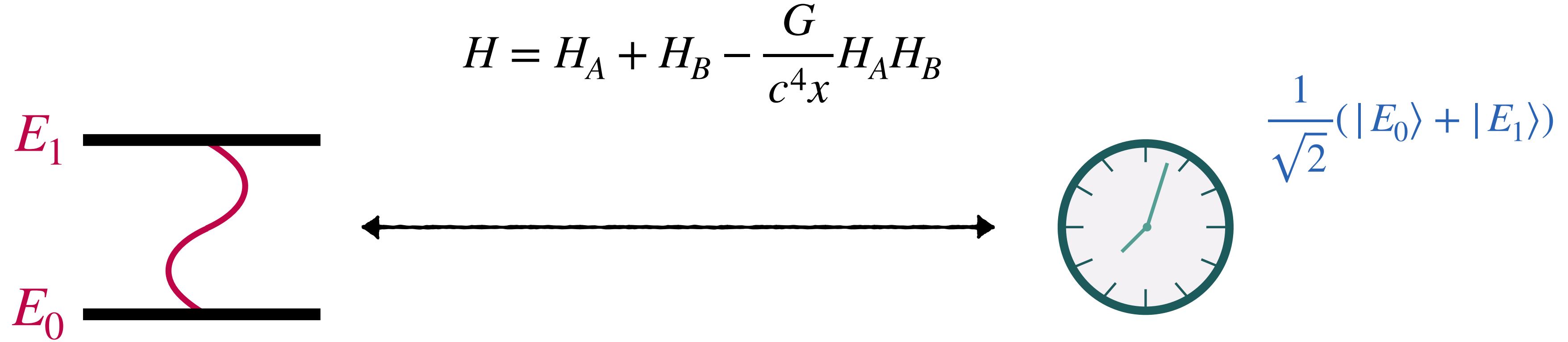
Difference between the ticks due to clock A

Ticking of clock B when A is in E_0

Ticking of clock B when A is in E_1

Castro Ruiz, Giacomini, Brukner, PNAS (2017)

ENTANGLEMENT OF QUANTUM CLOCKS THROUGH GRAVITY

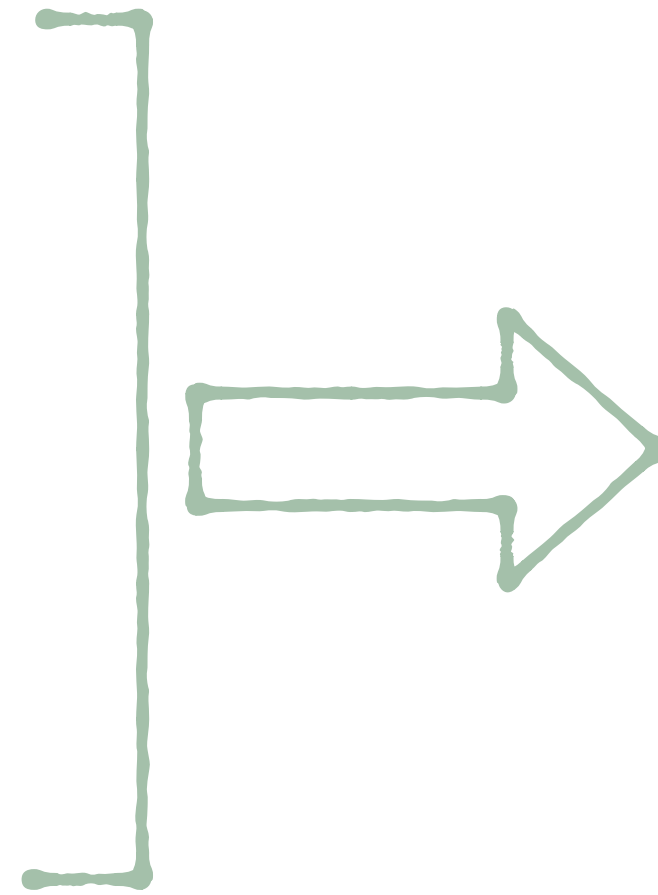


Precision of clock A

$$t_{\perp} = \frac{\pi \hbar}{(E_1 - E_0)}$$

Disturbance to clock B
due to clock A

$$\Delta t = \frac{G(E_1 - E_0)}{c^4 x} t$$



$$t_{\perp} \Delta t = \frac{\pi \hbar G t}{c^4 x}$$

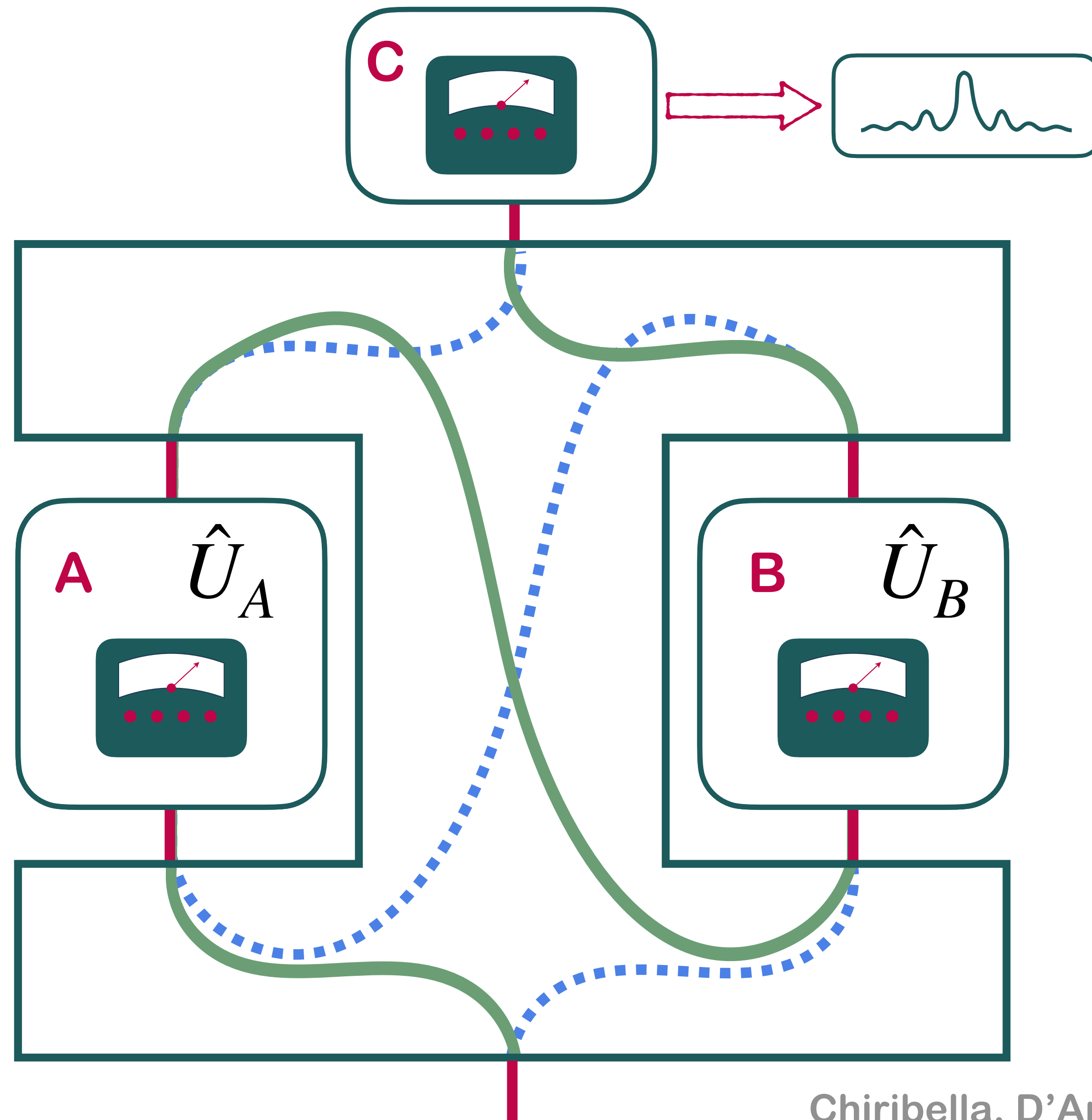
Limitation to the
measurability of time
for nearby clocks

THE QUANTUM SWITCH: SUPERPOSITION OF CAUSAL ORDERS

Process matrix with
3 parties: A, B, C

T: target qubit
C: control qubit

$$\frac{1}{\sqrt{2}} \left(|0\rangle_C \hat{U}_B \hat{U}_A |\psi\rangle + |1\rangle_C \hat{U}_A \hat{U}_B |\psi\rangle \right)$$



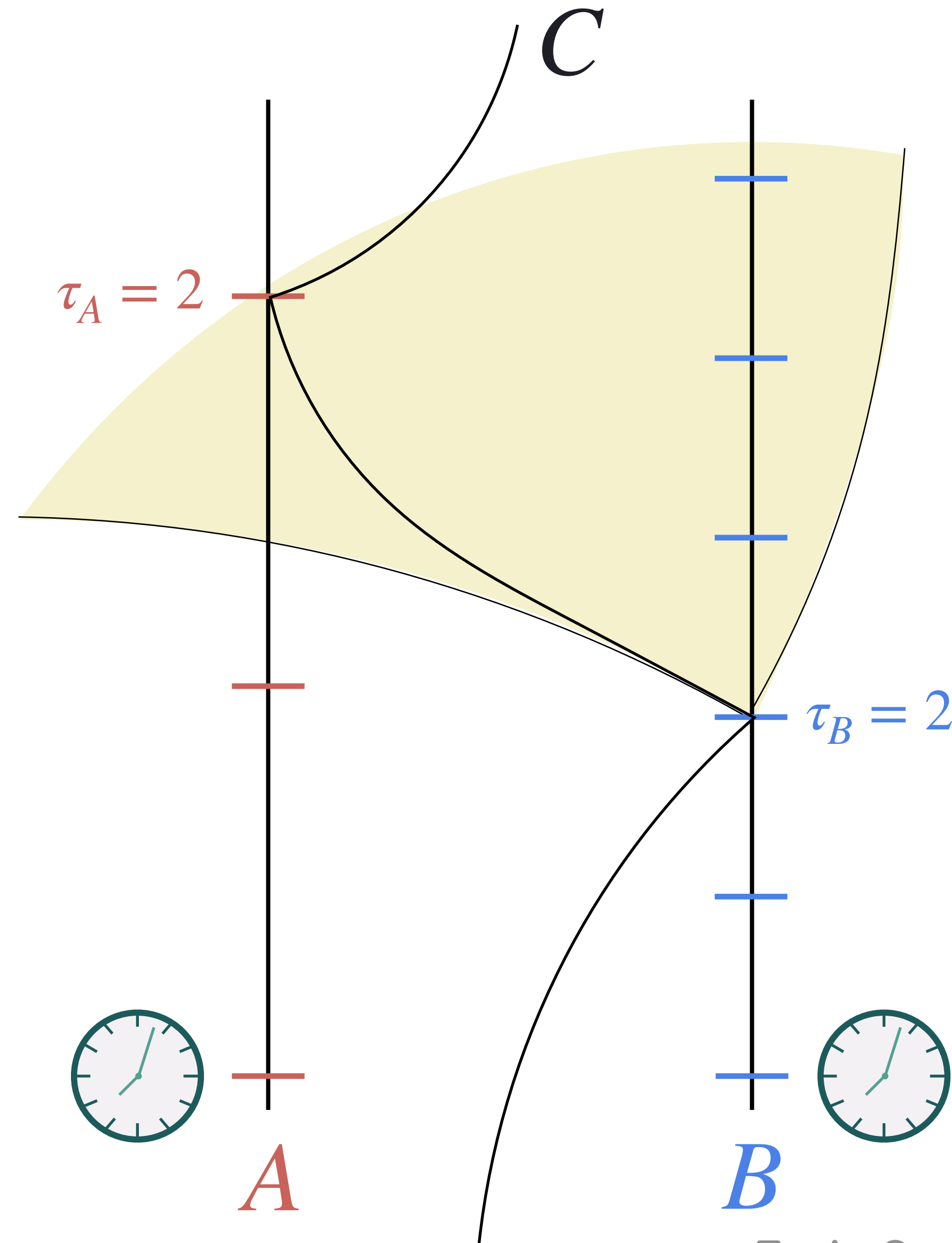
Measurement
in $|\pm\rangle_C \langle \pm|$

Chiribella, D'Ariano, Perinotti, Valiron, PRA (2013)
Araujo, Costa, Brukner, PRL (2014)

THE GRAVITATIONAL QUANTUM SWITCH



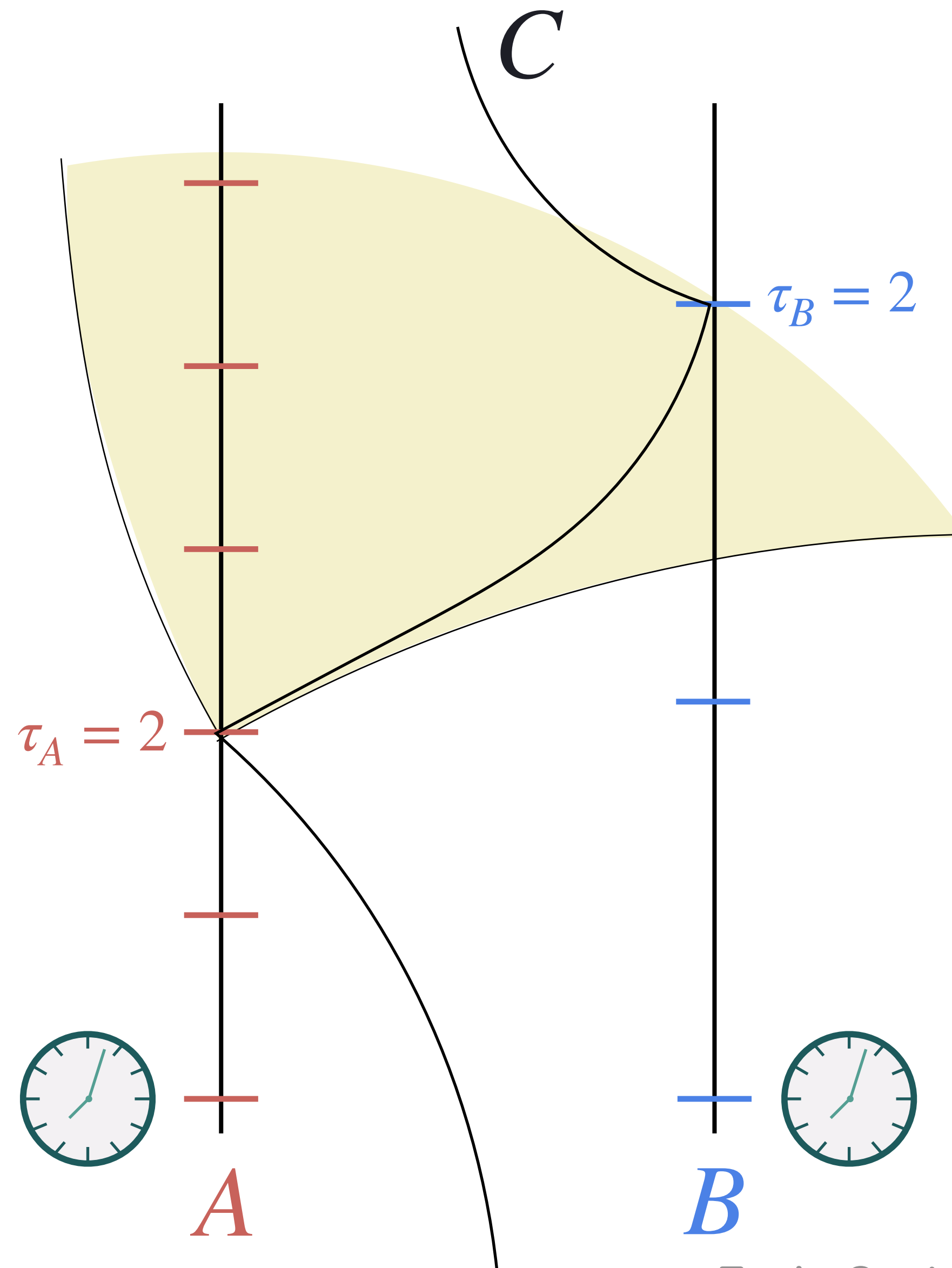
$$|L\rangle_M \hat{U}_A \hat{U}_B |\psi\rangle_S$$



Zych, Costa, Piovski, Brukner, Nat. Commun. (2019)

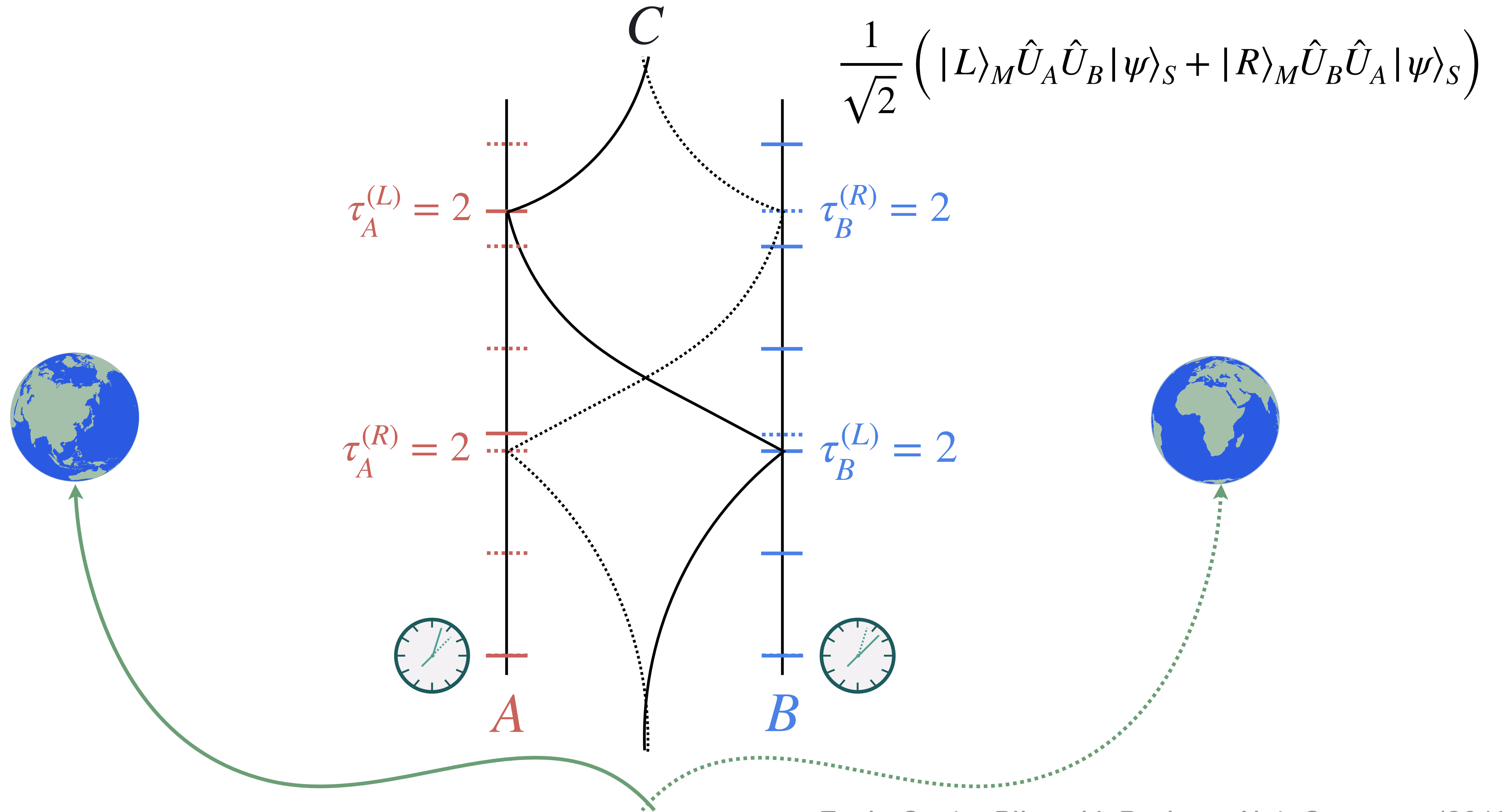
THE GRAVITATIONAL QUANTUM SWITCH

$$|R\rangle_M \hat{U}_B \hat{U}_A |\psi\rangle_S$$



Zych, Costa, Piovski, Brukner, Nat. Commun. (2019)

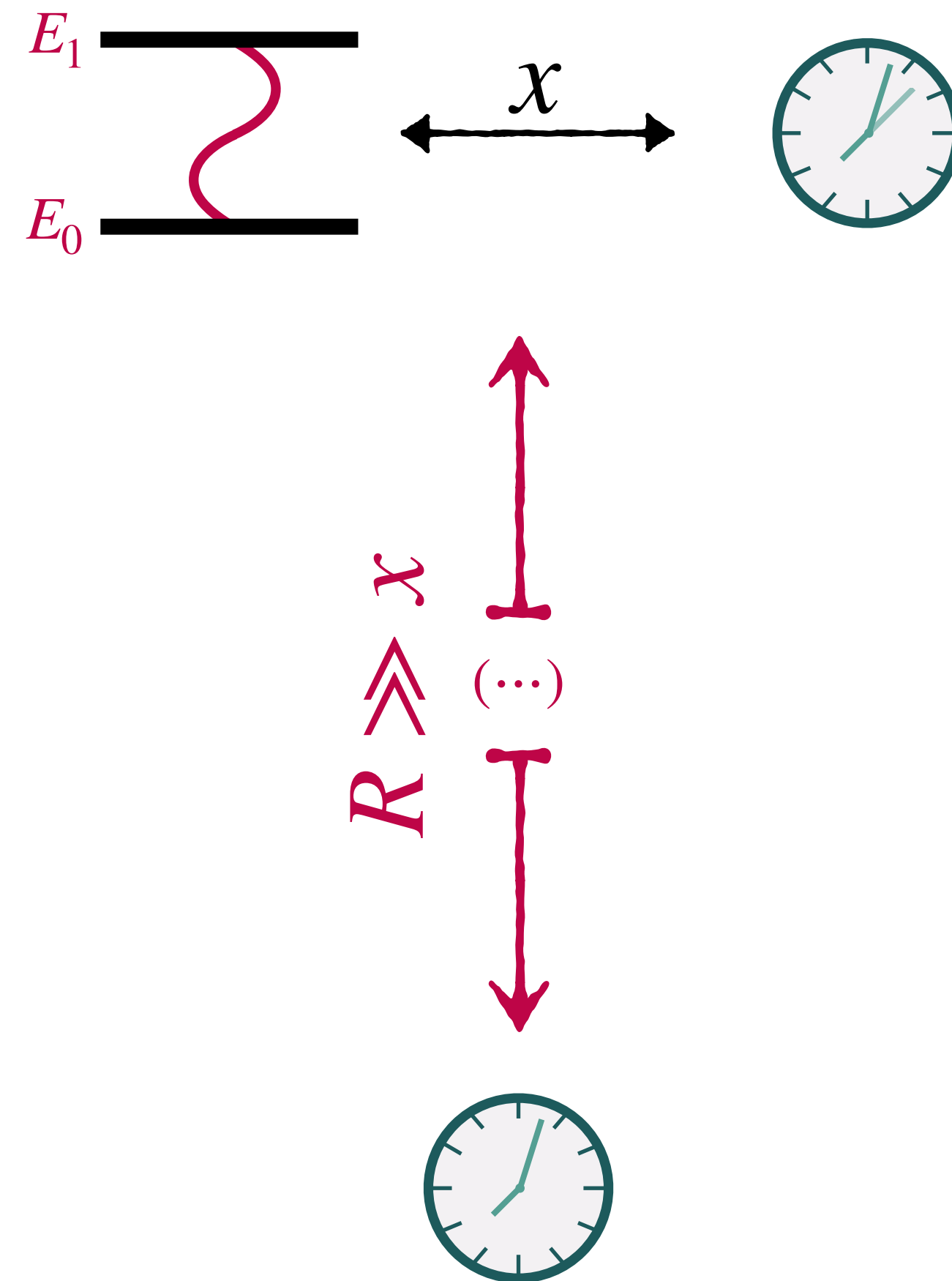
THE GRAVITATIONAL QUANTUM SWITCH



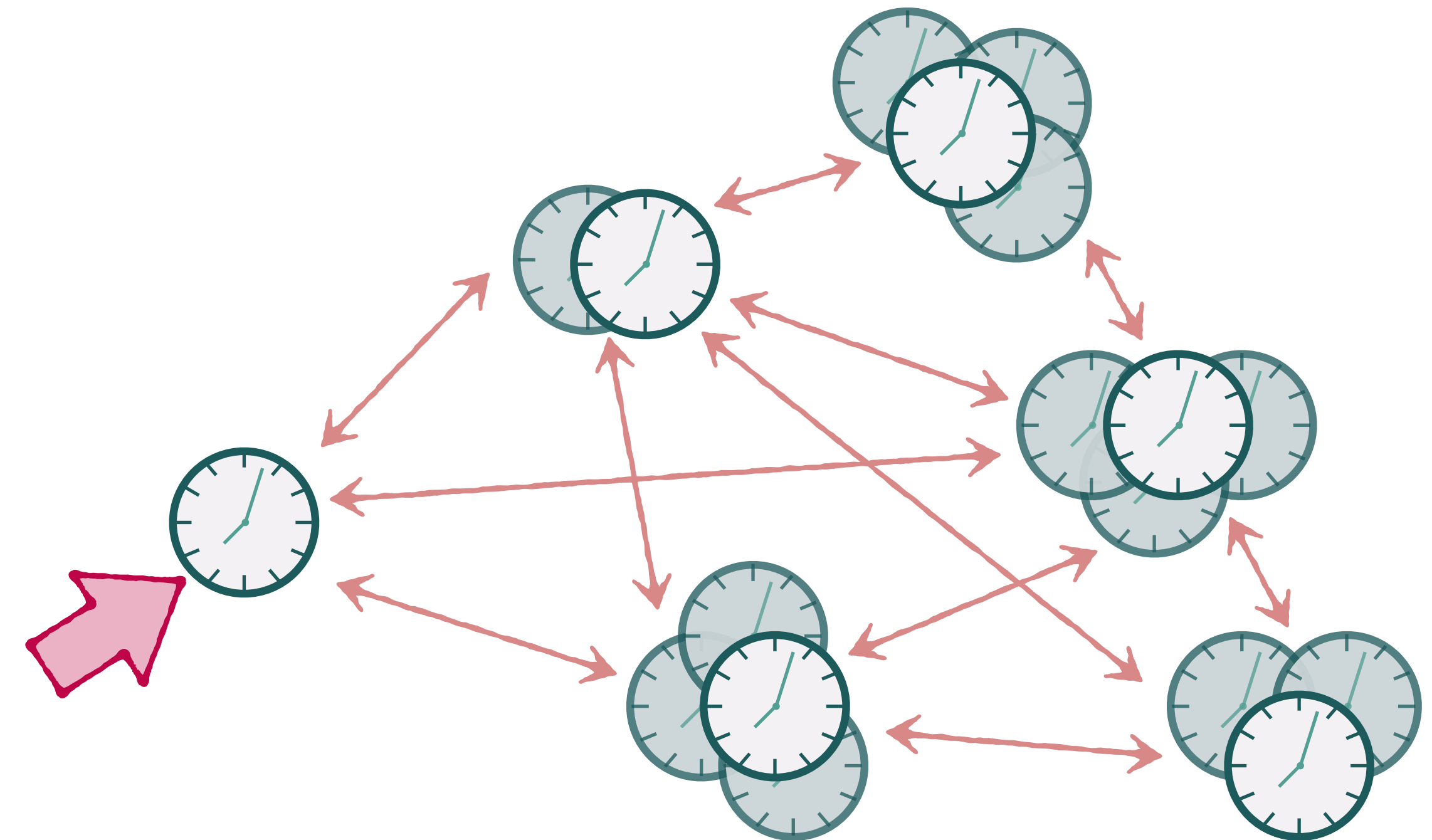
Zych, Costa, Piovski, Brukner, Nat. Commun. (2019)

WHO WRITES THE QUANTUM STATE?

The far-away observer picture



The quantum clock “local” picture



Can we “stand” on different clocks and describe the quantum dynamics from their point of view?

HOW DO QUANTUM CLOCKS TELL THE TIME?

PAGE-WOOTTERS FORMALISM

$$\hat{C} = \hat{H}_A \otimes \mathbb{I}_B + \mathbb{I}_A \otimes \hat{H}_B$$

$$[\hat{T}_i, \hat{H}_i] = i\hbar \quad \hat{C} |\Psi_{ph}\rangle = 0$$

$$|\Psi_{ph}\rangle = \int d\tau e^{\frac{i}{\hbar} \hat{C} \tau} |\phi\rangle_{AB}$$

Why is the state “timeless”?

$$e^{-\frac{i}{\hbar} \hat{C} \tau'} |\Psi_{ph}\rangle = \int d\tau e^{\frac{i}{\hbar} \hat{C} (\tau - \tau')} |\phi\rangle_{AB} = |\Psi_{ph}\rangle$$

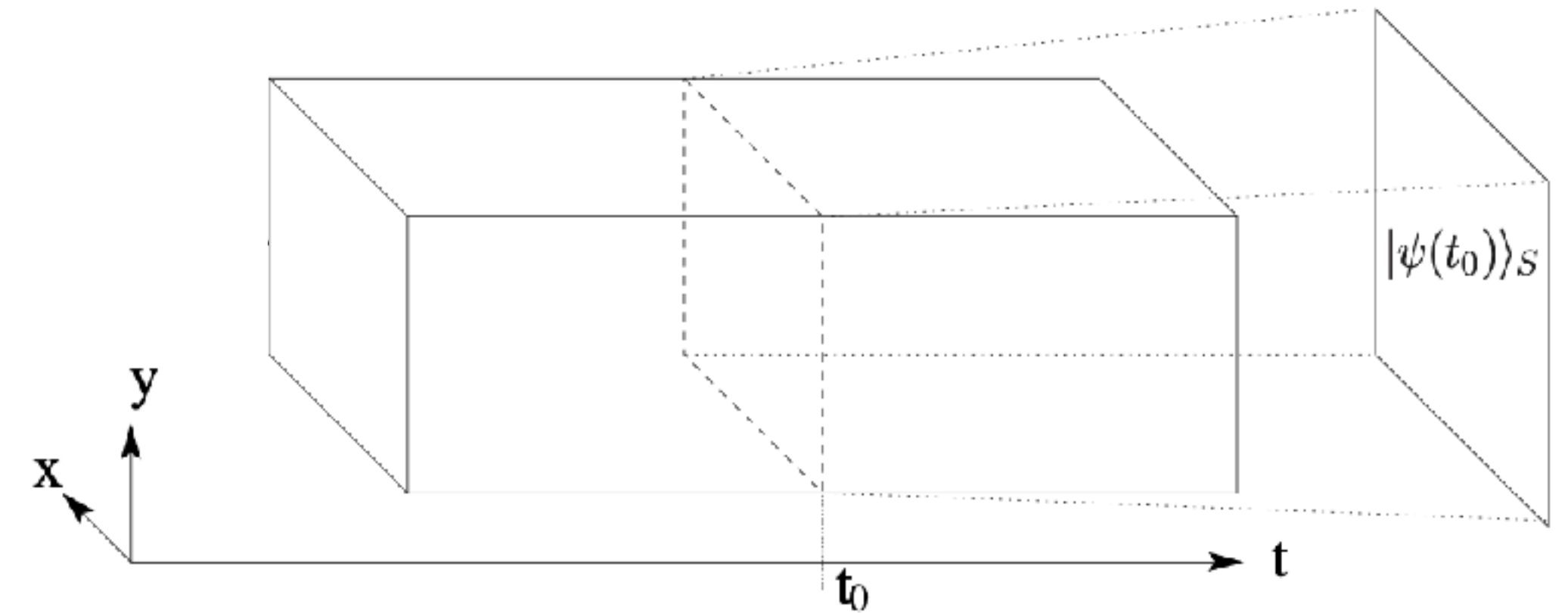


figure from Giovannetti, Lloyd, Maccone, PRD (2015)

Page, Wootters, PRD (1983)
Giovannetti, Lloyd, Maccone, PRD (2015)

PAGE-WOOTTERS FORMALISM

$$\hat{C} = \hat{H}_A \otimes \mathbb{1}_B + \mathbb{1}_A \otimes \hat{H}_B$$

$$[\hat{T}_i, \hat{H}_i] = i\hbar \quad \hat{C} |\Psi_{ph}\rangle = 0$$

How do we recover the dynamics?

$${}_A\langle t | \Psi_{ph} \rangle = |\psi(t)\rangle_B$$

$${}_A\langle t | (\hat{H}_A + \hat{H}_B) | \Psi_{ph} \rangle = 0$$

$$i\hbar \frac{d}{dt} |\psi(t)\rangle_B = \hat{H}_B |\psi(t)\rangle_B$$

Dynamics emerges from correlations between systems

$$|\Psi_{ph}\rangle = \int dt |t\rangle_A |\psi(t)\rangle_B$$

History state

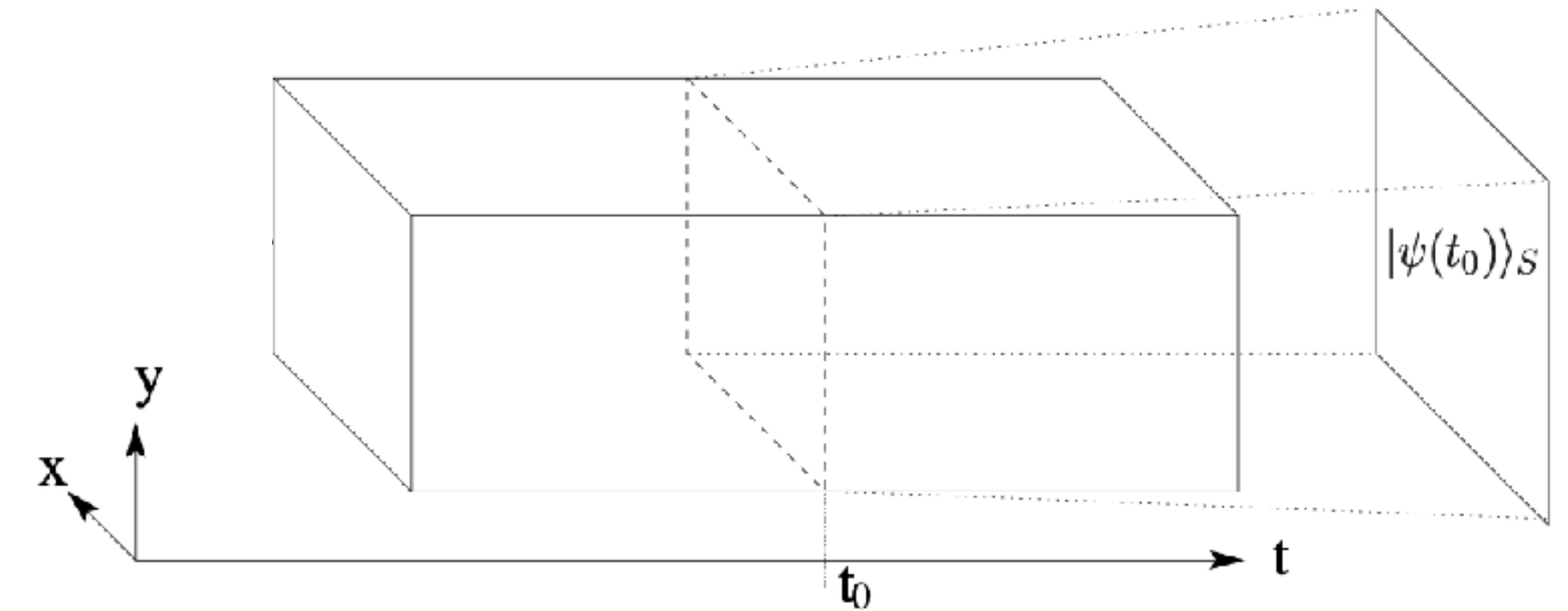


figure from Giovannetti, Lloyd, Maccone, PRD (2015)

Page, Wootters, PRD (1983)
Giovannetti, Lloyd, Maccone, PRD (2015)

MEASUREMENTS IN THE PAGE-WOOTTERS FORMALISM

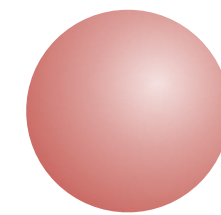
$$[\hat{C}, \hat{O}] = 0$$

Non-evolving quantities?
Restriction of observables?

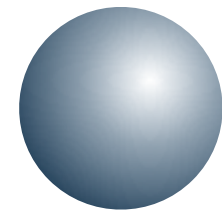
Solution: "Purify" the measurement



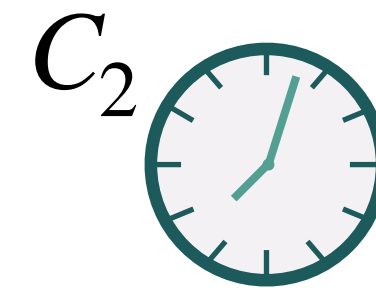
Clocks 1 and 2



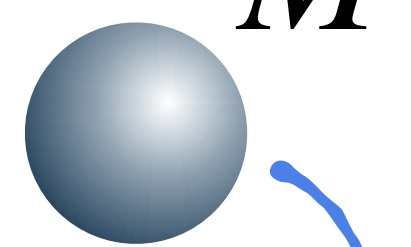
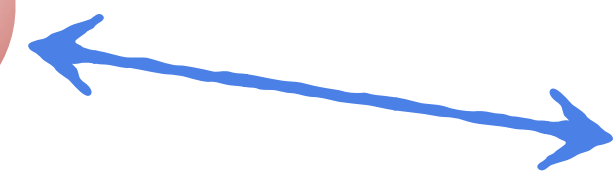
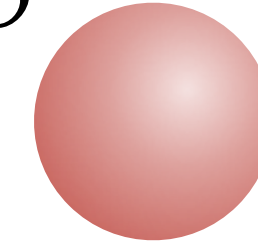
System S



Measurement system M



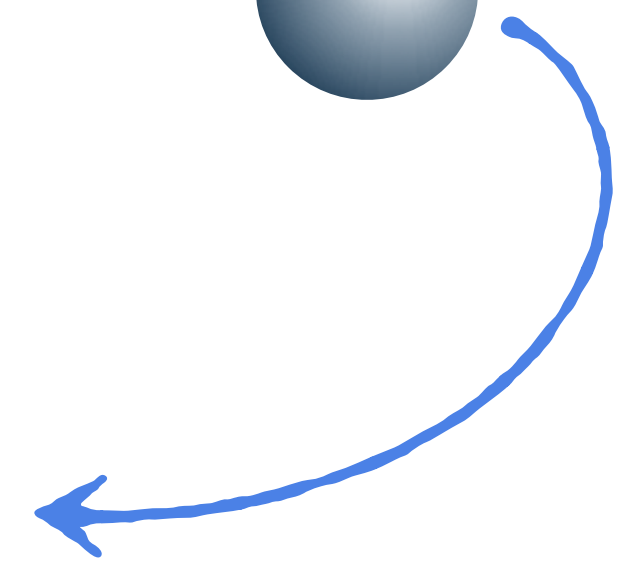
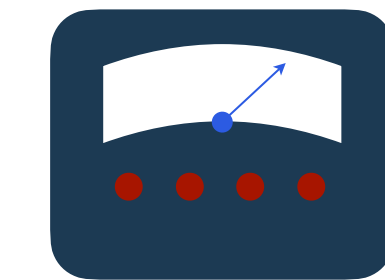
S



M



C₁



Previous Hamiltonian

$$\hat{C} = \hat{H}_1 + \hat{H}_2 + \hat{H}_S + (1 + \lambda \hat{H}_1) \sum_i \delta(\hat{T}_2 - t_i) \hat{K}_i^{MS}$$

Time dilation factor
due to clock 1

Time of measurement
controlled by clock 2

Observable on S and M

F Hellmann, M Mondragon, A Perez, C Rovelli PRD (2007)
V Giovannetti, S Lloyd, L Maccone, PRD (2015)

(GRAVITATIONALLY) INTERACTING QUANTUM CLOCKS

$$\hat{C} |\Psi\rangle_{ph} = 0$$

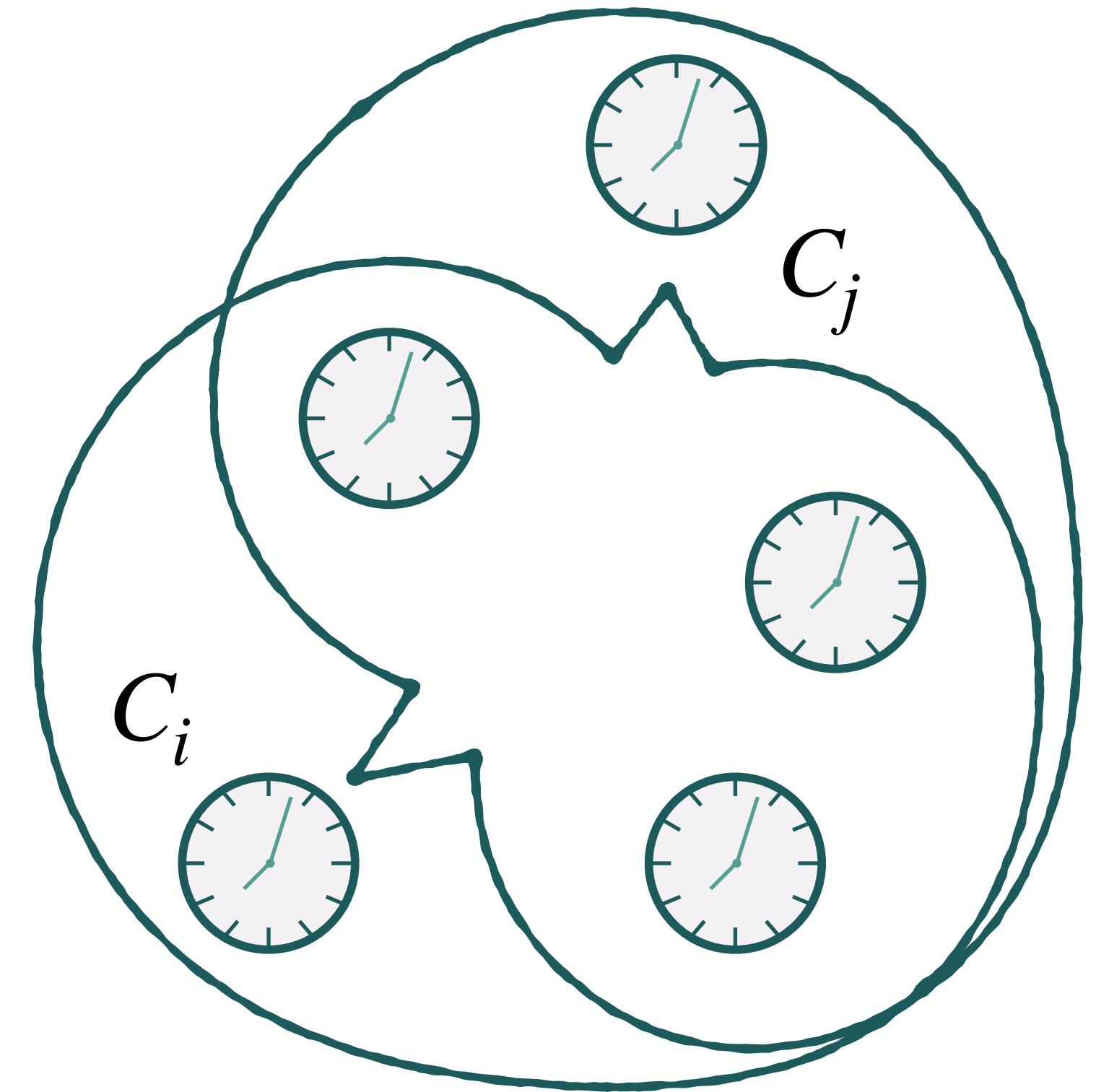
$$\hat{C} = \sum_{k=1}^N \hat{H}_k + \sum_{j<k} \lambda_{jk} \hat{H}_j \hat{H}_k$$

$$\lambda_{jk} = -\frac{G}{c^4 x_{jk}}$$

Perspective of clock i

$${}_i \langle t_i | \Psi \rangle_{Ph} = |\psi(t_i)\rangle^{(i)}$$

$$i\hbar \left(1 + \sum_{k \neq i} \lambda_{ik} \hat{H}_k \right) \frac{d|\psi(t_i)\rangle^{(i)}}{dt_i} = \left(\sum_{k \neq i} \hat{H}_k + \sum_{j < k} \lambda_{jk} \hat{H}_j \hat{H}_k \right) |\psi(t_i)\rangle^{(i)}$$



Check also
Smith Ahmadi, Quantum (2019)

Castro Ruiz, Giacomini, Belenchia, Brukner, Nat. Commun. (2020)

(GRAVITATIONALLY) INTERACTING QUANTUM CLOCKS

$$i\hbar \frac{d|\psi(t_i)\rangle^{(i)}}{dt_i} = \left(1 - \sum_{k \neq i} \lambda_{ik} \hat{H}_k\right) \left(\sum_{k \neq i} \hat{H}_k + \sum_{j < k} \lambda_{jk} \hat{H}_j \hat{H}_k\right) |\psi(t_i)\rangle^{(i)}$$

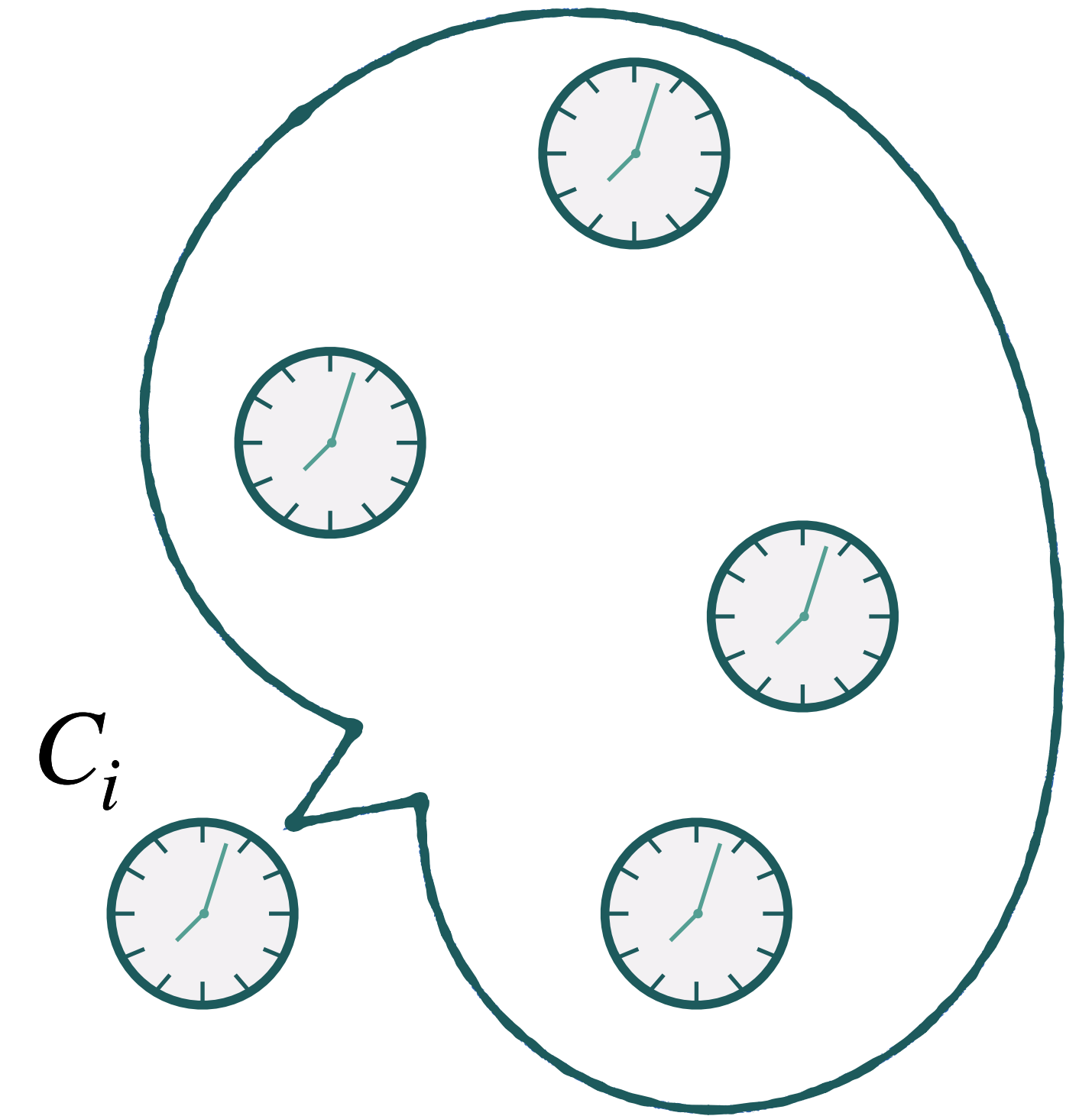
$$\lambda_{jk} = -\frac{G}{c^4 x_{jk}}$$

RESULTS

1. Recover clock Hamiltonian (far-away observer and clock description)

$$H = H_A + H_B - \frac{G}{c^4 x} H_A H_B$$

2. Relative localisation of events
3. Gravitational quantum switch



Castro Ruiz, Giacomini, Belenchia, Brukner, Nat. Commun. (2020)

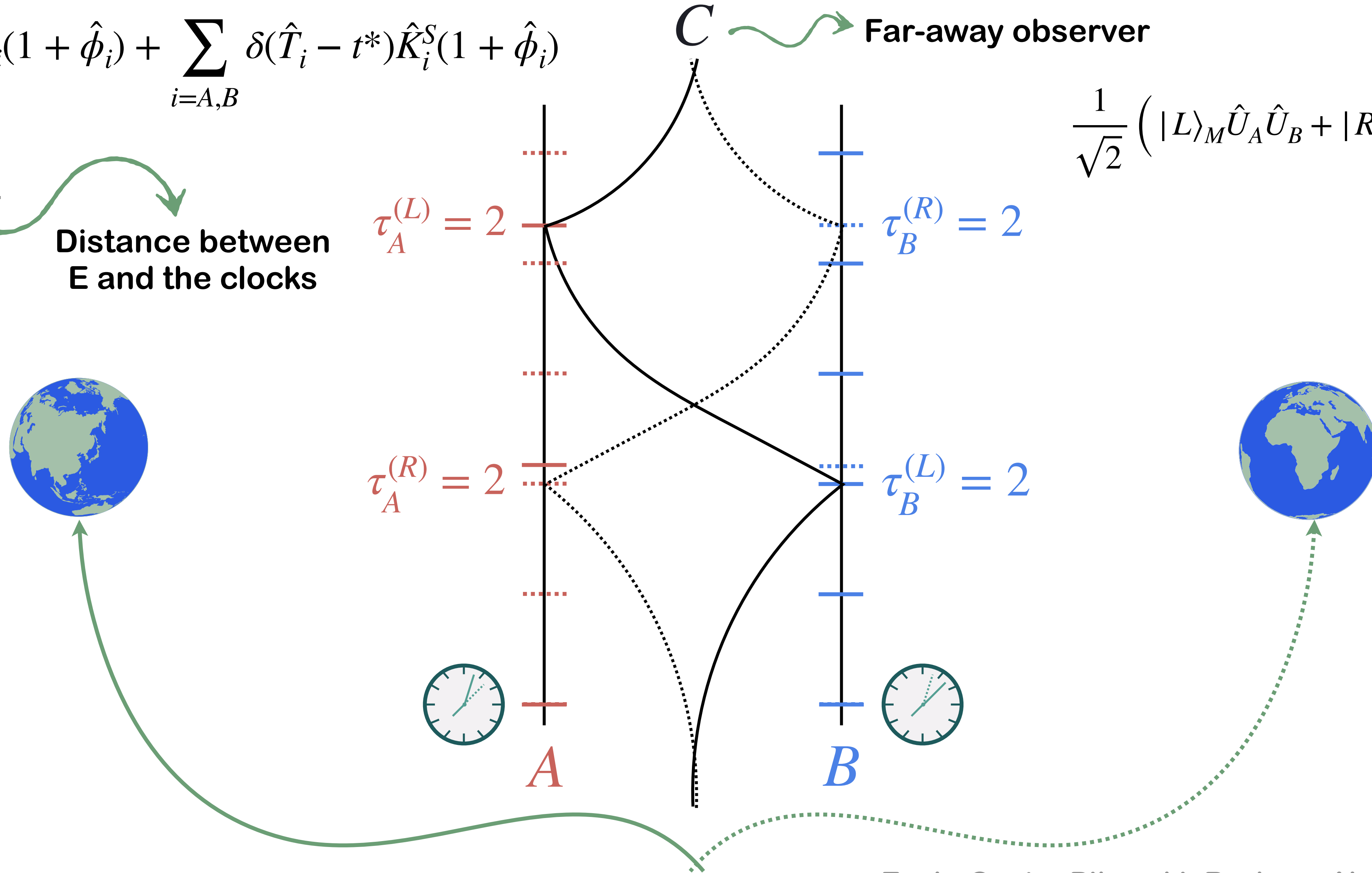
THE GRAVITATIONAL QUANTUM SWITCH

$$\hat{C} = \sum_{i=A,B,C} \hat{H}_i(1 + \hat{\phi}_i) + \sum_{i=A,B} \delta(\hat{T}_i - t^*) \hat{K}_i^S(1 + \hat{\phi}_i)$$

$$\hat{\phi}_i = -\frac{GM_E}{c^2 \hat{x}_i}$$

Distance between E and the clocks

$$\frac{1}{\sqrt{2}} \left(|L\rangle_M \hat{U}_A \hat{U}_B + |R\rangle_M \hat{U}_B \hat{U}_A |\psi\rangle_S \right)$$



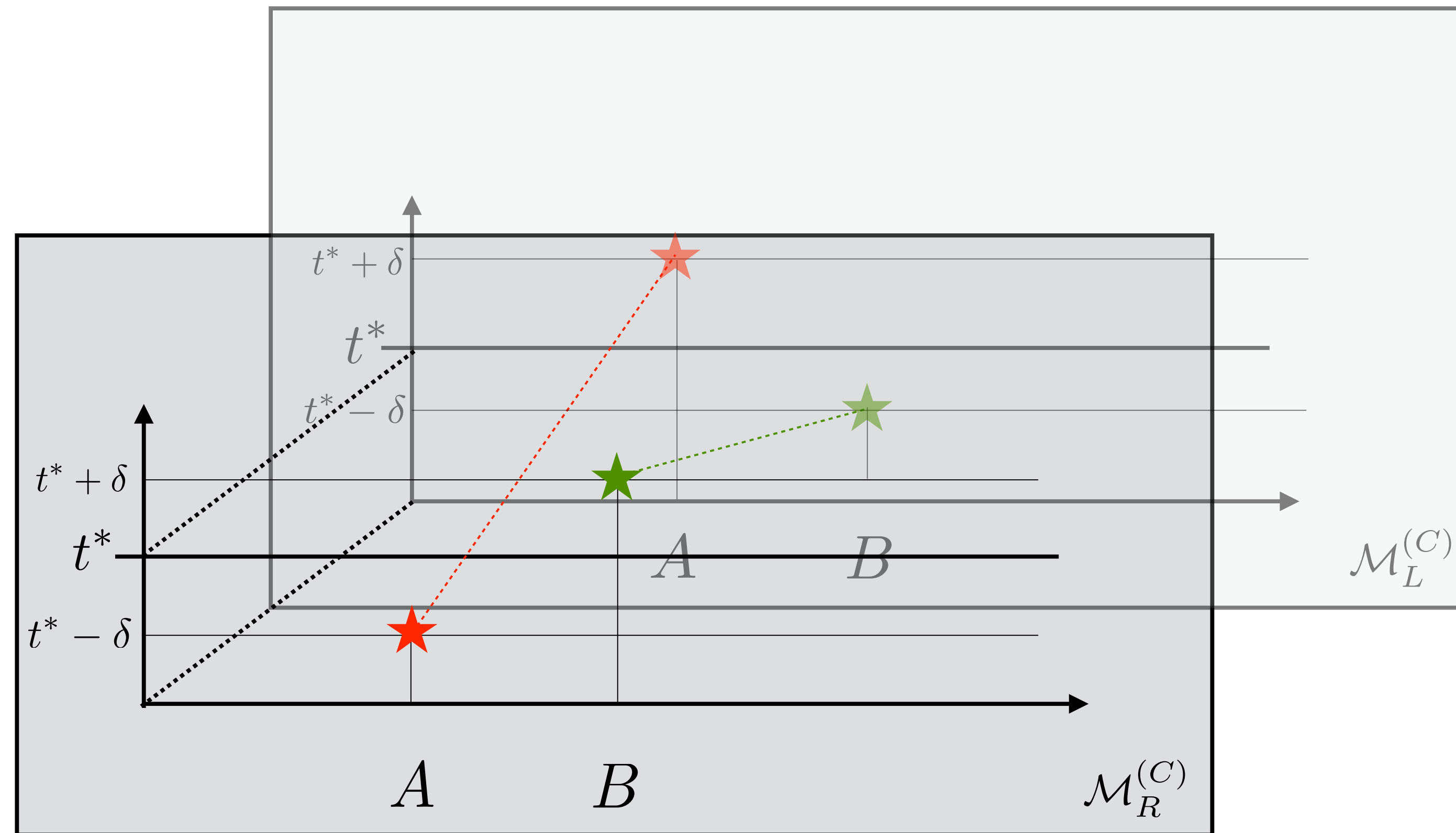
Zych, Costa, Piovski, Brukner, Nat. Commun. (2019)

RELATIVE LOCALISATION OF EVENTS IN THE GRAVITATIONAL SWITCH

$$\hat{C} = \sum_{i=A,B,C} \hat{H}_i(1 + \hat{\phi}_i) + \sum_{i=A,B} \delta(\hat{T}_i - t^*)\hat{K}_i^S(1 + \hat{\phi}_i)$$

$$\hat{\phi}_i = -\frac{GM_E}{c^2 \hat{x}_i}$$

From C's point of view



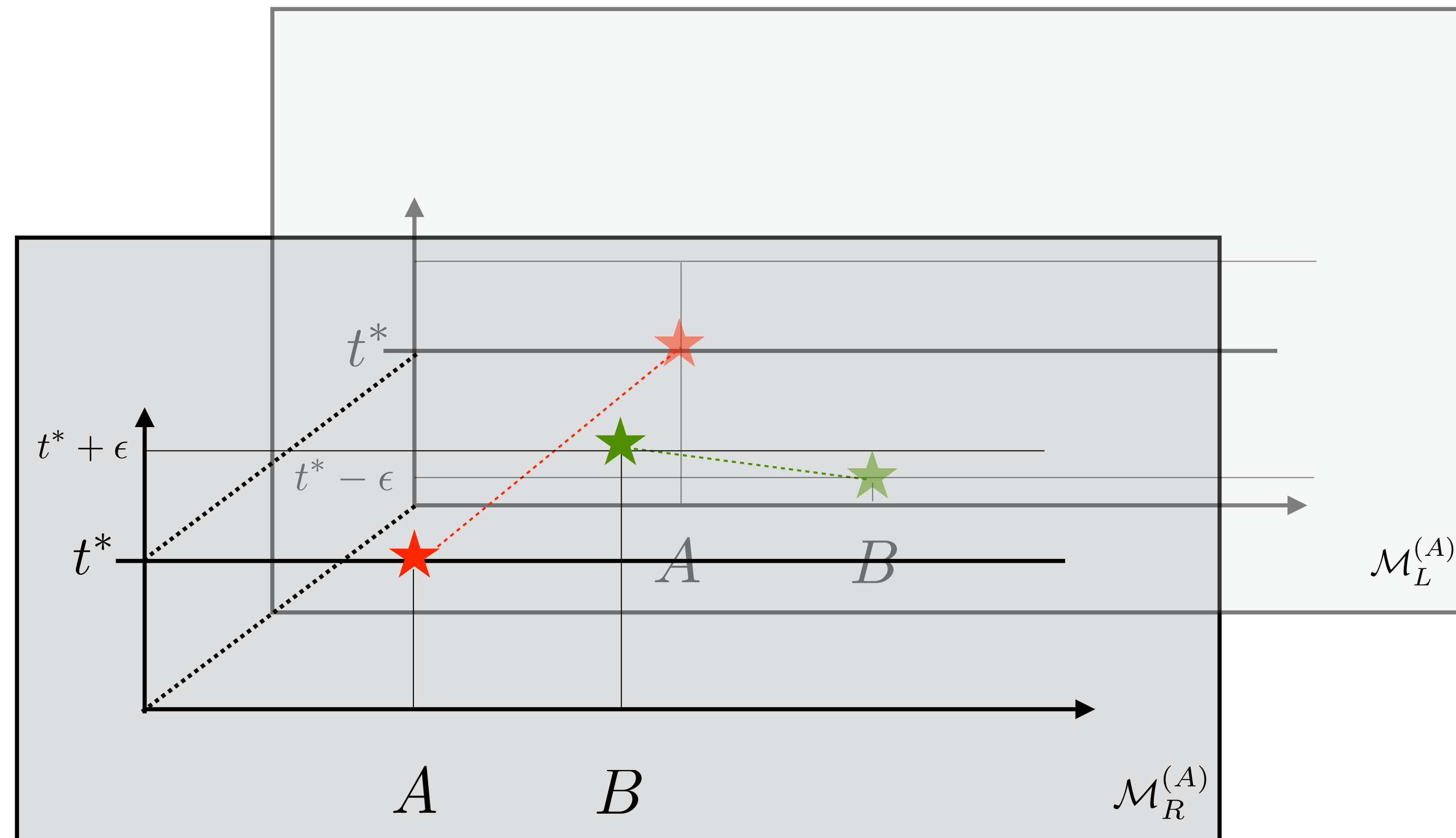
Castro Ruiz, Giacomini, Belenchia, Brukner, Nat. Commun. (2020)

RELATIVE LOCALISATION OF EVENTS IN THE GRAVITATIONAL SWITCH

$$\hat{C} = \sum_{i=A,B,C} \hat{H}_i(1 + \hat{\phi}_i) + \sum_{i=A,B} \delta(\hat{T}_i - t^*) \hat{K}_i^S(1 + \hat{\phi}_i)$$

$$\hat{\phi}_i = -\frac{GM_E}{c^2 \hat{x}_i}$$

From A's point of view



Castro Ruiz, Giacomini, Belenchia, Brukner, Nat. Commun. (2020)

SUMMARY

Quantum clocks are a powerful tool to test the interface between general relativity and quantum theory.

Simplest clock is a two-level system.

Fundamentally, mass-energy equivalence plays a role.

In classical spacetime, quantum clocks can encode information about their position.

In nonclassical spacetime:

- limits to measurability
- relative localisation of events
- relation to indefinite causal order



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