

GR as a Perturbative QFT

#6/6

Predictions, Lessons + Limitations John Donoghue
Sept 22, 2023

- trying it all together

GR - EFT

- know low E. d.o.f + interaction
- don't know "high energy"

\Rightarrow IR effects

In practice

- UV physics \rightarrow "local"
- IR physics \rightarrow "non local"

In calc.

$$\frac{1}{g^2 - M^2} \rightarrow -\frac{1}{M^2} + \frac{g^2}{M^4} + \dots$$

$$\rightarrow \frac{\mathcal{L}(x)}{M^2} + \frac{\mathcal{L}(x)}{M^4} + \dots$$

$\log g^2 \sim$ non analytic \rightarrow non local

Reframe the Program

- symmetric, $g_{\mu\nu}$
- Energy / derivative / curvature expansion

$$S = \int d^4x \sqrt{-g} \left[-\Lambda + \frac{2}{\kappa^2} R + C_1 R^2 + C_2 R_{\mu\nu} R^{\mu\nu} + \dots \right]$$

$$R \sim \partial^2$$

$$R^2 \sim \partial^4$$

Note: Gauss Bonnet

$$E = R_{\mu\nu\alpha\beta} R^{\mu\nu\alpha\beta} - 4 R_{\mu\nu} R^{\mu\nu} + R^2$$

- total deriv in 4D

$$\chi = \int d^4x \sqrt{-g} E$$

\Rightarrow don't need $R_{\mu\nu\alpha\beta} R^{\mu\nu\alpha\beta}$

Match or measure coefficients

$$\Lambda \sim (10^{-3} eV)^4 \quad ; \quad \kappa^2 = 32\pi G$$

Bounds on C_1, C_2

Roughly $(\nabla^2 + \frac{1}{\alpha} \nabla^4) \phi = \kappa^2 \phi$

$$\frac{1}{\alpha^2 - \frac{1}{\alpha} \nabla^4} = \frac{1}{\alpha^2} - \frac{1}{\alpha^2 - \frac{1}{\alpha} \nabla^4} \sim m^2$$

$$V(r) \sim \frac{Q}{r^2} \left[1 - e^{-M_2 r} \right]$$

$$\Rightarrow C_1, C_2 \leq 10^{65}$$

* Ostrogradsky comment: 1850

$$\begin{cases} \varphi_1 = \phi, & \varphi_2 = \dot{\phi} \\ \pi_1 \sim \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \dot{\phi}, & \pi_2 = -\frac{\partial \mathcal{L}}{\partial \phi} \end{cases} \quad \mathcal{L} = \dot{\phi} \ddot{\phi} + \phi \frac{\ddot{\phi}^2}{m^2} + V(\phi)$$

$\rightarrow H = \pi_1 \dot{\phi}_2 + \dots$

not positive definite \Rightarrow instability

But not problem in EFT:

QED

$$\mathcal{L} = F^2 + \frac{F \square F}{m^2}$$

\downarrow small

no instability
when treated
as perturbation

Renormalization

- EFT is wrong in UV
- UV physics is local - ~~local~~
- renormalize an "unrenormalizable" theory
- general \mathcal{L} _{local}
- if symmetry \rightarrow parameter in local \mathcal{L} preserve \sim^{**}

Background Field Method

't Hooft Veltman

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + \kappa h_{\mu\nu}(x)$$

Lesson 3

$$\mathcal{L} = \frac{-2}{\kappa^2} R = \sqrt{-\bar{g}} \left[\frac{-2}{\kappa^2} \bar{R} - \frac{1}{\kappa} [\bar{\nabla}_\alpha \bar{R} - 2\bar{R}^\alpha_\nu \bar{h}^\nu_\alpha] \right. \\ \left. + \frac{1}{2} \bar{D}_\alpha h_{\mu\nu} \bar{D}^\alpha h^{\mu\nu} - \frac{1}{2} \bar{D}_\alpha h^\lambda_\lambda \bar{D}^\alpha h^\sigma_\sigma \right. \\ \left. + \bar{D}_\nu h^\lambda_\lambda \bar{D}^\beta h^\nu_\beta - \bar{D}_\nu h_{\alpha\beta} \bar{D}^\alpha h^{\nu\beta} \right. \\ \left. - \bar{R} \left[\frac{1}{2} (h^\lambda_\lambda)^2 - \frac{1}{2} h^\alpha_\beta h^\beta_\alpha \right] + \bar{h}_\lambda h^\alpha_\nu \bar{R}^\nu_\alpha + 2\bar{h}^\alpha_\beta h^\beta_\alpha \bar{R}^\alpha_\nu \right]$$

gauge
invariant
wrt background

\Rightarrow Symmetry preserved \checkmark

Like massless scalar (Lesson 4)

$$\int_{\mathbb{R}} \sim \int \frac{d^d k}{k^2} \rightarrow 0$$

$$\int_{\mathbb{R}^d} \sim \int d^d k \frac{1}{k^2 (k^2 - \epsilon)^2} \rightarrow \mathbb{R} \left[\frac{1}{\epsilon} - \ln \pi \right] \mathbb{R}$$

\uparrow
Barvinsky-Vilkulinsky

Use heat kernel $t \rightarrow H - V$

$$\mathcal{L} = \frac{1}{16\pi^2} \frac{1}{\epsilon} \left[\frac{1}{120} R^2 + \frac{7}{20} R_{\mu\nu} R^{\mu\nu} \right]$$

\Rightarrow shift in c_1, c_2 (absorb into "measured value")

"Pure gravity is one loop finite"

- pure gravity satisfies $R_{\mu\nu} = 0$
only graviton $\leftarrow *$

- but matter exists

But, unreliable

\Rightarrow shift to low E

Gravitational potential

- Scattering potential $\Rightarrow NR \Rightarrow F.T.$

Leading

$$\cdot \underbrace{\{g\}}_{NR} \xrightarrow{NR} \frac{\chi^2 M m}{g^2} \xrightarrow{FT} -\frac{GMm}{R}$$

Next order local terms

$$c, R, c_2 h_{\mu\nu} R^{\mu\nu}$$

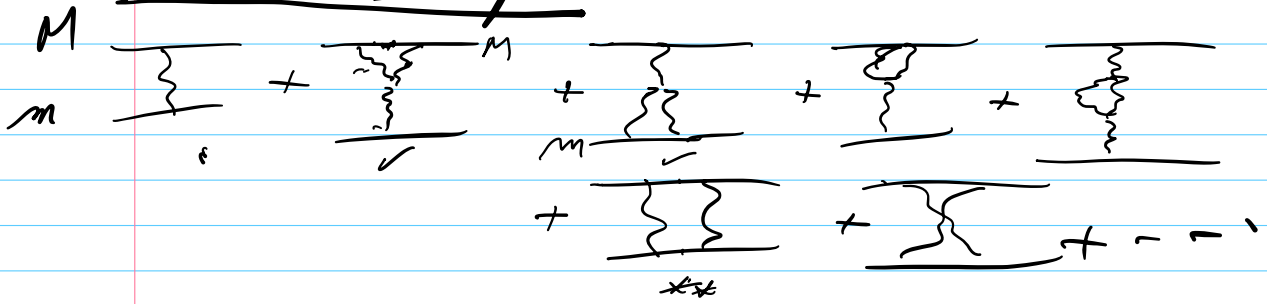
$$\sim \underbrace{c g^4}$$

$$\underbrace{\left\{ \begin{array}{l} \leftarrow \frac{1}{g^2} \\ \leftarrow \frac{1}{g^2} \end{array} \right\} c g^4}_{\sim \text{const}} \xrightarrow{FT} \delta^3(x)$$

$$\text{Coeff (+div)} \sim \delta^3(x)$$

$$\frac{e^{-mR}}{R} \rightarrow \frac{1}{m^2} \delta^3(x) \quad \text{rep of } \delta \text{ functions}$$

No do loops:



$$M = \frac{GMm}{r^2} \left[1 + \frac{4}{5} \frac{v^2}{c^2} + \frac{4}{5} \frac{v^2}{c^2} \ln \frac{r}{\lambda} + K \frac{v^2}{c^2} \right]$$

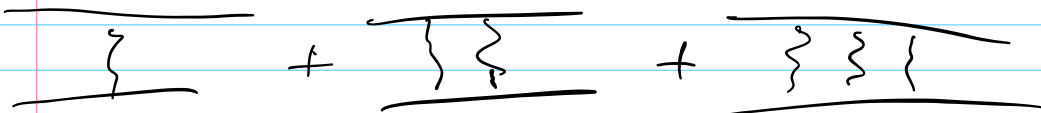
\downarrow FT

$$V(r) = \frac{GMm}{r} \left[1 + \frac{G(M+m)}{rc^2} + \frac{Gh}{5r^2c^3} \right] + \left(\frac{3}{5} \frac{GMm}{rc^2} \right)$$

\leftarrow quantum

Calculate:

$$V(r) = -\frac{GMm}{r} \left[1 + 3 \frac{G(M+m)}{rc^2} + \frac{41}{16\pi} \frac{Gh}{r^2} \right] \checkmark$$



\Rightarrow "Coulomb" phase } exponentials
 "Klein" phase } Weinberg
 $|e^{i\theta} \mathcal{M}_0|^2$

Recalculated "Unitarity" Technique

- Feynman's tree theorem



~ trees, bubble v...

$$GR \sim (\sqrt{M})^2$$

(Compton amplitude)

Low energy then

Bending of Light - massless particle



⇒ scattering amps

⇒ bending angle "Eikonal"

$$\theta = \frac{4GM}{b} + \frac{15 G^2 M^2 \pi}{4 b^2} + \frac{8c_b - 47 + 64 \log(2r_0/b)}{\pi} \frac{G^2 M \hbar}{b^3}$$

where $c_b = (371/120, 113/120, -29/8)$ for scalars, photons and gravitons

not universal

Graviton graviton

x

Lessons:

1) Potential is universal

- Ross + Weinstein
- understood "unitarity" technique
- Amp are universal \Rightarrow One loop amp is universal

2) Classical effects from loop

- "loop expansion is \hbar expansion" is false

$$\psi(i\hbar \not{\partial} - m)\psi \rightarrow \hbar \overline{\psi}(i\hbar \not{\partial} - \frac{m}{\hbar})\psi$$

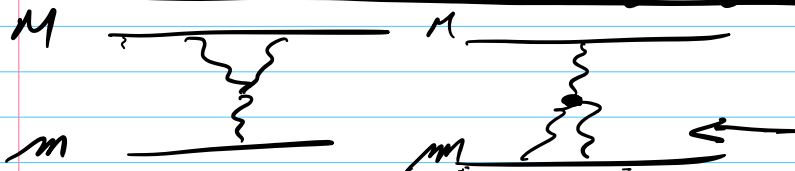
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$$\hbar \sqrt{\frac{g^2 m^2}{\hbar^2}} \sim \text{Classical}$$

* Beane subfield - gravity waves from QFT

3) No "test mass limit" for quantum effect



Class. $M\sqrt{g}$ $m\sqrt{g}$

Quant. $g^2 \ln$ $g^2 \ln$

4) No "quantum metric"

- no quant corr to Schw.
- Kinetic not invariant under field redef
=> Haag's theorem does not apply
- "off shell" partial amplitude
- not a good quantum question
not $\Delta^2 g_{\mu\nu}$

5) Massless particles not universal

6) Geodesics are ill defined:

- quantum \sim "tidal"



- small in EFT limit
- uncontrolled away from limit

7) G, Λ are ~~not~~ running parameters - in physical processes

Limits of EFT:

High energy / curvature

$$M \sim \frac{1}{g^2} \left[1 + \frac{K^2}{g^2} \ln + \frac{K^4}{g^4} \ln \right]$$

⇒ need some UV complete theory

IR issues:

- grav effects build up
- eg BH

$$g_{\mu\nu}^{(w)} = \gamma_{\mu\nu} + \frac{1}{3} R_{\mu\nu\alpha\beta}(w) \underbrace{(K'-N)^\alpha (K'-N)^\beta}$$

- technical
- quantum effects "feel" long distances
 $\sim \frac{1}{r^3}$

- Propagating past BH

⋮

Our core theory

$$Z = \int [d\phi d\psi dA dg] \exp i \int d^4x \sqrt{g} \left[\mathcal{L}_{SM} \right]$$

limit of knowledge

$$- \Lambda + \frac{2}{K^2} R + c_1 R^2 + c_2 R^{\mu\nu} R_{\mu\nu}$$

$$+ \mathcal{L}_{SM}$$

GR forms a QFT
- EFT types

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