

DOUBLE PARTON DISTRIBUTIONS ON THE LATTICE USING LARGE MOMENTUM EFFECTIVE THEORY

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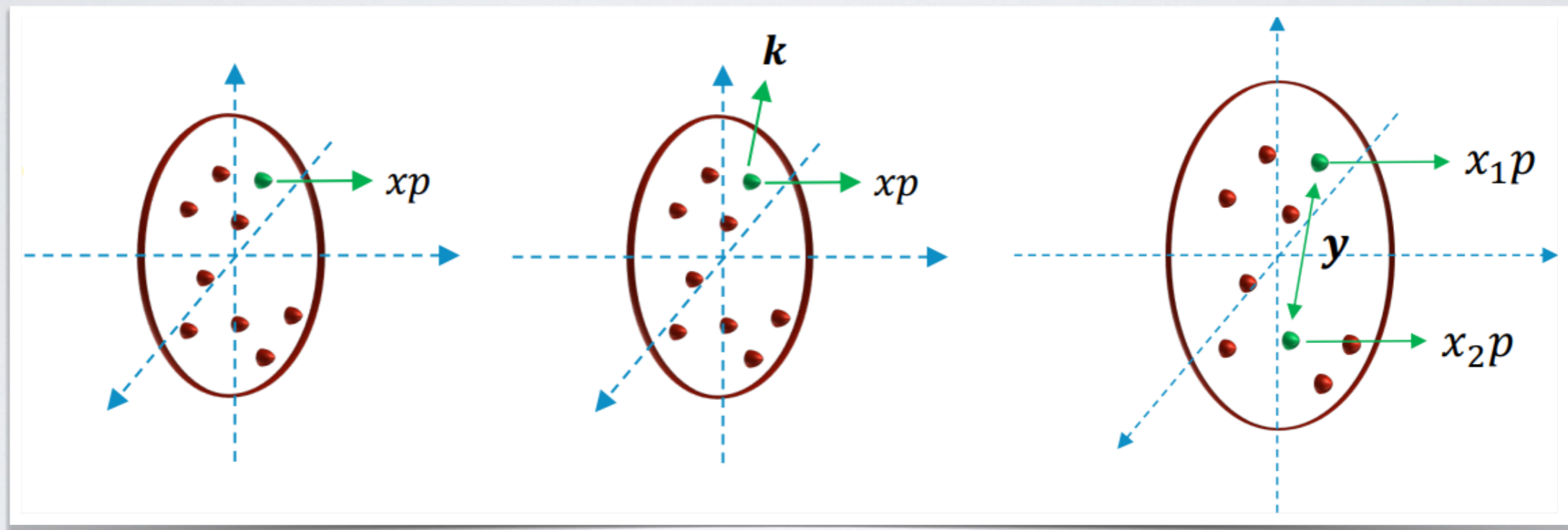
Based on work with Max Jaarsma
& Wouter Waalewijn

MPI@LHC, 22nd November 2023

OUTLINE

- Important features of PDFs and DPDs
- Time and the lattice
- Quasi-PDFs and LaMET
- Quasi-DPDs
- Results

PDFS, TMDs, AND DPDs



(J. Gaunt, via M. Jaarsma)

- What's in a proton?
- Both TMDs and DPDs introduce a transverse component

PARTON DISTRIBUTIONS

- Quantum mechanics: Probability amplitude:

$$\begin{aligned} f(x) &= \sum_X d\Pi_X \int d^4p \delta(xP^+ - p^+) \delta^4(P - p - p_X) |\langle X | \psi(0) | P \rangle|^2 \\ &= \int_{-\infty}^{\infty} \frac{dy^-}{2\pi} e^{-ixP^+ y^-} \langle P | \bar{\psi}(y^- n^\mu) \gamma^0 \psi(0) | P \rangle \end{aligned}$$

- Light cone separation tied to x-dependence
- Wilson line for gauge invariance:

$$f(x)_q = \int_{-\infty}^{\infty} \frac{dy^-}{2\pi} e^{-ixP^+ y^-} \langle P | \bar{\psi}(y^- n^\mu) \gamma^0 W_n(y^- \leftarrow 0) \psi(0) | P \rangle$$

DPDS

- Double the fun!

[Diehl, Ostermeier, Schaefer, '11]
[Manohar, Waalewijn, '12]

$$\begin{aligned} R_{F_{a_1 a_2}} &= -\pi P^+ \int \frac{db_1^-}{2\pi} \frac{db_2^-}{2\pi} \frac{db_3^-}{2\pi} e^{-ix_1 P^+ b_1^-} e^{-ix_2 P^+ b_2^-} e^{ix_1 P^+ b_3^-} \\ &\quad \times \langle P | T^\dagger \left[[\bar{\psi} W]_n(0^+, b_1^-, \mathbf{b}_\perp) \Gamma_{a_1} R_1 \right]_i \left[[\bar{\psi} W]_n(b_2^-) \Gamma_{a_2} R_2 \right]_j \\ &\quad \times T \left[[W^\dagger \psi]_n(0^+, b_3^-, \mathbf{b}_\perp) \right]_i \left[[W^\dagger \psi]_n(0) \right]_j |P\rangle \end{aligned}$$

- Colour and spin correlations between the extracted partons possible
- Like PDF assumes infinite momentum proton

THE LATTICE

- Idea: Monte Carlo integration for the path integral

$$\langle \phi(x_1)\phi(x_2) \rangle = \int \mathcal{D}\phi e^{iS[\phi]} \phi(x_1)\phi(x_2)$$

- Problem: Weight factor doesn't weigh

⇒ Wick rotate to imaginary time

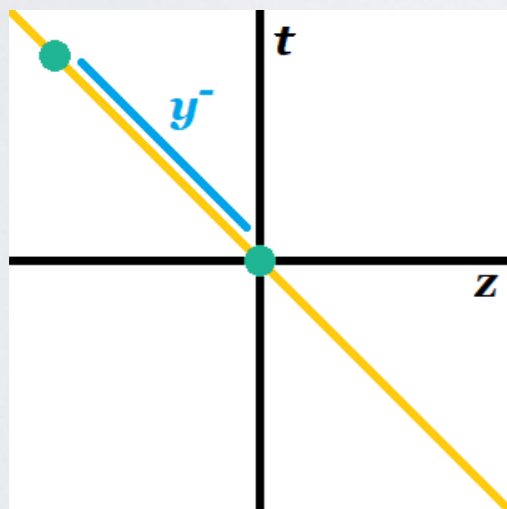
- But:

$$\bar{\psi}(y^{-n^\mu})\gamma^0\psi(0)$$

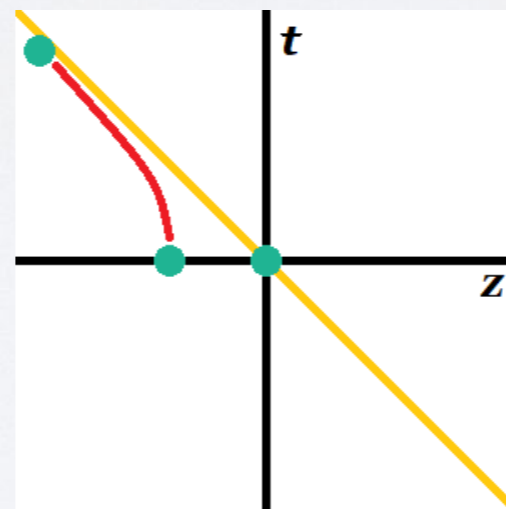
$$n^\mu = e_0^\mu + e_z^\mu$$

TOWARDS LAMET

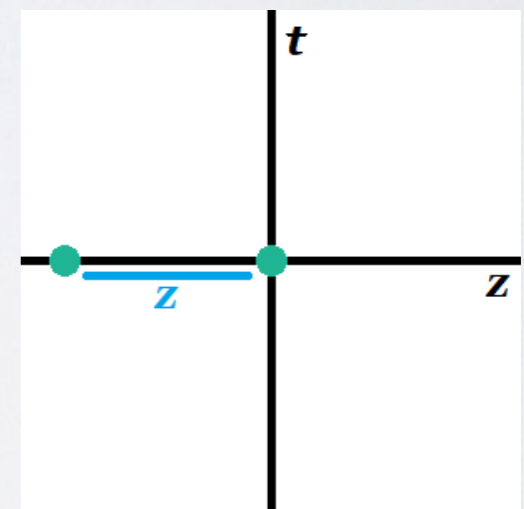
- Special relativity exists
- Boosted spacelike separations look lightlike



We want



We MacGyver



We have

THE QUASI PDF

[1, '13]

$$\tilde{f}(x, P^z) = \int \frac{dz}{4\pi} e^{-ixP^z z} \langle P | \bar{\psi}(z) \gamma^z W(z \leftarrow 0) \psi(0) | P \rangle$$

- Under a boost:

$$\frac{z}{\sqrt{2}}(-1^+, 1^-, \mathbf{0}_\perp) \rightarrow \frac{z}{\sqrt{2}}(-e^{-y_B}, e^{y_B}, \mathbf{0}_\perp)$$

- Proton momentum parametrises boost

- Expectation: $\lim_{P^z \rightarrow \infty} \tilde{f}(x, P^z) = f(x)$

THE QUASI PDF

[Ji, '13]

$$\tilde{f}(x, P^z) = \int \frac{dz}{4\pi} e^{-ixP^z z} \langle P | \bar{\psi}(z) \gamma^z W(z \leftarrow 0) \psi(0) | P \rangle$$

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- Proton momentum parametrises boost

- Expecta

$$\lim_{P^z \rightarrow \infty} \tilde{f}(x, P^z) \neq f(x)$$

A TALE OF TWO LIMITS

Light cone PDF

$$\langle P | \bar{\psi}(y^-) \gamma^+ W(y^- \leftarrow 0) \psi(0) | P \rangle$$

- Lightlike separation
- Infinite momentum P^+

\Rightarrow Lorentz invariant

$$\text{“ } \lim_{\epsilon \rightarrow 0} \lim_{P^+ \rightarrow \infty} \text{”}$$

Quasi PDF


$$\langle P | \bar{\psi}(z) \gamma^z W(z \leftarrow 0) \psi(0) | P \rangle$$

- Spacelike separation
- Finite momentum P^z

\Rightarrow Not Lorentz invariant

$$\text{“ } \lim_{P^z \gg 0} \lim_{\frac{1}{a} \rightarrow \infty} \text{”}$$

LARGE MOMENTUM EFFECTIVE THEORY

- HQET: $m_Q \rightarrow \infty$  After renormalising: No EFT
Before renormalising: EFT
- $P^+ \rightarrow \infty$ defines LaMET
 - P^+ isn't part of the Lagrangian \Rightarrow effective states

LARGE MOMENTUM EFFECTIVE THEORY

- HQET: $m_Q \rightarrow \infty$ $\left\{ \begin{array}{l} \text{After renormalising: No EFT} \\ \text{Before renormalising: EFT} \end{array} \right.$
- $P^+ \rightarrow \infty$ defines LaMET
 - P^+ isn't part of the Lagrangian \Rightarrow effective states
 - The Lightcone PDF is the effective theory object



[Ji, Liu, Liu, Zhang, Zhao, '20]

MATCHING

- Matching is UV physics, and QCD is **perturbative**

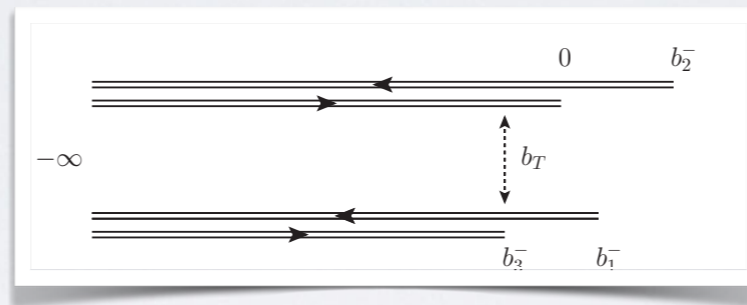
$$\tilde{f}(x, P^z, \mu) = \int \frac{dy}{y} \mathcal{C}\left(\frac{x}{y}, \frac{\mu}{yP^z}\right) f(y, \mu) + \mathcal{O}\left(\frac{M^2}{P_z^2}, \frac{\Lambda_{\text{QCD}}^2}{(xP^z)^2}\right)$$

[Ji, '13]
[Ma, Qiu '14]
[Izubuchi, Ji, Jin,
Stewart, Zhao '18]

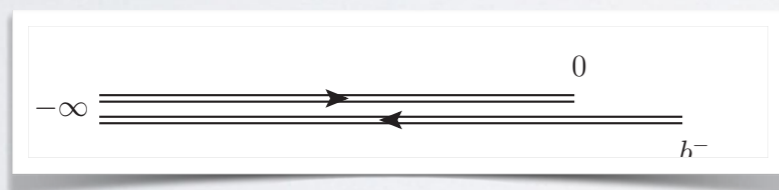
- So what about DPDs?
 - ▶ **Assumption:** The same, up to difficulties
 - ▶ See Jianhui Zhang's talk

WILSON LINE TROUBLE

- There are **four** (Wilson) lines!

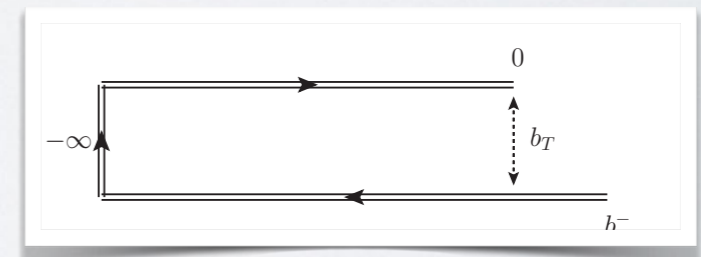


Colour-summed



Like **PDF** (x2)

Colour-correlated



Like **TMD** (x2)

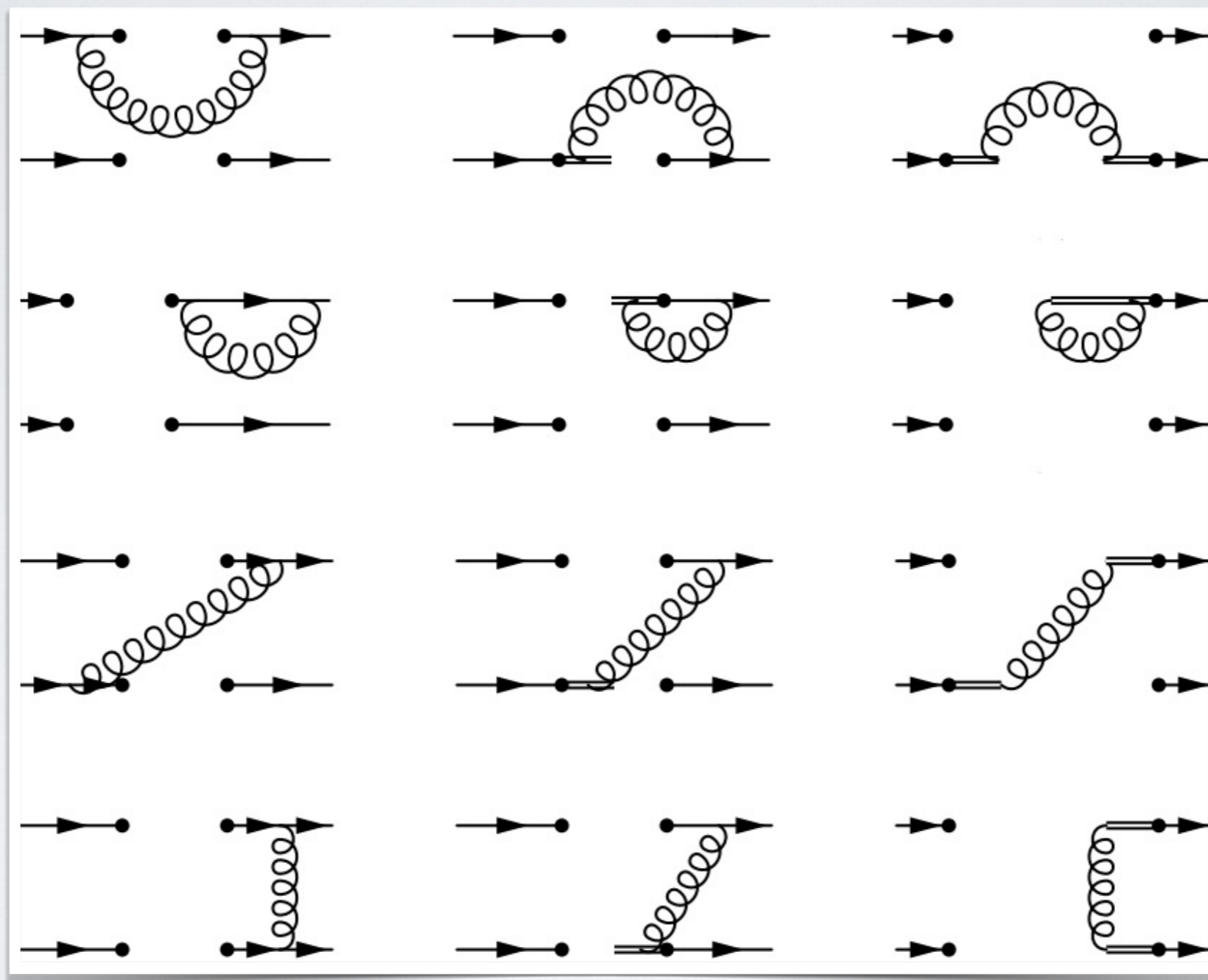
[Ebert, Stewart, Schindler, Zhao '22]

MATCHING CONJECTURE

$$\begin{aligned} & R\tilde{F}_{a_1 a_2}^{\text{sub}}(x_1, x_2, b_\perp, \mu, \tilde{\zeta}, P^z) \\ &= \int \frac{dy_1}{y_1} \frac{dy_2}{y_2} {}_{RR'}C_{a_1 a_2, a'_1 a'_2} \left(\frac{x_1}{y_1}, \frac{x_2}{y_2}, \frac{(x_1 P^z)^2}{\mu^2}, \frac{(x_2 P^z)^2}{\mu^2}, \frac{\tilde{\zeta}}{\mu^2} \right) \\ &\quad \times \exp \left[\frac{1}{2} {}_{R'}\gamma_\zeta(b_\perp, \mu) \log \left(\frac{\tilde{\zeta}}{\zeta} \right) \right] {}_{R'}F_{a'_1 a'_2}^{\text{sub}}(x_1, x_2, b_\perp, \mu, \zeta) \\ &\quad + \text{mixing with single PDFs} \\ &\quad + \mathcal{O} \left(\frac{M^2}{P_z^2}, \frac{1}{x_i b_\perp P^z}, \frac{\Lambda_{\text{QCD}}^2}{(x_i P^z)^2} \right) \end{aligned}$$

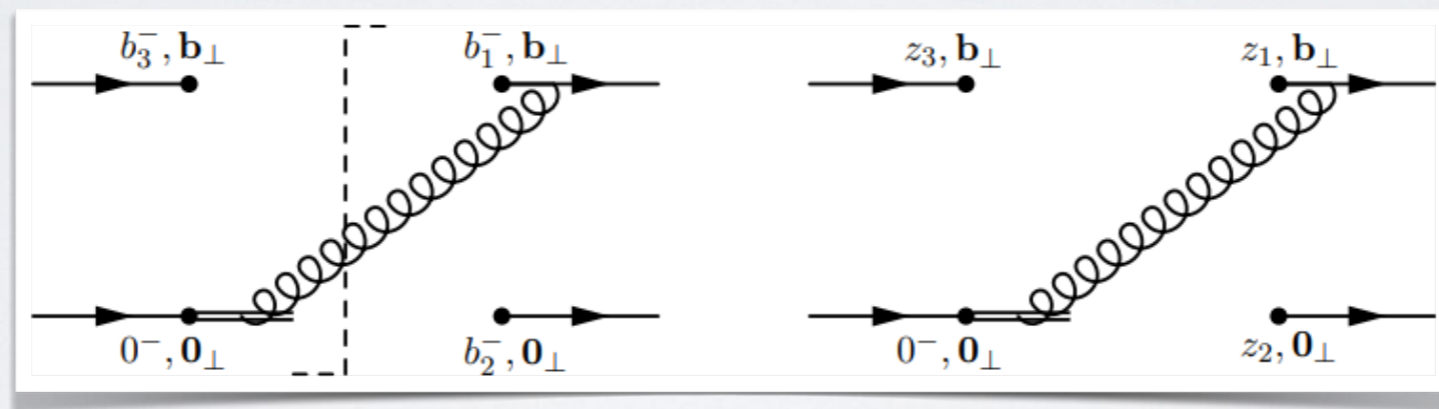
- We don't prove it, but we can **check** it

THE DIAGRAMS



AN EXAMPLE

- Calculate the diagrams and check if the matching is pure UV



$$\Delta^{\text{example}} = -4\pi\delta(1-x_1)\delta(1-x_2) \left[\frac{1}{\epsilon_{\text{ir}}} + \log\left(\frac{\mu^2 \mathbf{b}_\perp^2}{b_0^2}\right) \right] \left[2 + \log\left(\frac{\delta^2}{p_1^+ p_2^+}\right) \right] + \text{other terms}$$

- Individual diagrams have infrared dependence...

THE KERNELS

- ...but the sum of all is pure UV:

$${}^1C_{a_1 a_2} \left(\frac{x_1}{y_1}, \frac{x_2}{y_2}, \frac{\mu}{|y_1|P^z}, \frac{\mu}{|y_2|P^z} \right) = C_{a_1} \left(\frac{x_1}{y_1}, \frac{\mu}{|y_1|P^z} \right) C_{a_2} \left(\frac{x_2}{y_2}, \frac{\mu}{|y_2|P^z} \right)$$

$${}^8C_{a_1 a_2}^{(1)} = \left(1 - \frac{N_c}{2C_F} \right) {}^1C_{a_1 a_2}^{(1)} + \delta \left(1 - \frac{x_1}{y_1} \right) \delta \left(1 - \frac{x_2}{y_2} \right) \\ \times N_c \left[2 \log \left(\frac{\tilde{\zeta}}{\mu^2} \right) - \frac{1}{2} \log^2 \left(\frac{(2y_1 P^z)^2}{\mu^2} \right) - \frac{1}{2} \log^2 \left(\frac{(2y_2 P^z)^2}{\mu^2} \right) - \frac{5}{2} + \frac{\pi^2}{6} \right]$$

- Note also: No mixing of colour and spin structures at 1-loop

CONCLUSIONS AND OUTLOOK

- We verified the conjectured matching of Quasi and light cone DPDs at 1-loop order
- We find relations to the single-parton PDFs, and absence of colour and spin mixing
- Open questions: Proof, soft functions, lattice,...

THANK YOU!

