# ALICE measurements of particle production as a function of event topology in small collision systems

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University of Manchester 20–24 Nov 2023









 $\clubsuit$  Measurements at the LHC have revealed that small collision systems exhibit behaviors formerly thought to be achievable only in heavy-ion collisions, where the data support the formation of QGP.





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**Strangeness enhancement** in high multiplicity

Can this behavior be characterized by other event

- pp and p-Pb are similar to Pb-Pb collisions
- properties other than a difference in multiplicity?
- suppressed strangeness production?

Is it possible to find high-multiplicity events with



# Introduction and motivation



achievable only in heavy-ion collisions, where the data support the formation of QGP.





- **Strangeness enhancement** in high multiplicity pp and p-Pb are similar to Pb-Pb collisions
- Can this behavior be characterized by other event properties other than a difference in multiplicity?
- Is it possible to find high-multiplicity events with suppressed strangeness production?
- Similar features of **baryon-to-meson ratios** in pp, p-Pb and Pb-Pb collisions
- Is the origin the same in small and large collision systems?

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• Measurements at the LHC have revealed that small collision systems exhibit behaviors formerly thought to be achievable only in heavy-ion collisions, where the data support the formation of QGP.

• The origin of the QGP-like behavior in small systems is still unclear. One of the explanations is a multiple parton interactions (MPI) based picture with colour reconnection and ropes, however, MPI can not be

accessible directly in experiments.







MPI: strong interaction between strings overlapping on distance scales of O(1 fm)

# Introduction and motivation

 $N_{\text{mpi}}$  $\rangle$ dyd $p_{\text{T}}$ )

 $\langle mpi, MB \rangle d\mathsf{y} d\rho_T$ 



$$
R_{\rm pp} = \frac{\mathrm{d}^2 N_{\pi}^{\rm mpi} / (\langle I \rangle)}{\mathrm{d}^2 N_{\pi}^{\rm MB} / (\langle N_{\rm r} \rangle)}
$$

**PYTHIA 8** Ratio of yield in MPI-enhanced pp collisions to yield for minimum bias (MB) pp collisions:

20.11.2023 Sushanta Tripathy **7** Phys. Rev. D 102, 076014 (2020)





Up to 40% increase w.r.t. the binary partonparton scaling: "bump" structure in  $p_T = 1-6$ 





# Introduction and motivation

**PYTHIA 8**



GeV/c: The effect is driven by CR

MPI selection does not bias the high- $p_T$  yield

20.11.2023 Sushanta Tripathy **8** A. Ortiz, A. Paz, J. D. Romo, S.Tripathy, E. A. Zepeda, I. Bautista, Phys. Rev. D 102, 076014 (2020)

 $N_{\text{mpi}}\rangle$ dyd $p_{\text{T}}$ )

 $_{\text{mpi, MB}}\rangle$ dyd $p_{\text{T}}$ )

$$
R_{\rm pp} = \frac{\mathrm{d}^2 N_{\pi}^{\rm mpi} / (\langle \mathbf{I} \rangle)^2}{\mathrm{d}^2 N_{\pi}^{\rm MB} / (\langle \mathbf{N}_{\rm r} \rangle)^2}
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Ratio of yield in MPI-enhanced pp collisions to yield for minimum bias (MB) pp collisions:



# Introduction and motivation



$$
R_{\rm pp} = \frac{\mathrm{d}^2 N_{\pi}^{\rm mpi} / (\langle \mathbf{N}_{\rm r} \rangle)^2}{\mathrm{d}^2 N_{\pi}^{\rm MB} / (\langle \mathbf{N}_{\rm r} \rangle)^2}
$$







- and a selection bias is seen in high- $p_T$  yield
- with reduced selection bias

The "bump" structure is not seen in measurements as a function of multiplicity

**PYTHIA 8 DATA** Ratio of yield in MPI-enhanced pp collisions  $\begin{bmatrix} 0 & 0 \\ 0 & 10^6 \end{bmatrix}$ to yield for minimum bias (MB) pp collisions:

> Up to 40% increase w.r.t. the binary partonparton scaling: "bump" structure in  $p_T = 1-6$

> Explore event classifier: sensitivity to MPI

20.11.2023 Phys. Rev. D 102, 076014 (2020) Sushanta Tripathy **ALIUE, EFU U OU (2020) 695** A. Ortiz, A. Paz, J. D. Romo, S.Tripathy, E. A. Zepeda, I. Bautista, Phys. Rev. D 102, 076014 (2020)

 $N_{\text{mpi}}\rangle{dydp_T}$  $_{\text{mpi, MB}}\rangle$ dyd $p_{\text{T}}$ )

GeV/c: The effect is driven by CR

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 $\blacktriangleright$  Measurements at the LHC have revealed that small collision systems exhibit behaviors formerly thought to be



- achievable only in heavy-ion collisions, where the data support the formation of QGP.
- picture with colour reconnection and ropes, however, MPI can not be accessible directly in experiments.
- 
- can also isolate different physics regimes (soft and hard physics).
	- $\textsf{Transverse}\ \textsf{Spherosity}\ (S_{\mathit{O}}^{p_{\mathrm{T}}=1})$ *0*
	- Relative Transverse Activity Classifier (RT)  $\bullet$
	- Charged particle flattenicity (*⍴*ch) (A new classifier -> discussed later in slides)  $\bullet$

The origin of the QGP-like behavior in small systems is still unclear. One of the explanations is an MPI-based

Event selections based only on multiplicity have shown significant bias towards hard pp collisions (selection biases)

Based on MC studies, event topology classifiers have shown a significant reduction of the selection biases and one





# A Large Ion Collider Experiment

### **Inner Tracking System (ITS)**

Tracking, vertex and PID

### **Time Projection Chamber (TPC)**

Tracking and PID (d*E*/dx)

### **Time of Flight (TOF) detector**

PID via time-of-flight method

### **V0**

Trigger, multiplicity/ centrality estimator, event classification based on amplitude

### **Tracking and kinematics**

- ITS and TPC tracks
- $-$  |η|<0.8





# Transverse Spherocity







**Isotropic**: soft-QCD process

Identified particle ratios vs  $S_{0}^{p_T=1}$ *0*









Reduction of ratios relative to pion yields in jet-like events for all particle species -> significant strangeness suppression

Both PYTHIA Monash and Ropes fail to capture the absolute trends but the ratios to  $S_{0}^{p_T=1}$ -integrated events are well explained by the models *0*









- **Proton yield is not modified with spherocity**
- Approximately 20% effect for Ξ
- <sup>§</sup> Strength is ordered in strangeness









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- PYTHIA Monash is unable to capture the trend; also Herwig 7.2





![](_page_16_Picture_3.jpeg)

![](_page_16_Picture_4.jpeg)

ALICE, [arXiv:2310.10236](https://arxiv.org/abs/2310.10236)

![](_page_16_Figure_13.jpeg)

![](_page_16_Figure_14.jpeg)

![](_page_16_Figure_15.jpeg)

![](_page_16_Picture_16.jpeg)

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![](_page_16_Picture_0.jpeg)

![](_page_16_Figure_2.jpeg)

### **Enhanced strangeness production in highmultiplicity collisions seems to be the feature of isotropic events**

![](_page_17_Picture_0.jpeg)

![](_page_17_Figure_2.jpeg)

## $R_T = N_{ch}T / < N_{ch}T >$

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![](_page_17_Picture_14.jpeg)

- Using *R*T, one can vary the magnitude of the underlying event (UE)
- *R*T→0: Events with less UE (dominated by jets)
- Higher *R*<sup>T</sup> → Higher UE contribution
- $\triangle$  A minimum threshold on leading particle  $p_T$  is applied to ensure no bias on spectra vs  $R<sub>T</sub>$  measurements up to the minimum  $p_T$  of the leading particle

![](_page_17_Figure_12.jpeg)

T. Martin, P. Skands, and S. Farrington, Eur. Phys. J. C 76 no. 5, (2016) 299

P. Vargas, Tuesday at 16:00 "Charged-particle production as a function of the relative transverse activity classifier in pp, p-Pb, and Pb-Pb collisions"

![](_page_18_Picture_0.jpeg)

![](_page_18_Figure_1.jpeg)

**ALI-PUB-545303** 

![](_page_18_Figure_12.jpeg)

![](_page_18_Figure_13.jpeg)

![](_page_18_Figure_14.jpeg)

![](_page_18_Figure_15.jpeg)

![](_page_18_Picture_16.jpeg)

ALICE, [JHEP 06 \(2023\) 027](https://doi.org/10.1007/JHEP06(2023)027)

# Identified particle production vs.  $R_T$

- For the transverse region, the ratio of *p*T spectra to the *R*T-integrated spectra rises with increasing  $R_T$ .
- Toward and away regions' high- $p_T$ yields are independent of  $R_T$  (an artefact of the leading  $p_T$  requirement). However, at low- $p_T$ , the  $R_T$ dependence is more evident.
- No "bump" structure seen in these  $measurements \rightarrow selection bias$
- Explore event classifier to have sensitivity to MPI with reduced selection bias

**Motivation**: Search for observable highly sensitive to SOFT particle production (MPI) and CR effects without introducing a bias toward HARD production (multi-jets, high  $p_T$  yield)

![](_page_19_Picture_9.jpeg)

 $\{ {\rm body} \blacktriangleright {\rm soft} \ p{\rm p} \ {\rm collisions} \}$ logy ♦ hard pp collisions

![](_page_19_Figure_12.jpeg)

![](_page_19_Picture_13.jpeg)

![](_page_19_Picture_14.jpeg)

- Define a grid in the η-φ space covered by the V0 detector (10×8 cells)  $\mathcal{Q}$
- The particle multiplicity per cell is measured and flattenicity is calculated  $\mathcal{S}$

$$
\rho_{\rm nch} = \frac{\sqrt{\sum_i (N_{\rm ch}^{\rm cell,i} - \langle N_{\rm ch}^{\rm cell} \rangle)^2 / N_{\rm cell}^2}}{\langle N_{\rm ch}^{\rm cell} \rangle} \qquad \qquad \text{20 isotropic top} \\\text{1 jet-like topol}
$$

### **Charged-particle Flattenicity**

![](_page_19_Picture_0.jpeg)

![](_page_20_Figure_14.jpeg)

![](_page_20_Figure_15.jpeg)

![](_page_20_Figure_16.jpeg)

![](_page_20_Picture_17.jpeg)

A. Ortiz et. al, Rev.Mex.Fis.Suppl. 3 (2022) 4, 040911

To relate the types of events between Spherocity and Flattenicity, a change of variable is performed:  $\rho \rightarrow 1 - \rho$ 

![](_page_20_Picture_3.jpeg)

PYTHIA 8.303 (Monash 2013), pp  $\sqrt{s}$  = 13 TeV,  $N_{\text{mol}}$ =1,  $N_{\text{ch}}$ =235

![](_page_20_Figure_5.jpeg)

 $1 - \rho \rightarrow 1$ Soft pp collision  $1 - \rho \rightarrow 0$ Hard pp collision $\bullet$ æa ð Θ 8

![](_page_20_Picture_0.jpeg)

## Event classification with charged particle flattenicity

Thus, events with large number of MPI are selected when  $1 - \rho \rightarrow 1$ 

PYTHIA 8.303 (Monash 2013), pp  $\sqrt{s}$  = 13 TeV,  $N_{\text{mo}}$  = 24,  $N_{\text{ch}}$  = 325

Selection using flattenicity shows a **"bump" structure Peduced bias towards hard physics** 

![](_page_21_Picture_7.jpeg)

![](_page_21_Figure_2.jpeg)

# Event classification with charged particle flattenicity

A. Ortiz, A. Paz, J. D. Romo, S.Tripathy, E. A. Zepeda, I. Bautista,

![](_page_21_Figure_5.jpeg)

Ratio of yields to MB: **• "Bump"** structure: development of a peak for isotropic events and more evident for

- protons (flattenicity class (I))
- 

![](_page_22_Figure_2.jpeg)

**ALI-PREL-545666**

• Mass dependency: the maximum of the peak shows a mass-dependent ordering

![](_page_22_Figure_12.jpeg)

![](_page_22_Figure_13.jpeg)

![](_page_22_Figure_14.jpeg)

![](_page_22_Picture_15.jpeg)

![](_page_22_Picture_16.jpeg)

$$
Q_{\rm pp} = \frac{\mathrm{d}^2 N^{1-\rho \text{ class}}/(\langle \mathrm{d}N_{\rm ch} / \mathrm{d}\eta \rangle \mathrm{d}y \mathrm{d}p_{\rm T})}{\mathrm{d}^2 N^{\rm MB}/(\langle \mathrm{d}N_{\rm ch} / \mathrm{d}\eta \rangle \mathrm{d}y \mathrm{d}p_{\rm T})}
$$

![](_page_22_Picture_0.jpeg)

# Particle production vs charged particle flattenicity

![](_page_23_Figure_10.jpeg)

• PYTHIA 8 Monash 2013 with MPI and CR effects describes the data; sensitive to event selection due to CR • EPOS LHC describes the data partially (low-to-mid  $p_T$ ); opposite trend seen w.r.t. PYTHIA8 at high  $p_T$ 

# Particle production vs charged particle flattenicity

![](_page_23_Figure_1.jpeg)

**ALI-PREL-545686**

- $Q_{\rm pp} =$  $d^2N^{1-\rho}$  class/( $\langle dN_{ch}/d\eta \rangle$ dyd $p_T$ ) d2*N*MB/(⟨d*N*ch/d*η*⟩d*y*d*p*T)
- 
- 

![](_page_23_Figure_11.jpeg)

![](_page_23_Picture_12.jpeg)

![](_page_24_Figure_14.jpeg)

![](_page_24_Figure_15.jpeg)

![](_page_24_Picture_16.jpeg)

![](_page_24_Picture_0.jpeg)

• Jet-like events produce less strange hadrons than the average high-multiplicity event and the observed strangeness enhancement in high-multiplicity pp collisions is a feature of isotropic

• As suggested by MC studies, selections based on Flattenicity are sensitive to soft particle

- 
- events.
- production and less sensitive to a (jet-) bias.
- Isotropic events develop a bump-like structure with increasing multiplicity similar to the behavior seen as a function of MPI where it is attributed to CR.

![](_page_24_Picture_8.jpeg)

• Along with multiplicity, the event topology classifiers add a new dimension of separating jetlike and isotropic events for pp collisions. They significantly reduce the selection biases.

![](_page_25_Picture_0.jpeg)

**Outlook** 

20.11.2023 Sushanta Tripathy **26**

![](_page_25_Picture_12.jpeg)

### • Flattenicity is defined in the pseudorapidity regions covered by the new V0 and T0C detectors

![](_page_25_Picture_45.jpeg)

- in Run 3 of LHC.
- Stay tuned for new results!

![](_page_25_Picture_4.jpeg)

![](_page_25_Picture_8.jpeg)

![](_page_26_Picture_0.jpeg)

**Outlook** 

20.11.2023 Sushanta Tripathy **27**

![](_page_26_Picture_12.jpeg)

# Thank you for your attention!

![](_page_26_Picture_8.jpeg)

![](_page_26_Picture_4.jpeg)

### • Flattenicity is defined in the pseudorapidity regions covered by the new V0 and T0C detectors

![](_page_26_Picture_46.jpeg)

- in Run 3 of LHC.
- Stay tuned for new results!

![](_page_27_Picture_0.jpeg)

![](_page_27_Picture_3.jpeg)

![](_page_27_Picture_4.jpeg)

# Backup

Integrated yield and mean transverse momentum vs  $S_0^{p_T=1}$ *0*

![](_page_28_Picture_13.jpeg)

![](_page_28_Figure_14.jpeg)

![](_page_28_Figure_15.jpeg)

![](_page_28_Picture_16.jpeg)

![](_page_28_Picture_0.jpeg)

### **PYTHIA 8**

![](_page_28_Figure_3.jpeg)

**ALI-PUB-564123** 

Using mid-rapidity tracklets as an event classifier in conjunction with spherocity in MC shows a large shift in <*p*T> and a small change in <nMPI>.

High-multiplicity midrapidity measurements are biased towards jets -> Captured by jet-like events

<sup>2</sup> Reduced bias in isotropic events

![](_page_28_Picture_8.jpeg)

![](_page_29_Picture_14.jpeg)

Integrated yield and mean transverse momentum vs  $S_0^{p_T=1}$ *0*

![](_page_29_Picture_9.jpeg)

ALICE, [arXiv:2310.10236](https://arxiv.org/abs/2310.10236)

### **DATA**

![](_page_29_Figure_11.jpeg)

- Using mid-rapidity tracklets as an event classifier in conjunction with spherocity in data shows similar behavior as expected from studies as a function of <nMPI> in MC.
- High-multiplicity midrapidity measurements are biased  $\blacklozenge$ towards jets -> Captured by jet-like events
- Reduced bias in isotropic events

![](_page_29_Picture_0.jpeg)

 : Mid-rapidity multiplicity selection *N*|*η*|<0.8 V0M: Forward-rapidity multiplicity selection tracklets

![](_page_30_Picture_15.jpeg)

![](_page_30_Figure_16.jpeg)

![](_page_30_Figure_17.jpeg)

![](_page_30_Picture_18.jpeg)

ALICE, [arXiv:2310.10236](https://arxiv.org/abs/2310.10236)

- $\mathbf{s}_{0}^{\rho_{\tau} = 1}$  Integrated  $+ S_0^{p_{\tau} = 1}$ : 0-10%
- $+$   $S_0^{p_{\tau} = 1}$ : 90-100%
- PYTHIA 8.2 Monash
- $-$  PYTHIA 8.2 Ropes
- Particle production for jet-like events is suppressed at low- $p_T$ but enhanced at high- $p_T$ ; vice versa for isotropic events
- Indicates hardening of the spectra in Jet-like events
- $p_{\tau}$  (GeV/c)  $\approx$  Both PYTHIA Monash and Ropes describe the qualitative trends

![](_page_30_Picture_0.jpeg)

### Identified particle production vs  $S_{0}^{p_T=1}$ *0*

![](_page_30_Figure_2.jpeg)

![](_page_31_Figure_12.jpeg)

![](_page_31_Figure_13.jpeg)

![](_page_31_Figure_14.jpeg)

![](_page_31_Picture_15.jpeg)

Ratio of yields to MB:

$$
Q_{\rm pp} = \frac{\mathrm{d}^2 N^{1-\rho \text{ class}}/(\langle \mathrm{d}N_{\rm ch} / \mathrm{d}\eta \rangle \mathrm{d}y \mathrm{d}p_{\rm T})}{\mathrm{d}^2 N^{\rm MB}/(\langle \mathrm{d}N_{\rm ch} / \mathrm{d}\eta \rangle \mathrm{d}y \mathrm{d}p_{\rm T})}
$$

- 
- 
- 

• "Bump" structure: clear development of a peak for isotropic events (flattenicity class (I), 0–1% 1-ρ) • Mass dependency: the maximum of the peak shows a mass-dependent ordering

![](_page_31_Figure_2.jpeg)

![](_page_31_Picture_0.jpeg)

# Particle production vs charged particle flattenicity

![](_page_32_Picture_0.jpeg)

## Identified particle production vs. R<sub>T</sub>

![](_page_32_Figure_12.jpeg)

![](_page_32_Picture_13.jpeg)

A. Ortiz, A. Paz, J. D. Romo, S.Tripathy, E. A. Zepeda, I. Bautista, Phys. Rev. D 102, 076014 (2020)

### ALICE, [JHEP 06 \(2023\) 027](https://doi.org/10.1007/JHEP06(2023)027)

![](_page_32_Figure_3.jpeg)

**ALI-PUB-545303** 

No "bump" structure seen in these measurements → selection bias Explore event classifier: sensitivity to MPI with reduced selection bias

![](_page_32_Picture_8.jpeg)

![](_page_33_Picture_12.jpeg)

![](_page_33_Picture_13.jpeg)

![](_page_33_Picture_0.jpeg)

![](_page_33_Figure_3.jpeg)

**ALI-PUB-563329**

CERI

Relative Transverse activity classifier,  $R_T = N_{ch}$ Transverse /< $N_{ch}$ Transverse>

![](_page_34_Picture_0.jpeg)

![](_page_34_Figure_3.jpeg)

**ALI-PUB-563329**

The contribution from the jets dominate at low  $R<sub>T</sub>$  and the values are similar for all systems, as one would naively expect for  $R<sub>T</sub>→0$ 

![](_page_34_Picture_9.jpeg)

Relative Transverse activity classifier,  $R_T = N_{ch}$ Transverse /< $N_{ch}$ Transverse>

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![](_page_34_Picture_13.jpeg)

![](_page_34_Picture_14.jpeg)

![](_page_35_Picture_0.jpeg)

![](_page_35_Figure_3.jpeg)

![](_page_35_Picture_12.jpeg)

 $\bullet$  For large  $R_T$ , the  $\langle p_T \rangle$  approaches similar values in all three topological regions for a given system: dominant UE contribution

![](_page_35_Picture_8.jpeg)

Relative Transverse activity classifier,  $R_T = N_{ch}$ Transverse /< $N_{ch}$ Transverse>

![](_page_36_Figure_5.jpeg)

![](_page_36_Picture_0.jpeg)

Hard scattering: perturbative QCD

Soft QCD processes: low transverse momenta → non-perturbative QCD

![](_page_37_Figure_8.jpeg)

![](_page_37_Picture_0.jpeg)

Hard scattering: perturbative QCD

Includes:

Underlying Event (UE)

![](_page_38_Picture_14.jpeg)

Soft QCD processes: low transverse momenta → non-perturbative QCD

Includes:

![](_page_38_Figure_9.jpeg)

- Underlying Event (UE)
	- Multiparton interactions (MPI)

![](_page_38_Picture_0.jpeg)

Hard scattering: perturbative QCD

![](_page_39_Picture_13.jpeg)

Soft QCD processes: low transverse momenta → non-perturbative QCD

Includes:

![](_page_39_Figure_10.jpeg)

- Underlying Event (UE)
	- Multiparton interactions (MPI)
	- Initial- and final-state radiation

![](_page_39_Picture_0.jpeg)

Hard scattering: perturbative QCD

![](_page_40_Picture_14.jpeg)

Soft QCD processes: low transverse momenta → non-perturbative QCD

Includes:

![](_page_40_Figure_11.jpeg)

- Underlying Event (UE)
	- Multiparton interactions (MPI)
	- Initial- and final-state radiation
	- Beam remnants

![](_page_40_Picture_0.jpeg)

Hard scattering: perturbative QCD

![](_page_41_Picture_15.jpeg)

Includes:

![](_page_41_Picture_0.jpeg)

- Hard scattering: perturbative QCD
- Soft QCD processes: low transverse momenta → non-perturbative QCD

- Underlying Event (UE)
	- Multiparton interactions (MPI)
	- Initial- and final-state radiation
	- Beam remnants
- Hadronisation products
- Collective effects

![](_page_41_Figure_13.jpeg)

![](_page_42_Picture_8.jpeg)

![](_page_42_Picture_11.jpeg)

![](_page_42_Picture_0.jpeg)

To reduce the contribution from ISR and FSR, Transverse region is further sub-divided into two regions: Trans-min and Trans-max based on minimum and maximum number of charged particles

 $R_{\text{T,min}} = N_{\text{ch}}^{\text{T,min}} / < N_{\text{ch}}^{\text{T,min}}$  $R_{\text{T,max}} = N_{\text{ch}}^{\text{T,max}} / < N_{\text{ch}}^{\text{T,max}}$ 

G. Bencedi, A. Ortiz, and A. Paz, Phys. Rev. D 104, (2021) 016017