

Angular Correlations of Baryons as a Probe for Colour Reconnection and Hadronization

Stefan Kiebacher in collaboration with Stefan Gieseke and Simon Plätzer | November 21, 2023

Introduction

Hadronization:

- Hadronization is the non-perturbative transition from partons to hadrons and relies on heavy modelling
- Two types of hadronization models: The Lund string model (used by Pythia) and the cluster model (used in Sherpa and Herwig)
- Baryons are not straightforward to produce in a hadronization model and are often badly modelled
⇒ testing ground for Hadronization models

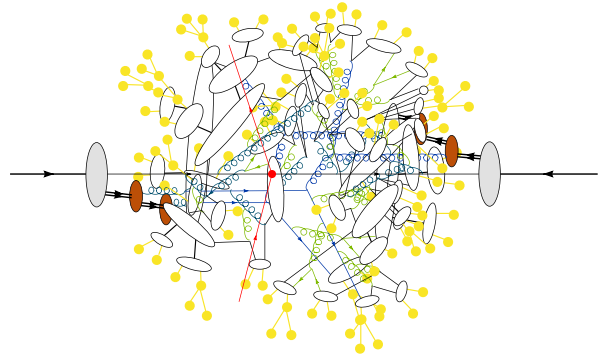
Open questions:

- How and with which kinematics are **Baryons** produced in high energy collisions?
- What knowledge can we extract from **Baryon** observables e.g. their **Angular Correlations**?

Sketch of Hadronization Model in Herwig

Stages of hadronization model in Herwig:

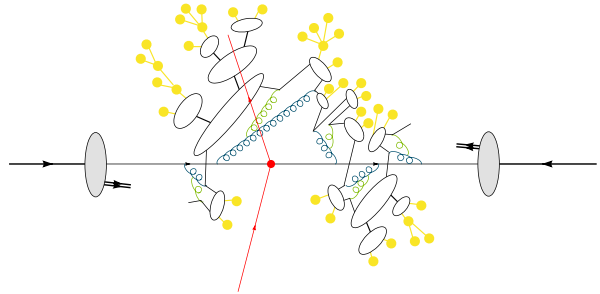
- Multiple Parton Interactions (MPIs)
- Primordial cluster formation
- Colour Reconnection (CR)
- Cluster Fission (CF)
- Cluster Decay (CD)
- Hadron Decay



Sketch of Hadronization Model in Herwig

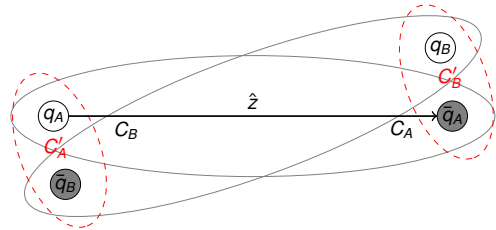
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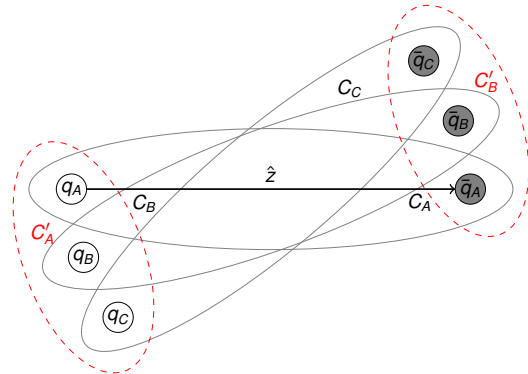
Colour Reconnection (CR)

- Select two clusters at random close in phase space
- Maximize closeness measure $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z}
- If $y_{q_B} > 0$ and $y_{\bar{q}_B} < 0$ for $y_{\text{sum}}^{\text{max}}$
 \Rightarrow Mesonic Colour Reconnection (MCR) accepted with probability P_M



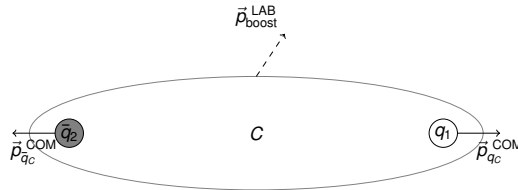
Colour Reconnection (CR)

- Select three clusters at random close in phase space
- Maximize closeness measure $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z}
- If $y_{q_B} < 0$ and $y_{\bar{q}_B} > 0$ for $y_{\text{sum}}^{\text{max}}$
- Find next to maximal y_{sum} cluster to make **baryon-antibaryon pair**
 \Rightarrow Baryonic Colour Reconnection (BCR) accepted with probability P_B [Gieseke, Kirchgaeßer, and Plätzer 2018]
- **Note:** Clusters can be light for Baryon Production However a lot of multiplicity is needed!



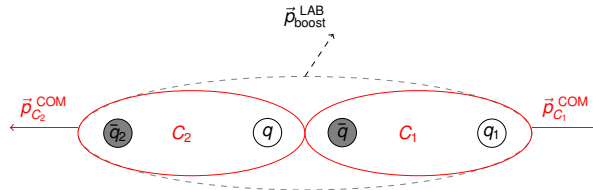
Cluster Fission (CF)

- Fission all clusters $M > M_{\max}(q_1, \bar{q}_2)$ above a threshold $M_{\max}(q_1, \bar{q}_2)$ **recursively**
- 1. Draw a light $q - \bar{q}$ pair from the vacuum with probability P_q (no diquarks currently allowed!)
- 2. Draw new masses M_1, M_2 for the fission products C_1, C_2
- 3. **Choose Direction** of decay \Rightarrow Currently aligned with the original constituent momenta



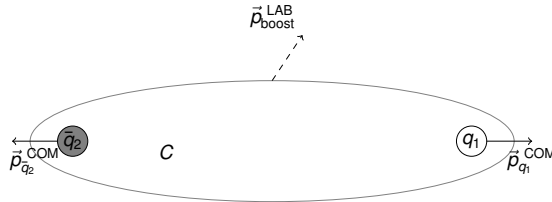
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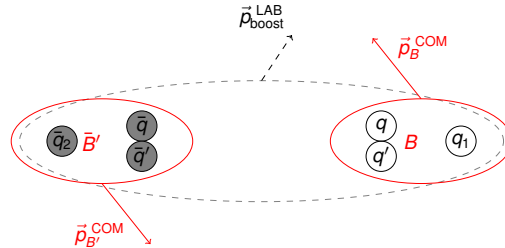
Cluster Decay (CD)

- Clusters decay to two hadrons
- Essentially the same as Cluster fission with a few differences:
 1. The masses are fixed by the hadron masses
 2. Mesonic clusters $C(q_1, \bar{q}_2)$ can decay to a **baryon-antibaryon pair** $B(q_1, (q, q'))$, $\bar{B}'(\bar{q}_2, (\bar{q}, \bar{q}'))$
 3. The direction of decay is chosen **isotropically** in the cluster rest frame
- **Note:** Not much multiplicity, but high-mass clusters are needed to produce **Baryons!**



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Baryon Angular Correlations

- Depletion of near-sided baryons only reproduced by **Baryonic Colour Reconnection (BCR)**
- **Cluster Decay (CD)** baryons are responsible for unphysical far-side peak
- BCR alone cannot produce enough baryons especially for low multiplicity events (e.g. at LEP)

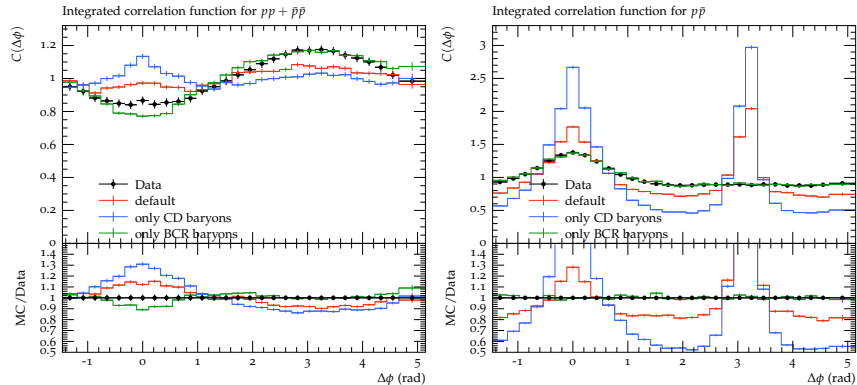


Figure: measured by ALICE [Adam et al. 2017]

Baryon Angular Correlations

Solutions:

1. Disable CD baryon production

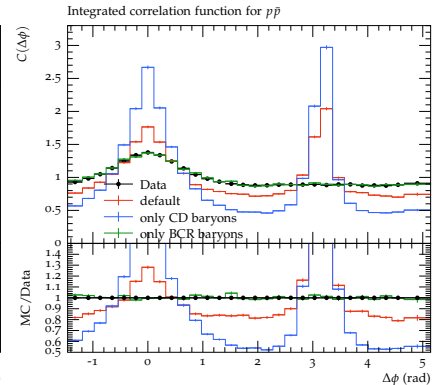
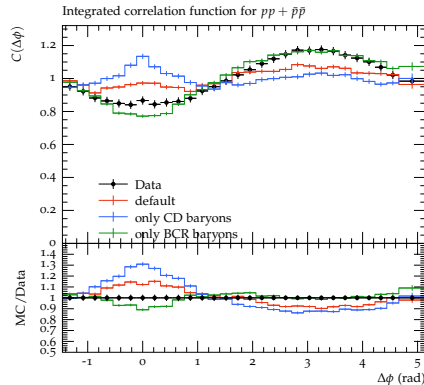


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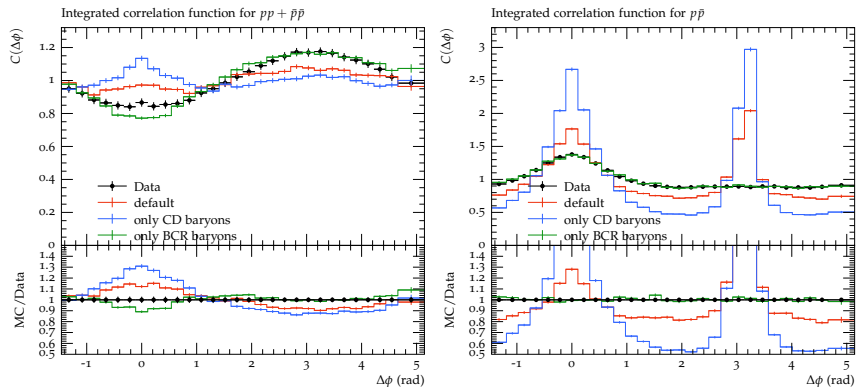


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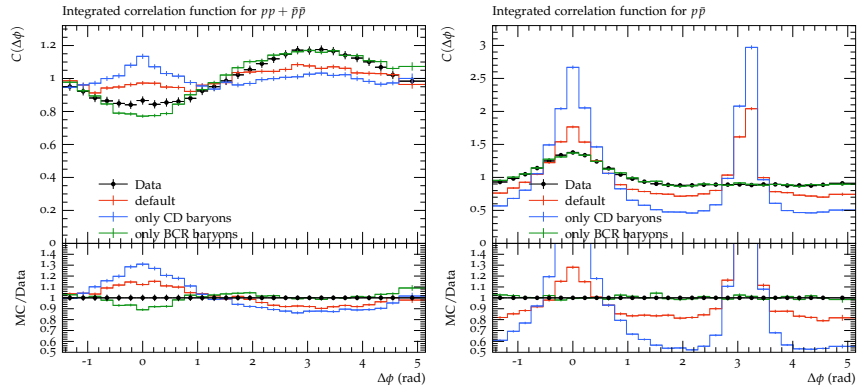


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2. Allow baryon production during Cluster Fission
3. New Diquark Colour Reconnection algorithm

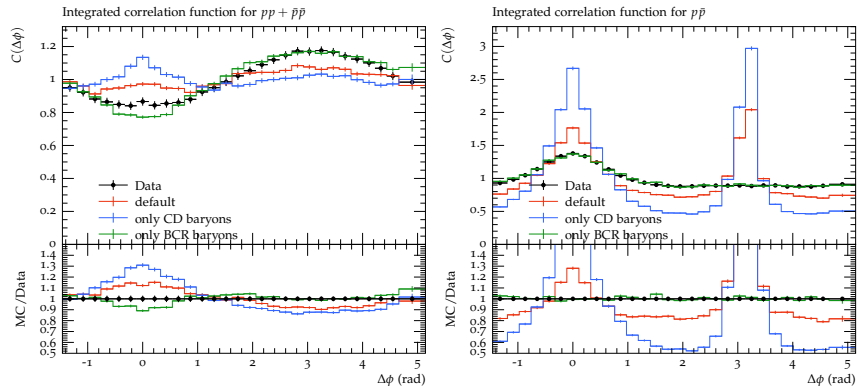


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Cluster Fission (CF) vs Cluster Decay (CD) Baryons

- Near-side depletion not reproduced
 \Rightarrow CD and CF are oblivious to other baryons
- Far-side peak is completely gone!
- Near-side still overshoots the data

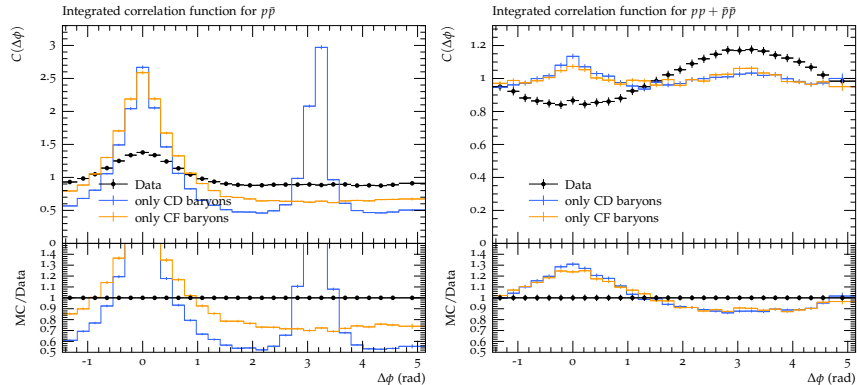
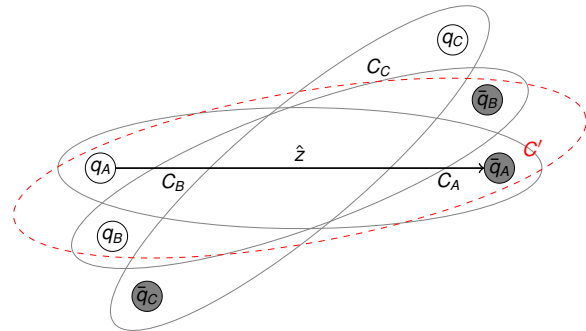


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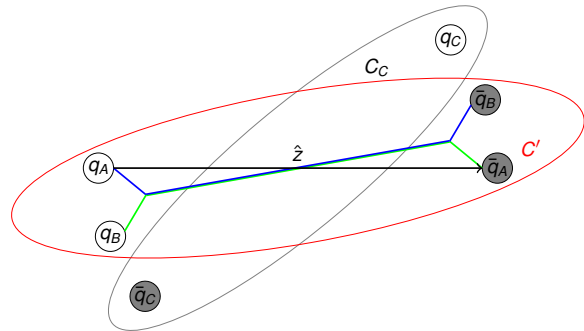
Diquark Colour Reconnection Algorithm

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- **Note:** Mixed need for multiplicity and existing mass for producing baryons
- Similar to Pythia's String Junction Colour Reconnections [Christiansen and Skands 2015]



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Diquark Colour Reconnection

- Purely Diquark-type CR with $P_D = 1$ has not enough depletion for pp correlations
- Near-sided peak reproduced!
- No far-sided peak for $p\bar{p}$
DCR and good phenomenology

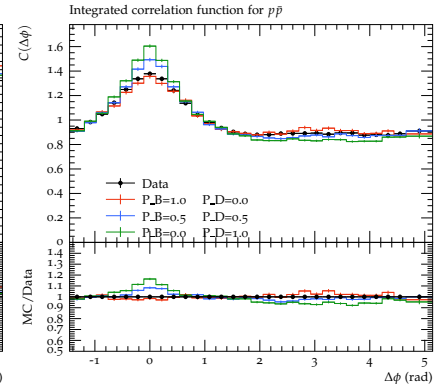
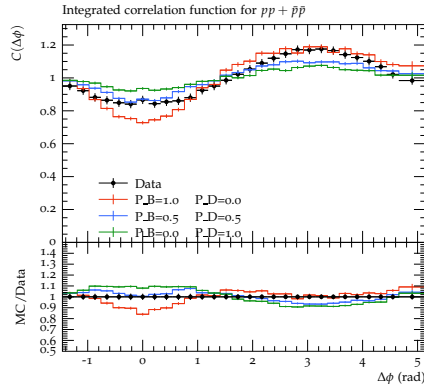


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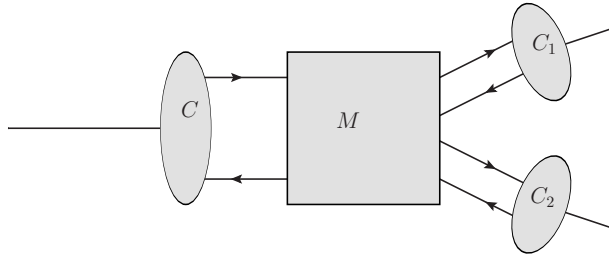
Revisiting Cluster Fission Kinematics

Idea: Cluster Fission is a partonic $2 \rightarrow 4$ process [Plätzer 2023].

- Factorize the process $C(p_i, p_j) \rightarrow C_1(q_i, q), C_2(q_j, \bar{q})$ (see Jan Priedigkeit's Bachelor thesis Graz):

$$d\Gamma(C \rightarrow C_1, C_2) = \int d^4\Phi_{q_i} d^4\Phi_q d^4\Phi_{\bar{q}} d^4\Phi_{q_j} (2\pi)^4 \delta^4(p_i + p_j - q_i - q - \bar{q} - q_j) |\mathcal{M}(p_i, p_j \rightarrow q_i, q, \bar{q}, q_j)|^2 \quad (1)$$

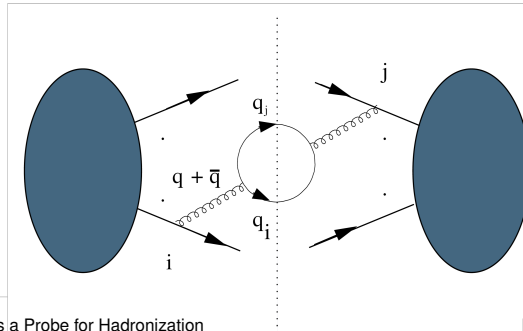
$$d\Gamma(C \rightarrow C_1, C_2) = \int dM_1 dM_2 d\Phi_2(P|Q_1, Q_2) d\Phi_2(Q_1|q_i, q) d\Phi_2(Q_2|q_j, \bar{q}) |\mathcal{M}(p_i, p_j \rightarrow q_i, q, \bar{q}, q_j)|^2 \quad (2)$$



Revisiting Cluster Fission Kinematics

- (Pre-)Sample Masses M_1, M_2 from flat Phase Space weight (Jan Priedigkeit's Bachelor thesis Graz)
 $d\Phi_4 \propto dM_1 dM_2 \sqrt{\lambda(M, M_1, M_2)} \sqrt{\lambda(M_1, m_1, m)} \sqrt{\lambda(M_2, m_2, m)} / (M_1 M_2)^2$
- Rejection sampling of soft $q\bar{q}$ emission diagram, which in the soft limit is given in [Catani and Grazzini 2000] by:

$$|\mathcal{M}(p_i, p_j \rightarrow q_i, q, \bar{q}, q_j)|^2 \propto \frac{2(q_i \cdot q_j)(q \cdot \bar{q}) + [q_i \cdot (q - \bar{q})][q_j \cdot (q - \bar{q})]}{2(q \cdot \bar{q})^2 [q_i \cdot (q + \bar{q})][q_j \cdot (q + \bar{q})]} \quad (3)$$



Summary and Outlook

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- Used the angular correlations of baryons to examine the kinematics of the cluster model
- Found and fixed the far-sided peak of baryon-antibaryon correlations
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Outlook for hadronization in Herwig:

- Dynamic gluon constituent masses (WIP by Daniel Samitz, S. Plätzer)
- Restructure the CF and CD to implement flexible kinematics (WIP with S. Plätzer, S. Gieseke)
- Make Colour Reconnection **dynamic** via soft gluon evolution [Gieseke, Kirchgaeßer, Plätzer, and Siodmok 2018; Plätzer 2023] (WIP with S. Plätzer, S. Gieseke) \Rightarrow only 2 free parameters

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Long term goals for hadronization in Herwig:






- Make shower+hadronization model less sensitive on the shower cutoff
- Generalisation of Hadronization to dark hadrons (by Simon Plätzer, Dominic Stafford et al.)
- Convince experimentalists to get more (identified) particle correlation data (also for LEP), since important modelling input

Short summary






TLDL: Lots of construction sites in the Hadronization model in Herwig ...



References I

-  Adam, Jaroslav et al. (2015). “Measurement of pion, kaon and proton production in proton–proton collisions at $\sqrt{s} = 7$ TeV”. In: Eur. Phys. J. C 75.5, p. 226. DOI: 10.1140/epjc/s10052-015-3422-9. arXiv: 1504.00024 [nucl-ex].
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-  Catani, Stefano and Massimiliano Grazzini (2000). “Infrared factorization of tree level QCD amplitudes at the next-to-next-to-leading order and beyond”. In: Nucl. Phys. B 570, pp. 287–325. DOI: 10.1016/S0550-3213(99)00778-6. arXiv: hep-ph/9908523.
-  Christiansen, Jesper R. and Peter Z. Skands (2015). “String Formation Beyond Leading Colour”. In: JHEP 08, p. 003. DOI: 10.1007/JHEP08(2015)003. arXiv: 1505.01681 [hep-ph].

References II

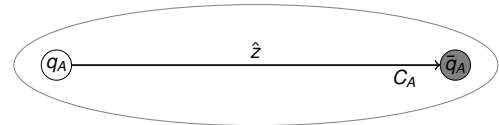
-  Gieseke, Stefan, Patrick Kirchga ber, and Simon Pl tzer (2018). “Baryon production from cluster hadronisation”. In: Eur. Phys. J. C 78.2, p. 99. DOI: 10.1140/epjc/s10052-018-5585-7. arXiv: 1710.10906 [hep-ph].
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-  Khachatryan, Vardan et al. (2011). “Strange Particle Production in pp Collisions at $\sqrt{s} = 0.9$ and 7 TeV”. In: JHEP 05, p. 064. DOI: 10.1007/JHEP05(2011)064. arXiv: 1102.4282 [hep-ex].
-  Pl tzer, Simon (2023). “Colour evolution and infrared physics”. In: JHEP 07, p. 126. DOI: 10.1007/JHEP07(2023)126. arXiv: 2204.06956 [hep-ph].
-  Seidl, R. et al. (2017). “Invariant-mass and fractional-energy dependence of inclusive production of di-hadrons in e^+e^- annihilation at $\sqrt{s} = 10.58$ GeV”. In: Phys. Rev. D 96.3, p. 032005. DOI: 10.1103/PhysRevD.96.032005. arXiv: 1706.08348 [hep-ex].

Thank You For Your Attention!

Questions? Remarks? Comments?

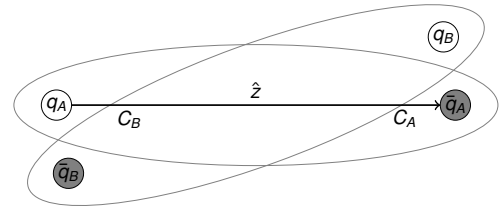
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- Boost in cluster rest frame



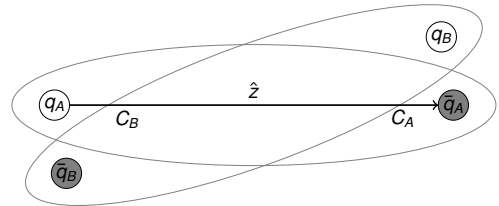
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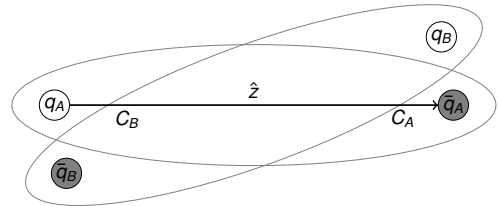
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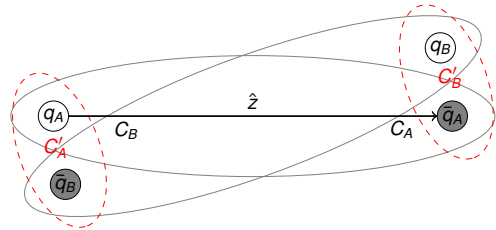
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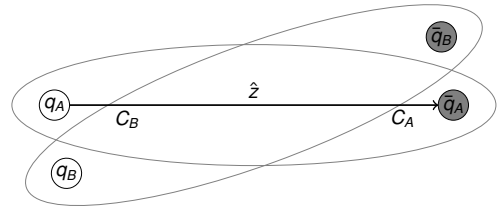
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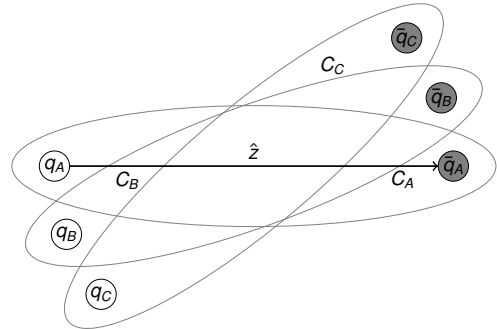
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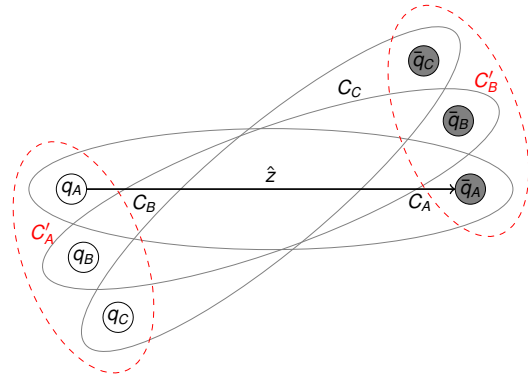
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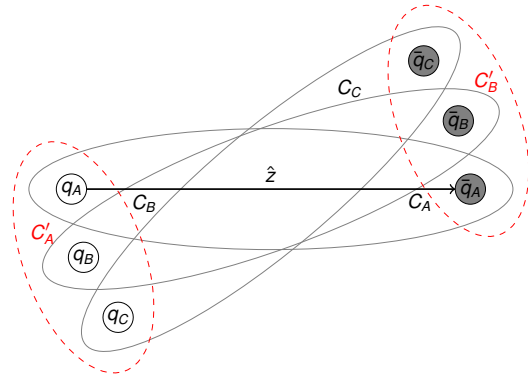
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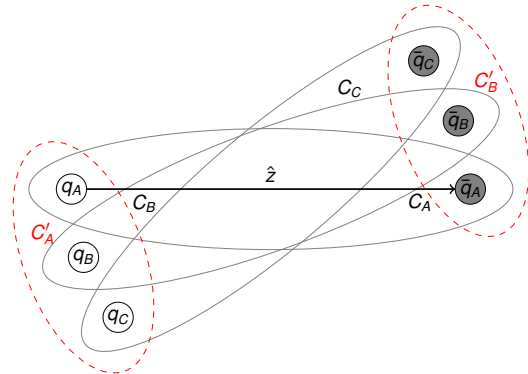
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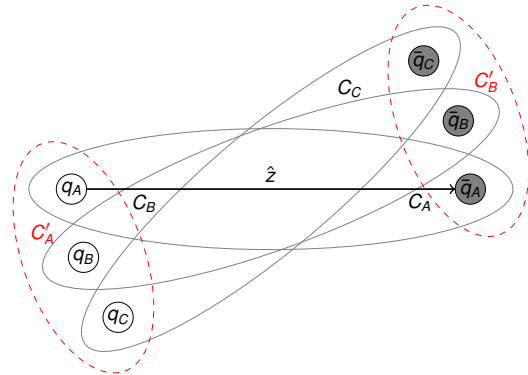
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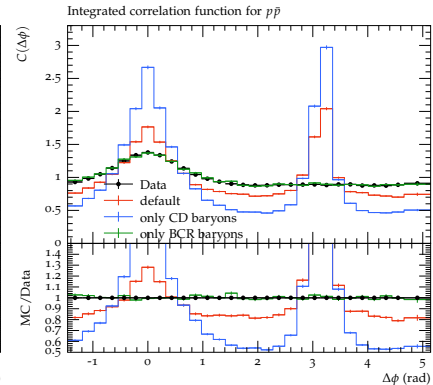
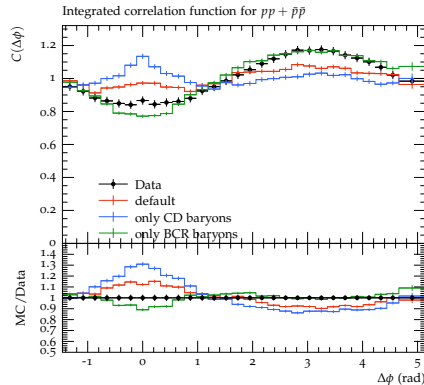
Backup: Colour Reconnection (CR)

- Boost in cluster rest frame
- Select next cluster at random
- Compute $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z}
 \Rightarrow find maximal y_{sum}
- If $y_{q_B} < 0$ and $y_{\bar{q}_B} > 0$ for $y_{\text{sum}}^{\text{max}}$
- Find next to maximal y_{sum} cluster to make **baryon-antibaryon pair**
 \Rightarrow Baryonic Colour Reconnection (BCR) accepted with probability P_B
- **Note:** Clusters can be light for Baryon Production
 However a lot of multiplicity is needed!
- In fact BCR regulated the over-abundance of high multiplicity events [Gieseke, Kirchga efer, and Pl atzer 2018]



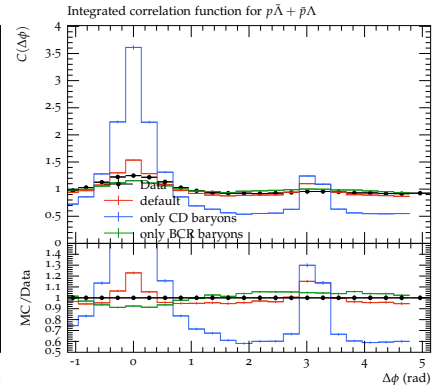
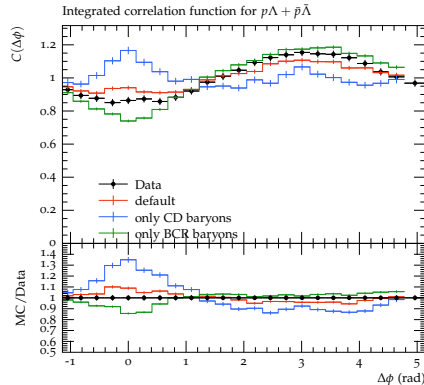
Backup: Baryon Angular Correlations

- Depletion of near-sided baryons only reproduced by **Baryonic Colour Reconnection (BCR)**
- **Cluster Decay (CD)** baryons are giving opposite features to data
- Cluster Decay baryons are solely responsible for unphysical far-side peak
- BCR alone cannot produce enough baryons especially for low multiplicity events (e.g. at LEP)



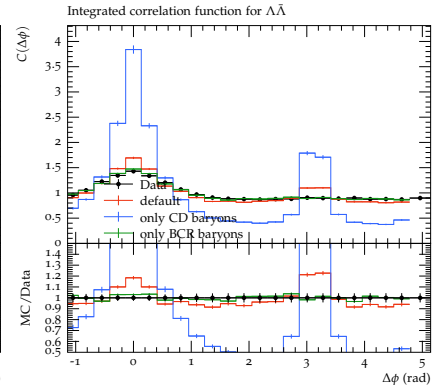
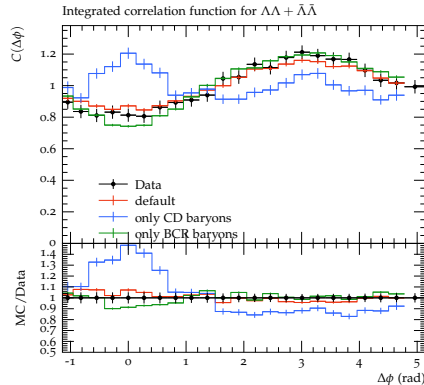
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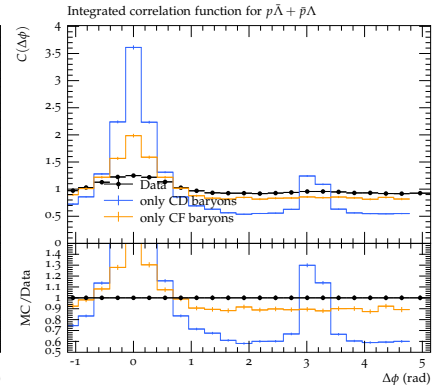
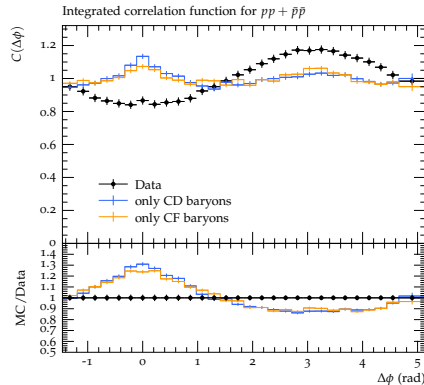
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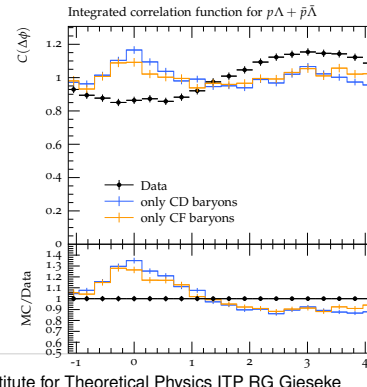
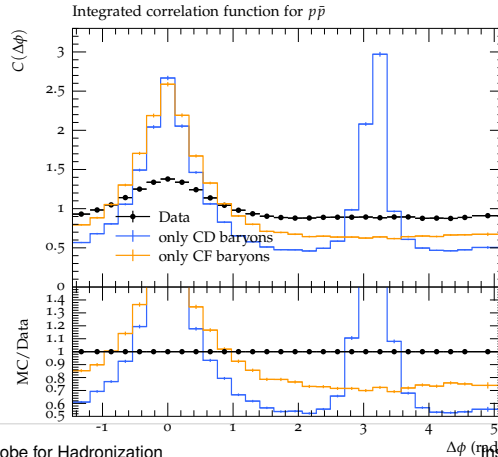
Backup: Cluster Fission vs Cluster Decay Baryons

- CD baryon mechanism vs new Cluster Fission (CF) mechanism
 - Near-side depletion not reproduced
- ⇒ CD and CF are oblivious to other baryons



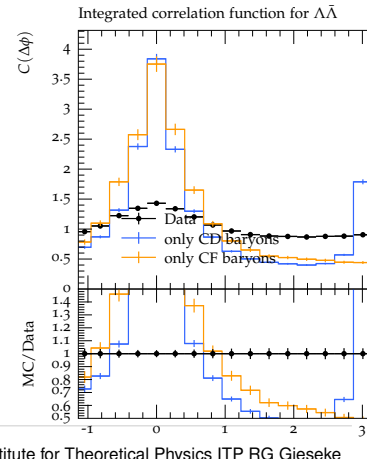
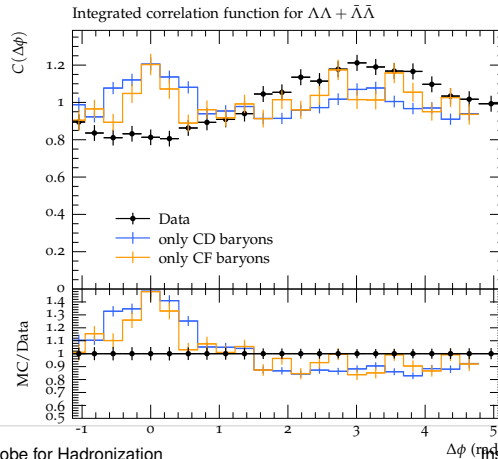
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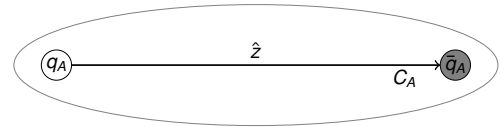
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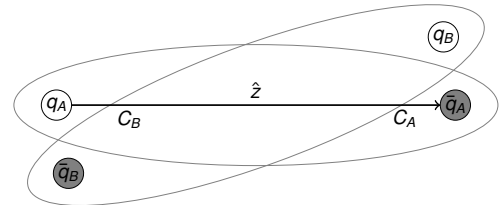
Backup: Diquark Colour Reconnection Algorithm

- Boost in cluster rest frame



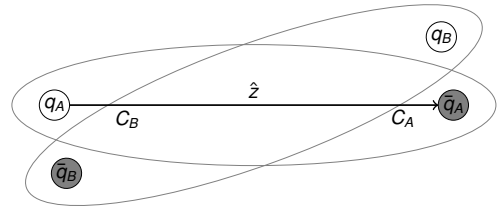
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- Boost in cluster rest frame
- Select next cluster at random



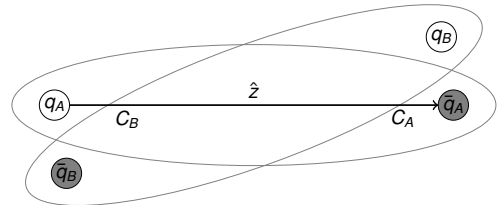
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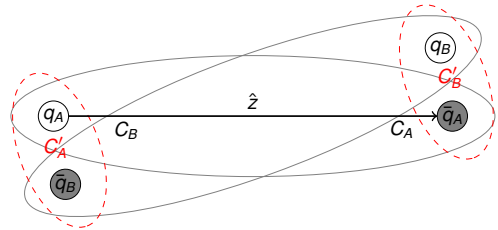
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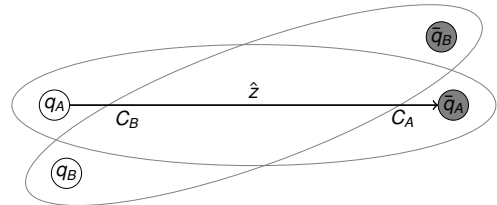
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 ⇒ find maximal y_{sum}
- If $y_{q_B} > 0$ and $y_{\bar{q}_B} < 0$ for $y_{\text{sum}}^{\text{max}}$
 ⇒ Mesonic Colour Reconnection
 accepted with probability P_M



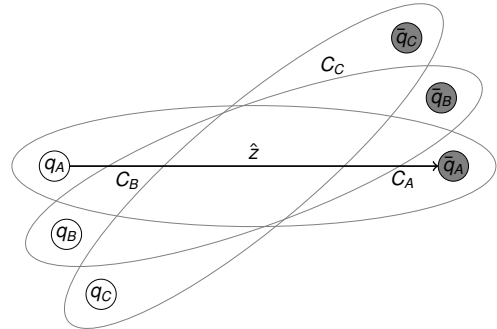
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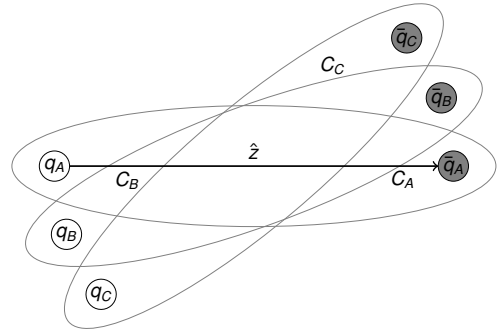
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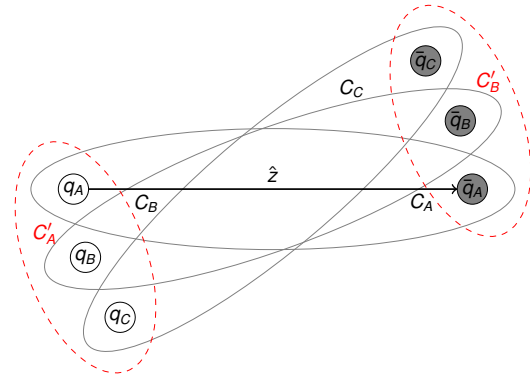
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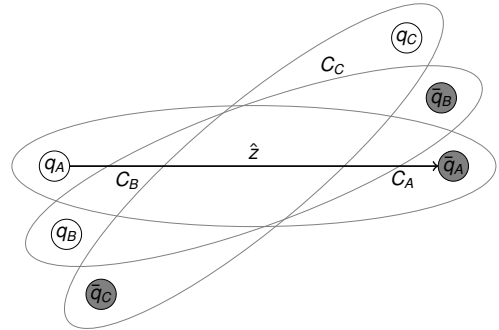
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- If $y_{q_B} < 0$ and $y_{\bar{q}_B} > 0$ for $y_{\text{sum}}^{\text{max},2}$
 \Rightarrow Baryonic Colour Reconnection (BCR)
 accepted with probability P_B



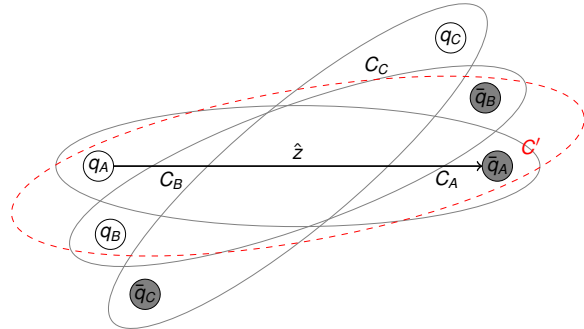
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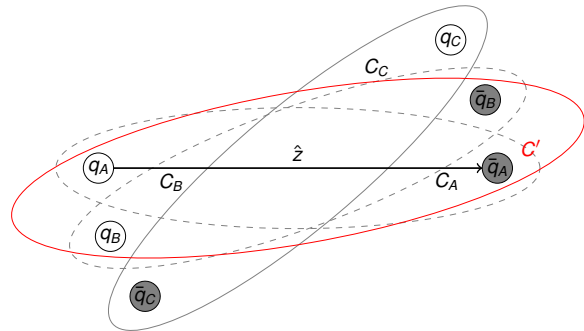
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 \Rightarrow Diquark Colour Reconnection (DCR)
 accepted with probability P_D



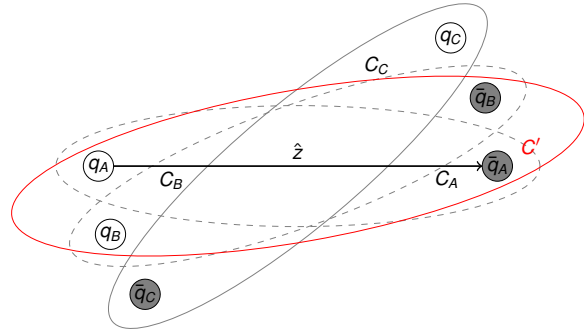
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 accepted with probability P_D if $M_{C'}$ > $M_{\text{Lightest Baryon Pair}}$



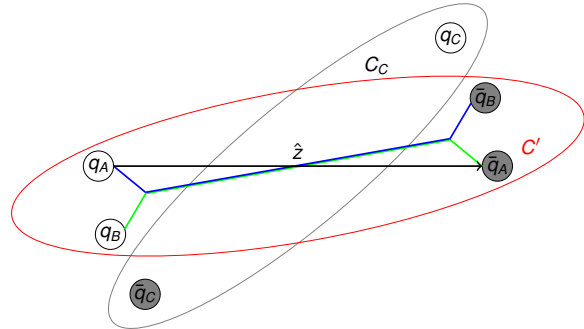
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- **Note:** Mixed need for multiplicity and existing mass for producing baryons



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- Similar to Pythia's String Junction Colour Reconnections [Christiansen and Skands 2015]



Preliminary: New Cluster Fission Kinematics applied

- BELLE data for di-hadron mass spectrum improve for large z with Preliminary results

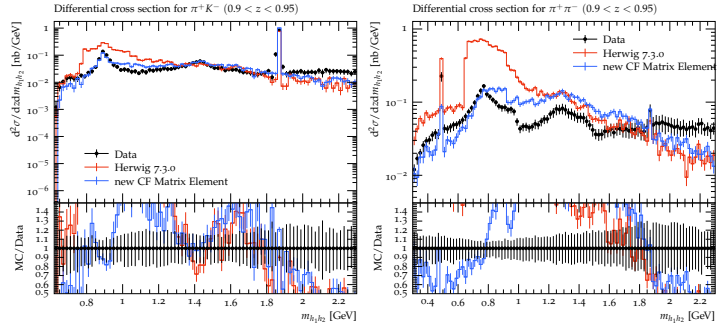


Figure: Thrust $T > 0.8$, $\vec{p}_1 \cdot \vec{p}_2 > 0$ and $z = \frac{2(E_1 + E_2)}{\sqrt{s}}$ [Seidl et al. 2017]

Preliminary: New Cluster Fission Kinematics applied

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- However still very inefficient rejection sampling for large mass clusters (LEP manageable; LHC way too slow)

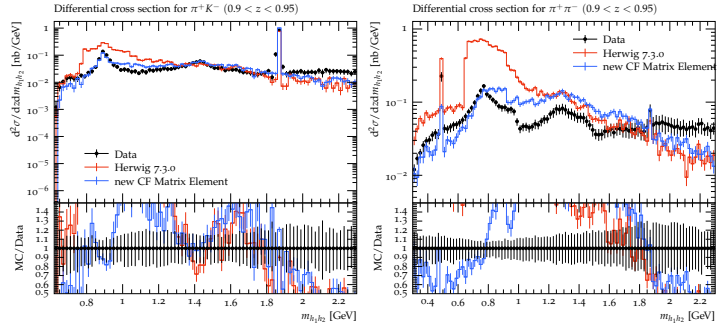


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Backup: Spectra of Protons

- Proton p_T -spectra are badly modelled

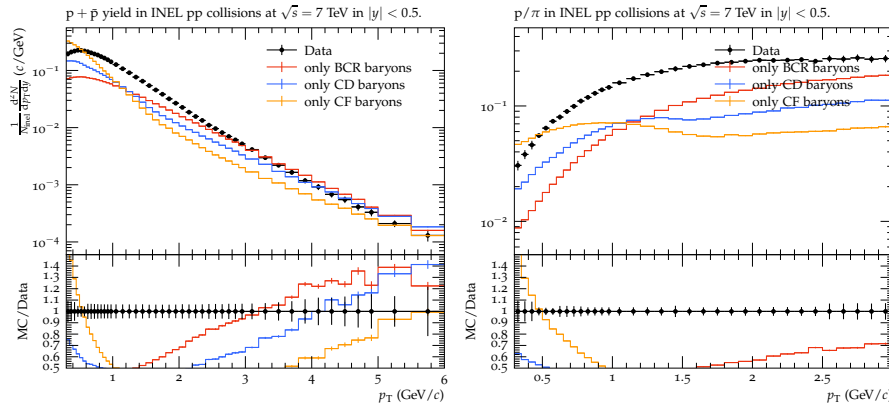


Figure: Compare p_T - spectra of p for only BCR, only CD or only CF baryon mechanisms [Adam et al. 2015]

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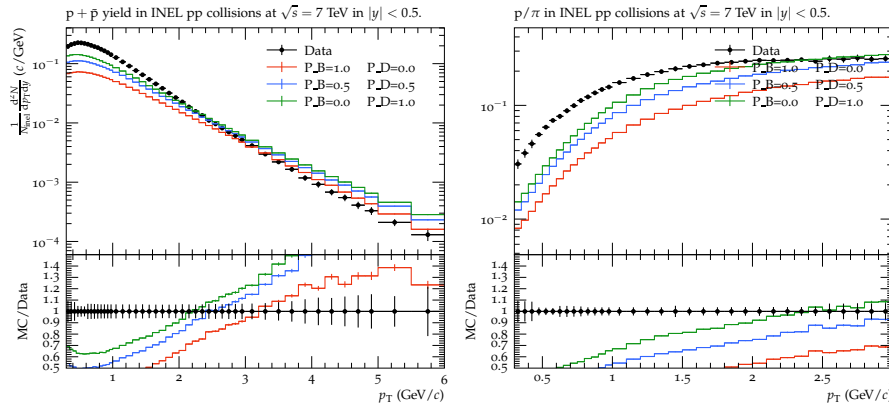


Figure: Compare p_T — spectra of p for only new DCR baryon mechanism with different probabilities [Adam et al. 2015]

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Possible solutions: Use of a different "Loss function" than regular χ^2 e.g. $\chi^2 \rightarrow \frac{\chi^2}{1+\chi^2}$ or $\tanh(\chi^2)$

Backup: Cluster Fission Details

- A cluster of mass M is fissioned if $M^{\text{Cl}_{\text{pow}}} \geq \text{Cl}_{\text{max}}^{\text{Cl}_{\text{pow}}} + (m_1 + m_2)^{\text{Cl}_{\text{pow}}}$, where m_1, m_2 are the masses of the constituents of the cluster
- Currently masses are sampled as follows, where $r_1, r_2 \in [0, 1]$ are uniform random numbers [Bahr et al. 2008]:

$$M_1 = m_1 + (M - m_1 - m_2) r_1^{\frac{1}{P_{\text{split}}}} \quad (4)$$

$$M_2 = m_2 + (M - m_1 - m_2) r_2^{\frac{1}{P_{\text{split}}}} \quad (5)$$

- Reject samples if $M_1 + M_2 > M$
- Problems: huge dependence on parameters Cl_{max} and especially P_{split}
- Work in progress: Sample masses according to phase space

Backup: Angular Correlations

- The shown plots are showing correlations integrated in $\Delta\eta$ up to $\Delta\eta_{\max} = 1.3$
- The angular correlations are measured via the event mixing [Adam et al. 2017]:

$$C_i(\Delta\phi, \Delta\eta) = \frac{S(\Delta\phi, \Delta\eta)}{B(\Delta\phi, \Delta\eta)} \quad (6)$$

$$S_i(\Delta\phi, \Delta\eta) = \frac{1}{N_{\text{pairs}}^{\text{same}}} \frac{d^2 N_{\text{pairs}}^{\text{same}}}{d\Delta\eta d\Delta\phi} \quad (7)$$

$$B_i(\Delta\phi, \Delta\eta) = \frac{1}{N_{\text{pairs}}^{\text{mixed}}} \frac{d^2 N_{\text{pairs}}^{\text{mixed}}}{d\Delta\eta d\Delta\phi} \quad (8)$$

$$C_i(\Delta\phi) = \int_0^{\Delta\eta_{\max}} C_i(\Delta\phi, \Delta\eta) d\Delta\eta \quad (9)$$

Spectra of Protons

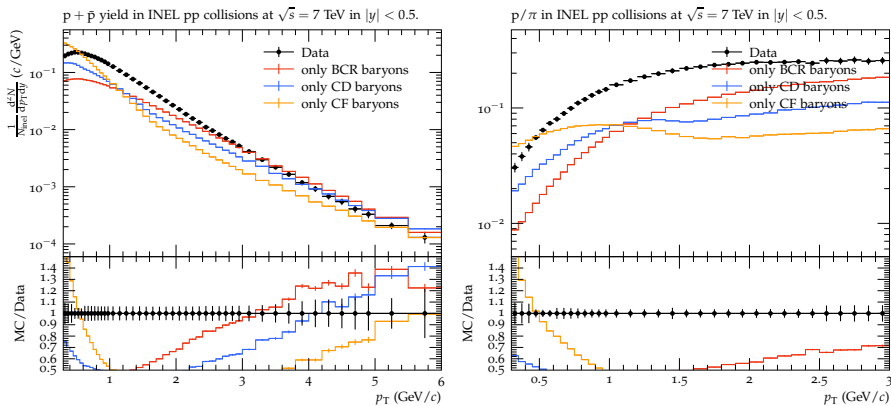


Figure: Compare p_T — spectra of p for only BCR, only CD or only CF baryon mechanisms [Adam et al. 2015]

Spectra of Protons

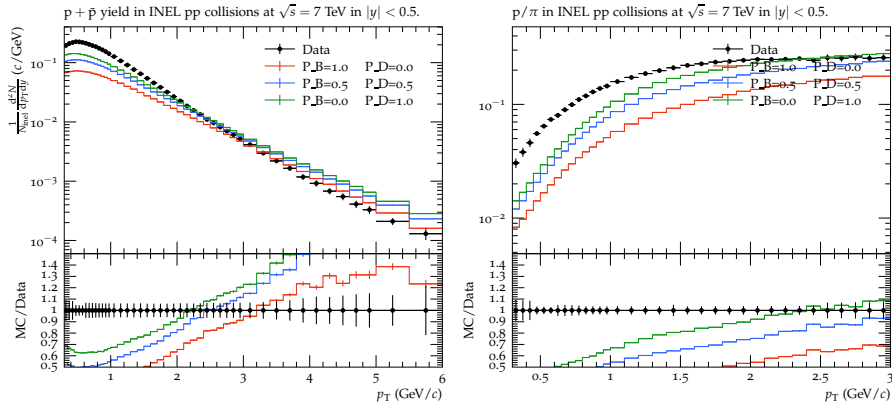


Figure: Compare p_T — spectra of p for only new DCR baryon mechanism with different probabilities [Adam et al. 2015]

Spectra of Strange Baryons

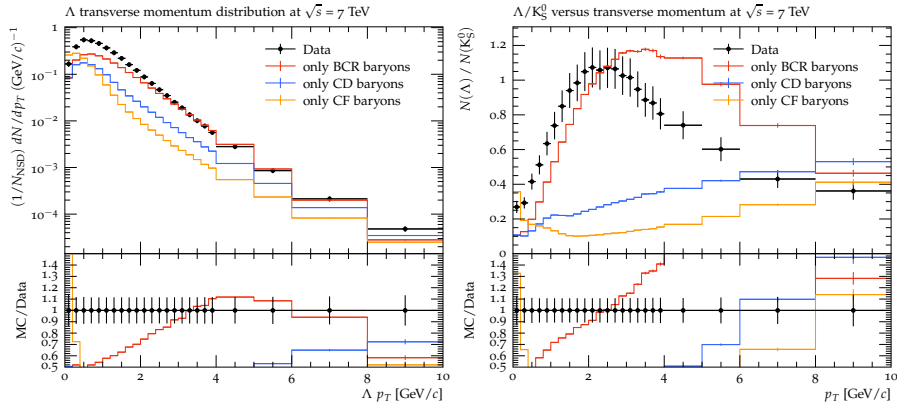


Figure: Compare p_T – spectra of Λ, Ξ for only BCR, only CD or only CF baryon mechanisms [Khachatryan et al. 2011]

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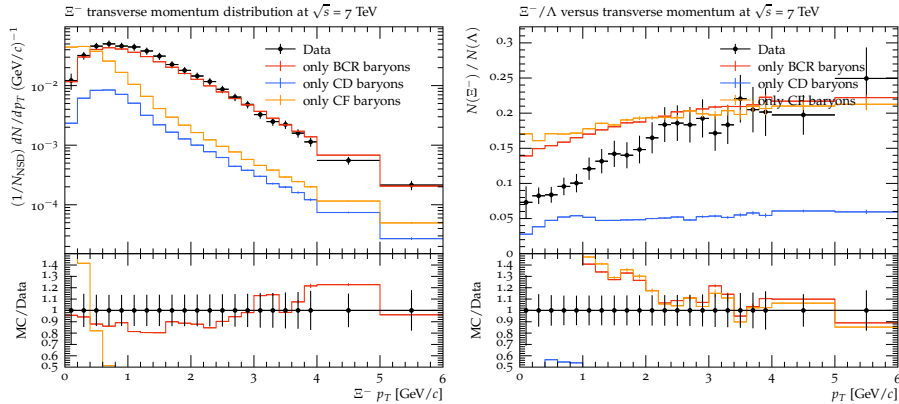


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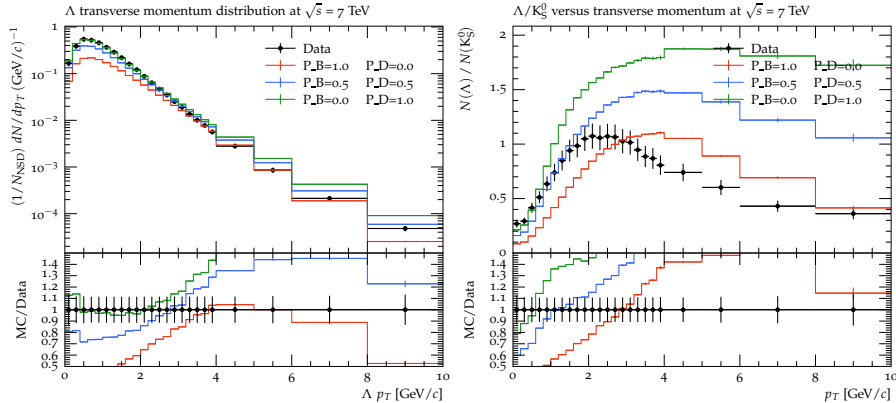


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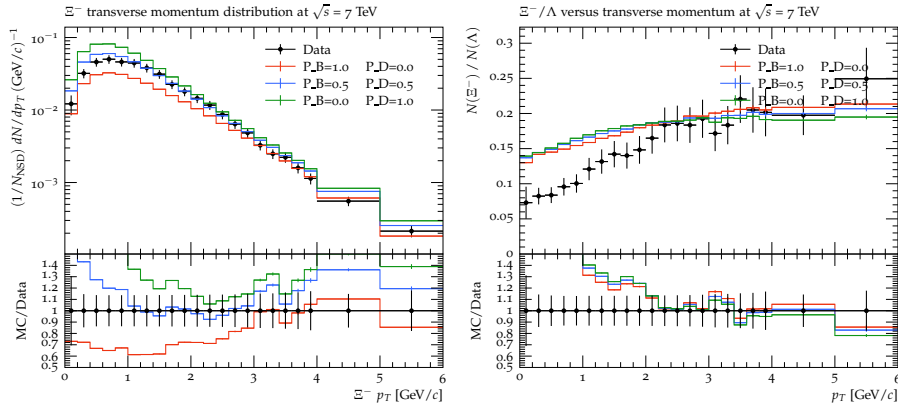


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Consistent Two Particle Boost

- If we boost a two particle system $P = (p_i + p_j)$ into its rest frame $\hat{P} = (\hat{p}_i + \hat{p}_j)$ one needs to be careful to transform the relative momentum correctly $\hat{P}_{\text{rel}} = (\hat{p}_i - \hat{p}_j)$
- The naive transformation would be to just use $\Lambda_{(-P)}$, but this would give in general $\hat{P}_{\text{rel}} = (\hat{p}_i - \hat{p}_j + 2k)$, because $\Lambda\hat{p}_i = p_i + k$ and $\Lambda\hat{p}_j = p_j - k$.
- Intuitively the momentum P is completely oblivious to its components and therefore Λ must depend on both the constituents p_i, p_j
- Want a Lorentz Transformation (matrix or tensor) $\Lambda(p_i, p_j | \hat{p}_i, \hat{p}_j)$ such that $\Lambda\hat{p}_i = p_i$ and $\Lambda\hat{p}_j = p_j$
- Found solution for $\Lambda(p_i, p_j | \hat{p}_i, \hat{p}_j)$, but numerically not very easy
- Work in Progress: Tensor for this trafo Λ^ν_{μ}