

New foundational experiments with quantum process tomography

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Michał Eckstein^{1,2} & Paweł Horodecki^{2,3}

¹ Institute of Theoretical Physics, Jagiellonian University, Kraków, Poland

² International Center for Theory of Quantum Technologies, University of Gdańsk

³ Gdańsk University of Technology, Poland



JAGIELLONIAN UNIVERSITY
IN KRAKÓW



University
of Gdańsk



GDAŃSK UNIVERSITY
OF TECHNOLOGY

West Lafayette, IN, Oct 2, 2023

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

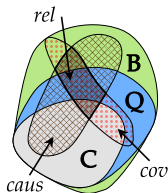
- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

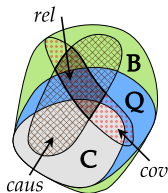


Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

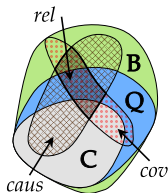


Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

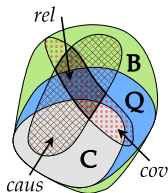


Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

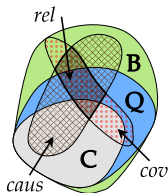


Routes towards New Physics

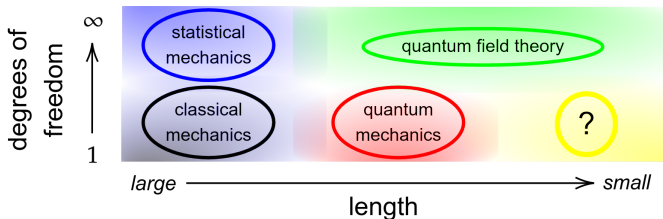
$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- 1 Beyond Standard Model, but still in QFT
 - SUSY, composite Higgs, dark sector, inflation, ...
- 2 Beyond Special Relativity, but assuming QM
 - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
 - quantum gravity
- 3 Beyond Quantum Mechanics, but assuming relativity
 - Super-quantum correlations, ...
 - Deviations from linearity in QM and/or QFT
 - Objective wave function collapse models

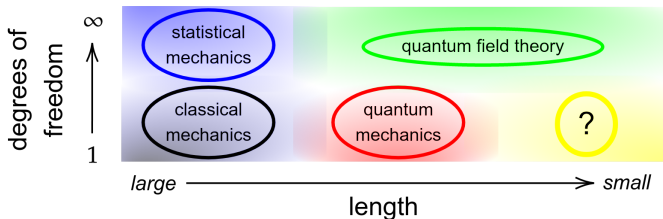


Beyond-quantum physics?



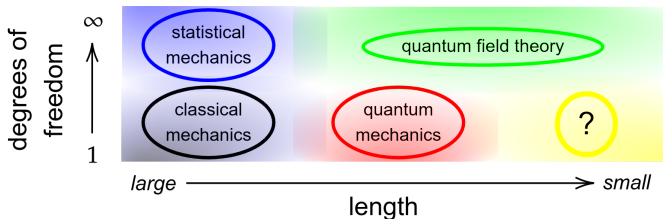
- Is there an 'objective collapse' in a decay process?
- Are correlations in QFT stronger than in QM?
- Are QM & QFT only effective descriptions of Nature?
- How to look for possible deviations from QM?

Beyond-quantum physics?



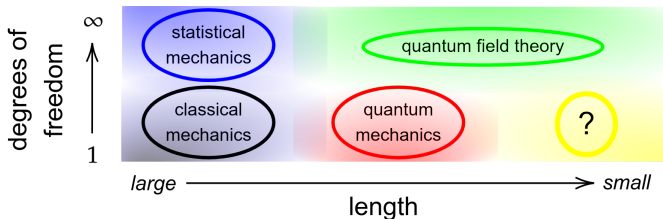
- Is there an 'objective collapse' in a decay process?
- Are correlations in QFT stronger than in QM?
- Are QM & QFT only effective descriptions of Nature?
- How to look for possible deviations from QM?

Beyond-quantum physics?



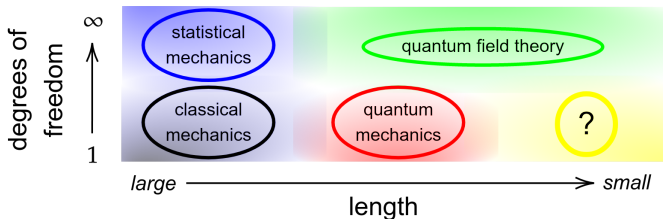
- Is there an 'objective collapse' in a decay process?
- Are correlations in QFT stronger than in QM?
- Are QM & QFT only effective descriptions of Nature?
- How to look for possible deviations from QM?

Beyond-quantum physics?



- Is there an 'objective collapse' in a decay process?
- Are correlations in QFT stronger than in QM?
- Are QM & QFT only effective descriptions of Nature?
- How to look for possible deviations from QM?

Beyond-quantum physics?



- Is there an 'objective collapse' in a decay process?
- Are correlations in QFT stronger than in QM?
- Are QM & QFT only effective descriptions of Nature?
- How to look for possible deviations from QM?

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)

The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

- No pre-correlations between the inputs (x, y) and the box (λ) .

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)

The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

- No pre-correlations between the inputs (x, y) and the box (λ) .

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)

The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

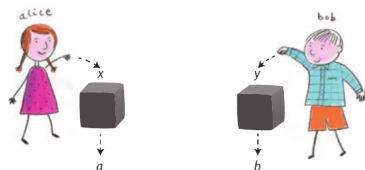
- No pre-correlations between the inputs (x, y) and the box (λ) .

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)



The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

[Sandu Popescu, *Nature Physics* 10, 264 (2014)]

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

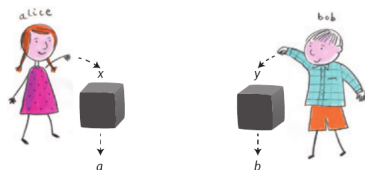
- No pre-correlations between the inputs (x, y) and the box (λ).

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)



The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

[Sandu Popescu, *Nature Physics* 10, 264 (2014)]

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

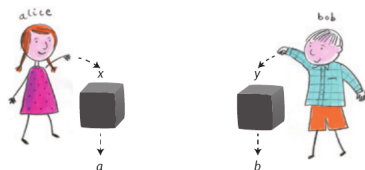
- No pre-correlations between the inputs (x, y) and the box (λ).

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)



The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

[Sandu Popescu, *Nature Physics* 10, 264 (2014)]

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

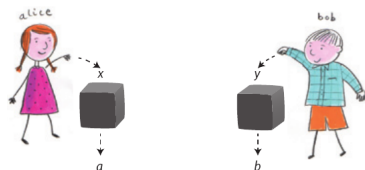
- No pre-correlations between the inputs (x, y) and the box (λ).

The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} | \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs (x, y) — 2 outputs (a, b)



The *experimental* (frequency) correlation function:

$$C_e(x, y) = P(a = b | x, y) - P(a \neq b | x, y)$$

[Sandu Popescu, *Nature Physics* 10, 264 (2014)]

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y | \lambda) = P(x) \cdot P(y)$$

- No pre-correlations between the inputs (x, y) and the box (λ).

Quantum-data boxes

- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{in}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{out} \rightarrow a$.

Quantum-data boxes

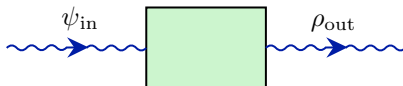
- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{in}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{out} \rightarrow a$.

Quantum-data boxes

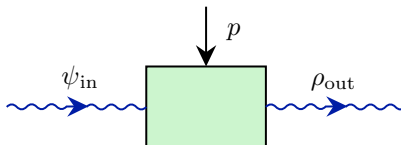
- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{\text{in}}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{\text{out}} \rightarrow a$.

Quantum-data boxes

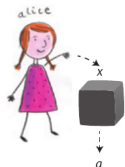
- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



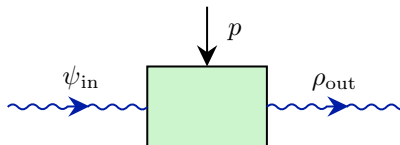
- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{in}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{out} \rightarrow a$.

Quantum-data boxes

- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



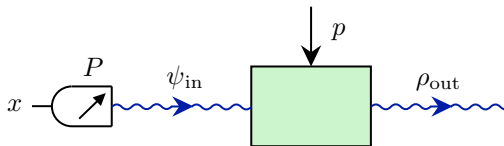
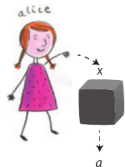
[Nat. Phys. 10, 264 (2014)]



- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{\text{in}}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{\text{out}} \rightarrow a$.

Quantum-data boxes

- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.

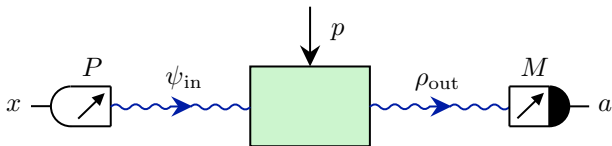
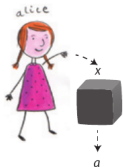


[Nat. Phys. 10, 264 (2014)]

- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{\text{in}}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{\text{out}} \rightarrow a$.

Quantum-data boxes

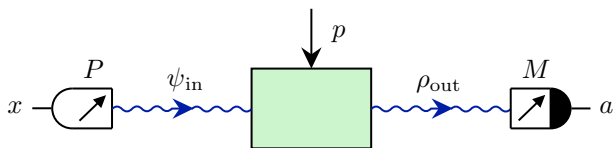
- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



[Nat. Phys. 10, 264 (2014)]

- p are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**, $P : x \rightarrow \psi_{\text{in}}$.
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements $M : \rho_{\text{out}} \rightarrow a$.

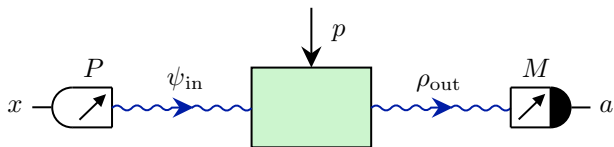
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

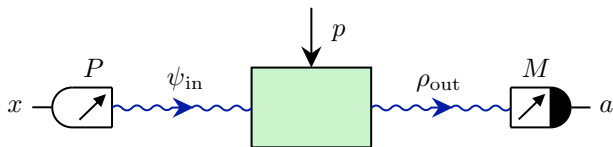
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

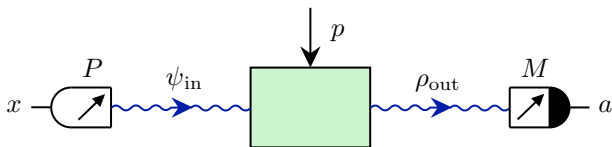
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

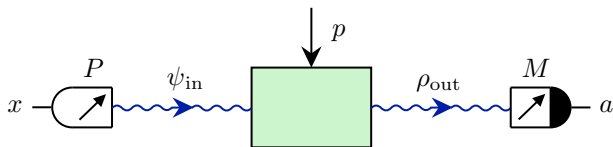
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

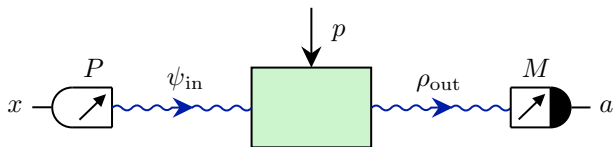
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

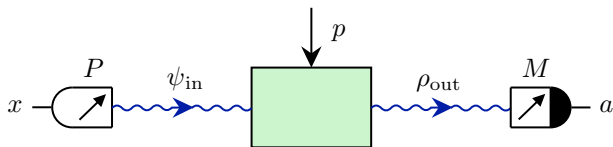
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

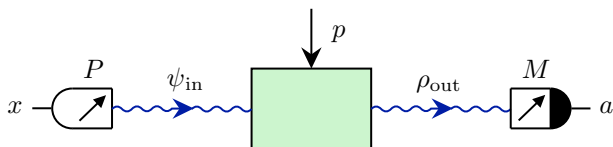
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

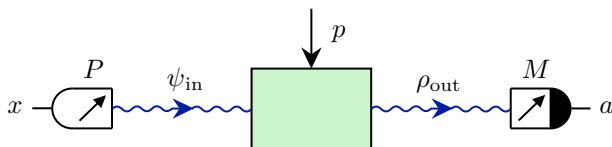
Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state ψ_{in} one performs the full tomography of ρ_{out} .
 - Take a complete set of projectors $\{M_i\}_{i=1}^{n^2-1}$ (e.g. $\{\sigma_x, \sigma_y, \sigma_z\}$).
 - Make *multiple* measurements and register $\{P(a_j | M_i)\}_{i,j}$.
 - The state ρ_{out} is estimated from $\text{Tr}(M_i \rho_{\text{out}}) = \sum_j a_j P(a_j | M_i)$.
- A Q-data test yields a dataset $\{\psi_{\text{in}}^{(k)}, p^{(\ell)}; \rho_{\text{out}}^{(k,\ell)}\}_{k,\ell}$.
- ψ_{in} is pure, initially **uncorrelated** with the box — **freedom of choice**.
- ρ_{out} is in general *mixed*, i.e. entangled with the box.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability 1/2).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

An example — the Helstrom test

- Suppose that we have two available inputs $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$.
- We choose randomly the input (with probability $1/2$).
- The task is to guess, which of the two states was input.
- Define the **success rate**: $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$.
- In quantum theory P_{succ} cannot exceed the **Helstrom bound**

$$P_{\text{succ}} \leq P_{\text{succ}}^{\text{QM}} := \frac{1}{2} \left(1 + \sqrt{1 - |\langle \psi_{\text{in}}^{(1)} | \psi_{\text{in}}^{(2)} \rangle|^2} \right).$$

- Make a Q-data test with $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$.
- If $P_{\text{succ}}(\rho_{\text{out}}^{(1)}, \rho_{\text{out}}^{(2)}) > P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)})$ then the Q-data box is **not** quantum.
- Violation of the Helstrom bound occurs in nonlinear modifications of QM.

Quantum process tomography

- In QM *any* dynamics $\mathcal{E} : S(\mathcal{H}_{\text{in}}) \rightarrow S(\mathcal{H}_{\text{out}})$ must be a **CPTP map**.
- \mathcal{E} is CPTP if and only if $\tilde{\mathcal{E}} := \frac{1}{m} \sum_{i,j=1}^m |i\rangle\langle j| \otimes \mathcal{E}(|i\rangle\langle j|)$ is a quantum state.
- \mathcal{E} is completely characterised by $m^2(n^2 - 1)$ real parameters, $m = \dim \mathcal{H}_{\text{in}}$, $n = \dim \mathcal{H}_{\text{out}}$.
- \mathcal{E} can be reconstructed from a **Q-data test** $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^{m^2}$.
- Overcomplete Q-data tests, $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^N$ with $N > m^2$ are sensitive to deviations from CPTP and linearity.

Quantum process tomography

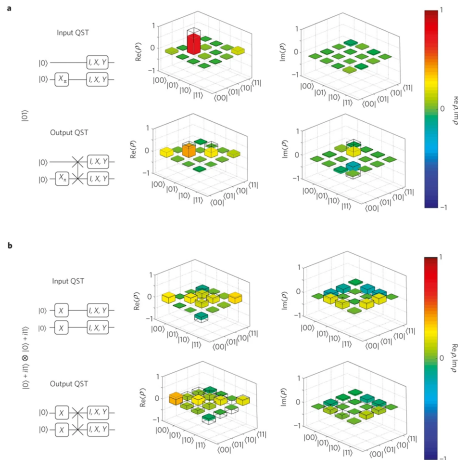
- In QM *any* dynamics $\mathcal{E} : S(\mathcal{H}_{\text{in}}) \rightarrow S(\mathcal{H}_{\text{out}})$ must be a **CPTP map**.
- \mathcal{E} is CPTP if and only if $\tilde{\mathcal{E}} := \frac{1}{m} \sum_{i,j=1}^m |i\rangle\langle j| \otimes \mathcal{E}(|i\rangle\langle j|)$ is a quantum state.
- \mathcal{E} is completely characterised by $m^2(n^2 - 1)$ real parameters, $m = \dim \mathcal{H}_{\text{in}}$, $n = \dim \mathcal{H}_{\text{out}}$.
- \mathcal{E} can be reconstructed from a **Q-data test** $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^{m^2}$.
- Overcomplete Q-data tests, $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^N$ with $N > m^2$ are sensitive to deviations from CPTP and linearity.

Quantum process tomography

- In QM *any* dynamics $\mathcal{E} : S(\mathcal{H}_{\text{in}}) \rightarrow S(\mathcal{H}_{\text{out}})$ must be a **CPTP map**.
- \mathcal{E} is CPTP if and only if $\tilde{\mathcal{E}} := \frac{1}{m} \sum_{i,j=1}^m |i\rangle\langle j| \otimes \mathcal{E}(|i\rangle\langle j|)$ is a quantum state.
- \mathcal{E} is completely characterised by $m^2(n^2 - 1)$ real parameters, $m = \dim \mathcal{H}_{\text{in}}$, $n = \dim \mathcal{H}_{\text{out}}$.
- \mathcal{E} can be reconstructed from a **Q-data test** $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^{m^2}$.
- Overcomplete Q-data tests, $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^N$ with $N > m^2$ are sensitive to deviations from CPTP and linearity.

Quantum process tomography

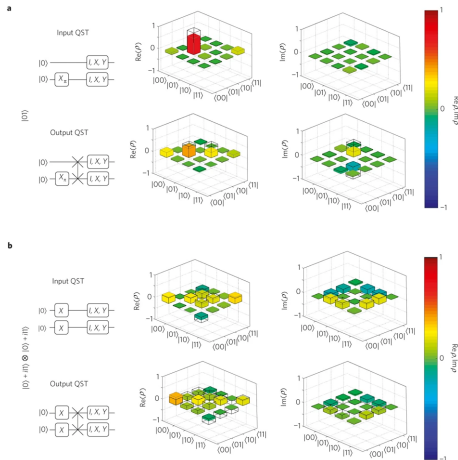
- In QM any dynamics $\mathcal{E} : \mathcal{S}(\mathcal{H}_{\text{in}}) \rightarrow \mathcal{S}(\mathcal{H}_{\text{out}})$ must be a **CPTP map**.
- \mathcal{E} is CPTP if and only if $\tilde{\mathcal{E}} := \frac{1}{m} \sum_{i,j=1}^m |i\rangle\langle j| \otimes \mathcal{E}(|i\rangle\langle j|)$ is a quantum state.
- \mathcal{E} is completely characterised by $m^2(n^2 - 1)$ real parameters, $m = \dim \mathcal{H}_{\text{in}}, n = \dim \mathcal{H}_{\text{out}}$.
- \mathcal{E} can be reconstructed from a **Q-data test** $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^{m^2}$.
- Overcomplete Q-data tests, $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^N$ with $N > m^2$ are sensitive to deviations from CPTP and linearity.



[R. Bialczak et al., *Nat. Phys.* **6**, 409 (2010)]

Quantum process tomography

- In QM *any* dynamics $\mathcal{E} : S(\mathcal{H}_{\text{in}}) \rightarrow S(\mathcal{H}_{\text{out}})$ must be a **CPTP map**.
- \mathcal{E} is CPTP if and only if $\tilde{\mathcal{E}} := \frac{1}{m} \sum_{i,j=1}^m |i\rangle\langle j| \otimes \mathcal{E}(|i\rangle\langle j|)$ is a quantum state.
- \mathcal{E} is completely characterised by $m^2(n^2 - 1)$ real parameters, $m = \dim \mathcal{H}_{\text{in}}, n = \dim \mathcal{H}_{\text{out}}$.
- \mathcal{E} can be reconstructed from a **Q-data test** $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^{m^2}$.
- **Overcomplete Q-data tests**, $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1}^N$ with $N > m^2$ are sensitive to deviations from CPTP and linearity.



[R. Bialczak et al., *Nat. Phys.* **6**, 409 (2010)]

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} ,
e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022),
Quantum State-Channel Duality applied to Particle Physics]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E}
associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Prospects for quantum process tomography in colliders

- 1 Prepare 'quantum-programmed' particles carrying ψ_{in} , e.g. electron's spin or photon's polarization. \rightsquigarrow polarized beams
 - 2 Collide them!
 - 3 Measure projectively the outgoing projectiles.
 - 4 Reconstruct the output states ρ_{out} . \rightsquigarrow weak decays (see Alan's talk)
- [Clelia Altomonte, Alan Barr (2022), *Quantum State-Channel Duality applied to Particle Physics*]
 - Spin dynamics in the $e^+e^- \rightarrow t\bar{t}$ process
 - $\mathcal{H}_{\text{in}} = \mathbb{C}_{e^+}^2 \otimes \mathbb{C}_{e^-}^2$, $\mathcal{H}_{\text{out}} = \mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2$
 - Calculations of the (diagonal part of) the 16×16 matrix \tilde{E} associated with the quantum channel $\mathcal{E} : \mathcal{H}_{\text{in}} \rightarrow \mathcal{H}_{\text{out}}$.
 - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!

Proc. R. Soc. A. **478**:20210806 (2022), arXiv:2103.12000

Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
 - Entanglement detection is not a Bell test!
 - Could we make **direct** projective measurements of spin??
 - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
 - Seek deviations from QM (unitarity, CPTP, linearity, ...)
 - Understand quantum *dynamics* in HEP.
 - Need polarised beams and targets.

Thank you for your attention!