Monte Carlo methods for parton energy loss



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JETSCAPE Summer School 2023



- What is a Monte Carlo event generator? (\mathbf{A})
 - How to generate a parton shower? *
- What do we need to worry about when trying to describe parton showers in heavy-ion events? \Rightarrow
 - In-medium interactions +
 - Parton shower modifications?
 - Medium re-scatterings?
 - Medium evolution? +
- What can we do with a Monte Carlo model for jet quenching? +

Outline





What is a Monte Carlo event generator?



Monte Carlo Event Generators

Physics event:

Quantum mechanics: amplitudes \Rightarrow probabilities *

Everything can happen, but more or less frequently



Monte Carlo Event Generators

- Physics event:
- Quantum mechanics: amplitudes \Rightarrow probabilities *
 - Everything can happen, but more or less frequently
- Monte Carlo Event generators: $\mathbf{+}$
 - Monte Carlo = Random numbers \approx Quantum Mechanical choices +
 - Event generator: trace evolution of the event structure



Monte Carlo Event Generators

- Physics event:
- Quantum mechanics: amplitudes \Rightarrow probabilities *
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 - Monte Carlo = Random numbers \approx Quantum Mechanical choices +
 - Event generator: trace evolution of the event structure

$\mathscr{L}_{int} \longleftrightarrow$ Final states

 $\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i F \mathcal{B} \mathcal{F} + h.c$ + $i F \mathcal{B} \mathcal{F} + h.c$ + $\mathcal{F} \mathcal{B} \mathcal{F}_{3} \mathcal{P} + h.c.$ $+\left|\mathcal{D}_{\mathcal{M}}\varphi\right|^{2}-V(\phi)$







How to describe such a process through an event generator? +

20000



How to describe such a process through an event generator? +

Factorising into simpler problems: *

) Q Q Q Q



How to describe such a process through an event generator? +

Factorising into simpler problems: *

Hard scattering +





How to describe such a process through an event generator? +

Factorising into simpler problems: *

Hard scattering +

Initial-state shower and final-state shower \blacklozenge





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Factorising into simpler problems: *

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MPI and Beam Remnants +





- How to describe such a process through an event generator?
 - Factorising into simpler problems:
 - Hard scattering
 - Initial-state shower and final-state shower
 - MPI and Beam Remnants
 - Hadronization



- How to describe such a process through an event generator? +
 - Factorising into simpler problems: *
 - Hard scattering +
 - Initial-state shower and final-state shower
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Separation in energy

scale

2000

• 2



Initial- and Final-State Showers

Two approaches to calculate additional radiation to the hard scattering:

Matrix elements (few particle corrections but higher order) *

+





The large logarithms are a symptom of interactions far away from the scale at which the coupling was fixed. Fortunately, there Lizzar core third can be an an a symptom of the base of the scale of the methods. As most of the collinear emissions are well separated in scale from the probe, q_0 , these ومومومومور الاموموموموم entise To a proaches to calculate additional radiation to the bard scattering rections to C_a . 0000000000 This results in a *Q*-dependence of the PDF that evolves the probe from the hard scale to lower Matrix elements (few particle corrections but higher order) momentum scales (indicated by the red sub-diagram in Fig 2.9). For the change $Q \to Q + \Delta Q$ the ·····ž····· 3 different Partorb show ever (more iparticle to ornection schieth L. Quandrads Quandrads when tum $Q < p_{\perp} <$ $Q + \Delta Q$ is given by Evolution equation based on splitting probabilities (SF) $\frac{2\pi}{2\pi} \frac{p_{\perp}^2}{p_{\perp}^2} P_{a \leftarrow b}(z) \simeq \frac{\pi}{\pi} \frac{Q}{Q} P_{a \leftarrow b}(z),$ (2.27)where $P_{a \leftarrow b}(z)$ is the splitting function for parton of b splitting into type a, and c_{abc} from the diagrams $\frac{\partial D_a^h(x, Q^2)}{\partial Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dz}{z} \sum_{z} d\hat{R}_z tril(u) iD_b^h \left(f_{z}^x a Q^2\right) a$ at momentum fraction ∂Q^2 . Within the diagram of other partons at $x' \stackrel{b}{=} x/z$, and can be written as BODG . q_0 $f_b(x',Q)\delta(x-zx')$ (2.28)Evolves parton system Splitting Function (SF)¹ $\frac{dz}{\pi} \int_{-\pi}^{1} \frac{dz}{z} f_b(\frac{x}{z}, Q^2) P_{a \leftarrow b}(z).$ from q₀ scale to lower scale

$$\Delta f_a(x,Q) = \sum_{b} \int_0^{\infty} \frac{2\pi}{dx} \int_0^{1} \frac{dz}{dx} \sum_{a} \frac{P_{a \leftarrow b}(z)D}{dz - \frac{b}{\pi}} \frac{\Delta Q}{Q} P_{a \leftarrow b}(z)$$

Probability of parton 'b' splits into parton 'a' with a fraction of energy z



Event generators = Monte Carlo techniques

Selection from a probability distribution function *

 t_0

+



Event generators = Monte Carlo techniques

Selection from a probability distribution function ✦

 t_0



Re-summation of multiple emissions

Sudakov Form factor:

$$\Delta_i(t) \equiv \exp\left[-\sum_j \int_{t_0}^t \frac{dt'}{t'} \int dx \frac{\alpha_s}{2\pi} P_{i \leftarrow j}(x) \right]$$





Event generators = Monte Carlo techniques

Selection from a probability distribution function



Probability of not decay between t₀ and t₁

Given a random number, R, what is t₁? At t_1 , it decays.

 \bigstar



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Just like a radioactive decay!

$$N(t) = \exp\left\{-\int_{t_0}^{t_1} dt f(t') dt'\right\}$$
$$\Rightarrow N(t) = N_0 e^{-\lambda t}$$







Parton Showers in pp

Probabilistic picture allows to build subsequent parton emissions: +





Parton Showers in pp

Probabilistic picture allows to build subsequent parton emissions: +



Jets in proton-proton



And now for something completely different...

... Heavy-ions!



Heavy-Ions Collision

PbPb collision: a complex multi-particle system +





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- Hot and dense medium (QGP)
 - Fluid with collectivity phenomena +
 - Also QCD system, but strongly interacting! +
 - How collectivity emerge from a QFT?
 - How does it evolve? +
 - How is thermalised? **+**
- Products from hard scattering: +
 - Particles modified w.r.t pp: +
 - Jet Quenching effects +



Hard Probes

- Hard probes: Heavy-flavour, Quarkonia, jets, ... See Monday lectures +
 - Produced in a high momentum transfer process (hard scattering)
 - Indirect observation of the QGP effects
 - Observe the evolution of the QGP (temperature, density,...)





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QGP-induced modifications on a proton-proton jet: +





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Medium-induced energy loss





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Medium: Strongly coupled fluid?

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Medium effects on Hadronization?



QGP-induced modifications on a proton-proton jet: +

Start with the building blocks:

Medium-induced energy loss

Collisional energy loss

Medium: Strongly coupled fluid?

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See Ismail Soudi (Fri)







Medium-induced radiation

Within a perturbative QCD perspective, the incoming quark will undergo multiple scatterings with the medium (QGP):





+



Medium-induced radiation

Within a perturbative QCD perspective, the incoming quark will undergo multiple scatterings with the medium (QGP):









In-medium propagators

Adapt Feynman rules to account for a hot and dense QCD medium: +

> Vacuum QCD Feynman rules

Medium longitudinal limits from $[x_{0+}, L_{+}]$






In-medium propagators

Adapt Feynman rules to account for a hot and dense QCD medium: +

> Vacuum QCD Feynman rules

$$= i\delta^{ab}\frac{\not\!\!\!\!\!\!/}{p^2 - m^2 + i\varepsilon}$$

Incoming parton's energy



In-medium propagators

Adapt Feynman rules to account for a hot and dense QCD medium:

Vacuum QCD Feynman rules

$$= i\delta^{ab} \frac{\not p + m}{p^2 - m^2 + i\varepsilon} \tag{(...)}$$

$$G(x_{0+}, \mathbf{x})$$

Incoming parton's energy





$$k_{+}\frac{dI}{dk_{+}d^{2}\mathbf{k}_{\perp}} = \frac{1}{k_{+}}\int_{x_{+}}^{L_{+}} d\bar{x}_{+} \ e^{-\frac{1}{2}\int_{x_{+}}^{L_{+}} d\xi n(\xi)\sigma(\mathbf{x})} \frac{\partial}{\partial \mathbf{y}} \cdot \frac{\partial}{\partial \mathbf{x}} \mathcal{K}(\mathbf{x})$$

 $\sigma(\boldsymbol{r}) \propto V(\boldsymbol{q}) = \frac{8\pi\mu^2}{(\boldsymbol{q}^2 + \mu^2)^2}$ Scattering rate (interaction potential)

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High energy approximation: \Rightarrow Decomposition with a fixed number of propagators \Rightarrow 3 different regions

picture Physical



 $(\mathbf{y}=0, x_+; \mathbf{x}, \bar{x}_+)$





Emission Kernel



 $\sigma(\boldsymbol{r}) \propto V(\boldsymbol{q}) = \frac{8\pi\mu^2}{(\boldsymbol{q}^2 + \mu^2)^2}$ Scattering rate (interaction potential)

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High energy approximation: \Rightarrow Decomposition with a fixed number of propagators \Rightarrow 3 different regions

Physical picture



$$\mathbf{y} = 0, x_+; \mathbf{x}, \bar{x}_+)$$





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High energy approximation: \Rightarrow Decomposition with a fixed number of propagators \Rightarrow 3 different regions

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$$\mathbf{y} = 0, x_+; \mathbf{x}, \bar{x}_+)$$





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Medium response

QGP part that become correlated with the jet:

Seen as (pQCD approach): \bigstar

Recoils from jet-medium interactions with a QGP particle distribution +

Dominated by small momentum transfers (close to non-perturbative region)

E.g: JEWEL
$$\frac{d\hat{\sigma}}{d\hat{t}}(\hat{s},|\hat{t}|) \simeq \frac{C_R 2\pi \alpha_s^2}{(|\hat{t}| + \mu_D^2)^2}$$
E.g:



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LBT: [Cao, Luo, Qin, Wang (16) He, Luo, Wang, Zhu (17)]

MARTINI: [Schenke, Gale, Jeon (09)]

JEWEL: [Elayavalli, Zapp (17)]







Description of a heavy-ion jet

- What is a jet in heavy-ion collisions? +
 - Multi-scale process: *
 - High momentum particles (typically from vacuum-like parton) shower)
 - "Semi-hard" & Soft medium-induced radiation \bullet
 - Soft jet-induced medium response





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- Space-temporal evolving structure:
 - parton fragmentation and parton re-scattering with medium constituents at some time



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Not as "easy" as in pp...

How to describe it?



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Analytical approaches

Based on first principle calculations that address elementary jet processes



Analytical approaches

Based on first principle calculations that address elementary jet processes

- Improvements beyond: \checkmark
 - static medium
 - limited kinematic approximations

- ...



Limited understanding for:

- lower momentum scales
- interplay between "vacuum" and
- "medium"-induced shower



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Monte Carlo approaches

Can consider the full jet shower evolution and evolving medium



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Rely on analytical results



... But lacking most recent analytical developments



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... But lacking most recent analytical developments







What is a jet quenching Monte Carlo?



Jet quenching Monte Carlo models

- N-particle system originated through a parton shower
- Vacuum radiation \bigstar
 - Medium-induced effects
 - Medium-induced radiation \bigstar
 - Jet-induced medium response
 - Medium response re-scattering







Hadronization

Medium-modified jet in all momentum scales?

























Two different approaches: +

Change in the jet evolution:

Modifications on a developed shower



Two different approaches: +

Change in the jet evolution:

Medium-induced modifications can take place throughout the parton evolution

Medium-modifications at all momentum scales



E.g: JEWEL, <u>MATTER</u>, Q-PYTHIA,...

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Modifications on a developed shower



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Modifications on a developed shower

Vacuum (hard and collinear) parton structure unmodified

Medium-modifications dominate low momentum scales



E.g:(Co-)LBT, Hybrid, MARTINI, JetMed..



Change in the jet evolution:

$$\omega, k_{\perp}$$

$$dP^{q \to qg} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{dk_{\perp}^2}{k_{\perp}^2}$$

Re-summation of multiple emissions

$$\Delta_i(t) \equiv \exp\left[-\sum_j \int_{t_0}^t \frac{dt'}{t'} \int dx \frac{\alpha_s}{2\pi} P_{i \leftarrow j}(x)\right]$$

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+



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+

$$k = zp = (k_+, k_-, \mathbf{k}_\perp)$$



$$q = (1 - z)p$$





Change in the jet evolution:

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+

$$k = zp = (k_+, k_-, \mathbf{k}_\perp)$$







Ansatz: $P_{\mathrm{tot}} = P_{\mathrm{vac}} + \Delta P$





Change in the jet evolution:

$$\omega, k_{\perp}$$

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0.35 -

0.05

0.00

 10^{-2}

 10^{-1}

Full Yukawa





q = (1 - z)p

Ansatz: $P_{\rm tot} = P_{\rm vac} + \Delta P$



10

 κ^2

 $\Delta P \simeq \frac{2\pi t}{\alpha_s} \, \frac{dI^{\rm med}}{dzdt}$

Modifications on a developed shower

Medium-induced emissions inside them medium: $t_f \leq \sqrt{2\omega/\hat{q}}$ (But no double logarithmic enhancement)

+

Parton formation time: $t_f \simeq 2\omega/k_{\perp}^2$

Transverse momentum acquired via multiple soft scatterings: $k_f^2 = \hat{q}t_f$.



Modifications on a developed shower

Medium-induced emissions inside them medium: $t_f \leq \sqrt{2\omega/\hat{q}}$ (But no double logarithmic enhancement)

(Vacuum emissions develop much faster than vacuum ones)

Parton formation time: $t_f \simeq 2\omega/k_{\perp}^2$

Transverse momentum acquired via multiple soft scatterings: $k_f^2 = \hat{q}t_f$.

Vacuum-like emissions inside them medium: $k_{\perp} \gg k_{\rm f} \Leftrightarrow \frac{2}{\omega \theta^2} \ll \sqrt{\frac{2\omega}{\hat{q}}} \Leftrightarrow t_f \ll \sqrt{2\omega/\hat{q}}$



Modifications on a developed shower

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Parton formation time: $t_f \simeq 2\omega/k_{\perp}^2$

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Vacuum-like emissions inside them medium: $k_{\perp} \gg k_{\rm f} \Leftrightarrow \frac{2}{\omega \theta^2} \ll \sqrt{\frac{2\omega}{\hat{q}}} \Leftrightarrow t_f \ll \sqrt{2\omega/\hat{q}}$ Vacuum-like structures **Medium-induced radiation** 000000





Comparison between the two: +

Change in the jet evolution:

Choose (or develop) a given vacuum parton shower (Fixed to the ordering variable and parton shower accuracy)

Modifications on a developed shower

Minimal changes to the vacuum parton shower (Easier to develop alongside vacuum physics)



Comparison between the two:

Change in the jet evolution:

Choose (or develop) a given vacuum parton shower (Fixed to the ordering variable and parton shower accuracy)

Medium-induced effects from in-medium radiation spectrum (inheriting kinematical restrictions)

Modifications on a developed shower

Minimal changes to the vacuum parton shower (Easier to develop alongside vacuum physics)

Medium-induced effects from transport equations (inheriting kinematical restrictions)



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Choose (or develop) a given vacuum parton shower (Fixed to the ordering variable and parton shower accuracy)

Medium-induced effects from in-medium radiation spectrum (inheriting kinematical restrictions)

Modifications done in momentum scales relatively above the non-perturbative region

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Minimal changes to the vacuum parton shower (Easier to develop alongside vacuum physics)

Medium-induced effects from transport equations (inheriting kinematical restrictions)

Modifications in the low-momentum particle distribution (close to non-perturbative region)



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No "correct" answer... All with their pros and cons...



Comparison between the two:

Change in the jet evolution:



Choose (or develop) a given vacuum parton shower (Fixed to the ordering variable and parton shower accuracy)

> Interplay between vacuum and medium shower

Modifications on a developed shower

Modifications in the low-momentum particle distribution (close to non-perturbative region)





Comparison between the two:

Change in the jet evolution:



Choose (or develop) a given vacuum parton shower (Fixed to the ordering variable and parton shower accuracy)

Interplay between vacuum and medium shower

MATTER (High-virtuality part of the shower)



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Modifications on a developed shower

Modifications in the low-momentum particle distribution (close to non-perturbative region)

QCD processes at lower momentum scales

LBT, MARTINI,... (Low-virtuality part of the shower)



Elastic Energy Loss

Need phase space density of scattering centres (sampled from hydro profile or Bjorken evolution model) +



Medium recoil

 $\frac{d\hat{\sigma}}{d\hat{t}}(\hat{s},|\hat{t}|) \simeq \frac{C_R 2\pi \alpha_s^2}{(|\hat{t}| + \mu_D^2)^2}$

Coupled jet-hydro evolution: +

$$\partial_{\mu}T^{\mu\nu}_{fluid} = J^{\nu}_{jet}(x)$$



Recoiling particles can further re-scatter

> Cooper-Frye for particles from medium response


Medium evolution model \blacklozenge

Bjorken 1D expansion

[Bjorken (1983)]



$$T = T_0 \left(\frac{\tau_0}{\tau}\right)^{v_s^2}$$

Longitudinal (1D) expansion (Energy density characterised by a power-law evolution)

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Medium evolution model

Bjorken 1D expansion

[Bjorken (1983)]



$$T = T_0 \left(\frac{\tau_0}{\tau}\right)^{v_s^2}$$

Longitudinal (1D) expansion (Energy density characterised by a power-law evolution)

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Event-by-event non-ideal hydrodynamics

[Molner et al(1407.8152), Shen et al (1409.8164),...]



$$\partial_{\mu}T^{\mu\nu} = j^{\nu}$$

3D expansion (Energy density characterised by relativistic hydrodynamic evolution)



Medium evolution model

Bjorken 1D expansion

[Bjorken (1983)]



$$T = T_0 \left(\frac{\tau_0}{\tau}\right)^{v_s^2}$$

Longitudinal (1D) expansion (Energy density characterised by a power-law evolution)

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+

See Mayank Singh (Wed)

Event-by-event non-ideal hydrodynamics

[Molner et al(1407.8152), Shen et al (1409.8164),...]



$$\partial_{\mu}T^{\mu\nu} = j^{\nu}$$

3D expansion (Energy density characterised by relativistic hydrodynamic evolution)



Uncertainty driven by the onset of medium-jet interactions... +



[Andrés, et al (1902.03231), Stojku et al (2008.08987), JETSCAPE (2102.11337), Adhya et al (2211.15803)]



Uncertainty driven by the onset of medium-jet interactions... +



Jet-medium interactions start at t₀? What happens before?

[Andrés, et al (1902.03231), Stojku et al (2008.08987), JETSCAPE (2102.11337), Adhya et al (2211.15803)]





Production point (for path-length dependence): +

Sampled from initial nuclei overlap *







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Production point (for path-length dependence): +

Sampled from initial nuclei overlap *

Needs to account for space-time structure of jets

In vacuum parton showers?

Parton formation time:







Production point (for path-length dependence): +

Sampled from initial nuclei overlap \bigstar

Needs to account for space-time structure of jets

In vacuum parton showers?

Parton formation time:

Hadronization: +

Usually taken from PYTHIA (might include recoiled particles)







State-of-the-art models

Several jet quenching Monte Carlo models: +

Q-PYTHIA

PYQUEN

JEWEL





MATTER



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See references in the backup slides

Jetmed(Saclay)





DREENA-A



Hybrid strong/weak coupling

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What can we do with a Monte Carlo model for jet quenching?

The successes of Monte Carlo approaches



Quantifying QGP properties

8

- Jets are formed in the beginning of the collision:
- Allow detailed imaging of the QGP \Rightarrow
 - QGP evolution (E.g: thermaligation process)
- Formed by collection of soft to hard particles \blacklozenge
 - Allow QGP probing by different scales
 - Scale dependent quantities (Eg.: "quasi-particles")







From the jet to the medium

Medium-induced radiation and momentum broadening closely connected (multiple soft-scattering approximation)

Accumulation of momenta enhances gluon radiation and partons undergo transverse momentum broadening $\mathbf{\mathbf{A}}$



+

Dipole cross-section (collision rate):

$$\sigma(\boldsymbol{r}) = \int_{\boldsymbol{q}} V(\boldsymbol{q}) \left(1 - e^{i\boldsymbol{q}\boldsymbol{r}} \right)$$



From the jet to the medium

Medium-induced radiation and momentum broadening closely connected (multiple soft-scattering approximation) +

Accumulation of momenta enhances gluon radiation and partons undergo transverse momentum broadening *







From the jet to the medium

Medium-induced radiation and momentum broadening closely connected (multiple soft-scattering approximation)

Accumulation of momenta enhances gluon radiation and partons undergo transverse momentum broadening



Transport coefficient:

$$\hat{q} = \frac{\langle k_T \rangle}{\lambda}$$
$$\hat{q} \propto \int d^2 \mathbf{q}^2 q^2 \frac{d\sigma(\mathbf{q})}{d^2 \mathbf{q}}$$



+

*



16

14

12

10

8

6

4

2

0

0.2

0.4

T (GeV)

â/T³

From single-particle or jet suppression recover \hat{q} +

M. Xie et. al, 2206.01340 LIDO, 2010.13680 C. Andres et. al, KLN LHC, 1606.04837



0.6



0.8

From single-particle or jet suppression recover \hat{q} +

Changing QGP initialisation conditions

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From single-particle or jet suppression recover \hat{q} +

Changing QGP initialisation conditions

Energy loss during all parton shower evolution vs energy loss during final stage (Compensation of effects with higher transport coefficient)

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From single-particle or jet suppression recover \hat{q} +

Changing QGP initialisation conditions

Energy loss during all parton shower evolution vs energy loss during final stage (Compensation of effects with higher transport coefficient)

Improved Bayesian analysis gives a stronger temperature dependence

See John Miller (

JETSCAPE

[LA, Y-J Lee, M. Winn (2203.163)



Т	ue	2	5)
52	2)]		
7			



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Hadron vs Jet measurements (model-dependent description of medium response on jets)

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JETSCAPE

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Quantitative assessment of QGP characteristics using hard probes

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See John Miller (

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[LA, Y-J Lee, M. Winn (2203.163)



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52	2)]		
7			



The elusive medium response

Soft components seem necessary for a better description of the jet radial profile and/or jet mass:

LBT



+



[Park, Jeon, Gale (1807.06550)]





The elusive medium response

Soft components seem necessary for a better description of the jet radial profile and/or jet mass: +

[Park, Jeon, Gale (807.06550)] 1.6 Pb-Pb @ 2.76TeV (0-10%) anti- $k_T R=0.3$ $(\Delta r)^{pp}$ $p_T^{\;jet}>\!\!100{
m GeV/c}$, $p_T^{trk}>\!\!1{
m GeV/c}$ < 2.0With recoil LBT 0-30% w. medium response $p^{assoc} > 1 \text{ GeV}$ ··· 2.5 Γ 0-30% w/o. medium response //o recoil $p_{-}^{assoc} > 1 \text{ GeV}$ With recoil $\rho_{Pb+Pb}^{}/\rho_{p+p}^{}$ 2 $p_{\tau}^{\gamma} > 80 \text{ GeV}$ 0.3 $\sqrt{s} = 2.76 \text{ TeV}$ 1.5 R = 0.3w/o recoil (b) 0.5 0.05 0.1 0.15 0.2 0.25 0.3 r [Luo, Cao, He, Wang, (1803.06785)]

Is the enhancement due to medium-response or to poorly known non-perturbative physics?

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QGP-wake signal

Jet-induced medium exceptions in Z+jet events: +



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QGP-wake signal

Jet-induced medium exceptions in Z+jet events: +



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Introduction of viscous hydro in MC \Rightarrow 3D Wake that depend on EoS



Comparison between quenched and unquenched made through some jet selection:

Impact of jet selection biases on jet substructure observables? +



How to compare unmodified with modified jets?

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Which jet selection results from applying grooming?



$$\frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{cut} \left(\frac{\Delta R_{12}}{R}\right)^{\rho}$$



Comparison between quenched and unquenched made through some jet selection:

Impact of jet selection biases on jet substructure observables? +

Hybrid





Jets passing the Soft Drop condition are more likely to have medium-induced/recoil effects





Comparison between quenched and unquenched made through some jet selection:

Impact of jet selection biases on jet substructure observables?

Hybrid



*



Jets passing the Soft Drop condition are more likely to have medium-induced/recoil effects



Grooming can be used to select different contributions of medium response







Comparison between quenched and unquenched made through some jet selection:

Impact of jet selection biases on jet substructure observables? +





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"Formation time" can also select jets with different degrees of quenching





Comparison between quenched and unquenched made through some jet selection:

Impact of jet selection biases on jet substructure observables?





*

"Formation time" can also select jets with different degrees of quenching

 $x_{j,Z}$









Towards the Future

. 1



Parton Showers in heavy-ions

Accuracy bounded by proton-proton +

But have a qualitatively different problem: quantum system developing on top of an evolving medium $\mathbf{\mathbf{A}}$



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Parton Showers in heavy-ions

Accuracy bounded by proton-proton +

But have a qualitatively different problem: quantum system developing on top of an evolving medium



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non pQCD

Heavy-ions are unique laboratory for:

- QGP tomography
- Interplay of parton showers with evolving medium
- Transition from perturbative to non-perturbative







Summary

and a

.



223.24

Summary

- Monte Carlo event generators widely used tools to probe QGP physics +
 - *
 - Ideal framework to probe novel QCD-related phenomena

Require phenomenological extensions that need to be constantly tested and refined by analytical input



Summary

- Monte Carlo event generators widely used tools to probe QGP physics

 - Ideal framework to probe novel QCD-related phenomena

- Invaluable instruments for: **+**
 - Testing new observables in more "realistic" conditions as compared to analytical approaches +
 - +
 - Understanding biases in our experimental results

+

Require phenomenological extensions that need to be constantly tested and refined by analytical input

Phenomenological studies targeting phenomena whose analytical description is (still) challenging

Thank you!







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Acknowledgments





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Backup Slides

.



Bibliography

- Monte Carlo models for jet quenching: (\mathbf{A})
 - CUJET: [Buzzatti, Liao, Gyulassy, Shi (14, 16, 18)] \blacklozenge
 - Dreena: [Zigic, Salom, Auvinen, Djordjevic, M. Djordjevic (19, 22)] +

 - JETSCAPE: [JETSCAPE Collab. (17)] \bigstar
 - [Krauss, Wiedemann, Zapp(13); Zapp (14); Elayavalli, Zapp (16;17)] JEWEL: +
 - LBT/Co-LBT: [Wang and Y. Zhu (16); Cao, Luo, Qin, Wang (15); He, Luo, Wang, Zhu (17);] +
 - MARTINI: [Schenke, Gale, Jeon (09); Park, Jeon, Gale (18)] +
 - [Majumder (13); Kordell, Majumder (17); Cao, Majumder (18)] MATTER: +
 - PYQUEN: [Lokhtin, Snigirev (06)]
 - ♦ Q-PYTHIA: [Armesto, Cunquero, Salgado (09)]
 - Jetmed(Saclay): [Caucal, Iancu, Mueller, Soyez (18)]

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Hybrid Strong/Weak coupling: [Casalderrey-Solana, Gulhan, Milhano, Pablos, Rajagopal (14;17); Helcher, Pablos, Rajagopal (18)]









Invariant mass distribution: $\tau = \frac{m^2}{E^2} = \sum_{i=\text{gluon}} z_i \theta_i^2 \longrightarrow \log \tau = \log z + \log \theta^2$ +









 $P(\text{emit in region } i) = \frac{\alpha_s C_F}{\pi} \cdot (\text{Area of region } i)$

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$$\log \tau = \log z + \log \theta^2$$



 $P(\text{no emit in region } i) = 1 - \frac{\alpha_s C_F}{\pi} \cdot (\text{Area of region } i)$

$$P(\text{no emissions}) = \left(1 - \frac{\frac{\alpha_s}{\pi} \frac{C_F}{2} \log^2 \tau}{N}\right)^N \qquad P(\text{no emissions}) = \exp\left[\frac{1 - \frac{\alpha_s}{\pi} \frac{C_F}{2} \log^2 \tau}{N}\right]^N$$







P(no emis)

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$$\log \tau = \log z + \log \theta^2$$

Area of region $i = \frac{\frac{1}{2} \log^2 \tau}{N}$

 $P(\text{no emit in region } i) = 1 - \frac{\alpha_s C_F}{\pi} \cdot (\text{Area of region } i)$

ssions) =
$$\left(1 - \frac{\frac{\alpha_s}{\pi} \frac{C_F}{2} \log^2 \tau}{N}\right)^N$$

Large N limit: Sudakov Form Factor

$$\Delta_i(t) \equiv \exp\left[-\sum_j \int_{t_0}^t \frac{dt'}{t'} \int dx \frac{\alpha_s}{2\pi} P_i\right]$$

ssions) = exp
$$\left[-\frac{\alpha_s}{\pi}\frac{C_F}{2}\log^2\tau\right]$$





Toy Parton Shower

No-emission probability:

Interpretations for the scale:



(Formation time)

$$s \rightarrow \tilde{m}^2 = 2p^+ t_{\rm f}^{-1}$$

(Virtuality)

$$s \rightarrow \tilde{\theta}^2 = \frac{|\boldsymbol{\ell}|^2}{(p^+)^2 [z(1-z)]^2}$$
 Angle)

$$\Delta(s_{\rm prev}, s) = \exp\left\{-\frac{\alpha C_R}{\pi} \int_s^{s_{\rm prev}} \frac{\mathrm{d}\mu}{\mu} \int_{z_{\rm cut}(\mu)}^1 \frac{\mathrm{d}z}{z}\right\}$$

To generate a splitting:



1. Sample a scale from $\Delta(s_{\rm prev}, s)$ 2. Sample a fraction from $\ \hat{P}(z) \propto 1/z$ 3. Retrieve the momenta $\{p_i^+, p_i\}$ Ensure that $|\boldsymbol{\ell}|^2 > k_{\text{had}}^2$



QCD Antenna setup: emission from a qqbar pair: +





QCD Antenna setup: emission from a qqbar pair: +



Probability of emitting "soft" (low-energy) gluons:

$$dN_q^{\omega \to 0} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{\sin\theta \, d\theta}{1 - \cos\theta} \Theta(\cos\theta_1 - \cos\theta_1)$$

 $(\cos \theta)$

QCD Angular ordering



QCD Antenna setup: emission from a qqbar pair: +



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For large angle emissions:



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QCD Angular ordering

For large angle emissions:





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Two in-medium emissions: emission from a qqbar pair: +

Vacuum QCD

$dN_q^{\omega \to 0} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{\sin\theta \, d\theta}{1 - \cos\theta} \Theta(\cos\theta_1 - \cos\theta)$



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In-medium QCD

$$dN_q^{\omega \to 0} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{\sin \theta d\theta}{1 - \cos \theta} \left[\Theta(\cos \theta_1 - \cos \theta)\right]$$



+

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[Mehtar-Tani, Tywoniuk, Salgado (1009.2965)]

$$\Delta_{med} \approx 1 - \mathrm{e}^{-\frac{1}{12}}$$

In-medium QCD

$$dN_q^{\omega \to 0} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{\sin \theta d\theta}{1 - \cos \theta} \left[\Theta(\cos \theta_1 - \cos \theta) \right]$$

 $\Delta_{med} \Theta(\cos \theta - \cos \theta)$

QCD Anti-Angular ordering

Antenna Transverse resolution: $r_{\perp} = \theta L$ Medium Transverse Scale: Qs⁻¹= $\sqrt{(q^{L})^{-1}}$









Two in-medium emissions: emission from a qqbar pair: +

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$$dN_q^{\omega \to 0} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{\sin\theta \, d\theta}{1 - \cos\theta} \Theta(\cos\theta_1 - \cos\theta)$$



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QCD Anti-Angular ordering









Color (de)coherence



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QGP Probes

- Soft probes: flow, hydrochemistry, ... +
 - Direct result of the QGP evolution $\mathbf{\mathbf{A}}$
 - Collective properties and hydrodynamical evolution of the medium





- Soft probes: flow, hydrochemistry, ... +
 - Direct result of the QGP evolution *
 - Collective properties and hydrodynamical evolution of the medium

- Hard probes: Heavy-flavour, Quarkonia, jets, ... +
 - Produced in a high momentum transfer process (hard scattering)
 - Indirect observation of the QGP effects
 - Observe the evolution of the QGP (temperature, density,...)





Adapt Feynman rules to account for a hot and dense QCD medium:

Vacuum QCD Feynman rules

$$= i\delta^{ab}\frac{\not\!\!\!\!\!/}{p^2 - m^2 + i\varepsilon}$$



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In-medium QCD Feynman rules





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> Vacuum QCD Feynman rules

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In-medium QCD Feynman rules



Wilson line (change in colour):

$$W(x_{0+}, L_{+}; \mathbf{x}_{\perp}) = \mathcal{P} \exp\left\{ ig \int_{x_{0+}}^{L_{+}} dx_{+} A_{-}(x_{+}, \mathbf{x}_{\perp}) \right\}$$

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58

Adapt Feynman rules to account for a hot and dense QCD medium: +

Due to Lorentz contraction one can further assume

$$A(x_{+}, x_{-}, x_{\perp}) = A(x_{+}, x_{\perp})$$



In-medium QCD Feynman rules Gluon fields Wilson line (change in colour): $W(x_{0+}, L_+; \mathbf{x}_\perp) = \mathcal{P} \exp\left\{ ig \int_{x_{0+}}^{L_+} dx_+ A_-(x_+, \mathbf{x}_\perp) \right\}$ Path-ordered gluon fields

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Adapt Feynman rules to account for a hot and dense QCD medium: +

> Vacuum QCD Feynman rules

Medium longitudinal limits from $[x_{0+}, L_{+}]$







Modifications on a developed shower

Medium-induced emissions inside them medium: $t_f \leq \sqrt{2\omega/\hat{q}}$ (But no double logarithmic enhancement)

+

Parton formation time: $t_f \simeq 2\omega/k_{\perp}^2$

Transverse momentum acquired via multiple soft scatterings: $k_f^2 = \hat{q}t_f$.



Modifications on a developed shower

Medium-induced emissions inside them medium: $t_f \leq \sqrt{2\omega/\hat{q}}$ (But no double logarithmic enhancement)

Vacuum-like emissions inside them medium: $t_f \ll \sqrt{2\omega/\hat{q}}$ (Vacuum emissions develop much faster than vacuum ones)

Parton formation time: $t_f \simeq 2\omega/k_{\perp}^2$

Transverse momentum acquired via multiple soft scatterings: $k_f^2 = \hat{q}t_f$.



Modifications on a developed shower

Medium-induced emissions inside them medium: $t_f \leq \sqrt{2\omega/\hat{q}}$ (But no double logarithmic enhancement)



At double logarithmic accuracy, medium-induced radiation is vetoed

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Modifications on a developed shower

(But no double logarithmic enhancement)



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