

# HADRONIZATION



**RAINER J FRIES**  
**TEXAS A&M UNIVERSITY**

- Hadronization in QCD
- How to make hadrons in five different ways
- Hybrid Hadronization
- Hadronization in JETSCAPE and XSCAPE

# PART I: HADRONIZATION IN QCD

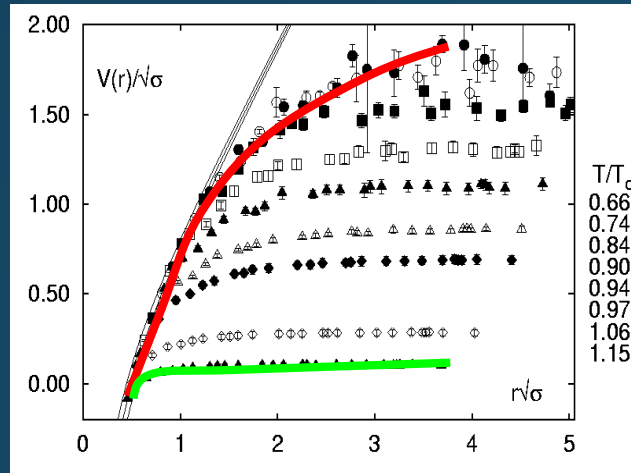
# WHAT'S THE DEAL WITH HADRONIZATION?

- The fundamental degrees of freedom in QCD (quarks and gluons) form bound states at low energies.
  - This true also in QED, but there is an added complication in QCD: Confinement.
  - Quarks and gluons *must* be in color singlet bound states, i.e. mesons or baryons.
- ⇒ Every simulation or event generator with parton physics must know how to get the partons from hadrons and how to turn partons back into hadrons (hadronization).
- Confinement is not understood from first principles.



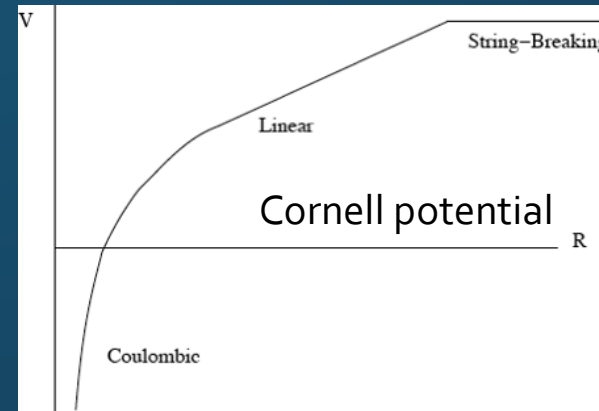
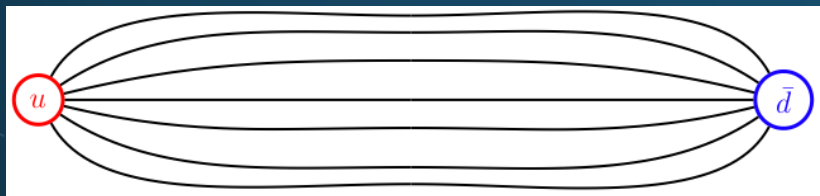
# HADRONIZATION $\leftrightarrow$ QCD AT LARGE DISTANCES

- Confinement can be seen on the lattice:



Static heavy quark potential at different temperatures (Karsch et al.)

- Phenomenology: Dual Meissner Effect in QCD vacuum leads to flux tubes  $\Rightarrow$  string picture!



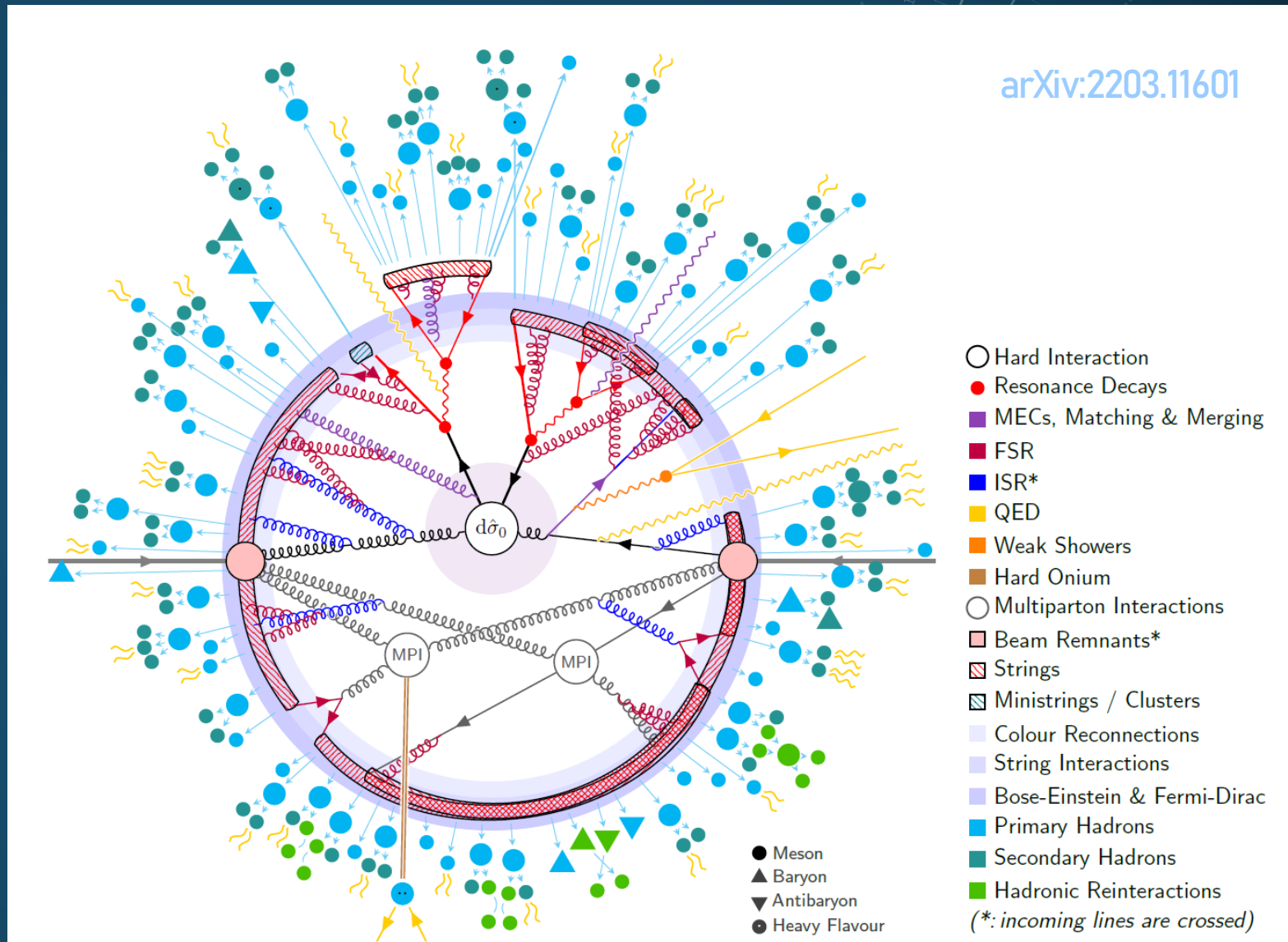
$$V(r) = -\frac{a}{r} + Kr$$

# THE CHALLENGE

In the absence of first principle calculations we need to model hadronization.

Models should:

- Obey relevant observation laws and symmetries.
- Enforce confinement.
- Implement as many phenomenological properties of QCD as possible.





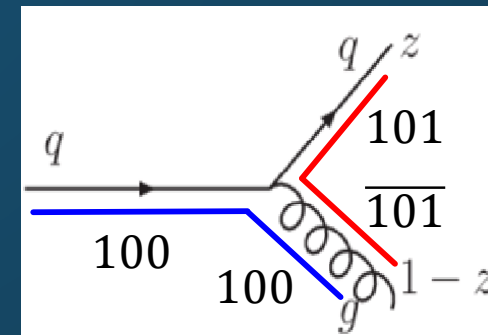
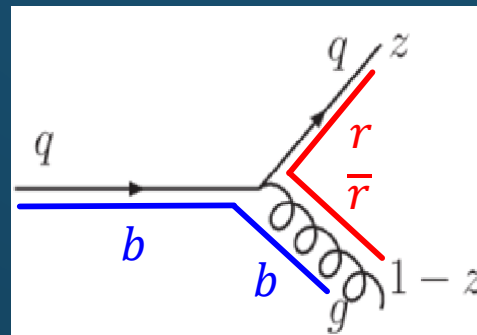
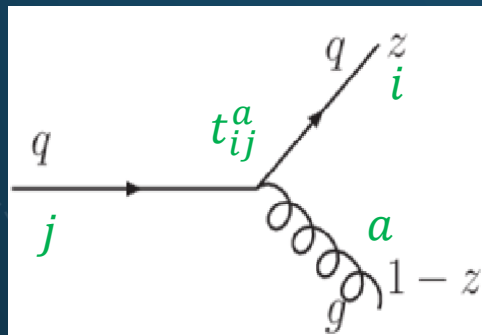
# THE ROLE OF COLOR

- QCD color is based on the group SU(3).
- We are not interested in color degrees of freedoms, our world is “white”. Calculations thus typically sum or average over color degrees of freedom → “color factors”.
- Parton based Monte Carlo’s typically employ the correct color factors for perturbative processes. In addition, there might be a need to track the “direction” of a parton is SU(3) as well: e.g. we might want to track which partons can form color singlets for hadronization.
- Instead of the full SU(3) information partons are assigned color tags in an  $N \rightarrow \infty$  approximation.

$$\lambda_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad \lambda_5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}$$

$$\lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad \lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} \quad \lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$



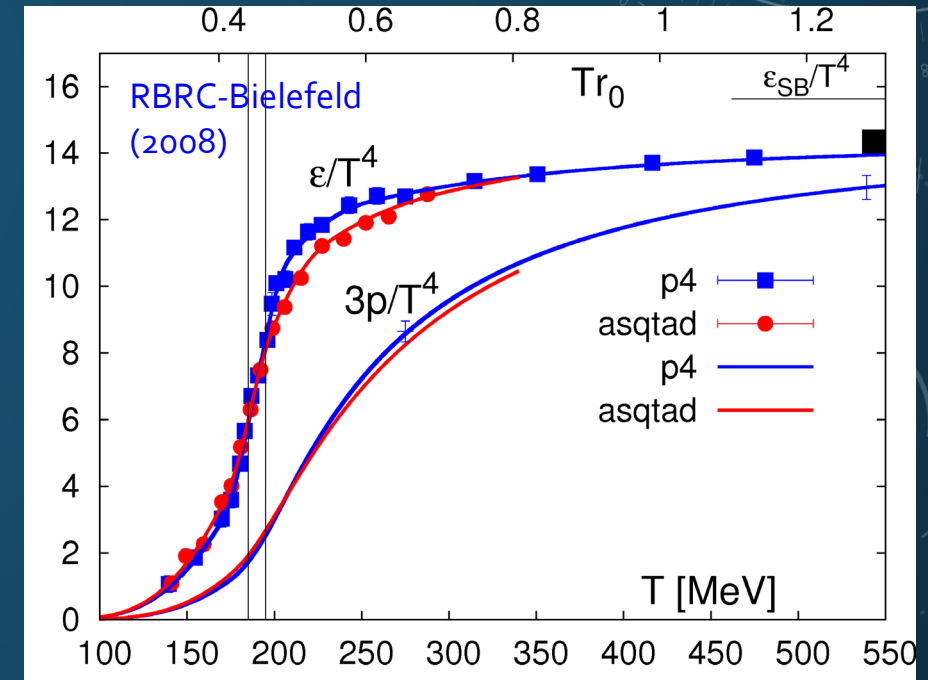


# PART II: HOW TO MAKE HADRONS IN FIVE DIFFERENT WAYS



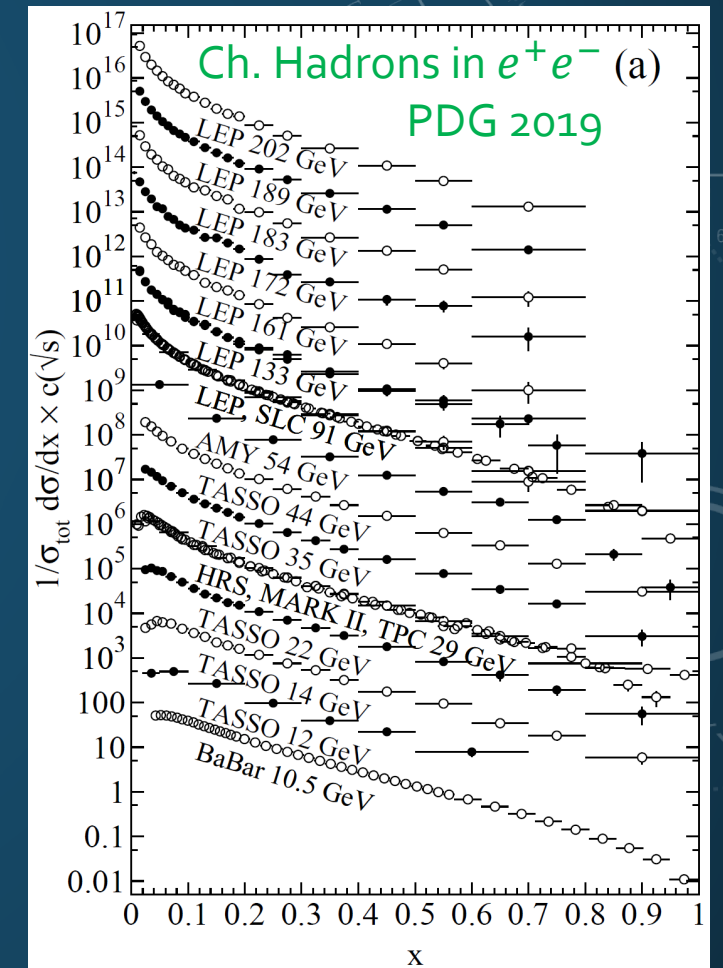
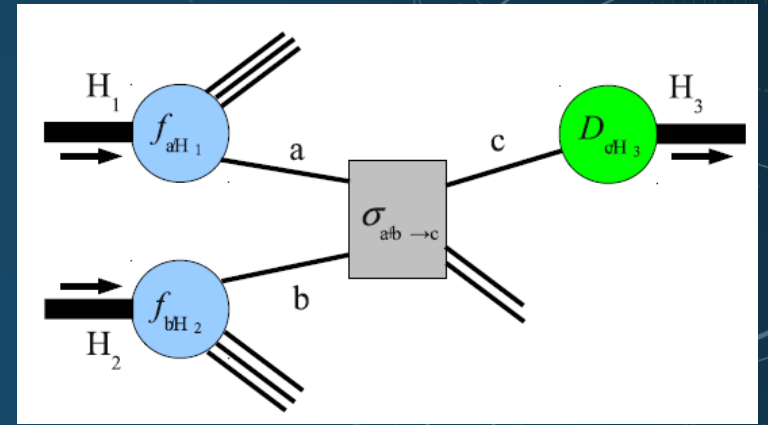
# THE HYDRO ANGLE

- In the long wave-length, large time, large volume limit of QCD we do not care about individual partons and hadrons. We only want to know average, thermodynamic quantities.
- Applies for fluid dynamic simulations for the bulk of heavy ion collisions!
- It is sufficient to know the equation of state (+ any necessary transport coefficients) around  $T_c$  from lattice QCD or experiment.
- For an event-by-event description a hypersurface at a lower temperature can be sampled to create hadrons.
- This avoids the much more challenging particle-by-particle modelling required in parton transport models!



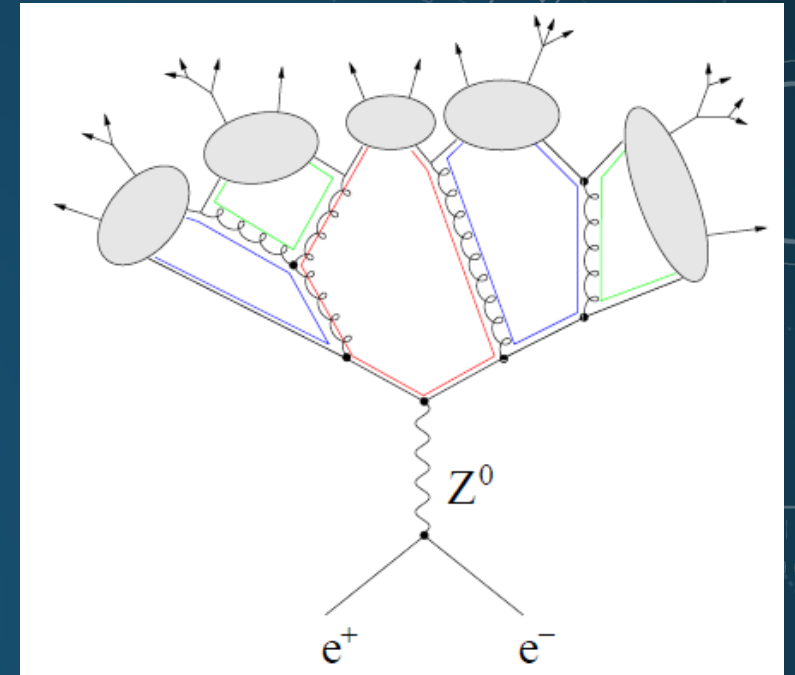
# INDEPENDENT FRAGMENTATION

- A first-principle approximation based on QCD factorization, but the necessary fragmentation functions are not calculable (even lattice is hard).
- Proper operator definition. E.g. quark fragmenting into a hadron H:
- $D_{q/H}(z) = z \int \frac{dy^-}{4\pi} e^{-iP^+ y^- / z} \left\langle 0 \left| q(0) a_H^+(P^+) a_H(P^+) \bar{q}(x^-) \right| 0 \right\rangle$   
 J. C. Collins and D. E. Soper, *Nucl.Phys.* 194, 445 (1982)
- Perfect for semi-inclusive pQCD but could be more widely applied as a model.
- The input are partons from a hard process without final state radiation!
- Fragmentation functions ~ final state radiation + hadronization



# CLUSTER HADRONIZATION

- Primarily for jet Monte Carlos: take fully developed parton showers.
- Force non-perturbative  $g \rightarrow q\bar{q}$  splitting
- Local color neutrality:  $q\bar{q}$  form color-neutral clusters.
- Clusters typically have masses of several GeV; decay into hadrons which can decay further into stable hadrons.
- Some similarities with recombination!

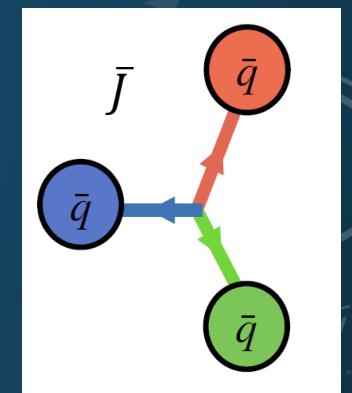
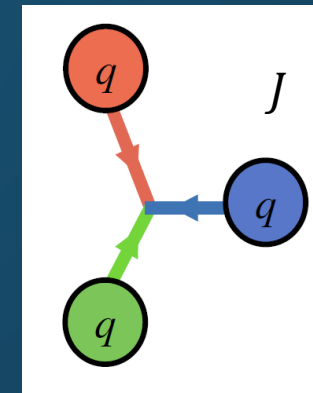
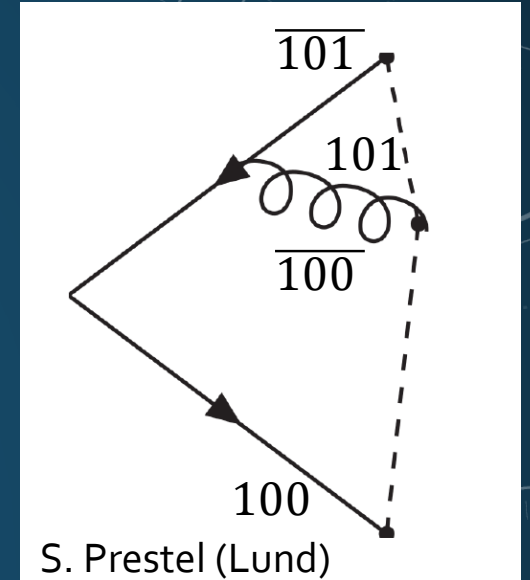


Implemented in HERWIG.

G. Marchesini, B.R. Webber et al.

# STRING FRAGMENTATION

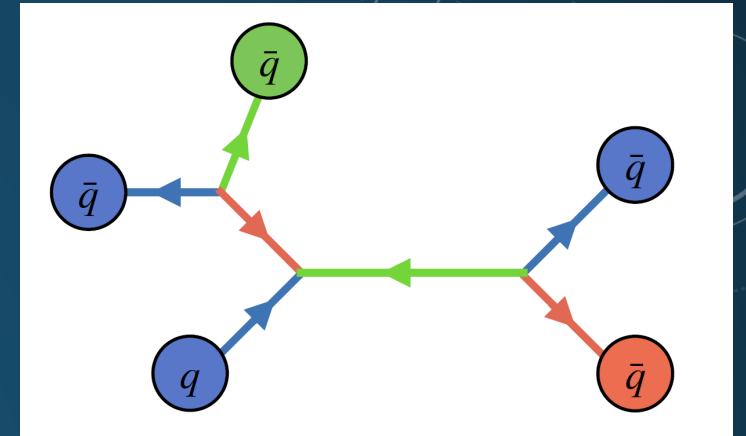
- Again primarily applied in jet Monte Carlos, most notably PYTHIA and JETSET.
- Keep track of color singlet configurations via color tags.
- In a color singlet string configuration each color tag (except those in junctions) has a matching anti-color tag and vice versa and they are connected by the string.
- Strings terminate in (anti)quarks, (anti)diquarks or can connect to a (anti)junction.
- Gluons are “kinks” in strings.
- (Anti)junctions are objects with three (anti)color tags in a singlet configuration that carry baryon number  $\pm 1$ .
- Simple examples (they could have many gluons or connect to other junctions):



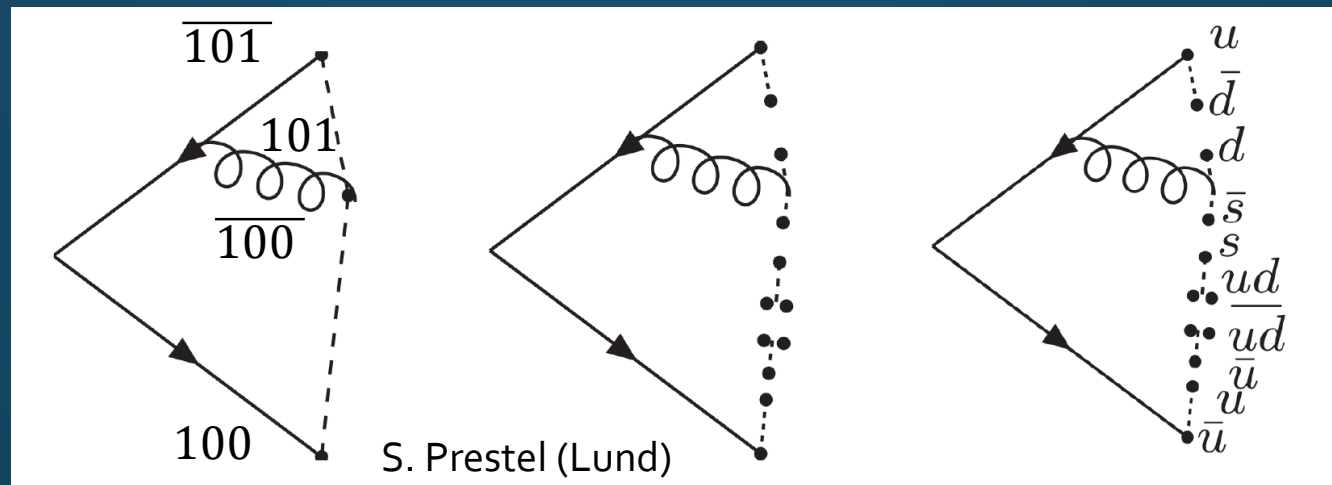
J. Altmann (Monash)

# STRING FRAGMENTATION

- The formalism allows for very complicated string topologies.
- Anything beyond di-junctions can usually not be handled by PYTHIA 8.
- Break strings by pair creation of quarks or diquarks.
- Let the emerging color singlet systems collapse into hadrons.



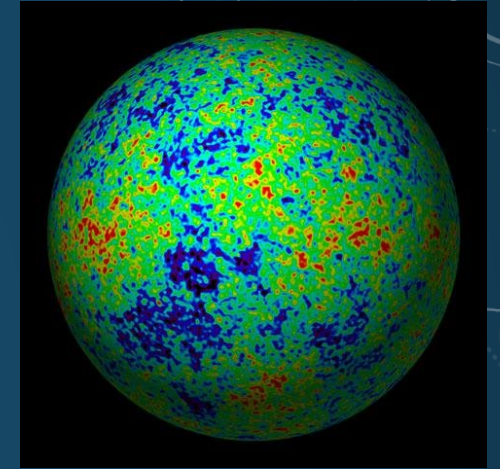
J. Altmann (Monash)





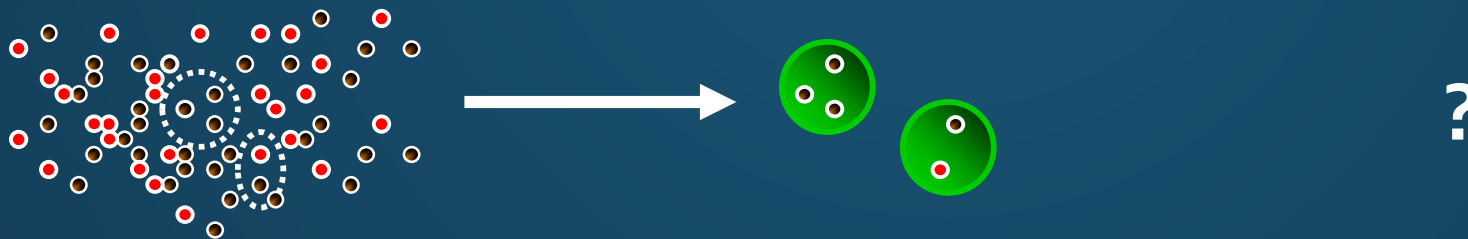
# QUARK RECOMBINATION

- Atomic physics: recombination of protons and electrons into hydrogen + photons



- Photons from the recombination event 300,000 years after the Big Bang.

- Nuclear/particle physics: recombination of quarks into mesons and baryons?

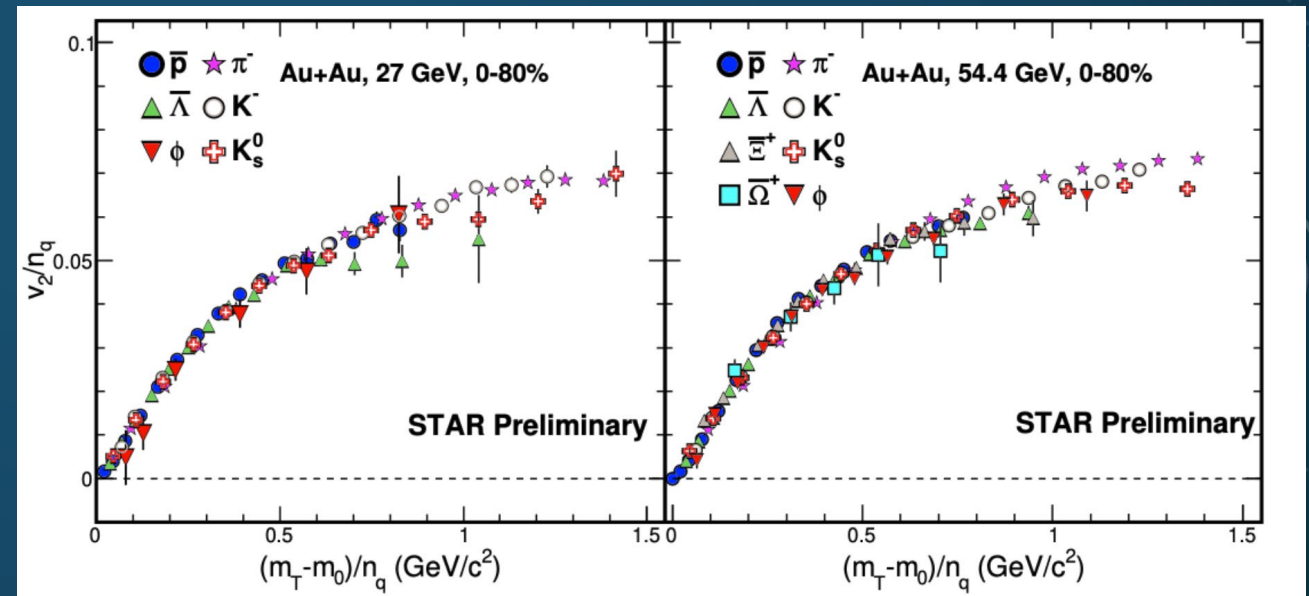
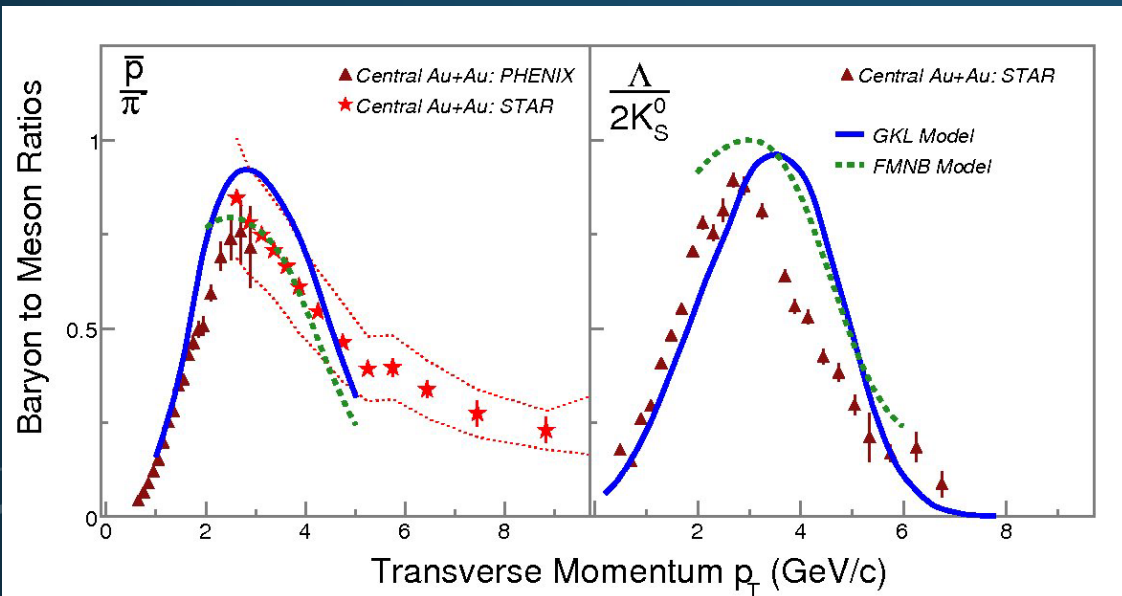


- Complications: gluons, confinement, relativistic dynamics, chiral symmetry breaking,...



# QUARK RECOMBINATION

- Quark recombination is as ancient as QCD, but has been in the shadows of string and cluster hadronization (because data was mostly from  $e^+e^-$ ,  $ep$  and  $pp$  machines!)
- Enter high energy heavy ion data ~ 2000.
- Quark recombination could explain anomalous large baryon numbers (“baryon puzzle”)
- It could also explain the elliptic flow scaling with valence quark number (“universal quark flow is translated to hadrons via recombination”)



# PART III: HYBRID HADRONIZATION

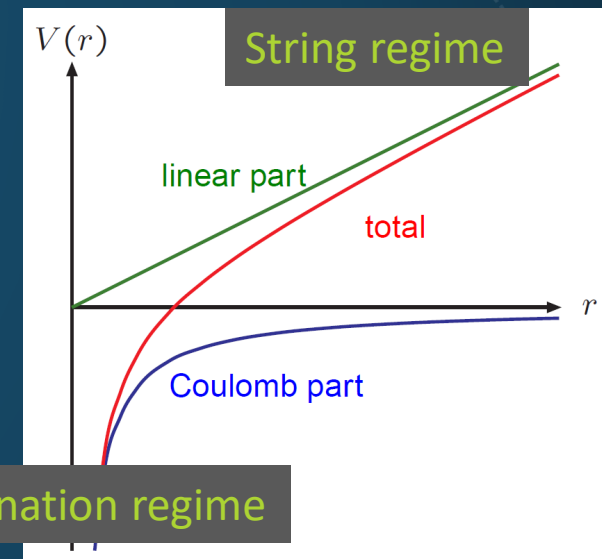
# THE NEED FOR A COMPREHENSIVE MODEL

- We need a model that can work for systems from  $e^+ + e^-$  to  $A + A$ .
- Excluding soft bulk hadronization in  $A + A$  for now (hydro does an excellent job)
- Strings work very well for jet-like systems (partons at large distances in phase space)
- Quark recombination works well if thermal/background partons are involved in large numbers.
- Enter Hybrid Hadronization, a hybrid of string fragmentation and recombination.

K. C. Han, R. J. Fries, C. M. Ko, Jet Fragmentation via Recombination of Parton Showers, Phys.Rev.C 93, 045207 (2016)

- Interpolates smoothly in between, two limits:
  - Dilute systems → Dominance of string fragmentation
  - Dense systems → Dominance of quark recombination

- Use a physics criterion to separate the domains: recombination probabilities vanish for large phase space distances



# HYBRID HADRONIZATION WORK FLOW

## Input:

Provide partons with virtualities below some cutoff, with space-time information and color tags

## Recombination Step:

Provisionally decay gluons into  $q\bar{q}$ . Go through the system sampling the recombination probabilities for all possible  $q$ - $q\bar{q}$  and  $q$ - $q$ - $q$  bound states.

## Intermediate Step:

Recombined hadrons and remnant partons in a string system (only color singlets were removed).

## Fragmentation Step:

Remnant partons tend to be farther apart in phase space. Fragmentation using PYTHIA 8.

# HYBRID HADRONIZATION WORK FLOW IN A MEDIUM

## Input:

Provide partons with virtualities below some cutoff, with space-time information and color tags

Bath of thermal partons

**Recombination Step:**  
Provisionally decay gluons into  $q\bar{q}$ . Go through the system sampling the recombination probabilities for all possible  $q$ - $q\bar{q}$  and  $q$ - $q$ - $q$  bound states.

Recombination with thermal partons

**Intermediate Step:**  
Recombined hadrons and remnant partons in a string system (only color singlets were removed).

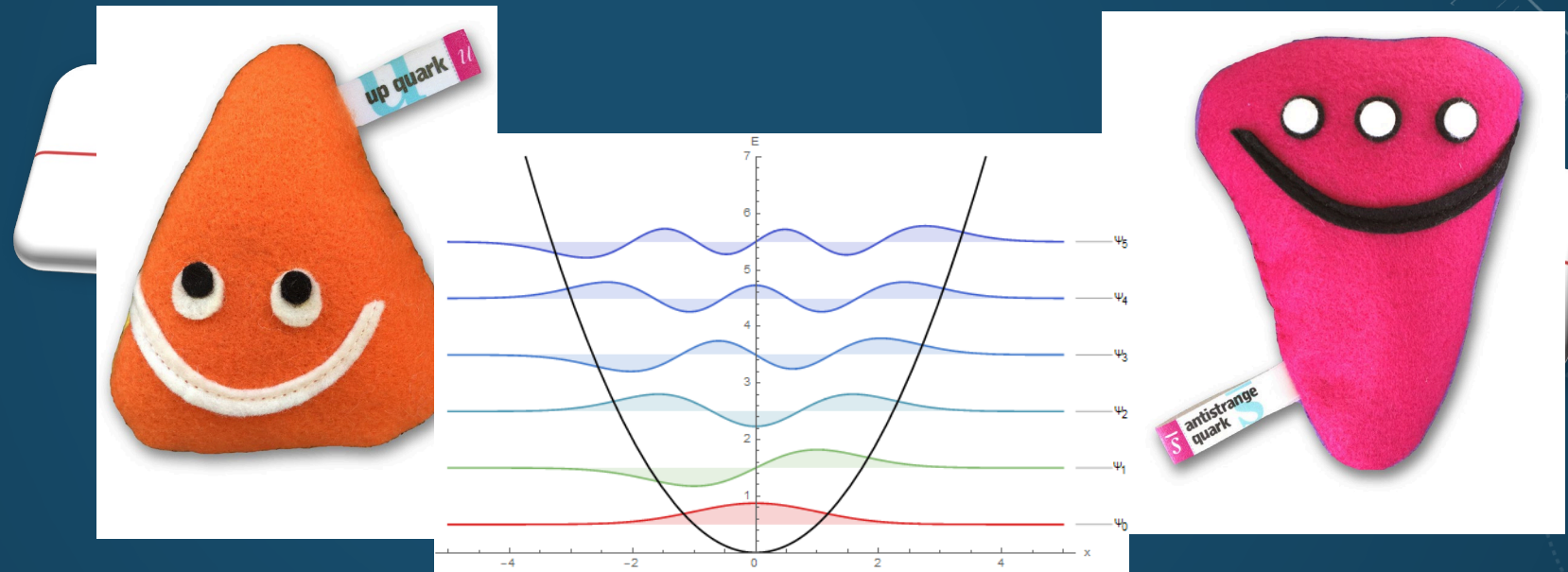
Remnant strings with thermal partons

**Fragmentation Step:**  
Remnant partons tend to be farther apart in phase space. Fragmentation using PYTHIA 8.



# SETTING UP THE RECOMBINATION PROBLEM

- Quarks/antiquarks = wave packets in phase space
- For simplicity: Gaussian wave packets around centroid phase space coordinates  $(\vec{r}_i, \vec{p}_i)$ , of given width  $\delta$ . Color and spin information might be available (otherwise treated statistically).

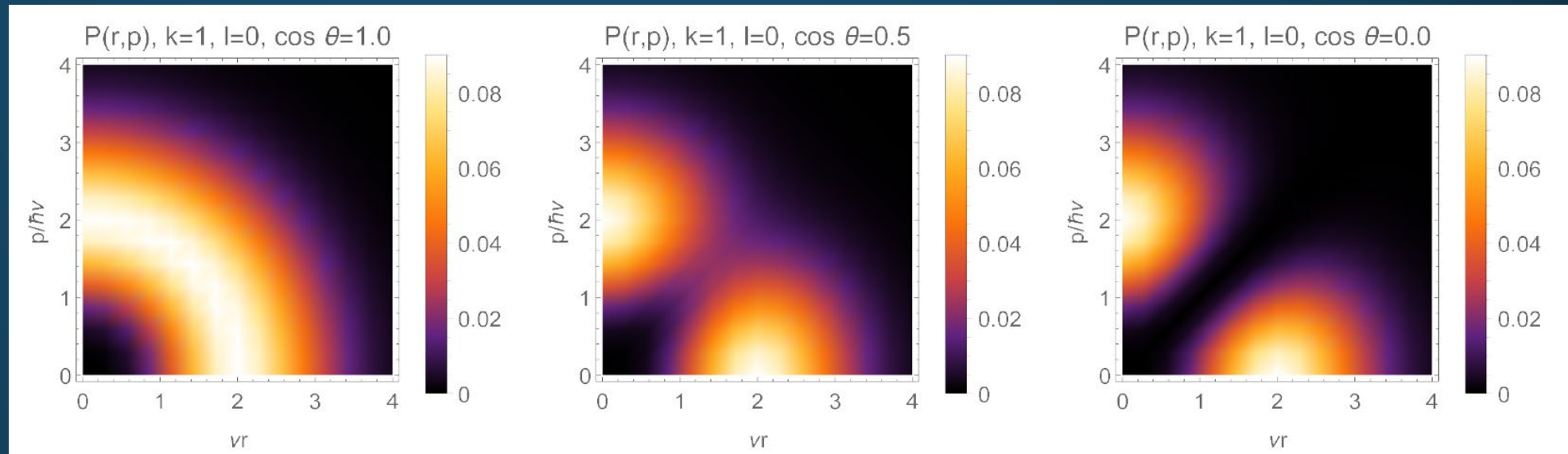
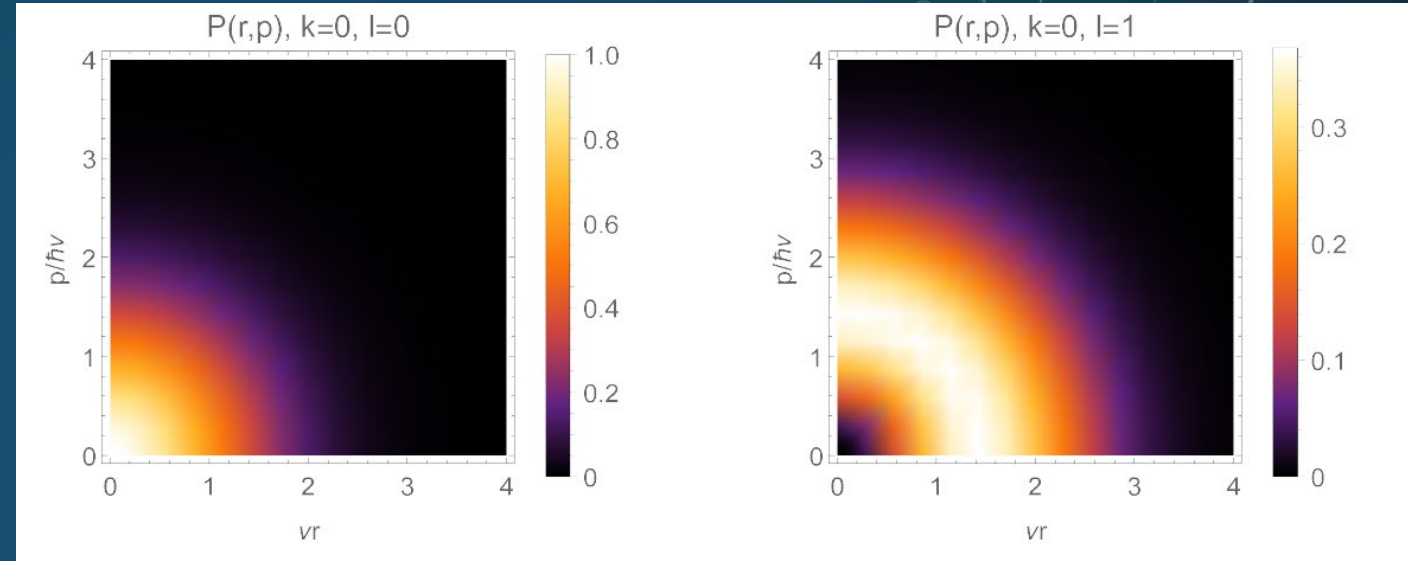


- Short range interaction modeled by isotropic harmonic oscillator potential of width  $1/\nu$ .
- Use the Wigner formalism in phase space. We need angular momentum eigenstates.



# COALESCENCE PROBABILITIES

- Example: Mesons
- We sum over magnetic quantum number  $m$ . (we do not track the spin polarization of hadrons)
- Probabilities depend on the relative coordinates of the wave packet centroids, called  $r$  and  $q$  here.
- $\theta$  = angle between  $r$  and  $q$ .



# COALESCENCE PROBABILITIES

- Probabilities can be written in terms of just two variables: total phase space distance squared  $v$  and total angular momentum squared  $t$ .

$$v = \frac{v^2 r^2}{2} + \frac{p^2}{2\hbar^2 v^2},$$

$$t = \frac{1}{\hbar^2} [p^2 r^2 - (\mathbf{p} \cdot \mathbf{r})^2] = \frac{1}{\hbar^2} L^2$$

$$\mathcal{P}_{00} = e^{-v},$$

$$\mathcal{P}_{01} = e^{-v} v,$$

$$\mathcal{P}_{02} = \frac{1}{2} e^{-v} \left( \frac{2}{3} v^2 + \frac{1}{3} t \right)$$

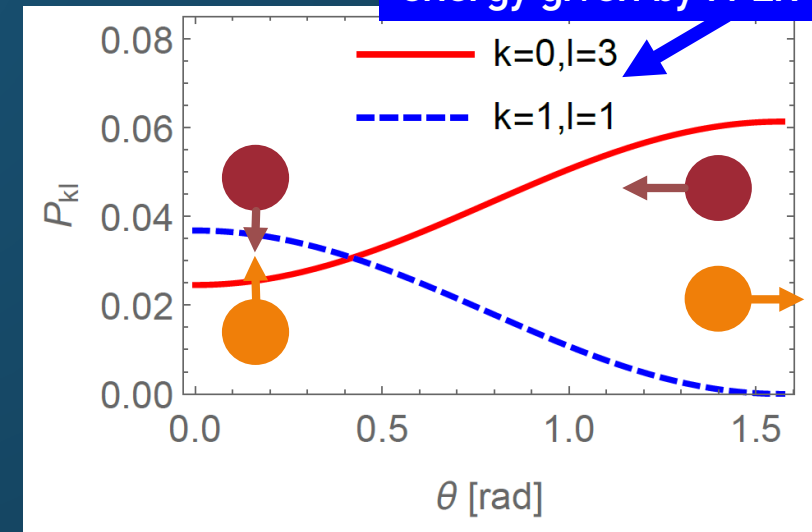
$$\mathcal{P}_{10} = \frac{1}{2} e^{-v} \left( \frac{1}{3} v^2 - \frac{1}{3} t \right)$$

- $t$  makes an intuitive connection between the relative angular momentum of the incoming quarks and the quantum number  $l$  of their bound state.

- The total recombination probability also takes into account quark spins (always statistically) and quark color.

- Color factors are determined by color tags. Thermal partons and shower partons with random color are initialized with color tag 0.

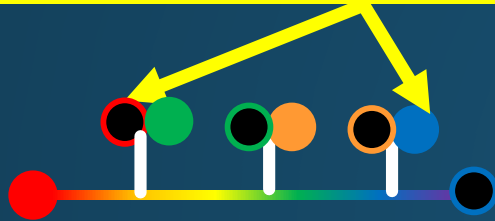
Both are states with the same energy given by  $N=2k+l=3$



# REMNANT PARTONS: STRING REPAIR

- Recombination only removes color singlets. Remaining strings “snap together” the right way automatically.

Suppose these two partons recombine.

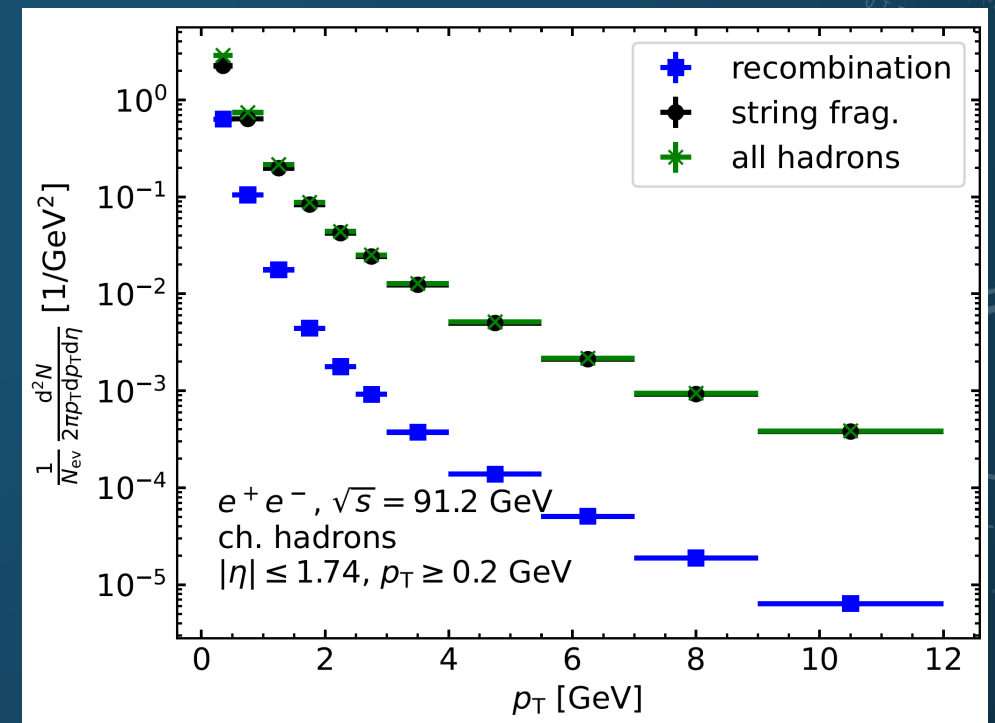


- Remnant partons with color tag 0 (e.g. from LBT) must be introduced into strings; unused gluons are restored.
- If the initial system was not a color singlet extra partons must be introduced to balance color (this could be beam partons, thermal partons, or fake partons with zero momentum).

# REMNANT STRINGS: FRAGMENTATION

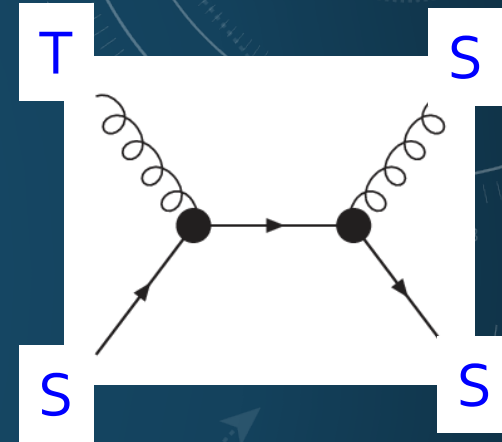
- Hand remaining string systems to PYTHIA 8 for fragmentation.
- Hadronic resonances are presumably important channels in hadronization (both through string fragmentation and recombination). In JETSCAPE decays of excited states can happen in PYTHIA or by invoking SMASH.

- Remnant partons tend to be farther removed from each other in phase space.
- $e^+e^-$  example:
- As intended fragmentation dominates this dilute system, in particular for the jet core.



# ADDING A MEDIUM

- The formalism stays the same, just include the additional partons!
- Some shower partons (e.g. LBT) arrive with randomized color (color tag 0)
- HH can process “negative partons” separately. They are the “holes” left in the medium through processes like  $q$  (shower) +  $g$  (medium)  $\rightarrow$   $q$  (shower) +  $g$  (shower). Important for background subtraction.
- Thermal partons can be added, or will be sampled from a  $T = T_c$  hypersurface or a brick.
- Recombination from only thermal partons, or strings with only thermal partons are currently disabled. On the flip side, all shower partons are always hadronized.



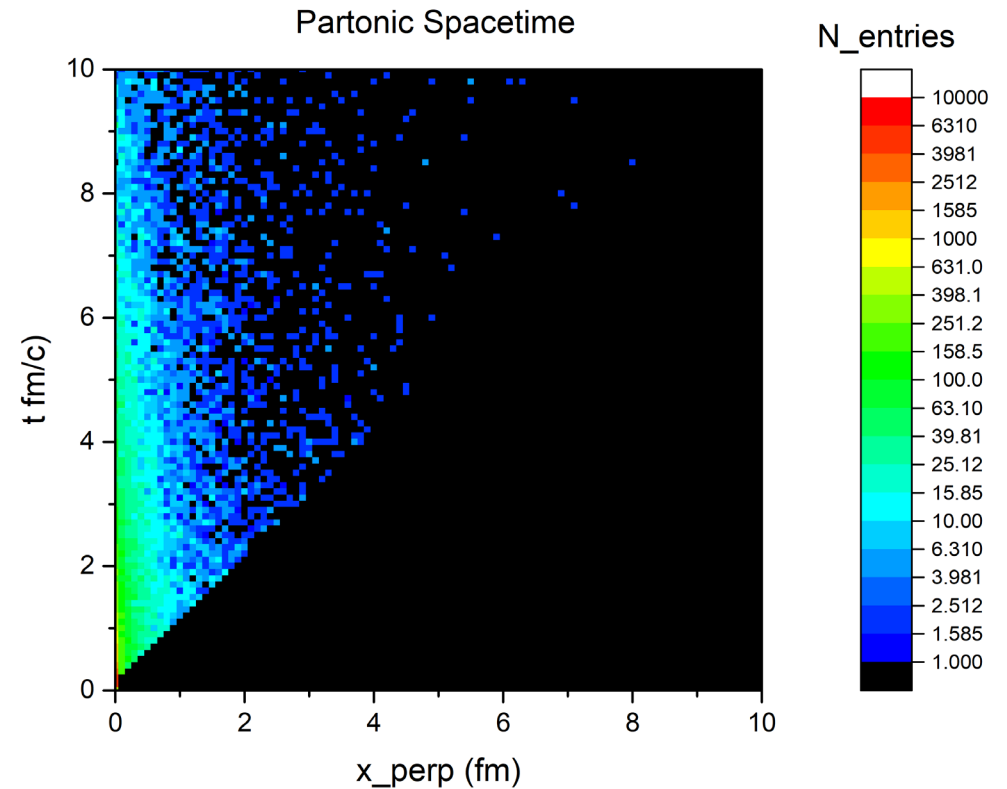


# IN-MEDIUM JETS: SPACE-TIME CONSIDERATIONS

- Sampled spatial positions of shower partons after shower evolution for 100 GeV jets (arb. normalization)
- Here: JETSCAPE:pGun+MATTER

100 GeV vacuum jet

Time since jet was created



Direction transverse to the jet

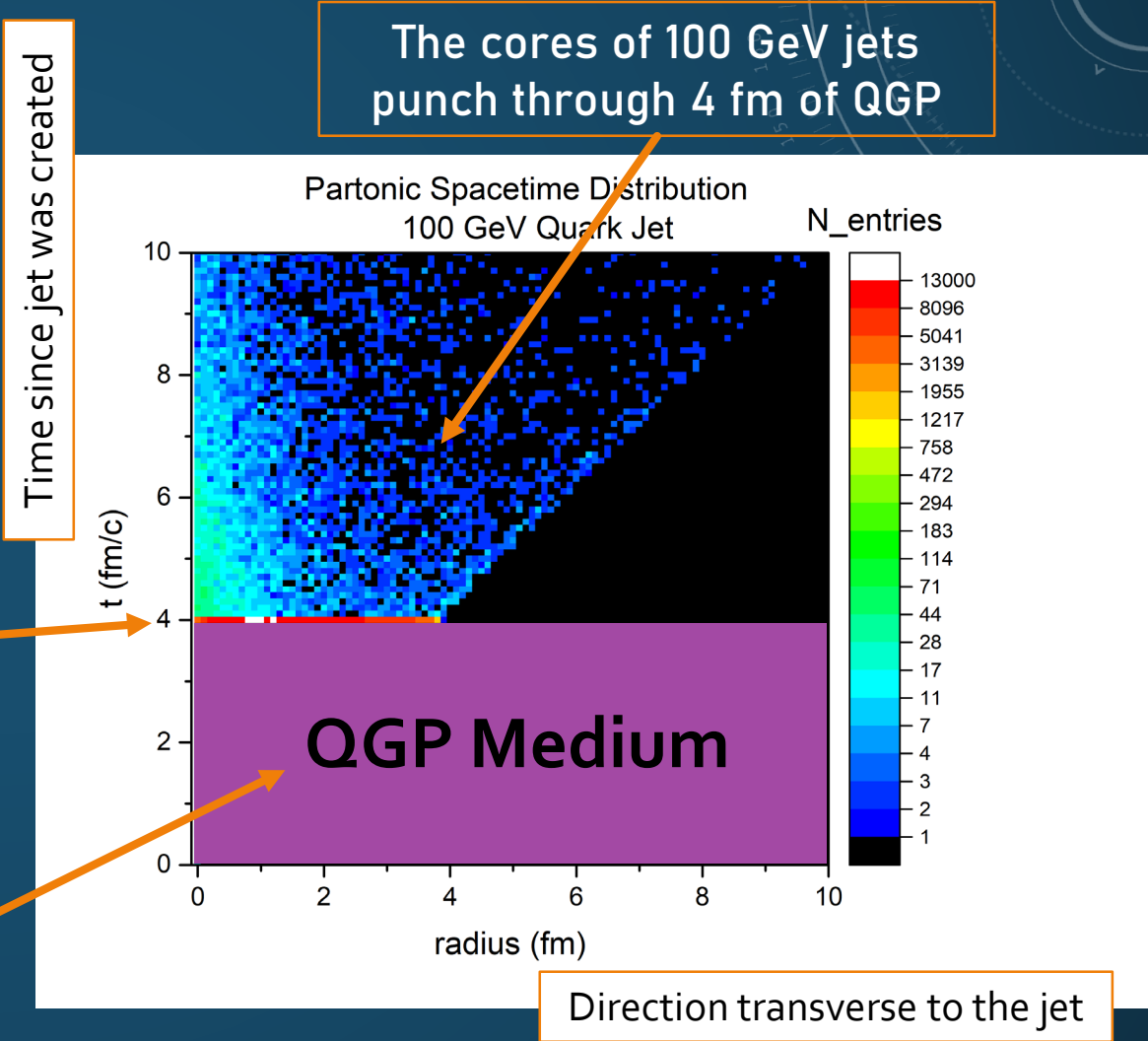


# IN-MEDIUM JETS: SPACE-TIME CONSIDERATIONS

- Sampled spatial positions of shower partons after shower evolution for 100 GeV jets (arb. normalization)
- Here:  
JETSCAPE:pGun+MATTER+LBT+Brick

Shower partons inside QGP are absorbed by the medium or accumulate on the hypersurface; color is randomized

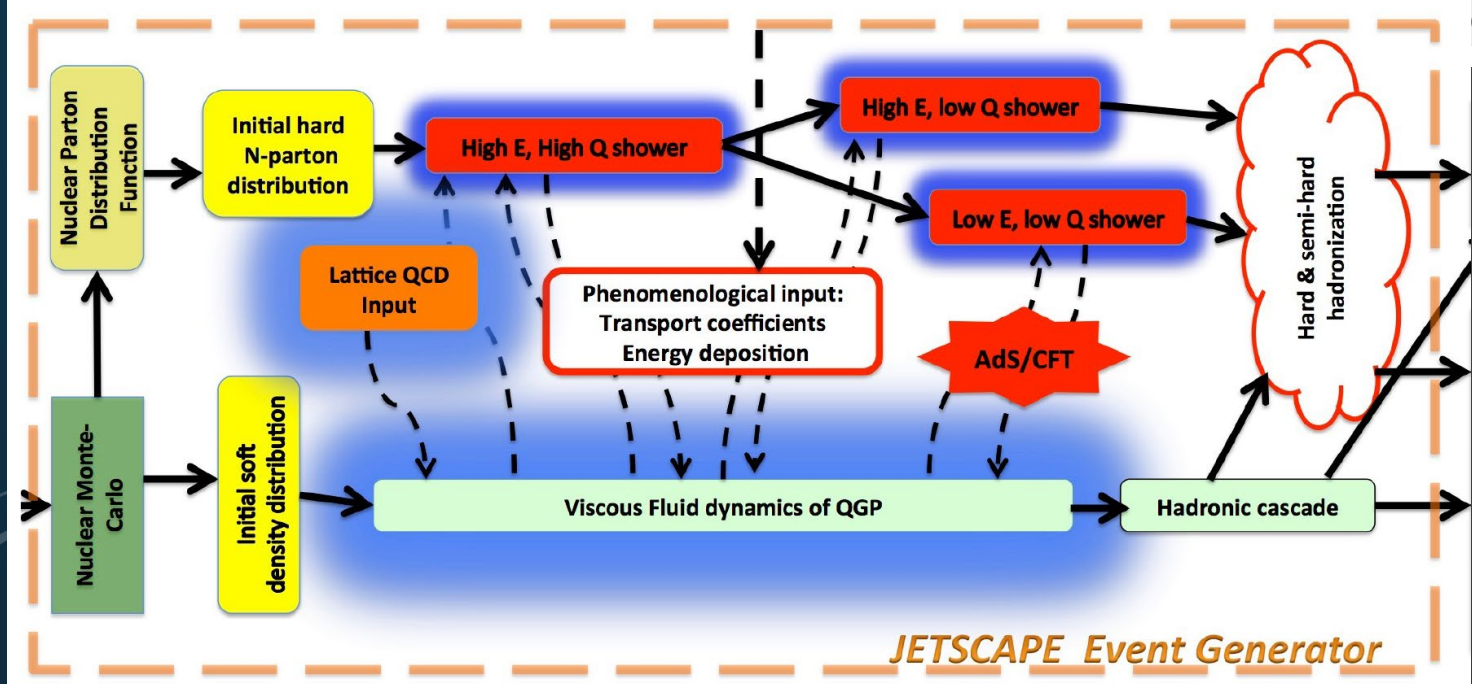
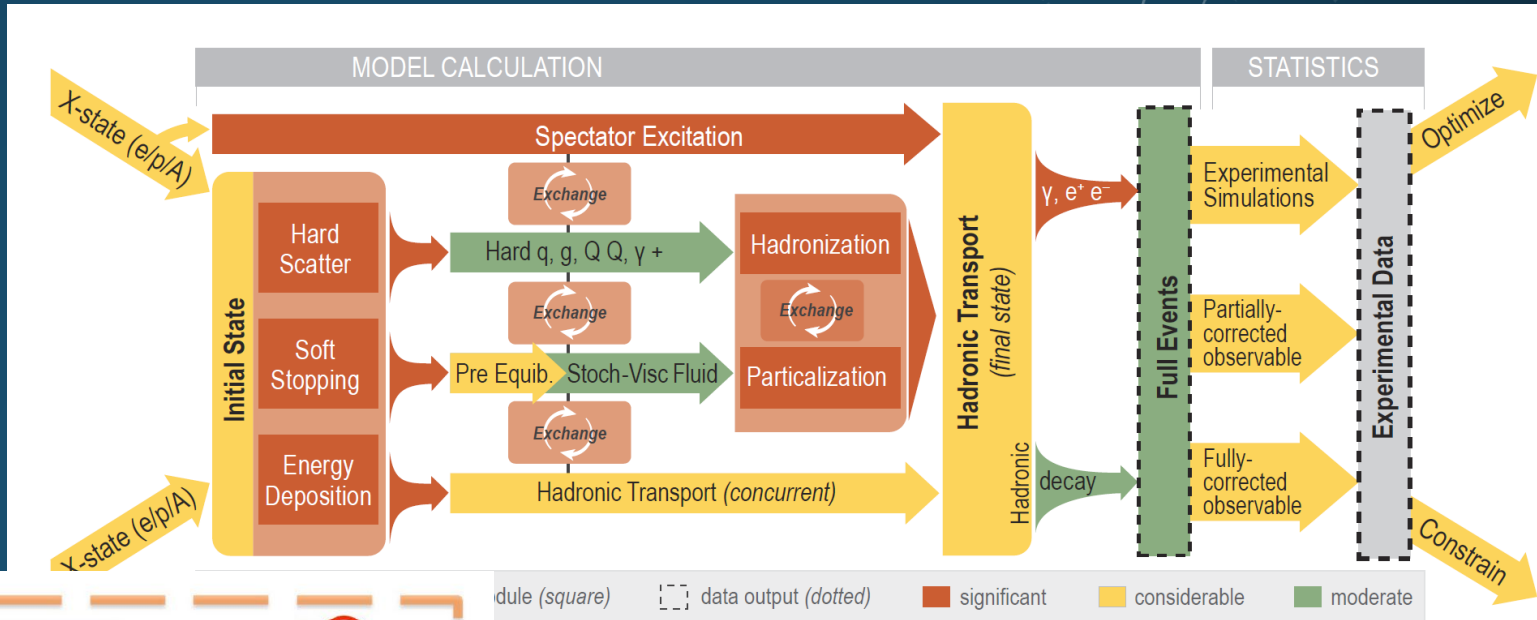
The jet starts in QGP; the temperature is set to drop below  $T_c$  after 4 fm/c





# PART IV: HADRONIZATION IN JETSCAPE

# THE PLACE OF HADRONIZATION IN JETSCAPE/XSCAPE



# JETSCAPE/XSCAPE OPTIONS

- There are three options for hadronization: colored, colorless, and hybrid hadronization.
- Colored and colorless hadronization are both pure PYTHIA 8 string fragmentation but with different implementations (honoring or discarding color tags). They can not add medium partons.
- In a nutshell:

	Uses color tags	Can deal with missing color tags	Can use or generate thermal partons	Can handle junctions/baryon
Colored	YES	NO	NO	NO
Colorless	NO	YES	NO	NO
Hybrid	YES	YES	YES	YES

- Colored hadronization is only recommended for vacuum systems, like p+p.

# JETSCAPE/XSCAPE OPTIONS

- Set the module you want in `<JetHadronization>` in the XML-file together with any parameters you want to customize.

```
<!-- Jet Hadronization Module -->  
<JetHadronization>  
  <name>colored/colorless/hybrid</name>
```

- You will play mostly with Hybrid Hadronization and a little bit with Colorless Hadronization later.
- Caution, the version of HH used at the school is not yet official, it will be in the JETSCAPE 3.6 release

# HYBRID HADRONIZATION OPTIONS

- A few options for HH:

```
<!-- Jet Hadronization Module -->
<JetHadronization>
  <name>hybrid</name>
  <!--eCMforHadronization only for pp collisions-->
  <!--in hybrid put the full eCM here, this is distributed to the beam partons-->
  <eCMforHadronization>0.0</eCMforHadronization>
  <had_postprop>0.0</had_postprop>
  <part_prop>1.0</part_prop>
  <pythia_decays>on</pythia_decays> <!-- lets the particles given to pythia decay-->
  <tau0Max>10.0</tau0Max> <!-- only particles with tau0 < tau0Max (given in mm/c) can decay, increase to include weak decays-->
  <reco_hadrons_in_pythia>1</reco_hadrons_in_pythia>
</JetHadronization>
```

Beam partons if extra partons are needed.

No need to touch these except when doing serious tuning.

Hadron decay settings for PYTHIA

```
<reco_Mlevelmax>1</reco_Mlevelmax>
<reco_Blevelmax>1</reco_Blevelmax>
<reco_goldstone>0</reco_goldstone>
```

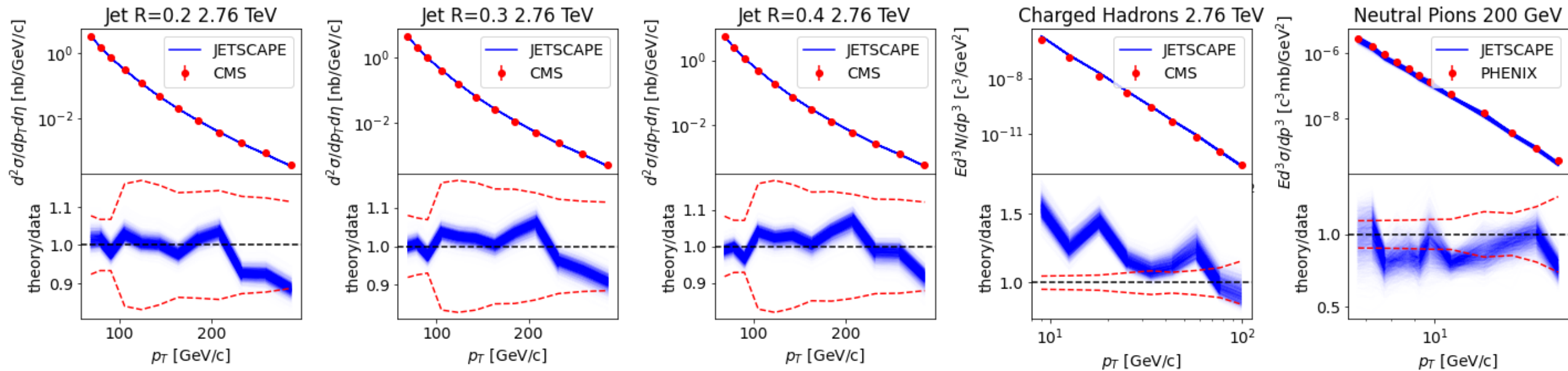
Highest excited energy level for mesons and baryons in recombination. (0=ground state; 1=p-wave allowed)

Pions and kaons have anomalously small masses and don't fit well into the recombination scheme.



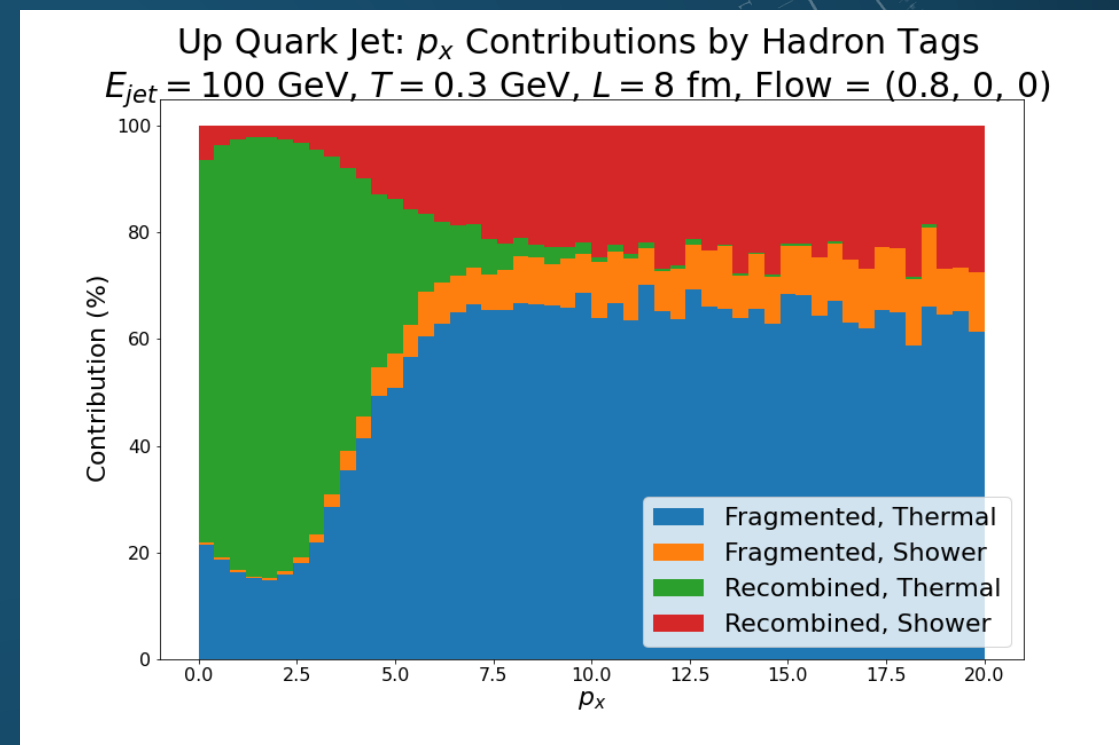
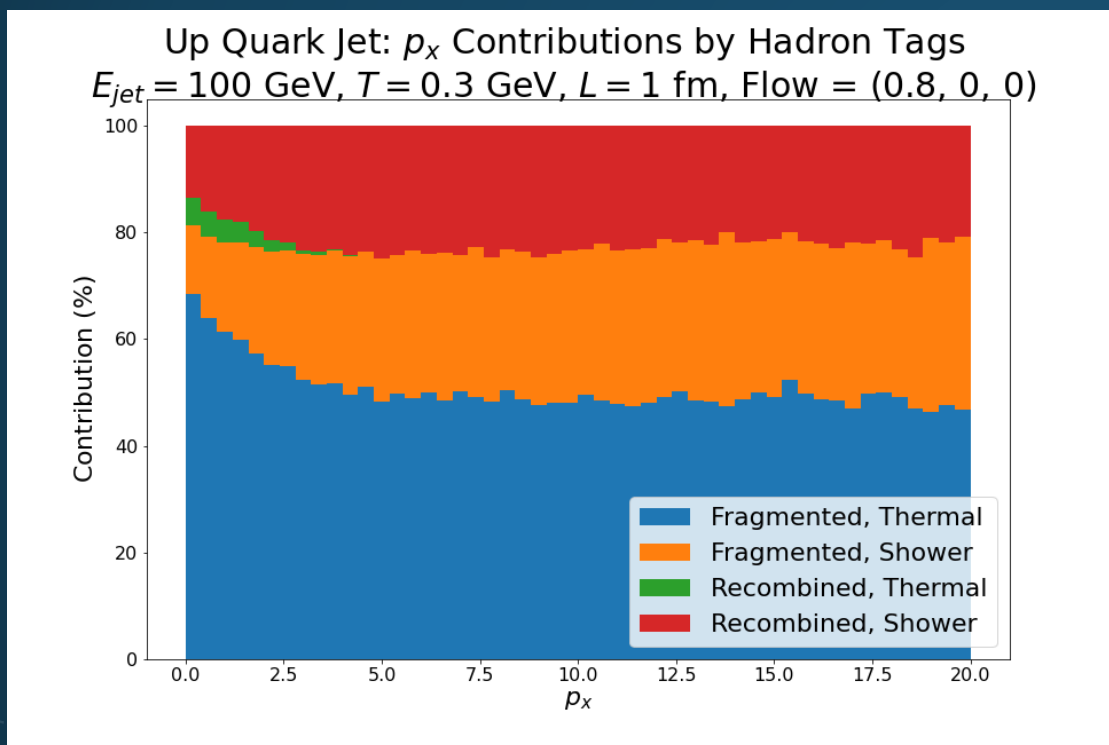
# A FEW RESULTS

- Most of these are obtained with the older, official version of Hybrid Hadronization.
- For vacuum systems, HH is designed to work as well as string fragmentation, with little changes through the addition of recombination. This mostly works out.



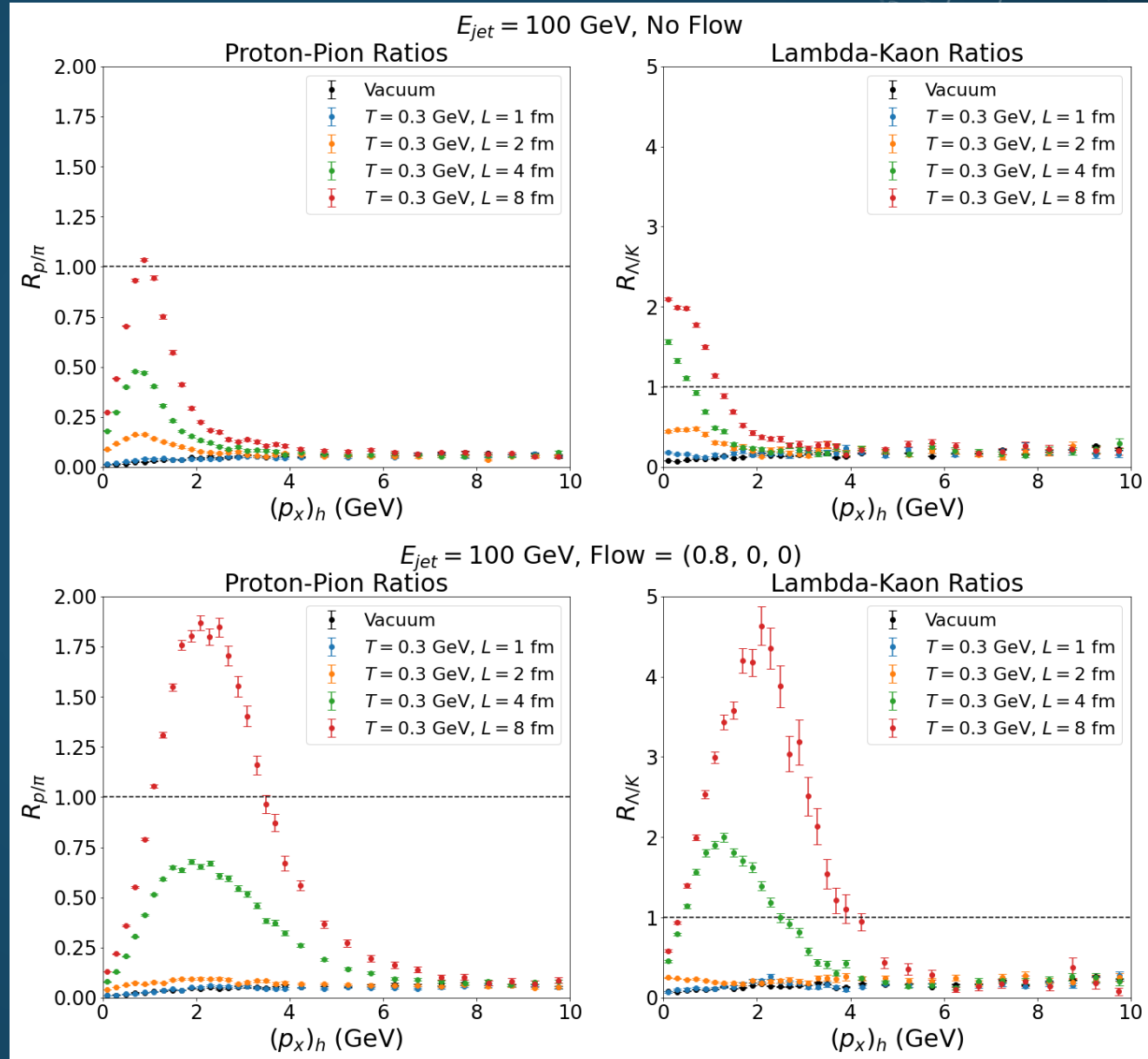
# A FEW RESULTS

- In systems with a medium recombination, with medium partons becomes a strong contribution at low to intermediate momentum, growing with medium size.



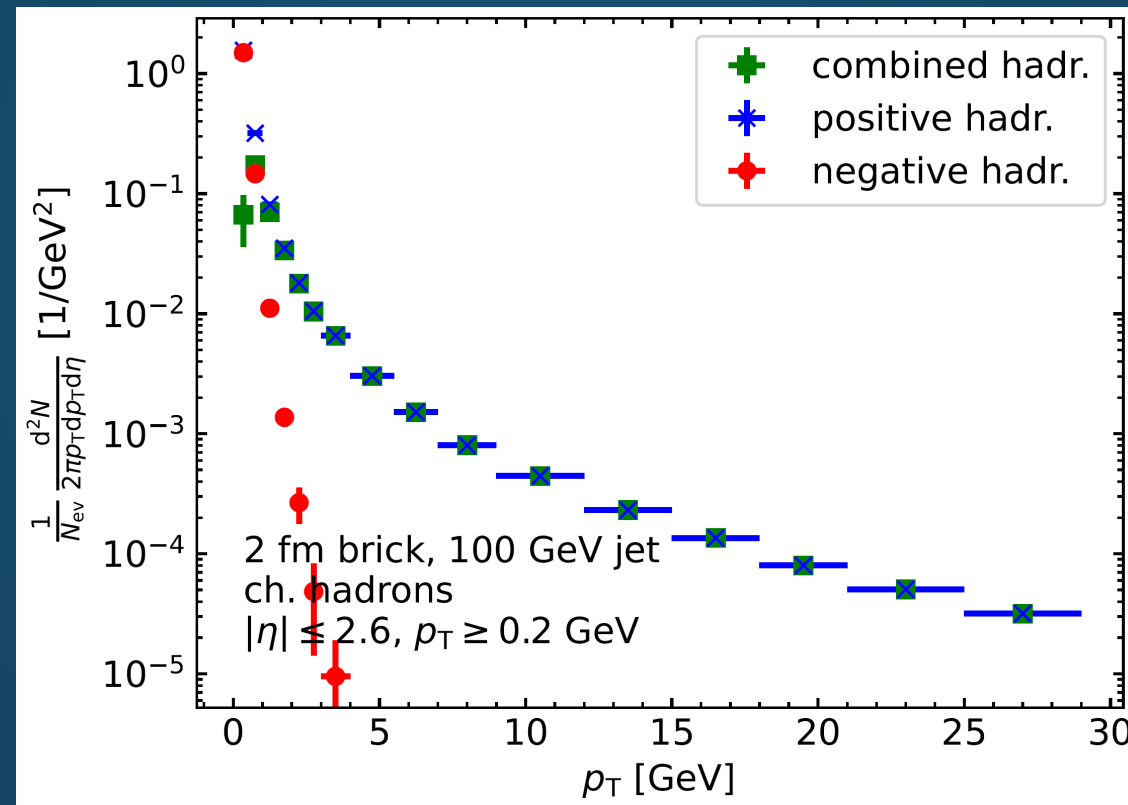
# A FEW RESULTS

- We recover a key signature of quark recombination: baryon/meson enhancement in a medium
- Hadronization is sensitive to medium flow.



# A FEW RESULTS

- This is a result with the new version: the background of negative hadrons can be successfully subtracted.



# ON TO THE HANDS-ON SESSION

- A few exercises running different models for  $e^+e^-$  and a jet in a brick.