



Transfer reactions to the continuum involving exotic nuclei

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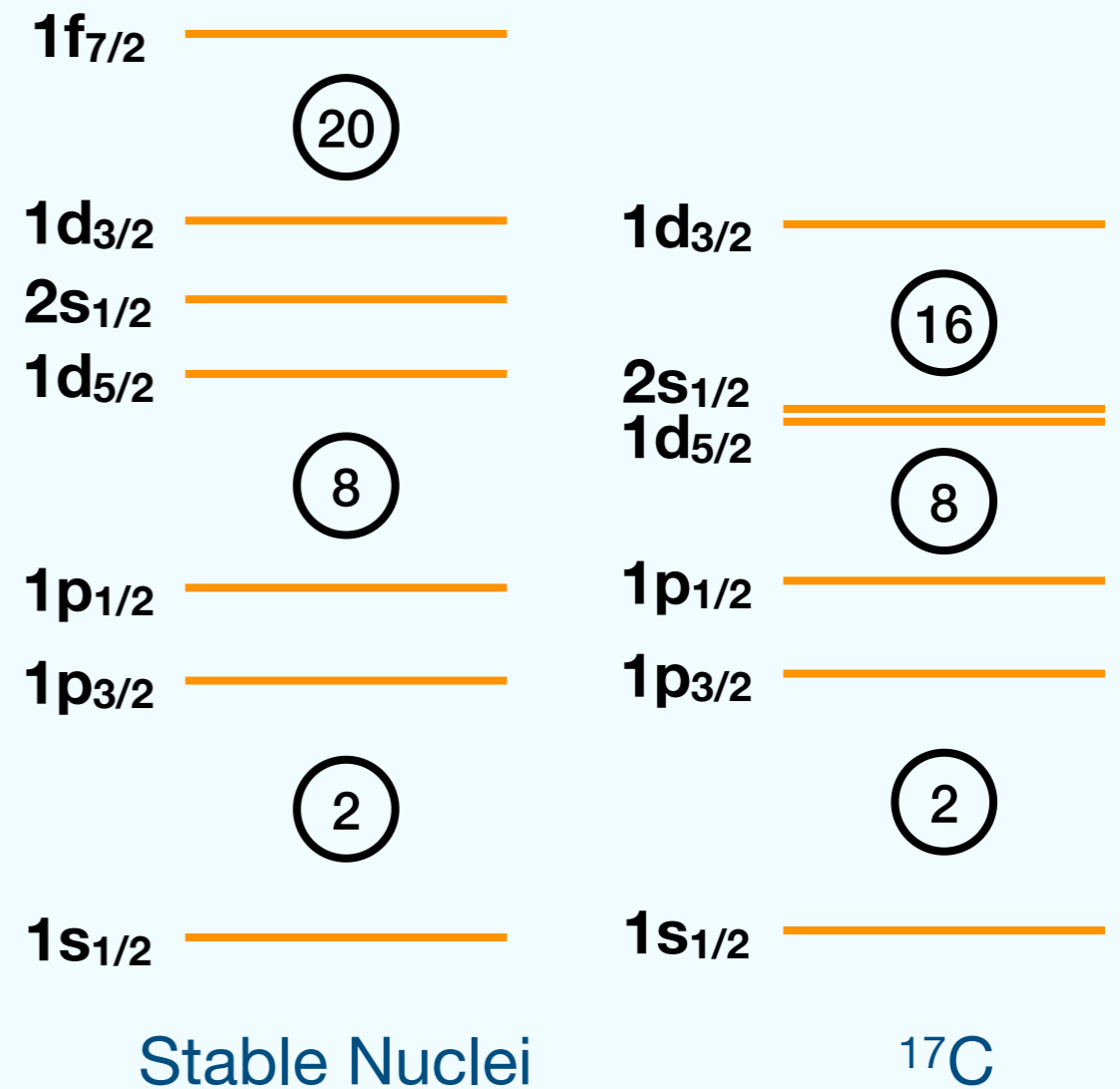
Introduction

- Exotic nuclei
Different properties
 - Shell closure
 - Halo nature
- Transfer reactions with exotic nuclei
Advances in radioactive beam facilities
- Few-body models
Significant effect of deformation
- ^{11}Be and ^{17}C
deformed core + neutron structure

										12O	13O	14O	15O	16O	17O	18O	19O	20O	21O			
										10N	11N	12N	13N	14N	15N	16N	17N	18N	19N	20N		
										8C	9C	10C	11C	12C	13C	14C	15C	16C	17C	18C	19C	
										6B	7B	8B	9B	10B	11B	12B	13B	14B	15B	16B	17B	18B
										5Be	6Be	7Be	8Be	9Be	10Be	11Be	12Be	13Be	14Be	15Be	16Be	
										3Li	4Li	5Li	6Li	7Li	8Li	9Li	10Li	11Li	12Li	13Li		
											3He	4He	5He	6He	7He	8He	9He	10He				
											1H	2H	3H	4H	5H	6H	7H					
												Neutron										

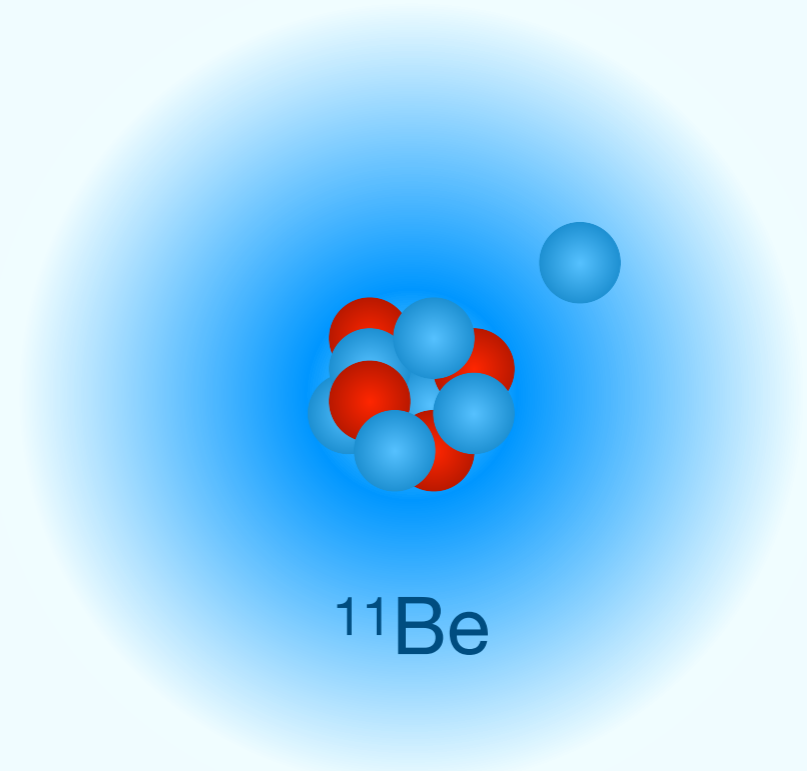
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Structure models

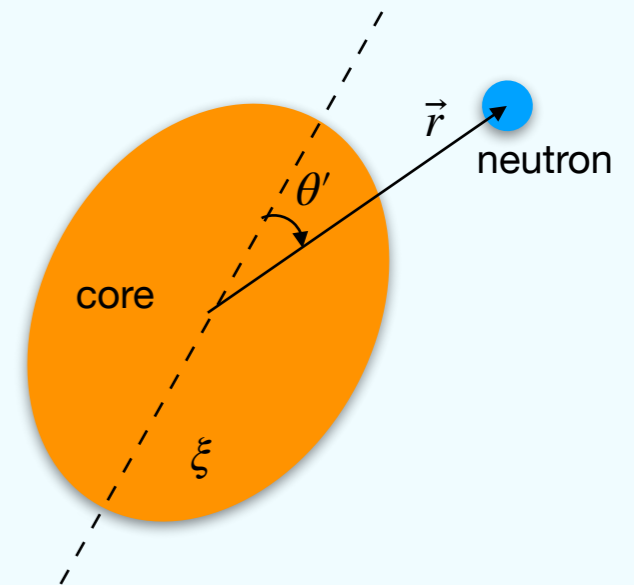
- Deformed 2-body models (neutron+core)

- Nilsson model

$$H = -\frac{\hbar^2}{2\mu}\nabla^2 + V_{ls}(r)(\vec{l} \cdot \vec{s}) + V_c(r) - r\beta \frac{dV_c(r)}{dr} Y_{20}(\theta') + \frac{\hbar^2}{2\mathcal{J}} \vec{I}^2$$

- PAMD [PRC89 (2014) 014333]

$$H = -\frac{\hbar^2}{2\mu}\nabla^2 + V_{ls}(r)(\vec{l} \cdot \vec{s}) + V_{vc}^{AMD}(\vec{r}, \xi) + h_{core}(\xi)$$

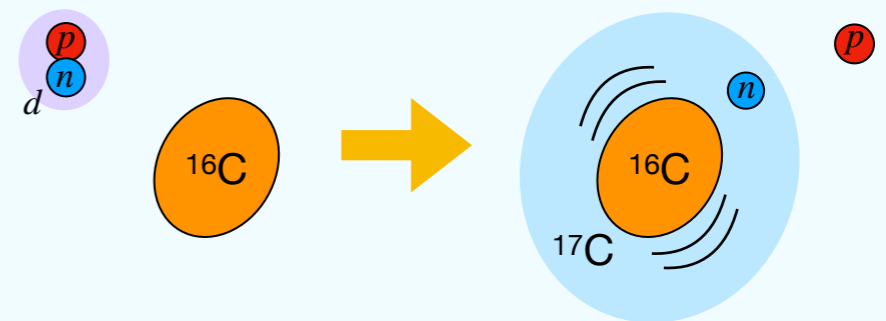
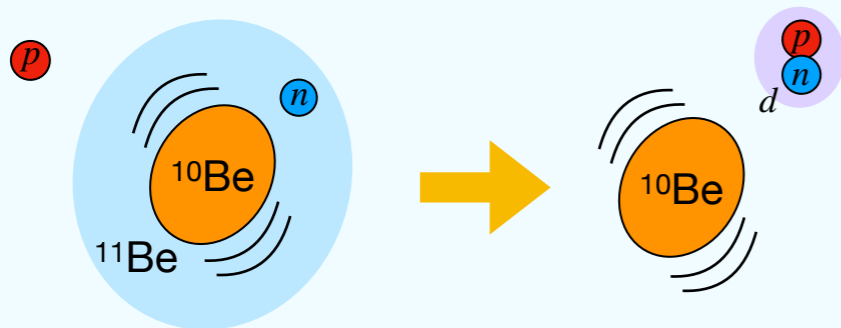


- Eigenstates from diagonalization in THO basis

$$\phi_{nl}^{THO}(r) = \sqrt{\frac{ds}{dr}} \phi_{nl}^{HO}[s(r)]$$

Neutron transfer reactions

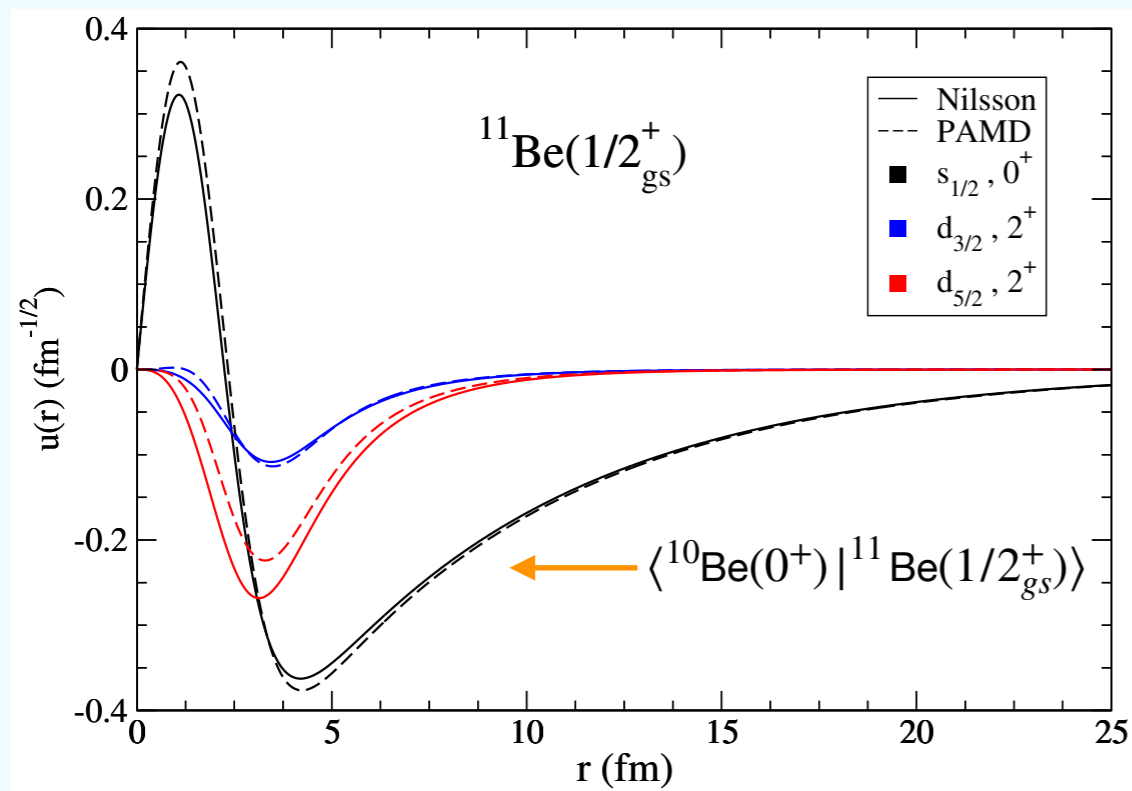
- (p,d) and (d,p) reactions are studied applying the Adiabatic Distorted Wave Approximation (ADWA).



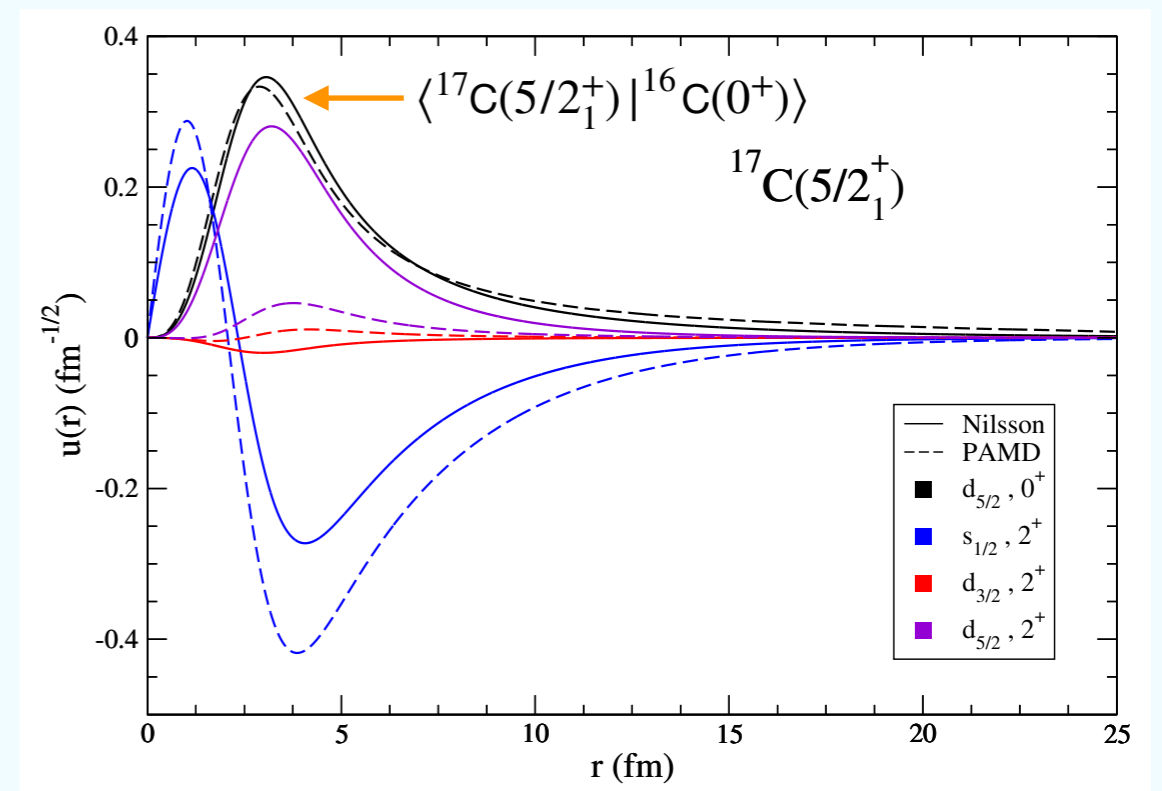
Neutron transfer reactions

- (p,d) and (d,p) reactions are studied applying the Adiabatic Distorted Wave Approximation (ADWA).
- These calculations require the $\langle {}^{10}\text{Be} | {}^{11}\text{Be} \rangle$ and $\langle {}^{17}\text{C} | {}^{16}\text{C} \rangle$ overlaps.

${}^{11}\text{Be}(p,d){}^{10}\text{Be}$



${}^{16}\text{C}(d,p){}^{17}\text{C}$



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- These calculations require the $\langle {}^{10}\text{Be} | {}^{11}\text{Be} \rangle$ and $\langle {}^{17}\text{C} | {}^{16}\text{C} \rangle$ overlaps.
- Calculations are compared with recent experimental data.

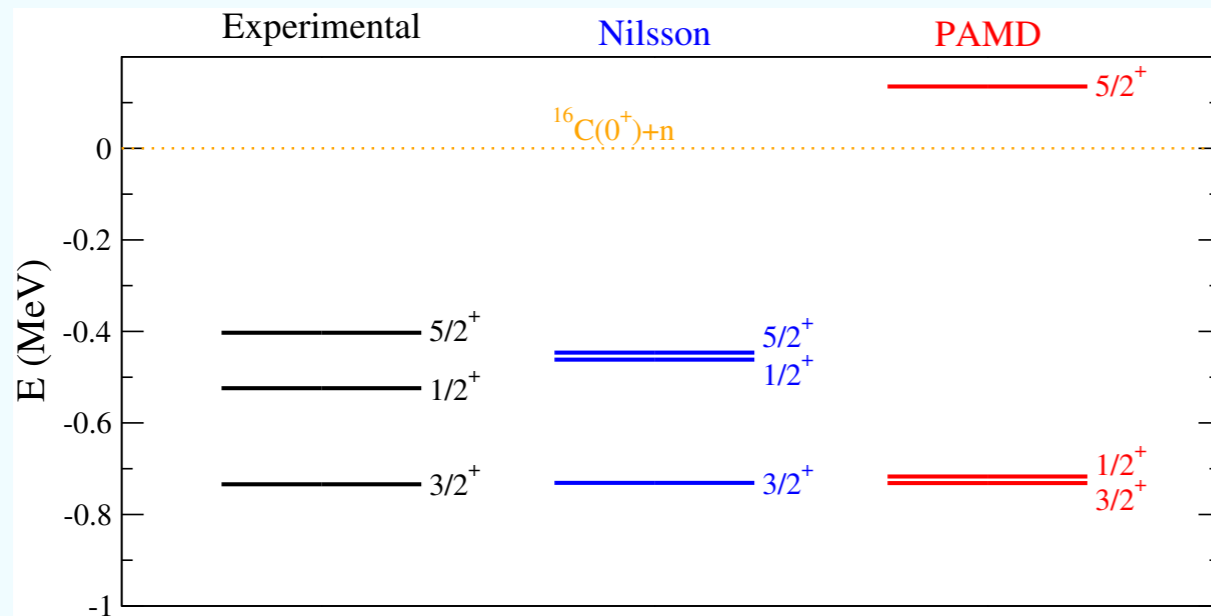


- GANIL, 35.3 MeV/nucleon beam
[NPA**683** (2001) 48]
- RCNP, 26.9 MeV/nucleon beam
[Chinese P. L. **35** (2018) 082501]

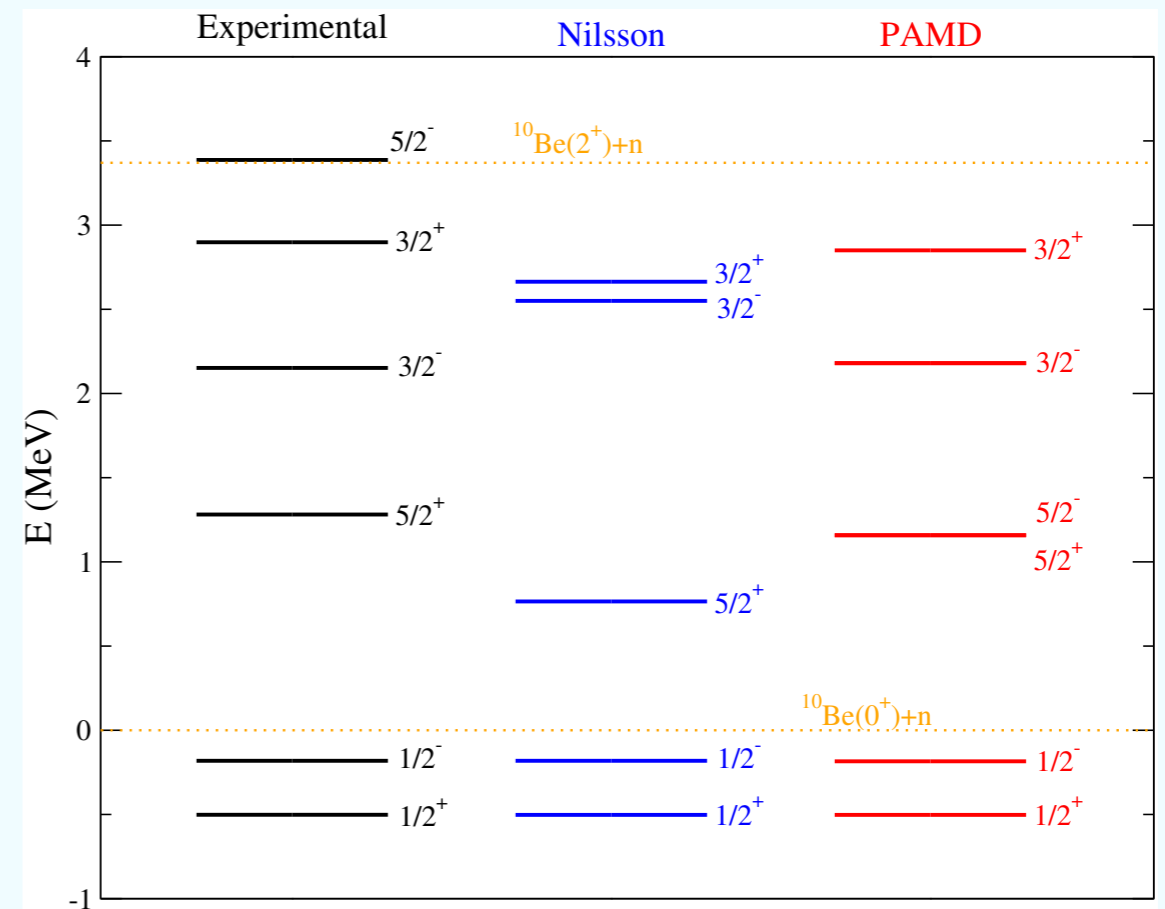
- GANIL, 17.2 MeV/nucleon beam
[PLB**811** (2020) 135939]

Application to ^{17}C and ^{11}Be

^{17}C

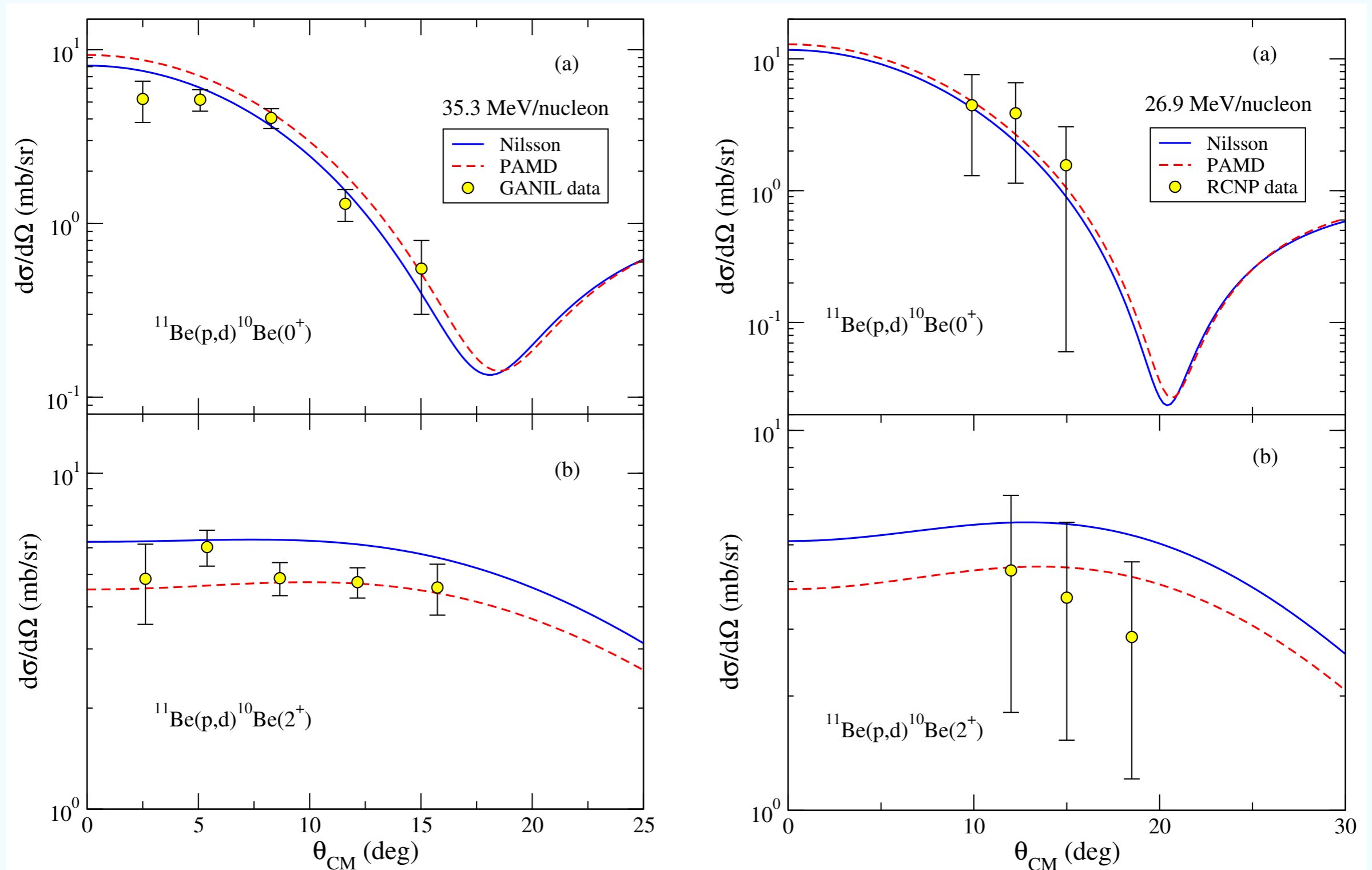


^{11}Be

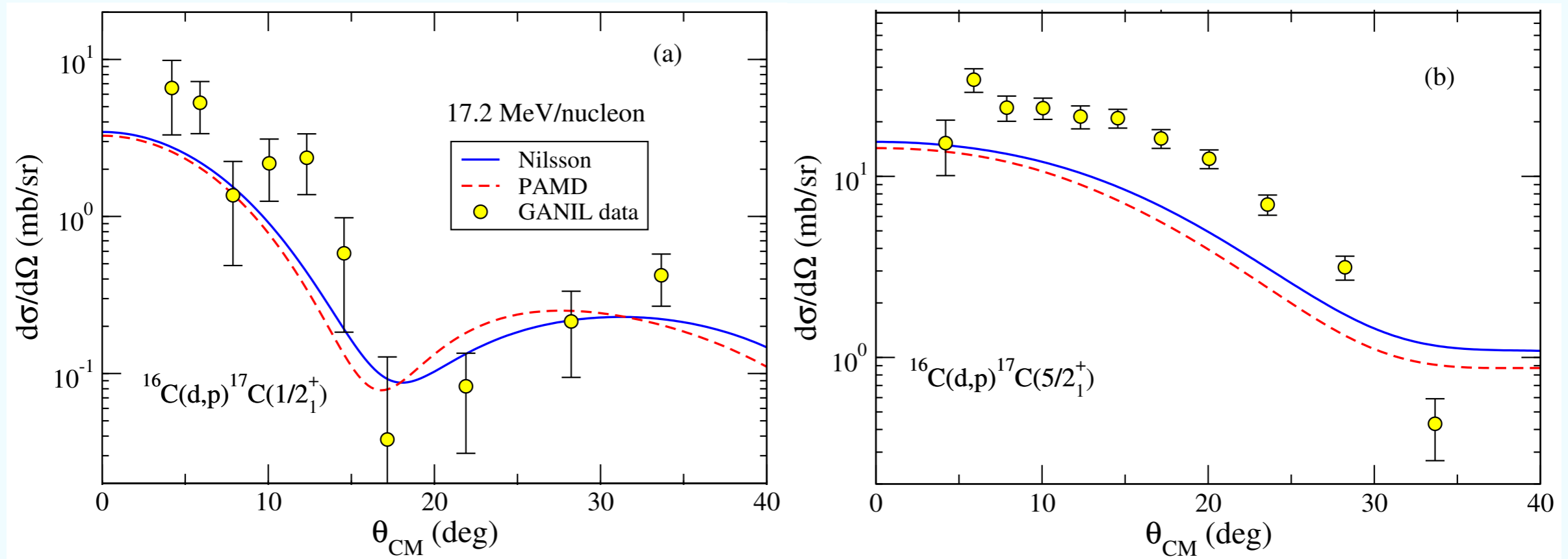


P. Punta, *et al.* PRC**108** (2023) 024613

Application to $^{11}\text{Be}(p,d)^{10}\text{Be}$



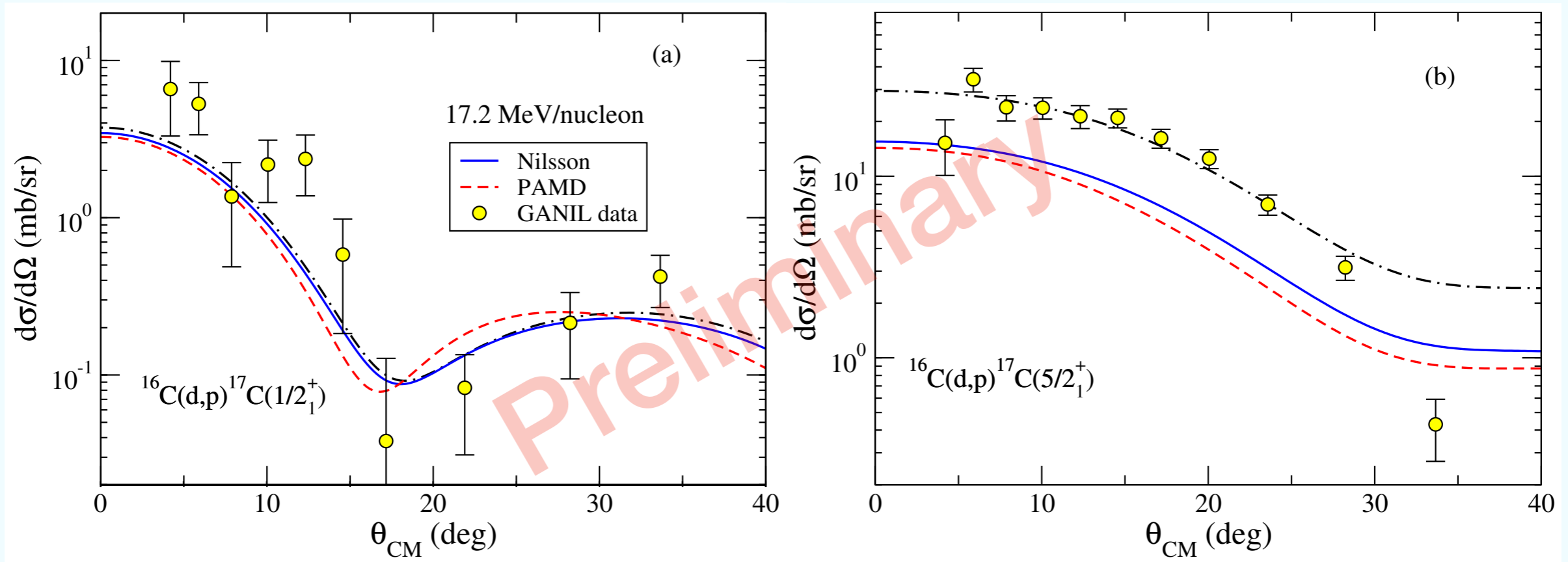
Application to $^{16}\text{C}(d,p)^{17}\text{C}$



Comparison with data suggest that:

- $\langle ^{17}\text{C}(1/2_1^+) | ^{16}\text{C}(0^+) \rangle$ overlaps are realistic in our models.
- $\langle ^{17}\text{C}(5/2_1^+) | ^{16}\text{C}(0^+) \rangle$ too small, suggesting a larger spectroscopic factor.

Application to $^{16}\text{C}(d,p)^{17}\text{C}$



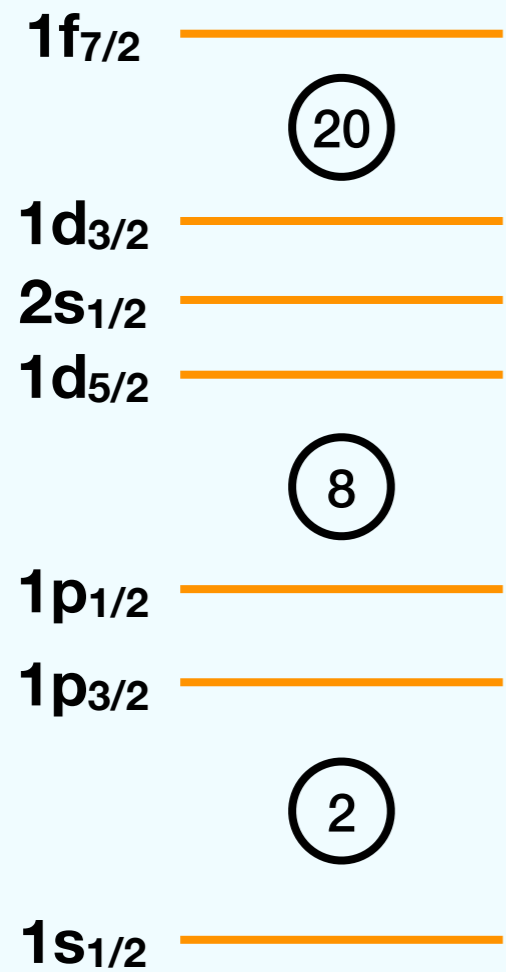
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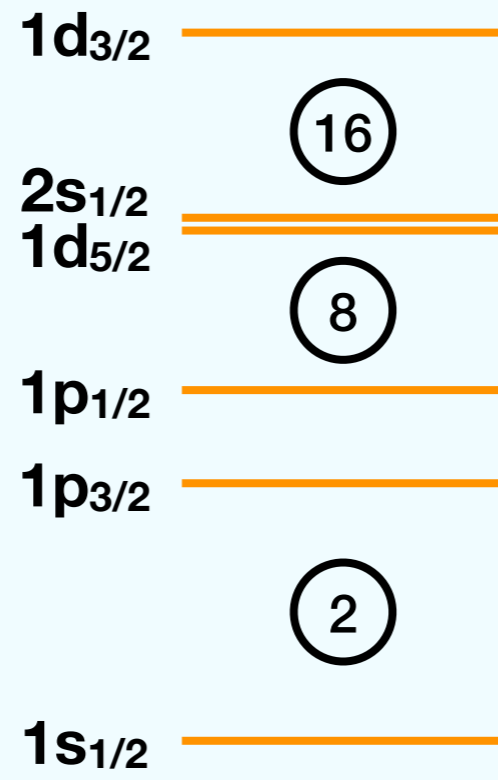


Pauli blocking

N=16 Shell-Gap

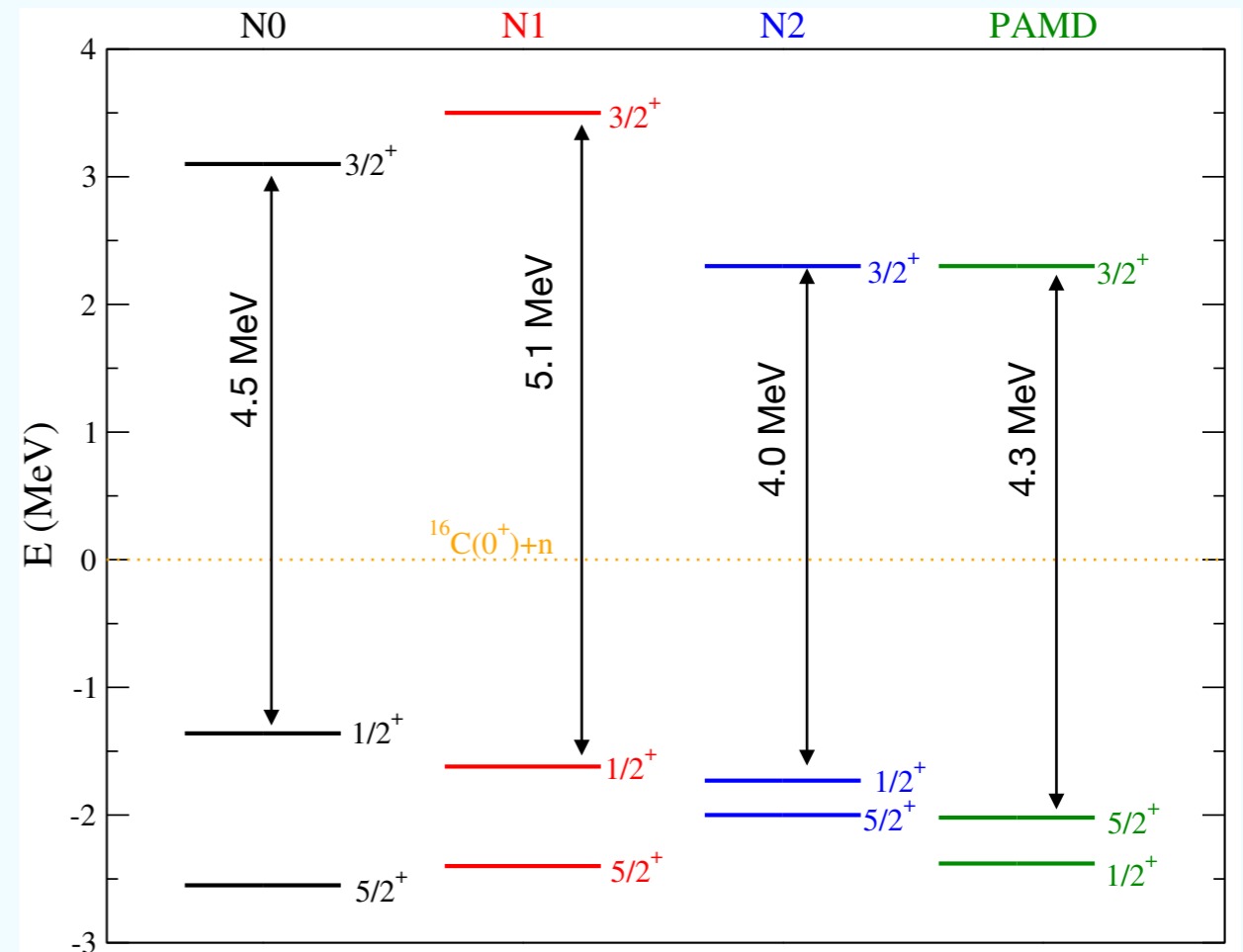


Stable Nuclei

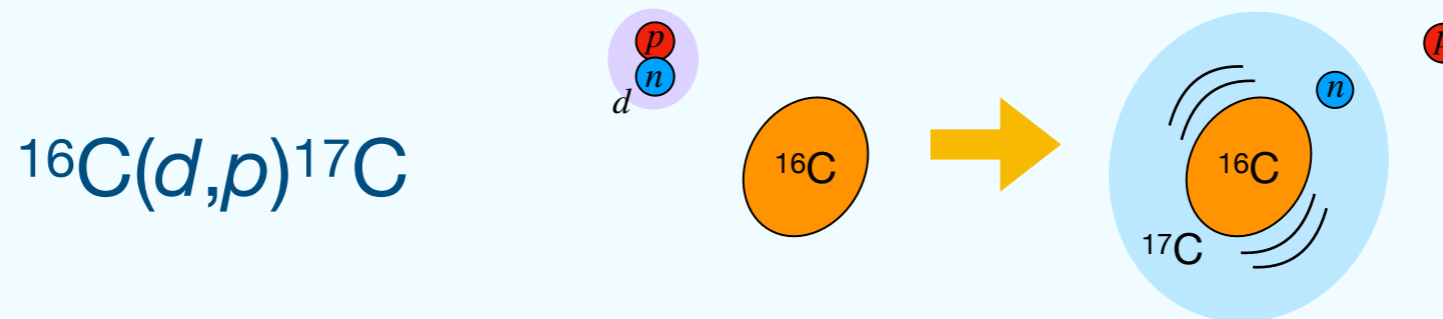


^{17}C

Single Particle Levels



Pseudo-states discretization method



$$\mathcal{F}(k, \theta) = \sqrt{\frac{\mu_i \mu_f k_f}{(2\pi\hbar^2)^2 k_i}} \langle \chi_f^{(-)} \varphi_p \Psi_{17\text{C}} | U | \chi_i^{(+)} \psi_d \phi_{16\text{C}} \rangle$$

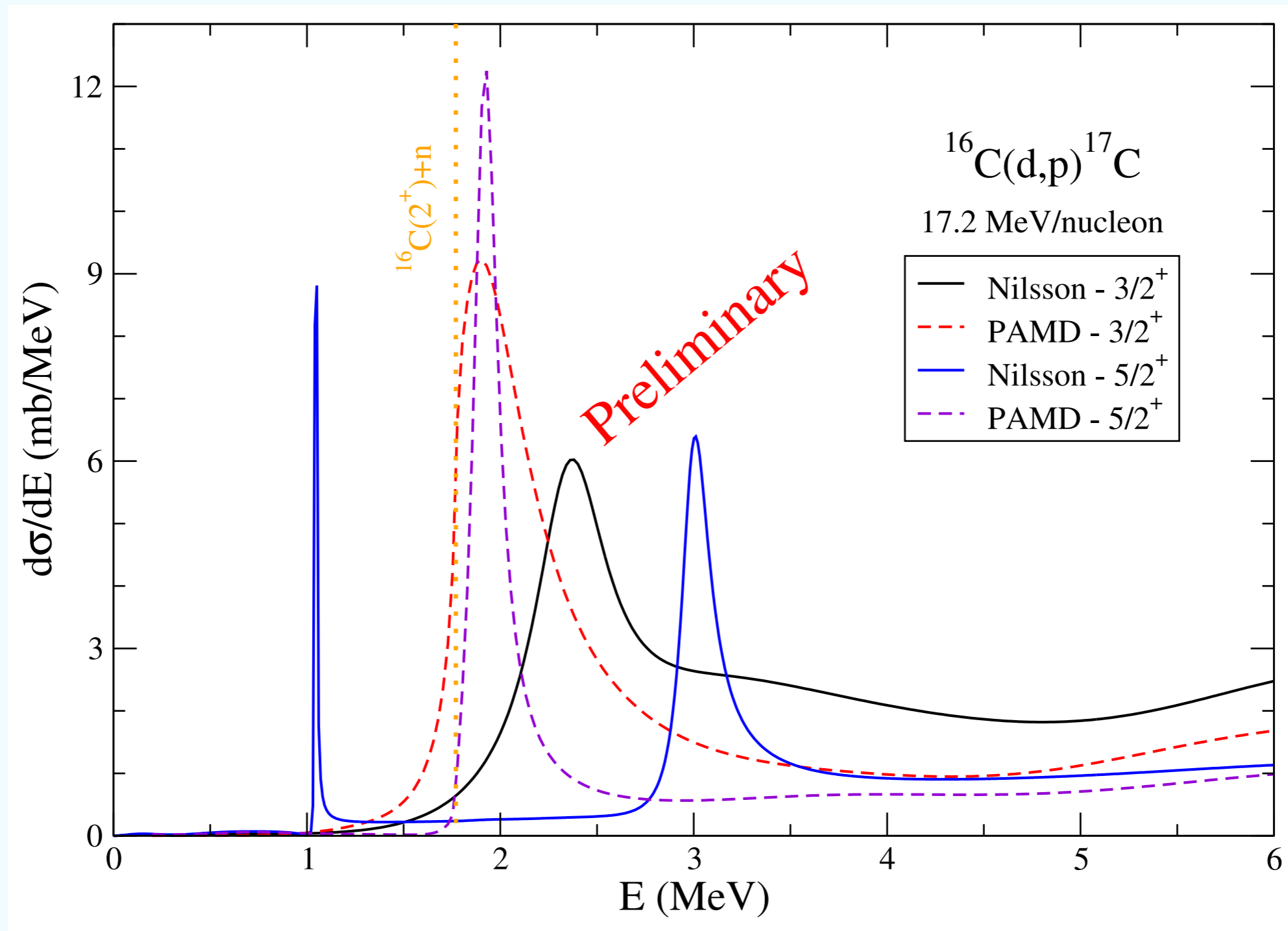
- Scattering amplitudes for a discrete number of pseudo-states

$$\mathcal{F}_n^{THO}(\theta) = \sqrt{\frac{\mu_i \mu_f k_f}{(2\pi\hbar^2)^2 k_i}} \langle \chi_f^{(-)} \varphi_p \Psi_n^{THO} | U | \chi_i^{(+)} \psi_d \phi_{16\text{C}} \rangle$$

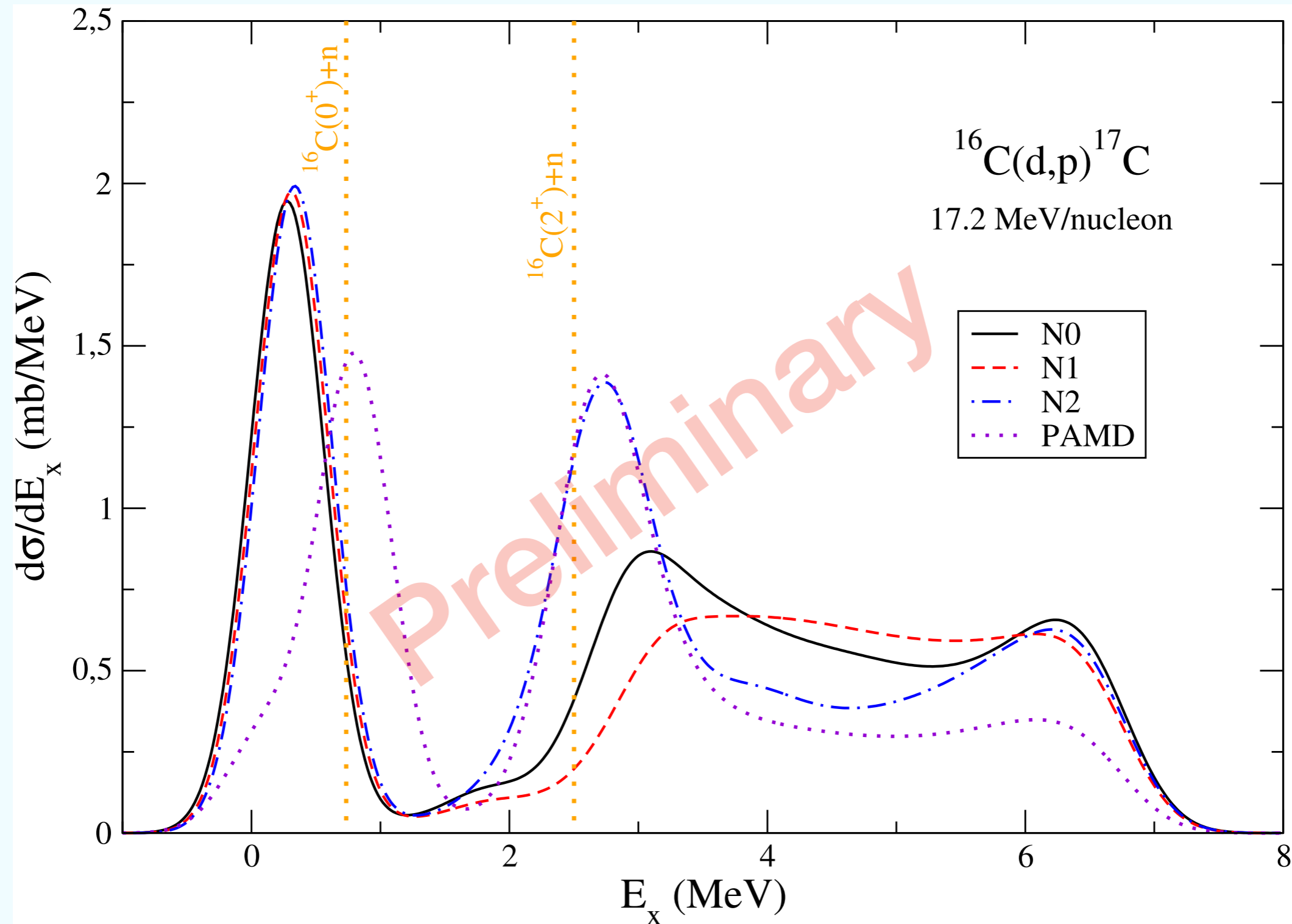
- Convolution with the *exact* n +core scattering states

$$\mathcal{F}(k, \theta) \approx \sum_n \langle \Psi_{17\text{C}}(k, r) | \Psi_n^{THO}(r) \rangle \mathcal{F}_n^{THO}(\theta)$$

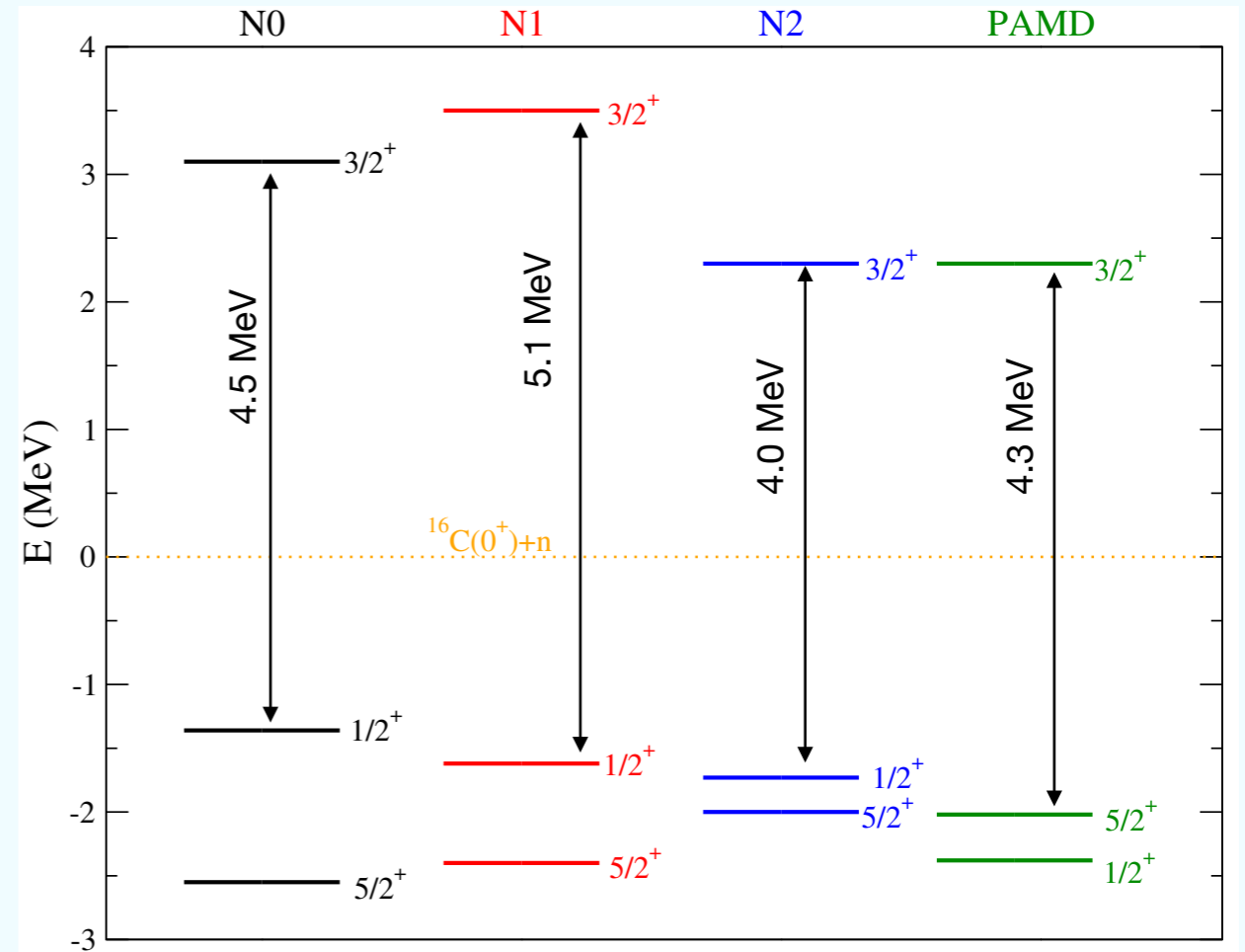
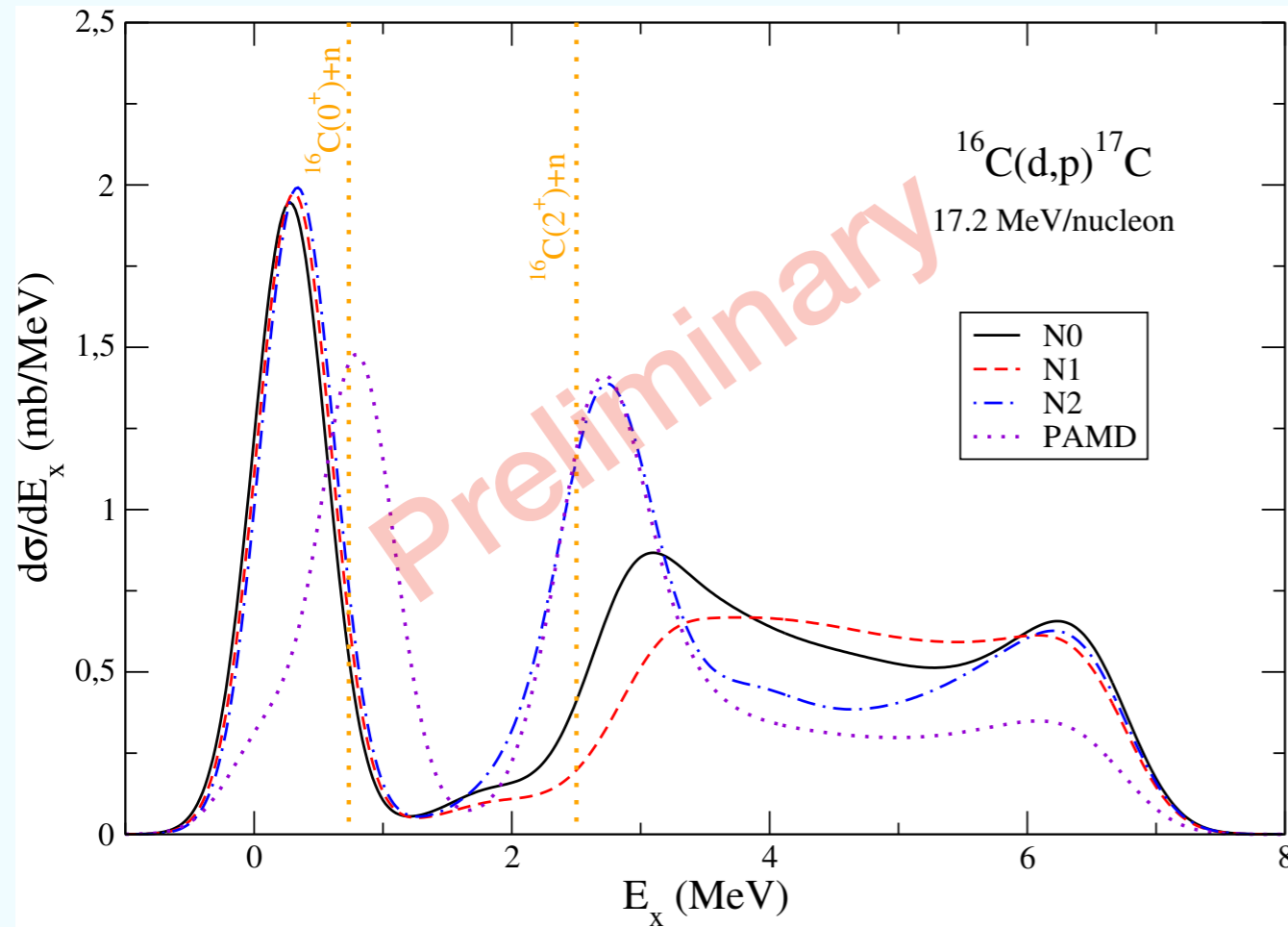
Transfer to the continuum: $^{16}\text{C}(d,p)^{17}\text{C}$



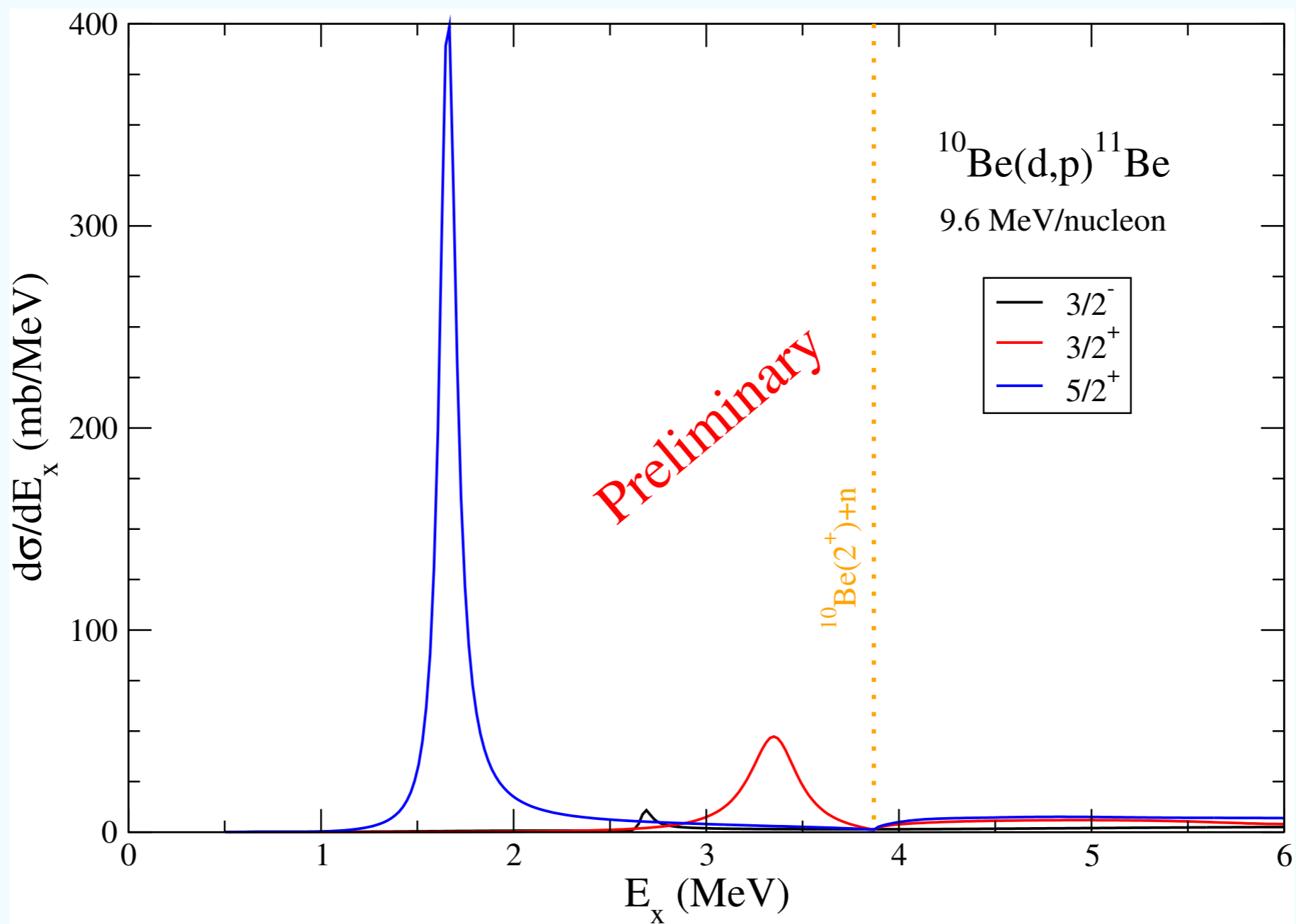
Transfer to the continuum: $^{16}\text{C}(d,p)^{17}\text{C}$



Transfer to the continuum: $^{16}\text{C}(d,p)^{17}\text{C}$



Transfer to the Continuum: $^{10}\text{Be}(d,p)^{11}\text{Be}$



Conclusions

- We have checked the suitability of the PAMD and Nilsson models to describe ^{11}Be and ^{17}C systems.
 - ▶ ^{11}Be : better with PAMD model (weak-coupling limit)
 - ▶ ^{17}C : better with Nilsson model (strong-coupling limit)
- These models have been applied to transfer reactions to bound states and to the continuum using pseudo-states.
- A good reproduction of the structure and transfer reactions have been found.
- The reaction $^{16}\text{C}(d,p)^{17}\text{C}$ to the continuum have been studied in order to confirm the existence of the N=16 shell-gap.
- Extensions to breakup reactions and other weakly bound nuclei are in progress.



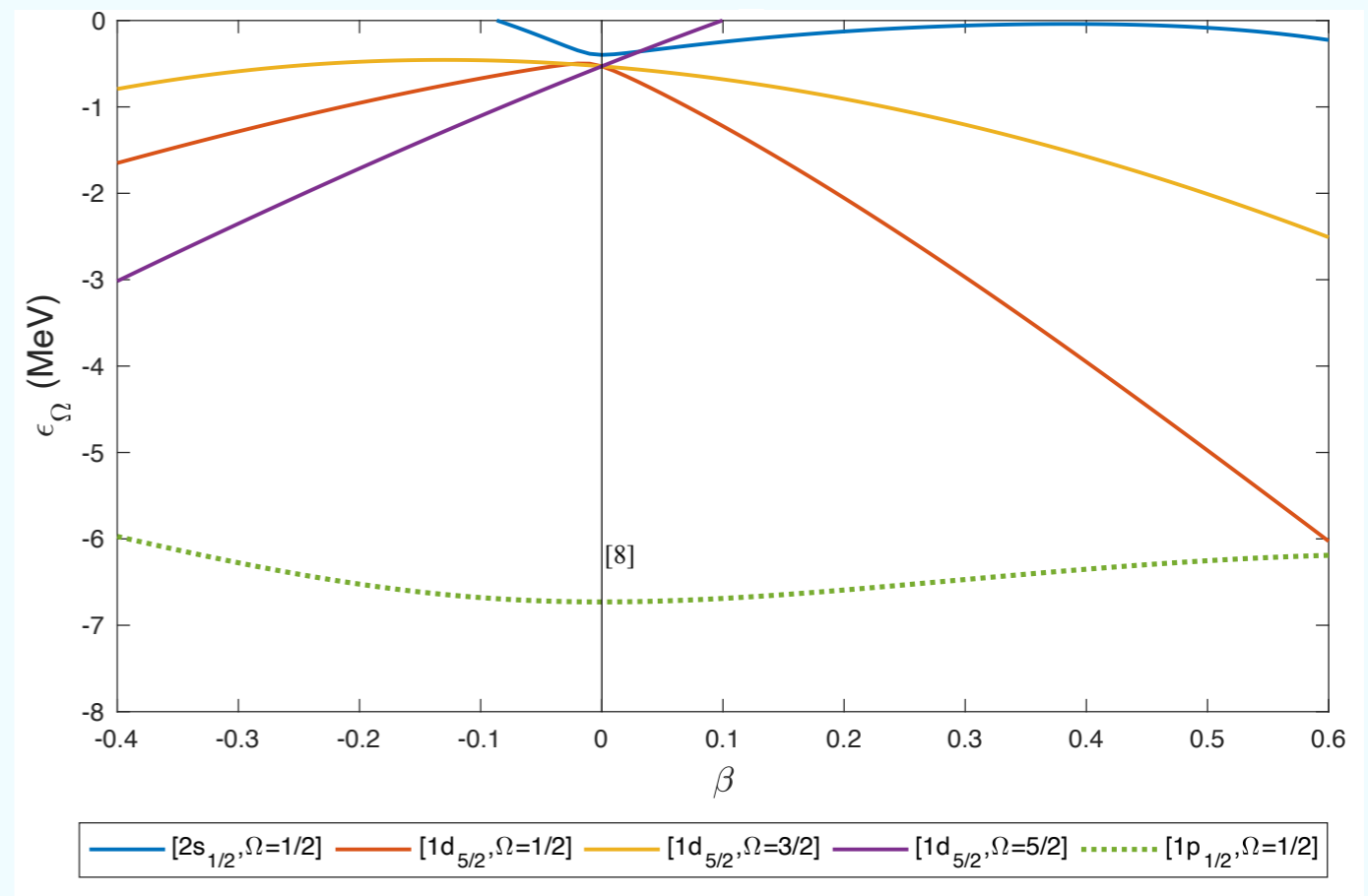
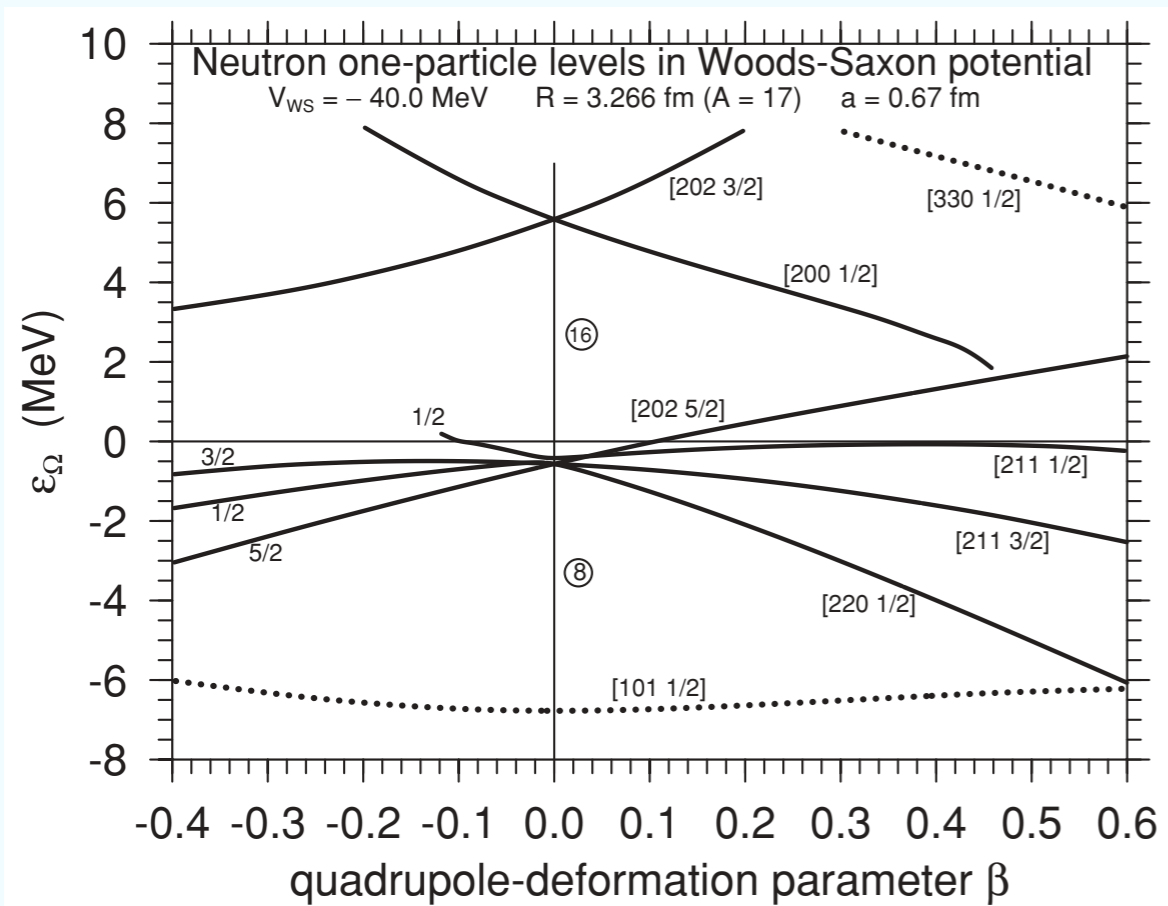
**Thanks for your
attention**



Backup

Nilsson Diagram for ^{17}C

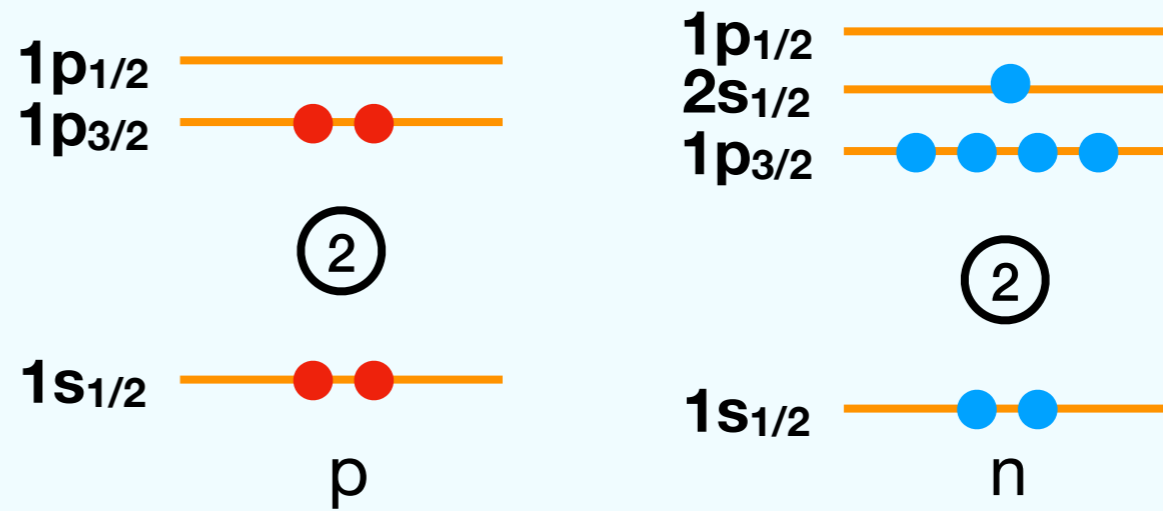
$$H_{sp} = -\frac{\hbar^2}{2\mu}\nabla^2 + V_c(r) + V_{ls}(r)(\vec{l} \cdot \vec{s}) - r\beta \frac{dV_c(r)}{dr} Y_{20}(\theta')$$



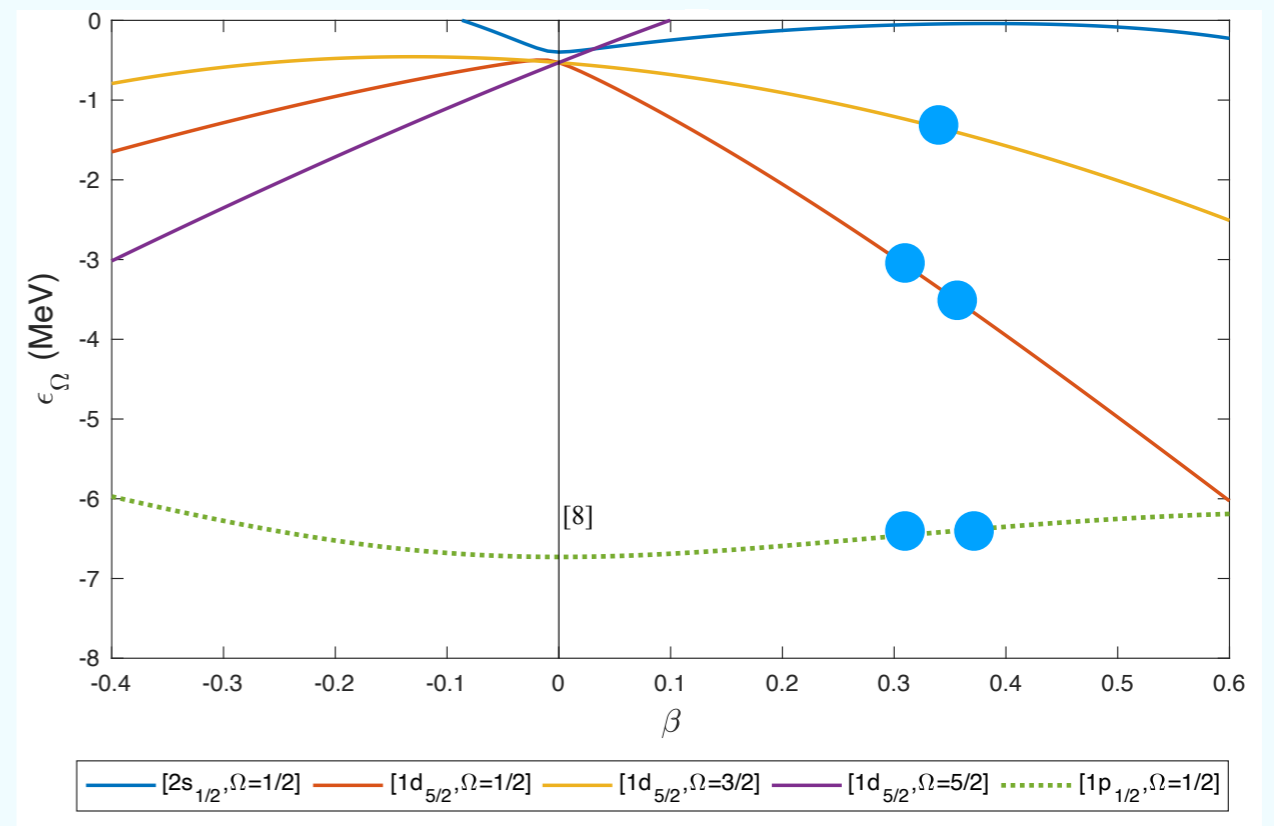
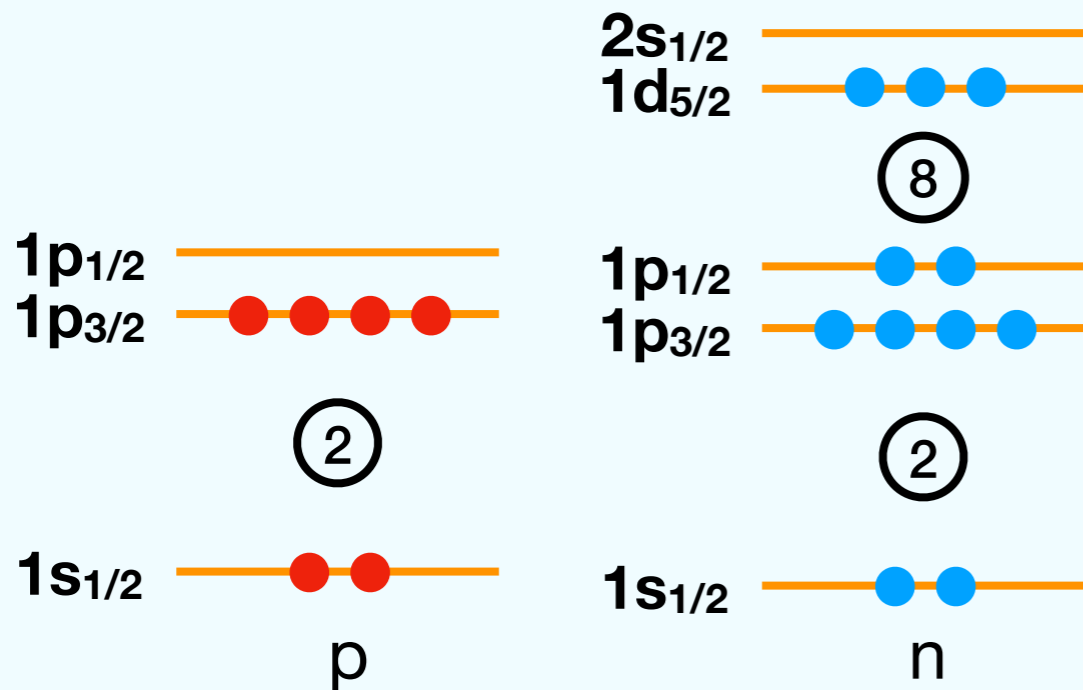
I. Hamamoto, PRC76 (2007) 054319

Occupation of single-particle levels

^{11}Be

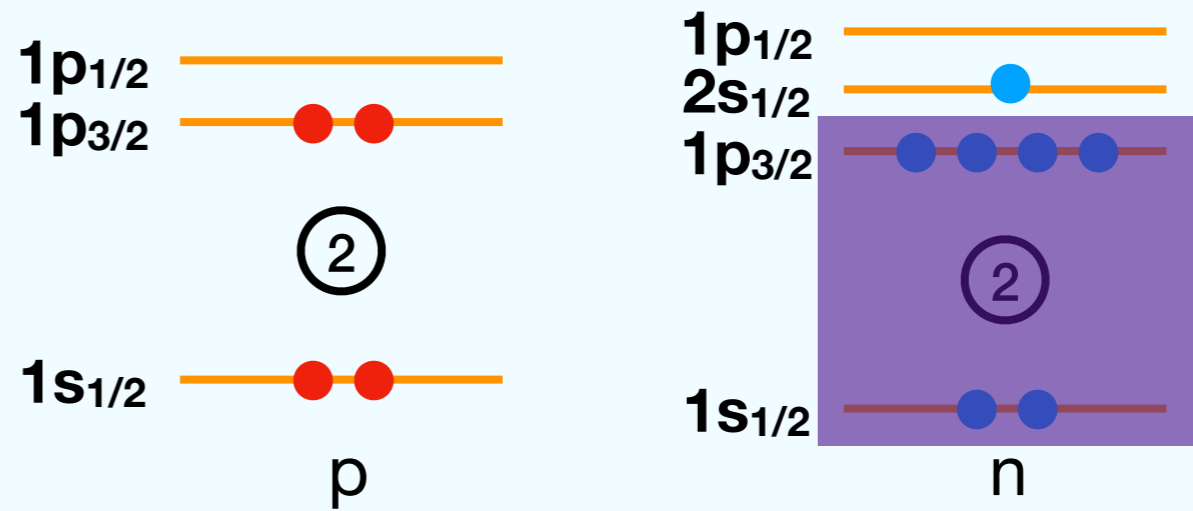


^{17}C

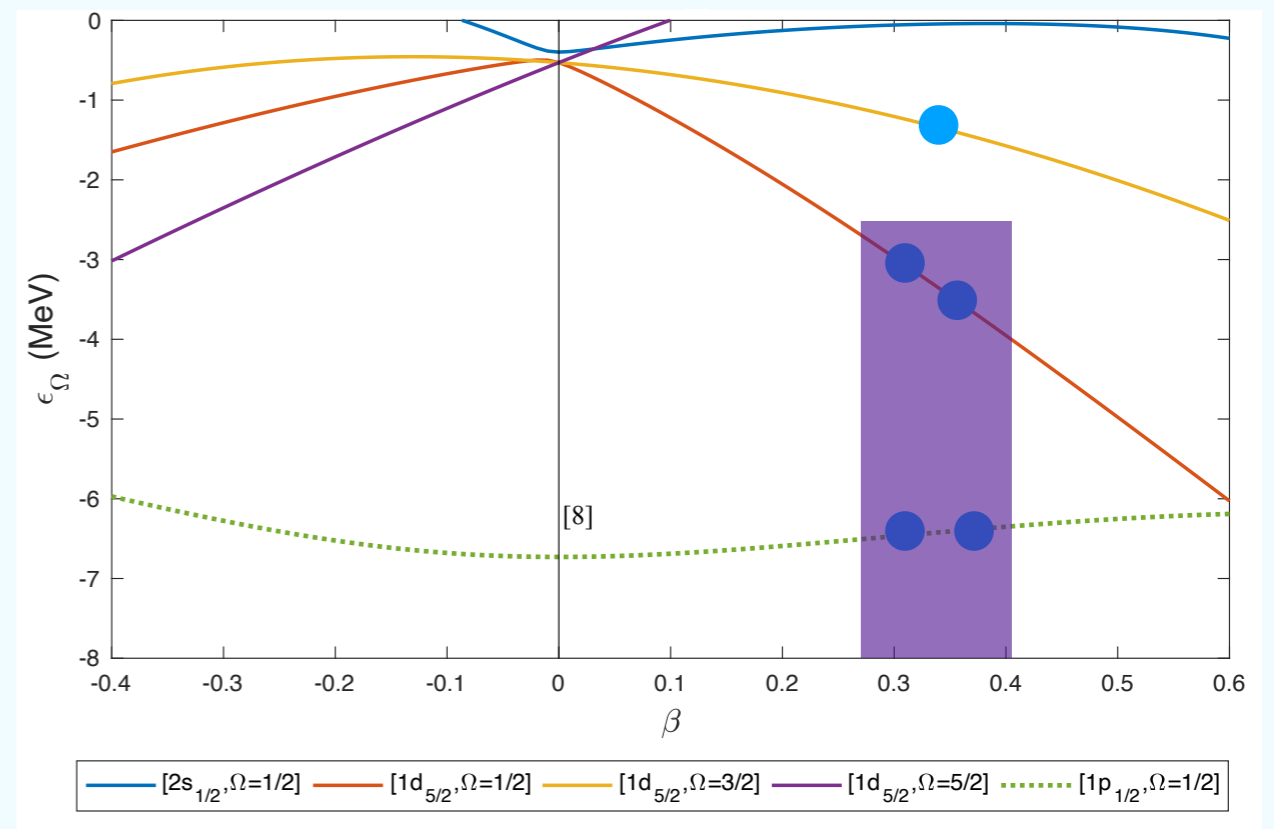
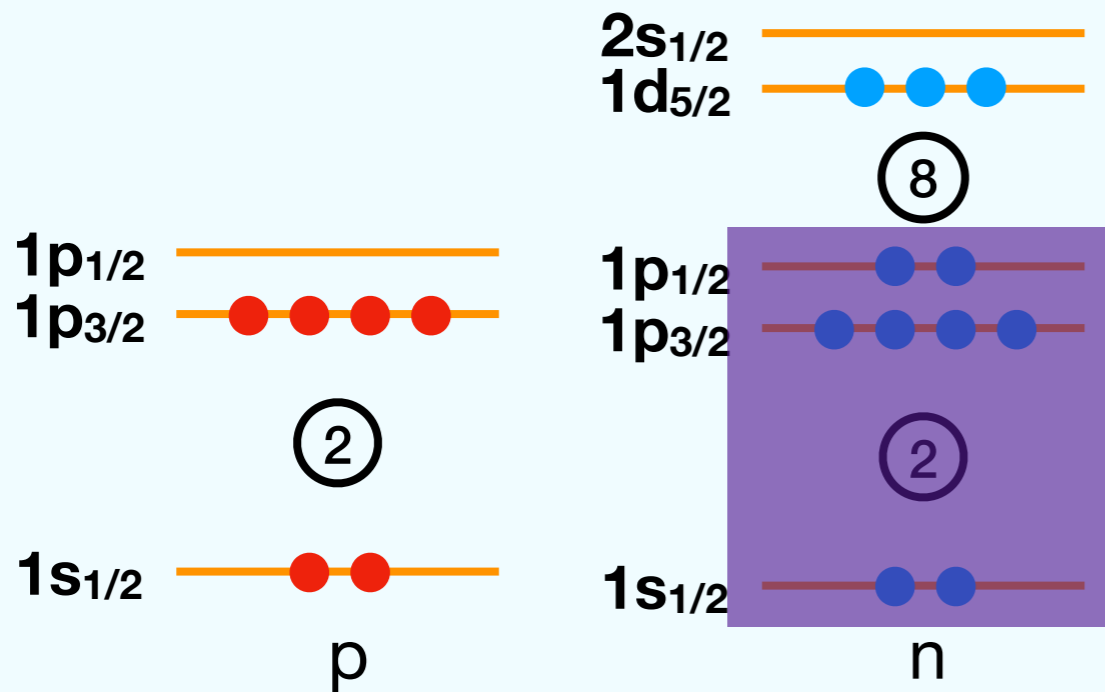


Pauli Blocking

^{11}Be

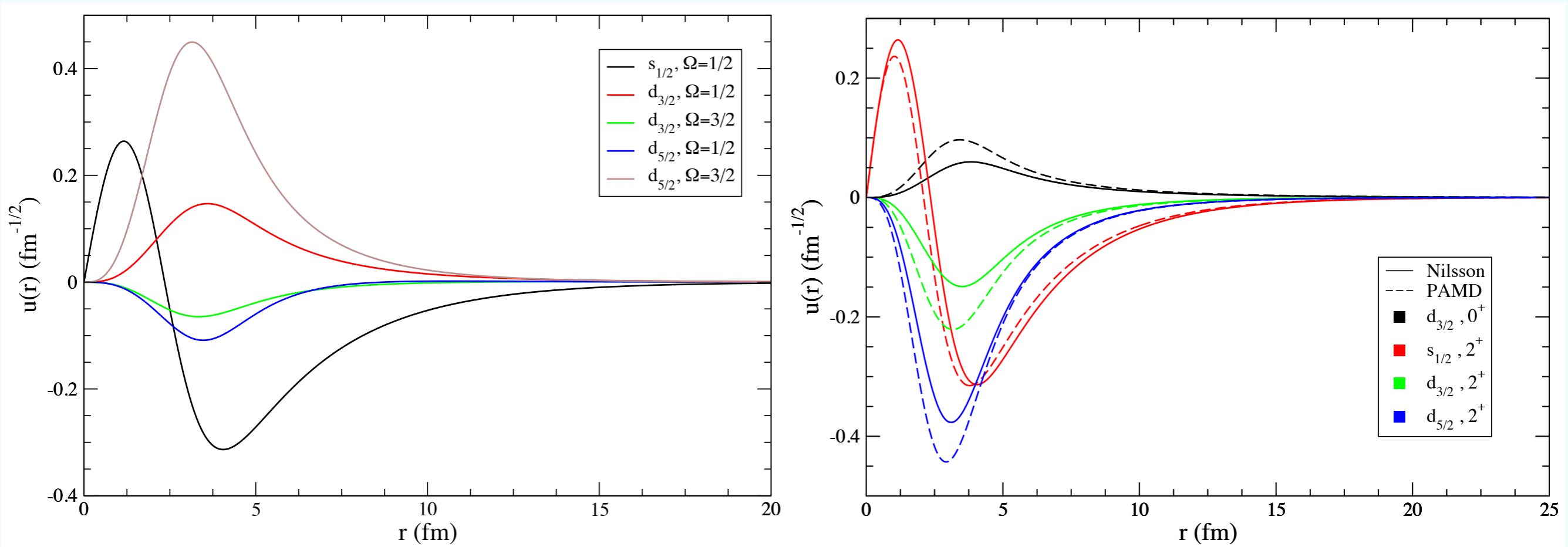


^{17}C



Wave Functions

^{17}C ground state $3/2^+$



Ω - axial projections

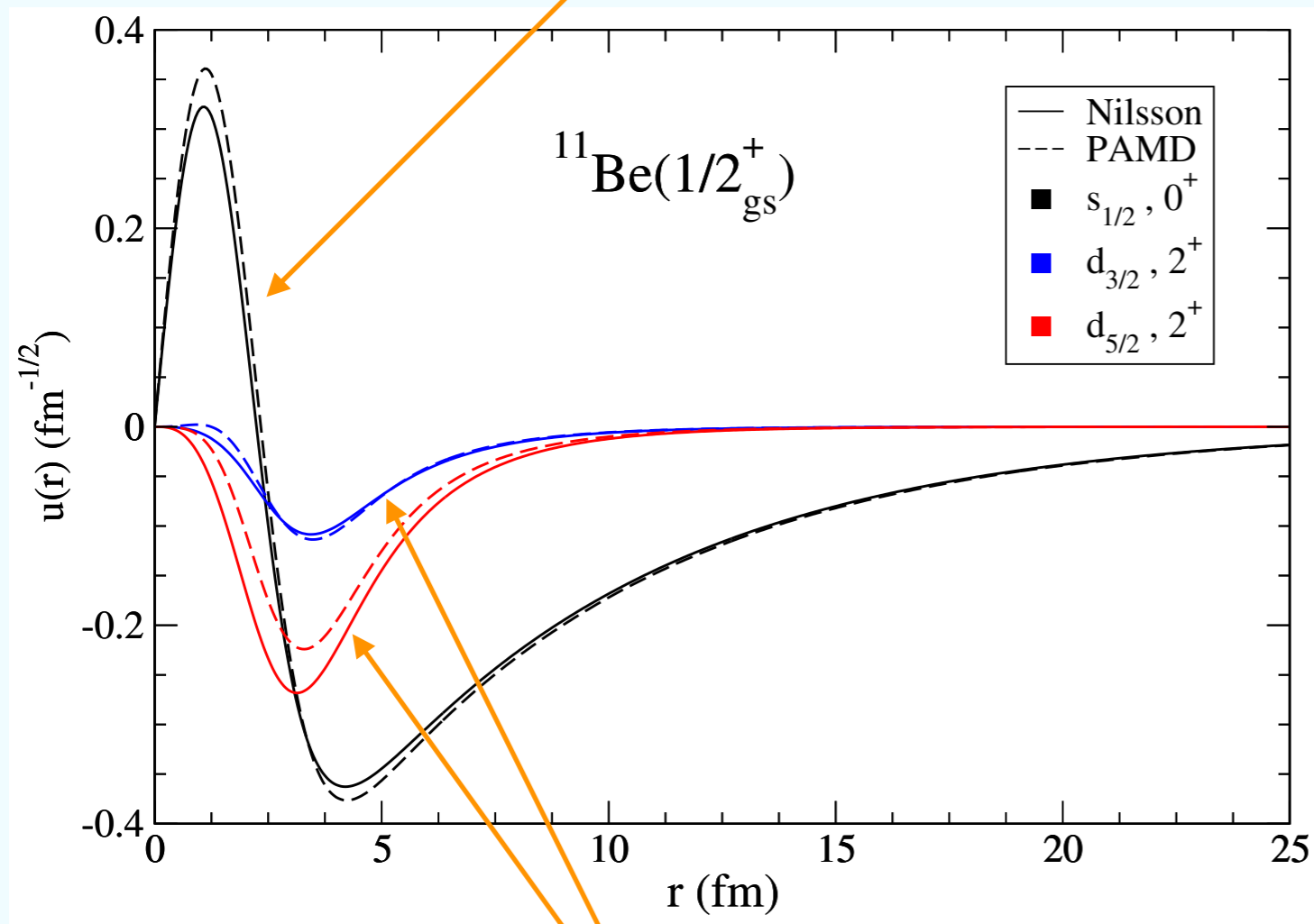


I - core states

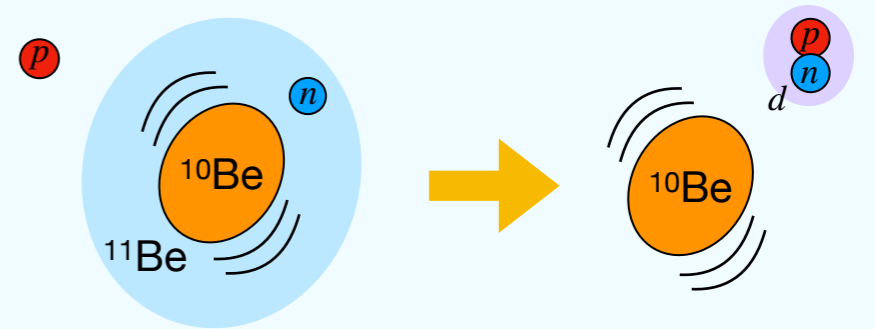
$$u_{jI}^{\lambda J\pi}(r) = \sqrt{2} \sum_{\Omega} (-1)^{J+\Omega} \langle I0 | j - \Omega J \Omega \rangle u_{j\Omega}^{\lambda J\pi}(r)$$

$^{11}\text{Be}(p,d)^{10}\text{Be}$

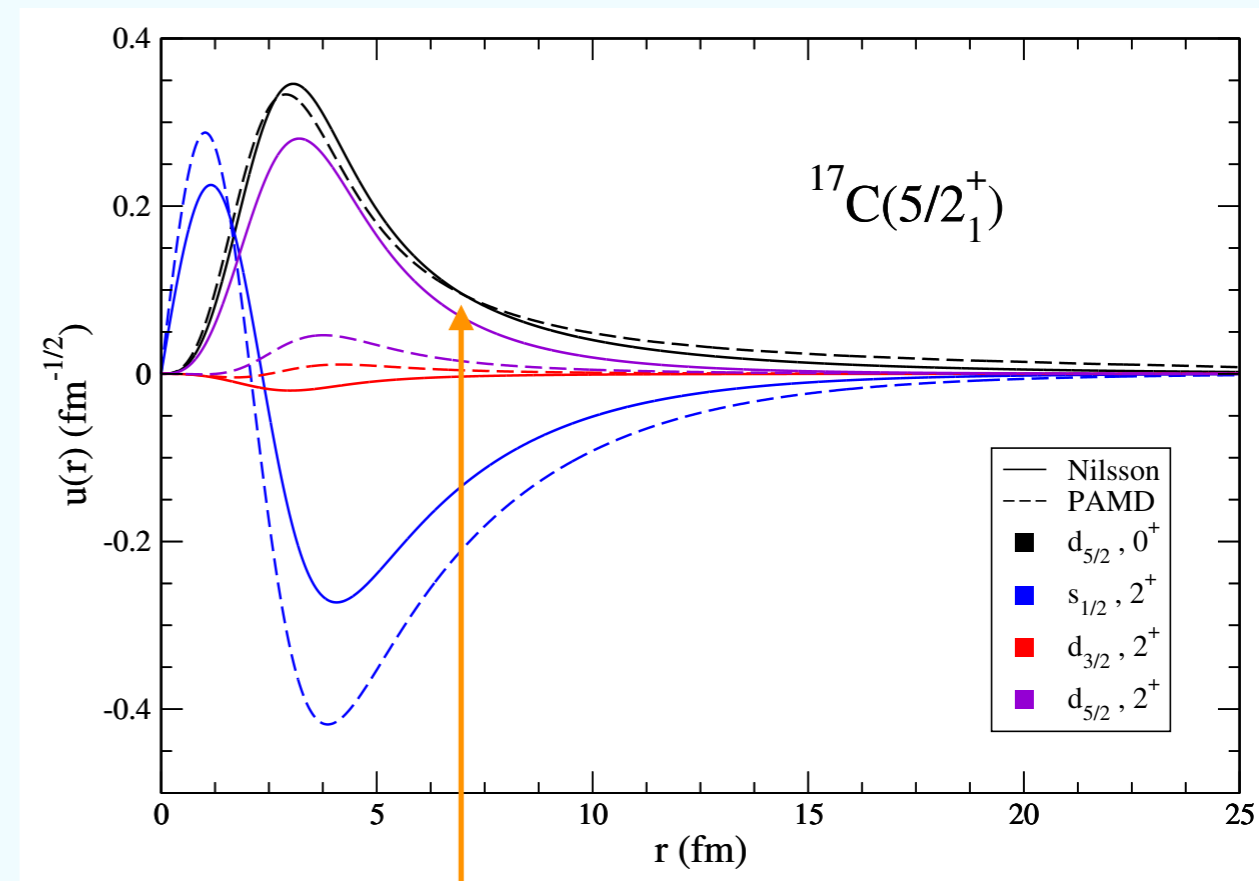
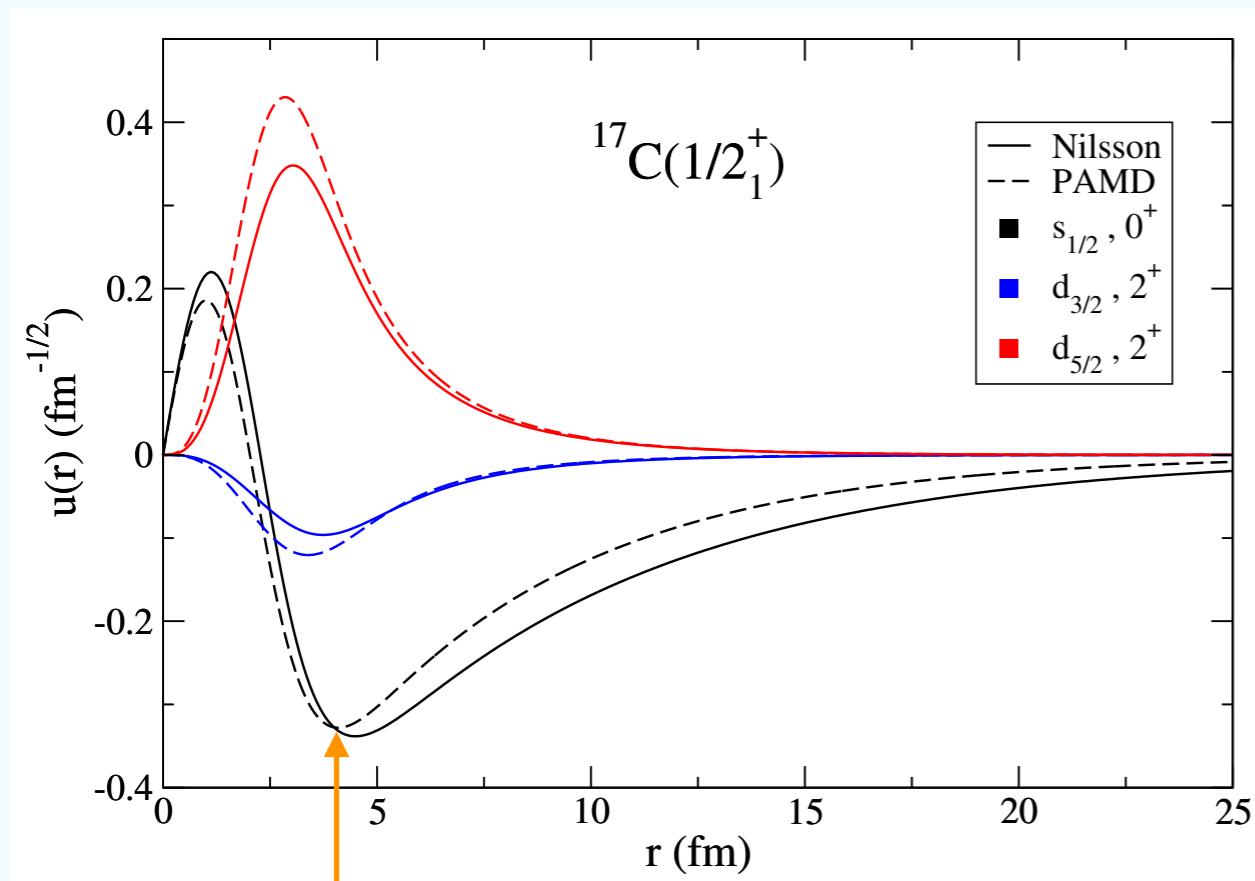
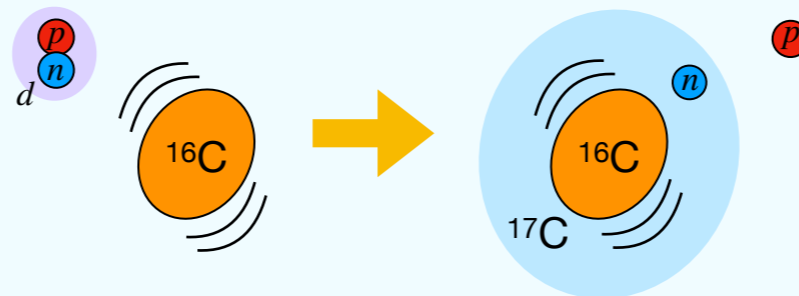
$$\langle ^{10}\text{Be}(0^+) | ^{11}\text{Be}(1/2_{gs}^+) \rangle \rightarrow \left(\frac{d\sigma}{d\Omega} \right)_{^{10}\text{Be}(0^+)}$$



$$\langle ^{10}\text{Be}(2^+) | ^{11}\text{Be}(1/2_{gs}^+) \rangle \rightarrow \left(\frac{d\sigma}{d\Omega} \right)_{^{10}\text{Be}(2^+)}$$



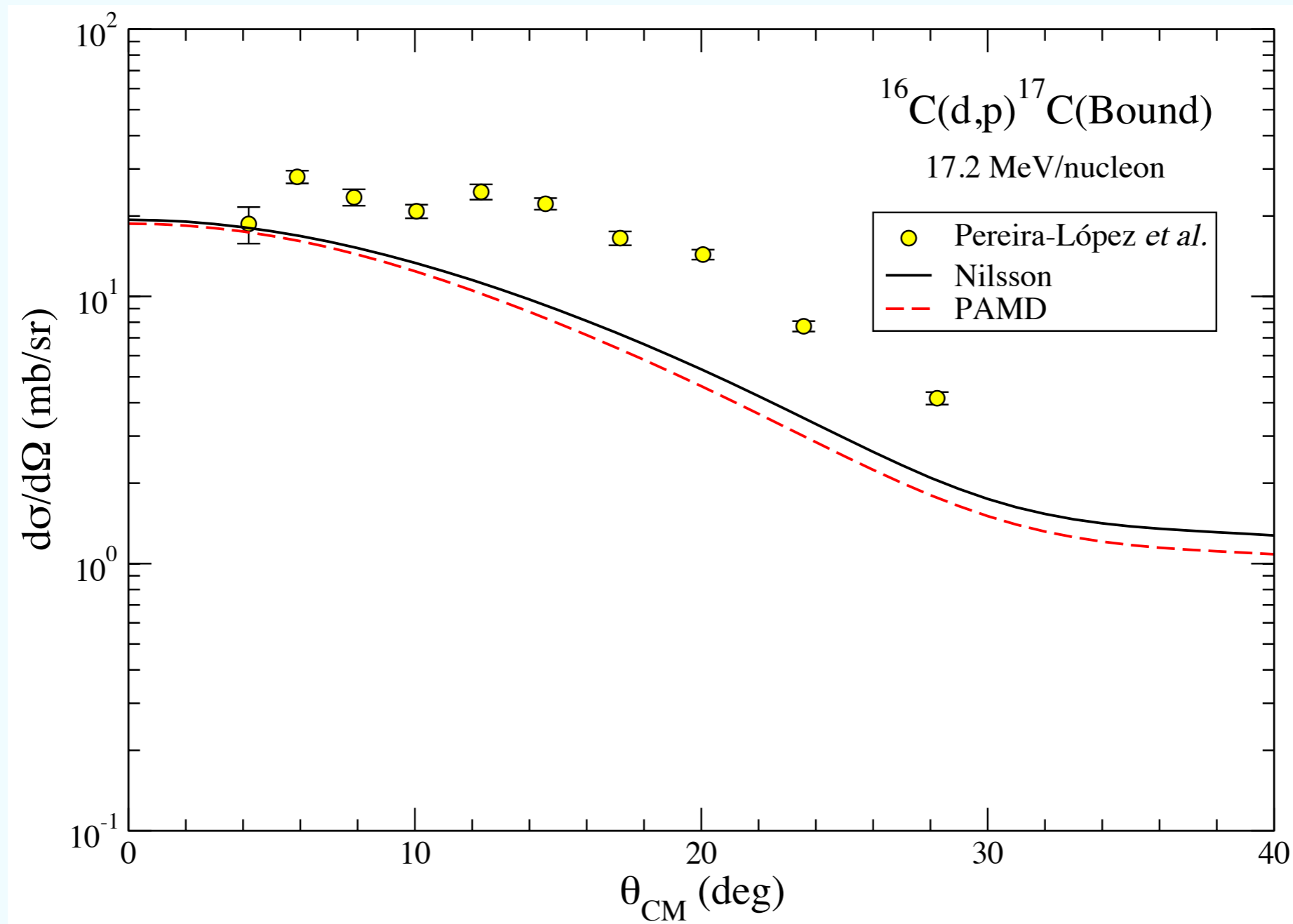
$^{16}\text{C}(d,p)^{17}\text{C}$



$$\langle ^{17}\text{C}(1/2_1^+) | ^{16}\text{C}(0^+) \rangle \rightarrow \left(\frac{d\sigma}{d\Omega} \right)_{^{17}\text{C}(1/2_1^+)} \quad \langle ^{17}\text{C}(5/2_1^+) | ^{16}\text{C}(0^+) \rangle \rightarrow \left(\frac{d\sigma}{d\Omega} \right)_{^{17}\text{C}(5/2_1^+)}$$

Transfer to bound states

Sum for the three bound states



Transfer to the continuum

- The *prior* form of the ADWA approximation is used

✗ $\langle {}^{17}\text{C} | {}^{16}\text{C} \rangle \rightarrow$ Too much spatial extension

✓ $\langle {}^{17}\text{C} | V_{n-16}\text{C} | {}^{16}\text{C} \rangle \rightarrow$ Little spatial extension

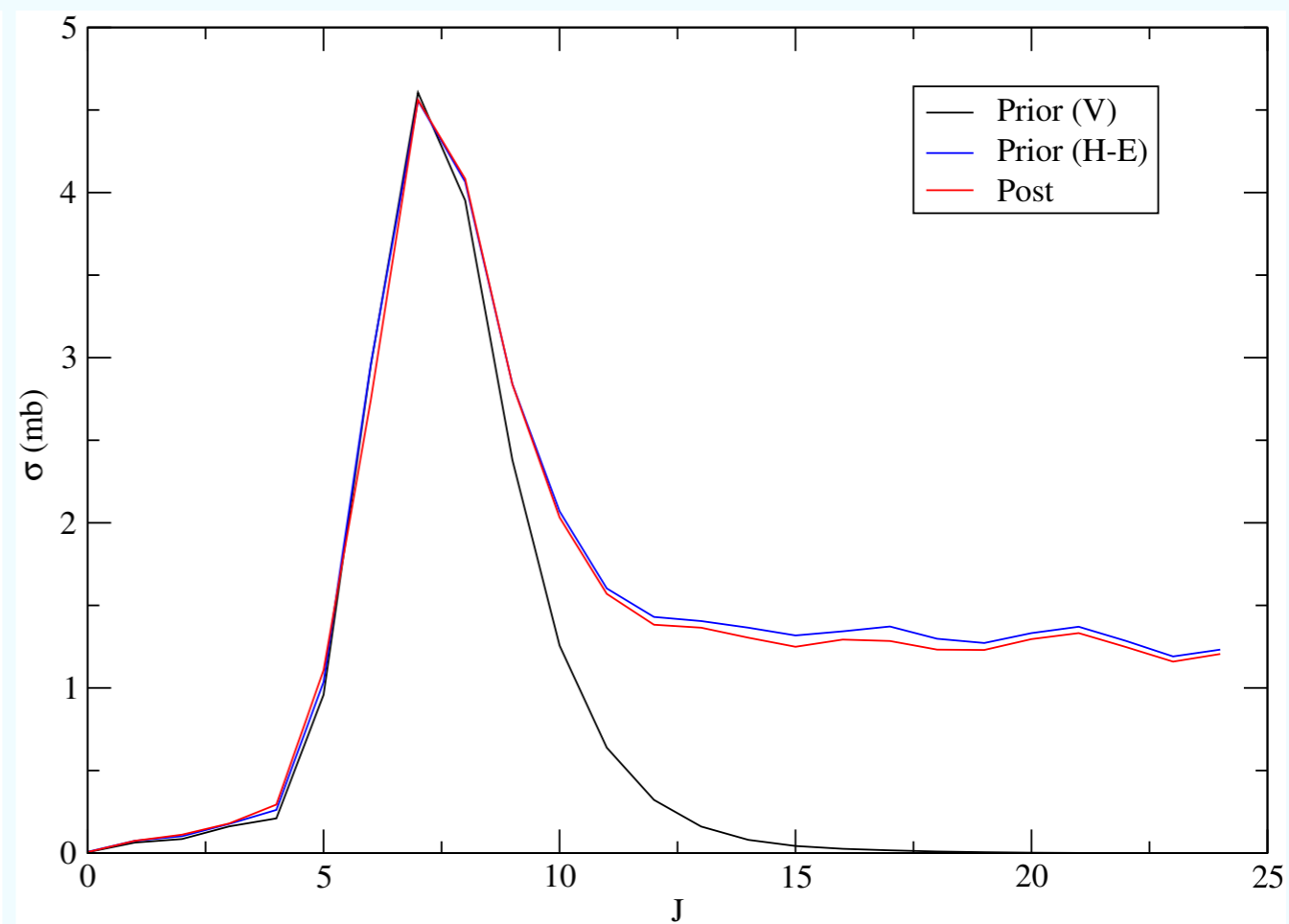
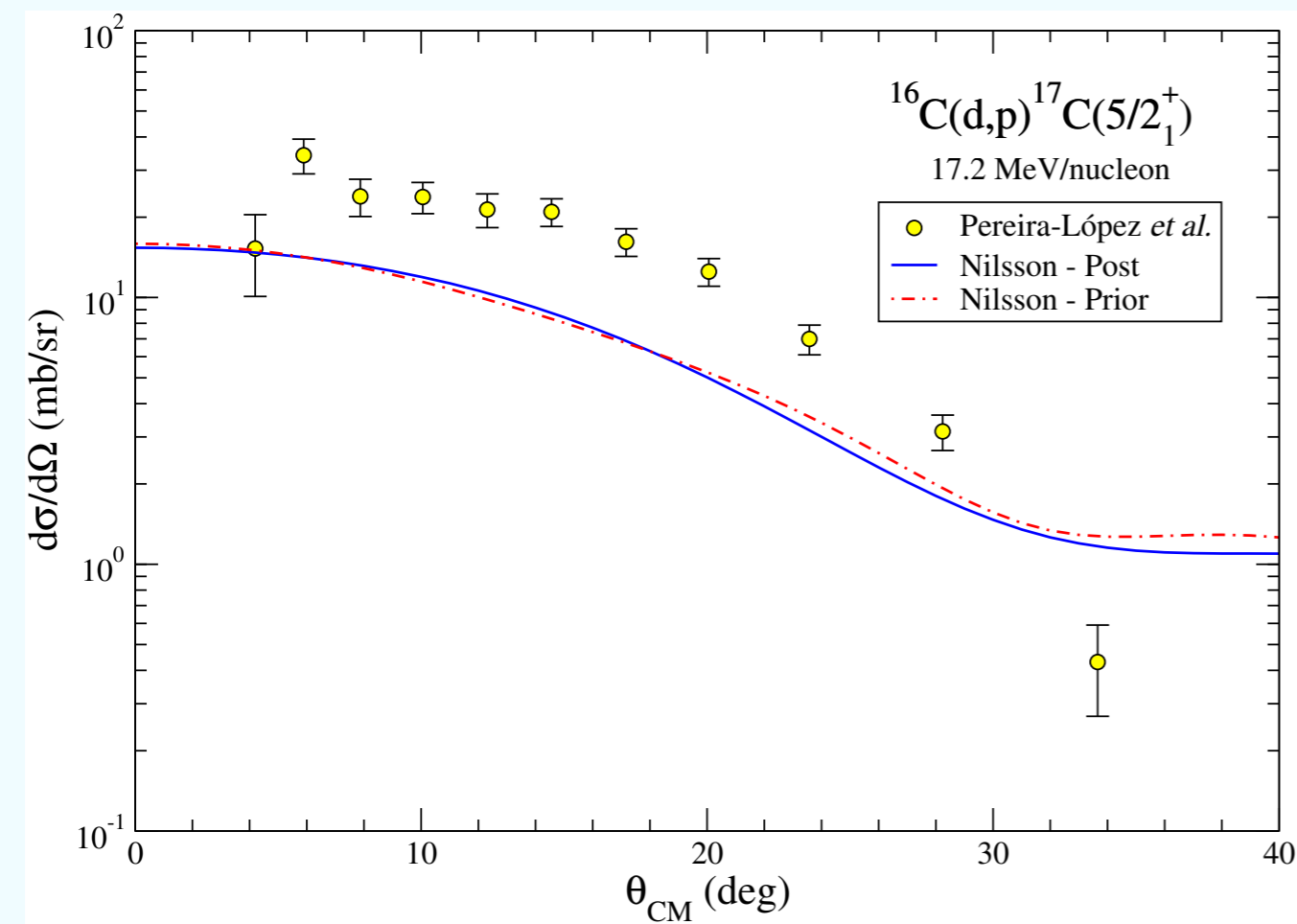
- The continuum is discretized with the pseudo-states method
 - Scattering amplitudes for a discrete number of pseudo-states
 - Convolution with the *exact* n +core scattering states

$$\mathcal{F}(k, \theta) \approx \sum_n \langle \Psi(k, r) | \Psi_n^{THO}(r) \rangle \mathcal{F}_n^{THO}(\theta)$$

$^{16}\text{C}(d,p)^{17}\text{C}$: Prior-Post

Bound states

Transfer to the continuum



Transfer to the continuum: Decay mode

