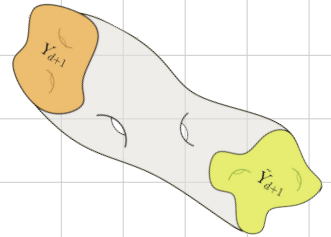


# ANOMALIES & BORDISMS OF NON-SUPERSYMMETRIC STRINGS

Matilda Delgado

Based on:  
[2310.06895] I. Basile, A. Debray, M.D., M. Montero



# BIG PICTURE

**Our world is non-supersymmetric**

# BIG PICTURE

**Our world is non-supersymmetric**

(at least at low energies)

# BIG PICTURE

**Our world is non-supersymmetric**

(at least at low energies)

**It is crucial for phenomenology to understand String Theory (QG) in setups without supersymmetry!**

# BIG PICTURE

**Our world is non-supersymmetric**

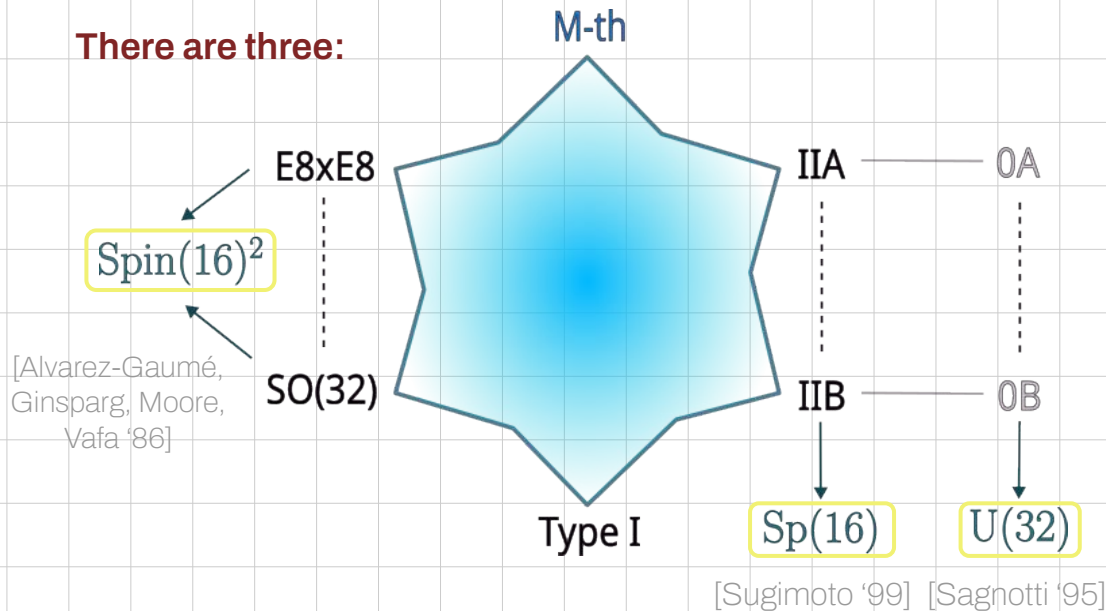
(at least at low energies)

**It is crucial for phenomenology to understand String Theory (QG) in setups without supersymmetry!**

## How?

# 10D NON-SUPERSYMMETRIC STRING THEORIES

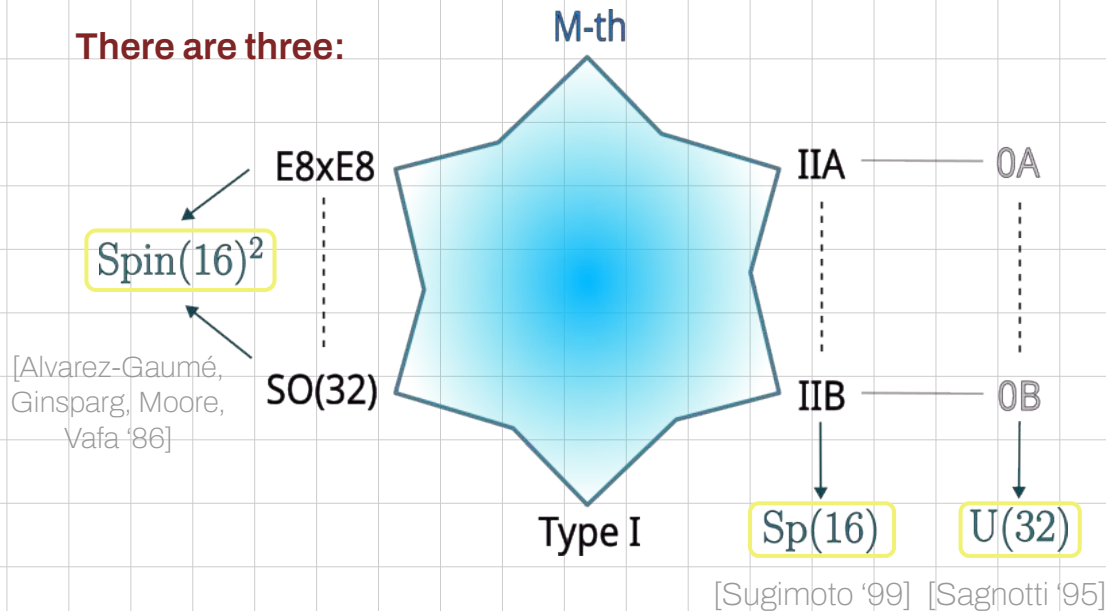
There are three:



Arguably the most natural way to study QG away from SUSY !!

# 10D NON-SUPERSYMMETRIC STRING THEORIES

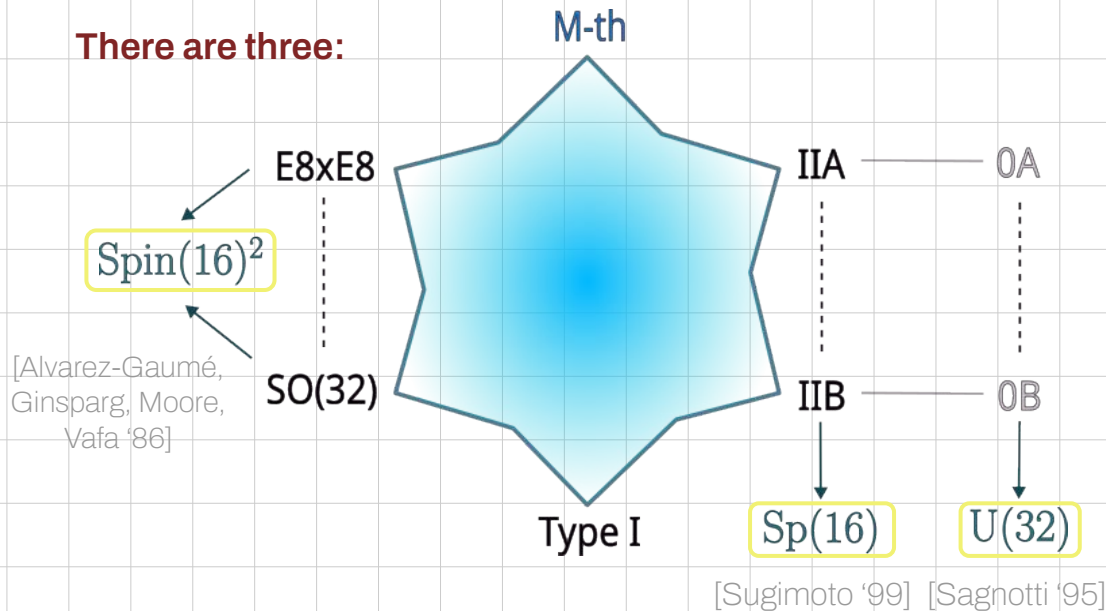
There are three:



Not a lot is known about these theories..

# 10D NON-SUPERSYMMETRIC STRING THEORIES

There are three:



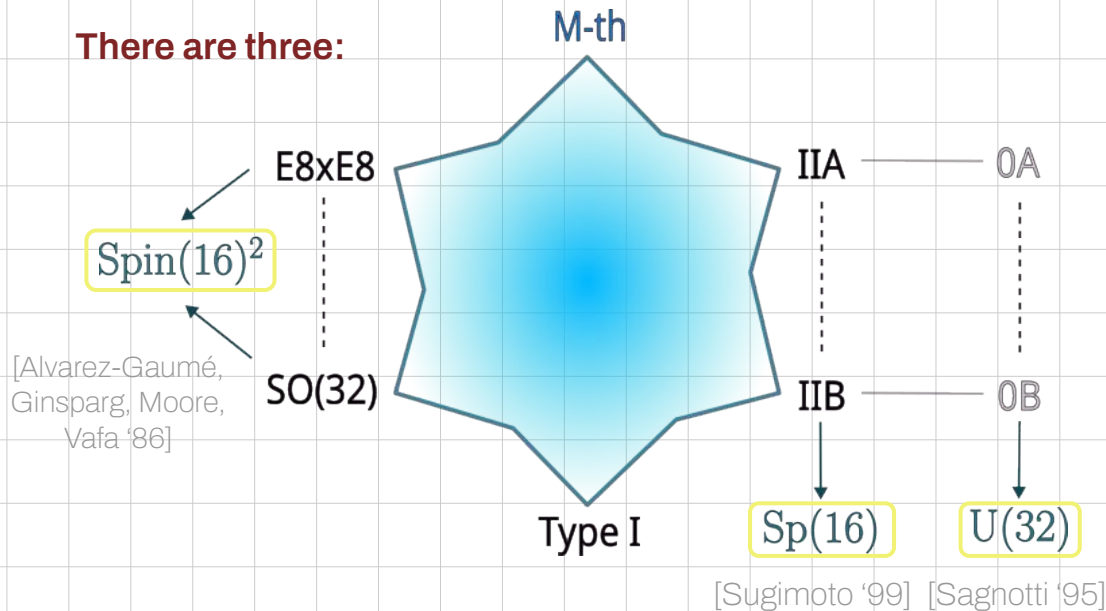
Not a lot is known about these theories..

e.g. do gauge/gravitational anomalies cancel?



# 10D NON-SUPERSYMMETRIC STRING THEORIES

There are three:



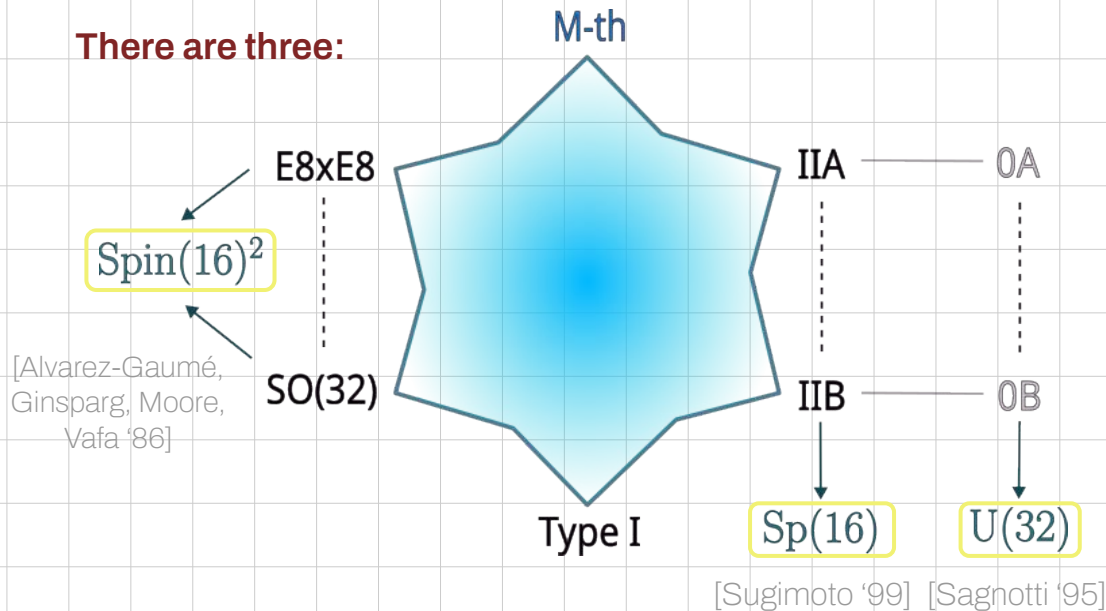
Not a lot is known about these theories..

e.g. do gauge/gravitational anomalies cancel?

- local anomaly cancellation ✓

# 10D NON-SUPERSYMMETRIC STRING THEORIES

There are three:



Not a lot is known about these theories..

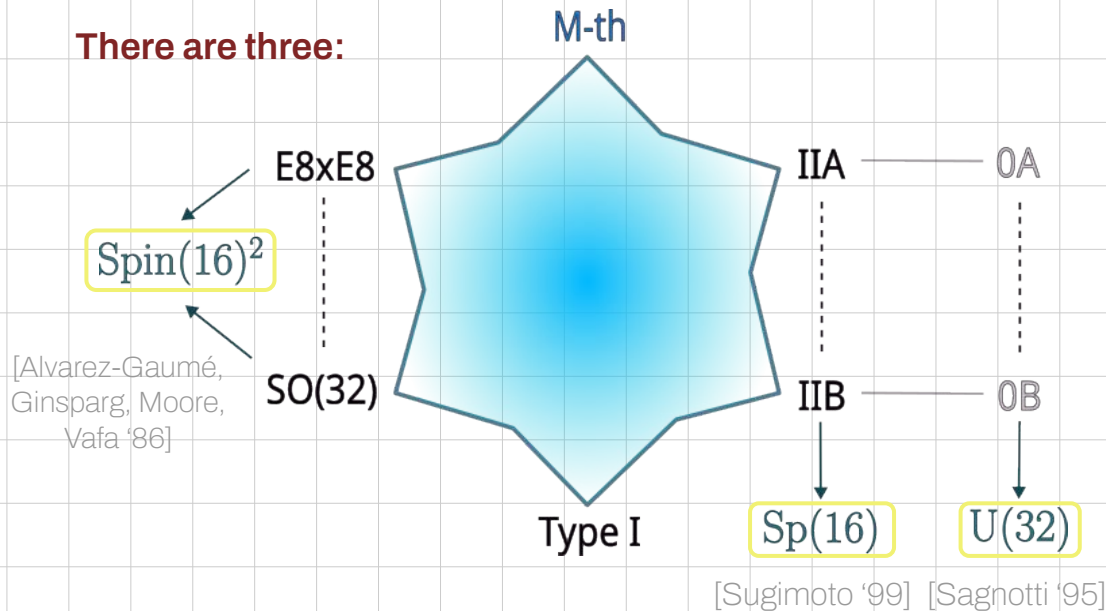
e.g. do gauge/gravitational anomalies cancel?

- local anomaly cancellation ✓

- global anomaly cancellation ?

# 10D NON-SUPERSYMMETRIC STRING THEORIES

There are three:



Not a lot is known about these theories..

e.g. do gauge/gravitational anomalies cancel?

- local anomaly cancellation ✓

- global anomaly cancellation ?

Our work answers this question

HOW?

by computing the relevant **bordism groups** for these theories

[e.g. Álvarez-Gaumé, Vázquez-Mozo '22]

# A CRASH COURSE ON GLOBAL ANOMALIES

An anomaly in a **gauge** transformation or diffeomorphism is:

$$Z[X_d] \implies \tilde{Z}[X_d] \neq Z[X_d]$$

Global anomalies = associated to a transformation that cannot be *deformed to the identity*

# A CRASH COURSE ON GLOBAL ANOMALIES

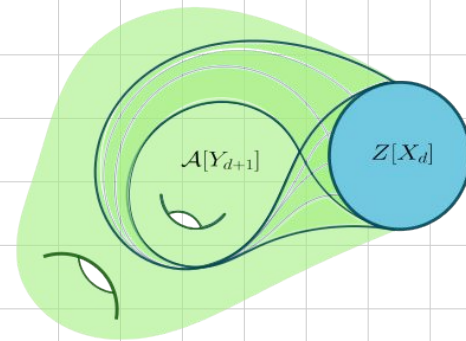
[eg. García-Etxebarria, Montero '18]

An anomaly in a **gauge** transformation or diffeomorphism is:

$$Z[X_d] \implies \tilde{Z}[X_d] \neq Z[X_d]$$

Global anomalies = associated to a transformation that cannot be *deformed to the identity*

The modern way of computing **global** anomalies of a theory on  $X_d$  is through a (d+1)-dimensional *anomaly theory* on  $Y_{d+1}$  such that  $\partial Y_{d+1} = X_d$ .



# A CRASH COURSE ON GLOBAL ANOMALIES

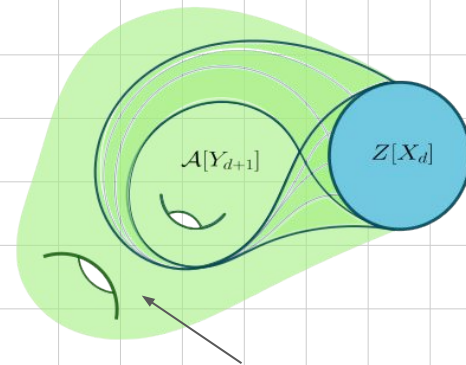
An anomaly in a **gauge** transformation or diffeomorphism is:

$$Z[X_d] \implies \tilde{Z}[X_d] \neq Z[X_d]$$

Global anomalies = associated to a transformation that cannot be *deformed to the identity*

The modern way of computing **global** anomalies of a theory on  $X_d$  is through a (d+1)-dimensional *anomaly theory* on  $Y_{d+1}$  such that  $\partial Y_{d+1} = X_d$ .

[eg. García-Etxebarria, Montero '18]



In QG, allow for topology-change  
 $\Rightarrow$  Dai-Freed anomalies

# A CRASH COURSE ON GLOBAL ANOMALIES

[eg. García-Etxebarria, Montero '18]

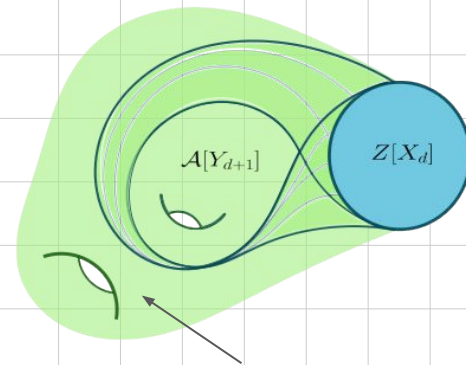
An anomaly in a **gauge** transformation or diffeomorphism is:

$$Z[X_d] \implies \tilde{Z}[X_d] \neq Z[X_d]$$

Global anomalies = associated to a transformation that cannot be *deformed to the identity*

The modern way of computing **global** anomalies of a theory on  $X_d$  is through a (d+1)-dimensional **anomaly theory** on  $Y_{d+1}$  such that  $\partial Y_{d+1} = X_d$ .

The anomaly theory is **engineered** to give the **exact (opposite) anomaly** of the one you started with.



In QG, allow for topology-change  
 $\implies$  Dai-Freed anomalies

# A CRASH COURSE ON GLOBAL ANOMALIES

[eg. García-Etxebarria, Montero '18]

An anomaly in a **gauge** transformation or diffeomorphism is:

$$Z[X_d] \implies \tilde{Z}[X_d] \neq Z[X_d]$$

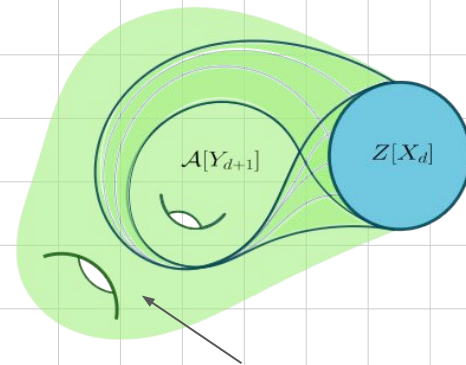
Global anomalies = associated to a transformation that cannot be *deformed to the identity*

The modern way of computing **global** anomalies of a theory on  $X_d$  is through a  $(d+1)$ -dimensional **anomaly theory** on  $Y_{d+1}$  such that  $\partial Y_{d+1} = X_d$ .

The anomaly theory is **engineered** to give the **exact (opposite) anomaly** of the one you started with.

—> the anomaly is much easier to compute this way because:

**The anomaly itself is a bordism invariant of these  $(d+1)$ -manifolds**



In QG, allow for topology-change  
=> Dai-Freed anomalies



# A CRASH COURSE ON GLOBAL ANOMALIES

[eg. García-Etxebarria, Montero '18]

An anomaly in a **gauge** transformation or diffeomorphism is:

$$Z[X_d] \implies \tilde{Z}[X_d] \neq Z[X_d]$$

Global anomalies = associated to a transformation that cannot be *deformed to the identity*

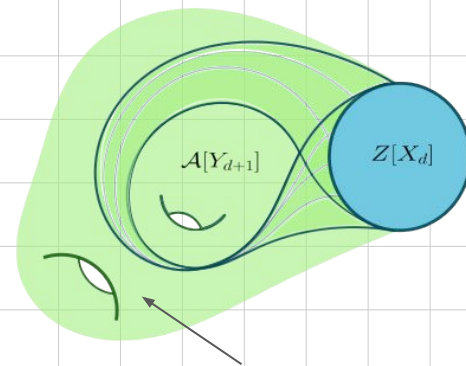
The modern way of computing **global** anomalies of a theory on  $X_d$  is through a (d+1)-dimensional **anomaly theory** on  $Y_{d+1}$  such that  $\partial Y_{d+1} = X_d$ .

The anomaly theory is **engineered** to give the **exact (opposite) anomaly** of the one you started with.

→ the anomaly is much easier to compute this way because:

**The anomaly itself is a bordism invariant of these (d+1)-manifolds**

→ we “just” have to compute 11D bordism groups for our three 10D theories



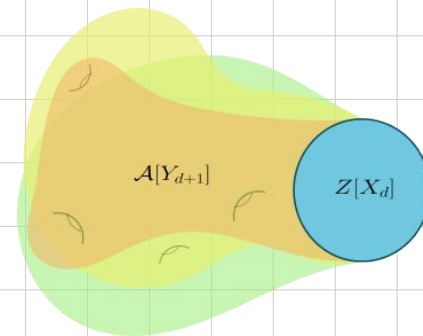
In QG, allow for topology-change  
 ⇒ Dai-Freed anomalies

# WHAT ARE THE BORDISM GROUPS?

What are the relevant bordism groups for these theories?

The background must satisfy:

$$dH \sim \text{tr}F^2 - \text{tr}R^2 = 0$$



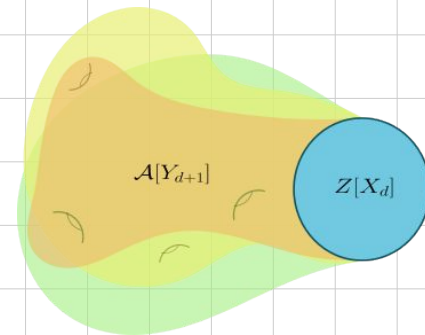
# WHAT ARE THE BORDISM GROUPS?

What are the relevant bordism groups for these theories?

The background must satisfy:

$$dH \sim \text{tr}F^2 - \text{tr}R^2 = 0$$

SO the relevant bordism groups for our 3 theories are **twisted-string bordism groups**



*(twisted-string) bordism groups are **not easy** to compute  
- Use Adams spectral sequence*

# WHAT ARE THE BORDISM GROUPS?

What are the relevant bordism groups for these theories?

The background must satisfy:

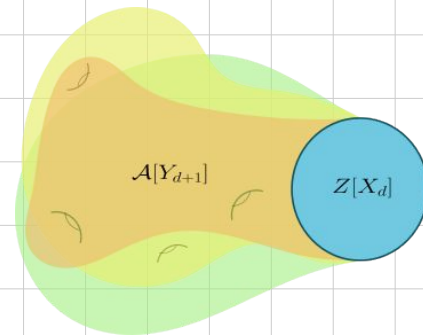
$$dH \sim \text{tr}F^2 - \text{tr}R^2 = 0$$

SO the relevant bordism groups for our 3 theories are **twisted-string bordism groups**

We find:

$$\underbrace{\Omega_{11}^{string-Sp(16)} = 0 \quad \Omega_{11}^{string-Spin(16)^2} = 0 \quad \Omega_{11}^{string-U(32)} = 0}_{\text{bordism groups are trivial} \Rightarrow \text{all global anomalies vanish!!!!}}$$

bordism groups are trivial  $\Rightarrow$  **all global anomalies vanish!!!!**



*(twisted-string) bordism groups are not easy to compute  
- Use Adams spectral sequence*

# WHAT ARE THE BORDISM GROUPS?

What are the relevant bordism groups for these theories?

The background must satisfy:

$$dH \sim \text{tr}F^2 - \text{tr}R^2 = 0$$

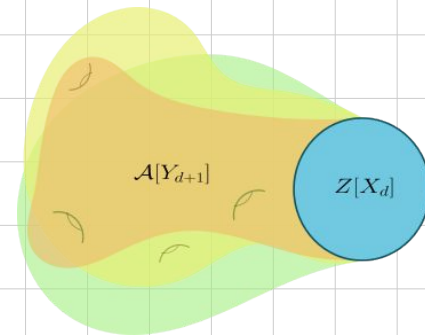
SO the relevant bordism groups for our 3 theories are **twisted-string bordism groups**

We find:

$$\underbrace{\Omega_{11}^{string-Sp(16)} = 0 \quad \Omega_{11}^{string-Spin(16)^2} = 0 \quad \Omega_{11}^{string-U(32)} = 0^*}_{\text{bordism groups are trivial} \Rightarrow \text{all global anomalies vanish!!!!}}$$

bordism groups are trivial  $\Rightarrow$  **all global anomalies vanish!!!!**

(up to a technical subtlety for the Sagnotti string)



*(twisted-string) bordism groups are not easy to compute  
- Use Adams spectral sequence*

# CONCLUSIONS

We showed there are **no global anomalies** for the three 10D **non-supersymmetric string theories**

**Huge consistency check!**

But that's not all...

- We used **anomaly inflow** to shed light on the chiral content of worldvolumes of branes in these theories
- We also computed **lower-dimensional cobordism** groups for these theories:

$$\begin{array}{ll}
 \Omega_0^{\text{String-Sp}(16)} \cong \mathbb{Z} & \Omega_6^{\text{String-Sp}(16)} \cong \mathbb{Z}_2 \\
 \Omega_1^{\text{String-Sp}(16)} \cong \mathbb{Z}_2 & \Omega_7^{\text{String-Sp}(16)} \cong \mathbb{Z}_4 \\
 \Omega_2^{\text{String-Sp}(16)} \cong \mathbb{Z}_2 & \Omega_8^{\text{String-Sp}(16)} \cong \mathbb{Z}^{\oplus 3} \oplus \mathbb{Z}_2 \\
 \Omega_3^{\text{String-Sp}(16)} \cong 0 & \Omega_9^{\text{String-Sp}(16)} \cong (\mathbb{Z}_2)^{\oplus 3} \\
 \Omega_4^{\text{String-Sp}(16)} \cong \mathbb{Z} & \Omega_{10}^{\text{String-Sp}(16)} \cong (\mathbb{Z}_2)^{\oplus 3} \\
 \Omega_5^{\text{String-Sp}(16)} \cong \mathbb{Z}_2 & \Omega_{11}^{\text{String-Sp}(16)} \cong 0.
 \end{array}$$

Example: Sugimoto

**No Global Symmetries** in QG implies that all of these bordism classes have to **trivialize in QG**.

[Cobordism Conjecture by McNamara, Vafa '19]

We can predict the existence of **new extended objects** that trivialize these classes!

... more info on arXiv:2310.06895

---

**THANKS!**

---