



# Gamma Factory and precision physics at LHC

WIESŁAW PŁACZEK



Jagiellonian University in Krakow

and

M. W. KRASNY



SORBONNE  
UNIVERSITÉ

Sorbonne University & CERN

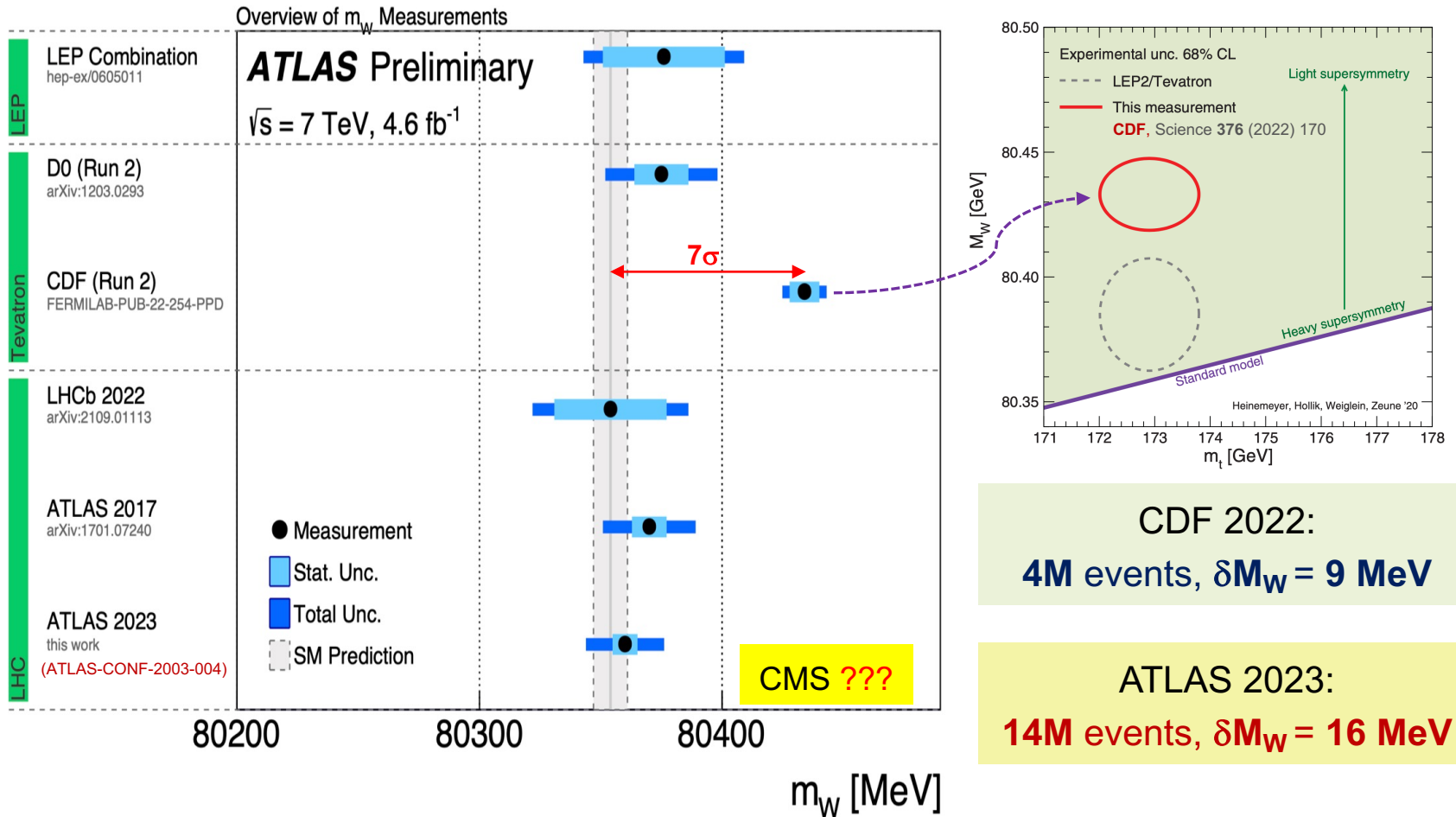


XXX Cracow EIPHANY Conference on Precision Physics  
at High Energy Colliders, 8–12 January 2024

dedicated to the memory of Staszek Jadach

**PRECISION MEASUREMENT  
OF STANDARD MODEL  
PARAMETERS AT LHC**

# W-boson mass measurement



CDF 2022:  
**4M events,  $\delta M_W = 9 \text{ MeV}$**

ATLAS 2023:  
**14M events,  $\delta M_W = 16 \text{ MeV}$**

- ATLAS 2023: improvement by  $\sim 15\%$  of ATLAS 2017  $M_W$  measurement with the **same data**, i.e. **2011** run at centre-of-mass energy of **7 TeV**.

# W-boson mass measurement

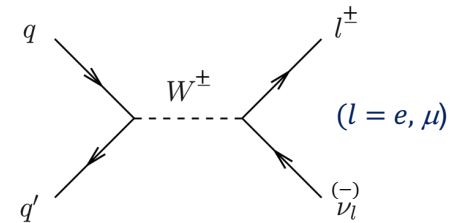
- **ATLAS** precision target:  $\delta M_W = 5 \text{ MeV}$  →  $\delta M_W/M_W < 0.01\%$   
(e.g. M. Boonecamp, Conference „Physics at LHC”, Krakow, 3–8 July 2006)

- Recent measurements:

- **CDF:**  $M_W = 80433.5 \pm 6.4 \text{ (stat.)} \pm 6.9 \text{ (syst.) MeV}$
- **ATLAS:**  $M_W = 80360 \pm 5 \text{ (stat.)} \pm 15 \text{ (syst.) MeV}$

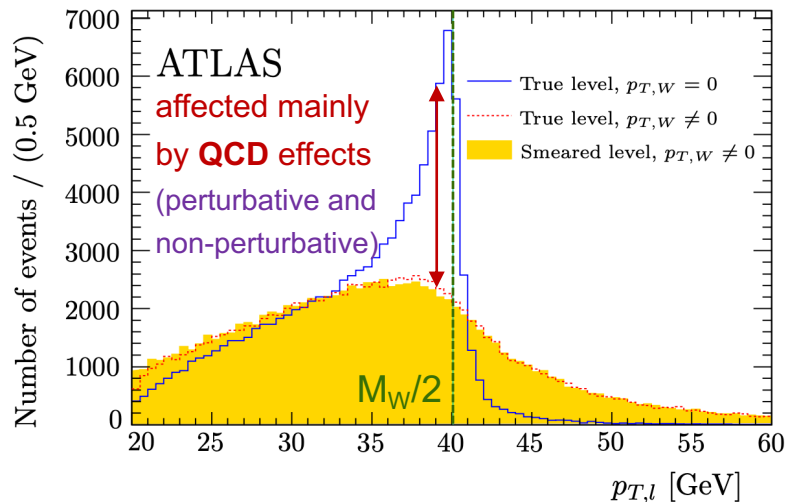
- Main observables (final-state neutrino not detected):

- charged-lepton transverse momentum  $p_{T,l}$
- leptons transverse mass  $m_T$

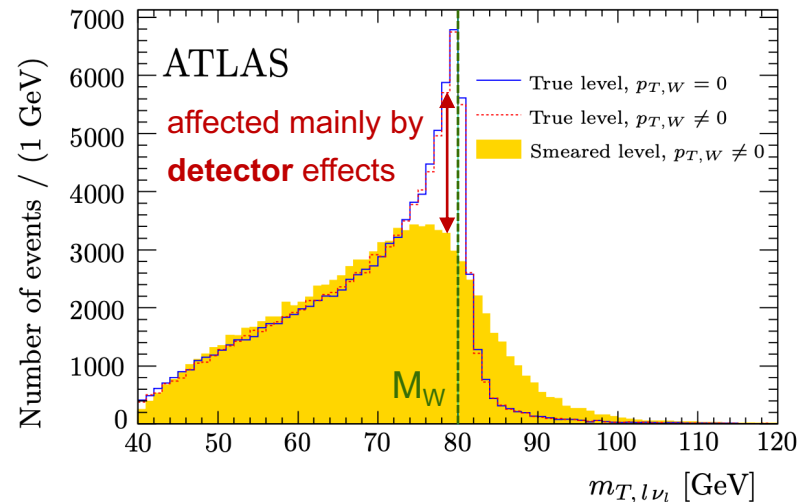


(a) - Charged lepton transverse momentum

(b) - Leptons transverse mass



→ theory uncertainties dominant!

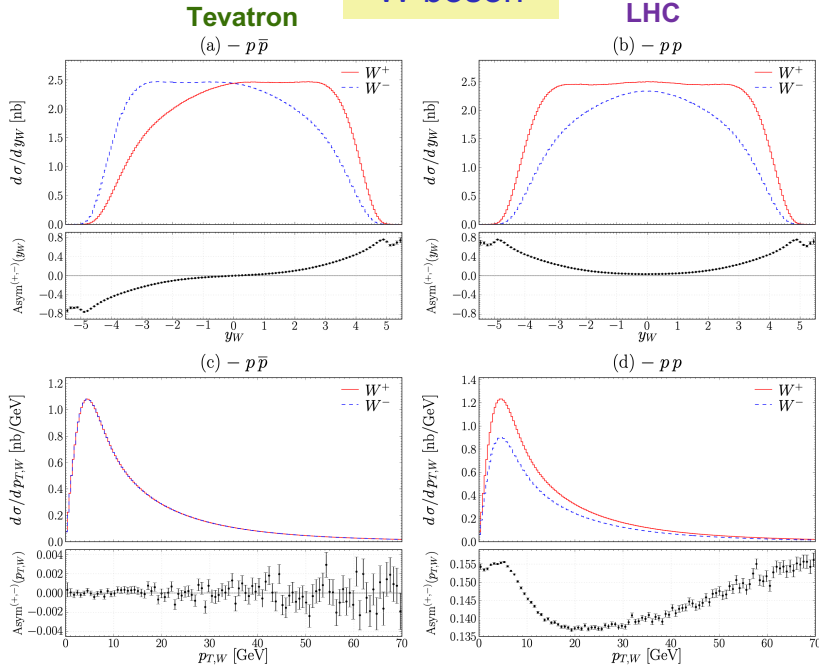


→ experimental uncertainties dominant!

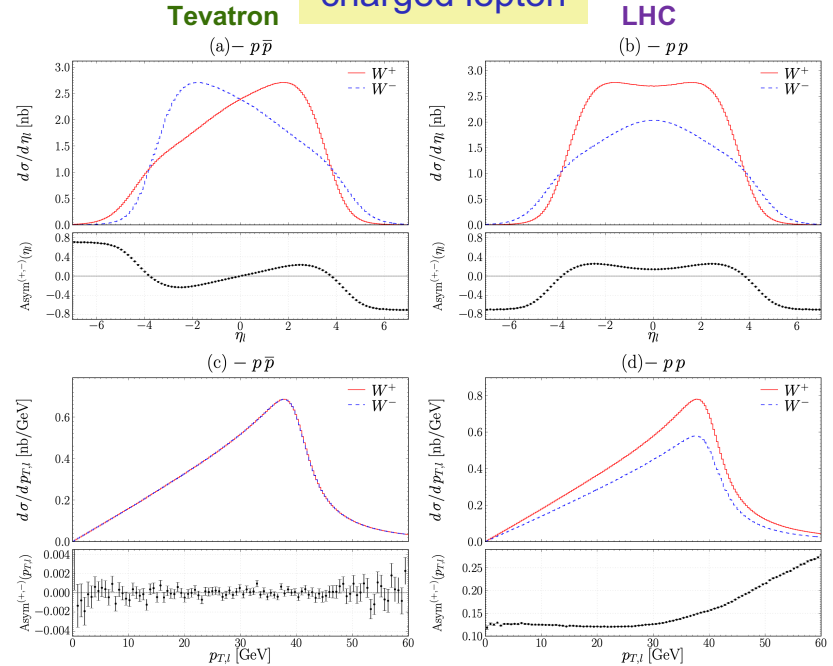


# Differences between Tevatron and LHC

## W boson



## charged lepton



	$W^+$					$W^-$				
Subprocesses	$u\bar{d}$	$u\bar{s}$	$c\bar{d}$	$c\bar{s}$	$c\bar{b}$	$d\bar{u}$	$s\bar{u}$	$d\bar{c}$	$s\bar{c}$	$b\bar{c}$
Tevatron II	90	2	1	7	0	90	2	1	7	0
LHC	74	4	1	21	0	67	2	3	28	0

	$Z$				
Subprocesses	$u\bar{u}$	$d\bar{d}$	$s\bar{s}$	$c\bar{c}$	$b\bar{b}$
Tevatron II	57	35	5	2	1
LHC	36	34	15	9	6

Flavour decomposition W and Z bosons in pp and p $\bar{p}$  collisions [%]

### □ Tevatron $\rightarrow$ distributions:

- $y_{W^+} + y_{W^-} \sim y_Z$  (V  $\leftrightarrow$  A couplings)
- $p_{T,W^+} + p_{T,W^-} \sim p_{T,Z}$  (quark flavours)
- Most of QCD/PDF uncertainties **cancel** in **ratios** of  $(W^+ + W^-)$  to Z distributions!

### □ LHC – the above not true!

- High precision theory predictions needed, particularly for  $p_{T,W}$

# Possible ways to reduce **systematic** uncertainties at LHC

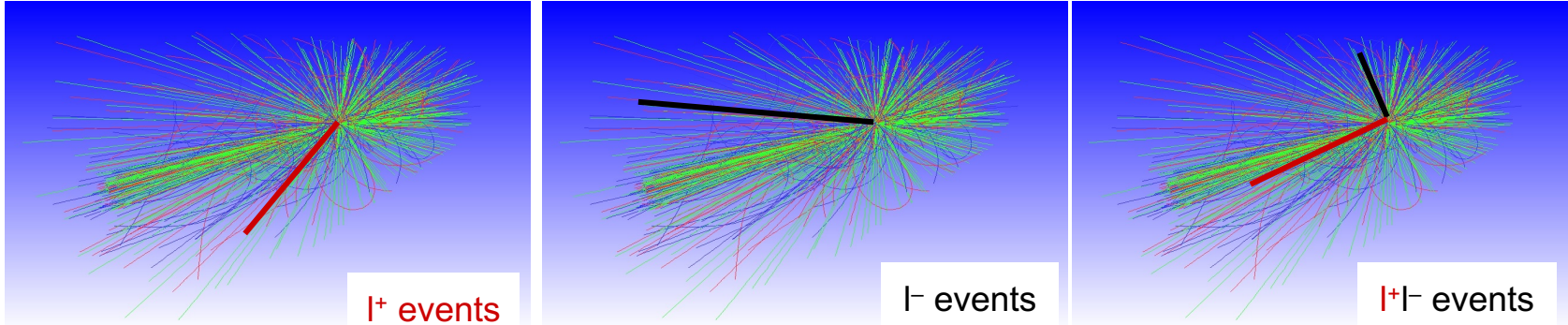
## ❖ Observables for simultaneous unfolding of partonic emittances and Standard Model parameters ( $W^+$ , $W^-$ and $Z$ boson production):

M.W. Krasny, [S. Jadach](#), WP, *Eur. Phys. J. C* 44 (2005) 333.

M.W. Krasny, F. Fayette, WP, A. Siódmok, *Eur. Phys. J. C* 51 (2007) 607.

F. Fayette, M.W. Krasny, WP, A. Siódmok, *Eur. Phys. J. C* 63 (2009) 33.

M.W. Krasny, F. Dydak, F. Fayette, WP, A. Siódmok, *Eur. Phys. J. C* 69 (2010) 379.



- $O_{W^+}(p_{T,l}, \eta_l) = \frac{d^2\sigma}{dp_{T,l^+}d\eta_{l^+}} ("l^+ \text{ events}"),$

- $O_{W^-}(p_{T,l}, \eta_l) = \frac{d^2\sigma}{dp_{T,l^-}d\eta_{l^-}} ("l^- \text{ events}"),$

- $O_{Z/\gamma}(M_{ll}, p_{T,u}, y_u) = \frac{d^3\sigma}{dM_{ll}dp_{T,u}dy_u} ("l^+l^- \text{ events}"),$

- $O_{(Z/\gamma)^+}(M_{ll}, p_{T,u}, y_u, p_{T,l}, \eta_l) = \frac{d^5\sigma}{dM_{ll}dp_{T,u}dy_u dp_{T,l^+}d\eta_{l^+}} ("l^+l^- \text{ events}"),$

- $O_{(Z/\gamma)^-}(M_{ll}, p_{T,u}, y_u, p_{T,l}, \eta_l) = \frac{d^5\sigma}{dM_{ll}dp_{T,u}dy_u dp_{T,l^-}d\eta_{l^-}} ("l^+l^- \text{ events}").$

} functions of  $u, d, s, c, \bar{u}, \bar{d}, \bar{s}, \bar{c}, M_W, \Gamma_W$

} functions of  $u, d, s, c, b, \bar{u}, \bar{d}, \bar{s}, \bar{c}, \bar{b}, M_Z, \Gamma_Z, \sin^2\theta_W, \alpha.$

➤ Partonic distributions:  $q(x, M_{W,Z}, k_T)$

□ MC simulations with WINHAC, WP and [S. Jadach](#) EPJC 29 (2003) 305.

# Precision observables

“almost” insensitive to strong-interaction effects

1. 
$$\text{Asym}_W^{(+,-)}(p_{T,l}, \eta_l) = \frac{O_{W^+}(p_{T,l}, \eta_l) - O_{W^-}(p_{T,l}, \eta_l)}{O_{W^+}(p_{T,l}, \eta_l) + O_{W^-}(p_{T,l}, \eta_l)},$$
  $\rightarrow$  sensitive to  $M_{W^+} - M_{W^-}$  and  $\Gamma_{W^+} - \Gamma_{W^-}$
2. 
$$\text{Asym}_Z^{(+,-)}(y_u, p_{T,u}, p_{T,l}, \eta_l) = \frac{O_{Z^+}(y_u, p_{T,u}, p_{T,l}, \eta_l) - O_{Z^-}(y_u, p_{T,u}, p_{T,l}, \eta_l)}{O_{Z^+}(y_u, p_{T,u}, p_{T,l}, \eta_l) + O_{Z^-}(y_u, p_{T,u}, p_{T,l}, \eta_l)},$$
  $\rightarrow$  sensitive to  $\sin^2\theta_W$
3. 
$$R_{WZ}(p_{T,l}, \eta_l) = \frac{O_{W^+}(p_{T,l}, \eta_l) + O_{W^-}(p_{T,l}, \eta_l)}{O_{Z^+}(p_{T,l}, \eta_l) + O_{Z^-}(p_{T,l}, \eta_l)},$$
  $\rightarrow$  sensitive to  $\alpha_s, M_{W^+} + M_{W^-}$ , and  $\Gamma_{W^+} + \Gamma_{W^-}$
4. 
$$R_Z^{norm}(p_{T,u}, y_u) = \frac{O_Z(p_{T,u}, y_u)}{O_{l^+l^-}^{norm}},$$

$$O_Z(p_{T,u}, y_u) = \int_{M_Z - 3\Gamma_Z}^{M_Z + 3\Gamma_Z} O_{(Z/\gamma)}(M_u, p_{T,u}, y_u) dM_u,$$

$$O_{Z^{+(-)}}(y_u, p_{T,u}, p_{T,l}, \eta_l) = \int_{M_Z - 3\Gamma_Z}^{M_Z + 3\Gamma_Z} [O_{(Z/\gamma)^{+(-)}}(M_u, y_u, p_{T,u}, p_{T,l}, \eta_l)] dM_u,$$

$\}$  only Z-peak region

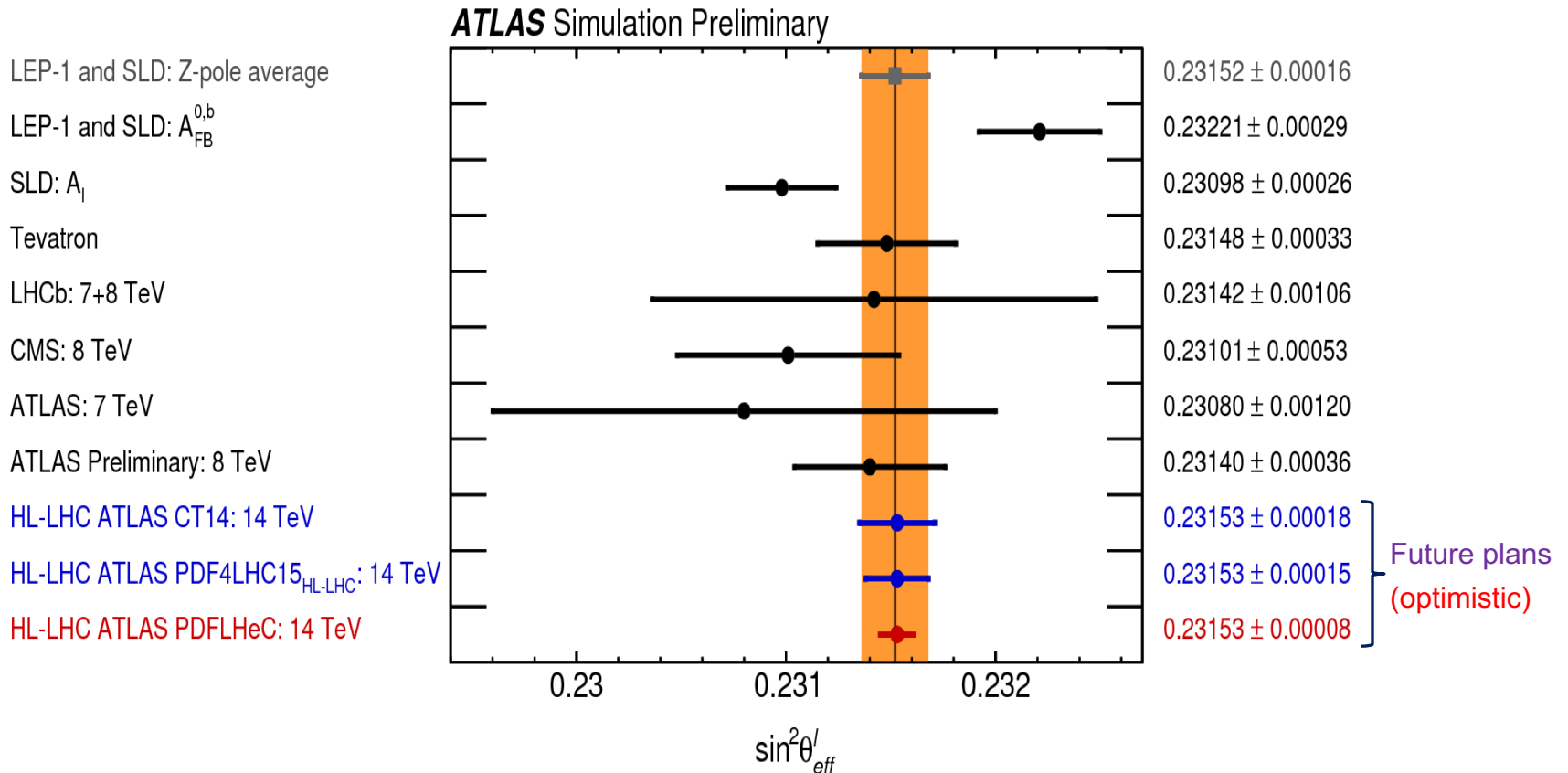
$$O_{l^+l^-}^{norm} = \int \int \int O_{Z/\gamma}(M_u, y_u, p_{T,u}) dM_u dy_u dp_{T,u} \longrightarrow \text{coplanar pairs – dedicated method of absolute normalization, M.W. Krasny et al., NIMA 584 (2008) 42}$$

- Measurements of **SM parameters** are based on **observables with reduced sensitivity to experimental systematic errors** (e.g. calibration of lepton momenta) and to **theoretical uncertainties of perturbative and non-perturbative QCD.**

# Unconstrained PDF degrees of freedom for pp collisions at LHC

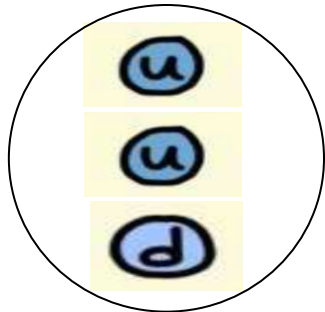
- ❖ Assuming  $s(x) = \bar{s}(x)$ ,  $c(x) = \bar{c}(x)$ ,  $b(x) = \bar{b}(x)$ :
  - ❑ 5 sea-quark flavours (u,d,s,c,b) + 2 valence-quark flavours ( $u_v, d_v$ ) = 7 unknown PDFs
  - ✓ 4 constraints coming from measurement of precision observables
  - $7 - 4 = 3$  degrees of freedom in flavour-dependent PDFs remain unconstrained at LHC:  $u_v(x) - d_v(x)$ ,  $c(x) - s(x)$ ,  $b(x)$ .
- Precision of  $\delta M_W < 12 \text{ MeV}$  cannot be achieved at LHC in pp runs (even in HL-LHC phase)
  - ✓ Missing external input: Tevatron data + dedicated muon-nucleon DIS experiment, e.g. at COMPASS (CERN)
- Similar conclusions for  $\sin^2\theta_W$  and  $\Gamma_W$ 
  - ❑ Note: in order to match  $\delta M_W \sim 5 \text{ MeV}$ , one needs  $\delta(\sin^2\theta_W) \sim 10^{-4}$  (indirect determination of  $M_W \rightarrow$  consistency test of SM).

# Measurement of $\sin^2\theta_W$



- **Most precise** measurement from **electron-positron** colliders LEP/SLD:  $\delta(\sin^2\theta_W) = 1.6 \times 10^{-4}$
- Precision at **hadron** colliders by factor **more than 2 worse** (the best from Tevatron).
- Main **limiting** factor at **LHC**: **PDF uncertainty** of  $u_v - d_v$ .

# Isoscalar ion beams



Proton

**u** and **d** quarks have different electric charges, weak isospin and vector and axial couplings.  
For EW-physics: proton beams are equivalent to beams of neutrinos and electrons mixed in not precisely known proportions.

In addition, relative distributions of valence and sea **u** and **d** quarks determine the effective W/Z boson polarisation. Proton beams → polarisation cannot be precisely controlled.

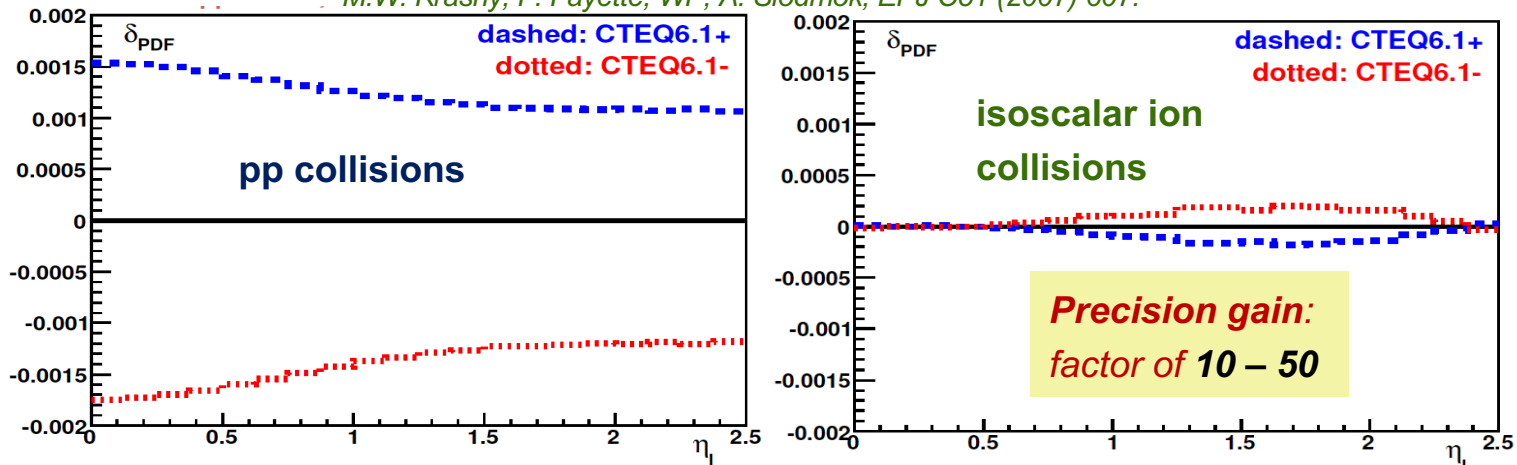
- ❑ **Isoscalar ion beams** (# of protons = # of neutrons:  $A = 2Z$ ) profit from **flavour symmetry** of strong interactions to **equalize** distributions of **u** and **d** quarks:

$$u_{v,s}^{A=2Z,Z}(x, k_t, Q^2) = d_{v,s}^{A=2Z,Z}(x, k_t, Q^2)$$

# Significantly **higher precision** in measuring **EW processes** by using **isoscalar-ion** rather than proton beams at LHC

- **Uncertainty** reflecting the limited knowledge of the **partonic** composition of **protons** and **isoscalar ions** in measurement of  $R_{WZ}(p_{T,l}, \eta_l)$  observable (**precision-optimised method**).

*M.W. Krasny, F. Fayette, WP, A. Siodmok, EPJ C51 (2007) 607.*



**Isoscalar beams provide a unique way for improving the present precision of SM parameters, such as  $M_W$ ,  $\sin^2\theta_W$ , ... at the LHC!**

- **High-statistics** ( $> 100 \text{ fb}^{-1}$ ) data from isoscalar-beam collisions needed!

**How to get such a luminosity at LHC?**

- M.W. Krasny, A. Petrenko and WP, *Prog. Part. Nucl. Phys.* 114 (2020) 103792.
- For nucleon-nucleon luminosity of **1000  $\text{fb}^{-1}$** :  **$\delta M_W < 5 \text{ MeV}$ ,  $\delta(\sin^2\theta_W) < 10^{-4}$**

# GAMMA FACTORY



# Gamma Factory studies

## The Gamma Factory proposal for CERN †

† An Executive Summary of the proposal addressed to the CERN management.

Mieczyslaw Witold Krasny\*

LPNHE, Universités Paris VI et VII and CNRS-IN2P3, Paris, France

e-Print: [1511.07794](https://arxiv.org/abs/1511.07794) [hep-ex]

**~100 physicists from ~40 institutions have contributed so far to the Gamma Factory studies**

A. Abramov<sup>1</sup>, A. Afanasev<sup>37</sup>, S.E. Alden<sup>1</sup>, R. Alemany Fernandez<sup>2</sup>, P.S. Antsiferov<sup>3</sup>, A. Apyan<sup>4</sup>, G. Arduini<sup>2</sup>, D. Balabanski<sup>34</sup>, R. Balkin<sup>32</sup>, H. Bartosik<sup>2</sup>, J. Berengut<sup>5</sup>, E.G. Bessonov<sup>6</sup>, N. Biancacci<sup>2</sup>, J. Bieroń<sup>7</sup>, A. Bogacz<sup>8</sup>, A. Bosco<sup>1</sup>, T. Brydges<sup>36</sup>, R. Bruce<sup>2</sup>, D. Budker<sup>9,10</sup>, M. Bussmann<sup>38</sup>, P. Constantin<sup>34</sup>, K. Cassou<sup>11</sup>, F. Castelli<sup>12</sup>, I. Chaikovska<sup>11</sup>, C. Curatolo<sup>13</sup>, C. Curceanu<sup>35</sup>, P. Czodrowski<sup>2</sup>, A. Derevianko<sup>14</sup>, K. Dupraz<sup>11</sup>, Y. Duthéil<sup>2</sup>, K. Dzierżęga<sup>7</sup>, V. Fedosseev<sup>2</sup>, V. Flambaum<sup>25</sup>, S. Fritzsche<sup>17</sup>, N. Fuster Martinez<sup>2</sup>, S.M. Gibson<sup>1</sup>, B. Goddard<sup>2</sup>, M. Gorshteyn<sup>20</sup>, A. Gorzawski<sup>15,2</sup>, M.E. Granados<sup>2</sup>, R. Hajima<sup>26</sup>, T. Hayakawa<sup>26</sup>, S. Hirlander<sup>2</sup>, J. Jin<sup>33</sup>, J.M. Jowett<sup>2</sup>, F. Karbstein<sup>39</sup>, R. Kersevan<sup>2</sup>, M. Kowalska<sup>2</sup>, M.W. Krasny<sup>16,2</sup>, F. Kroeger<sup>17</sup>, D. Kuchler<sup>2</sup>, M. Lamont<sup>2</sup>, T. Lefevre<sup>2</sup>, T. Ma<sup>32</sup>, D. Manglunki<sup>2</sup>, B. Marsh<sup>2</sup>, A. Martens<sup>12</sup>, C. Michel<sup>40</sup>, S. Miyamoto<sup>31</sup>, J. Molson<sup>2</sup>, D. Nichita<sup>34</sup>, D. Nutarelli<sup>11</sup>, L.J. Nevay<sup>1</sup>, V. Pascalutsa<sup>28</sup>, Y. Papaphilippou<sup>2</sup>, A. Petrenko<sup>18,2</sup>, V. Petrillo<sup>12</sup>, L. Pinard<sup>40</sup>, W. Placzek<sup>7</sup>, R.L. Ramjiawan<sup>2</sup>, S. Redaelli<sup>2</sup>, Y. Peinaud<sup>11</sup>, S. Pustelny<sup>7</sup>, S. Rochester<sup>19</sup>, M. Safronova<sup>29,30</sup>, D. Samoilenko<sup>17</sup>, M. Sapinski<sup>20</sup>, M. Schaumann<sup>2</sup>, R. Scrivens<sup>2</sup>, L. Serafini<sup>12</sup>, V.P. Shevelko<sup>6</sup>, Y. Soreq<sup>32</sup>, T. Stoehlker<sup>17</sup>, A. Surzhykov<sup>21</sup>, I. Tolstikhina<sup>6</sup>, F. Velotti<sup>2</sup>, A. Viatkina<sup>9</sup>, A.V. Volotka<sup>17</sup>, G. Weber<sup>17</sup>, W. Weiqiang<sup>27</sup>, D. Winters<sup>20</sup>, Y.K. Wu<sup>22</sup>, C. Yin-Vallgren<sup>2</sup>, M. Zanetti<sup>23,13</sup>, F. Zimmermann<sup>2</sup>, M.S. Zolotarev<sup>24</sup> and F. Zomer<sup>11</sup>



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Special Issue: Physics Opportunities with the Gamma Factory

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Issue Edited by: Dmitry Budker, Mikhail Gorchtein, Mieczyslaw Witold Krasny, Adriana Pálffy, Andrey Surzhykov

**Gamma Factory studies are anchored and supported by the CERN **Physics Beyond Colliders (PBC)** framework.**

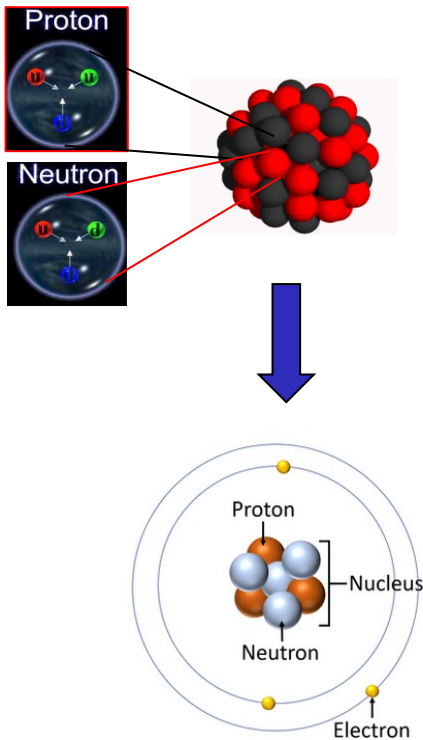
More info on all the GF group activities:

<https://indico.cern.ch/category/10874>

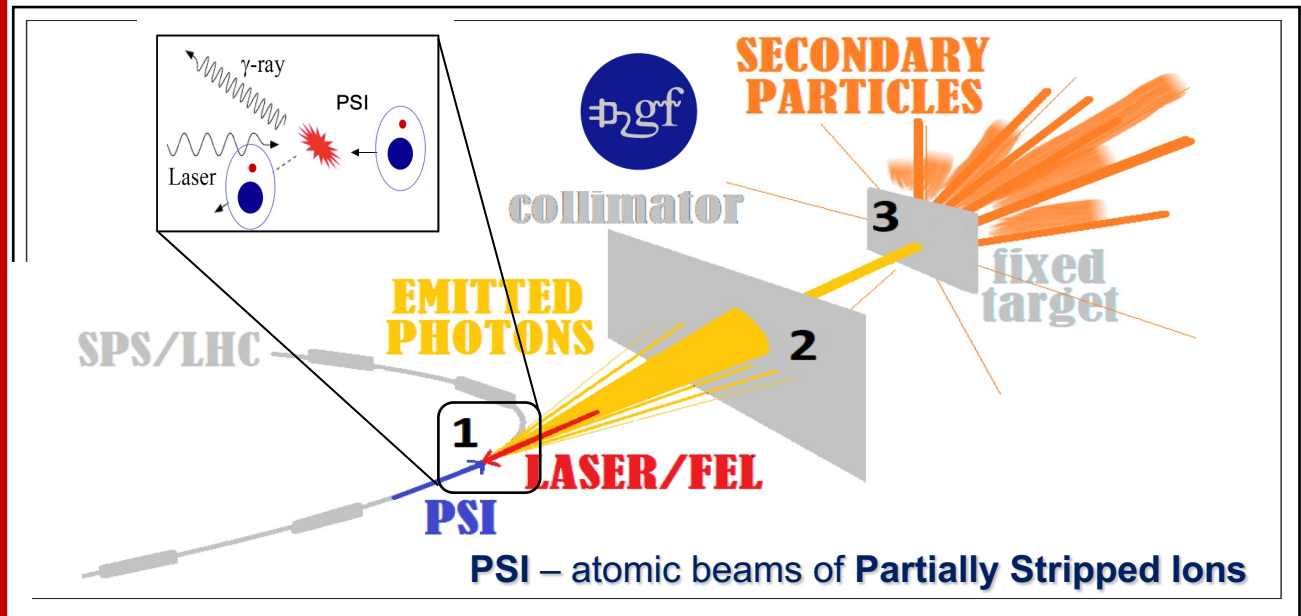
*We acknowledge the crucial role of the CERN **PBC** framework in bringing our accelerator tests, GF-PoP experiment design, software development and physics studies to their present stage!*

# Gamma Factory beams

## LHC beams



- Include **atomic beams of partially stripped ions** in the LHC menu.
- Collide them with laser pulses (**circulating in Fabry-Pérot resonators**) to **produce beams of polarised photons** and secondary beams of polarised electrons/positrons/muons, neutrons and radioactive ions.



# Extraordinary properties GF photon source

## 1. Point-like, small divergence:

- $\Delta z \sim l_{PSI-bunch} < 7 \text{ cm}$ ;  $\Delta x, \Delta y \sim \sigma^{PSI}_x, \sigma^{PSI}_y < 50 \text{ }\mu\text{m}$ ;  $\Delta(\theta_x), \Delta(\theta_y) \sim 1/\gamma_L < 1 \text{ mrad}$

## 2. Huge jump in intensity:

- **More than 7 orders of magnitude** w.r.t. existing (being constructed)  $\gamma$ -sources: up to  $10^{18} \text{ }\gamma/\text{s}$

## 3. Very wide range of tuneable energy photon beam:

- **10 keV – 400 MeV** – extending by a factor of **~1000** the energy range of FEL photon sources

## 4. Tuneable polarisation:

- **Polarisation transmission** from laser photons to  $\gamma$ -beams of **up to 99%**

## 5. Unprecedented plug power efficiency (energy footprint):

- **LHC RF power can be converted to the photon beam power.** Wall-plug power efficiency of the **GF photon source is by a factor of ~300 better** than that of **DESY-XFEL** (assuming power consumption: 200 MW at CERN and 19 MW at DESY)

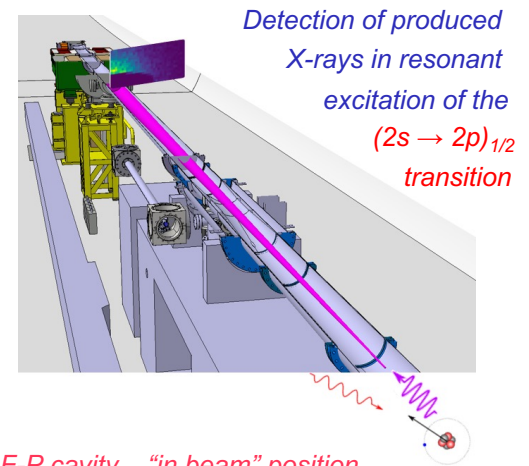
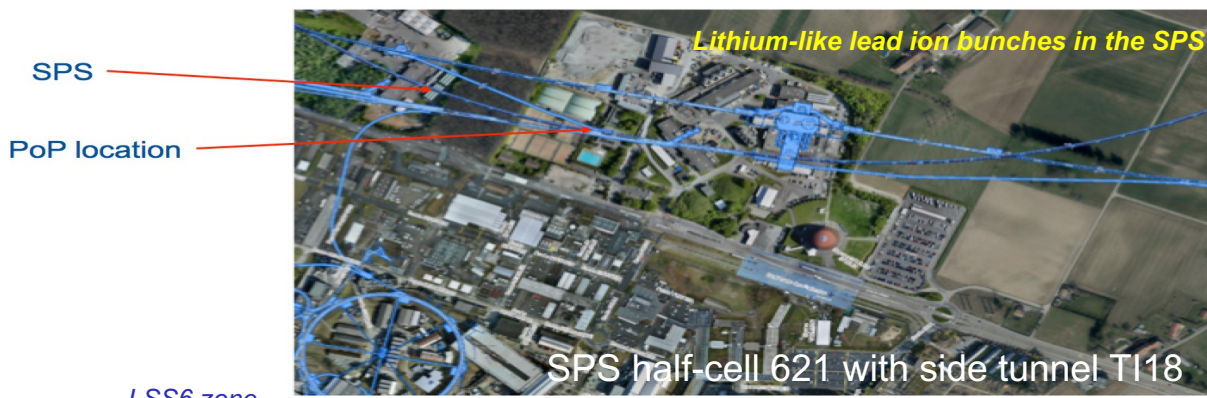
# Tertiary beams' sources – intensity/quality targets

- **Polarised positrons** – potential gain of up to a factor of  $10^4$  in intensity w.r.t. the KEK positron source, satisfying both the LEMMA (for muon collider) and the LHeC requirements; → A. Apyan, M.W. Krasny and WP, *Phys. Rev. Accel. Beams* **26**, 083401 (2023).
- **Pions** – potential, gain by a factor of  $10^3$ , gain in the spectral density ( $d^2N/dpdp_T$ ) with respect to proton-beam-driven sources at KEK and FNAL → *PRAB* **26**, 083401 (2023).
- **Muons** – potential gain by a factor of  $10^3$  in intensity w.r.t. the PSI (Villigen, Switzerland) muon source, charge symmetry ( $N_{\mu^+} \sim N_{\mu^-}$ ), polarisation control → *PRAB* **26**, 083401 (2023).
- **Neutrinos** – fluxes comparable to the proposed NuMAX (FNAL) but: (1) **Very Narrow-Band Beam**, driven by the small spectral density pion beam and (2) unique possibility of creating **flavour and CP-tuned beams** driven by the beams of polarised muons.
- **Neutrons** – potential gain of up to a factor of  $10^4$  in intensity of primary MeV-energy neutrons per 1 MW of the driver beam power.
- **Radioactive ions** – potential gain of up to a factor of  $10^4$  in intensity w.r.t. e.g. ALTO (Orsay).

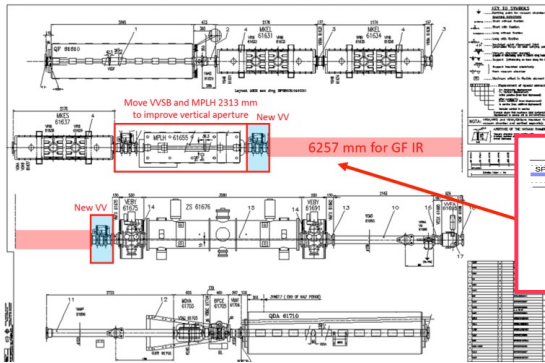
# Examples of potential applications domains of **Gamma Factory** research tools

- **Particle physics:** precision QED and **EW studies**, vacuum birefringence, **Higgs physics in  $\gamma\gamma$  collision mode with laser-cooled ion beams**, rare muon decays, precision neutrino physics, QCD-confinement studies, ...
- **Nuclear physics:** nuclear spectroscopy, cross-talk of nuclear and atomic processes, GDR, nuclear photo-physics, photo-fission research, gamma polarimetry, physics of rare radioactive nuclides,...
- **Atomic physics:** highly charged atoms, electronic, muonic, pionic and kaonic atoms, ...
- **Astrophysics:** dark matter searches, gravitational waves detection, gravitational effects of cold particle beams,  $^{16}\text{O}(\gamma,\alpha)^{12}\text{C}$  reaction and S-factors ...
- **Fundamental physics:** studies of the basic symmetries of the universe, atomic interferometry, ...
- **Accelerator physics:** beam cooling techniques, low-emittance hadronic beams, plasma wake field acceleration, high-intensity polarised positron and muon sources, beams of radioactive ions and neutrons, very narrow band and flavour-tagged neutrino beams, neutron sources, ...
- **Applied physics:** accelerator driven energy sources, nuclear fusion research, medical isotope and isomer production, ...

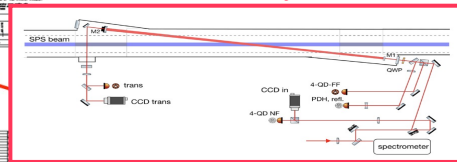
# Gamma Factory Proof-of-Principle (PoP) experiment at SPS



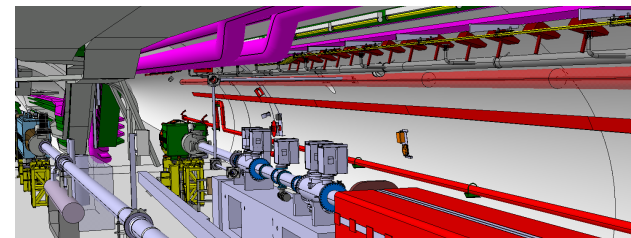
LSS6 zone



F-P cavity



F-P cavity – “in beam” position



F-P cavity length: 3.75 m – vertically tilted by 2.6 deg



# Purpose of GF SPS PoP experiment

1

Demonstrate that an adequate laser system (5 mJ @ 40 MHz) can be (remotely) operated in the high-radiation field of the SPS.

2

Demonstrate that very high rates of photons are produced: almost all ions are excited in a single collision of a PSI bunch with a laser pulse.

3

Demonstrate stable and repeatable operation.

4

Confront data and simulations.

5

Demonstrate **laser ion-beam cooling**: longitudinal and transverse.

6

New atomic physics measurements.

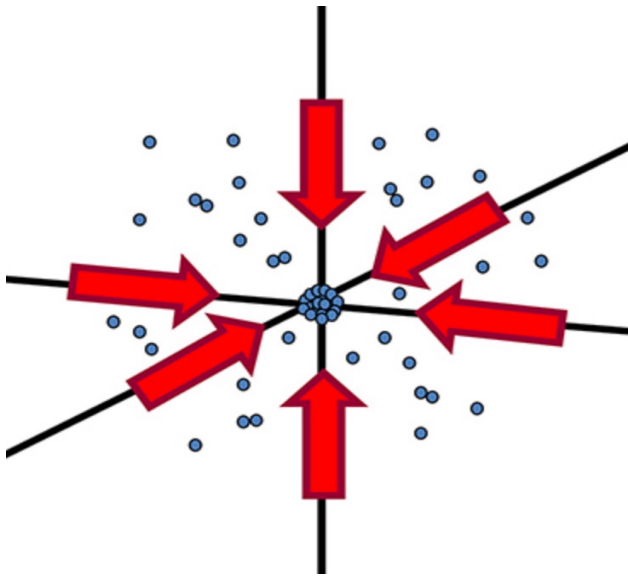
➤ *Target installation time: LS3 (2026–2027).*

➤ *Estimated cost: 2.5 MCHF*

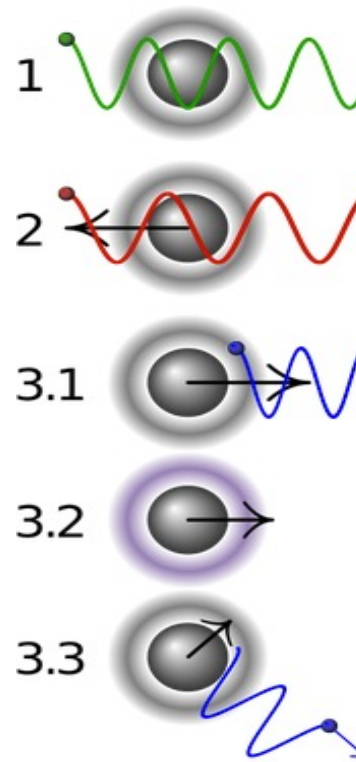
# LASER COOLING OF HIGH-ENERGY ION BEAMS



# Doppler laser-cooling of stationary atoms



Six “red-detuned” laser beams  
(optical molasses)



A stationary atom sees the laser neither red- nor blue-shifted and does not absorb the photon.

An atom moving away from the laser sees it red-shifted and does not absorb the photon.

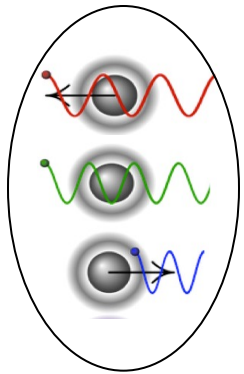
An atom moving towards the laser sees it blue-shifted and absorbs the photon, slowing the atom.

The photon excites the atom, moving an electron to a higher quantum state.

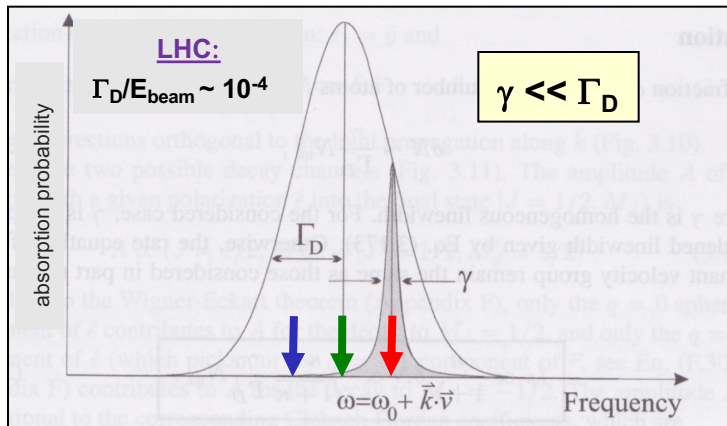
The atom re-emits a photon.

As its direction is random, there is no net change in momentum over many absorption-emission cycles.

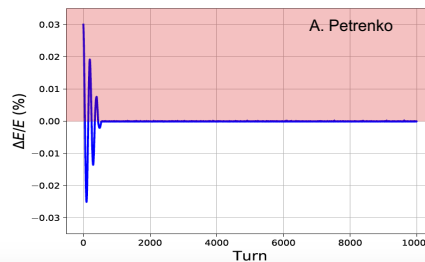
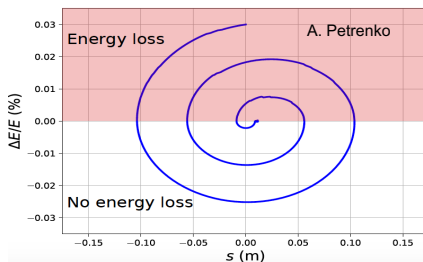
# Doppler laser-cooling methods of high-energy beams



Bunch



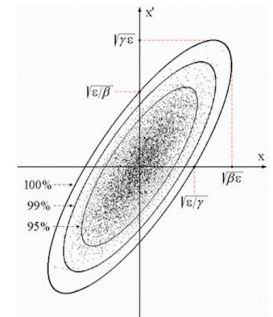
Opens a possibility of forming at CERN hadronic beams of the required **longitudinal** and **transverse emittances** within a seconds-long time scale.



**Beam cooling speed:** the laser wavelength band is chosen such that only the ions moving in the laser-pulse direction (in the bunch rest frame) can resonantly absorb photons.

**Beam emittance:** is the volume in the phase space occupied by the particle beam. The space is 6-dimensional  $(x, x', y, y', W, \phi)$  but are important the projections in the 3 dimensions: transverse emittance in x and y  $(\epsilon_x, \epsilon_y)$  and longitudinal emittance  $(\Delta W, \phi)$ .

**Liouville theorem:** under the action of linear forces, the beam emittance is constant. (BUT: in presence of non-linear forces, it increases). Consequence: in an accelerator, where non-linear forces are often present, the emittance will progressively increase.



# Gamma Factory path to high-luminosity LHC

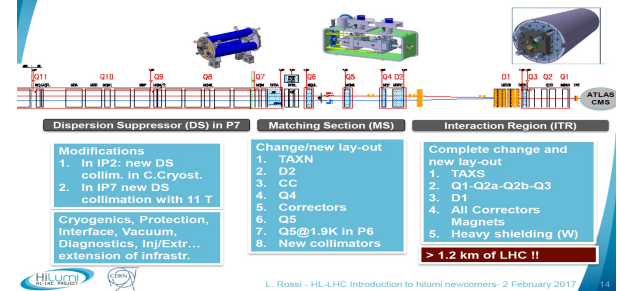
$$\mathcal{L} = f \frac{n_1 n_2}{4\pi \sqrt{\epsilon_x \beta_x^* \epsilon_y \beta_y^*}}$$

## Two complementary ways to increase collider luminosity:

1. Increase the focusing strength:  $\beta^*$
  2. Reduce the beam emittance:  $\epsilon$
- Both

A **low-emittance** particle beam is the beam where **particles** are confined within **small distances** and have nearly the **same momentum** vectors – **cold beams**.

## The on-going HL(pp)-LHC project



Levelled luminosity:  $2.5 (5) \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ , cost ~ 1 billion euro

Progress in Particle and Nuclear Physics  
 ELSEVIER Volume 114, September 2020, 103792

Review  
**High-luminosity Large Hadron Collider with laser-cooled isoscalar ion beams** ☆

M.W. Krasny <sup>a, b</sup>, A. Petrenko <sup>c, b</sup>, W. Płaczek <sup>d</sup>

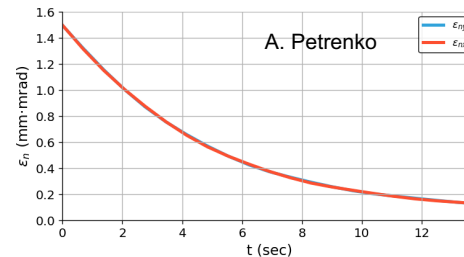
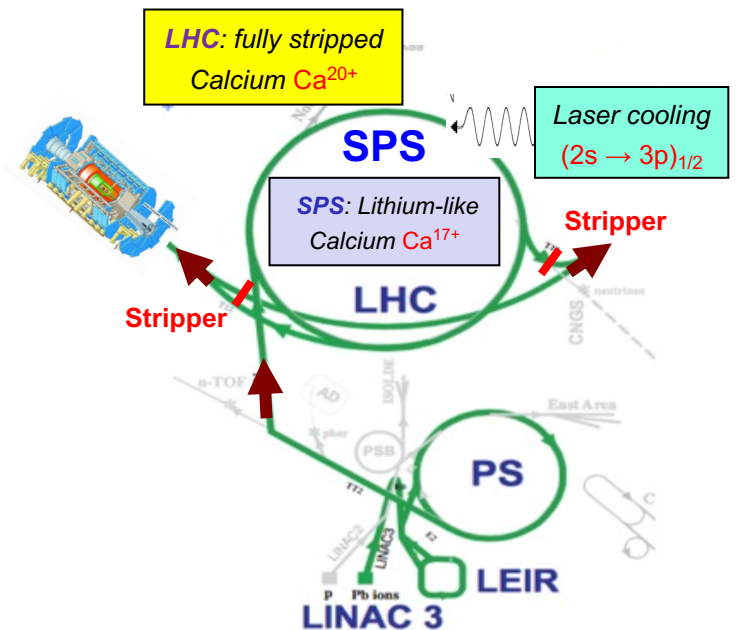
Show more ▼

<https://doi.org/10.1016/j.ppnp.2020.103792> Get rights and content

# GF scheme of reducing transverse beam emittance

1. **Produce highly-charged-ion** bunches (partially stripped ions) with existing CERN ion source.
2. **Leave a few electrons** attached to their parent nuclei for **SPS** acceleration phase (in the canonical SPS heavy-ion operation all electrons are already stripped off).
3. **Cool atomic beams** with specialised **laser system** at the top SPS energy to reduce its emittance (longitudinal and transverse cooling).
4. **Strip electrons** in **SPS-to-LHC** transfer line.
5. **Accelerate** and **collide fully stripped ion** beams in **LHC**.

# Gamma Factory path to HL(AA)-LHC: A concrete implementation scheme with Ca beams



**Reduction of transverse (x, y) emittances** by a factor of **5** can be achieved in **8 seconds** – sufficiently short to avoid  $\text{Ca}^{17+}$  beam losses in SPS.

Optical stochastic cooling time for Ca beam, if necessary, at the top energy ~ 1.5 hours (V. Lebedev).

Parameter	Value
$s^{1/2}$ [TeV]	7
$\sigma_{BFPP}(\text{Ca})/\sigma_{BFPP}(\text{Pb})$	$5 \times 10^{-5}$
$\sigma_{had}(\text{Ca})/\sigma_{tot}(\text{Ca})$	0.6
$N_b$	$3 \times 10^9$
$\epsilon_{(x,y)n}$ [ $\mu\text{m}$ ] <sup>(1)</sup>	<b>0.3</b>
IBS [h]	1–2
$\beta^*$ [m]	0.15
$L_{NN}$ [ $\text{cm}^{-2}\text{s}^{-1}$ ]	<b><math>4.2 \times 10^{34}</math></b>
No of bunches	1404
Collisions/beam crossing	<b>5.5</b>

**~700 for pp** for the same luminosity  $L_{NN}$

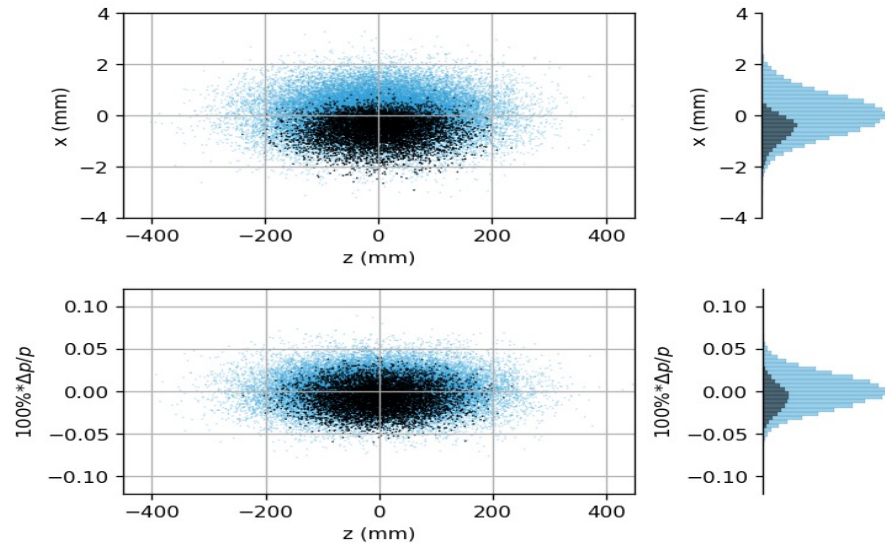
# Simulation of laser cooling of lithium-like $\text{Ca}^{17+}$ bunches in SPS – input parameters

Table 1: Parameters of the calcium-beam cooling configuration in the SPS.

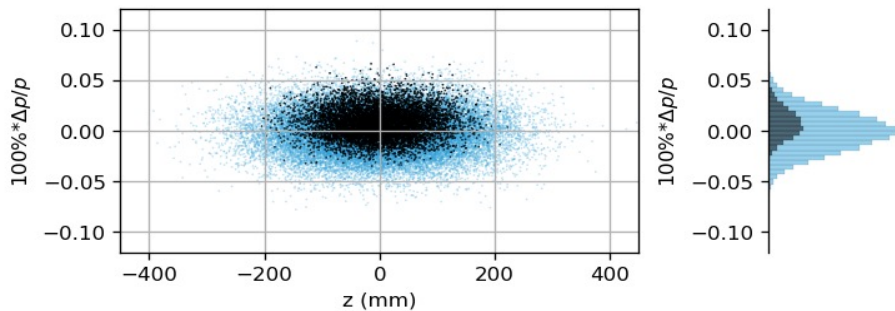
Ion beam	$^{40}\text{Ca}^{17+}$
$m$ – ion mass	37.21 GeV/ $c^2$
$E$ – mean energy	7.65 TeV
$\gamma_L = E/mc^2$ – mean Lorentz relativistic factor	205.62
$N$ – number ions per bunch	$4 \times 10^9$
$\sigma_E/E$ – RMS relative energy spread	$2 \times 10^{-4}$
$\epsilon_n$ – normalised transverse emittance	1.5 mm mrad
$\sigma_x$ – RMS transverse size	0.80 mm
$\sigma_y$ – RMS transverse size	0.57 mm
$\sigma_z$ – RMS bunch length	10 cm
Dispersion function	2.44 m
Laser	pulsed Ti:Sa (20MHz)
$\lambda$ – wavelength ( $\hbar\omega$ – photon energy)	768 nm (1.6 eV)
$\sigma_\lambda/\lambda$ – RMS relative band spread	$2 \times 10^{-4}$
$U$ – single pulse energy at IP	2 mJ
$\sigma_L$ – RMS transverse intensity distribution at IP ( $\sigma_L = w_L/2$ )	0.56 mm
$\sigma_t$ – RMS pulse duration	2.04 ps
$\theta_L$ – collision angle	1.3 deg
Atomic transition of $^{40}\text{Ca}^{17+}$	$2s \rightarrow 3p$
$\hbar\omega'_0$ – resonance energy	661.89 eV
$\tau'$ – mean lifetime of spontaneous emission	0.4279 ps
$\hbar\omega_1^{\text{max}}$ – maximum emitted photon energy	271 keV



# Simulation of laser cooling of lithium-like $\text{Ca}^{17+}$ bunches in SPS: laser pulse tuning



**Horizontal** positions and **momenta** of ions interacting with pulse of **first laser** in momentum-dispersive region (**black dots** – excited ions, **blue dots** – non-excited ions). **Shift** of laser pulse by **1.4 mm** provides **optimal coupling** of horizontal **betatron** oscillations to **synchrotron** oscillations (**dispersive coupling**, I. Lauer *et al.*, *Phys. Rev. Lett.* 81 (1998) 2052). **17%** of all ions are excited in each bunch crossing!



**Momentum** and **longitudinal** positions of the ions interacting with the photon-pulse of the **second laser**. The laser pulse focal point is aligned with the ion beam centre, but its **frequency band** is **shifted** to excite the **higher-momentum** ions.

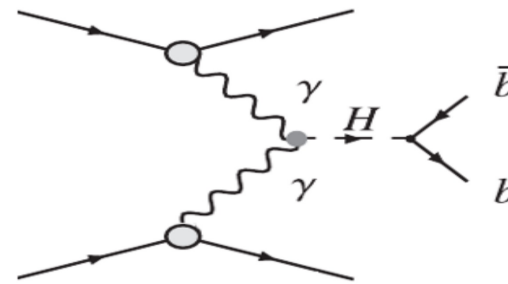
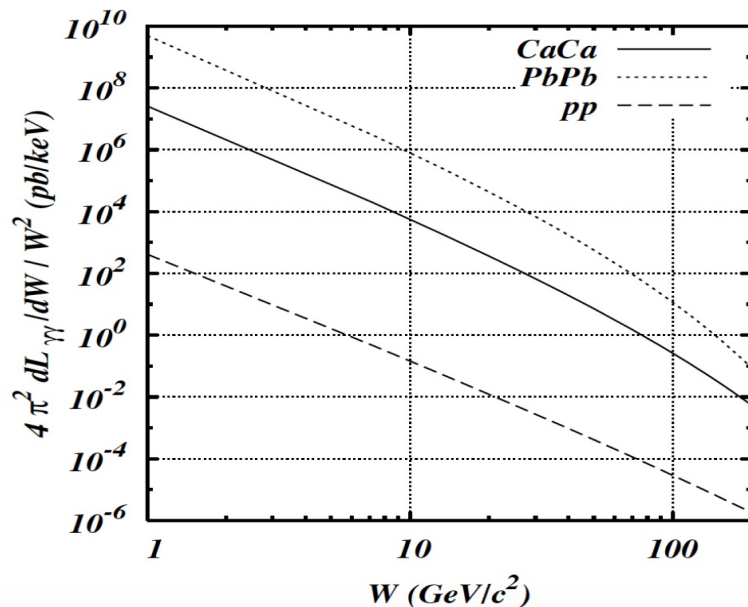
# Merits of cold isoscalar beams

- 1. Impact of modelling uncertainties of partonic emittances** (longitudinal and transverse) on achievable **measurement precision** can be drastically **reduced** and controlled by LHC data alone (no precision brick-walls coming from LHC-external data, PDFs, and PS models).
  - Significantly **higher systematic precision** in measuring **SM parameters** ( $M_W$ ,  $\sin^2\theta_W$ , ...) **than for proton beams**.
- 2.  $Z^4$  leap in photon fluxes** – access to **exclusive Higgs boson production in photon–photon collisions** – **unreachable for the pp running mode**.
- 3. Lower pile-up background** at the equivalent (high) nucleon-nucleon luminosity.
- 4. New research opportunities** for the EW-symmetry breaking sector.



# Leap in photon fluxes – access to **exclusive Higgs-boson production in photon–photon collisions** – beyond reach of pp running mode

Factor of  $Z^4$  gain in photon-photon luminosity



For collected **nucleon-nucleon luminosity** of  $1000 \text{ fb}^{-1}$ , expected **number of exclusively produced Higgs bosons in photon-photon collisions** is  **$\sim 420$  per experiment** → **unique possibility** for studies of  $\gamma\gamma \rightarrow H \rightarrow b\bar{b}$  in **clean environment**.

Higgs boson production in photon-photon interactions with proton, light-ion, and heavy-ion beams at current and future colliders

David d'Enterria (CERN), Daniel E. Martins (UFRJ, Rio de Janeiro), Patricia Rebello Teles (Rio de Janeiro State U.) (Apr 26, 2019)

Published in: *Phys.Rev.D* 101 (2020) 3, 033009 • e-Print: [1904.11936](https://arxiv.org/abs/1904.11936) [hep-ph]

[50] G. Baur, K. Hencken, D. Trautmann, S. Sadovsky and Y. Kharlov, Phys. Rept. **364** (2002) 359, doi:10.1016/S0370-1573(01)00101-6 [hep-ph/0112211].

# CONCLUSIONS

# Gamma Factory in general

- ❑ **Gamma Factory** can create, at CERN, a variety of **novel research tools** which could open novel research opportunities in a **very broad domain of basic and applied science**.
- ❑ Its **“quest for diversity of research subjects and communities”** is of particular **importance** in the **present** phase of **accelerator-based** research, as we **neither** have any solid **theoretical** guidance for a **new physics** “just around the corner”, accessible by FCC or ILC/CLIC, **nor** an established **“reasonable-cost” technology** for a leap into very high-energy “terra incognita”.
- ❑ The **Gamma Factory** research programme can be **largely based on the existing CERN accelerator infrastructure** – it requires “relatively” minor infrastructure investments.

# Gamma Factory for HL-LHC

- ❖ **HL(AA)-LHC with laser-cooled isoscalar ion beams** appears to be attractive option to **significantly enlarge research potential of LHC**.
- ❖ Achievable nucleon-nucleon **luminosities** in **HL-LHC** modes:
  - $L_{NN}$  in **pp** mode  $2.5 - 5 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
  - $L_{NN}$  in **AA** mode  $\sim 4 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
  - *Optimal running-time sharing: 2/3 in **pp** mode, 1/3 in **AA** mode.*
- ❖ Of particular **importance**, if **no new phenomena** are discovered over **initial** phase of **HL(pp)-LHC** operation.
- ❖ It provides also **a new way to maximise partonic luminosities** at **future high-energy frontier hadron colliders**, such as **FCC-hh**.

# BACKUP SLIDES

# Sources of uncertainties in $M_W$ measurements

## ATLAS 2023

Obs.	Mean [MeV]	Elec. Unc.	PDF Unc.	Muon Unc.	EW Unc.	PS & $A_i$ Unc.	Bkg. Unc.	$\Gamma_W$ Unc.	MC stat. Unc.	Lumi Unc.	Recoil Unc.	Total sys.	Data stat.	Total Unc.
$p_T^\ell$	80360.1	8.0	7.7	7.0	6.0	4.7	2.4	2.0	1.9	1.2	0.6	15.5	4.9	16.3
$m_T$	80382.2	9.2	14.6	9.8	5.9	10.3	6.0	7.0	2.4	1.8	11.7	24.4	6.7	25.3

## CDF 2022

Table 2. Uncertainties on the combined  $M_W$  result.

Source	Uncertainty (MeV)
Lepton energy scale	3.0
Lepton energy resolution	1.2
Recoil energy scale	1.2
Recoil energy resolution	1.8
Lepton efficiency	0.4
Lepton removal	1.2
Backgrounds	3.3
$p_T^\ell$ model	1.8
$p_T^W/p_T^\ell$ model	1.3
Parton distributions	3.9
QED radiation	2.7
$W$ boson statistics	6.4
Total	9.4

## Collisions of deuteron (isoscalar) beams at LHC

### ➤ Main uncertainties:

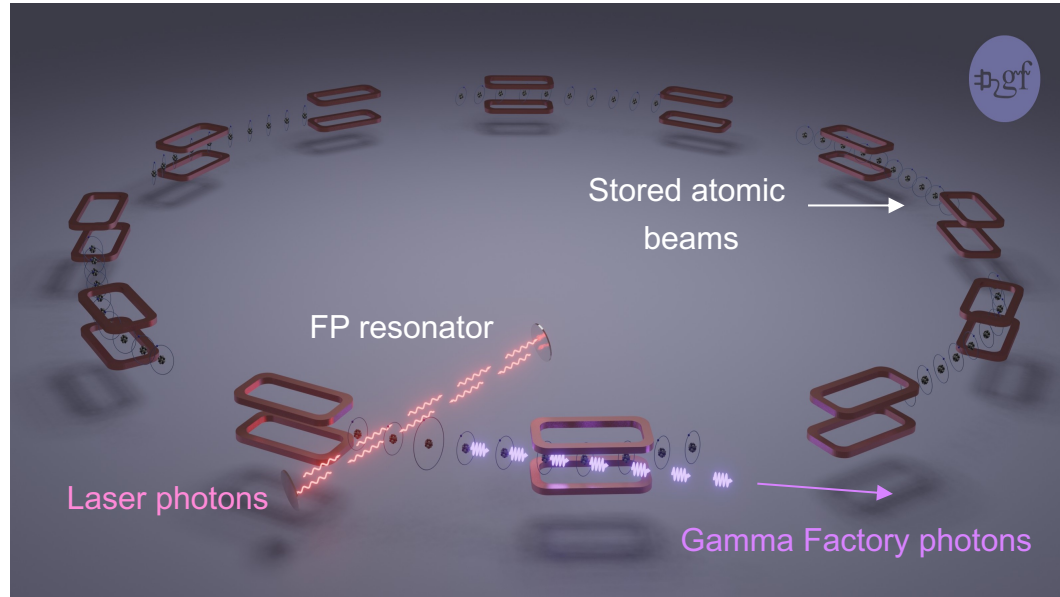
Systematic $\xi$	Expected precision [%]	$\Delta M$ [MeV]
“ $u^{(v)} - d^{(v)}$ ”	0.2	< 5
“ $s - c$ ” $\mathcal{L}_{int} = 2.5 fb^{-1}$	2	25
“ $s - c$ ” $\mathcal{L}_{int} = 25 fb^{-1}$	0.7	8
“ $b$ ” $\mathcal{L}_{int} = 50 fb^{-1}$	40	< 10

( $L_{int}$  –  $dd$  luminosity,  $L_{NN} = 4L_{int}$ )

A. Siódmok, PhD thesis, 2010.

# Gamma Factory photon source: intensity leap

**Requirements:** Accelerated bunches of  $\sim 10^8$ – $10^9$  **partially stripped atoms**, delivered with  $\sim 20$  MHz frequency,  $\sim 5$  mJ laser photon pulses stacked in 20 MHz Fabry-Perot resonator

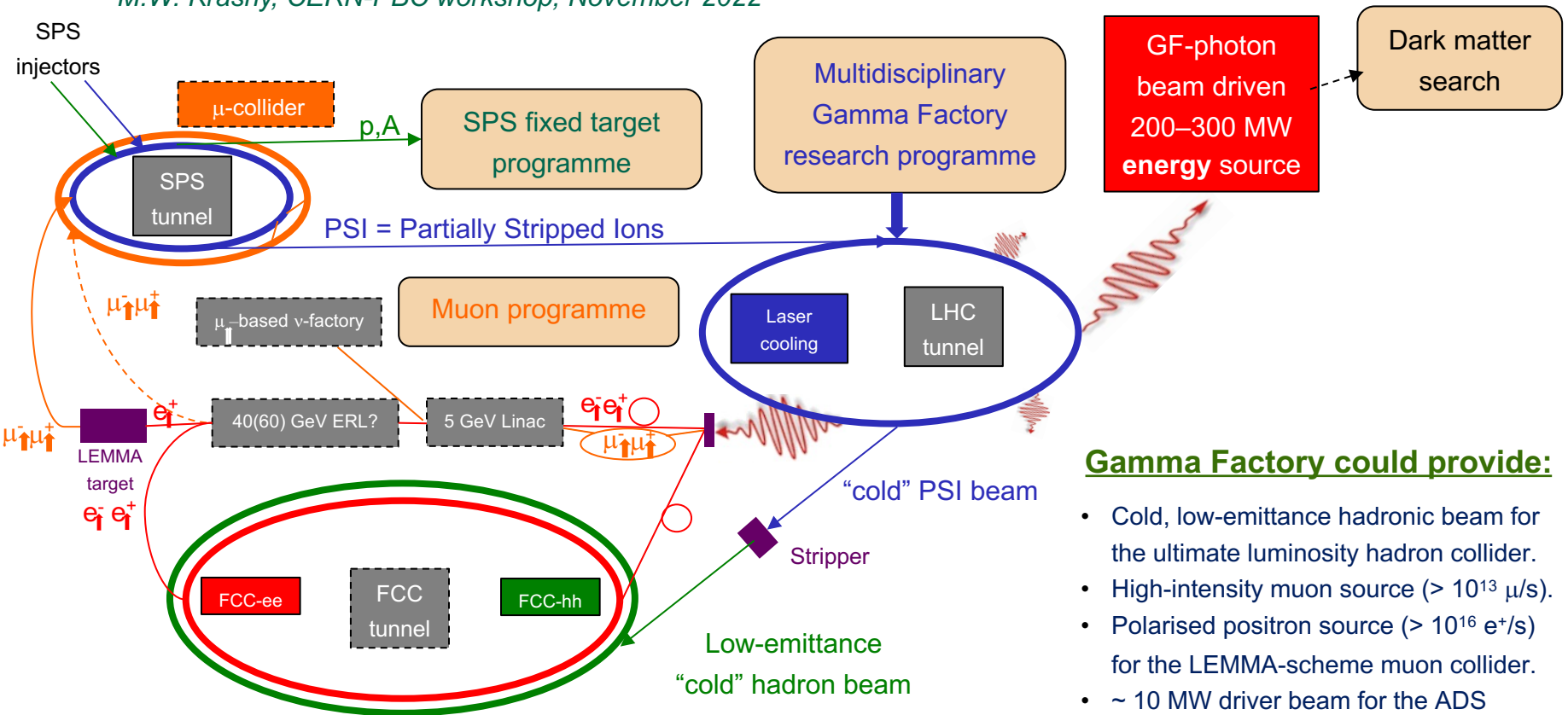


## Novel technology:

Resonant scattering of laser photons on ultra-relativistic atomic beam

# Potential **Gamma Factory** role in the **incremental**, **sustainable** and **multi-disciplinary** development of the research infrastructure at **CERN**

M.W. Krasny, CERN-PBC workshop, November 2022



## Gamma Factory could provide:

- Cold, low-emittance hadronic beam for the ultimate luminosity hadron collider.
- High-intensity muon source ( $> 10^{13} \mu/s$ ).
- Polarised positron source ( $> 10^{16} e^+/s$ ) for the LEMMA-scheme muon collider.
- $\sim 10$  MW driver beam for the ADS energy source (under study).



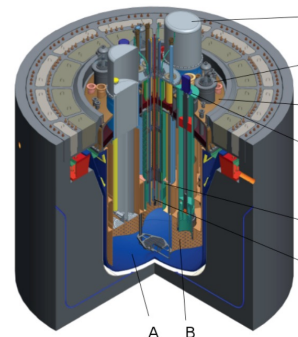
# High-intensity (MW) photon beams

Best use of the CERN expertise to produce rather than buy the plug-power:

## GF photon-beam-driven energy source (ADS)

Satisfying three conditions:

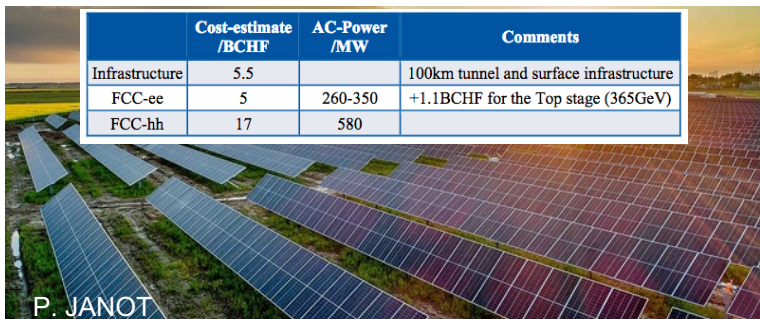
- requisite power for the present and future CERN scientific programme
- operation safety (**a subcritical reactor**)
- efficient transmutation of the nuclear waste (**very important societal impact if demonstrated at CERN – given its reputation**)



**APS April Meeting 2023**  
 Minneapolis, Minnesota (Apr 15-18)

M06 **Invited** [Accelerate Solving Energy Crisis: From Fission to Fusion](#)  
 Room: MG Salon F - 3rd Floor **Sponsor:** DPB FIP **Chair:** Christine Darve, European Spallation Source  
**Invited Speakers:** Hamid Ait Abderrahmane, Mieczyslaw Witold Krasny, Ahmed Diallo, Alireza Haghighat

	Cost-estimate /BCHF	AC-Power /MW	Comments
Infrastructure	5.5		100km tunnel and surface infrastructure
FCC-ee	5	260-350	+1.1BCHF for the Top stage (365GeV)
FCC-hh	17	580	



P. JANOT

→ Would require a 500m-wide band of solar panel along the FCC ring

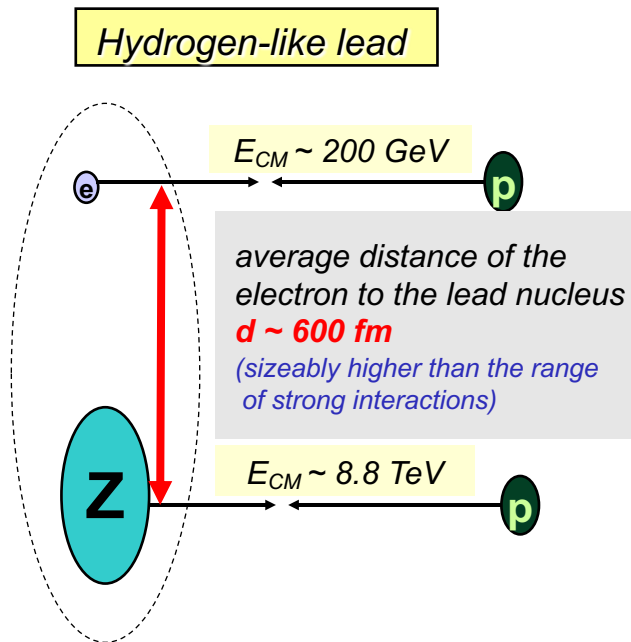


P. JANOT

→ Would require 500 such turbines (one every 200m) along the FCC ring

# Electron beam for ep collisions at LHC

(in the ATLAS, CMS, ALICE and LHCb interaction points)



Atomic beams can be considered as **independent electron and nuclear beams** as long as the incoming proton scatters with the momentum transfer  $q \gg 300 \text{ KeV}$ !

Opens the possibility of collecting by each of the LHC detectors over one day of the  $\text{Pb}^{81+}-\text{p}$  operation the effective ep-collision luminosity comparable to the HERA integrated luminosity in the first year of its operation (1992) – **in-situ diagnostic of the emittance of partonic beams at the LHC!**



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Nuclear Instruments and Methods in Physics Research A 540 (2005) 222–234



[www.elsevier.com/locate/nima](http://www.elsevier.com/locate/nima)

Electron beam for LHC

Mieczyslaw Witold Krasny

LPNHE, Université Pierre et Marie Curie, 4 Pl. Jussieu, Tour 33, RDC, 75025 Paris, France

Received 14 September 2004; received in revised form 19 November 2004; accepted 23 November 2004

Available online 22 December 2004

Initial studies:

Recent development:

PHYSICAL REVIEW ACCELERATORS AND BEAMS **23**, 101002 (2020)

Editors' Suggestion

Collimation of partially stripped ions in the CERN Large Hadron Collider

A. Gorzawski<sup>1,2,\*</sup>, A. Abramov<sup>1,3,†</sup>, R. Bruce<sup>1</sup>, N. Fuster-Martinez<sup>1</sup>, M. Krasny<sup>1,4</sup>, J. Molson<sup>1</sup>, S. Redaelli<sup>1</sup>, and M. Schaumann<sup>1,‡</sup>

<sup>1</sup>CERN European Organization for Nuclear Research, Esplanade des Particules 1, 1211 Geneva, Switzerland

<sup>2</sup>University of Malta, Msida, MSD 2080 Malta

<sup>3</sup>JAI, Egham, Surrey, United Kingdom

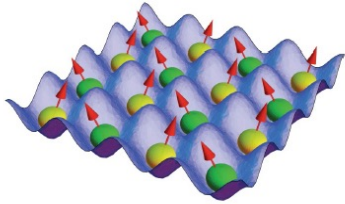
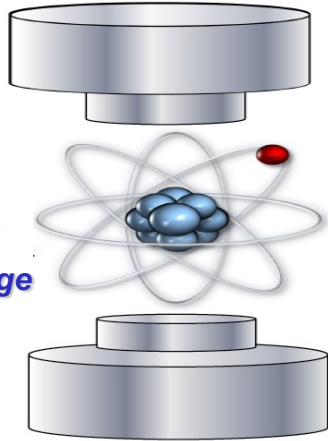
<sup>4</sup>LPNHE, Sorbonne University, CNRS/INP2P3, Tour 33, RdC, 4, pl. Jussieu, 75005 Paris, France

✉ (Received 3 August 2020; accepted 5 October 2020; published 23 October 2020)

# Atomic traps of highly-charged “small-size” atoms

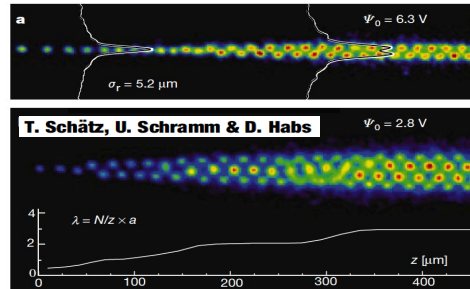
## Atomic rest-frame

Trapped stationary atoms  
exposed to pulsed magnetic  
and electric fields of the storage  
ring



*Crystalline beams?*

## letters to nature



## Opening new research opportunities in atomic physics:

- Highly-charged atoms – very strong ( $\sim 10^{16}$  V/cm) electric field (QED-vacuum effects)
- Small size atoms (electroweak effects)
- Hydrogen-like and helium-like atomic structure (calculation precision and simplicity)
- Atomic degrees of freedom of trapped highly-charged ions can be resonantly excited by lasers



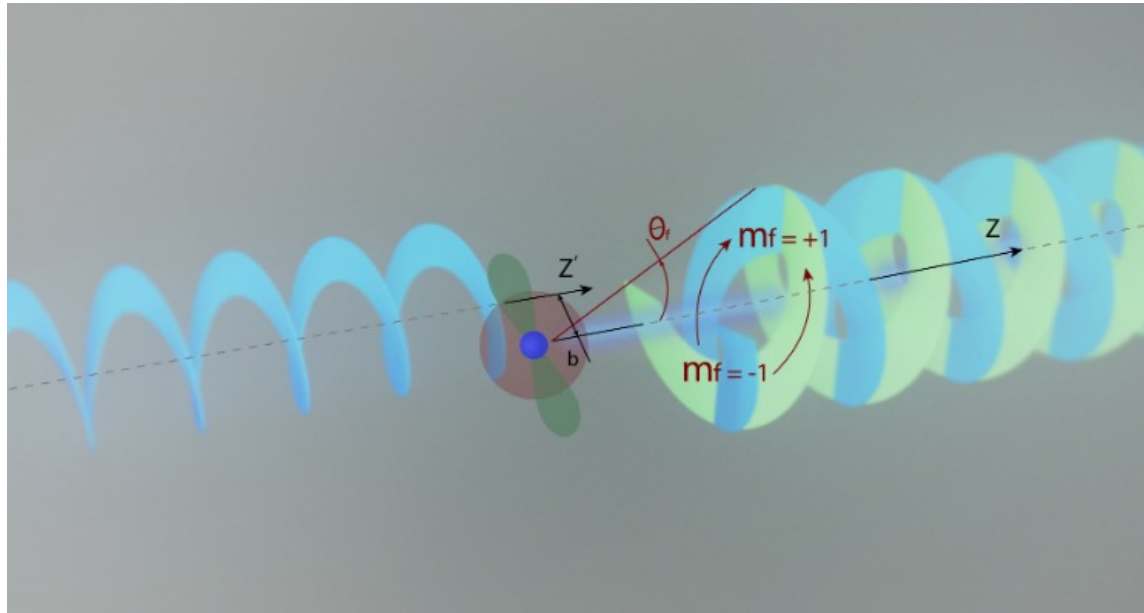
Feature Article | Open Access |

## Atomic Physics Studies at the Gamma Factory at CERN

Dmitry Budker , José R. Crespo López-Urrutia, Andrei Derevianko, Victor V. Flambaum, Mieczysław Witold Krasny, Alexey Petrenko, Szymon Pustelny, Andrey Surzhykov , [Vladimir A. Yerokhin](#), Max Zolotarev ... See fewer authors

First published: 09 July 2020 | <https://doi.org/10.1002/andp.202000204>

# Gamma Factory twisted photons



## Resonant scattering of plane-wave and twisted photons at the Gamma Factory

Valeriy G. Serbo

*Novosibirsk State University, RUS-630090, Novosibirsk, Russia and  
Sobolev Institute of Mathematics, RUS-630090, Novosibirsk, Russia*

Andrey Surzhykov

*Physikalisch-Technische Bundesanstalt, D-38116 Braunschweig, Germany  
Institut für Mathematische Physik, Technische Universität Braunschweig, D-38106 Braunschweig, Germany and  
Laboratory for Emerging Nanometrology Braunschweig, D-38106 Braunschweig, Germany*

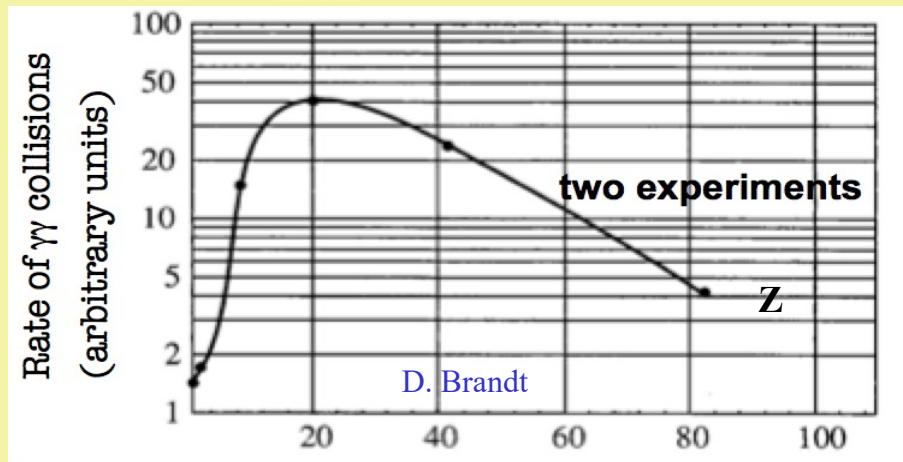
Andrey Volotka

*School of Physics and Engineering, ITMO University, RUS-199034, Saint-Petersburg, Russia*

# Gamma Factory path is restricted to a narrow range of nuclei

1. Ion source.
2. Space-charge tune shift  
 $\sim Z^2 N_{\text{ion}} / A \gamma_L^3$ .
3. Electron stripping pattern.
4. Laser constraints.
5. Lifetime of atomic beams in SPS.
6. Intra beam scattering (IBS):  
 $\alpha_{\text{IBS}}^A = Z^3 / A^2 \alpha_{\text{IBS}}^p$
7. Beam collimation:  $\xi_{x,y} = \frac{r_0 N \beta_{x,y}^*}{2\pi \gamma \sigma_{x,y}^* (\sigma_x^* + \sigma_y^*)}$
8. Beam-beam parameter.
9. Parasitic beam burning:  
 $e^+e^-$  production  $\sim Z^7$  and  
nucleus dissociation  $\sim [(A - Z)Z^3/A^{2/3}]$ .

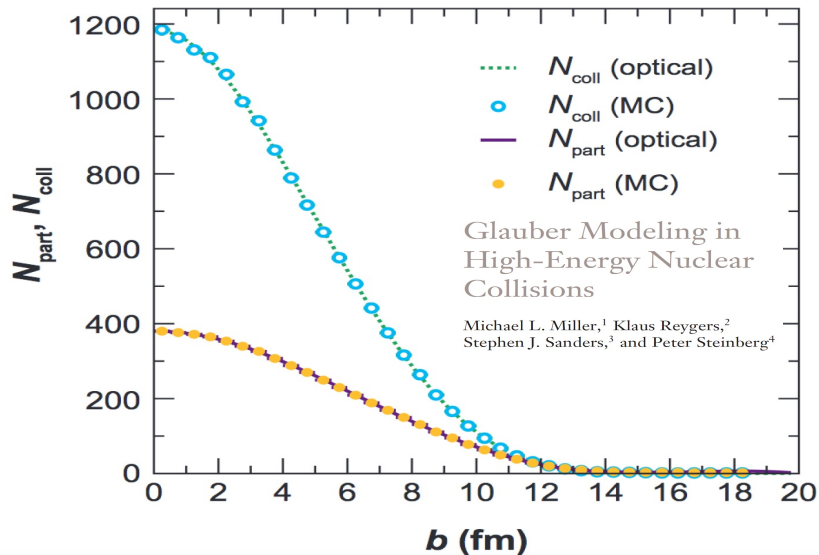
**Ca<sup>20+</sup> beam** turns out to be the best candidate for **HL(AA)-LHC “parton carrier”** – it has a stable **isoscalar isotope** and **maximises photon-photon luminosity** in **AA** collisions.



# Lower pile-up background at equivalent (high) partonic luminosity

$$\nu_{AA} = \nu_{pp} \times \frac{\sigma_{AA}}{\sigma_{pp}} \times A^{-2}$$

Example: Au-Au collisions



**Reduction** of the average number of **beam-particle collisions** per bunch-crossing at the same partonic luminosity by a factor of 40, **136**, 650 and 1260 for O-O, **Ca-Ca**, Xe-Xe and Pb-Pb collisions – opens a possibility of studying **exclusive processes**.

**Reduction** of the multiplicity of soft, **pile-up particles** – opens a possibility to run the LHC detectors (trackers) at **higher partonic luminosity** w.r.t. HL(pp)-LHC.

$$N_{pile-up}(pp) \sim N_{coll}, N_{pile-up}(AA) \sim N_{part}/2$$

(stripping-off of soft partons and nucleus surface)