THE HENRYK NIEWODNICZAŃSKI
INSTITUTE OF NUCLEAR PHYSICS
POLISH ACADEMY OF SCIENCES

NLO matching with KrkNLO

theory and progress

James Whitehead (IFJ PAN, Kraków)
with Wiesław Płaczek, Pratixan Sarmah, Andrzej Siódmok (UJ, Kraków)

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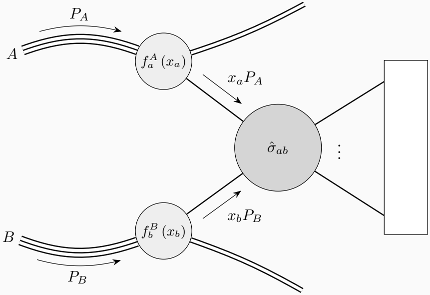
Epiphany XXX, Kraków

Matching at NLO

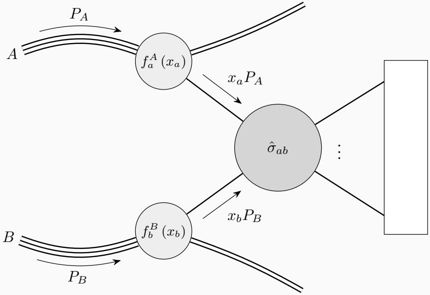
NLO matching



Interested in some[1] function $\mathcal{O}$ of phase-space:

$$\mathrm{d}\sigma_{AB}\left[\mathcal{O}\right](\mu_\mathrm{F}, \mu_\mathrm{R}) = \sum_{a,b} f_a^A(\xi_1; \mu_\mathrm{F}) \otimes_{\xi_1} \mathrm{d}\hat{\sigma}_{ab}\left[\mathcal{O}\right](\xi_1, \xi_2; \mu_\mathrm{F}, \mu_\mathrm{R}) \otimes_{\xi_2} f_b^B(\xi_2; \mu_\mathrm{F}).$$

NLO matching



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[1] infrared and collinear safe

Building a cross-section

For fixed-order calculations: expand perturbatively (and subtract)

$$\mathrm{d}\sigma_{ab}^{\mathrm{NLO}}\left[\mathcal{O}\right](\xi_1,\xi_2) = \left(\frac{\alpha_s}{2\pi}\right)^k \left\{\mathrm{d}\Phi_m(\xi_1 P_1, \xi_2 P_2)\left[\mathrm{B}(\Phi_m)\right]\mathcal{O}(\Phi_m)\right\}$$
$$+\left(\frac{\alpha_s}{2\pi}\right)^{k+1}\left\{\mathrm{d}\Phi_m(\xi_1 P_1, \xi_2 P_2)\left[\mathrm{V}(\Phi_m)\right]\mathcal{O}(\Phi_m)\right.$$
$$\left. + \mathrm{d}\Phi_{m+1}(\xi_1 P_1, \xi_2 P_2)\left[\mathrm{R}(\Phi_{m+1})\right]\mathcal{O}(\Phi_{m+1})\right\}$$

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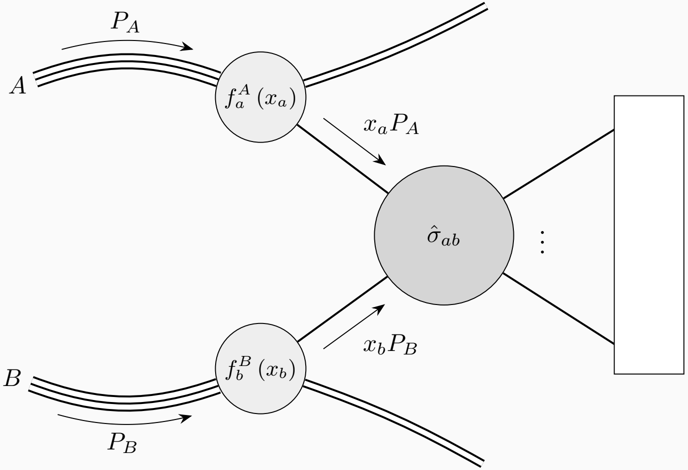
& + \left(\frac{\alpha_s}{2\pi}\right)^{k+1} \Bigg\{ \mathrm{d}\Phi_m \left[\mathrm{V}(\Phi_m) + \mathrm{I}(\Phi_m)\right] \mathcal{O}(\Phi_m) \\
& \qquad + \mathrm{d}x \, \mathrm{d}\Phi_m(x\xi_1 P_1, \xi_2 P_2) \left[(\mathrm{P}+\mathrm{K})_{ac}(x) \cdot \mathrm{B}_{cb}(\Phi_m)\right] \mathcal{O}(\Phi_m) \\
& \qquad + \mathrm{d}x \, \mathrm{d}\Phi_m(\xi_1 P_1, x\xi_2 P_2) \left[(\mathrm{P}+\mathrm{K})_{bc}(x) \cdot \mathrm{B}_{ac}(\Phi_m)\right] \mathcal{O}(\Phi_m) \\
& \qquad + \mathrm{d}\Phi_{m+1} \left[\mathrm{R}(\Phi_{m+1}) - \mathrm{S}(\Phi_{+1}) \cdot \mathrm{B}(\tilde{\Phi}_m)\right] \mathcal{O}(\Phi_{m+1}) \Bigg\}
\end{aligned}
$$

What is a parton shower?



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...look familiar?

,

New legs from old[2]

What is a parton shower?

At its heart:

$$\mathrm{PS}\left[\mathcal{O}\right](\Phi_m) = \Delta_{\mu_s}^{Q(\Phi_m)}\, \mathcal{O}(\Phi_m)$$

[2] Based on ongoing work with Andrzej Siódmok and Simon Plätzer.

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\begin{aligned}
\mathrm{PS}\left[\mathcal{O}\right](\Phi_m) = \Delta_{\mu_s}^{Q(\Phi_m)}\, \mathcal{O}(\Phi_m) \\
+ \sum_\alpha \int \mathrm{d}\Phi_{+1}^{(\alpha)}\, \Theta[\mu_s < \mu(\Phi_{+1}) < Q(\Phi_m)]\, \Delta_{\mu(\Phi_{+1})}^{Q(\Phi_m)} \\
\times P_m^{(\alpha)}(\Phi_{+1})\, \Theta_{\mathrm{PS}}^{(\alpha)}\left[\Phi_{m+1}^{(\alpha)}\right]\, \mathrm{PS}\left[\mathcal{O}\right](\Phi_{m+1}^{(\alpha)})
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\end{aligned}
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where the Sudakov form factor is

$$
\Delta_{\mu_s}^{Q(\Phi_m)} = \exp\left[-\sum_\alpha \int \mathrm{d}\Phi_{+1}^{(\alpha)}\, \Theta[\mu_s < \mu(\Phi_{+1}) < Q(\Phi_m)]\, P_m^{(\alpha)}(\Phi_{+1})\, \Theta_{\mathrm{PS}}^{(\alpha)}\right]
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To be concrete: choose

1. ordering variable $\mu(\Phi_{+1})$
2. radiative splitting phase-space $\Phi_{m+1}^{(\alpha)}$
3. radiative splitting kernels $P_m^{(\alpha)}(\Phi_{m+1}^{(\alpha)})$
4. renormalisation (II: factorisation) scale choice $\mu_{\text{R,F}}(\Phi_{m+1}^{(\alpha)})$
5. shower starting-scale $Q(\Phi_m)$, cut-off scale $\mu_s(\Phi_m)$

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What is a parton shower?

$$\mathrm{PS}\left[\mathcal{O}\right](\Phi_m) = \Delta_{\mu_s}^{Q(\Phi_m)}\,\mathcal{O}(\Phi_m) + \sum_\alpha \int \mathrm{d}\Phi_{+1}^{(\alpha)}\,\Theta[\mu_s < \mu(\Phi_{+1}) < Q(\Phi_m)]\,\Delta_{\mu(\Phi_{+1})}^{Q(\Phi_m)} \times P_m^{(\alpha)}(\Phi_{+1})\,\Theta_{\mathrm{PS}}^{(\alpha)}\left[\Phi_{m+1}^{(\alpha)}\right]\,\mathrm{PS}\left[\mathcal{O}\right](\Phi_{m+1}^{(\alpha)})$$

To be concrete: choose

1. ordering variable $\mu(\Phi_{+1})$ — $p_{\mathrm{T}}^{(\alpha)}$
2. radiative splitting kinematics $\Phi_{m+1}^{(\alpha)}$ — Catani–Seymour II/(IF/FI)/FF
3. radiative splitting kernels $P_m^{(\alpha)}(\Phi_{m+1}^{(\alpha)})$ — Catani–Seymour $\mathrm{D}^{(\alpha)}(\Phi_{+1}^{(\alpha)})$
4. renormalisation (II: factorisation) scale choice $\mu_{\mathrm{R,F}}(\Phi_{m+1}^{(\alpha)})$ — $p_{\mathrm{T}}^{(\alpha)}$
5. shower starting-scale $Q(\Phi_m)$, cut-off scale $\mu_s(\Phi_m)$ — ∞ ('power'), 1 GeV

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Colour singlet: one-emission example

$$\mathrm{d}\hat{\sigma}_{q\bar{q}}^{\mathrm{LO+PS_1(0)}'}[\mathcal{O}] = \mathrm{d}\Phi_m \, \frac{1}{2\hat{s}_{12}} \left[\mathrm{B}_{q\bar{q}}(\Phi_m)\right] \Theta_{\mathrm{cut}}\left[\Phi_m\right] \, \Delta^{(0)}\Big|_{p_\mathrm{T}^{\mathrm{cut}}}^{Q(\Phi_m)} \mathcal{O}(\Phi_m)$$

$$\mathrm{d}\hat{\sigma}_{q\bar{q}}^{\mathrm{LO+PS_1(1)}'}[\mathcal{O}] = \mathrm{d}\Phi_m \, \frac{1}{2\hat{s}_{12}} \left[\mathrm{B}_{q\bar{q}}(\Phi_m)\right] \Theta_{\mathrm{cut}}\left[\Phi_m\right] \, \Delta^{(1)}\Big|_{p_\mathrm{T}^{\mathrm{cut}}}^{Q(\Phi_m)} \mathcal{O}(\Phi_m)$$

$$+ \, \mathrm{d}\Phi_{m+1} \frac{1}{2x\hat{s}_{12}} \Bigg[\mathrm{B}_{q\bar{q}}(\tilde{\Phi}_m^{\mathrm{II}_1}) \, \Theta_{\mathrm{cut}}\left[\tilde{\Phi}_m^{\mathrm{II}_1}\right] \, \Theta_{p_\mathrm{T}^{\mathrm{cut}}}^{Q(\tilde{\Phi}_m^{\mathrm{II}_1})}(\tilde{\Phi}_m^{\mathrm{II}_1}) \, \mathrm{D}^{qg}(x) \, \Delta^{(0)}\Big|_{p_{\mathrm{T},1}}^{Q(\tilde{\Phi}_m^{\mathrm{II}_1})}$$

$$+ \, \mathrm{B}_{q\bar{q}}(\tilde{\Phi}_m^{\mathrm{II}_2}) \, \Theta_{\mathrm{cut}}\left[\tilde{\Phi}_m^{\mathrm{II}_2}\right] \, \Theta_{p_\mathrm{T}^{\mathrm{cut}}}^{Q(\tilde{\Phi}_m^{\mathrm{II}_2})}(\tilde{\Phi}_m^{\mathrm{II}_2}) \, \mathrm{D}^{qg}(x) \, \Delta^{(0)}\Big|_{p_{\mathrm{T},1}}^{Q(\tilde{\Phi}_m^{\mathrm{II}_2})}\Bigg] \Delta^{(0)}\Big|_{p_\mathrm{T}^{\mathrm{cut}}}^{p_{\mathrm{T},1}} \mathcal{O}(\Phi_{m+1})$$

$$\mathrm{d}\hat{\sigma}_{qg}^{\mathrm{LO+PS_1(0)}'}[\mathcal{O}] = 0$$

$$\mathrm{d}\hat{\sigma}_{qg}^{\mathrm{LO+PS_1(1)}'}[\mathcal{O}] = \mathrm{d}\Phi_{m+1} \, \frac{1}{2x\hat{s}_{12}} \, \mathrm{B}_{q\bar{q}}(\tilde{\Phi}_m^{\mathrm{II}_2}) \, \Theta_{\mathrm{cut}}\left[\tilde{\Phi}_m^{\mathrm{II}_2}\right]$$

$$\times \, \Theta_{p_\mathrm{T}^{\mathrm{cut}}}^{Q(\tilde{\Phi}_m^{\mathrm{II}_2})} \, \mathrm{D}^{gq}(x) \, \Delta^{(0)}\Big|_{p_{\mathrm{T},1}}^{Q(\tilde{\Phi}_m^{\mathrm{II}_2})} \, \Delta^{(0)}\Big|_{p_\mathrm{T}^{\mathrm{cut}}}^{p_{\mathrm{T},1}} \mathcal{O}(\Phi_{m+1})$$

NLO matching criterion

The matched NLO cross-section shouldn't spoil the fixed-order result:

$$\hat{\sigma}^{\mathrm{NLO+PS}}[\mathcal{O}] = \hat{\sigma}^{\mathrm{NLO}}[\mathcal{O}]$$

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The matched NLO cross-section shouldn't spoil the fixed-order result (up to power corrections and higher-order terms) :

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	$\mathcal{O}(\Phi_m)$	$\mathcal{O}(\Phi_{m+1})$
LO (α_s^k)	$\mathrm{B}(\Phi_m)$	
NLO (α_s^{k+1})	$\mathrm{V}(\Phi_m)$	$\mathrm{R}(\Phi_{m+1})$
shower (α_s^{k+1})	$-\mathrm{B} \cdot \int \mathrm{d}\Phi_{+1} P_m^{(\alpha)}(\Phi_{+1})$	$+\mathrm{B} \cdot \mathrm{d}\Phi_{+1} P_m^{(\alpha)}(\Phi_{+1})$
factorisation scheme (α_s^{k+1})	$\Delta f_a \otimes_{\xi_1} \mathrm{B} + \mathrm{B} \otimes_{\xi_2} \Delta f_b$	

Satisfying the matching condition

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- need full phase-space coverage from shower!
- but weight is always positive.
- no subtraction...

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Satisfying the matching condition

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$\mathcal{O}(\Phi_m)$: restore the cancellation required by the matching condition by modifying the PDF factorisation scheme

- collinear convolution terms can only go into the PDF
- where to put end-point contributions $\propto \delta(1-x)$?

Familiar ingredients

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For dipoles, we already know the answer from dipole subtraction:[5]

$$-\sum_{\alpha} \int \mathrm{d}\Phi_{+1}\, \Theta[\mu_s < \mu(\Phi_{+1}) < Q(\Phi_m)]\, P_m^{(\alpha)}(\Phi_{+1})\, \Theta_{\mathrm{PS}}^{(\alpha)} = \sum_{(\alpha)} \mathrm{I}^{(\alpha)} + \mathrm{d}x \left(\mathrm{P}^{(\alpha)} + \mathrm{K}^{(\alpha)}\right)$$

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This provides the recipe for the PDF transformation (more later).

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Final recipe[6]

$$\mathrm{d}\Phi_m\, \Theta_{\text{cut}}\left[\Phi_m\right] \quad \Bigg[\Bigg\{\mathrm{B}(\Phi_m) + \mathrm{V}(\Phi_m) + I(\Phi_m) + \Delta_0^{\text{FS}}\Bigg\}\Delta_{\mu_s}^{Q(\Phi_m)}\,\mathcal{O}(\Phi_m)$$

$$+ \sum_\alpha \mathrm{d}\Phi_{+1}^{(\alpha)} \left\{\frac{\mathrm{R}^{(\alpha)}(\Phi_{m+1}^{(\alpha)})}{\mathrm{PS}^{(\alpha)}(\Phi_{m+1}^{(\alpha)})}\,\Theta_{\text{PS}}^{(\alpha)}\left[\Phi_{m+1}^{(\alpha)}\right]\,\mathrm{PS}^{(\alpha)}\left[\Phi_{m+1}^{(\alpha)}\right]\,\Theta_{\mu_s}^{(\alpha)}\,\mathcal{O}(\Phi_{m+1}^{(\alpha)})\right\}\Bigg]$$

1. generate a Born phase-space point, ME and shower:
   - if an emission is generated, reweight to R
   - if not, reweight to B + V
2. matching complete; allow the shower to proceed!

6 S. Jadach et al. "Matching NLO QCD with parton shower in Monte Carlo scheme — the KrkNLO method". arXiv: 1503.06849 [hep-ph], Stanislaw Jadach et al. "New simpler methods of matching NLO corrections with parton shower Monte Carlo". arXiv: 1607.00919 [hep-ph].

Final recipe[6]

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\begin{aligned}
\mathrm{d}\Phi_m\,\Theta_{\text{cut}}\left[\Phi_m\right] \quad & \Bigg[\Bigg\{\mathrm{B}(\Phi_m)+\mathrm{V}(\Phi_m)+I(\Phi_m)+\Delta_0^{\mathrm{FS}}\Bigg\}\Delta_{\mu_s}^{Q(\Phi_m)}\,\mathcal{O}(\Phi_m) \\
& +\sum_\alpha \mathrm{d}\Phi_{+1}^{(\alpha)}\left\{\frac{\mathrm{R}^{(\alpha)}(\Phi_{m+1}^{(\alpha)})}{\mathrm{PS}^{(\alpha)}(\Phi_{m+1}^{(\alpha)})}\,\Theta_{\mathrm{PS}}^{(\alpha)}\left[\Phi_{m+1}^{(\alpha)}\right]\,\mathrm{PS}^{(\alpha)}\left[\Phi_{m+1}^{(\alpha)}\right]\,\Theta_{\mu_s}^{(\alpha)}\,\mathcal{O}(\Phi_{m+1}^{(\alpha)})\right\}\Bigg]
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$$

1. generate a Born phase-space point, ME and shower:
   - if an emission is generated, reweight to R
   - if not, reweight to B + V
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This is NLO accurate, but differs from other methods at higher orders.

6 S. Jadach et al. "Matching NLO QCD with parton shower in Monte Carlo scheme — the KrkNLO method". arXiv: 1503.06849 [hep-ph], Stanislaw Jadach et al. "New simpler methods of matching NLO corrections with parton shower Monte Carlo". arXiv: 1607.00919 [hep-ph].

Krk PDF scheme

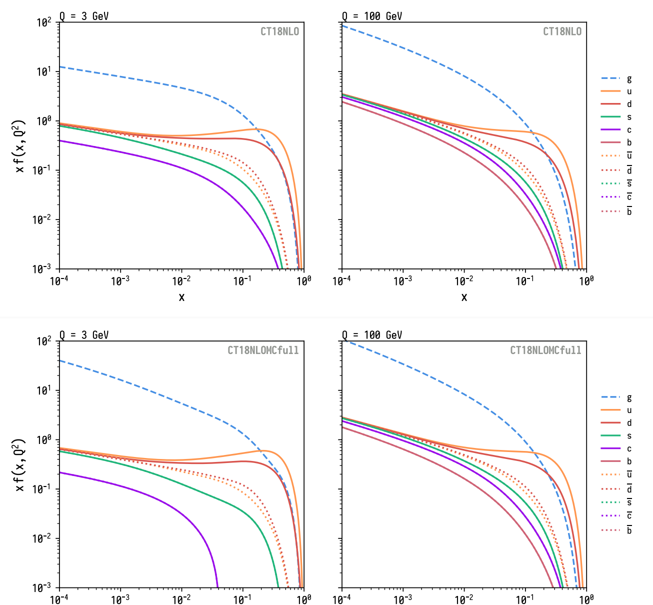
Krk (/MC/CS) factorisation scheme[7]

From the dipole operators, we can write down the convolution terms:

$$
\begin{aligned}
f_q^{\mathrm{Krk}}(x,\mu_{\mathrm{F}}) = f_q^{\overline{\mathrm{MS}}}(x,\mu_{\mathrm{F}}) \\
&- \frac{\alpha_s(\mu_{\mathrm{F}})}{2\pi}\,\frac{3}{2} C_{\mathrm{F}}\; f_q^{\overline{\mathrm{MS}}}(x,\mu_{\mathrm{F}}) \\
&+ \frac{\alpha_s(\mu_{\mathrm{F}})}{2\pi}\; C_{\mathrm{F}}\left[\int_x^1 \frac{\mathrm{d}z}{z} f_q^{\overline{\mathrm{MS}}}\left(\frac{x}{z},\mu_{\mathrm{F}}\right) \quad \left[\frac{1+z^2}{1-z}\log\frac{(1-z)^2}{z} + 1 - z\right]_+\right] \\
&+ \frac{\alpha_s(\mu_{\mathrm{F}})}{2\pi}\; C_{\mathrm{A}}\left[\int_x^1 \frac{\mathrm{d}z}{z} f_g^{\overline{\mathrm{MS}}}\left(\frac{x}{z},\mu_{\mathrm{F}}\right)\left[z^2+(1-z)^2\right]\log\frac{(1-z)^2}{z} + 2z(1-z)\right] \\
f_g^{\mathrm{Krk}}(x,\mu_{\mathrm{F}}) = f_g^{\overline{\mathrm{MS}}}(x,\mu_{\mathrm{F}}) \\
&- \frac{\alpha_s(\mu_{\mathrm{F}})}{2\pi}\; C_{\mathrm{A}}\left[\frac{\pi^2}{3} + \frac{341}{72} - \frac{59}{36}\frac{N_f T_{\mathrm{R}}}{C_{\mathrm{A}}}\right]\; f_g^{\overline{\mathrm{MS}}}(x,\mu_{\mathrm{F}}) \\
&+ \frac{\alpha_s(\mu_{\mathrm{F}})}{2\pi}\; C_{\mathrm{A}}\left[\int_x^1 \frac{\mathrm{d}z}{z} f_g^{\overline{\mathrm{MS}}}\left(\frac{x}{z},\mu_{\mathrm{F}}\right)\left[4\left[\frac{\log(1-z)}{1-z}\right]_+ - 2\frac{\log z}{1-z}\right.\right. \\
&\qquad \left.\left. +2\left(\frac{1}{z} - 2 + z(1-z)\right)\ln\frac{(1-z)^2}{z}\right]\right] \\
&+ \frac{\alpha_s(\mu_{\mathrm{F}})}{2\pi}\; C_{\mathrm{F}}\sum_{q_f,\bar{q}_f}\left[\int_x^1 \frac{\mathrm{d}z}{z} f_q^{\overline{\mathrm{MS}}}\left(\frac{x}{z},\mu_{\mathrm{F}}\right)\left[\frac{1+(1-z)^2}{z}\log\frac{(1-z)^2}{z} + z\right]\right]
\end{aligned}
$$

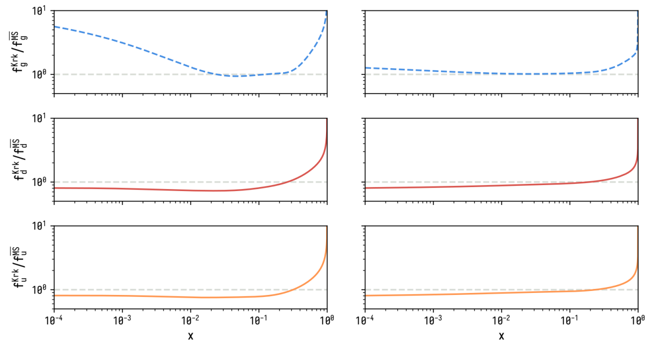
7 S. Jadach et al. "Parton distribution functions in Monte Carlo factorisation scheme". arXiv: 1606.00355 [hep-ph].

Applied to LHAPDF6 grids:



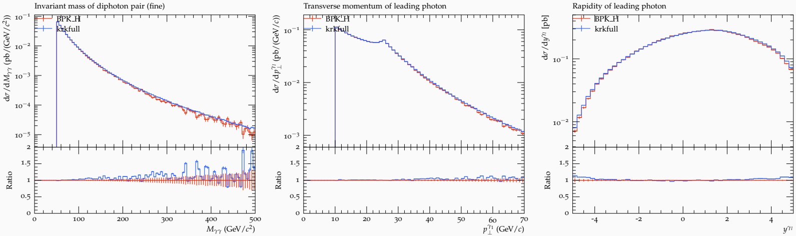
PDFs in MC scheme

Applied to LHAPDF6 grids:



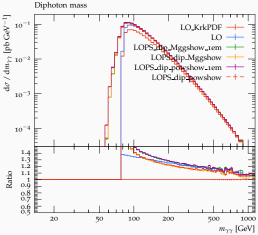
Validation

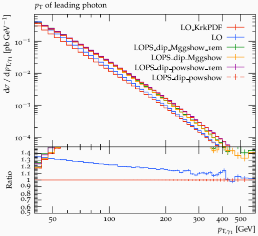
Do we reproduce the Herwig (Matchbox) automated P and K operators?

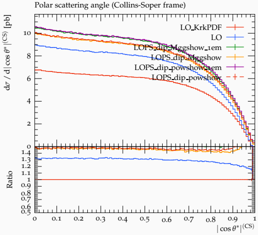


Normalisation

What is the numerical impact of the Krk scheme?







Validation

Seeing Sudakovs

To verify the real weight, we must *unweight* the Sudakov:

- numerical integration of dipole kernels considered in shower algorithm;
- over the same splitting phase-space/kinematic region used in the shower algorithm;
- with the same scales, PDF arguments, α_s etc

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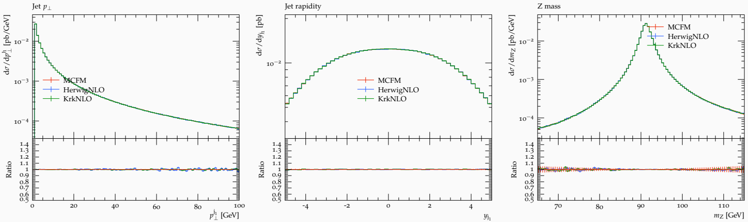
$$\Delta_{\mu_s}^{Q(\phi_m)} = \exp\left[-\sum_{\alpha}\int \mathrm{d}q(\phi_m)\,\Theta[\mu_s < \mu(q) < Q(\phi_m)]\,P_m^{(\alpha)}(q)\,\Theta_{\mathrm{PS}}^{(\alpha)}\right]$$

This is non-trivial!

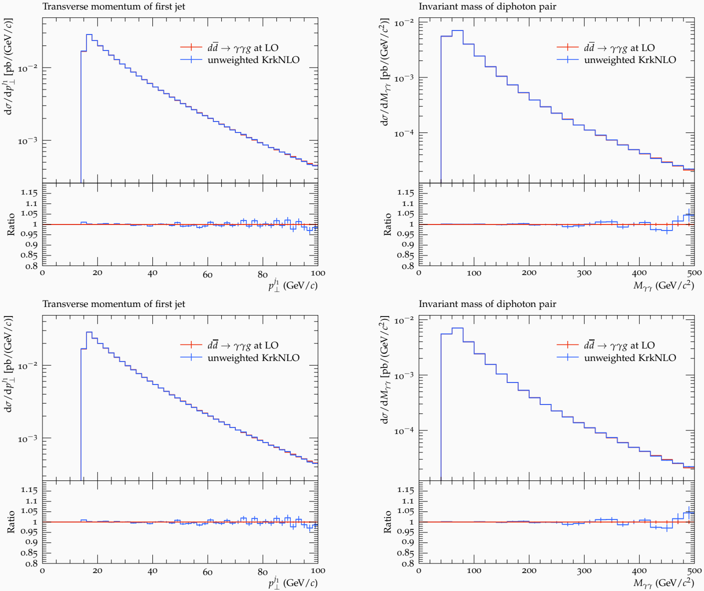
Does it work?

Drell-Yan:



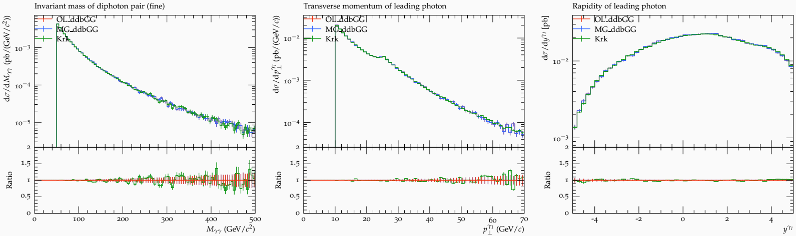
Does it work?

Diphoton:



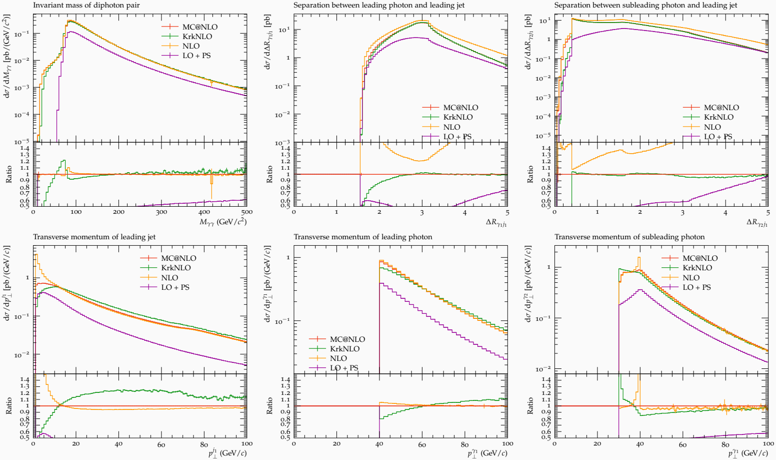
What about the virtuals?

Diphoton:

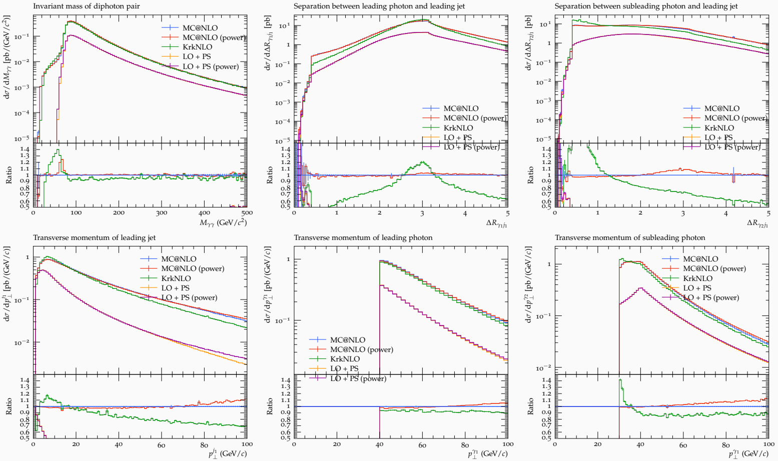


Results

Single emision

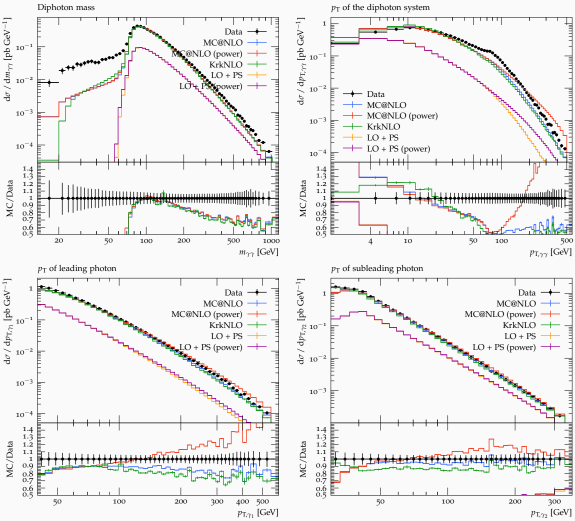


Full shower



Invariant mass of diphoton pair
Separation between leading photon and leading jet
Separation between subleading photon and leading jet
Transverse momentum of leading jet
Transverse momentum of leading photon
Transverse momentum of subleading photon
MC@NLO
MC@NLO (power)
KrkNLO
LO + PS
LO + PS (power)
Ratio
$d\sigma/dM_{\gamma\gamma}$ [pb/(GeV/c^2)]
$M_{\gamma\gamma}$ (GeV/c^2)
$d\sigma/d\Delta R_{\gamma_1 j_1}$ [pb]
$\Delta R_{\gamma_1 j_1}$
$d\sigma/d\Delta R_{\gamma_2 j_1}$ [pb]
$\Delta R_{\gamma_2 j_1}$
$d\sigma/dp_\perp^{j_1}$ [pb/(GeV/c)]
$p_\perp^{j_1}$ (GeV/c)
$d\sigma/dp_\perp^{\gamma_1}$ [pb/(GeV/c)]
$p_\perp^{\gamma_1}$ (GeV/c)
$d\sigma/dp_\perp^{\gamma_2}$ [pb/(GeV/c)]

Full shower: with data



Outlook

- more new processes in the pipeline
- PDF factorisation scheme[8]
- logs?
- automation!
- …+jet?

8 S. Jadach et al. “Parton distribution functions in Monte Carlo factorisation scheme”. arXiv: 1606.00355 [hep-ph], S. Jadach. “On the universality of the KRK factorization scheme”. arXiv: 2004.04239 [hep-ph].

Thank you!

Momentum mappings

require: four-momentum conservation & all particles remain on-shell

final-final

$$\tilde{p}_i = p_i + p_j - \frac{s_{ij}}{s_{ik} + s_{jk}} p_k \qquad \tilde{p}_k = \left(1 + \frac{s_{ij}}{s_{ik} + s_{jk}}\right) p_k$$

initial-final & final-initial

$$\tilde{p}_a = \left(1 - \frac{s_{jk}}{s_{aj} + s_{ak}}\right) p_a \qquad \tilde{p}_k = p_j + p_k - \frac{s_{jk}}{s_{aj} + s_{ak}} p_a$$

initial-initial

$$\tilde{p}_a = \left(1 - \frac{s_{aj} + s_{bj}}{s_{ab}}\right) p_a \qquad \tilde{p}_b = p_b$$

(in this case we further need to boost all FS particles)

Details of KrkNLO

Krk PDF scheme within KrkNLO

Krk PDFs compensate for the integrated shower radiation at $\mathcal{O}(\alpha_s)$ within the Sudakov factor. Schematically:

$$
\begin{aligned}
&\mathrm{d}\xi_1\, \mathrm{d}\xi_2 && \left\{\mathbf{f}^{\overline{\mathrm{MS}}} \otimes (\mathbb{I} + \mathbf{P} + \mathbf{K})\right\}_a \left\{\mathbf{f}^{\overline{\mathrm{MS}}} \otimes (\mathbb{I} + \mathbf{P} + \mathbf{K})\right\}_b \\
&\Bigg\{ \mathrm{d}\phi_m\, \Theta_{\mathrm{cut}}\left[\phi_m\right] && \Bigg[u(\phi_m)\, \mathrm{B}(\phi_m) \left\{ 1 + \frac{\mathrm{V}}{\mathrm{B}} + \sum_\alpha \mathrm{I}^{(\alpha)} - \mathrm{I}^{\mathrm{FS}}_{ab} \right\} \Delta_{\mu_s}^{Q_{\max}(\phi_m)} \\
& && + \sum_\alpha \mathrm{d}q^{(\alpha)}\; u(\Phi^{(\alpha)}_{m+1}) \left\{ \frac{\mathrm{R}}{\mathrm{PS}}\, \Theta^{(\alpha)}_{\mathrm{PS}} \left[\Phi^{(\alpha)}_{m+1}\right] \mathrm{PS}^{(\alpha)} \left[\Phi^{(\alpha)}_{m+1}\right] \Theta^{(\alpha)}_{\mu_s} \right\} \Bigg] \Bigg\}
\end{aligned}
$$

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&\Bigg\{\mathrm{d}\phi_m\, \Theta_{\mathrm{cut}}\,[\phi_m] \quad \Bigg[u(\phi_m)\,\mathrm{B}(\phi_m)\left\{1 + \frac{\mathrm{V}}{\mathrm{B}} + \sum_\alpha \mathrm{I}^{(\alpha)} - \mathrm{I}^{\mathrm{FS}}_{ab}\right\} \Delta_{\mu_s}^{Q_{\max}(\phi_m)} \\
&\qquad + \sum_\alpha \mathrm{d}q^{(\alpha)}\; u(\Phi^{(\alpha)}_{m+1}) \left\{\frac{\mathrm{R}}{\mathrm{PS}}\, \Theta^{(\alpha)}_{\mathrm{PS}}\left[\Phi^{(\alpha)}_{m+1}\right]\, \mathrm{PS}^{(\alpha)}\left[\Phi^{(\alpha)}_{m+1}\right]\, \Theta^{(\alpha)}_{\mu_s}\right\}\Bigg]\Bigg\}
\end{aligned}
$$

Additional convolutions define a PDF factorisation scheme: the ‘Krk scheme’.

Krk PDF transformation

Full details:[9]

Parton distribution functions in Monte Carlo factorisation scheme

S. Jadach[1], **W. Płaczek**[2], **S. Sapeta**[1,3], **A. Siódmok**[1,3,a], **M. Skrzypek**[1]

[1] Institute of Nuclear Physics, Polish Academy of Sciences, ul. Radzikowskiego 152, 31-342 Kraków, Poland
[2] Marian Smoluchowski Institute of Physics, Jagiellonian University, ul. Łojasiewicza 11, 30-348 Kraków, Poland
[3] Theoretical Physics Department, CERN, Geneva, Switzerland

$$\begin{bmatrix} q(x,Q^2) \\ \bar{q}(x,Q^2) \\ g(x,Q^2) \end{bmatrix}_{\mathrm{MC}} = \begin{bmatrix} q \\ \bar{q} \\ g \end{bmatrix}_{\overline{\mathrm{MS}}} + \int dz dy \begin{bmatrix} K_{qq}^{\mathrm{MC}}(z) & 0 & K_{qg}^{\mathrm{MC}}(z) \\ 0 & K_{\bar{q}\bar{q}}^{\mathrm{MC}}(z) & K_{\bar{q}g}^{\mathrm{MC}}(z) \\ K_{gq}^{\mathrm{MC}}(z) & K_{g\bar{q}}^{\mathrm{MC}}(z) & K_{gg}^{\mathrm{MC}}(z) \end{bmatrix} \begin{bmatrix} q(y,Q^2) \\ \bar{q}(y,Q^2) \\ g(y,Q^2) \end{bmatrix}_{\overline{\mathrm{MS}}} \delta(x-yz), \tag{4.2}$$

$$\begin{aligned}
K_{gq}^{\mathrm{MC}}(z) &= \frac{\alpha_s}{2\pi} C_F \left\{ \frac{1+(1-z)^2}{z} \ln\frac{(1-z)^2}{z} + z \right\}, \\
K_{gg}^{\mathrm{MC}}(z) &= \frac{\alpha_s}{2\pi} C_A \Bigg\{ 4\left[\frac{\ln(1-z)}{1-z}\right]_+ + 2\left[\frac{1}{z} - 2 + z(1-z)\right] \\
&\quad \times \ln\frac{(1-z)^2}{z} - 2\frac{\ln z}{1-z} - \delta(1-z) \\
&\quad \times \left(\frac{\pi^2}{3} + \frac{341}{72} - \frac{59}{36}\frac{T_f}{C_A}\right) \Bigg\}, \\
K_{qq}^{\mathrm{MC}}(z) &= \frac{\alpha_s}{2\pi} C_F \Bigg\{ 4\left[\frac{\ln(1-z)}{1-z}\right]_+ - (1+z)\ln\frac{(1-z)^2}{z} \\
&\quad -2\frac{\ln z}{1-z} + 1 - z - \delta(1-z)\left(\frac{\pi^2}{3} + \frac{17}{4}\right) \Bigg\}, \\
K_{qg}^{\mathrm{MC}}(z) &= \frac{\alpha_s}{2\pi} T_R \left\{ \left[z^2 + (1-z)^2\right] \ln\frac{(1-z)^2}{z} + 2z(1-z) \right\}, \\
K_{\bar{q}\bar{q}}^{\mathrm{MC}}(z) &= K_{qq}^{\mathrm{MC}}(z), \quad K_{\bar{q}g}^{\mathrm{MC}}(z) = K_{qg}^{\mathrm{MC}}(z).
\end{aligned} \tag{4.3}$$

Note additional imposition of sum rules:

$$\begin{aligned}
\int_0^1 \mathrm{d}z\, z \left[K_{qq}^{\mathrm{MC}}(z) + K_{gq}^{\mathrm{MC}}(z) \right] &= 0\,, \\
\int_0^1 \mathrm{d}z\, z \left[K_{gg}^{\mathrm{MC}}(z) + 2n_f K_{qg}^{\mathrm{MC}}(z) \right] &= 0\,.
\end{aligned} \tag{4.5}$$

[9] S. Jadach et al. "Parton distribution functions in Monte Carlo factorisation scheme". arXiv: 1606.00355 [hep-ph].