Electroweak Precision Measurements at LEP and FCCee



F.J. Hasert et al., Neutrino-like interactions

In 1973 – 50 years ago: Gargamelle discovery of Neutral currents, The Standard Model (unified electro-weak theory) was born to the experimental world



The Nobel Prize in Physics 1979



Photo from the Nobel Foundation archive. Sheldon Lee Glashow Prize share: 1/3





Photo from the Nobel Foundation archive. Abdus Salam Prize share: 1/3

Photo from the Nobel Foundation archive. Steven Weinberg Prize share: 1/3

The Nobel Prize in Physics 1979 was awarded jointly to Sheldon Lee Glashow, Abdus Salam and Steven Weinberg "for their contributions to the theory of the unified weak and electromagnetic interaction between elementary particles, including, inter alia, the prediction of the weak neutral current" The measurement of the NC/CC ratio led to the determination of $\sin^2\theta_w$ and immediately, with G_F and α_{QED} predicted the mass of W and Z (within ± 10 GeV)

Veltman in 1975 (w. D. Ross),1977 explained that relationship bw $G_{F_{r}} \sin^2 \theta_w m_z m_w$ would be modified by heavy physics if it violates the SU(2)_L symmetry

$$\rho = \left(\frac{m_W}{m_Z \cos \theta_w}\right)^2 = 1 + \alpha / \pi \left(\frac{m_{top}}{m_W}\right)^2$$

→ EW precision measurements as a way to investigate the existence of heavy physics inaccessible directly at contemporary accelerators

This is **not** why the 26.7 km circumference LEP was built.

Veltman *lui-même* who was part of the committee, insisted that it should be large enough to verify that W pair production was not divergent. (TGC)

The construction of LEP was (pushed since 1976)decided by CERN council in 1981, before the W and Z were observed at the proton-antiproton collider! Construction started in 1983.

A big scare of the time was the unknown **number of neutrinos** too many neutrinos would wash out the Z peak (J. Ellis)











on EW measurements



Table 1.2: The $q\overline{q}$ and $\ell^+\ell^-$ event statistics, in units of 10³, used for Z analyses by the experiments ALEPH (A), DELPHI (D), L3 (L) and OPAL (O).



GENRAP (1975) MUSTRAHL(1982)

.... KKMC KORALZ TAULA BHLUMI KORALWW

•••



Staszek Jadach

FCC

... All collaborations use BHLUMI 4.04 [23], the best available Monte-Carlo generator for small angle Bhabha scattering, to calculate the acceptance of their low angle luminosity counters.... ... All collaborations use TAULA, KORALZ

LEP Collaborations Phys.Rept.427:257-454,2006

- [20] S. Jadach, B.F.L. Ward and Z. Was, Comput. Phys. Commun. 79 (1994) 503, (KORALZ 4.0).
- [21] S. Jadach, B.F.L. Ward and Z. Was, Comput. Phys. Commun. 130 (2000) 260, (KK Monte Carlo).
- [22] F.A. Berends, R. Kleiss and W. Hollik, Nucl. Phys. B304 (1988) 712, (BABAMC).
- [23] S. Jadach, W. Placzek, E. Richter-Was, B.F.L. Ward and Z. Was, Comput. Phys. Commun. 102 (1997) 229, (BHLUMI 4.04).

FCC from S. Jadach list of 500+ papers	The Monte Carlo program KORALZ, version 4.0, for the lepton or quark pair production at #7 LEP / SLC energies S. Jadach (CERN and Cracow, INP), B.F.L. Ward (Tennessee U. and SLAC), Z. Was (CERN and Cracow, INP) (Nov, 1993) Published in: <i>Comput.Phys.Commun.</i> 79 (1994) 503-522					
	∂ DOI	[→ cite	🗟 claim	বি reference search	→ 739 citations	
	The tau o	decay libr	ary TAUOLA: Versio	n 2.4	#8	
	S. Jadach (CERN and Cracow, INP), Z. Was (CERN and Cracow, INP), R. Decker (Karlsruhe U.), Johann H. Kuhn (Karlsruhe U.) (Feb, 1993)					
	Published in: Comput.Phys.Commun. 76 (1993) 361-380					
	∂ DOI	[→ cite	🗟 claim	c reference search	➔ 1,108 citations	
	Toward a	a model ir	dependent analysis	s of electroweak data	#9	
	Guido Altarelli (CERN), Riccardo Barbieri (Pisa U.), S. Jadach (CERN) (Jun, 1991)					
	Published in: Nucl.Phys.B 369 (1992) 3-32, Nucl.Phys.B 376 (1992) 444 (erratum)					
	ℓ DOI	[∃ cite	🗟 claim	C reference search	→ 536 citations	

TAUOLA: A Library of Monte Carlo programs to simulate decays of polarized tau leptons #10

ELECTROWEAK RADIATI	VE CORRE CTIONS		
relation betweens observable.	s and $\{G_{F}, \alpha(H_{z}), M_{z}\}$		Altarelli, Barbieri, Jadach 91
DEFINITIONS			
$\sin^2 \Theta_w^{\text{eff}} \cdot \cos^2 \Theta_w^{\text{eff}} = \frac{\pi d(M_z)}{\sqrt{z} G_F M_z^2 (1 + \Delta q) (1 - \epsilon_3/\omega_5 \Theta_w)}$	$\Delta \rho = \frac{\alpha}{\pi} \frac{m_{E}^{2}}{m_{Z^{2}}} - \frac{\alpha}{4\pi} \frac{\ell_{n} \frac{m_{H}^{2}}{m_{Z}^{2}}}{m_{Z}^{2}}$		this is $\epsilon_1 \propto \alpha T$
$\Gamma_{e} = \frac{G_{F} M_{z}^{3}}{24\sqrt{z}\pi} \left(1 + \Delta \rho\right) \left(1 + \frac{3}{4}\frac{\alpha}{\pi}\right)$	$E_3 = + \frac{\alpha}{12\pi} \ln \frac{m_H^2}{m_Z^2}$		and this is $\propto \alpha S$
2	$\mathcal{E}_{2} \approx \frac{1}{9\pi} \frac{\kappa}{m_{z}^{2}} \ln \frac{m_{H}}{m_{z}^{2}}$		
$M_{W}^{2} = \frac{\pi \alpha(M_{z})}{\sqrt{z} G_{F} \left(1 - \frac{M_{W}^{2}}{M_{z}^{2}}\right) \left(1 - \Delta_{r}^{ew}\right)}$	SVB = - 20 x m2 13 TT m2		this is $\epsilon_{\rm b}$ (too often forgotten!)
$\Gamma_{b} = \Gamma_{d} \left(\lambda + \delta_{vb} \right)$	$\Delta r^{ew} = -\frac{G^2 \Theta_{uv}}{S u^2 \Theta_{uv}} \Delta \rho + 2 \varepsilon_3 - \frac{G \sigma_3}{C \sigma_3}$	-su ² 22	this is (part of) Marciano's Δ r
$\sin^2 \theta^{eff}_{W} \equiv 1/4 (1-g_v/g_a)_{electron}$			

 in principle using different observables, it is possible to disentangle the effects of top, Higgs and even something else. (3 most elevant parameters ε₁, ε₃, ε_b + more in BSM)
 the Z mass and width are measured using all Z decays and thus faster statistically and very easy systematically rather than asymmetries or partial widths that require final state selection
 QED corrections constitute a gauge inv. set that factorizes out and has little sensitivity to heavy physics

Polarization at LEP

As a side effect of synchrotron radiation emission, e^+/e^- beams polarize spontaneously (align their spins) in the transverse (vertical) direction, i.e. along the direction of the bending field.

Polarization is however a slow and delicate process which requires a lot of care in machine setup and special conditions.



Resonant Depolarization

The interest of P_T : magnetic moments precess in B-fields

The **number of precessions/turn** v, called **the spin tune**, is **proportional to the energy** :

To determine the energy





Principle :

 ❑ Sweep the B-field of a fast pulsing magnet ("kicker") in frequency and observe P_T,

 $\nu = \frac{g_e - 2}{2} \frac{E}{mc^2} =$

Measure v

□ If kicker frequency and v match, P_T is rotated away from the vertical axis.

 $\frac{E[\text{MeV}]}{440.6486(1)[\text{MeV}]}$

Resonant depolarization

Resonant Depolarization II

In practice :

- □ The kicker frequency is swept over a selected interval (~ 22 Hz).
- P_{T} can be destroyed or flipped when the kicker is in resonance.



Intrinsic accuracy at LEP :



This technique is over an order of magnitude more accurate than any other method !

But it required a large amount of DEDICATED beam time as polarization was not considered compatible with physics data taking. Done at end of physics fills → bias For instance, solenoids were not spin-compensated Only e- (not e+) was measured (AB: in hindsight this was a big mistake) → in the future need pilot bunches, compensation, and both e+ and e- polarimeters The measurements were very precise but not reproducible! no correlation with temperature or time of day.



and indeed the measurements correlated nicely with the calculated amplitude of the earth tides.



The scatter is reduced to about ±1 MeV.

in 1992 we stopped scanning and spent some time understanding things better...









beginning.

Success in the Press !



In 1993 the Z peak was scanned very thoroughly with a sequence of data points **LEP and the top quark** at spin tunes of 101.5 (peak '-2') , 103.5 (peak) , 105.5 (peak'+2') Nature was kind because these points were both ar away from spin resonances, and very near optimal for the Z width determination with precision of ± 3 MeV. At the same time the muon forward-backward asymmetries (this also depends strongly on energy) were measured as well as tau polarization and all things that measure $\sin^2\theta_w^{eff}$.

At the end of year the cross-section and asymmetry data were analysed and put together by the LEP electroweak working group to obtain a prediction for the top quark mass of

 m_{top} = 177 ±11 (+18-19 for m_{H} =1000,100, 30), as kindly referred to in the CDF paper of April 1994 who reported an excess of 2.8 σ in that same mass range with best mass of 176±16 GeV.

CDF and D0 went on to discover the top in 1995, and LEP and SLC went on to predicting the Higgs mass using the top quark mass from the Tevatron.

....LHC discoverd the Higgs boson with m=125 GeV



The Nobel Prize in Physics 1999





Photo from the Nobel Foundation archive. Gerardus 't Hooft Prize share: 1/2

Photo from the Nobel Foundation archive. Martinus J.G. Veltman Prize share: 1/2

The Nobel Prize in Physics 1999 was awarded jointly to Gerardus 't Hooft and Martinus J.G. Veltman "for elucidating the quantum structure of electroweak interactions in physics"

in the 'advanced information'...written by C. Jarlskog For example, the mass of the top quark could be predicted, using high precision data from the accelerator LEP (Large Electron Positron) at the Laboratory CERN, Switzerland, several years before it was discovered, in 1995 at the Fermi National Laboratory in USA. The top quark, in spite of being too heavy to be produced at the LEP accelerator, contributed through quantum corrections by a measurable amount to several quantities that could be measured at LEP. Similarly, comparison of theoretical values of quantum corrections involving the Higgs Boson with precision measurements at LEP gives information on the mass of this as yet undiscovered particle.

FCG he final LEPI and SLC and Tevatron results can be found in arXiv:hep-ex/0509008 providing spectacular agreement of data with the Standard Model EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH STANFORD LINEAR ACCELERATOR CENTER

> CERN-PH-EP/2005-041 SLAC-R-774 hep-ex/0509008 7 September 2005

Precision Electroweak Measurements

on the Z Resonance

and a prediction for the Higgs boson mass of 129 +74-69 GeV.

The ALEPH, DELPHI, L3, OPAL, SLD Collaborations,¹ the LEP Electroweak Working Group,² the SLD Electroweak and Heavy Flavour Groups

Parameter	Value	Correlations				
		$\Delta \alpha_{\rm had}^{(5)}(m_{\rm Z}^2)$	$\alpha_{\rm S}(m_{\rm Z}^2)$	$m_{\rm Z}$	m_{t}	$\log_{10}(m_{\rm H}/{\rm GeV})$
$\Delta \alpha_{\rm had}^{(5)}(m_{\rm Z}^2)$	$0.02767 {\pm} 0.00034$	1.00				
$\alpha_{ m S}(m_{ m Z}^2)$	$0.1188 {\pm} 0.0027$	-0.02	1.00			
$m_{\rm Z} \; [{\rm GeV}]$	$91.1874 {\pm} 0.0021$	-0.01	-0.02	1.00		
$m_{\rm t} [{\rm GeV}]$	$178.5 {\pm} 3.9$	-0.05	0.11	-0.03	1.00	
$\log_{10}(m_{\rm H}/{\rm GeV})$	$2.11{\pm}0.20$	-0.46	0.18	0.06	0.67	1.00
$m_{\rm H} \; [{\rm GeV}]$	$129\pm^{74}_{49}$	-0.46	0.18	0.06	0.67	1.00

Table 8.3: Results for the five SM input parameters derived from a fit to the Z-pole results and $\Delta \alpha_{\rm had}^{(5)}(m_{\rm Z}^2)$, plus $m_{\rm t}$, $m_{\rm W}$, and $\Gamma_{\rm W}$ from Tevatron Run-I and LEP-II. The fit has a $\chi^2/{\rm dof}$ of 18.3/13, corresponding to a probability of 15%. See Section 8.4 for a discussion of the theoretical uncertainties not included here. The results on $m_{\rm H}$, obtained by exponentiating the fit results on $\log_{10}(m_{\rm H}/{\rm GeV})$, are also shown. The direct measurements of $m_{\rm W}$ and $\Gamma_{\rm W}$ used here are preliminary.



At the end of LEP: Phys.Rept.427:257-454,2006

 $N_v = 2.984 \pm 0.008$

This is determined from the Z line shape scan and dominated by the measurement of the hadronic cross-section at the Z peak maximum →

The dominant systematic error is the theoretical uncertainty on the Bhabha cross-section (0.06%) which represents an error of ± 0.0046 on N_v



Improving on N_v by more than a factor 10 would require a large effort to improve on the Bhabha cross-section calculation!

See Patrick's talk for a further avatar on this important measurement.



Verified a number of fundamentals of the SM at the 10⁻³ level

- -- universality of $e/\mu/\tau/\nu$ couplings @10⁻³ for NC (Z p. widths and asymetries) as well as CC (using tau decays)
- -- lepton quark universality (from ratio of hadron to lepton decay width)
- -- Measured essential input (Z mass at 10⁻⁵ level)
- -- observed effect of top mass and predicted the mass

before and independently of first direct observations

- -- constrained the Higgs boson mass [114.7(direct) →285 GeV(Rad.Cor.)]
- -- measured W mass, WW production and gauge couplings

Assuming NO BSM physics modifies the SM predictions

More lessons learned...

1. LEP exceeded expectations in almost every aspect that involved « systematicsdominated » measurements!

examples of precisions: Z mass and width (exp:20-50 MeV, achieved 2 MeV) (EPOL) luminosity measurement (exp ~2%, achieved ~6 10⁻⁴) $\rightarrow N_v$ $sin^2 \theta^{eff}_w$ (exp 0.001, achieved 0.00016) R_b (exp 2-5%? achieved 0.3%) etc etc

+ corresponding improvements in theory calulations and superb MC codes

2. also realized that things could have been better if better prepared difficulty in measuring the leptonic B.R. (end-cap design) did not expect all the difficulties in beam energy calibration still some measurements limited by th. unertainties Partly because spelled out 'expectations' were too conservative! its not because things are difficult that we dont dare... <u>its because we dont dare that things are difficult</u>



What happened since LEP?

LHC of course, and much more

The Nobel Prize in Physics 2013





© Nobel Media AB. Photo: A. Mahmoud François Englert Prize share: 1/2

Mahmoud Peter W. Higgs Prize share: 1/2

The Nobel Prize in Physics 2013 was awarded jointly to François Englert and Peter W. Higgs "for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"



The Nobel Prize in Physics 2015 Takaaki Kajita, Arthur B. McDonald

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The Nobel Prize in Physics 2015



Photo © Takaaki Kajita Takaaki Kajita Prize share: 1/2



Photo: K. MacFarlane. Queen's University /SNOLAB

Arthur B. McDonald Prize share: 1/2

The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald *"for the discovery of neutrino oscillations, which shows that neutrinos have mass"*

Particle physics after the discovery of the Higgs boson

Two facets:

The Higgs boson is very special

It generates (couples to) mass. Alone?

-- W,Z masses \Leftrightarrow Higgs coupling to WW, ZZ? FCC-ee -- (all) fermion masses \Leftrightarrow Higgs couplings? FCC-ee FCC-hh(+ee) decays ($\gamma\gamma$, gg, Z γ) \Leftrightarrow SM particle content? -- are all elementary particles given mass this way? FCC-ee(?) even electrons? and even the neutrinos? Yukawa ($\rightarrow v_R$, sterile \rightarrow Majorana HNL) FCC-ee

Higgs couples to itself!

- -- Spontaneous Symmetry Breaking
- -- What is the value of the self-coupling?
 - -- impact on σ_{HZ} near threshold
- FCC-ee

FCC-h

- -- HH production
 - FCC-hh, high energy lepton colliders

-- Higgs and top masses predicted from EWPOs assuming no new SM coupled particles exist

The SM is « complete »

- -- $m_H \approx 125 \text{GeV} \rightarrow \text{SM}$ extrapolates to the Plank scale assuming no new SM coupled particles exist
- -- SM works wonderfully... So why continue?

-- SM does not explain everything

Baryon Asymmetry of the Universe Dark Matter Neutrino masses and more.... → require new particles!

e.... → require new particles! FCC-ee: nature and mass scale is unknown EWPO Flavour

FCC-hh

Are there any further SM-coupled particles?

- -- no guarantee or exp. indication that any exist
 - -- but many BSM solutions include them...
 - -- DARK SECTOR \rightarrow possibly light, sterile particles

NC50 Orsay

FCC-ee: LLP EWPO Flavour

Particle physics after the discovery that neutrinos have mass

« Beyond the Standard Model » because SM is defined as having massless neutrinos

Neutrinos oscillate

3x3 oscillation → possibility of CP violation → T2K, HyperK, DUNE

'near future' (2030-2040..) and after that?

New degrees of freedom

Fermion number is no longer a conserved quantity
Neutrino coupling with Higgs boson
→ right handed neutrinos
minimal see-saw (see slides 44-51)
→ Heavy Majorana, sterile, Neutral Leptons



Sakharov condition for generation of the **Baryon Asymmetry of the Universe:**

- -- Fermion number violation
- -- CP or T violation and
- -- out-of-equilibrium universe (Big Bang)
- → Baryogenesis or Leptogenesis + sphalerons

Massive neutrinos are THE natural candidate to explain the dominance of matter over antimatter in the universe.



a hard look at the situation...

Since the NC discovery we have been relying on increasing collider energies for the next SM particle to show up... ... or else a drama would happen (t-less models, no-lose theorem, etc...)

This is no longer necessarily the case.

The SM-coupled particles predicted by the SM have all been found, yet unexplained phenomena are observed. (DM. BAU) It is quite possible that no more SM-coupled particle exist!

The question 'are there any more particles with SM couplings?' must be tested by all possible means!

→ Any <u>solid set of SM deviations would be a big discovery</u>

→ EW+Flavours at colliders and high precision facilities with several orders of magnitude increase of precision.

The new physics there is : Higgs boson and massive neutrinos.

What is predicted are sterile particles with couplings many many orders of magnitude smaller than SM and whose mass can vary between few keV and 10¹⁰ GeV... still must look for them wherever we can!

→ High precision, huge intensities and more energy are required.

A NEW ERA OF EXPLORATION

The FCC integrated program at CERN

Comprehensive cost-effective program inspired by successful LEP – LHC success story

- Stage 1: FCC-ee (Z, W, H, tt) as first generation Higgs EW and top factory at highest luminosities.
- Stage 2: FCC-hh (~100 TeV) as natural continuation at energy frontier, with ion and eh options.
- Maximizes physics output with strong complementarity
- Integrating an ambitious high-field magnet R&D program
- Common civil engineering and technical infrastructures, building on and reusing CERN's existing infrastructure.
- FCC-INT project plan is fully integrated with HL-LHC exploitation -> seamless continuation of HEP
- Feasibility study approved and funded at CERN (100MCHF/5yrs) + magnet R&D (120 MCHF/6yrs)

*** GLOBAL COLLABORATION ***





6. 10^{12} e+e- \rightarrow Z (qq)

1.5 10⁶ e+e- → ZH

O(5000) e+e- → H

2yrs 2. 10^8 e+e- \rightarrow WW

Event statistics (4IP)

WW threshold

ZH maximum

s-channel H

E_{cm}: 91 GeV

E_{cm} ≥ 157-161

E_{cm} : 240 GeV

 $E_{cm}: m_H$

4yrs

3yrs

(3yrs?)

 E_{cm} : 340-365 GeV 5yrs 2. 10⁶ e+e- \rightarrow tt

Z peak

tt

	E _{CM} errors:	
LEP x 3.10 ⁵	<100 keV	
LEP x 2.10 ³	<300 keV	
Never done	1 MeV	
Never done	<< 1 MeV	
Never done	2 MeV ₃₀	



FCC-ee: main machine parameters



Parameter	Z	ww	H (ZH)	ttbar
beam energy [GeV]	45	80	120	182.5
beam current [mA]	1280	135	26.7	5.0
number bunches/beam	10000	880	248	36
bunch intensity [10 ¹¹]	2.43	2.91	2.04	2.64
SR energy loss / turn [GeV]	0.0391	0.27	1.869	10.0
total RF voltage 400/800 MHz [GV]	0.120/2	CFRN	2.08/0	4.0/7.25
long. damping time [turns]	sight to		64.5	18.5
horizontal beta* [m]	roject	0.2	0.3	1
vertical beta* [mm]	0.8	1	1	1.6
horizontal g	0.71	2.17	0.64	1.49
vertical geon	1.42	4.34	1.29	2.98
horizontal rms IP spot size [µm]	8	21	14	39
vertical rms IP spot size [nm]	34	66	36	69
luminosity per IP [10 ³⁴ cm ⁻² s ⁻¹]	182	19.4	7.3	1.33
total integrated luminosity / year [ab ⁻¹ /yr] 4 IPs	87	9.3	3.5	0.65
beam lifetime (rad Bhabha + BS+lattice)	8	18	6	10
	4 years 5 x 10 ¹² Z LEP x 10 ⁵	2 years > 10 ⁸ WW LEP x 10 ⁴	3 years 2 x 10 ⁶ H	5 years 2 x 10 ⁶ tt pairs

F. Gianotti, P5 meeting 2023-04-15

□ x 10-50 improvements on all EW observables

- □ up to x 10 improvement on Higgs coupling (model-indep.) measurements over HL-LHC
- □ x10 Belle II statistics for b, c, т
- □ indirect discovery potential up to ~ 70 TeV
- □ direct discovery potential for feebly-interacting particles over 5-100 GeV mass range

Up to 4 interaction points \rightarrow robustness, statistics, possibility of specialised detectors to maximise physics output



Physics at FCC

L. HIGGS FACTORY



Higgs provides a very good reason why we need e+e- (or $\mu\mu$) collider

2. ELECTROWEAK PRECISION (10^{-3} today \rightarrow $10^{-5/6}$)

Z + WW + top required! ***CHALLENGES***

3. Z FACTORY

(6 10^{12} Z i.e. 1.5 10^{11} ee, $\mu\mu$, $\tau\tau$; ~0.7 10^{12} uu,dd,ss,cc,bb; 10^{12} vv)



High statistics for Heavy Flavours, QCD Search for Feebly-Coupled/sterile Particles (HNL, ALPS, etc) "Dark Sector" <u>Place for 'direct discovery'</u>

4. {90-120} TeV FCC-hh

The most powerful high energy exploration machine esp for any gluon-mediated process including W (>> 10¹³ from top decay) and Higgs (2 10¹⁰ from gg-> H) and direct searches in the multi TeV region (up to 50 TeV)

06.12.2023

NC50 Orsay

Motivation for the precision measurements *and* precision calculations

- 1. Given that the minimal SM is complete with the Higgs discovery, how do we find out:
- -- if the Higgs boson is exactly what is foreseen by the standard model? $(\rightarrow$ Higgs Factory)
- -- where/what are the new physics phenomena that must be present to explain:
 - baryon asymmetry dark matter, neutrino masses (and other mysteries we don't understand) (\rightarrow EW/top factory)

2. A powerful and broadly efficient method is to perform precision EW measurements
 -- many observables contain sensitivity to new phenomena, either by loops, direct long distance propagator effects, or mixing with SM coupled particles. Having many observables is essential to provide redundancy and point to the origin.

➔ are there any more weakly coupled particles?

«Т»

«S»

Weltman: Δρ =ΔT.α=
 α/π. (m²_{top}-m²_b)/m²_W
 any custodial SU(2)-violating effect appears <u>regardless of mass scale</u>
 -- is there mixing ? Z-Z' active-sterile neutrino mixing
 -- high mass SM-coupled and custodial SU(2)-respecting → (ex: Z' or degenerate SuSy)
 Emphasis on different observables depending on the question asked → different patterns of effects

26.04.2023

·				
Observable	present	FCC-ee	FCC-ee	Comment and
	value $\pm \text{ error}$	Stat.	Syst.	leading exp. error
$m_{\rm Z} ~({\rm keV})$	91186700 ± 2200	4	100	From Z line shape scan
				Beam energy calibration
$\Gamma_{\rm Z} \ ({\rm keV})$	2495200 ± 2300	4	25	From Z line shape scan
				Beam energy calibration
$\sin^2 \theta_{\rm W}^{\rm eff}(\times 10^6)$	231480 ± 160	2	2.4	from $A_{FB}^{\mu\mu}$ at Z peak
				Beam energy calibration
$1/\alpha_{\rm QED}({\rm m}_{\rm Z}^2)(\times 10^3)$	128952 ± 14	3	small	from $A_{FB}^{\mu\mu}$ off peak
				QED&EW errors dominate
$\mathbf{R}_{\ell}^{\mathbf{Z}}$ (×10 ³)	20767 ± 25	0.06	0.2 - 1	ratio of hadrons to leptons
				acceptance for leptons
$\alpha_{\rm s}({\rm m_Z^2})~(\times 10^4)$	1196 ± 30	0.1	0.4-1.6	from $\mathbf{R}_{\ell}^{\mathbf{Z}}$ above
$\sigma_{\rm had}^0 \ (\times 10^3) \ ({\rm nb})$	41541 ± 37	0.1	4	peak hadronic cross section
				luminosity measurement
$N_{\nu}(\times 10^3)$	2996 ± 7	0.005	1	Z peak cross sections
				Luminosity measurement
$R_b (\times 10^6)$	216290 ± 660	0.3	< 60	ratio of bb to hadrons
				stat. extrapol. from SLD
$A_{FB}^{0}, 0 \ (\times 10^{4})$	992 ± 16	0.02	1-3	b-quark asymmetry at Z pole
				from jet charge
$A_{FB}^{pol,\tau}$ (×10 ⁴)	1498 ± 49	0.15	<2	au polarization asymmetry
				au decay physics
τ lifetime (fs)	290.3 ± 0.5	0.001	0.04	radial alignment
$\tau \text{ mass (MeV)}$	1776.86 ± 0.12	0.004	0.04	momentum scale
τ leptonic $(\mu\nu_{\mu}\nu_{\tau})$ B.R. (%)	17.38 ± 0.04	0.0001	0.003	e/μ /hadron separation
$m_W (MeV)$	80350 ± 15	0.25	0.3	From WW threshold scan
	0005 1 40	1.0	0.0	Beam energy calibration
$\Gamma_{\mathbf{W}}$ (MeV)	2085 ± 42	1.2	0.3	From WW threshold scan
$(-2)(-10^4)$	1170 400	0	. 11	Beam energy calibration
$\alpha_{\rm s}({\rm m}_{\rm W}^{-})(\times 10^{-})$	1170 ± 420	3	small	from R_{ℓ}
$N_{\nu}(\times 10^{\circ})$	2920 ± 50	0.8	small	ratio of invis. to leptonic
(20.001/2)				in radiative Z returns
$m_{top} (MeV/c^2)$	172740 ± 500	17	small	From tt threshold scan
				QCD errors dominate
$\Gamma_{\rm top} ({\rm MeV/c}^{-})$	1410 ± 190	45	small	From tt threshold scan
i SM				QCD errors dominate
$\lambda_{ m top}/\lambda_{ m top}^{ m SW}$	1.2 ± 0.3	0.10	small	From tt threshold scan
		0.0		QCD errors dominate
ttZ couplings	$\pm 30\%$	0.5 - 1.5%	small	From $\sqrt{s} = 365 \mathrm{GeV}$ run

Precision EW measurements: is the SM complete?



- Higgs and EWPOs are complementary
- top quark mass and couplings essential!
 (the 100km circumference is optimal for this)
- preliminary systematics!

aim at reducing to the level of systematics

- many observables still to be added (flavours)
- complemented by high energy FCC-hh
- Theory work is critical and initiated
- 1809.01830 and several follow ups
- e c Target precision is statistics column

NC50 Orsay

FCCOverview of loop correction relationships and examples of new physics effects



uncertainty on $\Delta \alpha_{QED}$ impacts sin² θ^{eff}_{W} sensitivity. Direct measurement of $\Delta \alpha_{QED}$ by FCCee from $A_{FB}^{\mu\mu}$ (s) is UNIQUE ³⁵

Precision Natural SUSY at CEPC, FCC-ee, and ILC, JJ Fan, M. Rees and Liantao Wang

arXiv:1412.3107v2 figure 5 (top row)



Figure 5. Regions in the stop physical mass plane that are/will be excluded at 2σ by EWPT with oblique corrections (left column), R_b at FCC-ee (mid column) and Higgs couplings (right column) for different choices of $X_t/\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2}$: 0 (first row), 0.6 (2nd row), 1.0 (3rd row) and 1.4 (last row). We chose the mass eigenstate with $m_{\tilde{t}_1}$ to be mostly left-handed while the mass eigenstate with $m_{\tilde{t}_2}$ to be mostly right-handed. For non-zero choices of X_t , there are regions along the diagonal line which cannot be attained by diagonalizing a Hermitian mass matrix [32]. Also notice that the vacuum instability bound constrains $X_t/\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \lesssim \sqrt{3}$ [76].

" also, b \rightarrow sy could be useful"

Alain Blondel Le future des courants neutres

FCC



Huge statistics will help with systematics.... but detectors must be designed for this!

With statistics 300 000 times those of LEP → statistical precision improves by 500 even more for b, tau and charm observables because of improvement of vertex detector and smaller beam pipe

Target: improve systematics to the level of the statistical precision

Examples(all work in progress, or to be done)

- -- monitor luminosity with 10¹⁰ $\gamma\gamma$ events in addition to low angle Bhabha.
 - \rightarrow 1.5 10⁻⁵ precision provided fiducial volume limit (10-20 degrees) can be known to 10-15 microns at 2.5m.
- -- perform beam energy calibration continuously with pilot polarized bunches RDP and solenoid compensation
- -- use several tags (inclusive double tag, exclusive B decay modes tag, etc. etc.) for b-jet efficiency for R_b
- -- use huge samples of muon pairs (1.5 10¹¹) to perform exquisite detector alignment (for tau life time, luminosity etc...

--- etc etc

Success requires to be proactive and design detectors with these precision targets in mind.

→ a new era of detector design!

→ Considerable opportunities in physics studies during phase towards preparation of detector collaborations!



Huge statistics will help with systematics.... and theory must be ready for it!

Staszek Jadach got interested early in TLEP/FCCee

attracted by the physics potential of the electroweak measurements .

- ➔ Made many contributions at FCC weeks and physics workshops/meetings + publications
- → worked with/attracted many collaborators to follow on his pioneering ideas and the huge challenges



arbores serit, quorum poma non viderit

He plants trees, the fruits of which he will not see.

F Snapsot of ongoing efforts on theory side

Precision calculations for the Z line shape at the FCC-ee

S. Jadach et al, FCC week 2018 (Amsterdam)

+ I.Dubovyk A. Freitas, J. Gluza K. Grzankac T. Riemann J. Usovitsch

Next decade: complete 3-loop calculations [3]

$Z ightarrow e^+ \; e^-,$						
Number of	1 loop	2 loops	3 loops			
topologies	1	$14 \rightarrow^{(A)} 7 \rightarrow^{(B)} 5$	211 → ^(A) 84 → ^(B) 51			
Number of diagrams	14	2012 → ^(A,B) 880	397690 → ^(A,B) 91472			
Fermionic loops	0	114	13104			
Bosonic loops	14	766	78368			
Planar / Non-planar	14 / 0	782 / 98	65487 / 25985			
QCD / EW	0 / 14	0 / 880	144 / 91328			

Table 3: Presents the number of Z decay Feynman diagrams needed to be calculated for TH3 of Table 2. Tadpoles, products of lower loop diagrams (A) and symmetrical diagrams (B) are not included.

A first tackle might concentrate on the 13,104 electroweak 2-loop diagrams with closed internal fermionic loops, to be determined with a net accuracy of two relevant digits.



Needs for substantially improved theoretical analysis software:

- QED Monte Carlo code of the KKMC-type [S. Jadach et al.]
- Unfolding code of the SMATASY type [M. Grünewald et al.]
- Electroweak library of the ZFITTER type [T. Riemann et al.]

 σ^{meas} KKMC,

$$\xrightarrow{\cdots} \sigma^{real} \xrightarrow{\text{SMATASY,\cdots}}_{\text{ZFITTER,\cdots}} \begin{cases} \sigma_0 \equiv \sigma^{\text{eff},f} \\ M_Z, \Gamma_Z, \Gamma_f \\ A_{FB}^{\text{eff},f}, A_{LR}^{\text{eff},f} \\ R_b, R_\ell, R_{had} \\ \cdots \end{cases}$$

first thing: accept that the target precision is the statistical error!

C Centre of mass Energy Calibration: the cornerstone of the precision programme

Oide optics with Qx=0.1, Qv=0.2, Qs=0.05 Large ring \rightarrow transverse polarization of e[±] up to E_{beam} > 80 GeV $\sigma_{\rm F} \propto E^2/\sqrt{\rho}$ *E. Gianfelice* 100 Resonant depolarization provides high precision $E_{\text{beam}} v_s = \frac{g-2}{2} \frac{E_b}{m_a} = \frac{E_b}{0.4406486(1)}$ Linear SITROS 80 60 Polariza Unique to circular machines (ee or $\mu\mu$) 40 E [MeV] 20 Improve over LEP by using pilot bunches + both e- and e+ polarimeter 44717 5 ∕∂, Relationship between v_s and E_{CM} \rightarrow CM boost, σ_{ECM} , α_{coll} determined from 10⁶ $\mu\mu$ /5min C=97.75 km, 45.59 GeV, Q s=0.025, σ δ =0.000 Beamstahlung monitor under study etc... FP First round of studies (arxiv 1909.12245) FCC-ee simulation of 101.48 101.481 101.482 101.483 m_z , Γ_z , $\sin^2\theta_w^{eff}$, $\alpha_{OED}(m_z)$, m_w resonant depolarization next target: bring syst. closer to stat. errors, esp. pt-to-pt errors - 0.002 -0.0015-0.001-0.00050 0.0005 0.001 0.0015 0.002 Flipper frequency detuning: $\nu - \gamma a$ statistics ΔE_{CMabs} $\Delta E_{CMSyst-ptp}$ σE_{CM} Quantity calib. stats. $|200 \text{ keV}/\sqrt{(N^i)}|(84) \pm 0.05 \text{ MeV}$ 100 keV 40 keV To⁵ Spread (no BS) m_{Z} (keV) 100 28 Spread (BS) 4 $\sigma_{a_a} = 0.1 \text{ mrad}$ $\Gamma_{\rm Z}$ (keV) 22 2.5 10 With ISB $sin^2 \theta_W^{\text{eff}} \times 10^6 \text{ from } A_{FB}^{\mu\mu}$ Asymmetry = ± 0.1 2.4 2 0.1 $\frac{\Delta \alpha_{QED}(M_Z)}{\alpha_{QED}(M_Z)} \times 10^5$ 10^{4} 0.9 3 0.10.05 10^{3} At our luminosity level, longitudinal polarization brings nothing that cannot be done otherwise. future des courants ne 10²_5 NC50 Orsav -2 0 2 3 4 5 Longitudinal Boost, x

FCC

From FCC week 2022 and FCC EPOL workshop:





-- Resonant depolarization measures energy every 15 minutes at < 50 keV/beam level at Z, 100 keV/beam at W

→ syst will be reduced to < 100keV on m_w at Z point to point uncertainties remain to be understood

-- Only one RF station around the ring, + the Energy losses of the two beams are strongly constrained from the direct measurement of boost at the IPs O(5 keV level) every 8hrs shift in 2/4 experiments

-- beam-beam deflection measurement is extremely sensitive to beam beam offset and local opposite-sign dispersion (previously large point-to-point error): still lots to do but O(20 keV) per measurement every 3 second

→ targeting to match or go below statistical errors of 4 keV EPOL working group (Keintzel, Wilkinson, et al) establishing requirements

FCC

Electroweak Physics

Z factory + WW + top The realm of FCC-ee

Highest luminosities at 91, 160 and 350 GeV Transverse pol. at 91 and 160 GeV \rightarrow Ecm calibration m_z (100 keV) Γ_z (25 keV), m_w (<500 keV), $\alpha_{\rm QED}$ (m_z) (3.10⁻⁵) and sin² $\theta_w^{\rm eff}$ (1.5 10⁻⁶⁾

Complete set of EW observables can be measured Precision unique to FCC-ee + new physics sensitivity → a lot more potential to exploit with detector design than present treatment suggests

The reach for new physics depends on the new physics: -- new non-degenerate SU(2) doublet should not have mass splitting greater than ~5 GeV

- -- Heavy Neutrino mixing limit ~ 10^{-5} mixing up to 500-1000 TeV
- -- $1/\Lambda^2$ new physics \rightarrow 30-70 TeV



FCC The Flavour Factory

More on TeraZ

Progress in flavour physics wrt SuperKEKb/BELLEII requires > 10¹¹ b pair events,

FCC-ee(Z): will provide ~10¹² b pairs. "Want at least 5 10¹² Z..."

- necision of CKM matrix elements
- Push forward searches for FCNC, CP violation and mixing
- $● Study rare penguin EW transitions such as b → s <math>\tau_+ \tau_-$, spectroscopy (produce b-baryons, B_s ...)
- Test lepton universality with $10^{11} \tau$ decays (with τ lifetime, mass, BRs) at 10^{-5} level, LFV to 10^{-10}

-- all very important to constrain / (provide hints of) new BSM physics.

need special detectors (PID); a story to be written!

The 5 × 10¹² hadronic Z decays also provide precious input for QCD studies

High-precision measurement of $\alpha_s(mz)$ with Re in Z and W decay, jet rates, τ decays, etc. : 10⁻³ \rightarrow 10⁻⁴ huge \sqrt{s} lever-arm between 30 GeV and 365 GeV, fragmentation, baryon production Testing running of α_s to excellent precision with hadron production from low energy ($\gamma^*/Z^* + \gamma$) to 365 GeV

And... H→gg is a pure gluon factory (100'000 H→ gg events)! Alain Blondel Le future des courants neutres NC50 Orsay

Is *not* in itself a law or a symmetry of the Standard Model

For charged fermions (e/mu/tau and the quarks) it is not possible to transform a fermion into an antifermion because of charge conservation

For neutrinos, which are neutral, the SM assumes they are massless. neutrino is left-handed (identical if massless to negative helicity) and the antineutrino has positive helicity neutrino <-> antineutrino transition is forbidden by **angular momentum conservation**

This results in practice in apparent, accidental, conservation of fermion number

The existence of massive neutrinos allows for spin flip and thus in principle a neutrino-antineutrino transition since a left-handed field (EW eigenstate) has a component of the opposite helicity (EW state \neq physical state) $v_{L} \approx v_{-} + v_{+} m/E$ (mass is what allows to flip the helicity)

for the allowed masses of light neutrinos this is very, very small: for $m_v = 50$ meV and $P_{\pi}^* = 30$ MeV \rightarrow (m/E)² = 10⁻¹⁸

This can be observed in neutrino less double beta decay or by searching directly for the right-handed neutrinos



NEUTRINO MASSES

Electroweak eigenstates



Along with 'Antimatter,' and 'Dark Matter, we've recently discovered the existence of "Doesn't Matter,' which appears to have no." Alain Blondel Neutrino Physics II effect on the universe whatsoever.

NB unlike for v_1 , no interaction distinguishes particle and antiparticle of v_R which is a singlet (no 'charge') \rightarrow naturally a Majorana particle 45

Neutrino masses occur via processes which are intimately related to the Higgs boson what are the couplings of the H(125) to neutrinos?

Let us follow the steps of the Standard Model to construct a minimal neutrino mass model

Adding neutrino masses to the Standard model 'simply' by adding a Dirac mass \rightarrow right-handed neutrino

 $m_{D}\overline{\nu_{L}}\nu_{R}$

m_D is the Higgs **Yukawa coupling** (like everybody else). Then the right handed neutrinos are sterile, (**except** that they couple to both the Higgs boson and gravitation). Things become more interesting: **a Majorana mass term** arises (So-called **Weinberg Operator**) using the Higgs boson and the neutrino Yukawa coupling:

Alain Blondel Neutrino Physics II

Origin of neutrino mass:





Pilar Hernandez,

Granada 2019-05

Majorana mass term is extremely interesting as this is the particle-to-antiparticle transition that we want in order to explain the Baryon asymmetry of the Universe (+ CP violation in e.g. neutrinos)

mD



B. Kayser 1989)

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Having two mass terms per family , neutrinos undergo level splitting -> Mass eigenstates

See-saw type I :

$$\mathcal{L} = \frac{1}{2} (\bar{\nu}_L, \ \bar{N}_R^c) \begin{pmatrix} 0 & m_D \\ m_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ N_R \end{pmatrix}$$

 $M_{R} \neq 0$ $m_{D} \neq 0$ <u>Dirac + Majorana</u> <u>mass terms</u>

$$\tan 2\theta = \frac{2 m_D}{M_R - 0} \ll 1$$

$$m_{\nu} = \frac{1}{2} \begin{bmatrix} (0 + M_R) - \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ M = \frac{1}{2} \begin{bmatrix} (0 + M_R) - \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ (0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \end{bmatrix} \simeq -m_D^2/M_R$$

$$m_L = \frac{M_R = 0}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R = 0}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R \neq 0}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R \neq 0}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R \neq 0}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R \neq 0}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R} + \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4 m_D^2} \\ m_L = \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R} + \frac{M_R}{\frac{1}{2} \left[(0 + M_R) + \frac{M_R}{\frac{1}{2}$$

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Manifestations of right handed neutrinos

one family see-saw :	$v = v_L \cos\theta - N^c_R \sin\theta$	v = light mass eigenstate N = heavy mass eigenstate HNL
$m_v \approx \frac{m_D^2}{M}$	$N = N_R \cos\theta + v_L^{c} \sin\theta$	$ eq oldsymbol{v}_L$, active neutrino which couples to weak inter.
$m_{\rm N} \approx {\rm M}$ $ {\rm U} ^2 \propto \theta^2 \approx m_v / m_{\rm N}$	what is produced in W, Z decays is: $v_L = v \cos\theta + N \sin\theta$	and $\neq N_R$, which does'nt.

- -- mixing with active neutrinos leads to various observable consequences
 - -- if very light (eV) , possible effect on neutrino oscillations ('eV sterile neutrino'

(LSND/miniBooNE/reactor anomalies etc... but ruled out since PLANCK mission

MINOS/ICECUBE/DAYABAY/microBooNE. Search still ongoing in broader region)

- -- if in 5-100 keV region (dark matter), monochromatic photons from galaxies with $E=m_N/2$, KATRIN
- -- possibly measurable effects at High Energy
 - \rightarrow If N is heavy it will decay in the detector \rightarrow spectacular
 - \rightarrow Higgs, Z, W visible exotic decays H \rightarrow v_i \overline{N}_i and Z \rightarrow v_i \overline{N}_i , W-> I_i \overline{N}_i
 - → also in K, charm and b decays via W^{*}-> $I_i \pm N$, N → $I_j \pm$ with any of six sign and lepton flavour combination
 - \rightarrow violation of unitarity and lepton universality in Z, W or τ decays
 - → PMNS matrix unitarity violation and deficit in Z «invisible» width N_v < 3 (C. Jarlskog 1990)</p>
 -- etc... etc...

-- Couplings are very small ($|U|^2 = m_v / m_N$) for one family. For three families they can be somewhat larger but most interesting region is near the one-family see-saw limit ⁴⁸

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FCC

Heavy Neutral Leptons -- recent litterature

The Present and Future Status of Heavy Neutral Leptons

Asli M. Abdullahi, Pablo Barham Alzas, Brian Batell, Alexey Boyarsky, Saneli Carbajal, Animesh Chatterjee, Jose I. Crespo-Anadon, Frank F. Deppisch, Albert De Roeck, Marco Drewes, Alberto Martin Gago, Rebeca Gonzalez Suarez, Evgueni Goudzovski, Athanasios Hatzikoutelis, Marco Hufnagel, Philip Ilten, Alexander Izmaylov, Kevin J. Kelly, Juraj Klaric, Joachim Kopp, Suchita Kulkarni, Mathieu Lamoureux, Gaia Lanfranchi, Jacobo Lopez-Pavon, Oleksii Mikulenko, Michael Mooney, Miha Nemevsek, Maksym Ovchynnikov, Silvia Pascoli, Ryan Plestid, Mohamed Rashad Darwish, Federico Leo Redi, Oleg Ruchayskiy, Richard Ruiz, Mikhail Shaposhnikov, Ian M. Shoemaker, Robert Shrock, Alex Sousa, Nick Van Remortel, Vsevolod Syvolap, Volodymyr Takhistov, Jean-Loup Tastet, Inar Timiryasov, Aaron C. Vincent, Jaehoon Yu

The existence of non-zero neutrino masses points to the likely existence of multiple SM neutral fermions. When such states are heavy enough that they cannot be produced in oscillations, they are referred to as Heavy Neutral Leptons (HNLs). In this white paper we discuss the present experimental status of HNLs including colliders, beta decay, accelerators, as well as astrophysical and cosmological impacts. We discuss the importance of continuing to search for HNLs, and its potential impact on our understanding on key fundamental questions, and additionally we outline the future prospects for next-generation future experiments or upcoming accelerator run scenarios.

Comments: 82 pages, 34 figures. Contribution to Snowmass 2021
Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - E
Cite as: arXiv:2203.08039 [hep-ph]
(or arXiv:2203.08039v1 [hep-ph] for this version)
https://doi.org/10.48550/arXiv.2203.08039 1

High Energy Physics - Experiment

[Submitted on 10 Mar 2022 (v1), last revised 11 Mar 2022 (this version, v2)]

Searches for Long-Lived Particles at the Future FCC-ee

J. Alimena, P. Azzi, M. Bauer, A. Blondel, M. Drewes, R. Gonzalez Suarez, J. Klaric, S. Kulkarni, M. Neubert, C. Rizzi, R. Ruiz, L. Rygaard, A. Sfyrla, T. Sharma, A. Thamm, C. B. Verhaaren

The electron-positron stage of the Future Circular Collider, FCC-ee, is a frontier factory for Higgs, top, electroweak, and flavour physics. It is designed to operate in a 100 km circular tunnel built at CERN, and will serve as the first step towards \geq 100 TeV proton-proton collisions. In addition to an essential and unique Higgs program, it offers powerful opportunities to discover direct or indirect evidence of physics beyond the Standard Model. Direct searches for long-lived particles at FCC-ee could be particularly fertile in the high-luminosity Z run, where $5 \times 10^{12} Z$ bosons are anticipated to be produced for the configuration with two interaction points. The high statistics of Higgs bosons, W bosons and top quarks in very clean experimental conditions could offer additional opportunities at other collision energies. Three physics cases producing long-lived signatures at FCC-ee are highlighted and studied in this paper: heavy neutral leptons (HNLs), axion-like particles (ALPs), and exotic decays of the Higgs boson. These searches motivate out-of-the-box optimization of experimental conditions and analysis techniques, that could lead to improvements in other physics searches.

Comments: Contribution to Snowmass 2021

Subjects: High Energy Physics - Experiment (hep-ex); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th)

- Cite as: arXiv:2203.05502 [hep-ex] (or arXiv:2203.05502v2 [hep-ex] for this version)
 - https://doi.org/10.48550/arXiv.2203.05502

777 references!



F This picture from the briefing book is relevant to Neutrino, Dark sectors and High Energy Frontiers. FCC-ee (Z) compared to the other machines for right-handed (sterile) neutrinos How close can we get to the 'see-saw limit'?



-- the purple line shows the 95% CL limit if no HNL is observed. (here for 10^{12} Z), -- the horizontal line represents the sensitivity to **mixing of neutrinos** to the dark sector, using EWPOs (G_F vs sin² θ_W^{eff} and m_z, m_W, tau decays) which extends sensitivity ^{06.12.} to 10⁻⁵ mixing all the way to very high energies (500-1000 TeV at least). arxiv:2011.04725



We absolutely need a next accelerator but the next facility must be versatile (and feasible!) with as broad and powerful reach as possible, as there is no precise target

more Sensitivity, more Precision, more Energy

FCC@CERN, thanks to local synergies and internal complementarities, offers the most versatile and adapted response to today's physics landscape

The huge step in statistics (and precision) is extremely challenging on all accounts Accelerator, Detectors, Theory must plan proactively to match the challenges!



THANK YOU STASZEK!

for your immense contributions and insight for the extraction of physics from LEP data For your irreplaceable enthusiasm and pragmatism in planning the FCCee 'impossible' precision For your vision!





arbores serit, quorum poma non viderit

He plants trees, the fruits of which he will not see.