

Possibilities of light-by-light scattering measurement in FoCal detector

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$$\gamma\gamma \rightarrow \gamma\gamma$$

O. Halpern, *Scattering processes produced by electrons in negative energy states.*, Physical Review 44:10 855-85, 1933.

- In the limit of weak e-m fields:

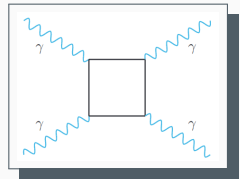
$$\sigma \sim \left(\frac{e^2}{mc^2}\right)^4 \left(\frac{\hbar}{mc}\right)^2 \cdot \frac{1}{\lambda^2}$$

H. Euler, B. Kockel, *Über die Streuung von Licht an Licht nach der Diracschen Theorie*, 1935*

- For visible light:

$$\sigma \approx 10^{-76} \text{cm}^2 = 10^{-52} \text{b} = 10^{-43} \text{nb}$$

* "The experimental test of the deviation from the Maxwell theory is difficult since the noteworthy effects are extraordinarily small."



Fermionic box.

Equivalent Photon Approximation

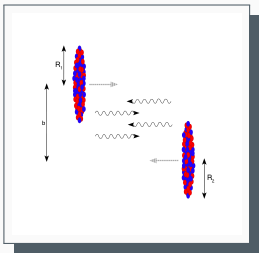
- Nuclear cross section:

$$\sigma_{A_1 A_2 \rightarrow A_1 A_2 X_1 X_2} = \int \frac{d\sigma_{\gamma\gamma \rightarrow X_1 X_2}(W_{\gamma\gamma})}{d\cos\theta} \times N(\omega_1, b_1) N(\omega_2, b_2) S_{abs}^2(b)$$

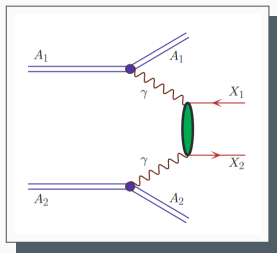
$$\times \frac{W_{\gamma\gamma}}{2} dW_{\gamma\gamma} dY_{X_1 X_2} d\bar{b}_x d\bar{b}_y d^2b \times \frac{d\cos\theta}{dy_{X_1} dy_{X_2} dp_t} \times dy_{X_1} dy_{X_2} dp_t$$

- Photon flux:

$$N(\omega, b) = \frac{Z^2 \alpha_{em}}{\pi^2 \beta^2} \frac{1}{\omega b^2} \times \left| \int d\chi \chi^2 \frac{F(\frac{\chi^2 + u^2}{b^2})}{\chi^2 + u^2} J_1(\chi) \right|^2 \quad u = \frac{\omega b}{\gamma\beta} \quad \chi = k_{\perp} b$$



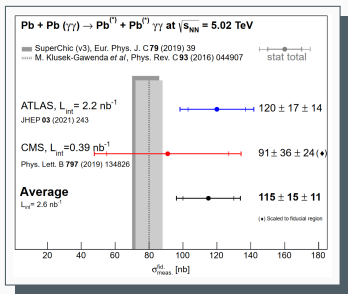
Collision geometry.



Photon fusion.

Year	Experiment	$p_{t,min}^\gamma$ [GeV]	$M_{\gamma\gamma,min}$ [GeV]	$\sigma_{tot.}^{exp.}$ [nb]	$\sigma_{tot.}^{theo.}$ [nb]
2017	ATLAS	3	6	70 ± 29	51 ± 5
2018	CMS	2	5	120 ± 55	103 ± 10
2019	ATLAS	2.5	5	120 ± 22	80 ± 8

Total cross section for light-by-light scattering in collisions with energy $\sqrt{s_{NN}} = 5.02$ TeV, in range of photon rapidity $|y| < 2.4$; $p_{t,min}^\gamma$ is a minimal measured value of photon transverse momentum of single photon, $M_{\gamma\gamma,min}$ is a diphoton invariant mass.



Current experiments have a high minimum threshold for transverse momentum of photon and diphoton invariant mass:

$$p_{t,min}^\gamma > 2 \text{ GeV}$$

$$M_{\gamma\gamma} > 5 \text{ GeV}$$

G. K. Krintiras, I. Grabowska-Bold, M. Klusek-Gawenda, É. Chapon, R. Chudasama and R. Granier de Cassagnac, *Light-by-light scattering cross-section measurements at LHC*. arXiv:2204.02845, 2022.

The average light-by-light scattering cross section value along with the individual cross section measurements at 5.02 TeV by ATLAS and CMS.

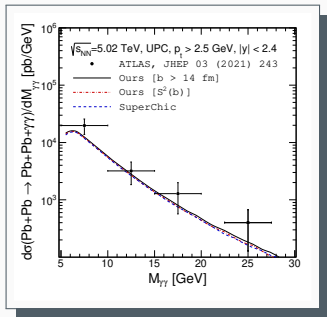
Results for ATLAS experiment

- Sharp edge of nucleus:

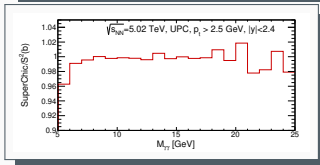
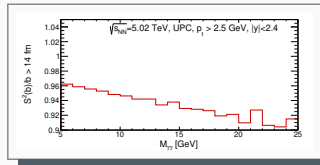
$$b_{min} = R_{A_1} + R_{A_2} \approx 14 \text{ fm}$$

- Smooth edge of nucleus:

$$S_{abs}^2(b) = \exp\left(-\sigma_{NN} \int d^2\rho T_A(\vec{\rho} - \vec{b}) T_A(\rho)\right), \quad T_A(\vec{\rho}) = \int \rho_A(\vec{r}) dz$$



Differential cross section as function of diphoton invariant mass.



Ratios between different theoretical approaches.

Results for ATLAS experiment

Kinematical cuts for new ATLAS measurement:

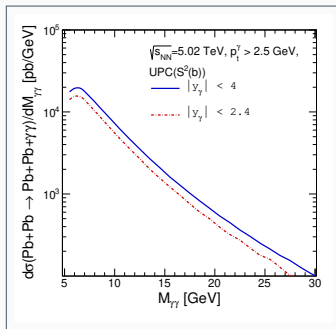
$$|y_\gamma| < 4$$

$$p_t^\gamma > 2.5 \text{ GeV}$$

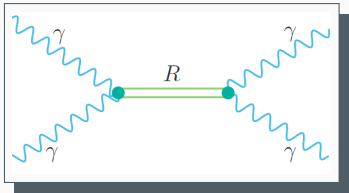
The ATLAS collaboration, *Expected tracking and related performance with the updated ATLAS Inner Tracker layout at the High-Luminosity LHC*, ATL-PHYS-PUB-2021-024.

$ y_\gamma <$	$\sigma_{tot.}^{theo.}$ [nb]	$\sigma_{tot.}^{exp.}$ [nb]
2.4	77.084 ± 0.005	120 ± 22
4	100.444 ± 0.027	planned

Total cross section for light-by-light scattering in collisions with energy $\sqrt{s_{NN}} = 5.02$ TeV, in range of photon transverse momentum $p_t^\gamma < 2.5$ GeV and diphoton invariant mass $M_{\gamma\gamma} > 5$ GeV; y_γ is a rapidity of measured single photon.

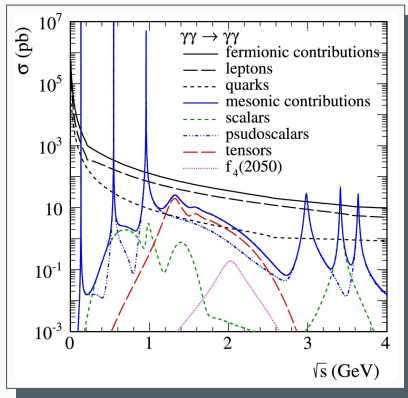


Differential cross section as function of diphoton invariant mass for future and recent ATLAS measurement.

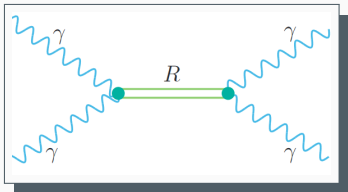


Mesonic resonance.

P. Lebedowicz, A. Szczurek, *The role of meson exchanges in light-by-light scattering*, Physics Letters B, 772:330-335, 2017.



Elementary cross section for different $\gamma\gamma \rightarrow \gamma\gamma$ processes.



Mesonic resonance.

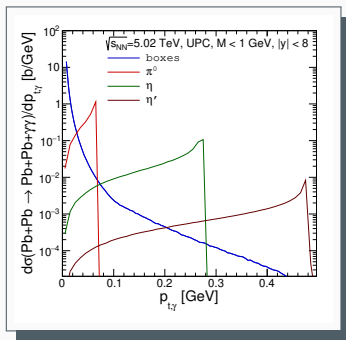
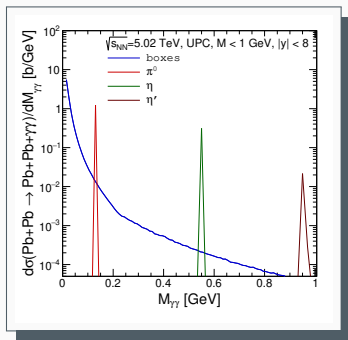
P. Lebedowicz, A. Szczurek, *The role of meson exchanges in light-by-light scattering*, Physics Letters B, 772:330-335, 2017.

$$\frac{d\sigma_{\gamma\gamma \rightarrow R \rightarrow \gamma\gamma}(W_{\gamma\gamma})}{d\cos\theta} = \frac{1}{32\pi W_{\gamma\gamma}^2} \frac{1}{4} \sum_{\lambda_1, \lambda_2} |\mathcal{M}_{\gamma\gamma \rightarrow R \rightarrow \gamma\gamma}(\lambda_1, \lambda_2)|^2$$

$$\mathcal{M}_{\gamma\gamma \rightarrow R \rightarrow \gamma\gamma}(\lambda_1, \lambda_2) = \frac{\sqrt{64\pi^2 W_{\gamma\gamma}^2 \Gamma_R^2 B r^2 (R \rightarrow \gamma\gamma)}}{\hat{s} - m_R^2 - im_R \Gamma_R} \times \frac{1}{\sqrt{2\pi}} \delta_{\lambda_1 - \lambda_2}$$

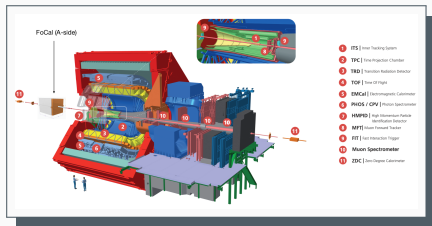
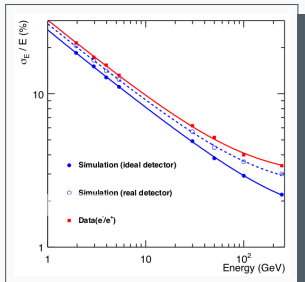
Relativistic Breit-Wigner distribution

The distributions of nuclear cross section as function of diphoton invariant mass and single photon transverse momentum for fermionic loops and resonances.



- Acceptance:
 $3.4 < y_\gamma < 5.8$
 $p_t^\gamma > 200 \text{ MeV}$
- Position resolution:
 $\sigma_x = \sigma_y = 1 \text{ mm}$
- Energy resolution:

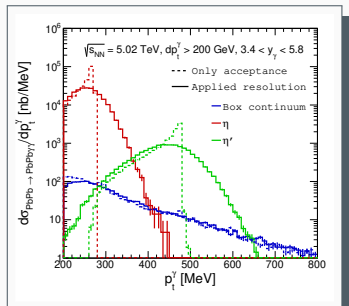
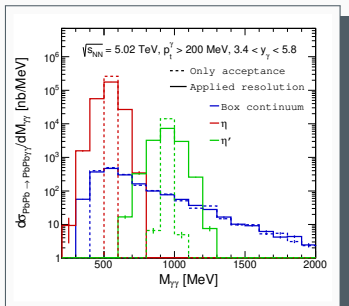
$$\frac{\sigma_E}{E} = \frac{28.5\%}{\sqrt{E(\text{GeV})}} + \frac{6.3\%}{E(\text{GeV})} + 2.95\%$$



A.P. de Haas et al. (ALICE Collaboration), *The FoCal prototype – an extremely fine-grained electromagnetic calorimeter using CMOS pixel sensors*, JINST 13 P01014, 2018.

C. Loizides, *The Forward Calorimeter project in ALICE*, EF06 meeting 2020,
https://indico.fnal.gov/event/44126/contributions/191953/attachments/132434/162766/20200805_focal_snowmass.pdf

Results of combined theoretical results for light-by-light scattering and Monte Carlo simulation of energy and position resolution for diphoton invariant mass and transverse momentum of measured photon for FoCal detector.



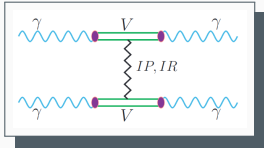
- Light-by-light scattering is a fundamental prediction of QED.
- Ultrarelativistic, ultraperipheral collisions of heavy ions allow observations of photon-photon processes hitherto not accessible.
- Light-by-light scattering not only provides evidence of the quantum nature of the electromagnetic interaction, but is also a tool to test the limits of theoretical models, and it provides a basis for the exploration of so-called "New Physics" (axions, dark photons).
- Future experiments such as FoCal and ATLAS will improve statistics and extend the kinematic ranges of measurements, allowing theoretical predictions to be tested



P. Jucha, M. Klusek-Gawenda, A. Szczurek, *Light-by-light scattering in ultraperipheral collisions of heavy ions at two future detectors*, Physical Review D 109, 014004, 2024.

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BACKUP



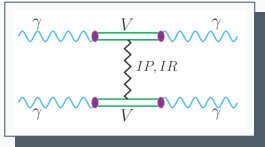
M. Klusek-Gawenda, P. Lebiedowicz, A. Szczurek
Light-by-light scattering in ultraperipheral PbPb collisions at the Large Hadron Collider,
 Phys. Rev. C, 93:4, 2016.

VDM-Regge: Fluctuation of γ into virtual ρ, ω, ϕ .

$$\frac{d\sigma_{\gamma\gamma\rightarrow\gamma\gamma}}{d\Omega} = \frac{1}{64\pi^2s} |\mathcal{A}_{\gamma\gamma\rightarrow\gamma\gamma}|^2$$

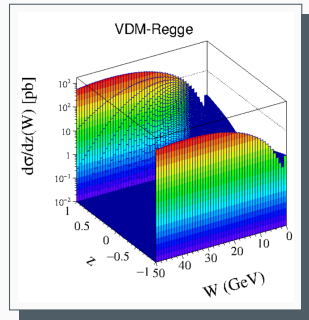
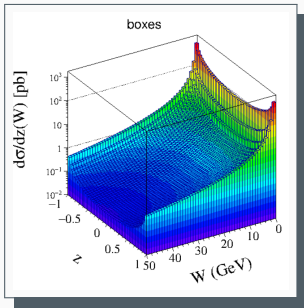
$$\mathcal{A}_{\gamma\gamma\rightarrow\gamma\gamma}(s, t) \approx \left(\sum_{i=1}^3 C_{\gamma\rightarrow V_i}^2 \right) \mathcal{A}(s, t) \exp\left(\frac{B}{2}t\right) \left(\sum_{j=1}^3 C_{\gamma\rightarrow V_j}^2 \right)$$

$$\mathcal{A}(s, t) \approx s \left((1+i)C_R \left(\frac{s}{s_0}\right)^{\alpha_R(t)-1} + iC_P \left(\frac{s}{s_0}\right)^{\alpha_P(t)-1} \right),$$



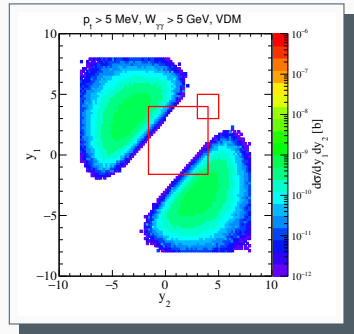
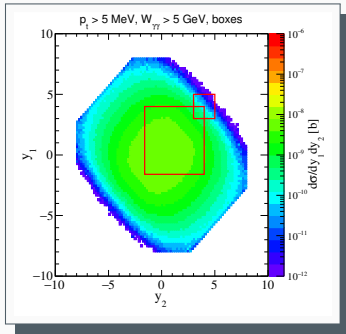
M. Klusek-Gawenda, P. Lebedowicz, A. Szczurek
Light-by-light scattering in ultraperipheral PbPb collisions at the Large Hadron Collider,
 Phys. Rev. C, 93:4, 2016.

VDM-Regge: Fluctuation of γ into virtual ρ, ω, ϕ .



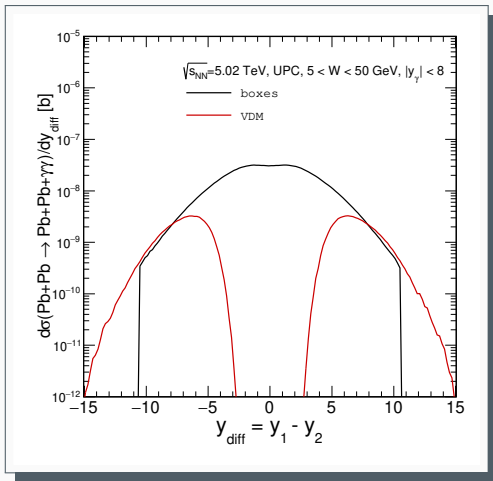
VDM-Regge - nuclear cross section

The differential cross-section for the VDM-Regge process in UPC, Pb-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV.



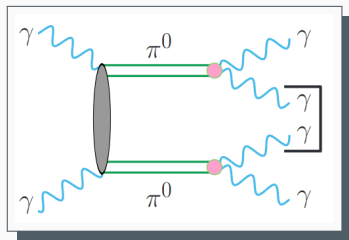
The red frames mark the acceptance range of the ALICE 3 detectors.

Wide rapidity range is key to the experimental observation of VDM.



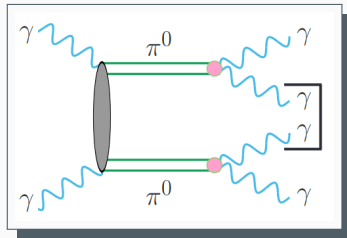
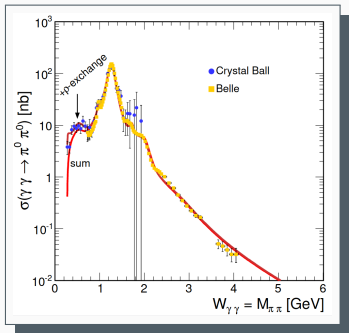
Introduction of the new variable: difference of photons rapidity.

$$\frac{d\sigma_{\gamma\gamma \rightarrow \pi^0\pi^0}(W_{\gamma\gamma})}{d\cos\theta} = \frac{\sqrt{\frac{W_{\gamma\gamma}^2}{4} - m_\pi^2}}{\frac{W_{\gamma\gamma}}{2}} \frac{4\pi}{4 \times 64\pi^2 W_{\gamma\gamma}^2} \sum_{\lambda_1, \lambda_2} |\mathcal{M}_{\gamma\gamma \rightarrow \pi^0\pi^0}(\lambda_1, \lambda_2)|^2$$



Two pion production.

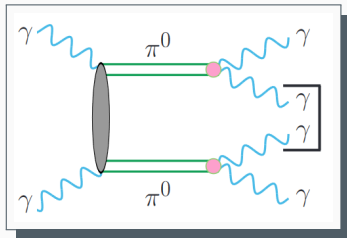
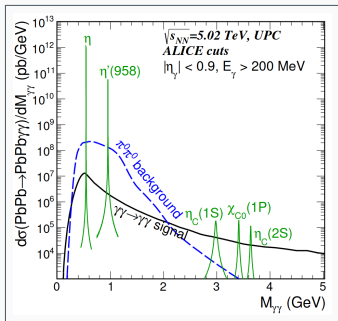
M. Klusek-Gawenda, A. Szczurek, $\pi^+\pi^-$ and $\pi^0\pi^0$ pair production in photon-photon and in ultraperipheral ultrarelativistic heavy-ion collisions, Phys. Rev. C87 (2013) 054908;



Two pion production.

Elementary cross section for $\gamma\gamma \rightarrow \pi^0\pi^0$ for $|\cos\theta| < 0.8$.

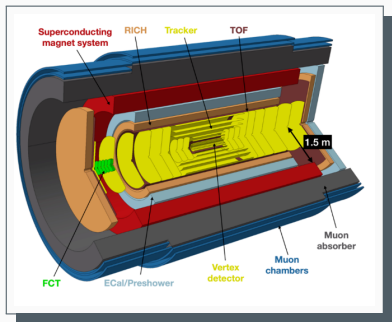
M. Klusek-Gawenda, A. Szczurek, $\pi^+\pi^-$ and $\pi^0\pi^0$ pair production in photon-photon and in ultraperipheral ultrarelativistic heavy-ion collisions, Phys. Rev. C87 (2013) 054908;



Two pion production.

Differential cross section as a function of invariant diphoton mass within the ALICE fiducial region.

M. Klusek-Gawenda, R. McNulty, R. Schicker, and A. Szczurek. *Light-by-light scattering in ultraperipheral heavy-ion collisions at low diphoton masses*, Phys. Rev. D, 99(9):093013, 2019.

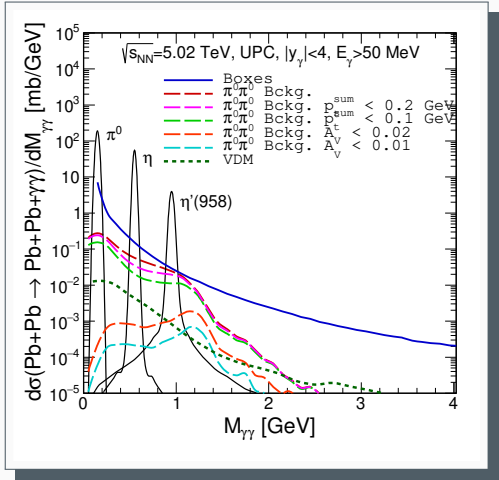


$p_{t,min}$ [GeV]	$p_{t,max}$ [GeV]	y_{min}	y_{max}
0.001	0.1	3	5
0.2	50	-1.6	4

Assumed kinematic limits in ALICE 3 experiment for photon measurement.

ALICE 3 - a next-generation heavy-ion detector for the LHC Runs 5-6.

L. Musa, W. Riegler, *Letter of intent for ALICE 3: A next generation heavy-ion experiment at the LHC*, arXiv:2211.02491, 2022.



Differential cross section as function of diphoton invariant mass for future ALICE 3 experiment.