

# Novel approach to measure quark/gluon jets at the LHC

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11.1.2024

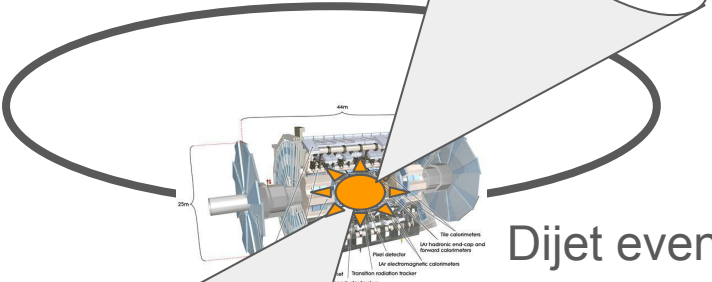


1. Introduction
2. Motivation
3. Results
4. Conclusion

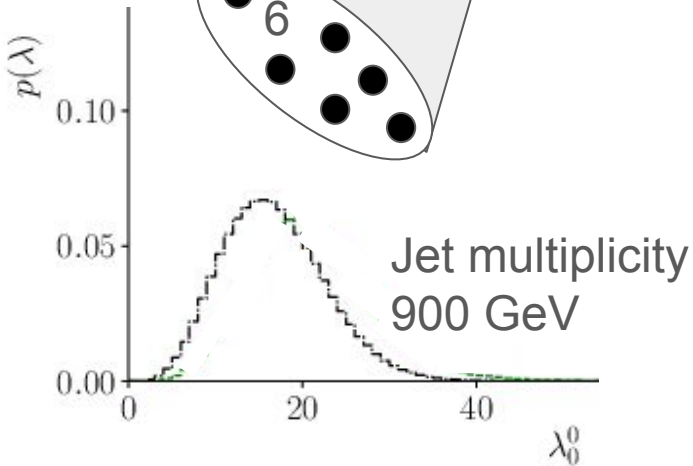
Back-up

# Introduction

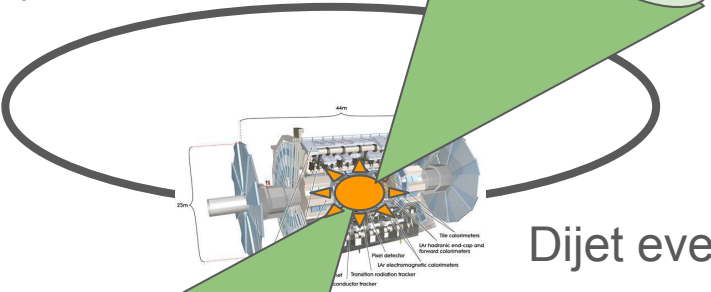
$\sqrt{s} = 900 \text{ GeV}$



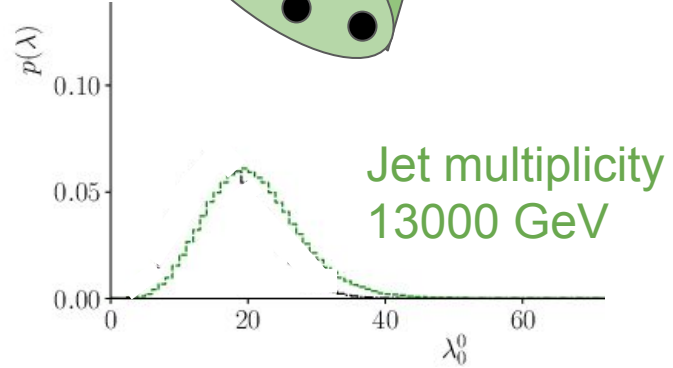
Dijet events



$\sqrt{s} = 13000 \text{ GeV}$

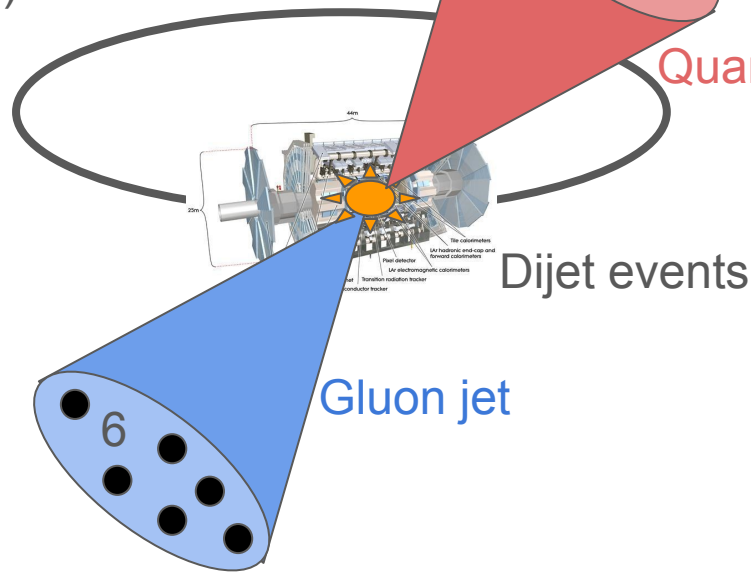


Dijet events

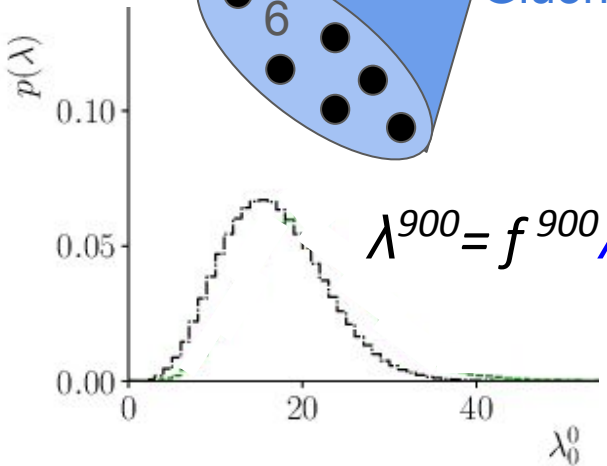
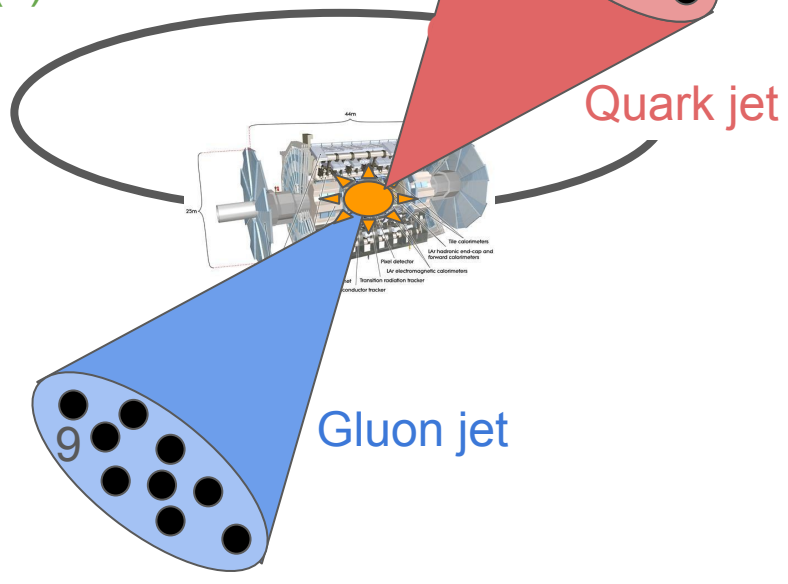


# Introduction

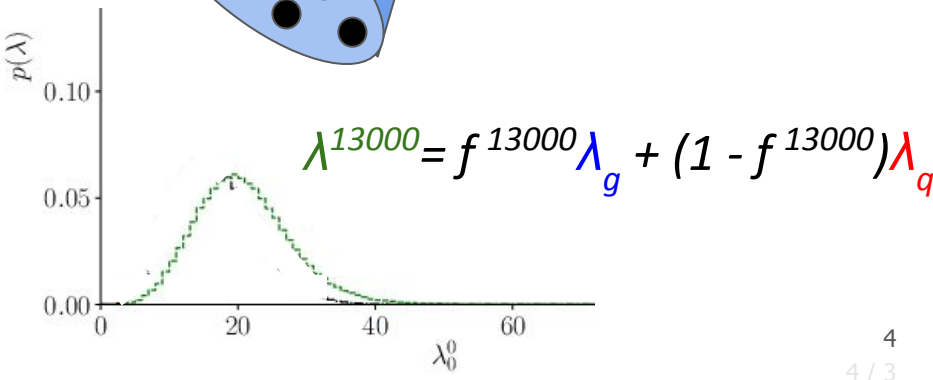
sqrt(s) = 900 GeV



sqrt(s) = 13000 GeV



$$\lambda^{900} = f^{900} \lambda_g + (1 - f^{900}) \lambda_q$$

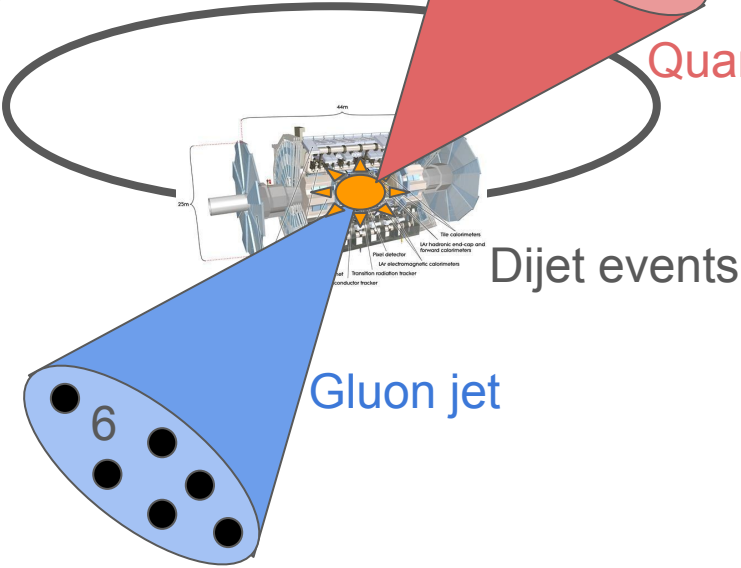


$$\lambda^{13000} = f^{13000} \lambda_g + (1 - f^{13000}) \lambda_q$$

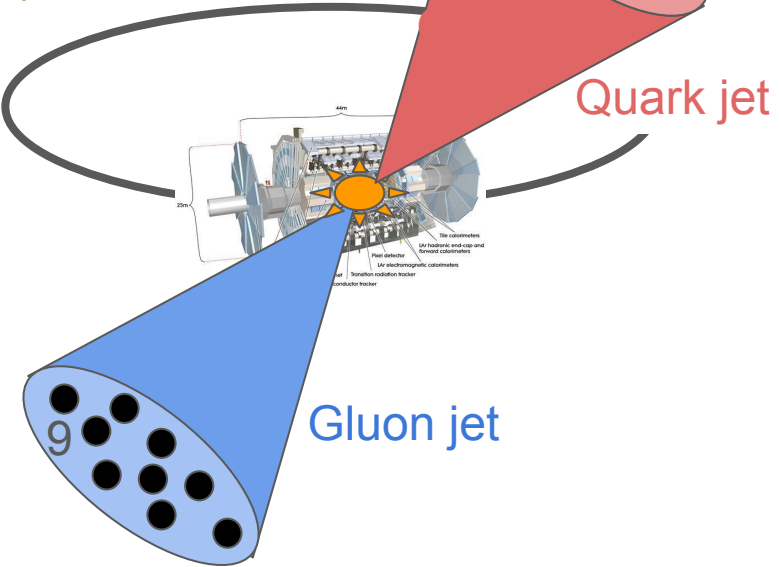


# Introduction

$\sqrt{s} = 900 \text{ GeV}$

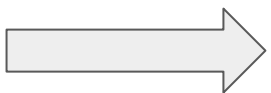


$\sqrt{s} = 13000 \text{ GeV}$

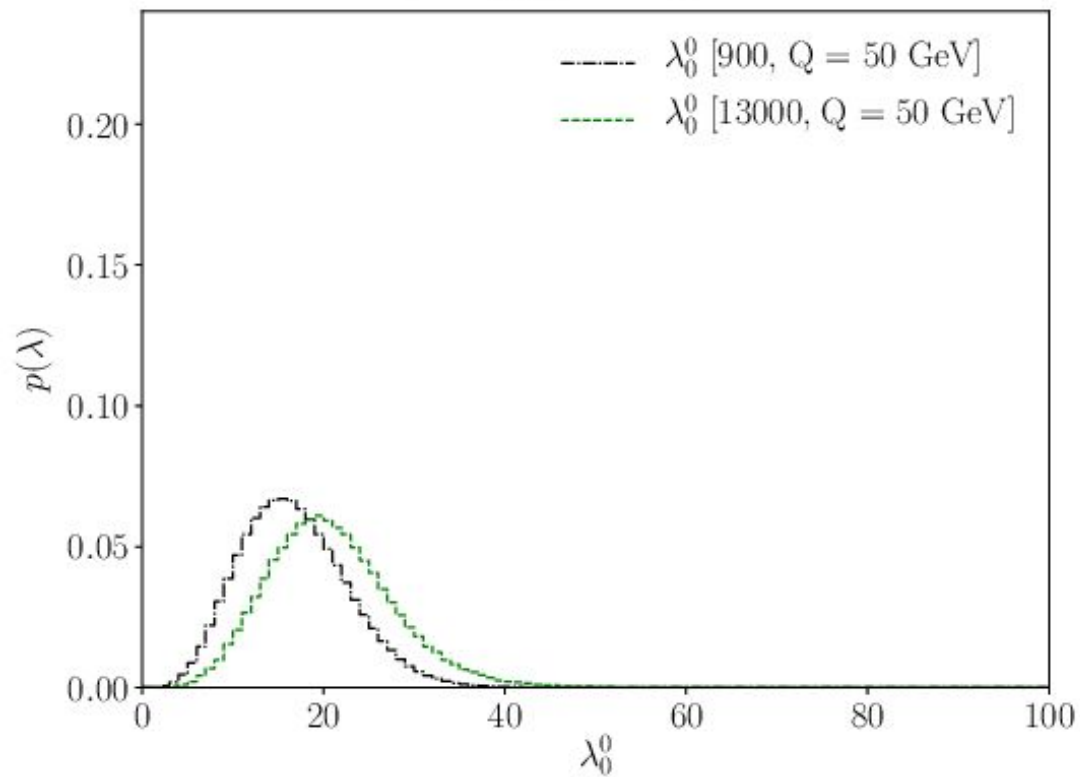


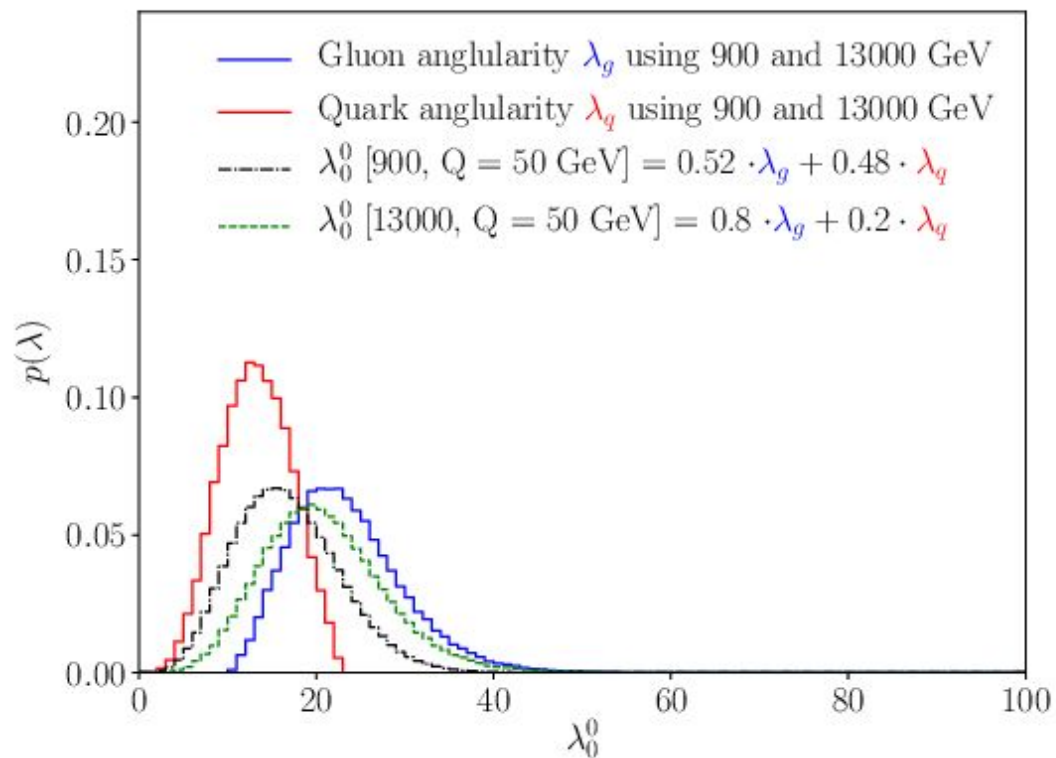
$$\lambda^{900} = f^{900} \lambda_g + (1 - f^{900}) \lambda_q$$

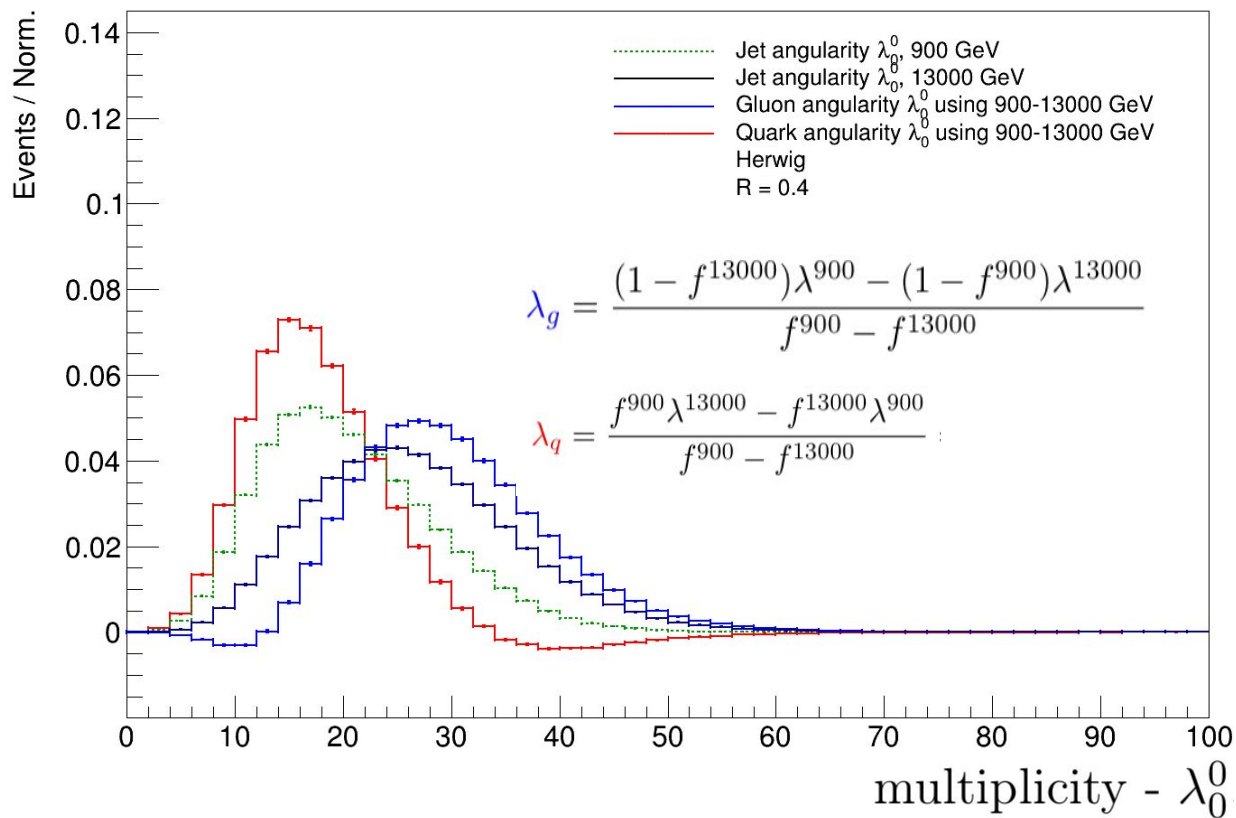
$$\lambda^{13000} = f^{13000} \lambda_g + (1 - f^{13000}) \lambda_q$$



$\lambda_q = \frac{f^{900} \lambda^{13000} - f^{13000} \lambda^{900}}{f^{900} - f^{13000}}$	$\lambda_g = \frac{(1 - f^{13000}) \lambda^{900} - (1 - f^{900}) \lambda^{13000}}{f^{900} - f^{13000}}$
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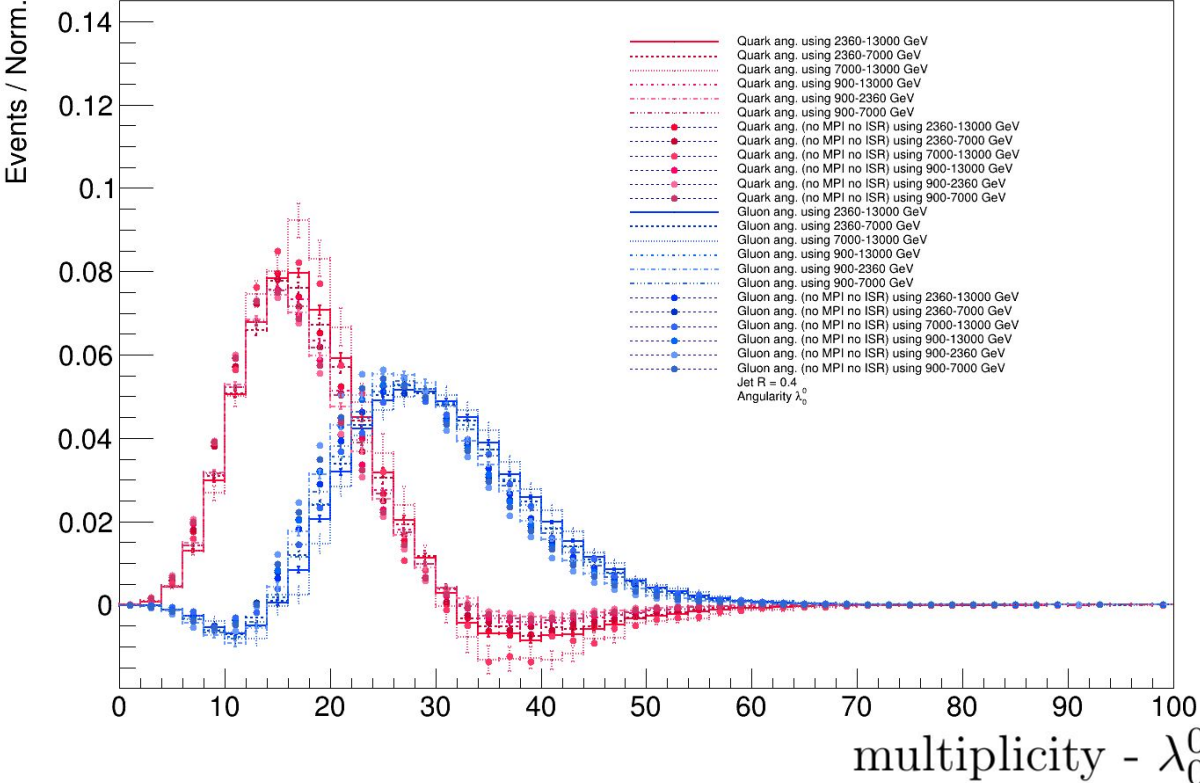
Multiplicity,  $pp \rightarrow 2j$ ,  $R = 0.4$ 

Multiplicity,  $pp \rightarrow 2j$ ,  $R = 0.4$ 



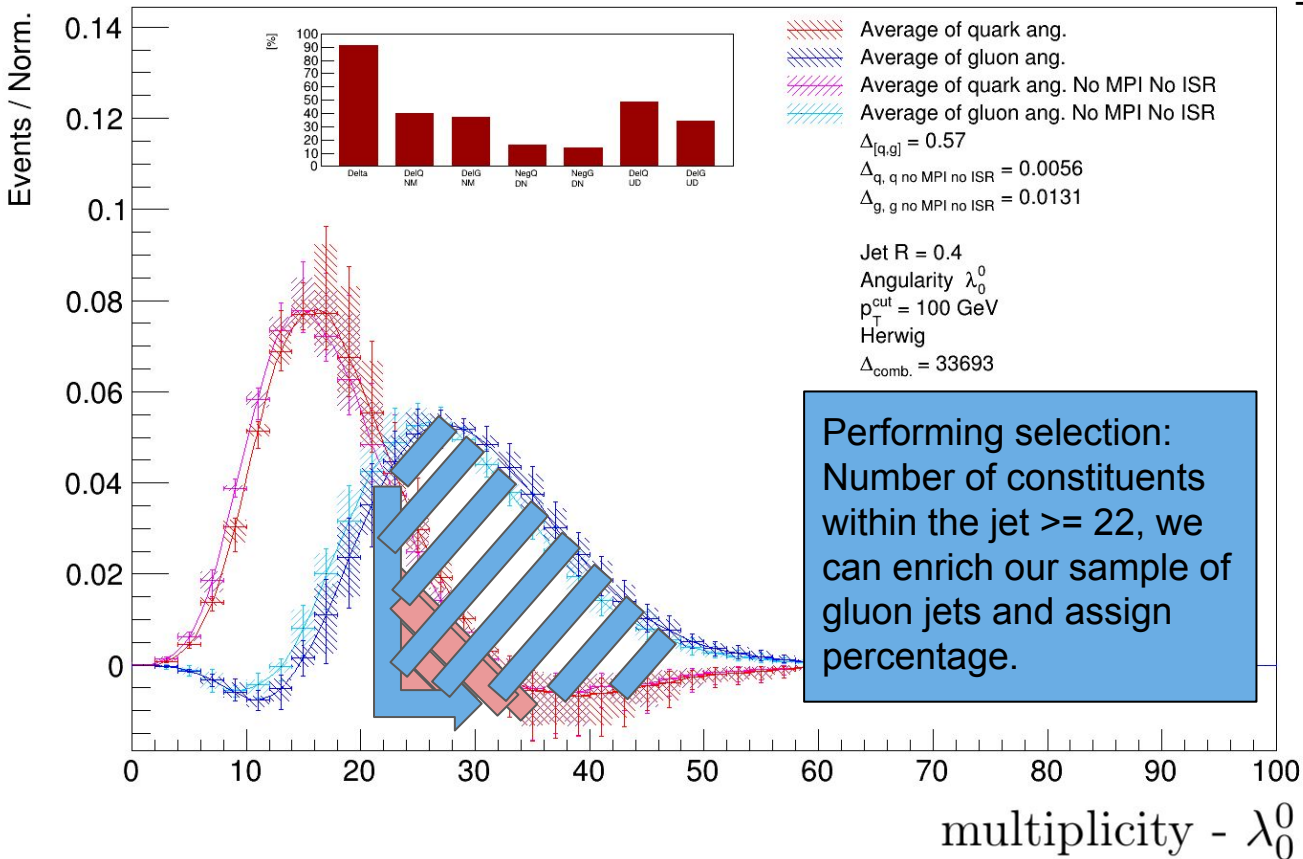
Let's use more 6 energy combinations:

900-2360, 900-7000, 900-13000, 2360-7000, 2360-13000, 7000-13000 GeV

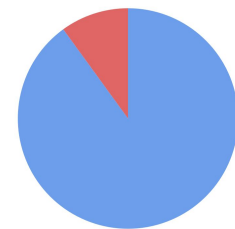


Dotted lines test the robustness to Multi Parton Interactions MPI and Initial State Radiation ISR

Simplified averaged plot over 6 energy combinations: - filled area (energy comb. variation),  
 - ticks - stat. Unc.

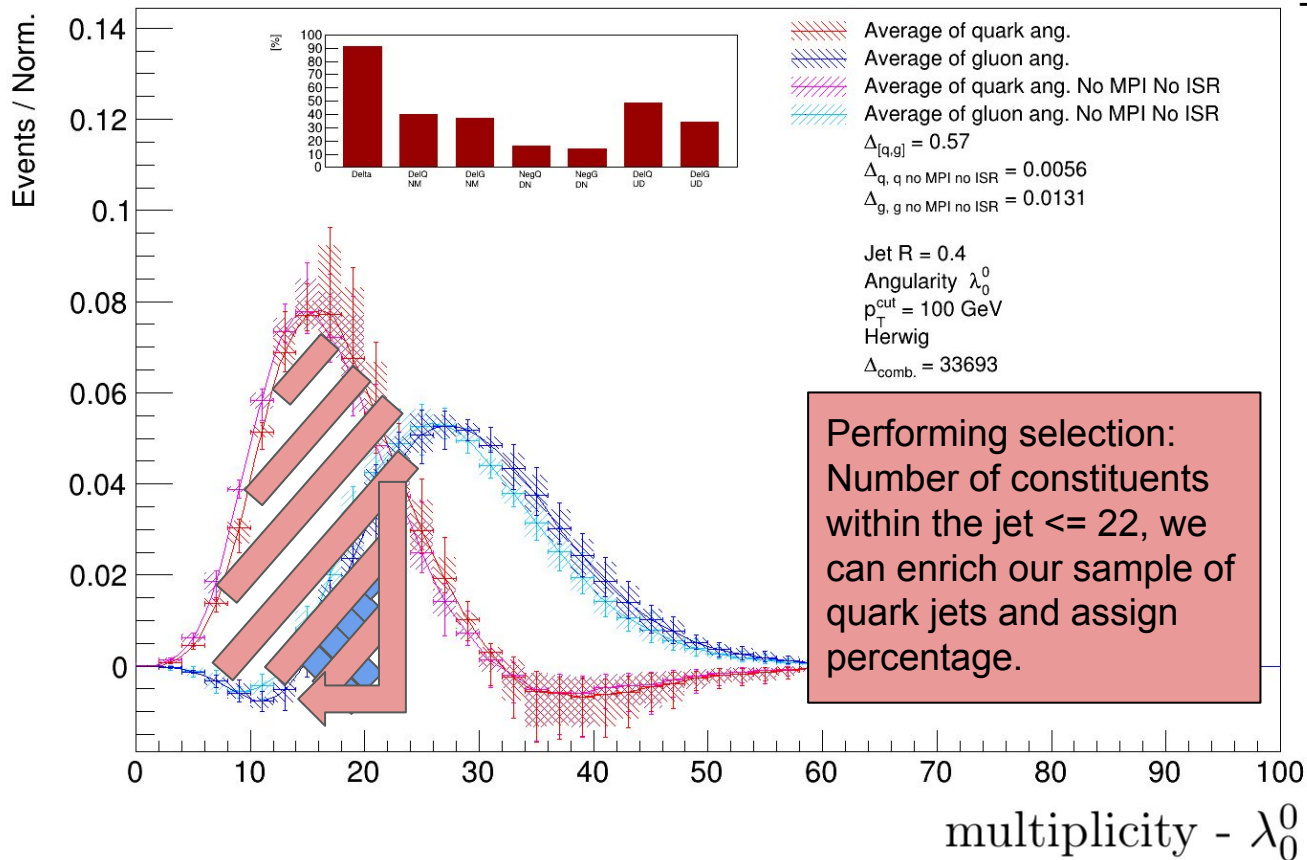


Quark jets

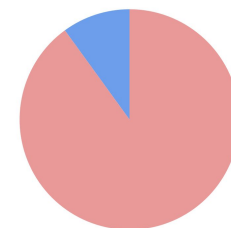


Gluon jets

Simplified averaged plot over 6 energy combinations: - filled area (energy comb. variation),  
 - ticks - stat. Unc.



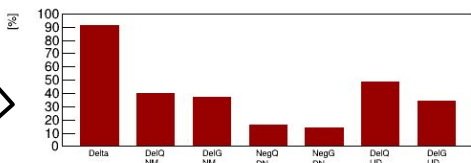
Glue jets



Quark jets

Simplified averaged plot over 6 energy combinations: - filled area (energy comb. variation), - ticks - stat. Unc.

**WARNING: Fake plot!**



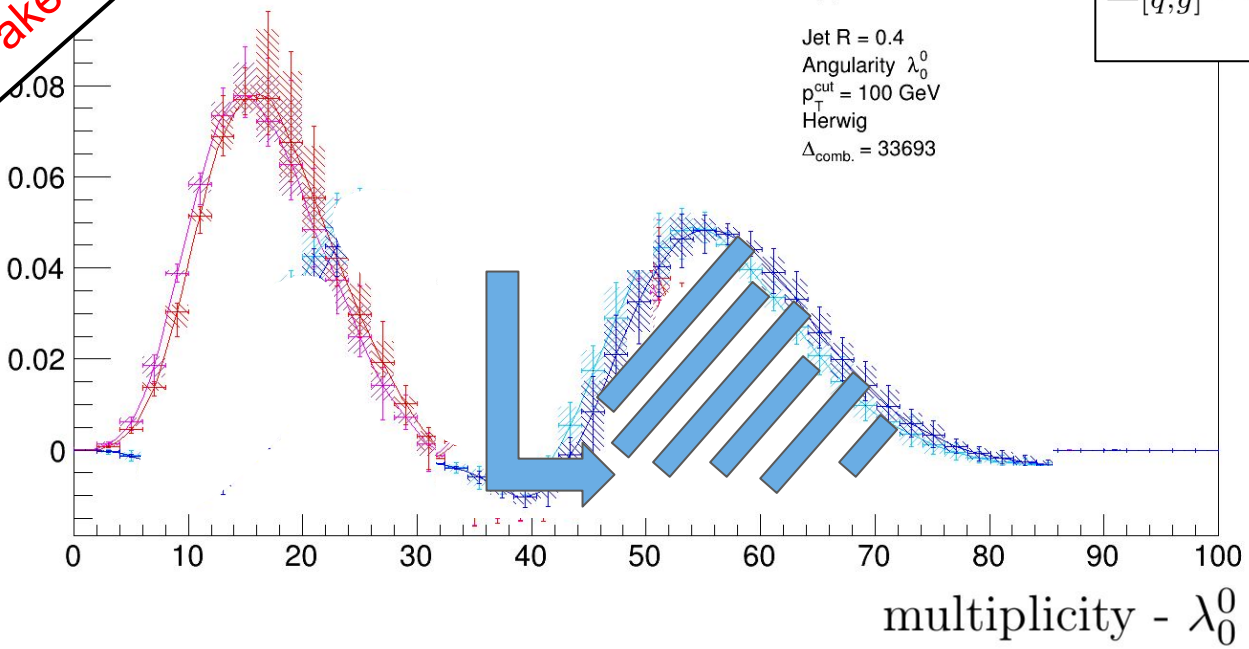
- Average of quark ang.
- Average of gluon ang.
- Average of quark ang. No MPI No ISR
- Average of gluon ang. No MPI No ISR

$\Delta_{[q,g]} = 0.57$   
 $\Delta_{q, q \text{ no MPI no ISR}} = 0.0056$   
 $\Delta_{g, g \text{ no MPI no ISR}} = 0.0131$

Jet R = 0.4  
 Angularity  $\lambda_0^0$   
 $p_T^{\text{cut}} = 100 \text{ GeV}$   
 Herwig  
 $\Delta_{\text{comb.}} = 33693$

**Separation power:**

$$\Delta_{[q,g]} = \frac{1}{2} \sum_{i=1}^N \frac{(\lambda_{q_i} - \lambda_{g_i})^2}{\lambda_{q_i} + \lambda_{g_i}}$$

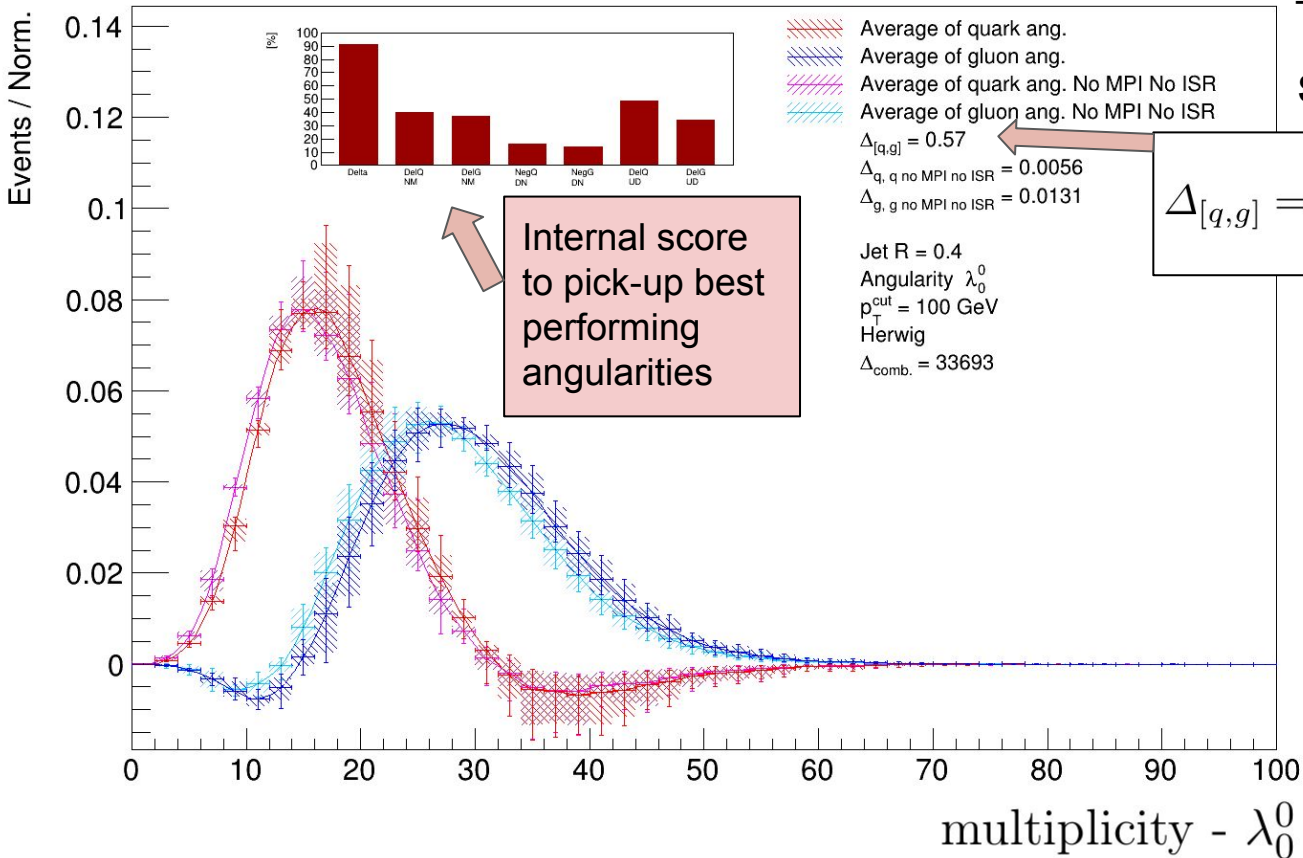


Aim is to find variable with highest separation power





Simplified averaged plot over 6 energy combinations: - filled area (energy comb. variation),  
 - ticks - stat. Unc.



**Separation power:**

$$\Delta_{[q,g]} = \frac{1}{2} \sum_{i=1}^N \frac{(\lambda_{q_i} - \lambda_{g_i})^2}{\lambda_{q_i} + \lambda_{g_i}}$$

Average of quark ang.  
 Average of gluon ang.  
 Average of quark ang. No MPI No ISR  
 Average of gluon ang. No MPI No ISR  
 $\Delta_{[q,g]} = 0.57$   
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 Jet R = 0.4  
 Angularity  $\lambda_0^0$   
 $p_T^{\text{cut}} = 100 \text{ GeV}$   
 Herwig  
 $\Delta_{\text{comb.}} = 33693$

[Gras, Hoeche, Kar, Larkoski, Lönnblad, Plätzer, AS, Skands, Soyez, Thaler, JHEP 1707 (2017) 091]

Cartoon:

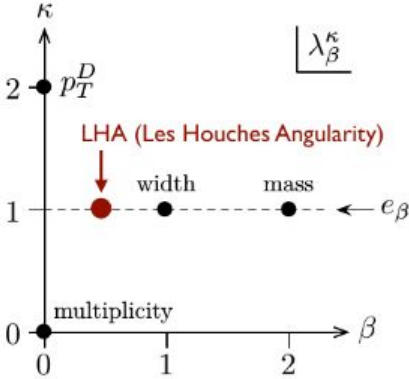


Quark:  $C_F = 4/3$  vs. Gluon:  $C_A = 3$

Probe radiation pattern with e.g. Generalized Angularities

$$\lambda_\beta^\kappa = \sum_{i \in \text{jet}} z_i^\kappa \theta_i^\beta$$

$z_i$ : momentum fraction  
 $\theta_i$ : angle to recoil-free axis  
 $(\lambda_\beta^\kappa)_{\text{quark}} < (\lambda_\beta^\kappa)_{\text{gluon}}$



multiplicity -  $\lambda_0^0$   
 $p_T^D$  -  $\lambda_0^2$ ,  
 LHA -  $\lambda_{0.5}^1$ ,  
 width -  $\lambda_1^1$ , and  
 mass -  $\lambda_2^1$ .

[Larkoski, Salam, Thaler, 13]  
 [Larkoski, Thaler, Waalewijn, 14]

Selection of dijet events:

$N_{\text{jets}} = 2$   
veto neutrinos

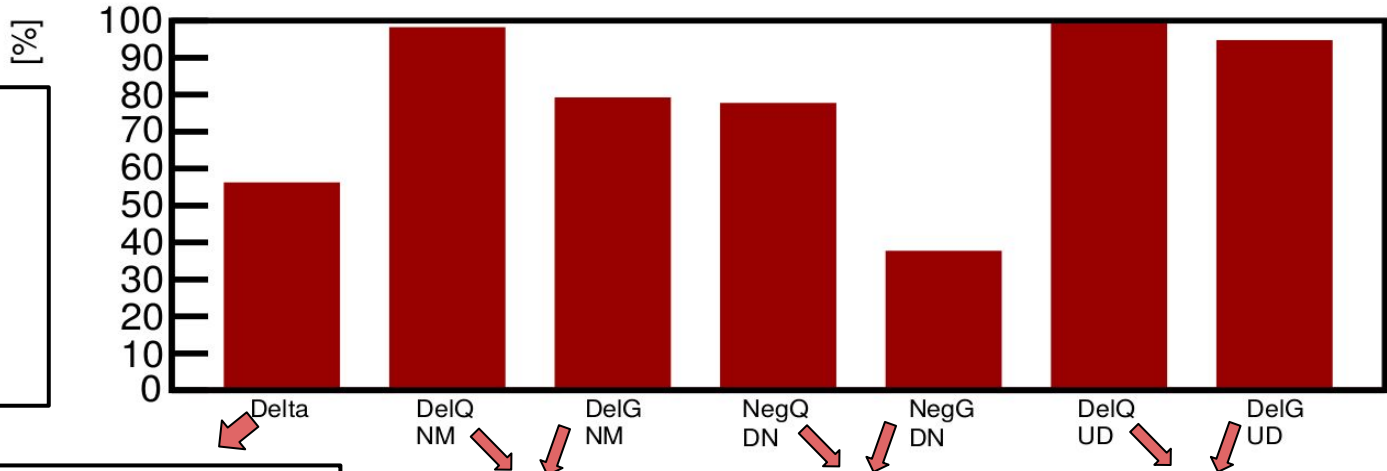
$$p_T \text{ sublead} / p_T \text{ lead} > 0.8$$

We considered all combinations of:

- 5 – angularities  $\lambda_0^0, \lambda_{0.5}^1, \lambda_1^1, \lambda_0^2, \lambda_2^1$
- 2 – using groomed (MMDT) / not groomed jets
- 5 – jet radii  $R = 0.2, 0.4, 0.6, 0.8, 1.0$
- 4 – regions - dijet average  $p_T^{\text{cut}} = 50 \text{ GeV}, 100, 200,$   
and  $400 \text{ GeV}$   $(p_T \text{ lead} + p_T \text{ sublead})/2 > p_T^{\text{cut}}$
- 2 – quark/gluon
- 2 – MPI and ISR switched on/off
- 6 – energy combinations:  $900\text{--}2360, 900\text{--}7000, 900\text{--}13000, 2360\text{--}7000, 2360\text{--}13000, 7000\text{--}13000 \text{ GeV}$
- 2 – event generators HERWIG and PYTHIA

# Results:

Each column represents percentile of given feature of all our studied plots.



In **separation between Q/G jets** about 55 % of variations have lower performance

Separation power between noMPI variations - **robustness to noMPI no ISF**

Negativity (percentage of negative area to whole area) **negative bins**

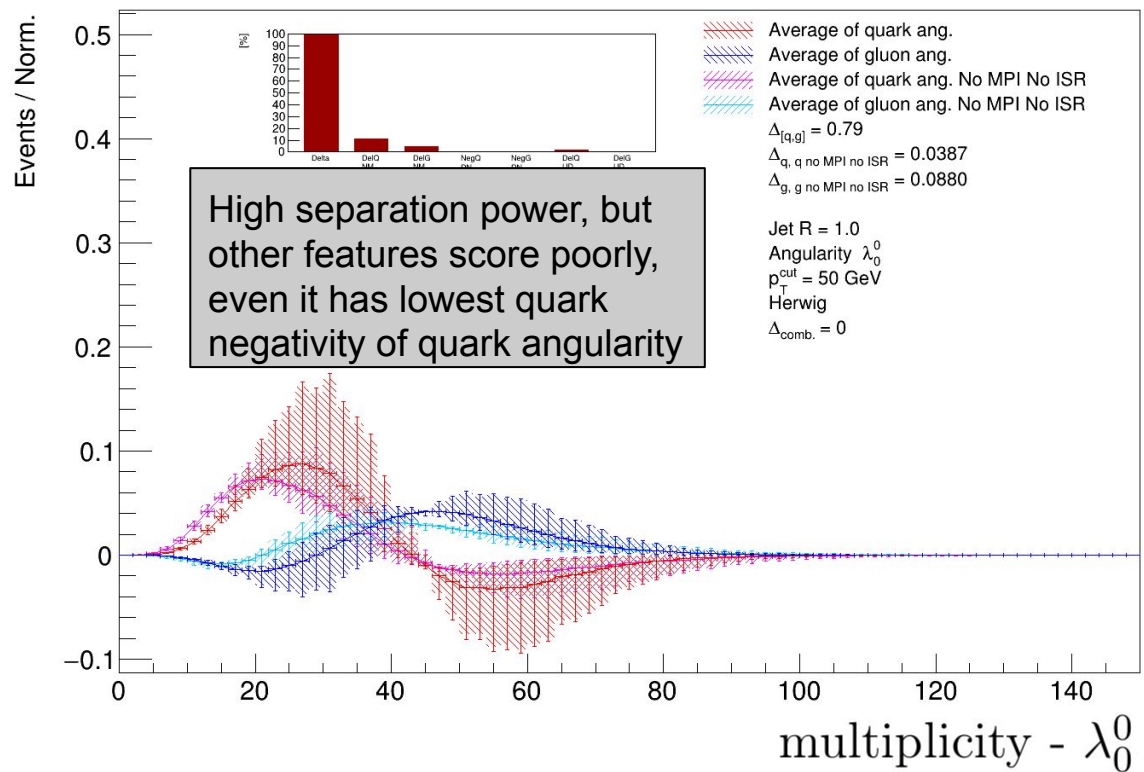
Separation power to UP and DOWN energy combination variations **robustness to different energy combination used**

Combining all columns gives us our internal score :

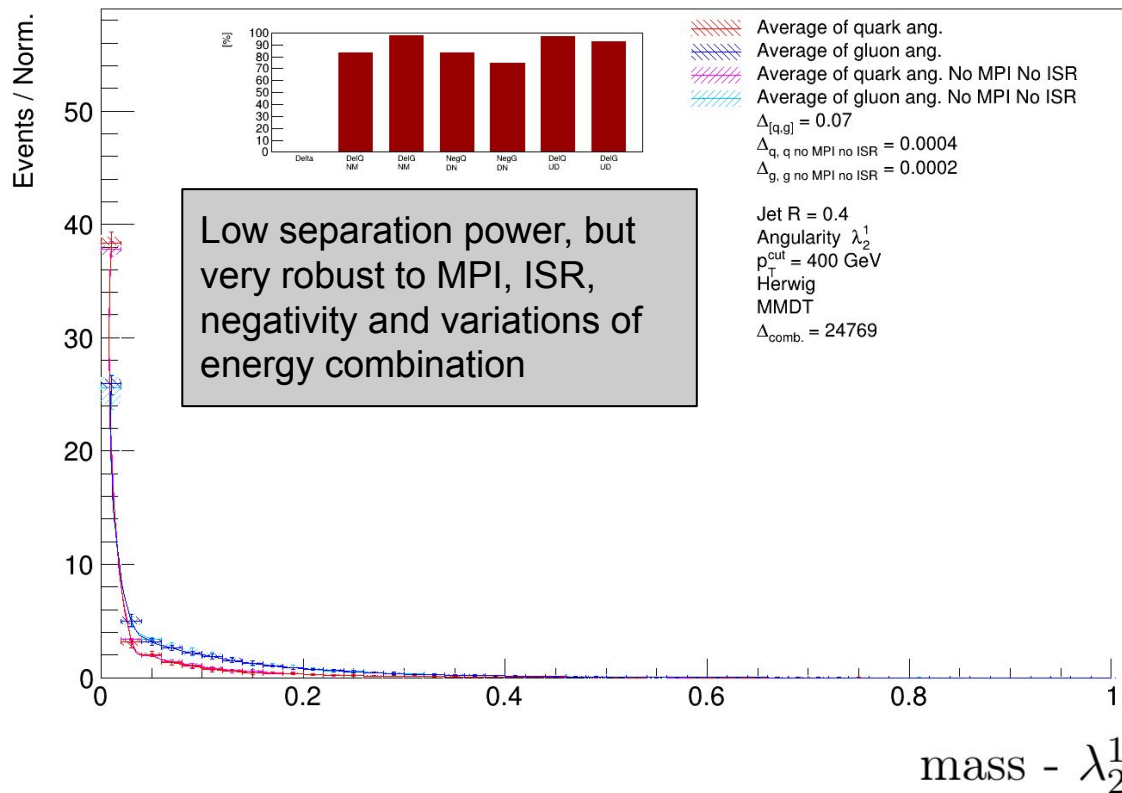
$$\Delta_{\text{comb}} = 1000 \cdot \ln \left[ 1 + (\text{Delta})^3 \cdot (\text{DelQ NM}) \cdot (\text{DelG NM}) \cdot (\text{NegQ DN}) \cdot (\text{NegG DN}) \cdot (\text{DelQ UD}) \cdot (\text{DelG UD}) \right]$$

**min-max 0-41447**

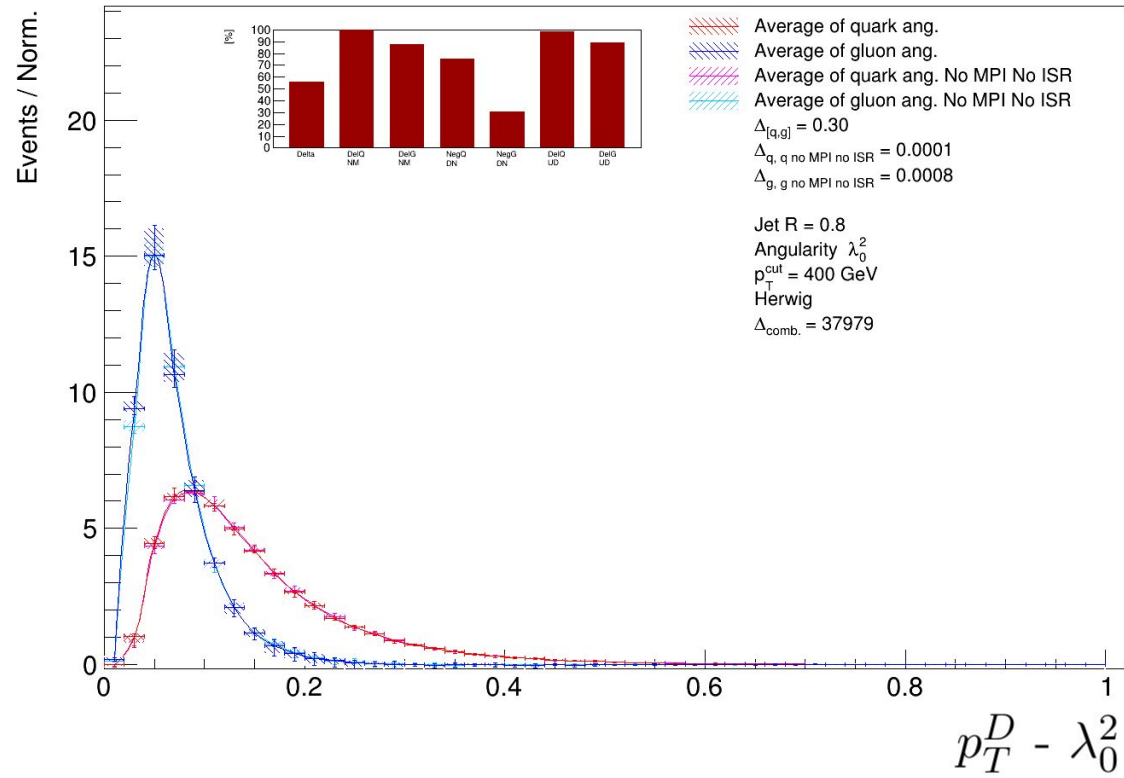
Why we looking into other features then separation power (bad examples):



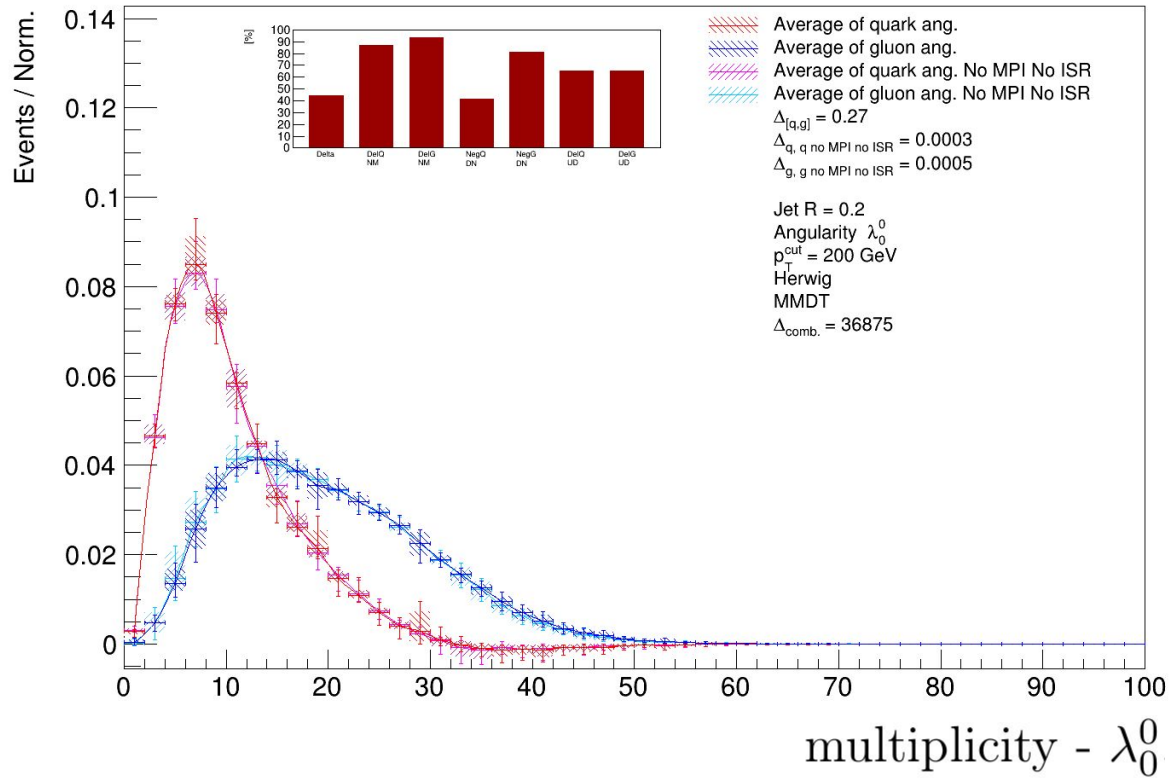
Why we looking into other features then separation power (bad examples):



Best performing angularities:  $p_T^D$

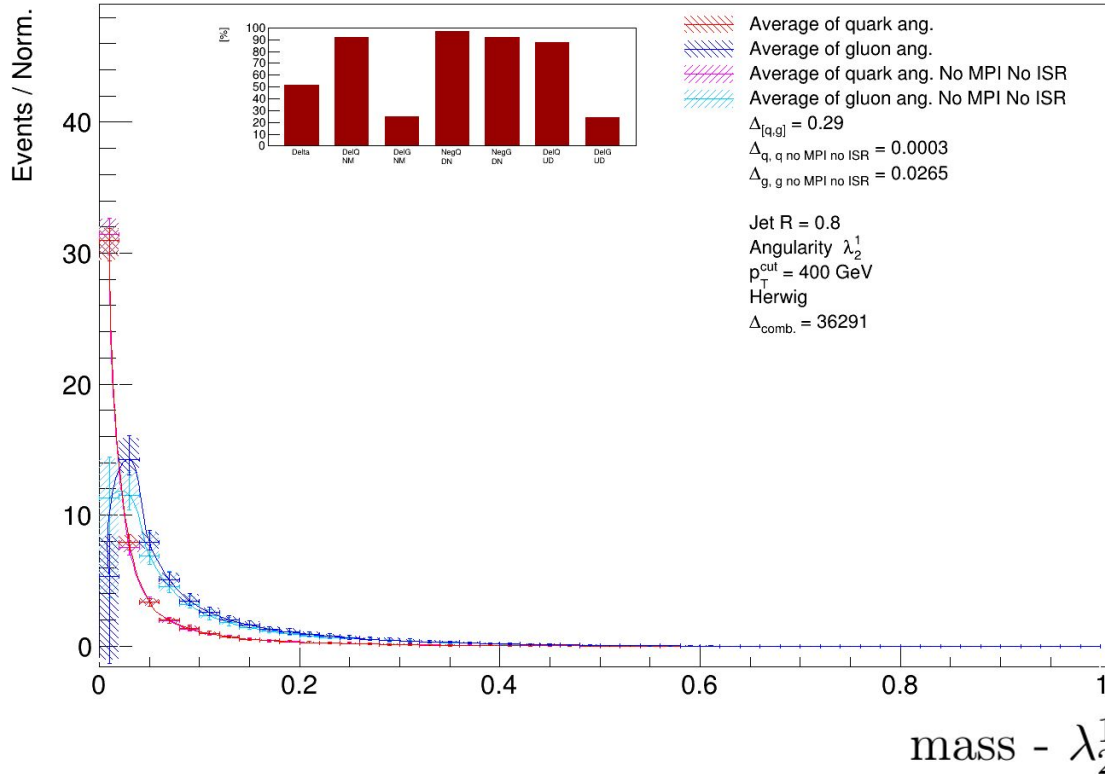


### Best performing angularities: Multiplicity

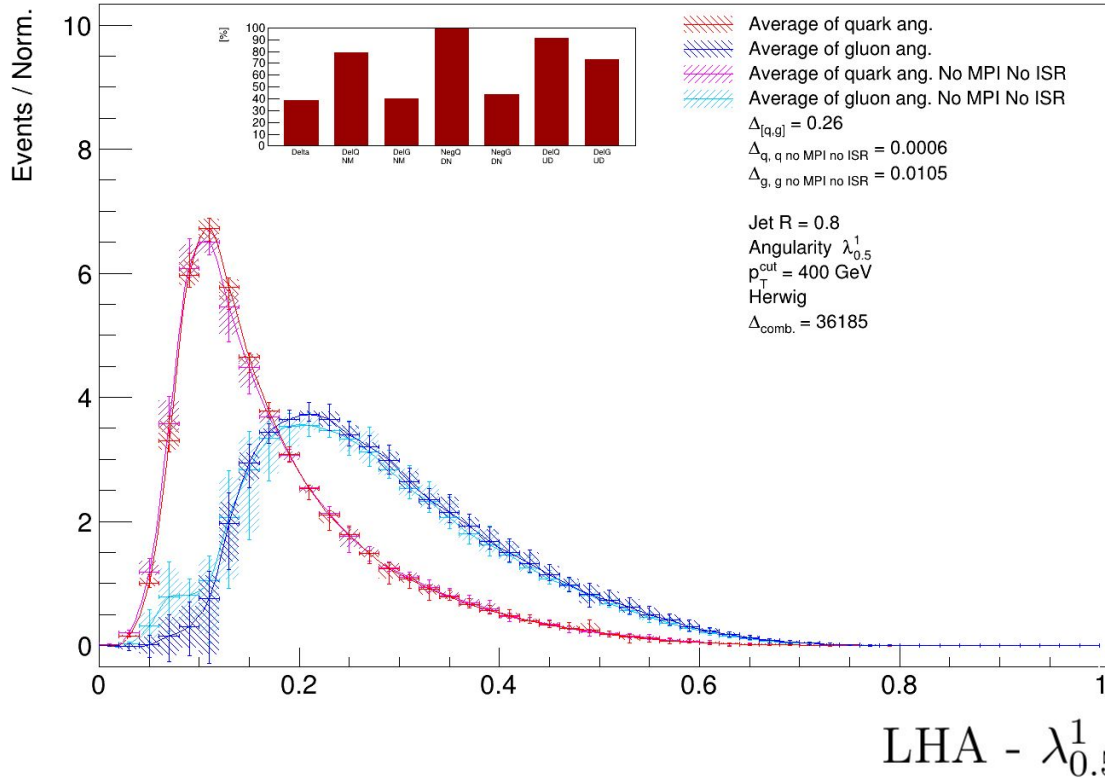




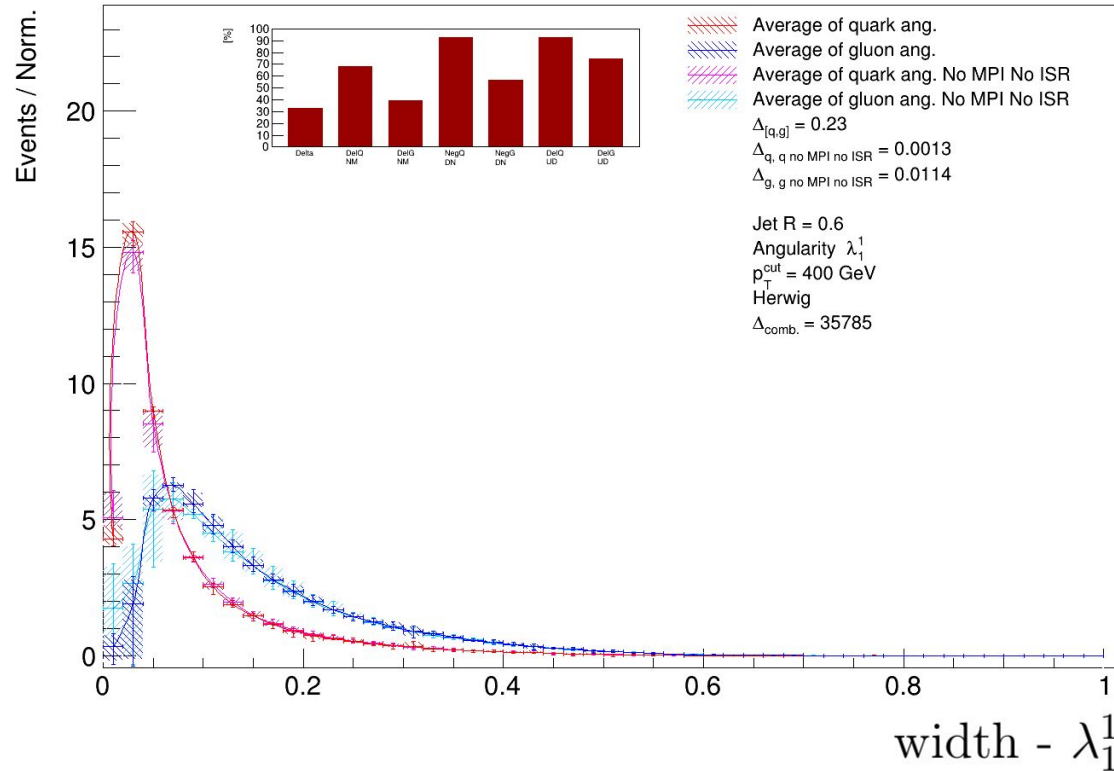
### Best performing angularities: Mass



### Best performing angularities: LHA



## Best performing angularities: Width



## 4. Conclusion

- Main idea it that properties of jets of a given flavour and transverse momentum, are almost entirely independent of the jet's production mechanism.

Thus, the energy-dependence can be used to extract the flavour-dependent properties on a statistical basis.

- We proposed selection of best performing angularities
- More details:

<https://doi.org/10.1140/epjc/s10052-023-12363-4>

- Plans:
  - Multidim angularities (2D, 3D) with machine learning approach to enhance separation power
  - Perform measurement at LHC

Eur. Phys. J. C (2024) 84:28  
<https://doi.org/10.1140/epjc/s10052-023-12363-4>

**THE EUROPEAN  
PHYSICAL JOURNAL C**

Regular Article - Theoretical Physics

## Novel approach to measure quark/gluon jets at the LHC

Petr Baroň<sup>1,a</sup>, Michael H. Seymour<sup>2,b</sup>, Andrzej Siódmok<sup>3,c</sup>

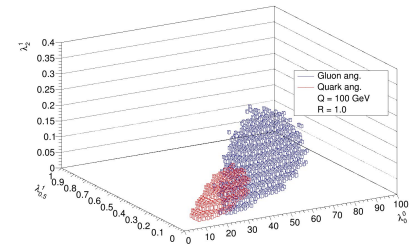
<sup>1</sup> Institute of Nuclear Physics, Polish Academy of Sciences, ul. Radzikowskiego 152, 31-342 Kraków, Poland  
<sup>2</sup> Department of Physics and Astronomy, University of Manchester, Manchester M13 9PL, UK  
<sup>3</sup> Jagiellonian University, ul. prof. Stanisława Łojasiewicza 11, 30-348 Kraków, Poland

Received: 1 September 2023 / Accepted: 18 December 2023  
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**Abstract** In this paper, we present a new proposal on how to measure quark/gluon jet properties at the LHC. The measurement strategy takes advantage of the fact that the LHC has collected data at different energies. Measurements at two or more energies can be combined to yield distributions of any jet property separated into quark and gluon jet samples on a statistical basis, without the need for an independent event-by-event tag. We illustrate our method with a variety of different angularity observables, and discuss how to narrow down the search for the most useful observables.

will also propose how to calibrate that observable by independently tagging quark and gluon jet samples. In some studies, this has been done by calibrating against Monte Carlo samples in which the “truth” flavour of the jet is known. However, one might worry about whether event generators make sufficiently reliable predictions of these flavour-dependent properties [21–23] and, indeed, this is something one would like to test against the data. In other studies, another method is used to tag the jet flavour, for example the hard process dependence [24, 25], and used to calibrate the measurement of the proposed observable. Here, one would worry that the



- Main idea it that properties of jets of a given flavour and transverse momentum, are almost entirely independent of the jet's production mechanism.

Thus, the energy-dependence can be used to extract the flavour-dependent properties on a statistical basis.

- We proposed selection of best performing angularities
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**THE EUROPEAN  
PHYSICAL JOURNAL C**

Regular Article - Theoretical Physics

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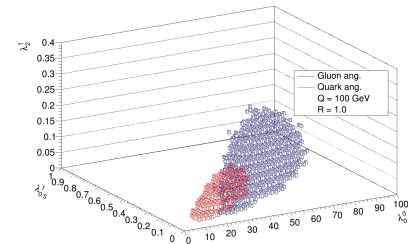
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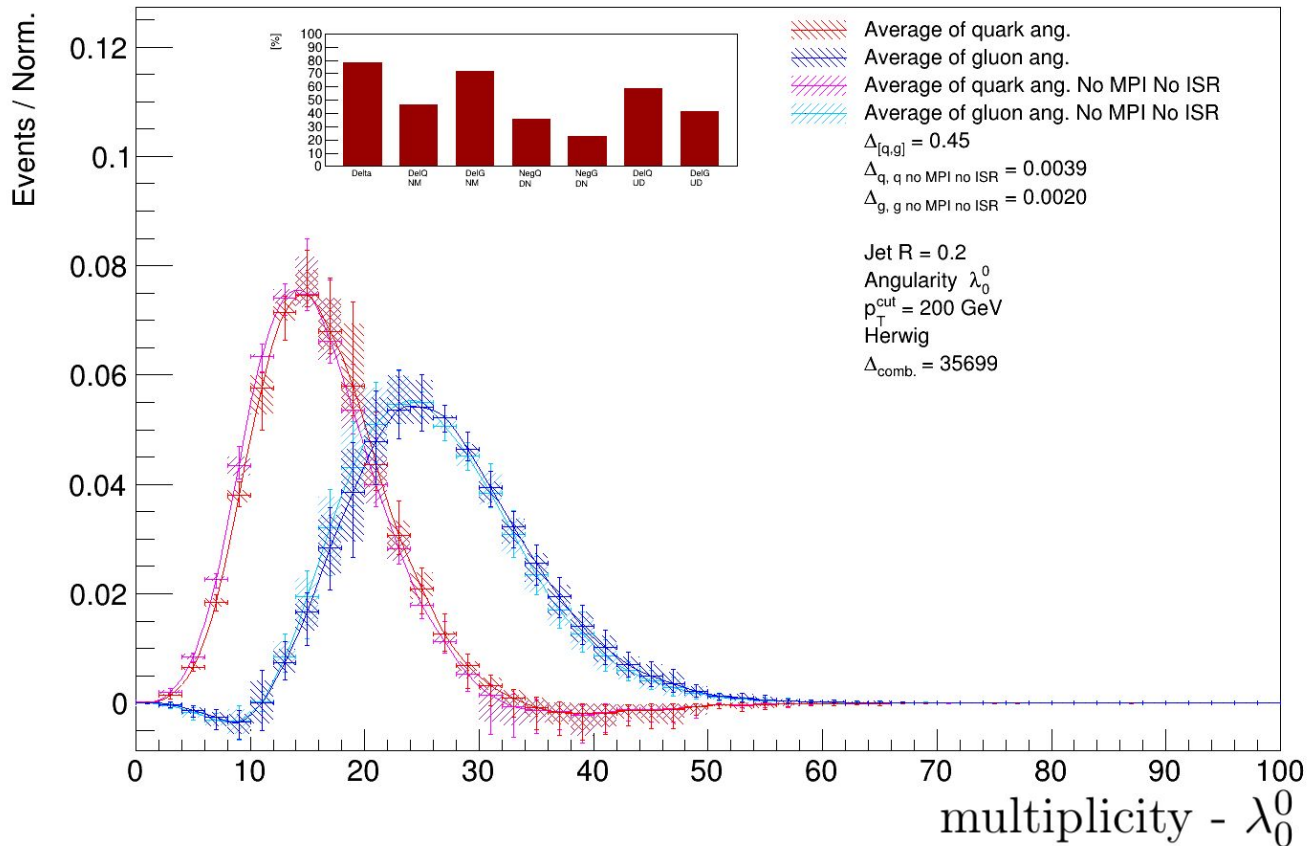
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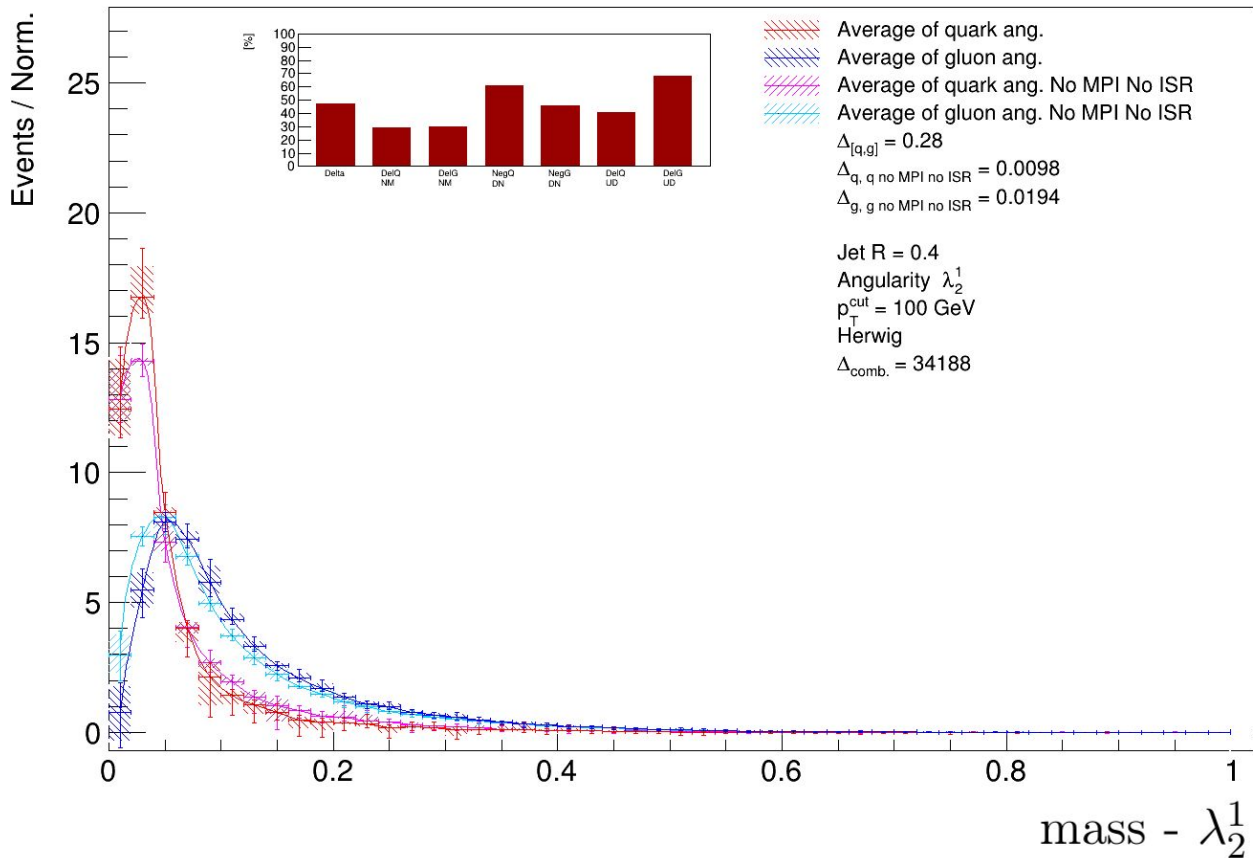


Thank you!

### Wild cards (chosen by "eye"): multiplicity

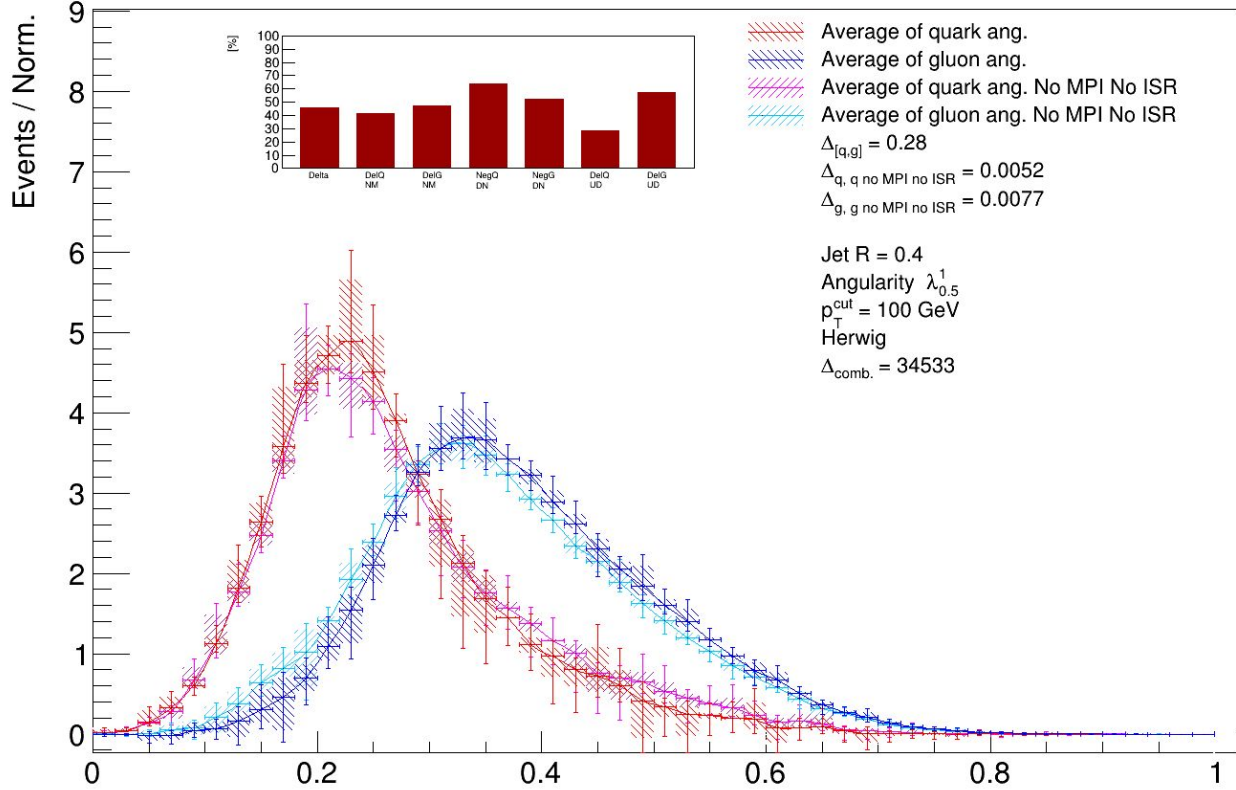


### Wild cards (chosen by "eye"): mass



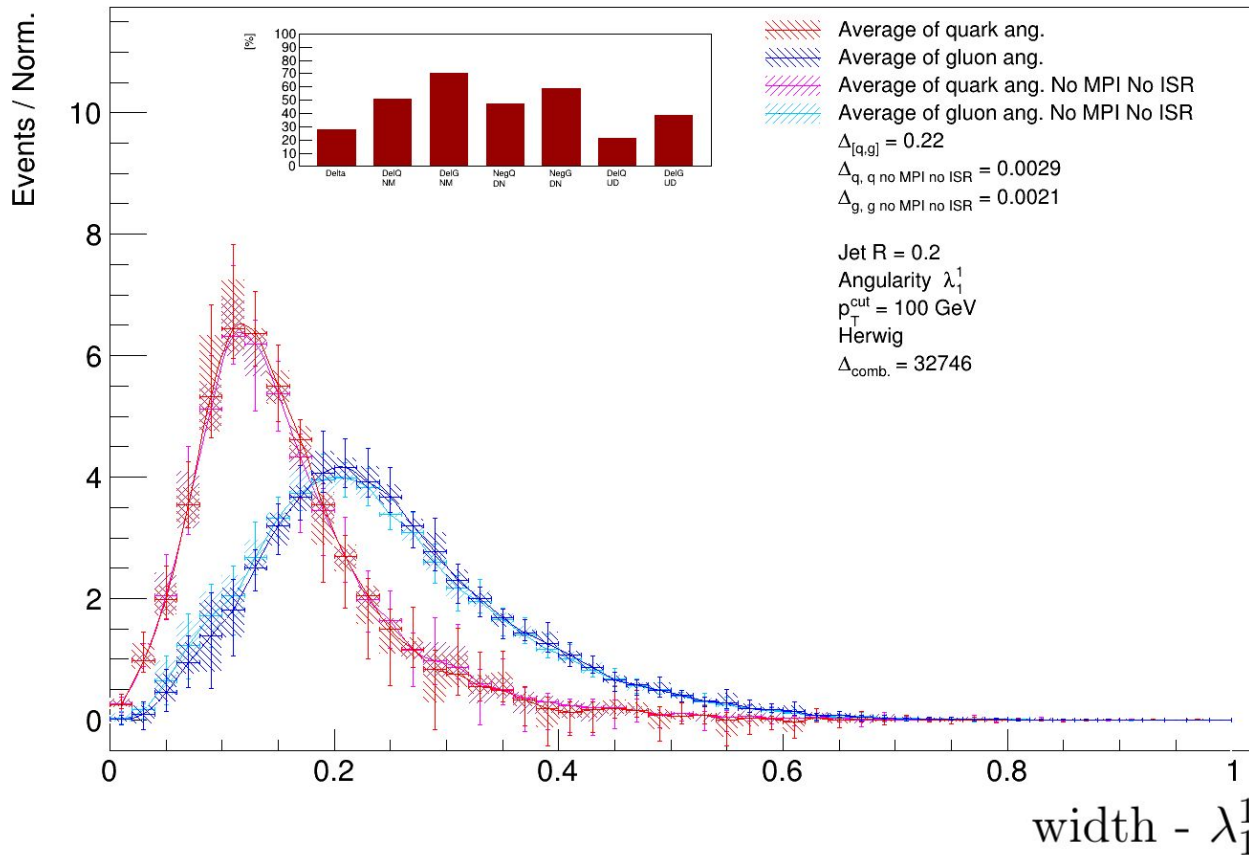


### Wild cards (chosen by "eye"): LHA

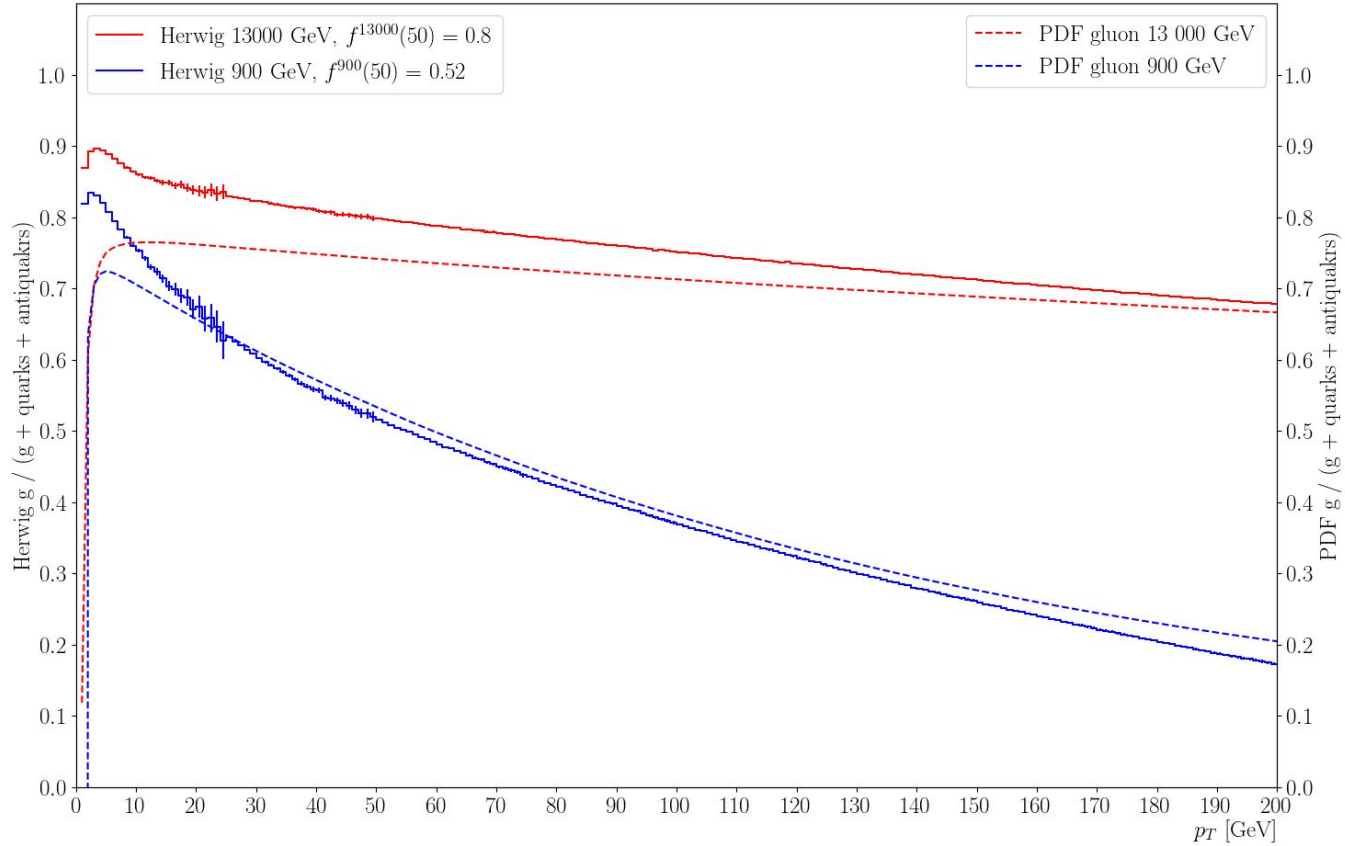


LHA -  $\lambda_{0.5}^1$

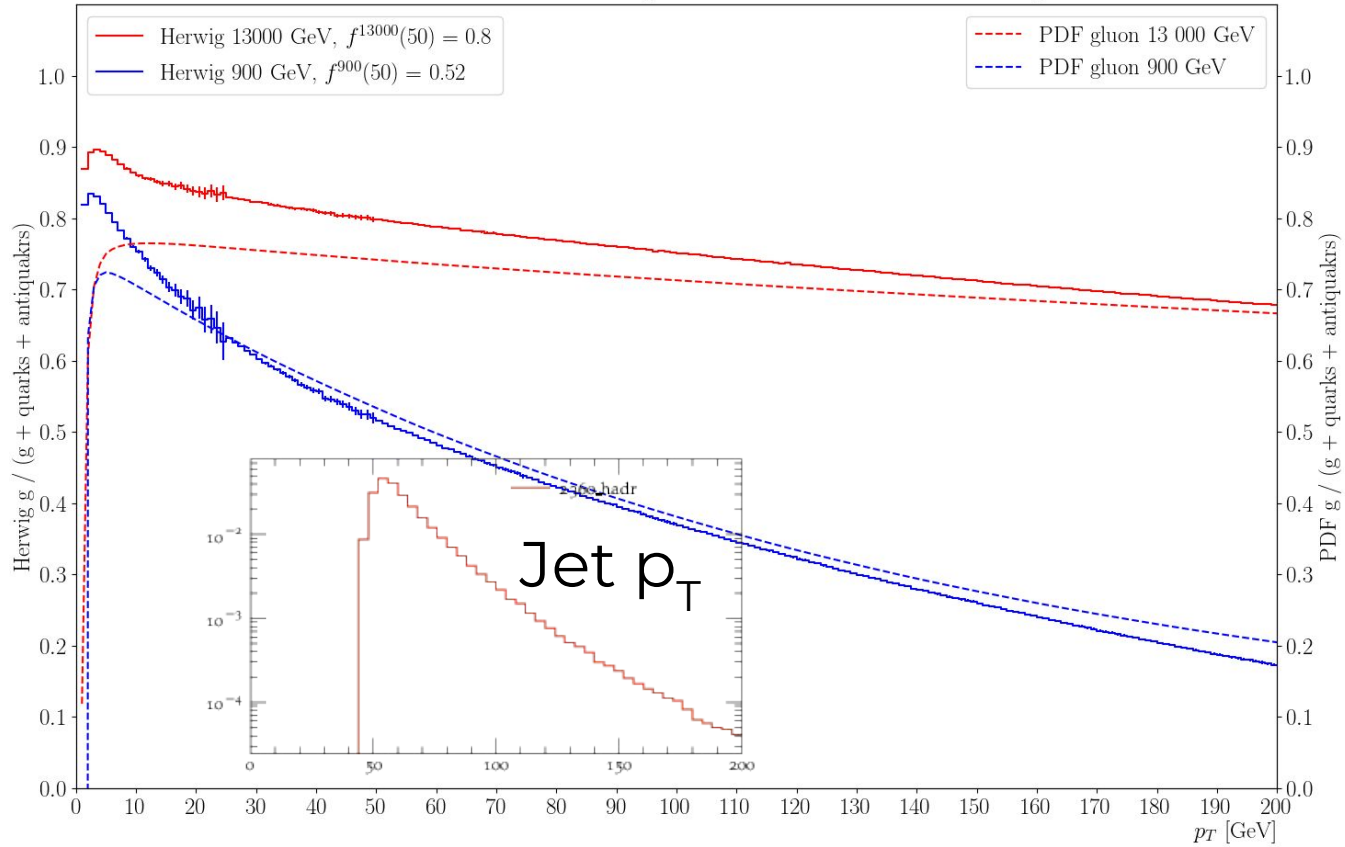
### Wild cards (chosen by "eye"): width



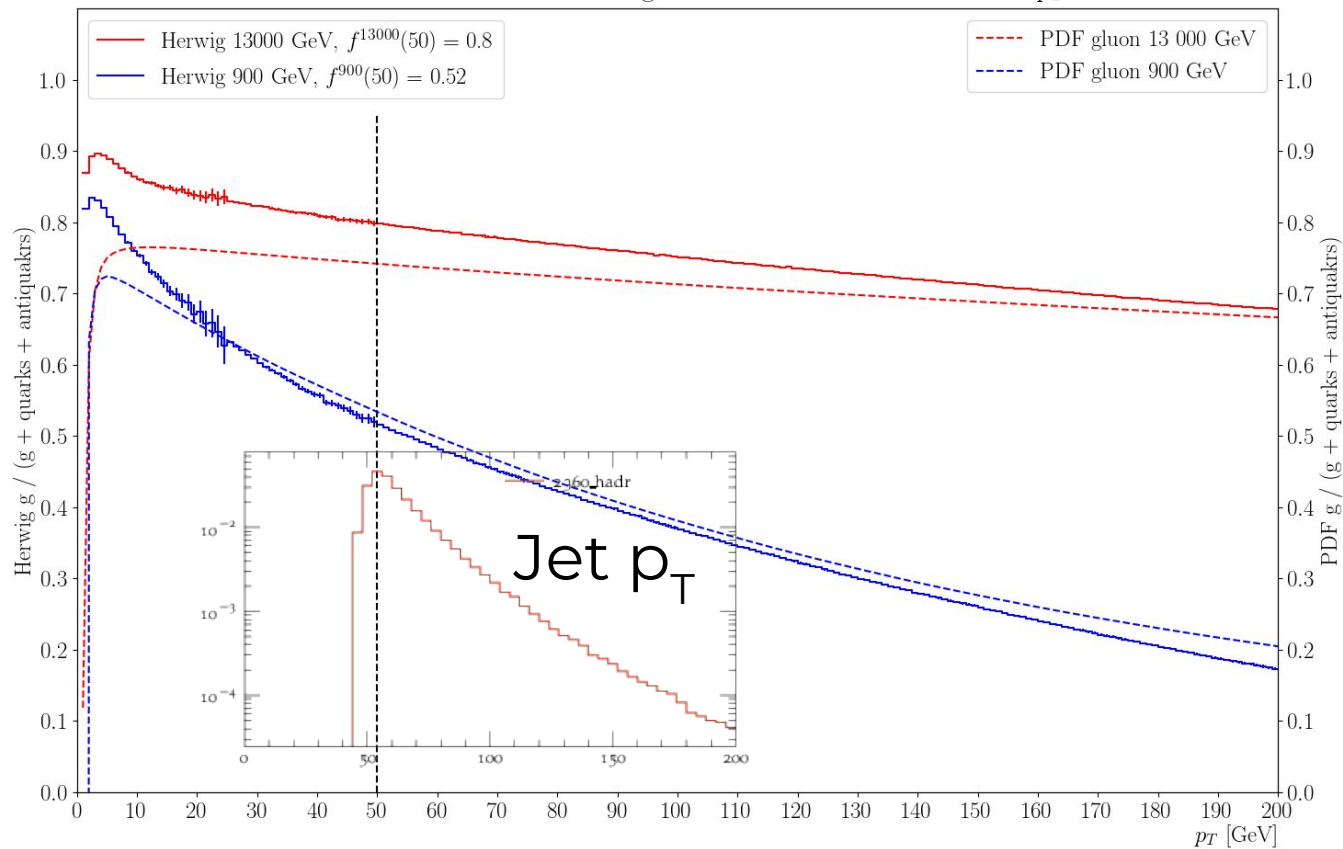
Gluon Fraction PDF and Herwig MHT2014nlo68cl as a function of  $p_T$



Gluon Fraction PDF and Herwig MHT2014nlo68cl as a function of  $p_T$

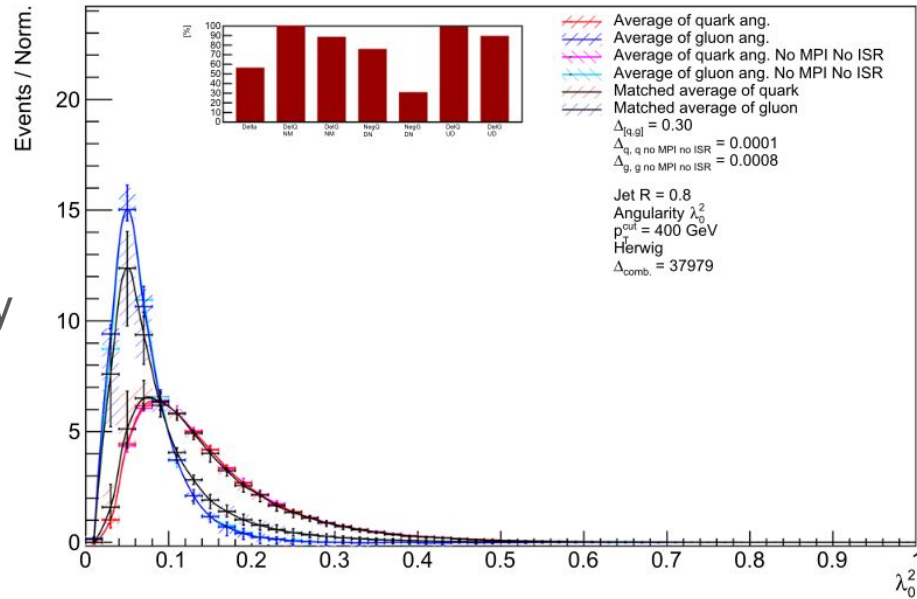


Gluon Fraction PDF and Herwig MHT2014nlo68cl as a function of  $p_T$



Matching to the truth parton level by minimal dR

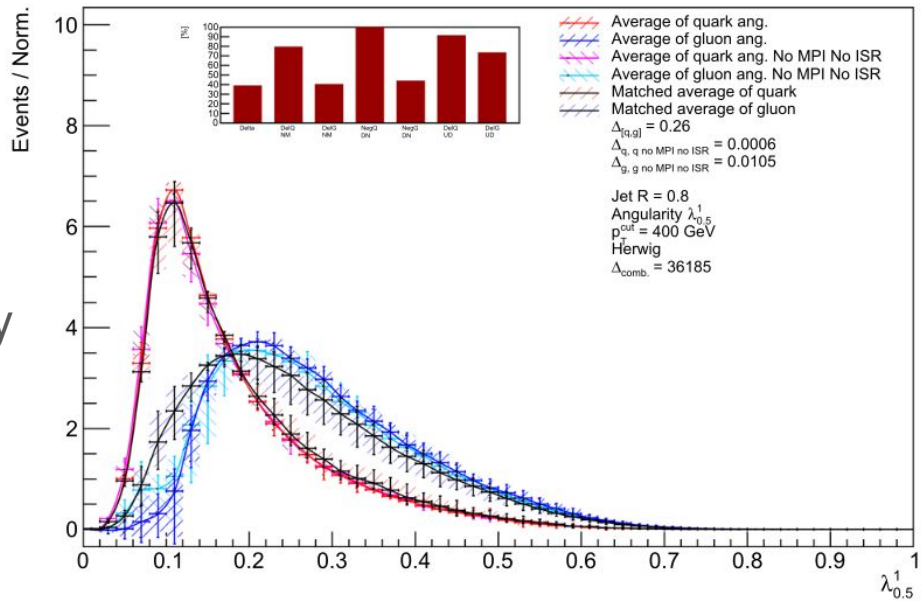
- Black lines



**Fig. 19** Quark and gluon averaged angularities  $\lambda_0^2$ ,  $R = 0.8$  with highest score  $\Delta_{\text{comb}} = 37979$ . Using HERWIG event generator, with  $p_T^{\text{cut}} = 400 \text{ GeV}$ , using the average of 6 energy combinations 900–2360, 900–7000, 900–13000, 2360–7000, 2360–13000, 7000–13000 GeV

Matching to the truth parton level by minimal dR

- Black lines



**Fig. 22** Quark and gluon averaged angularities MMDT  $\lambda_{0.5}^1$ ,  $R = 0.8$  with score  $\Delta_{\text{comb}} = 36185$ . Using HERWIG event generator, with  $p_T^{\text{cut}} = 400$  GeV, using the average of 6 energy combinations 900–2360, 900–7000, 900–13000, 2360–7000, 2360–13000, 7000–13000 GeV

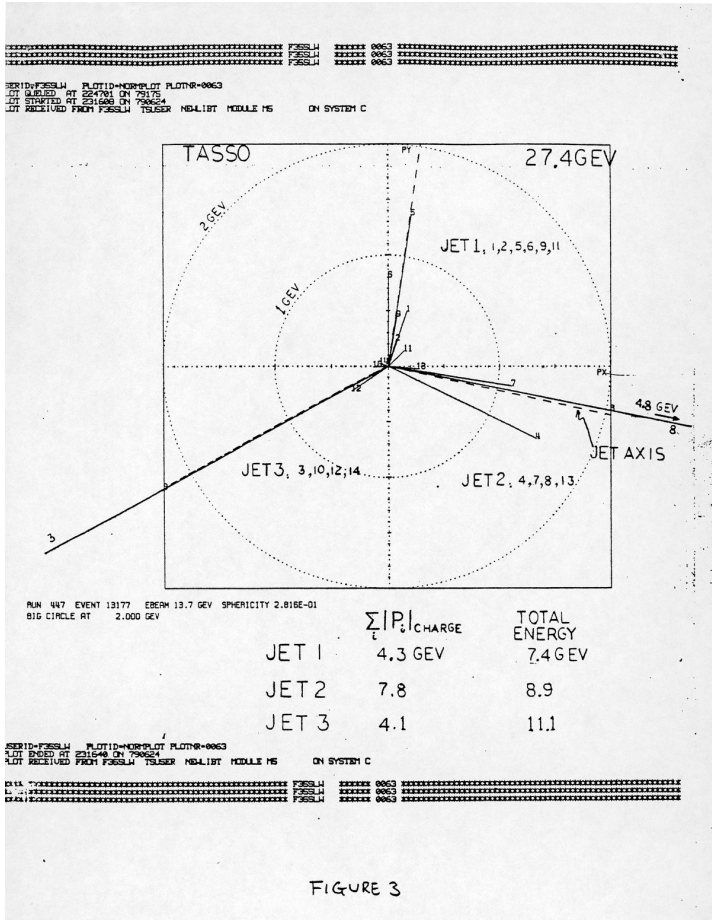
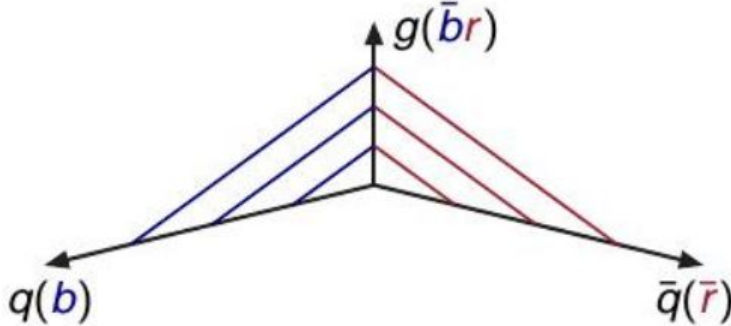


FIGURE 3

This collision event recorded in **1979**, provided the first evidence of the gluon.

Recorded as event 13177 of run 447 of the TASSO experiment at the Deutsches Elektronen-Synchrotron (DESY), the graphic shows three jets of particles produced in an electron-positron collision.





## Quark - Gluon Separation in Three Jet Events #1

Hans Peter Nilles (SLAC), K.H. Strömgren (SLAC) (Aug 1, 1980)

Published in: *Phys.Rev.D* 23 (1981) 1944

 pdf  links  DOI  cite

 32 citations

## A Monte Carlo Program for Quark and Gluon Jet Generation #2

Torbjorn Sjostrand (Lund U., Dept. Theor. Phys.) (Apr 1, 1980)

 pdf  cite

 1 citation

## Quark and gluon jet separation: Conventional and neural network methods #2

Z. Fodor (Eotvos U.) (Jul, 1991)

Published in: *Conf.Proc.C* 910725V1 (1991) 438 · Contribution to: [Joint International Lepton Photon Symposium at High Energies \(15th\) and European Physical Society Conference on High-energy Physics](#), 438



## Quark versus Gluon Jet Tagging Using Charged Particle Multiplicity with the ATLAS Detector #7

ATLAS Collaboration (Apr 11, 2017)

BSM searches: often signature for a BSM signals: many quark, backgrounds: QCD gluons

- 8-jet Gluino event:  $pp \rightarrow \tilde{g}\tilde{g}$  and each  $\tilde{g}$  decays to 4 quarks:

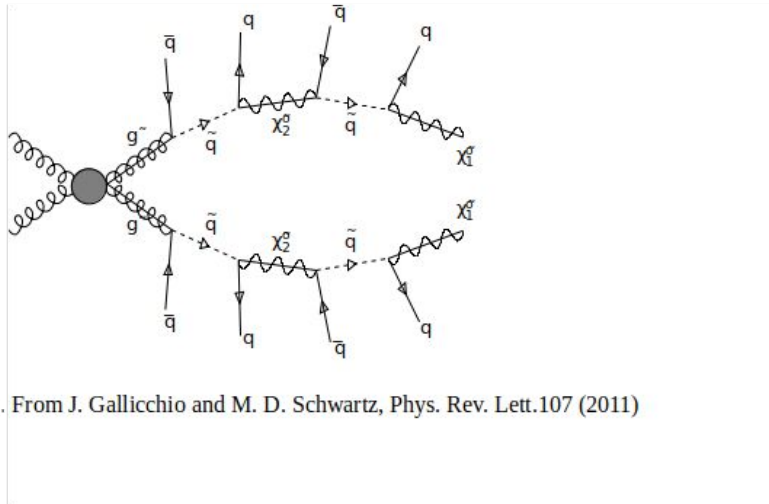


Fig. From J. Gallicchio and M. D. Schwartz, Phys. Rev. Lett.107 (2011)

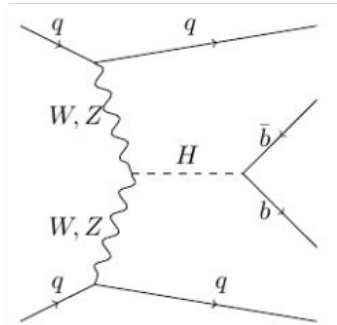
- Higgs  $H^+ \rightarrow c\bar{s}$  (for charged Higgs mass between  $\tau$  and  $t$  mass)
- Measure  $Z'$  coupling to hadrons (or find a leptophobic  $Z'/W'$ )

Interesting standard model physics also tends to be quark-heavy

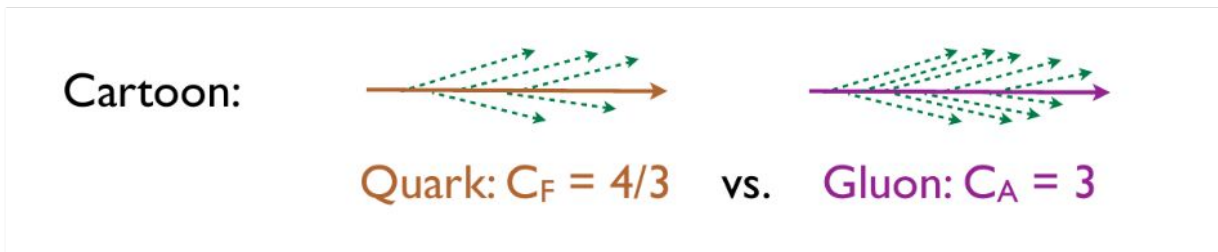
Examples:

- $W$ 's decaying hadronically (no b's!):  $W^+ \rightarrow u\bar{d}$  or  $c\bar{s}$
- Tops ( $t\bar{t} \rightarrow b\bar{b} + 0, 2, \text{ or } 4$  light quarks)
- Vector Boson Scattering/Fusion (forward 'tag' jets are quarks)

**QCD background:** mainly composed by gluons  
**Signal:** mainly composed by quarks

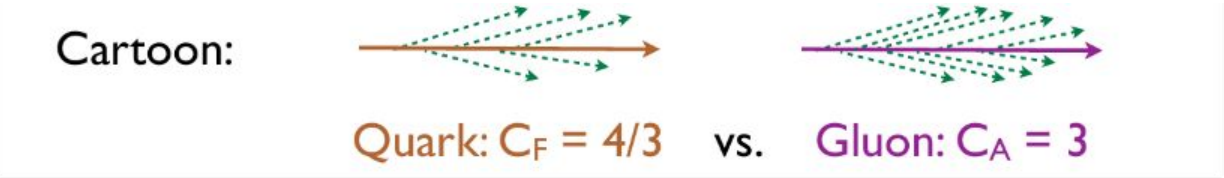


Gluon has a greater effective color charge (squared) than quark:

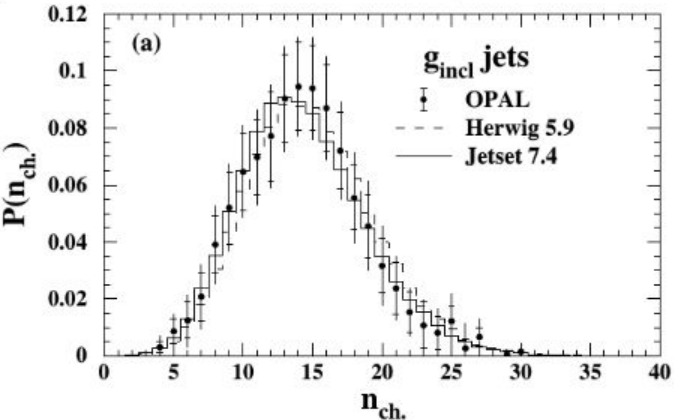
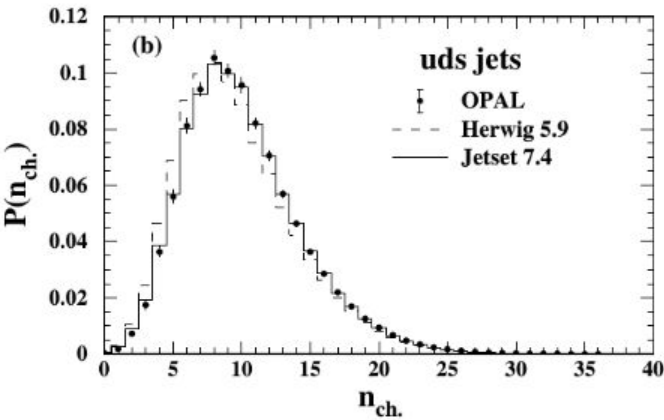


Expectation:

- Gluon will radiate more
- Gluon will radiate wider
- Multiple radiation  $\rightarrow$  effect will exponentiate



Gluon will radiate more, gluon will radiate wider  $\frac{\langle N_g \rangle}{\langle N_q \rangle} = \frac{C_A}{C_F} = 3/(4/3) = 9/4 = 2.25$



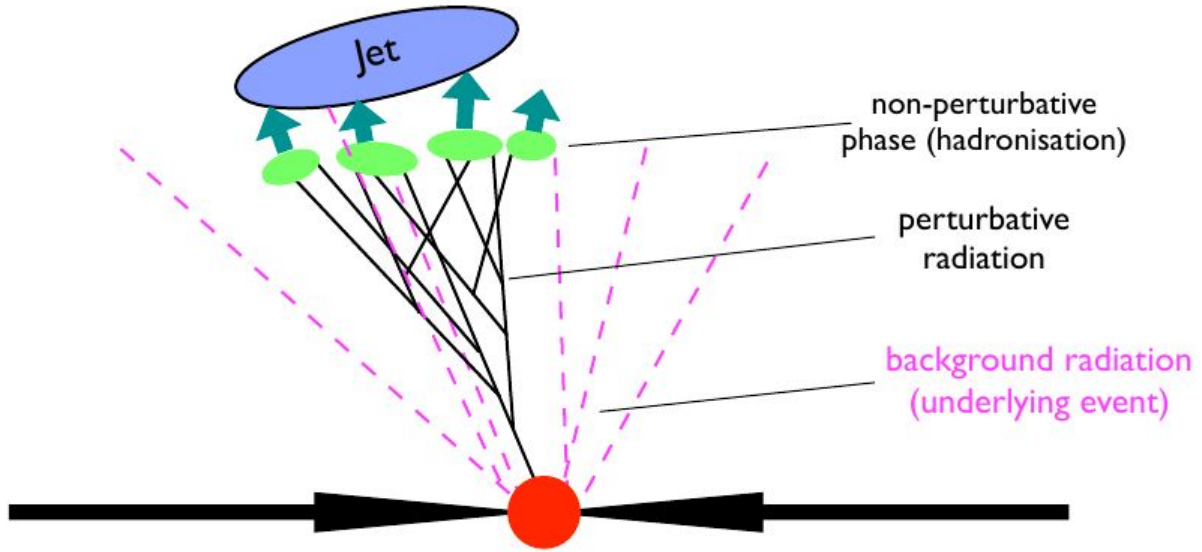
13.82  
 Sig 4.37  
 14

“Multiplicity distributions of gluon and quark jets and tests of QCD analytic predictions”  
 [hep-ex/9708029]

Mean 10.6

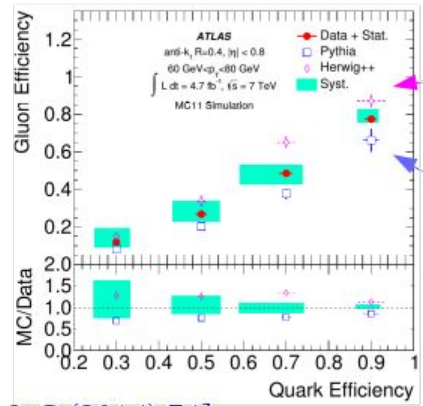
Sig 4.297

8.5



- hadronisation:  $-\Lambda / R$
- MPI:  $+\Lambda * R^2$

Efficiency is simply the ratio of the number of jets selected by a discriminant over the total number in the sample.



Herwig++ is too pessimistic, Quark and gluon jets look more the same than in the data.

Pythia is too optimistic, Quark and Gluon jets are too different compared to data.

[ATLAS, Eur. Phys. J. C (2014) 74]

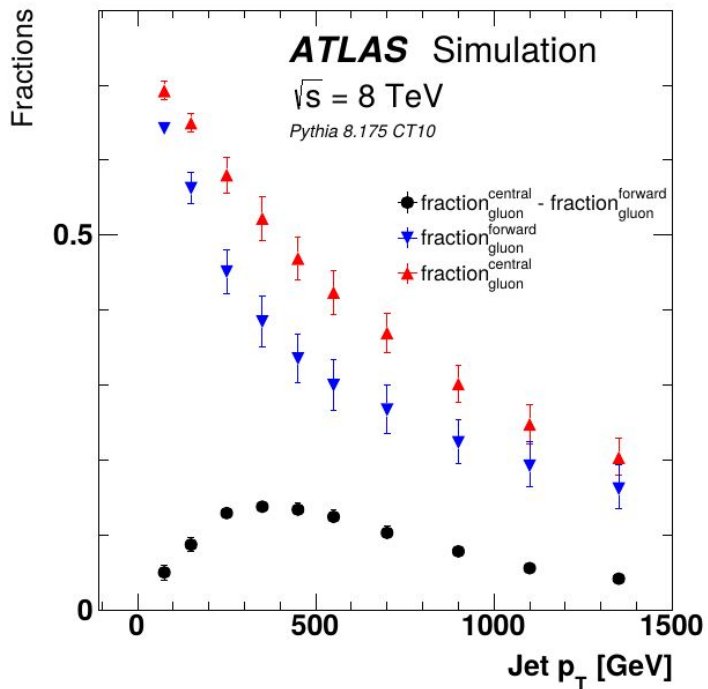
### Conclusion:

*“A detailed study of the jet properties reveals that quark-and gluon-jets look more similar to each other in the data than in the Pythia 6 simulation and less similar than in the Herwig++ simulation.”*

**Problem:** Q/G jets LHC data show discrepancy with the predictions from MC generators

# Quark versus Gluon Jet Tagging Using Charged Particle Multiplicity with the ATLAS Detector

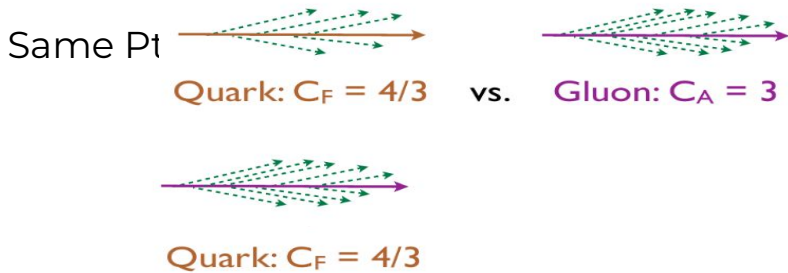
ATLAS Collaboration (Apr 11, 2017)



Using phase space cuts, for example:

- $p_T$  - jet transverse momentum
- $\eta$  - jet rapidity (central/forward)

But then we will have quark and gluon sample jets with different  $(p_T, \eta)$ .

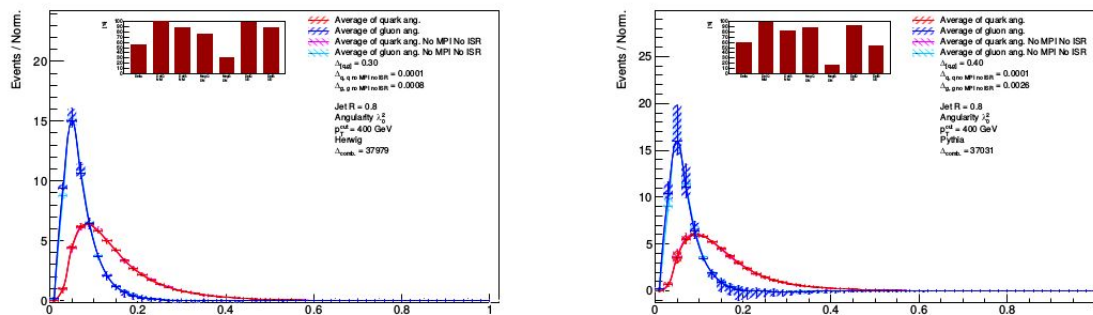


But high  $p_T$  Q will radiate more and look like a G

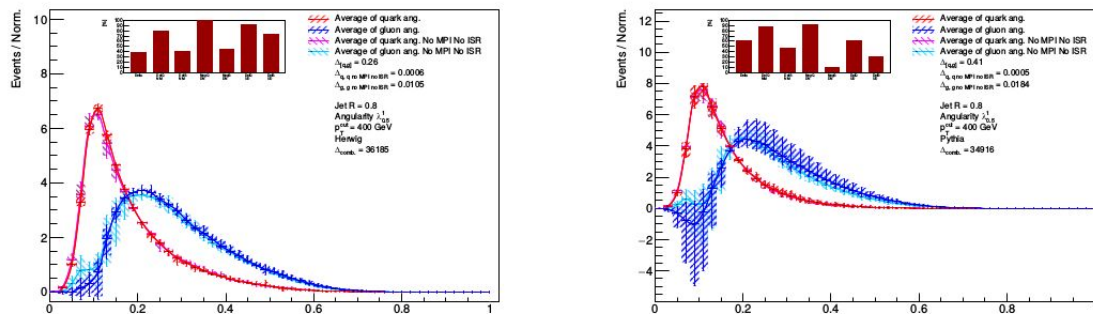
Can we find a way to get enhanced Q/G with the same  $p_T, \eta$ ?



## Herwig vs Pythia

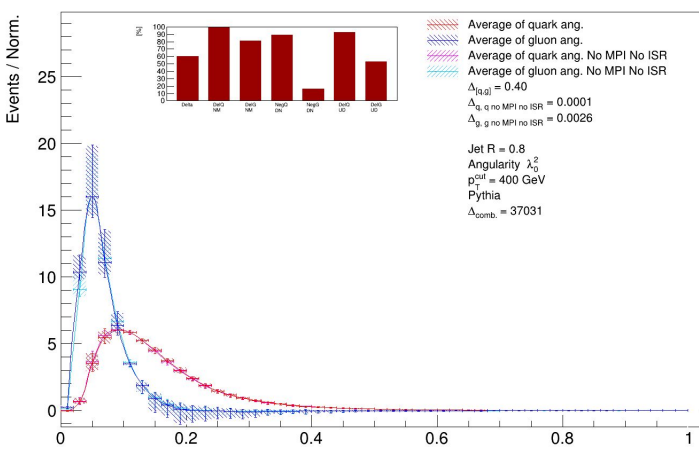
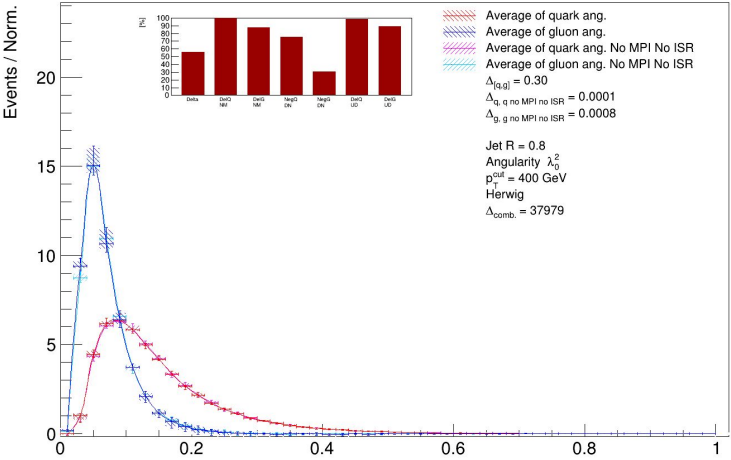


**Fig. 17** Quark and gluon averaged angularities  $\lambda_0^2$ ,  $R = 0.8$  with highest score  $\Delta_{\text{comb}} = 37979$  using HERWIG event generator (left) and  $\Delta_{\text{comb}} = 37031$  using PYTHIA event generator (right), with  $p_T^{\text{cut}} = 400$  GeV, using the average of 6 energy combinations 900–2360, 900–7000, 900–13000, 2360–7000, 2360–13000, 7000–13000 GeV.

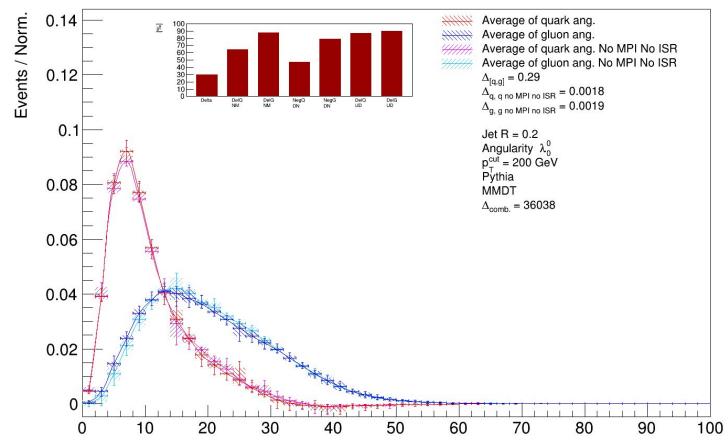
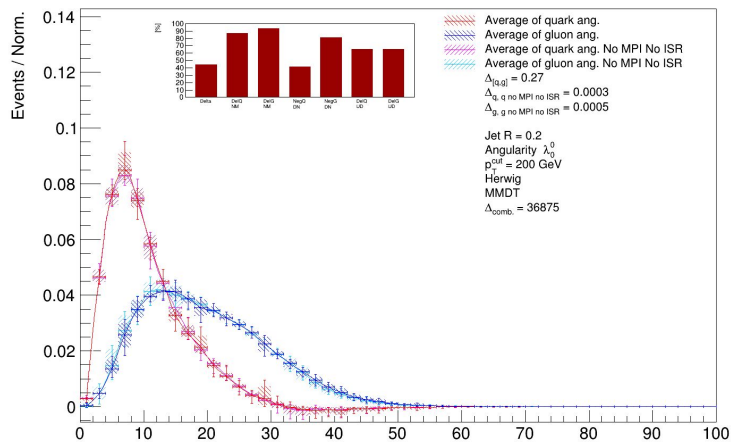


**Fig. 18** Quark and gluon averaged angularities MMDT  $\lambda_{0.5}^1$ ,  $R = 0.8$  with score  $\Delta_{\text{comb}} = 36185$  using HERWIG event generator (left) and  $\Delta_{\text{comb}} = 34916$  using PYTHIA event generator (right), with  $p_T^{\text{cut}} = 400$  GeV, using the average of 6 energy combinations 900–2360, 900–7000, 900–13000, 2360–7000, 2360–13000, 7000–13000 GeV.

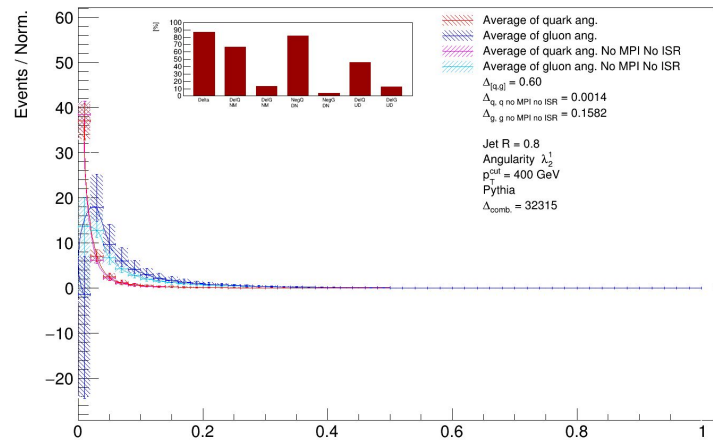
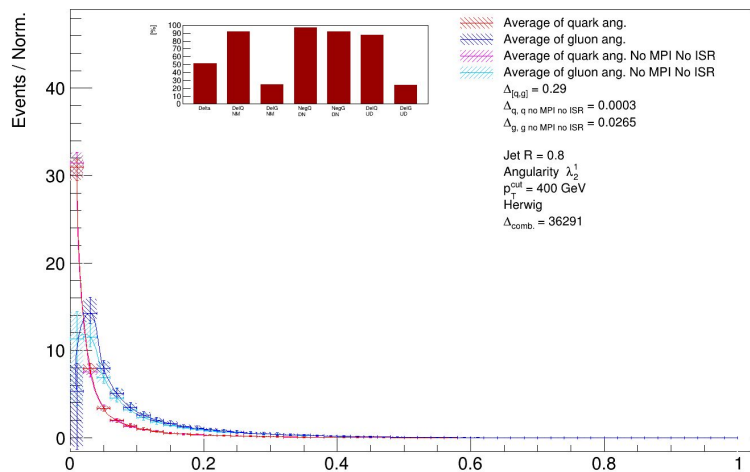
# Best performing angularities: $p_T^D$



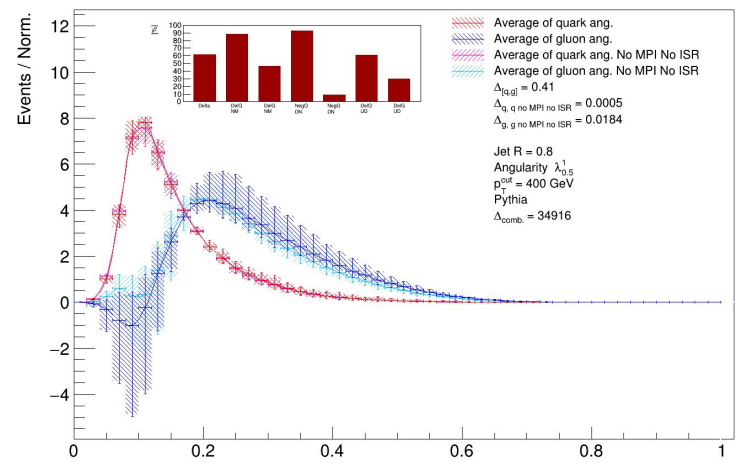
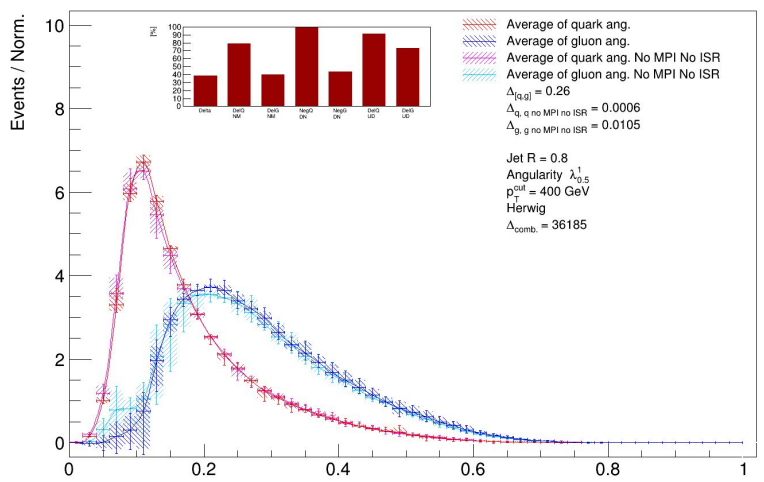
# Best performing angularities: Multiplicity



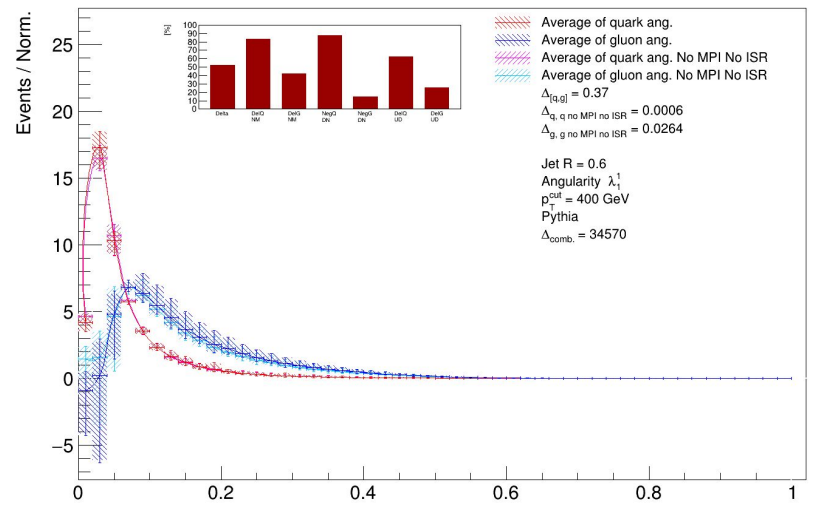
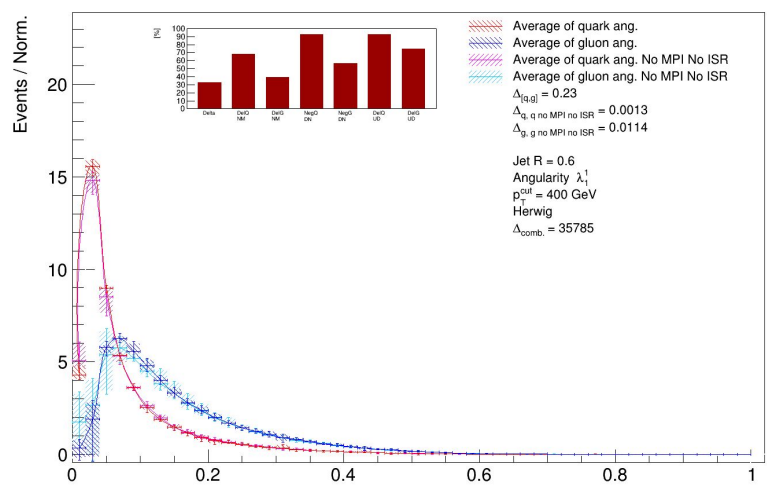
## Best performing angularities: Mass



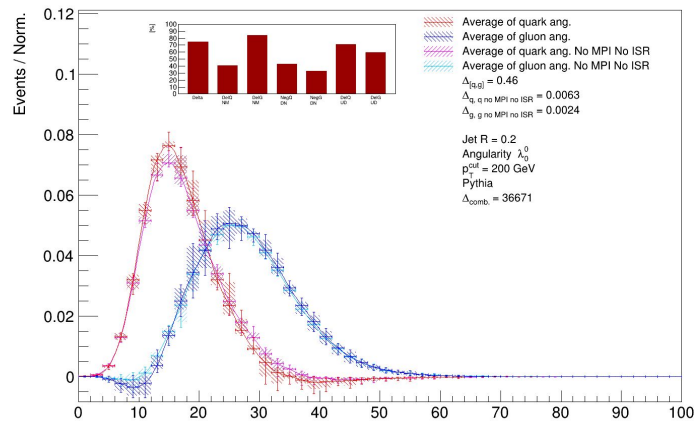
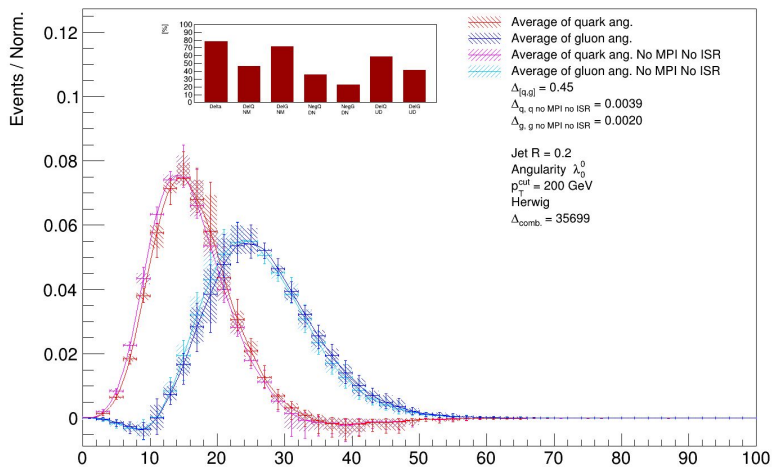
# Best performing angularities: LHA



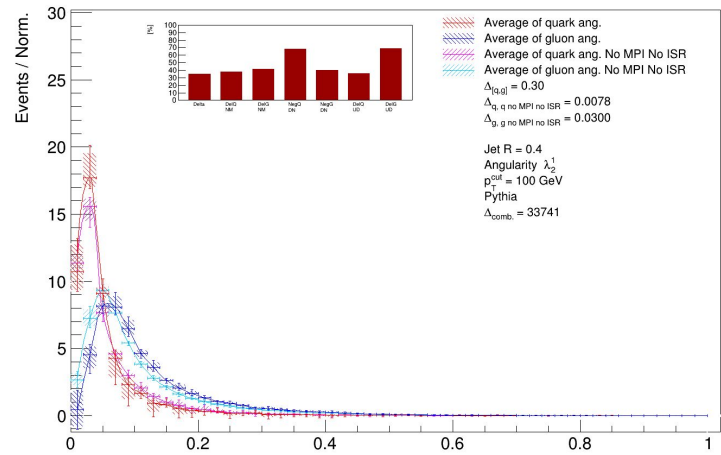
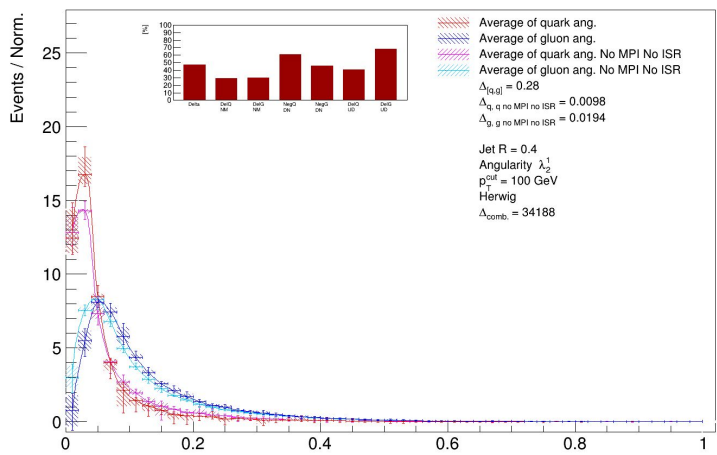
## Best performing angularities: Width



### Wild cards (chosen by "eye"): multiplicity

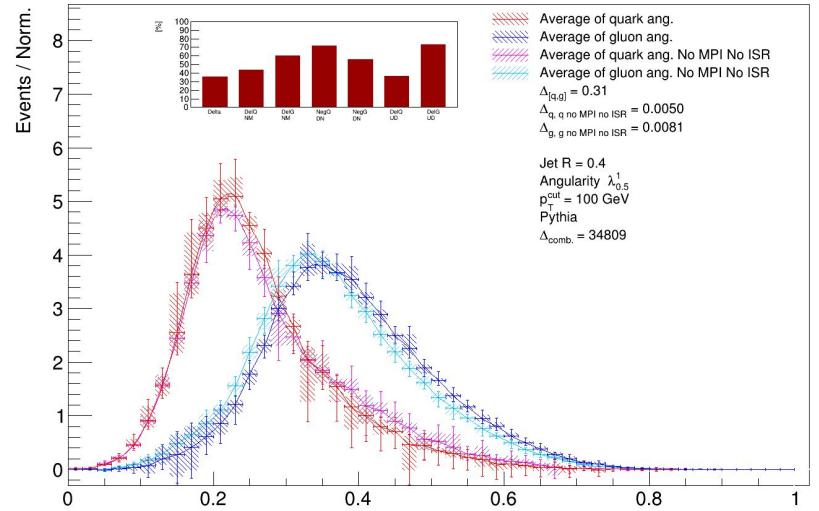
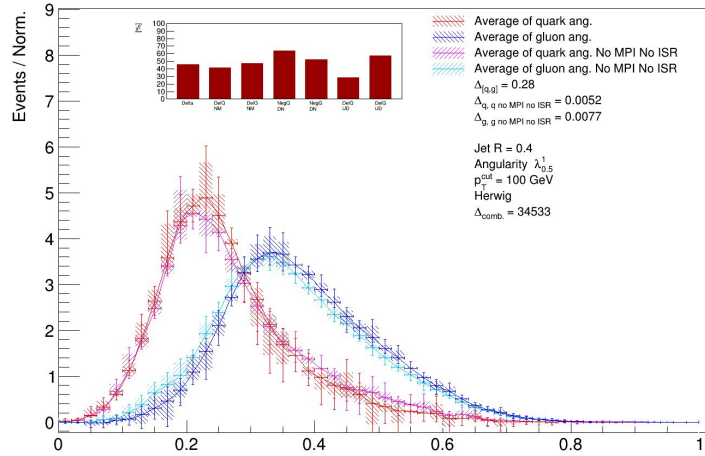


# Wild cards (chosen by "eye"): mass

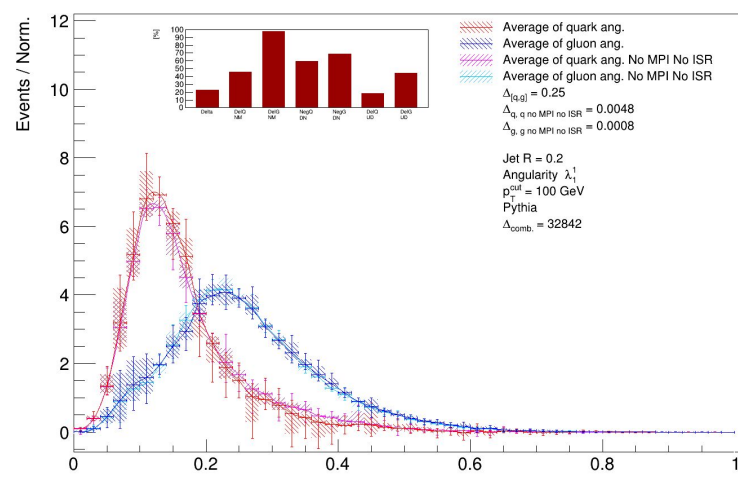
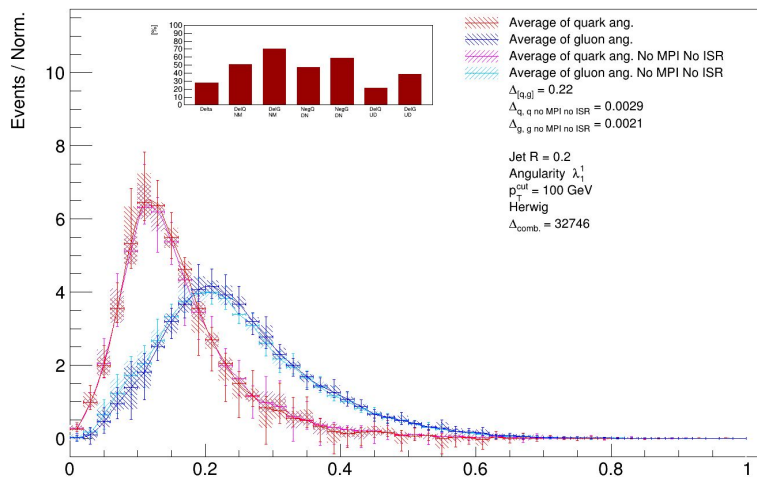




## Wild cards (chosen by “eye”): LHA



# Wild cards (chosen by "eye"): width



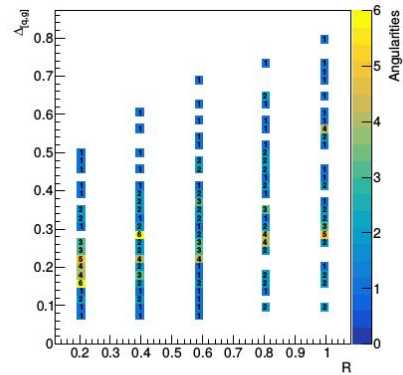


Fig. 19 First column scatter plot of  $\Delta_{[q,g]}$  as a function of jet radius.

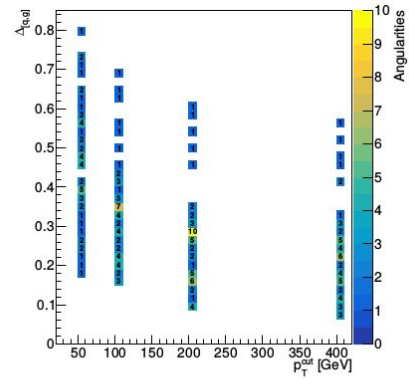


Fig. 21 First column scatter plot of  $\Delta_{[q,g]}$  as a function of  $p_T^{\text{cut}}$ .

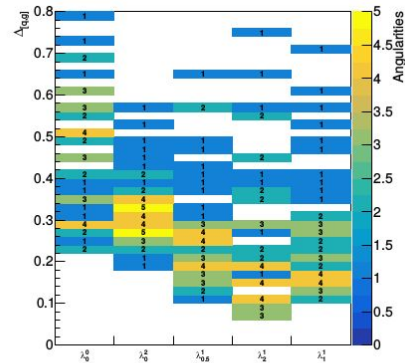


Fig. 20 First column scatter plot of  $\Delta_{[q,g]}$  as a function of jet angularity.

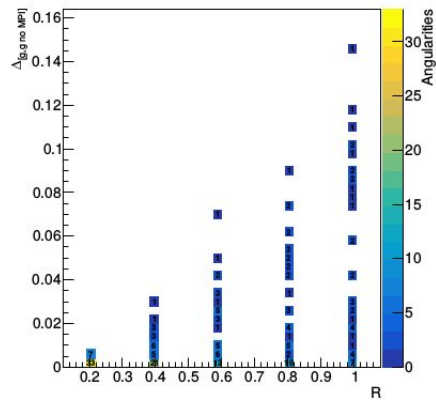
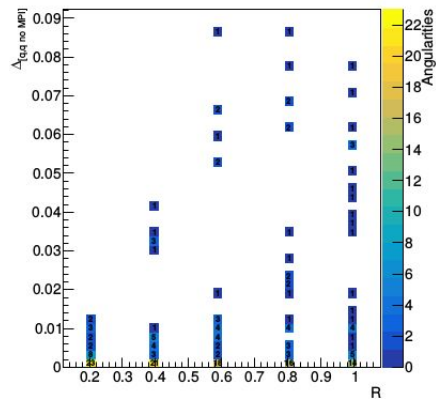


Fig. 22 Second/third column quark  $\Delta_{[q,q] \text{ noMPI}}$  (top) and gluon  $\Delta_{[g,g] \text{ noMPI}}$  (bottom) as a function of jet radius.

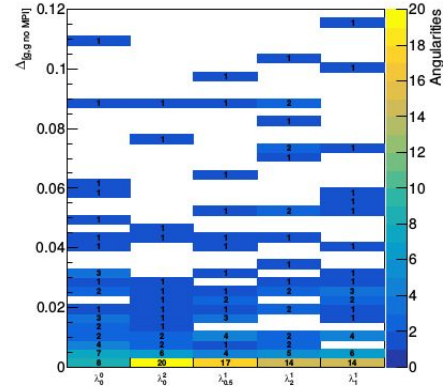
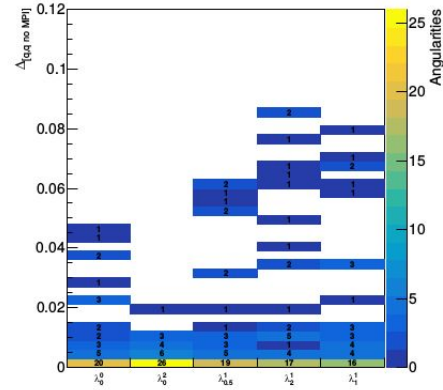
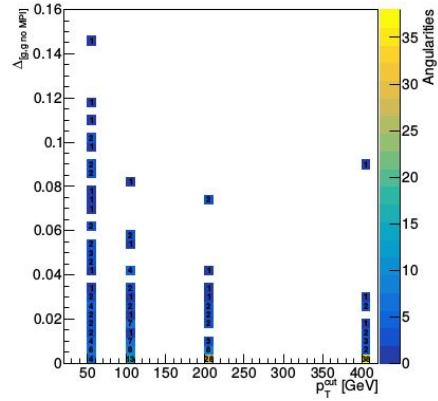
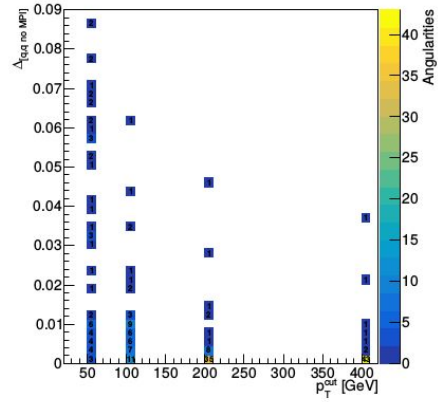
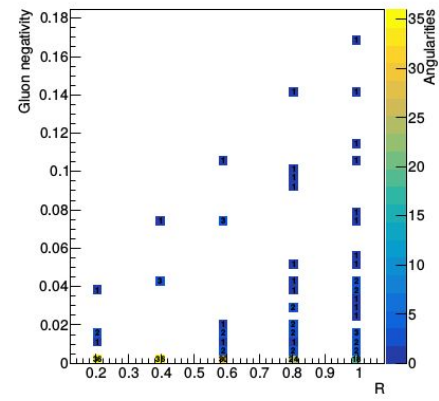
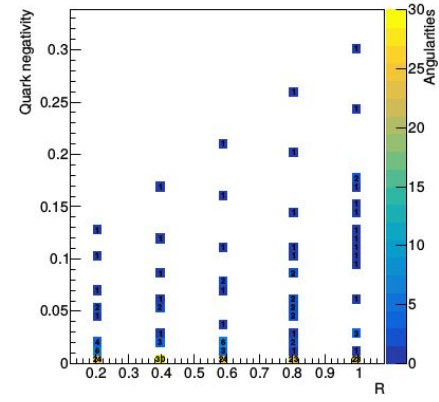


Fig. 23 Second/third column quark  $\Delta_{[q,q] \text{ noMPI}}$  (top) and gluon  $\Delta_{[g,g] \text{ noMPI}}$  (bottom) as a function of angularities.



**Fig. 24** Second/third column quark  $\Delta_{[q,q] \text{ no MPI}}^{(2,3)}$  (top) and gluon  $\Delta_{[g,g] \text{ no MPI}}^{(2,3)}$  (bottom) as a function of  $p_T^{\text{jet}}$ .



**Fig. 25** Fourth/fifth column quark (top) and gluon negativity (bottom) as a function of jet radius.

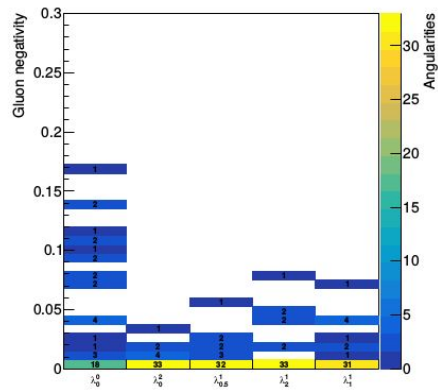
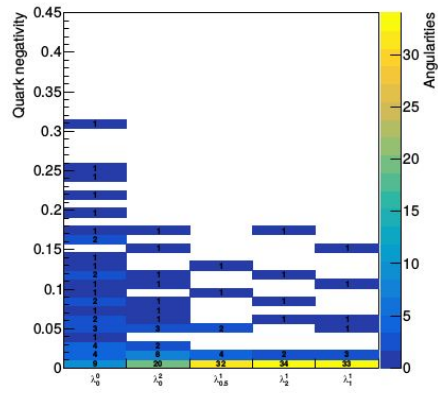


Fig. 26 Fourth/fifth column quark (top) and gluon negativity (bottom) as a function of angularities.

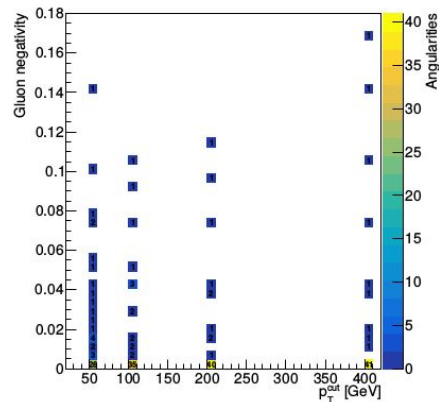
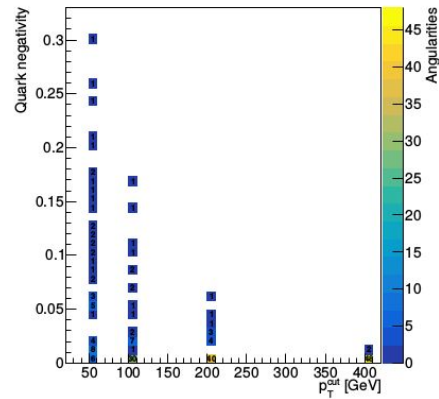


Fig. 27 Fourth/fifth column quark (top) and gluon negativity (bottom) as a function of  $p_T^{c\bar{c}}$ .

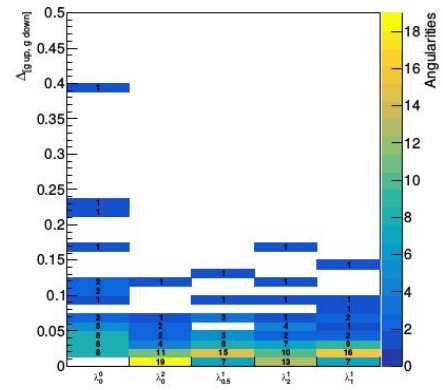
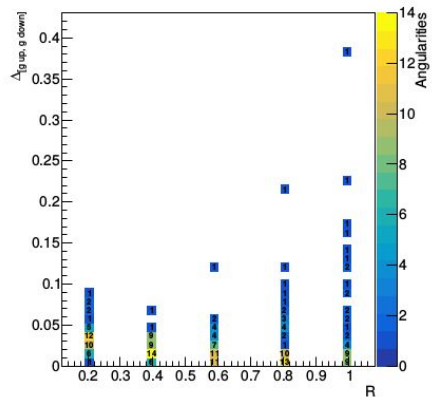
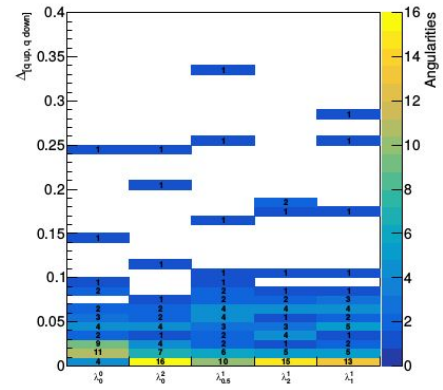
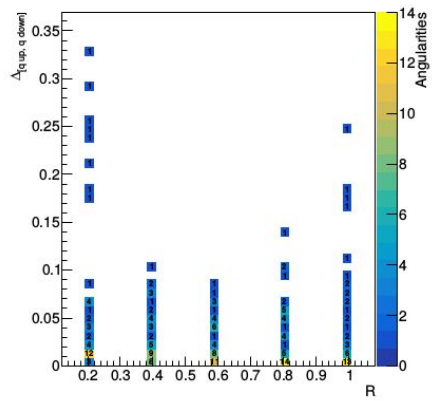


Fig. 28 Sixth/Seventh column quark  $\Delta_{[q \text{ down}, q \text{ up}]}$  (top) and gluon  $\Delta_{[g \text{ down}, g \text{ up}]}$  (bottom) as a function of jet radius.

Fig. 29 Sixth/Seventh column quark  $\Delta_{[q \text{ down}, q \text{ up}]}$  (top) and gluon  $\Delta_{[g \text{ down}, g \text{ up}]}$  (bottom) as a function of angularities.

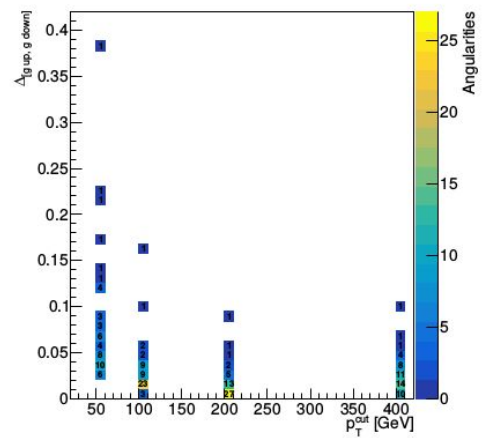
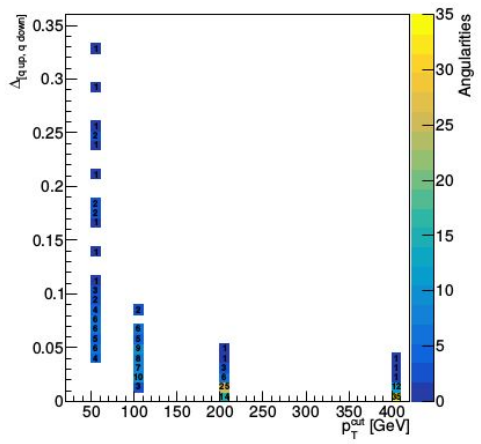


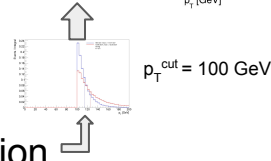
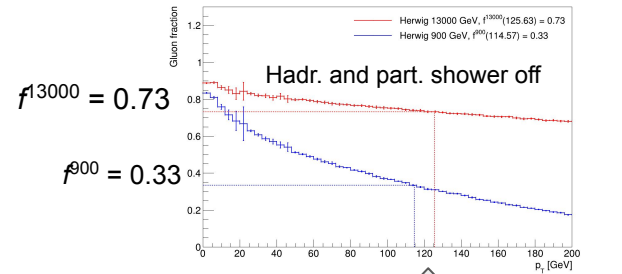
Fig. 30 Sixth/Seventh column quark  $\Delta_{[q_{down}, q_{up}]}$  (top) and gluon  $\Delta_{[g_{down}, g_{up}]}$  (bottom) as a function of  $p_T^{cut}$ .



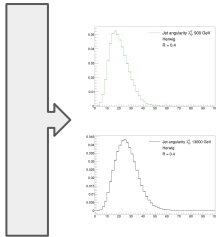
Each angularity  $\lambda$  is composed of gluon  $\lambda_g$  and quark  $\lambda_q$  angularities

$$\lambda = f \lambda_g + (1 - f) \lambda_q$$

where  $f$  ... gluon fraction  
 $(1 - f)$  ... quark fraction  
 taken from simulation



Let's write equations for measurement at energy 900 GeV and 13 000 GeV (with the same event selection)



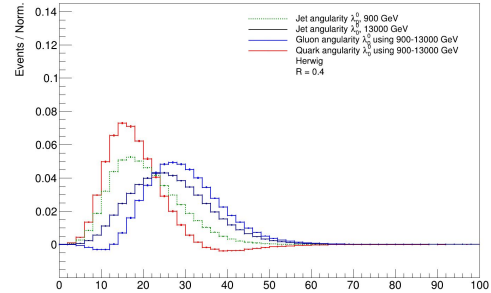
$$\lambda^{900} = f^{900} \lambda_g + (1 - f^{900}) \lambda_q$$

$$\lambda^{13000} = f^{13000} \lambda_g + (1 - f^{13000}) \lambda_q$$

Assuming  $\lambda_g$  and  $\lambda_q$  are energy independent.

$$\lambda_g = \frac{(1 - f^{13000}) \lambda^{900} - (1 - f^{900}) \lambda^{13000}}{f^{900} - f^{13000}}$$

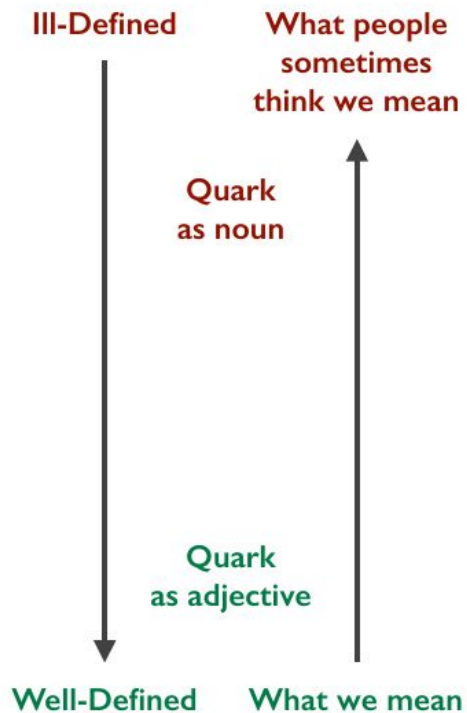
$$\lambda_q = \frac{f^{900} \lambda^{13000} - f^{13000} \lambda^{900}}{f^{900} - f^{13000}}$$



# What is a Quark Jet?

*From lunch/dinner discussions*

[slide by Jesse Thaler]



A quark parton

A Born-level quark parton

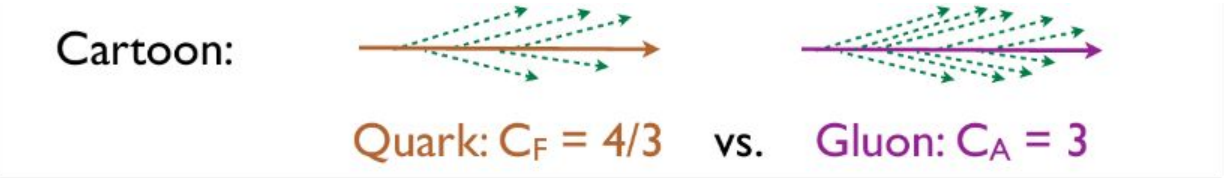
The initiating quark parton in a final state shower

An eikonal line with baryon number  $1/3$  and carrying triplet color charge

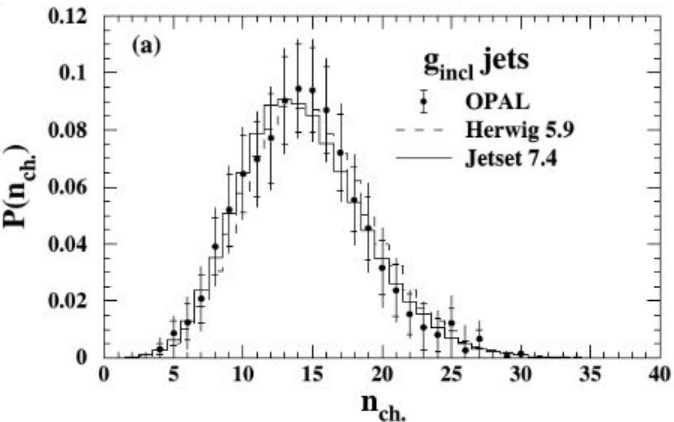
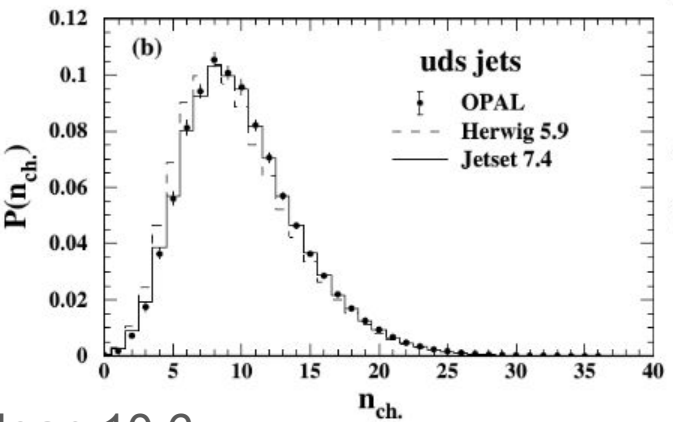
A quark operator appearing in a hard matrix element in the context of a factorization theorem

A parton-level jet object that has been quark-tagged using a soft-safe flavored jet algorithm (automatically collinear safe if you sum constituent flavors)

A phase space region (as defined by an unambiguous hadronic fiducial cross section measurement) that yields an enriched sample of quarks (as interpreted by some suitable, though fundamentally ambiguous, criterion)



Gluon will radiate more, gluon will radiate wider  $\frac{\langle N_g \rangle}{\langle N_q \rangle} = \frac{C_A}{C_F} = 3/(4/3) = 9/4 = 2.25$

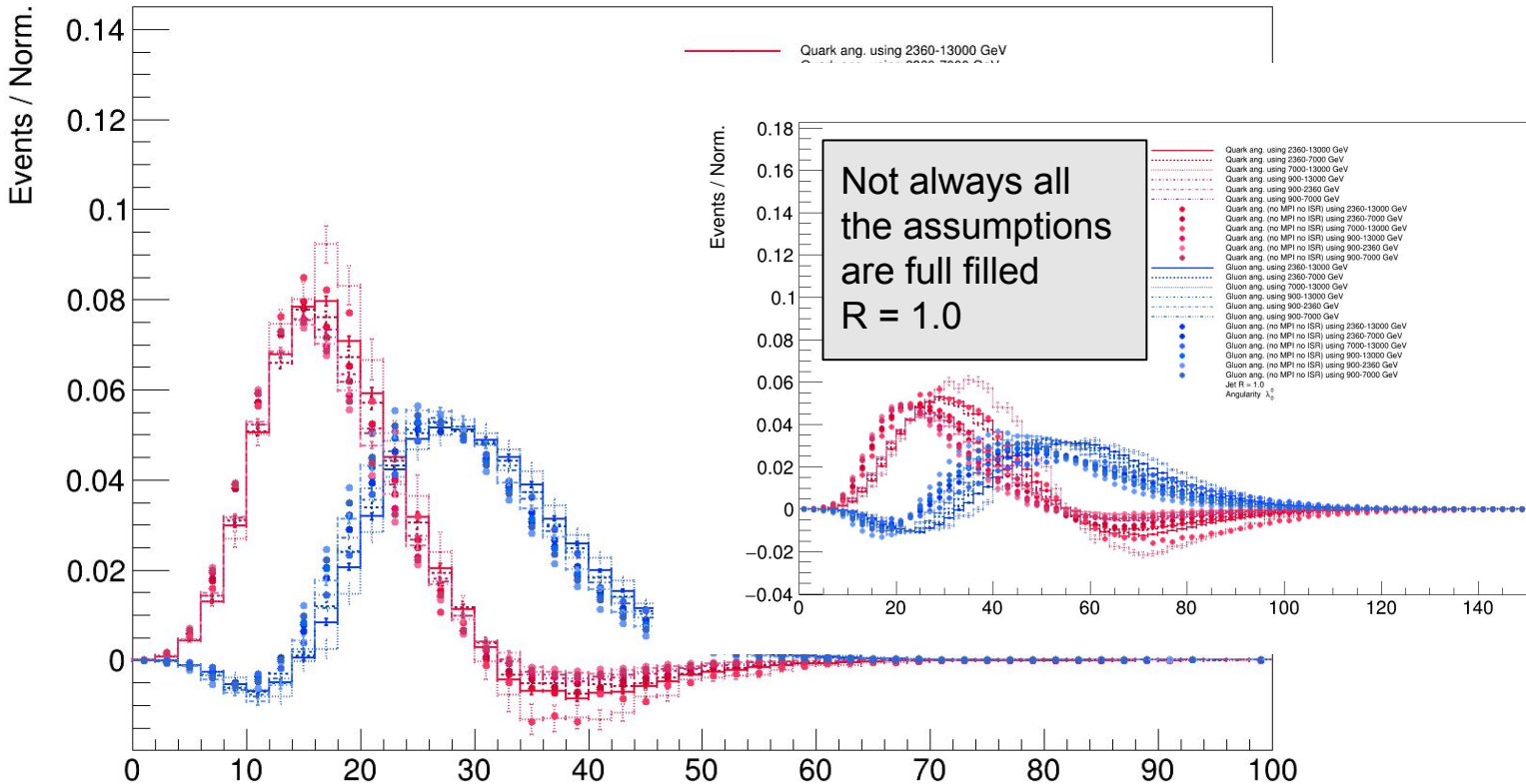


13.82  
 Sig 4.37  
 14  
 41.8 GeV

Mean 10.6  
 Sig 4.297  
 8.5  
 45.6 GeV

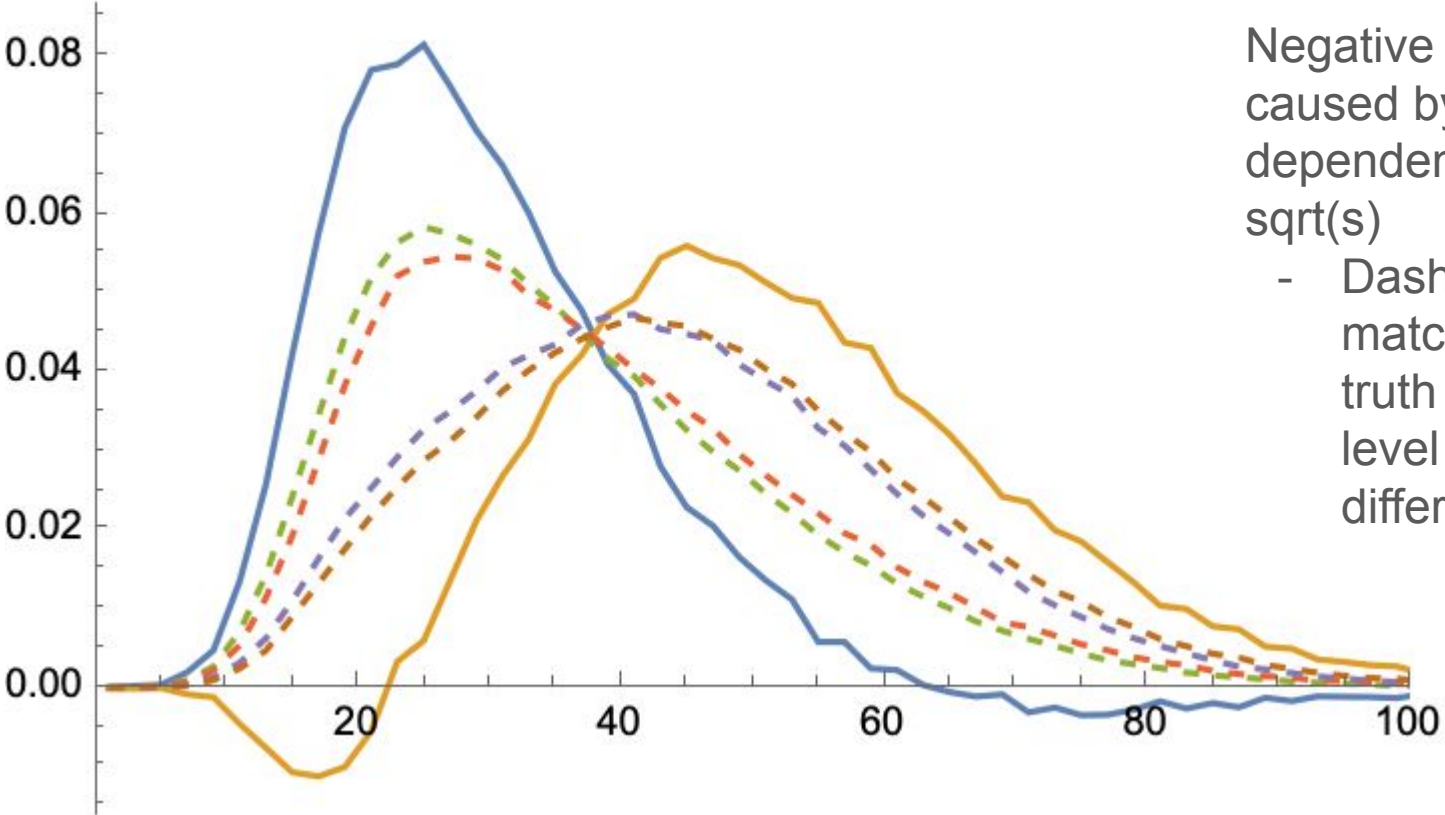
Multiplicity distributions of gluon and quark jets and tests of QCD analytic predictions”  
 [hep-ex/9708029]

Let's use more 6 energy combinations:  
900-2360, 900-7000, 900-13000, 2360-7000, 2360-13000, 7000-13000 GeV



Dotted lines test the robustness to Multi Parton Interactions MPI and Initial State Radiation ISR

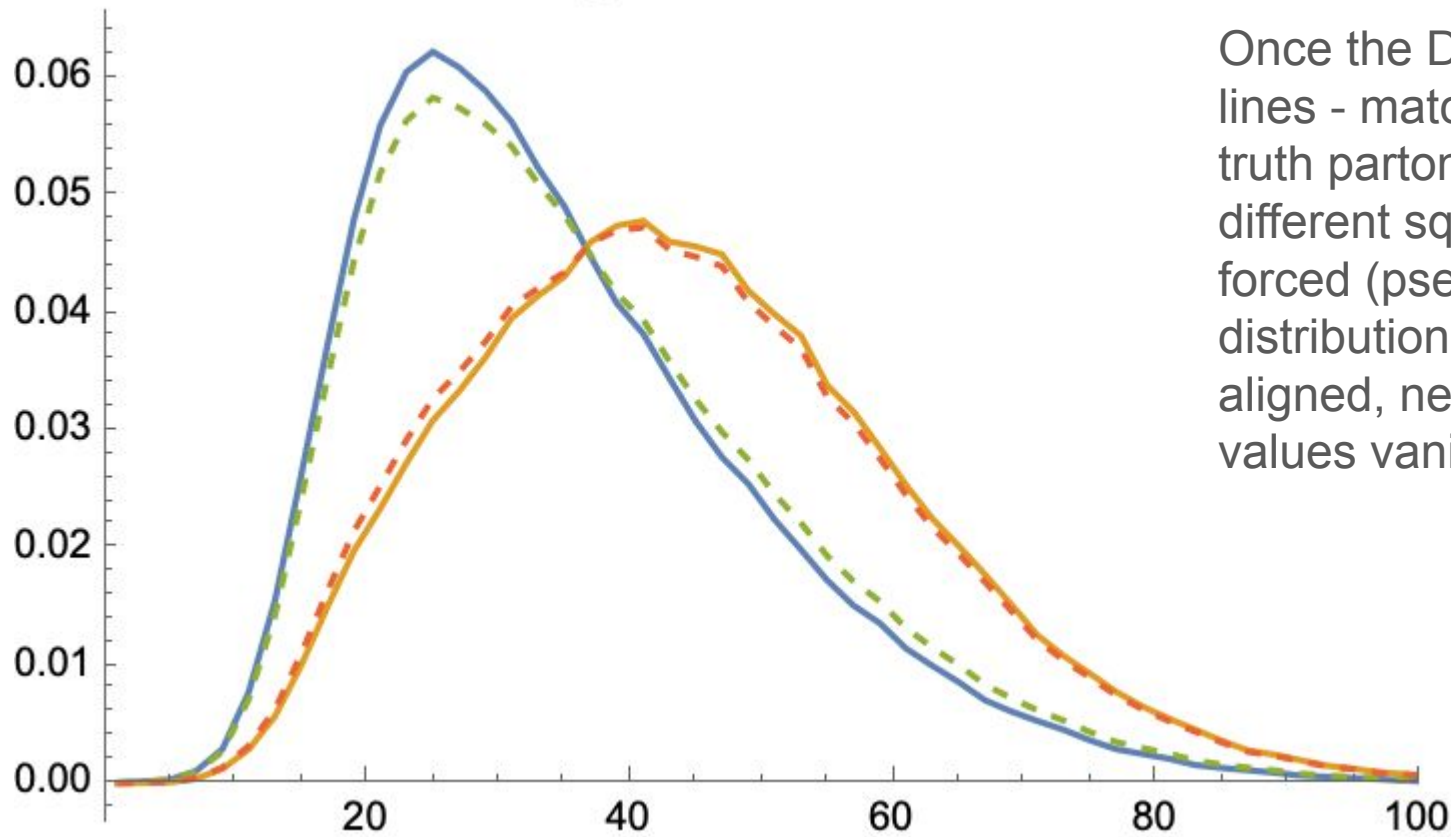
### Unfolded & matched



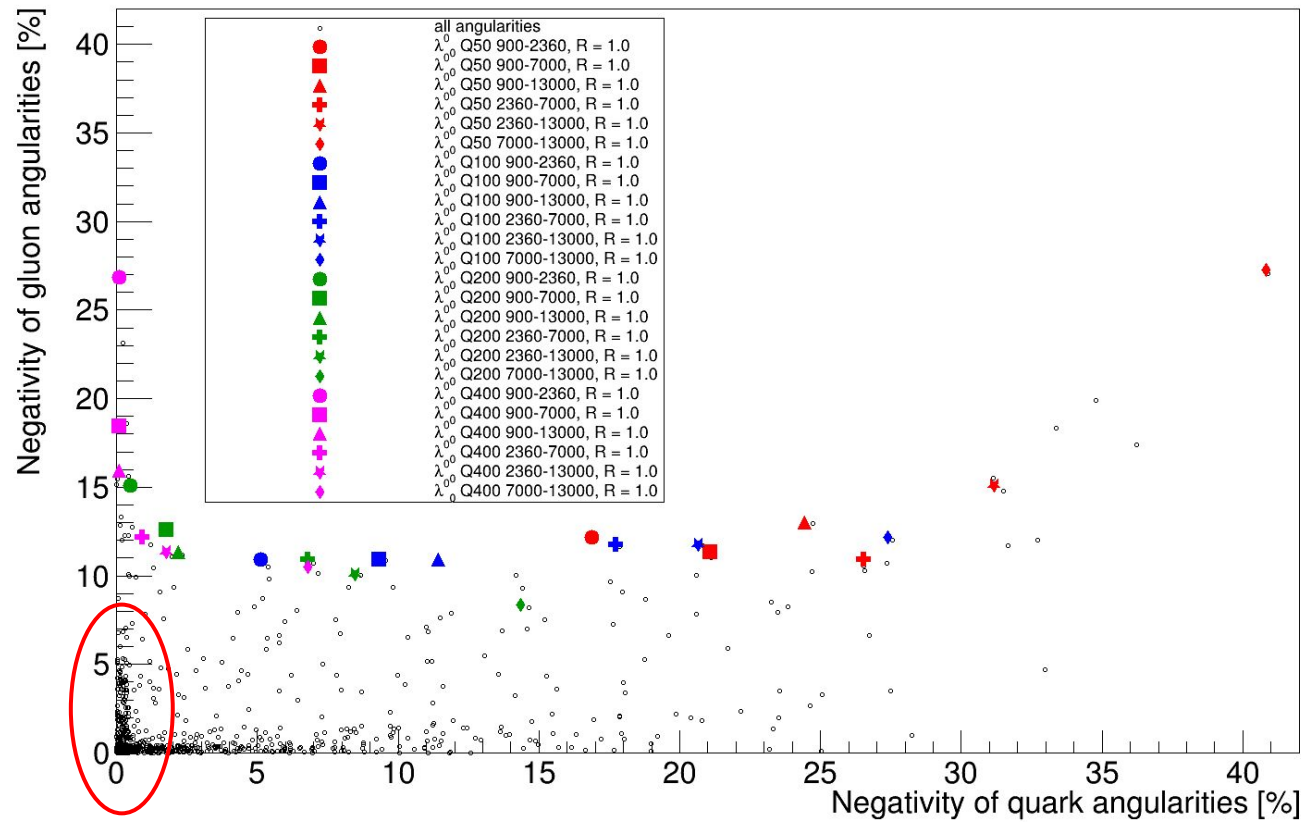
Negative values caused by dependency on  $\sqrt{s}$

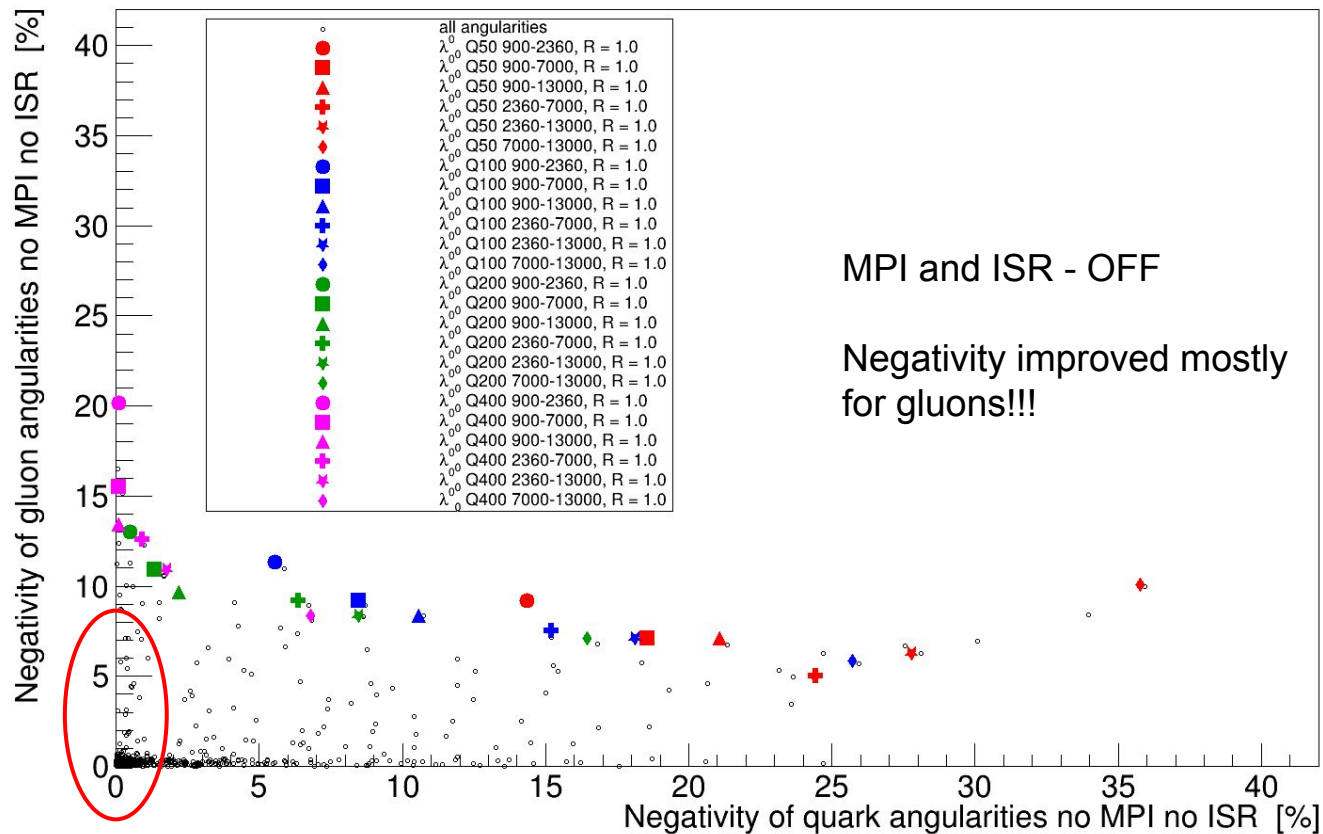
- Dashed lines matched to truth parton level at different  $\sqrt{s}$

Unfolded using fake 13 TeV data & matched

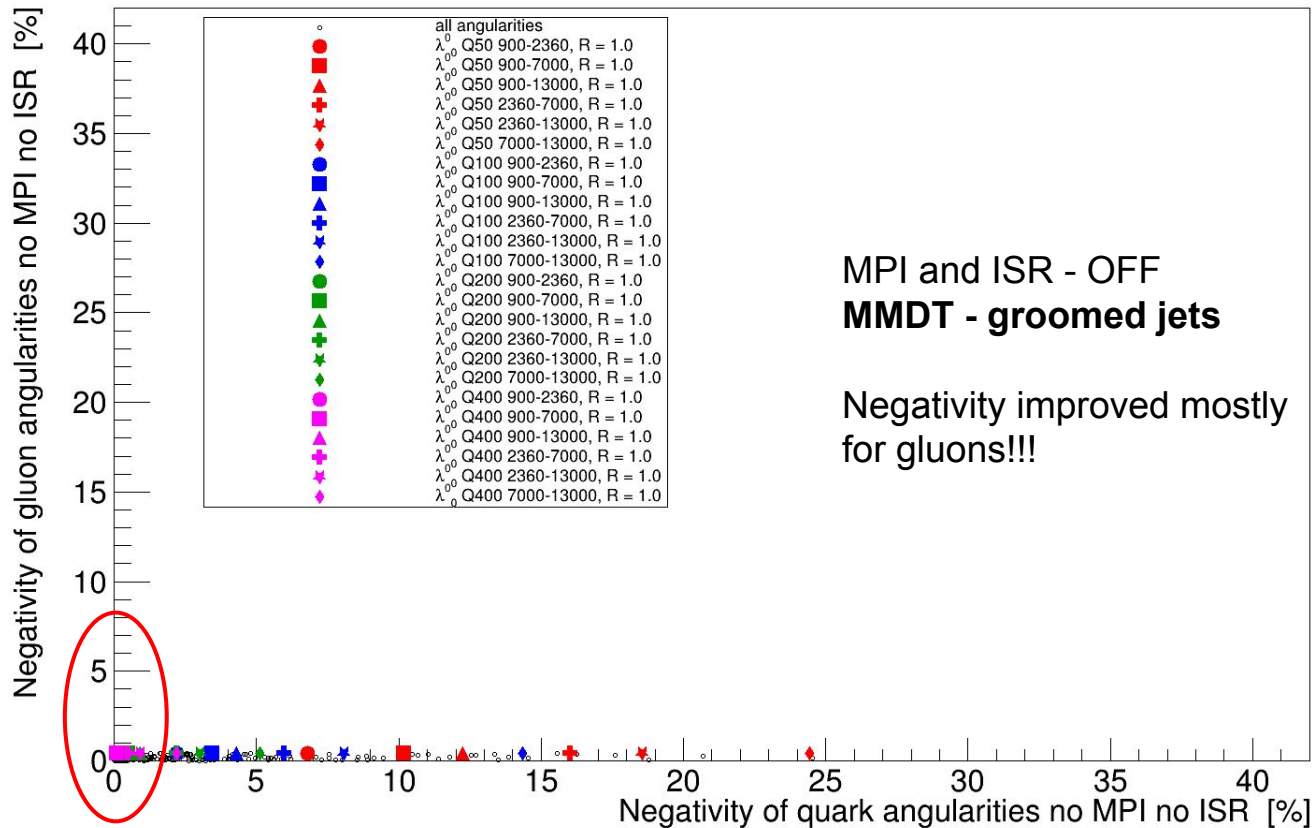


Once the Dashed lines - matched to truth parton level at different sqrt(s) are forced (pseudo-distribution) to be aligned, negative values vanish









ZOOMED

