NLO for hybrid k_T-factorization

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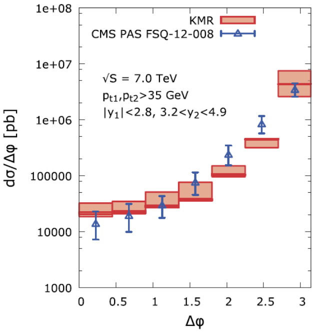
Motivation for study k_T factorization in QCD

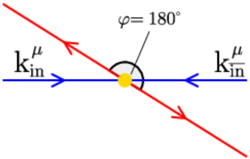
Following the work by M. Deak, F. Hautmann, H. Jung and K. Kutak: Forward jet production at the Large Hadron Collider

- kT-factorization is relevant for hadron collision processes in which one of the hadrons delivers a much smaller momentum fraction x to the partonic process than the other hadron. An example is the production of forward jets.
- in kT-factorization there is a momentum imbalance in the final state allowing for non-trivial distributions already at LO where collinear factorization requires at least NLO. An example is the angle between two jets in dijet production.
- so far, calculations in kT-factorization are mostly performed at LO, and we want to move towards a level of automation at NLO like it exists in collinear factorization.



Distribution of the azimuthal angle.

From A. v. Hameren, P. Kotko, K. Kutak and S. Sapeta: Small-xdynamics in forward–central dijet correlations at the LHC



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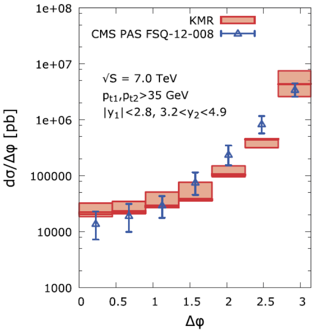

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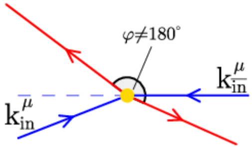


Collinear factorization in QCD at NLO

initial states:

$$k_{in}^{\mu} = x_{in} P^{\mu}$$
$$k_{\overline{in}}^{\mu} = \overline{x}_{\overline{in}} \overline{P}^{\mu}$$

$$d\sigma^{LO} = \int \frac{dx_{in}}{x_{in}} \frac{d\overline{x}_{\overline{in}}}{\overline{x}_{\overline{in}}} f_{in}(x_{in}) f_{\overline{in}}(\overline{x}_{\overline{in}}) dB(x_{in}, \overline{x}_{\overline{in}})$$

$$\begin{aligned} d\sigma^{NLO} = & \int \frac{dx_{in}}{x_{in}} \frac{d\overline{x}_{\overline{in}}}{\overline{x}_{\overline{in}}} \Big\{ f_{in}(x_{in}) f_{\overline{in}}(\overline{x}_{\overline{in}}) \left[dV(x_{in}, \overline{x}_{\overline{in}}) + dR(x_{in}, \overline{x}_{\overline{in}}) \right]_{\text{cancelling}} \\ + & \Big[f_{in}(x_{in}) \frac{-\alpha_s}{2\pi\epsilon} \int_{\overline{x}_{\overline{in}}}^{1} d\overline{z} \mathcal{P}_{\overline{in}}(\overline{z}) f_{\overline{in}}(\overline{x}_{\overline{in}}/\overline{z}) \\ & + f_{\overline{in}}(\overline{x}_{\overline{in}}) \frac{-\alpha_s}{2\pi\epsilon} \int_{x_{in}}^{1} dz \mathcal{P}_{in}(z) f_{in}(x_{in}/z) \Big] dB(x_{in}, \overline{x}_{\overline{in}}) \\ + & \left[f_{in}^{(1)}(x_{in}) f_{\overline{in}}(\overline{x}_{\overline{in}}) + f_{in}(x_{in}) f_{\overline{in}}^{(1)}(\overline{x}_{\overline{in}}) \right] \frac{\alpha_s}{2\pi} dB(x_{in}, \overline{x}_{\overline{in}}) \Big\} \end{aligned}$$

Finite at all

$$f_{\overline{in}}^{(1)}(\overline{x}_{\overline{in}}) - \frac{1}{\epsilon} \int_{\overline{x}_{\overline{in}}}^{1} d\overline{z} \mathcal{P}_{\overline{in}}(\overline{z}) f_{\overline{in}}(\overline{x}_{\overline{in}}/\overline{z}) = finite$$

$$f_{in}^{(1)}(x_{in}) - \frac{1}{\epsilon} \int_{x_{in}}^{1} dz \mathcal{P}_{in}(z) f_{in}(x_{in}/z) = finite$$

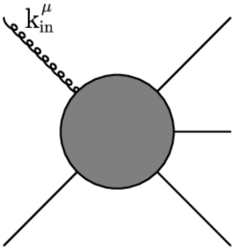
Objective

Hybrid k_T factorization in QCD

Establish the same within hybrid k_T-factorization, for which the LO cross section formula is:

$$d\sigma^{LO} = \int \frac{dx_{in}}{x_{in}} \frac{d^2k_T}{\pi} \frac{d\overline{x}_{\overline{in}}}{\overline{x}_{\overline{in}}} F_{in}(x_{in}, k_T) f_{\overline{in}}(\overline{x}_{\overline{in}}) dB^*(x_{in}, k_T, \overline{x}_{\overline{in}}) \quad (1)$$

- The amplitudes inside $B^*(x_{in}, k_T, \overline{x}_{\overline{in}})$ depend explicitly on k_T.
- They involve a space-like initial-state gluon with momentum $k^\mu_{in} = x_{in}P^\mu + k^\mu_T$



We define k_T as:
$P \cdot k_T = 0$
$\overline{P} \cdot k_T = 0$

- Such amplitudes need care to be well-defined, to be gauge invariant
- We apply the auxiliary-parton method, and our objective is within this constraint

Auxiliary parton method

Introduced by A. v. Hameren, P. Kotko and K. Kutak in Helicity amplitudes for high-energy scattering.

We put our interest on process with one space-like gluon. $\omega(p_1) = g(p_1)/q(P_1)$

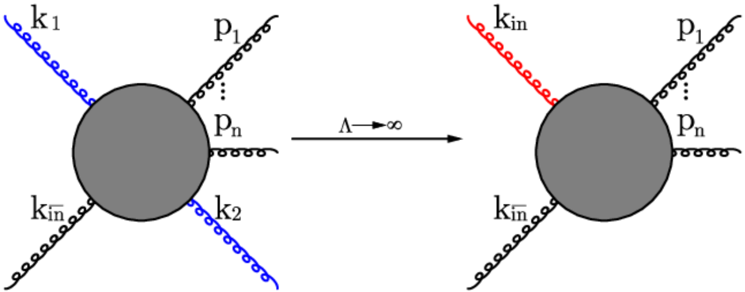
$$g^*(k_{in})\omega_{\overline{in}}(k_{\overline{in}}) \to \omega_1(p_1)\omega_2(p_2)\cdots\omega_n(p_n).$$

This process is obtained via named auxiliary parton method from process

$$q(k_1(\Lambda))\omega_{\overline{in}}(k_{\overline{in}}) \to q(k_2(\Lambda))\omega_1(p_1)\omega_2(p_2)\cdots\omega_n(p_n)$$

with light-like momenta parametrized with Λ

$$k_1^\mu = \Lambda P^\mu,\ k_2^\mu = p_\Lambda^\mu = (\Lambda - x_{in})P^\mu - k_T^\mu + \frac{|k_T|^2}{2P^\mu \cdot \overline{P}^\mu(\Lambda - x_{in})}\overline{P}^\mu.$$



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Their difference is

$$k_1^\mu - k_2^\mu = k_{in}^\mu + O(\Lambda^{-1}) = x_{in}P^\mu + k_T^\mu + O(\Lambda^{-1})$$

Taking $\Lambda \to \infty$ one will obtain the matrix element with space-like gluon

$$\frac{x_{in}^2|k_T|^2}{g_s^2 C_{aux}\Lambda^2}|\overline{M}^{aux}|^2(\Lambda P, k_{\overline{in}}; p_\Lambda, \{p_i\}_{i=1}^n) \xrightarrow{\Lambda\to\infty} |\overline{M}^*|^2(k_{in}, k_{\overline{in}}; \{p_i\}_{i=1}^n) \quad (2)$$

As auxiliary partons we can choose quarks as well as gluons. Then

$$C_{aux-q} = \frac{N_c^2 - 1}{N_c},\ C_{aux-g} = 2N_c.$$

The NLO virtual contribution

Virtual contributions

$$dV^* = dV^{*fam} + dV^{*unf}$$

- Familiar contribution conserve smooth on-shell $k_T \to 0$
- Unfamiliar contribution $dV^{*unf} = a_\epsilon N_c Re(\mathcal{V}_{aux}) dB^*$

$$a_\epsilon = \frac{\alpha_s}{2\pi}\frac{(4\pi)^\epsilon}{\Gamma(1-\epsilon)}; \quad \epsilon = \frac{4-dim}{2}$$

$$\mathcal{V}_{aux} = \left(\frac{\mu^2}{|k_T|^2}\right)^\epsilon \left[\frac{2}{\epsilon} ln\frac{\Lambda}{x_{in}} - i\pi + \overline{\mathcal{V}}_{aux}\right] + \mathcal{O}(\epsilon) + \mathcal{O}(\Lambda^{-1})$$

$$\overline{\mathcal{V}}_{aux-q} = \frac{1}{\epsilon}\frac{13}{6} + \frac{\pi^2}{3} + \frac{80}{18} + \frac{1}{N_c^2}\left[\frac{1}{\epsilon^2} + \frac{3}{2}\frac{1}{\epsilon} + 4\right] - \frac{n_f}{N_c}\left[\frac{2}{3}\frac{1}{\epsilon} + \frac{10}{9}\right]$$

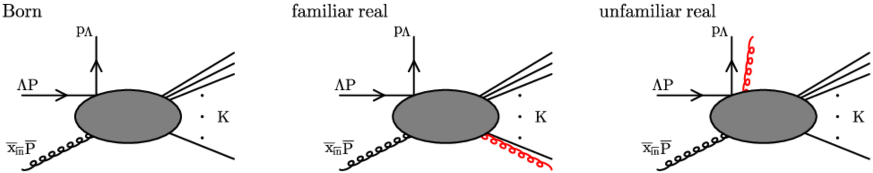
$$\overline{\mathcal{V}}_{aux-g} = -\frac{1}{\epsilon^2} + \frac{\pi^2}{3}$$

Details in E. Blanco, A. Giachino, A. v. Hameren, P. Kotko: One-loop gauge invariant amplitudes with a space-like gluon.

Unfamiliar real contribution

$$dR^* = dR^{*fam} + dR^{*unf}$$

In the unfamiliar case the radiative gluon participates in the consumption of Λ $\quad k_T = q_T + r_T$



The phase space also factorizes, we can perform analytical integration, the result is:

$$dR^{*unf}(k_{in}, k_{\overline{in}}; \{p_i\}_{i=1}^{n+1}) = \left\{ a_\epsilon N_c \left(\frac{\mu^2}{|k_T|^2}\right)^\epsilon \left[-\frac{2}{\epsilon} ln \frac{2P\cdot\overline{P}\Lambda}{|k_T|^2} + \overline{R}_{aux} \right] + \mathcal{O}(\epsilon, \Lambda^{-1}) \right\} dB^*(k_{in}, k_{\overline{in}}; \{p_i\}_{i=1}^{n})$$

- depends of type of auxiliary partons
- violates the smooth on-shell limit and smooth large Λ limit

Unfamiliar contributions - completed

Collection of virtual and real unfamiliar contribution brings

$$\Delta_{unf} dB^* = dR^{*unf} + dV^{*unf}$$

general unfamiliar contribution is given by

$$\Delta_{unf} = \frac{a_\epsilon N_c}{\epsilon}\left(\frac{\mu^2}{|k_T|^2}\right)^\epsilon \left[\mathcal{J}_{aux} + \mathcal{J}_{univ} + \mathcal{J}_{univ} - 2ln\frac{2P\cdot \overline{P}x_{in}}{|k_T|^2}\right]$$

where

$$\mathcal{J}_{univ} = \frac{11}{6} - \frac{n_f}{3N_c} - \frac{\mathcal{K}}{N_c}(-\epsilon) \qquad \text{with} \qquad \mathcal{K} = N_c\left(\frac{67}{18} - \frac{\pi^2}{6}\right) - \frac{5n_f}{9}$$

$$\mathcal{J}_{aux-g} = \frac{11}{6} + \frac{n_f}{3N_c^3} + \frac{n_f}{6N_c^3}(-\epsilon), \qquad \mathcal{J}_{aux-q} = \frac{3}{2} - \frac{1}{2}(-\epsilon)$$

- No $ln\Lambda$
- Target impact factor corrections as in Ciafaloni, Colferai 1999.
- Other terms also known in literature (Regge trajectory, renormalization of the coupling constant)

Familiar real collinear singularities

The dR^{*fam} has a singularity when a radiative gluon becomes collinear to $\overline{P}$ which leads to divergence $\Delta_{\overline{coll}}$ with splitting as $\dfrac{1}{z(1-z)} - 2 + z(1-z)$ included.

Despite k$_T$, tree-level matrix elements with a space-like gluon still have a singularity when a radiative gluon becomes collinear to P.

$$|\overline{M}^*|^2\left(x_{in}P + k_T, k_{\overline{in}}; r, \{p_i\}_{i=1}^n\right) \xrightarrow{r \to x_r P} \frac{2N_C}{\color{red}{P \cdot r}} \frac{x_{in}^2}{x_r(x_{in}-x_r)^2} |\overline{M}^*|^2\left((x_{in}-x_r)P + k_T, k_{\overline{in}}; \{p_i\}_{i=1}^n\right) \tag{3}$$

Similar to usual collinear gluon splitting with only the $\dfrac{1}{z(1-z)}$ part.

This leads to a non-cancelling divergence similar to the collinear case given by

$$\Delta^*_{coll}(x_{in}, k_T) = -\frac{\alpha_\epsilon}{\epsilon}\int_{x_{in}}^1 dz \left[\frac{2N_C}{[1-z]_+} + \frac{2N_C}{z} + \gamma_g \delta(1-z)\right] F\left(\frac{x_{in}}{z}, k_T\right) \tag{4}$$

Completed cross section formula

General NLO formula

$$d\sigma^{NLO} = \int \frac{dx_{in}}{x_{in}} \frac{d^2 k_T}{\pi} \frac{d\overline{x}_{\overline{in}}}{\overline{x}_{\overline{in}}} \Big\{ F_{in}(x_{in}, k_T) f_{\overline{in}}(\overline{x}_{\overline{in}}) \left[dR^*(x_{in}, k_T, \overline{x}_{\overline{in}}) + dV^*(x_{in}, k_T, \overline{x}_{\overline{in}}) \right]_{cancelling}$$
$$+ \left[F_{in}^{NLO}(x_{in}, k_T) + F_{in}(x_{in}, k_T)\Delta_{unf}(x_{in}, k_T) + \Delta^*_{coll}(x_{in}, k_T) \right] f_{\overline{in}}(\overline{x}_{\overline{in}}) dB^*(x_{in}, k_T, \overline{x}_{\overline{in}})$$
$$\left[f^{NLO}\overline{in}(\overline{x}_{\overline{in}}) + \Delta_{\overline{coll}} \right] F_{in}(x_{in}, k_T) dB^*(x_{in}, k_T, \overline{x}_{\overline{in}}) \Big\} \qquad (5)$$

Details in A. v. Hameren, L. Motyka, G. Ziarko: Hybrid kT-factorization and impact factors at NLO. J. High Energ. Phys. 2022, 103 (2022). https://doi.org/10.1007/JHEP11(2022)103 [SPRINGER]

The collinear divergences Δ^*_{coll} and $\Delta_{\overline{coll}}$
$f^{NLO}\overline{in}(\overline{x}_{\overline{in}}) + \Delta_{\overline{coll}} \rightarrow$ finite as in collinear factorization
$F_{in}^{NLO}(x_{in}, k_T) + F_{in}(x_{in}, k_T)\Delta_{unf}(x_{in}, k_T) + \Delta^*_{coll}(x_{in}, k_T) \rightarrow$ still necessity for scheme for renormalization of PDFs

Progress with subtraction method for Δ^*_{coll} in A. Giachino, A. v. Hameren, G. Ziarko: A new subtraction scheme at NLO exploiting the privilege of k$_T$-factorization. [arxiv.org/2312.02808]

Sumary

- We established the framework for calculating cross section within k_T factorization.
- We showed the consistency in our scheme (no $ln(\Lambda)$).
- All divergences are recognized.

Thank you for listening!

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Back-Up

Real radiation - familiar contribution

Real contribution we defined as

$$dR^{*\,fam}(k_{in}, k_{\overline{in}}; \{p_i\}_{i=1}^{n+1}) = \frac{a_\epsilon \mu^{2\epsilon}}{\pi_\epsilon} \frac{1}{|k_T|^2} d\Sigma^*_{n+1}(k_{in}, k_{\overline{in}}; \{p_i\}_{i=1}^{n+1}) J_R(\{p_i\}_{i=1}^{n+1}) \quad (6)$$

$$a_\epsilon = \frac{\alpha_s}{2\pi} \frac{(4\pi)^\epsilon}{\Gamma(1-\epsilon)}; \quad \pi_\epsilon = \frac{\pi^{1-\epsilon}}{\Gamma(1-\epsilon)}$$

- One parton more in a final state (compared to Born)
- One collinear pair and / or one soft parton
- The singularities look the same as if the initial-state gluon were on-shell

- Independent of the type of auxiliary partons
- No $ln\Lambda$

There is also unfamiliar component:

$$dR^* = dR^{*\,fam} + dR^{*\,unf}$$

Subtraction method

$$\int_0^1 dx \frac{1}{x} f(x) \tag{7}$$

is straightforwardly divergent.

$$\int_0^1 dx \frac{x^\epsilon}{x} f(x) \tag{8}$$

To demonstrate this, we can apply the Taylor expansion:

$$\begin{aligned}\int_0^1 dx \frac{x^\epsilon}{x} f(x) &= \int_0^1 dx \frac{x^\epsilon}{x} [f(x) - f(0)] + \int_0^1 dx \frac{x^\epsilon}{x} f(0) \\ \int_0^1 dx \frac{x^\epsilon}{x} f(x) &= \int_0^1 dx \frac{x^\epsilon}{x} \left[f(0) + x \cdot f'(0) + \frac{1}{2} x^2 \cdot f''(0) + \ldots - f(0) \right] + \int_0^1 dx \frac{x^\epsilon}{x} f(0)\end{aligned} \tag{9}$$

Doing the second integral on the right hand side we get:

$$\int_0^1 dx \frac{x^\epsilon}{x} f(x) = \int_0^1 dx \cdot x^\epsilon \left[\cdot f'(0) + \frac{1}{2} x \cdot f''(0) + \ldots \right] + \frac{1}{\epsilon} f(0) x^\epsilon \Big|_0^1 \tag{10}$$