The signal and the noise

Yihui Quek **MIT**

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Can we compute with noisy devices = can we extract any signal from the noise?

Answer: depends on the noise!

A popular noise model

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\text{Depolarizing noise: } \mathcal{D}_p(\rho) = p \cdot \rho + (1 - p) \cdot \frac{\mathbb{I}}{2}
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Noisy circuit: Every gate is followed by ${\cal D}_{p}$

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But what about other sources (T1, decay, readout error etc) that can decrease entropy? \rightarrow drives towards non m.m. state

Depolarizing noise is usually bad news

Limitations of variational quantum algorithms: a quantum optimal transport approach

Giacomo De Palma, Milad Marvian, Cambyse Rouzé, Daniel Stilck França

The impressive progress in guantum hardware in the last years has raised the interest of the guantum computing community in harvesting the computational power of such devices. However, in the absence of error correction, these devices can only reliably implement very shallow circuits or comparatively deeper circuits at the expense of a nontrivial density of errors. In

this work, we obtain extremely tight limitation bounds for standard NISQ proposals in both the noisy with or without error-mitigation tools. The bounds limit the performance of both circuit model algorit and also continuous-time algorithms, such as quantum annealing. In the noisy regime with local depo prove that at depths $L = \Box(p^{-1})$ it is exponentially unlikely that the outcome of a noisy quantum circ classical algorithms for combinatorial optimization problems like Max-Cut. Although previous results

A Polynomial-Time Classical Algorithm for Noisy Random Circuit **Sampling**

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ABSTRACT

We give a polynomial time classical algorithm for sampling from the output distribution of a noisy random quantum circuit in the regime of anti-concentration to within inverse polynomial total variation distance. The algorithm is based on a quantum analog of noise induced low degree approximations of Boolean functions, which takes the form of the truncation of a Feynman path integral in the Pauli basis.

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are collected from the experimental implementation of RCS (though this number must necessarily scale exponentially in d), followed by a classical verification of these samples, using a statistical measure such as linear cross entropy (XEB), which requires classical postprocessing time that is much larger and scales exponentially in n. Moreover in the experiments the depth d is sufficiently large that the output distribution of the ideal random quantum circuit (Fig. 1 (a)) is anti-concentrated², and indeed the output distribution tends to the Porter-Thomas distribution.

Can noise be mitigated?

Short answer: no.

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Not-so-short answer: Error mitigation is impractical (in the worst case)

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Vignette I: Error mitigation can fail badly

Exponentially tighter bounds on limitations of quantum error mitigation

Yihui Quek,¹ Daniel Stilck França,^{2,3,1} Sumeet Khatri,¹ Johannes Jakob Meyer,¹ and Jens Eisert^{1,4}

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Quantum error mitigation has been proposed as a means to combat unwanted and unavoidable errors in near-term quantum computing by classically post-processing outcomes of multiple quantum circuits. It does so in a fashion that requires no or few additional quantum resources, in contrast to fault-tolerant schemes that come along with heavy overheads. Error mitigation leads to noise reduction in small schemes of quantum computation. In this work, however, we identify strong limiuhigh ang prima ngiga gan ha affactivaly (n

arXiv: 2210.11505

In a world with noiseless quantum computers:

In the real world, C is noisy:

- qubit decoherence
- gate errors

Solution 1: **fault-tolerance** (requires lots of machinery – mid-circuit measurements, auxiliary qubits, pumping out entropy)

- qubit decoherence
- gate errors

Zero-noise extrapolation:

1) Run the circuit of interest at amplified noise level $λ$ (call this). C_{λ}

Plot taken from Giurgica-Tiron et al, 2020 IEEE International Conference on Quantum Computing and Engineering (QCE)

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 $E(\lambda) = \text{Tr}(C_{\lambda}(\rho_{\text{in}})O)$

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- 3) Repeat steps 1, 2 for different λ.
- 4) Output the extrapolated value $E(0)$

Plot taken from Giurgica-Tiron et al, 2020 IEEE International Conference on Quantum Computing and Engineering (QCE)

Our question: sample complexity of error mitigation?

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Relevant parameters: n, D (circuit width/depth); p (depolarizing noise strength)

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For depolarizing noise:

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p^{-\Omega(n\,D)} \text{ copies for } \;\; D = \Omega(\log\log(n))
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For non-unital noise (toy model): $c^{-\Omega(n D)}$ copies

Relevant parameters: n, D (circuit width/depth); p (depolarizing noise strength)

Previously proven: $\Omega(\exp(D))$

Our result: $\Omega(\exp(nD)) \to \text{Exponentially stronger}$

Intuition

We show: $exp(\Omega(nD))$ runs of a noisy circuit are required for good error mitigation.

Proof intuition:

1. Depolarizing noise + rapidly mixing circuit drives *every* **input state towards the maximally-mixed state**

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- **1. Depolarizing noise + rapidly mixing circuit drives** *every* **input state towards the maximally-mixed state**
- 2. This makes circuit output states non-distinguishable
- 3. All information lost after circuit \rightarrow cannot error mitigate

The wrong conclusion:

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Let's now look at a different type of noise.

Canonical example: **amplitude damping noise!**

$$
\mathcal{N}(\rho)=\left(\begin{array}{cc}1&0\\0&\sqrt{1-p}\end{array}\right)\rho\left(\begin{array}{cc}1&0\\0&\sqrt{1-p}\end{array}\right)+\left(\begin{array}{cc}0&\sqrt{p}\\0&0\end{array}\right)\rho\left(\begin{array}{cc}0&0\\ \sqrt{p}&0\end{array}\right)
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 $\vert \mathbf{1} \rangle$ $\vert \mathbf{0} \rangle$

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 \mathbf{H}

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Example of amplitude damping: T1 decay, readout error etc. Dominant source of noise in superconducting qubits.

Amplitude-damping noise ~ **reset**-to-all-0s

Non-unital noise: fault-tolerance "for free"?

Quantum refrigerator, Ben-Or, Gottesman, Hassidim (arXiv 1301.1995)

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What happens when both amplitude-damping and depolarizing noise are present?

Depolarizing noise tends to "scramble" the distribution by increasing entropy.

A layer of random gates also tries to"scramble" the distribution!

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Who wins this fight?!

Vignette II: Non-unital noise induces absence of barren plateaus

With Antonio Mele, Armando Angrisani, Soumik Ghosh, Daniel Stilck Franca, Jens Eisert

Barren plateaus are bad for variational algorithms

Gradient vanishes in all directions; can't figure out where to go!

https://www.eurekalert.org/multimedia/739167

Barren plateaus make optimization hard

Math translation of 'vanishing gradient':

$$
\hbox{Var}_{U_1,...U_m}[C]=O(\exp(-n))
$$

C = gradient of cost function

Barren plateaus make optimization hard

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Recall: We are optimizing over $U_1,\ldots U_m$

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Math translation of 'vanishing gradient':

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\hbox{Var}_{U_1,...U_m}[C]=O(\exp(-n))
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Recall: We are optimizing over $U_1,\ldots U_m$

BP implies: if we initialize the optimization randomly, gradient looks the same in all directions.

BP: $Var_{U_1,...U_m}[C] = O(\exp(-n))$

What if the circuit is affected by noise \bullet ?

[1] Barren plateaus in quantum neural network training landscapes. *J. R. McClean et al. Nature Comm. (2018).*

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I.e. we show noise-induced *absence of* barren plateaus! (No guarantees that the minimum is still in the right place, though…)

Intuition Think about $\text{Var}_{\mathcal{C}}[\text{Tr}(Z_1\mathcal{C}(\rho))]$

Will show this is at least some constant.

Intuition Claim: $\mathsf{Var}_{\mathcal{C}}[\mathsf{Tr}(Z_1\mathcal{C}(\rho))] \ge \Omega(1)$

1. In Heisenberg picture:

 $Tr(Z_1C(\rho)) = Tr(C^{\dagger}(Z_1)\rho)$

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- 2. After constant # layers: - noise resets entire lightcone of Z1 to some fixed state with constant probability
	- equivalent to running a constant-depth circuit
- 3. Constant depth circuits have no barren plateaus (not enough randomness)

Who wins this fight: us (with some thought)

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BUT: Non-unital noise decreases entropy. Can we take advantage of this in the intermediate term?

