

Deep quantum neural networks form Gaussian processes

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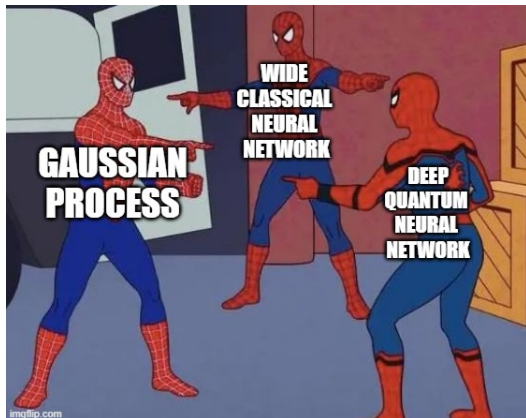
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November 21st

Overview

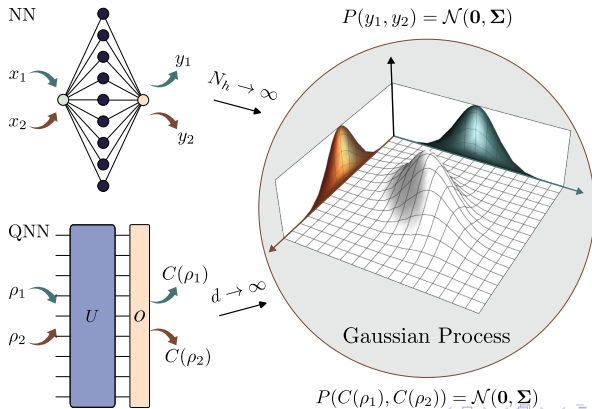
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Gaussian processes

Definition (Gaussian process)

A collection of random variables $\{X_1, X_2, \dots\}$ is a GP if and only if, for every finite set of indices $\{1, 2, \dots, m\}$, the vector (X_1, X_2, \dots, X_m) follows a multivariate Gaussian distribution, $\mathcal{N}(\vec{\mu}, \Sigma)$.



GPs in classical ML

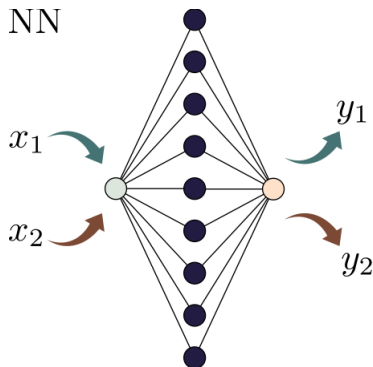
Theorem (Central limit theorem)

Let $\{X_1, X_2, \dots, X_S\}$ be i.i.d. random variables. The sum $X_1 + X_2 + \dots + X_S$ converges to a Gaussian distribution when $S \rightarrow \infty$.

The input of the network is $x \in \mathbb{R}$, and the output is given by

$$f(x) = b + \sum_{l=1}^{N_h} v_l h_l(x)$$

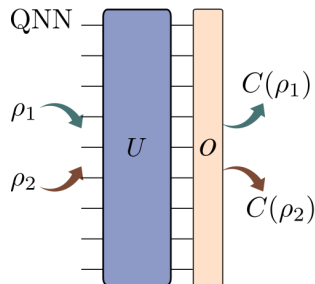
where $h_l(x) = \phi(a_l + u_l x)$ models the action of each neuron in the hidden layer.



GPs in QML

GPs in QML

- Quantum neural network (i.e. variational circuit) U
- Set of initial pure states $\{\rho_i\}_i$
- Traceless observable from a set $\{O_j\}_j$ such that $O_j^2 = \mathbb{1} \quad \forall j$
- Cost function of the form $C_j(\rho_i) = \text{Tr}[U\rho_i U^\dagger O_j]$



$$\mathcal{C} = (C_j(\rho_i), \dots, C_j(\rho_{i'}), \dots) \xrightarrow{d \rightarrow \infty} \text{GP}$$

We cannot leverage the central limit theorem (or its variants) because of the **unitarity constraint**.

$$C_j(\rho_i) = \sum_{k,k',r,r'=1}^d u_{kk'} \rho_{k'r} u_{r'r}^* o_{r'k}$$

The entries of Haar random unitaries over the unitary group $\mathbb{U}(d)$ are **i.d.**, but **correlated** as $\frac{-1}{d-1}$ (same row) or $\frac{1}{(d-1)^2}$ (different row).

Strategy

- i) Computing all the moments over the Haar measure of certain groups (unitary and orthogonal).

- ii) Showing they match those of a multivariate Gaussian.

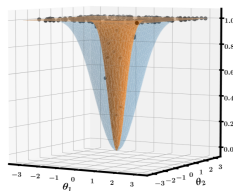
First moments

$$\mathbb{E}_{\mathbb{U}(d)}[C_j(\rho_i)] = \mathbb{E}_{\mathbb{O}(d)}[C_j(\rho_i)] = 0$$

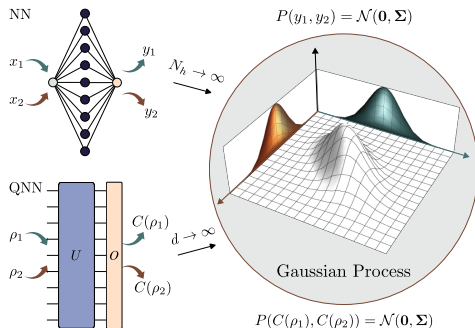
Second moments

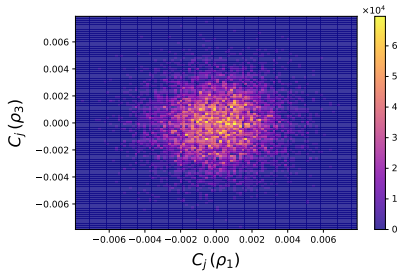
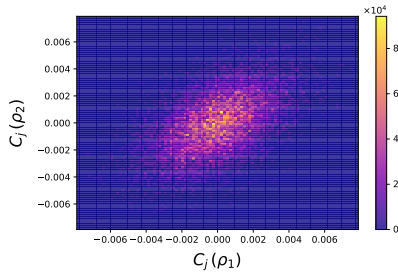
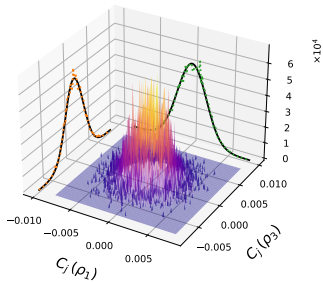
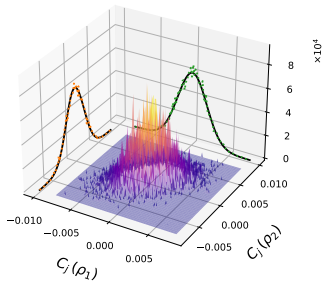
$$\Sigma_{i,i'}^{\mathbb{U}} = \frac{d}{d^2 - 1} \left(\text{Tr}[\rho_i \rho_{i'}] - \frac{1}{d} \right)$$

$$\Sigma_{i,i'}^{\mathbb{O}} = \frac{2(d+1)}{(d+2)(d-1)} \left(\text{Tr}[\rho_i \rho_{i'}] \left(1 - \frac{1}{d+1} \right) - \frac{1}{d+1} \right)$$



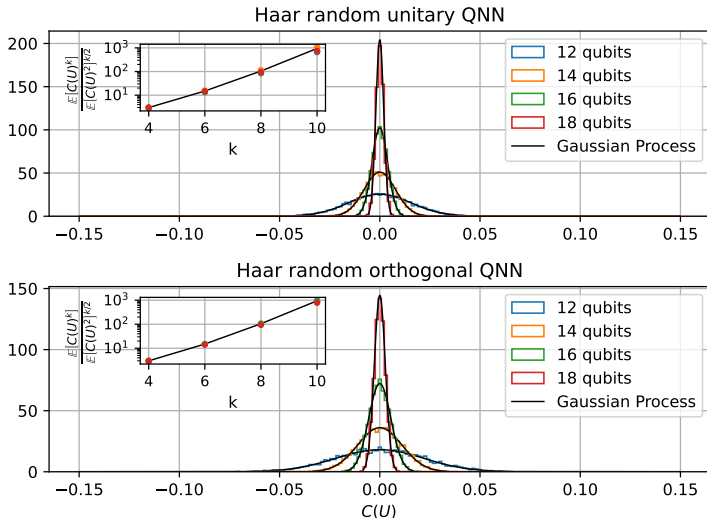
Dataset. For all $\rho_i, \rho_{i'} \in \mathcal{D}$:	GP	Correlation	Statement
$\text{Tr}[\rho_i \rho_{i'}] \in \Omega\left(\frac{1}{\text{poly}(\log(d))}\right)$	Yes	Positive	Theorem 1
$\text{Tr}[\rho_i \rho_{i'}] = \frac{1}{d}$	Yes	Null	Theorem 2
$\text{Tr}[\rho_i \rho_{i'}] = 0$	Yes	Negative	Theorem 3





Corollary

For any ρ_i and O_j , we have $P(C_j(\rho_i)) = \mathcal{N}(0, \sigma^2)$, where $\sigma^2 = \frac{1}{d}, \frac{2}{d}$ when U is Haar random over $\mathbb{U}(d)$ and $\mathbb{O}(d)$, resp.



Corollary

Assuming that there exists a parametrized gate in U of the form $e^{-i\theta H}$ for some Pauli operator H , then

$$P(|C_j(\rho_i)| \geq c), P(|\partial_\theta C_j(\rho_i)| \geq c) \in \mathcal{O}\left(\frac{1}{ce^{dc^2}\sqrt{d}}\right).$$

**Concentration of measure is doubly-exponential
in the number of qudits!**

Corollary

Let U be drawn from a t -design. Then, under the same conditions for which Theorems 1, 2 and 3 hold, the vector \mathcal{C} matches the first t moments of a GP.

An extension of Chebyshev's inequality to higher order moments leads to

$$P(|C_j(\rho_i)| \geq c), P(|\partial_\theta C_j(\rho_i)| \geq c) \in \mathcal{O} \left(\frac{(2 \lfloor \frac{t}{2} \rfloor)!}{2^{\lfloor \frac{t}{2} \rfloor} (\lfloor \frac{t}{2} \rfloor)! (dc^2)^{\lfloor \frac{t}{2} \rfloor}} \right)$$

New and tighter bounds for t -designs.

Theorem

Given the set of observations $y(\rho_1), \dots, y(\rho_m)$ obtained from $N \in \mathcal{O}(\text{poly}(\log d))$ measurements, then the predictive distribution of the GP is trivial:

$$P(C_j(\rho_{m+1}) | C_j(\rho_1), \dots, C_j(\rho_m)) = \mathcal{N}(0, \sigma^2)$$

We cannot use the GPs to efficiently predict new outputs of the QNN using Bayesian statistics.

Proof technique

We need to compute quantities of the form (for arbitrary k)

$$\mathbb{E}_G \left[\text{Tr} \left[U^{\otimes k} \Lambda (U^\dagger)^{\otimes k} O^{\otimes k} \right] \right] = \text{Tr} \left[\mathbb{E}_G \left[U^{\otimes k} \Lambda (U^\dagger)^{\otimes k} \right] O^{\otimes k} \right],$$

where $\Lambda = \rho_{i_1} \otimes \cdots \otimes \rho_{i_k}$ and $G = \mathbb{U}(d), \mathbb{O}(d)$.

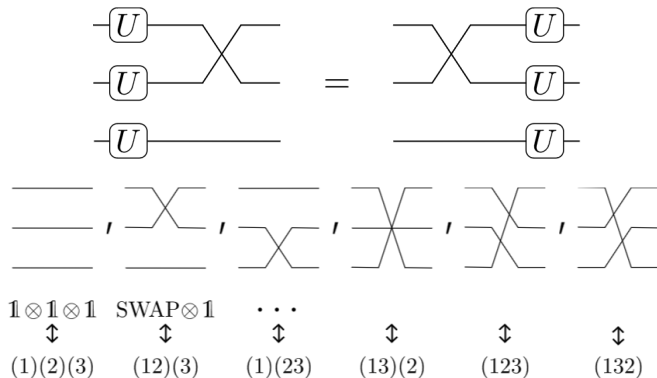
$\mathcal{T}_G^{(k)}[\Lambda] \equiv \int d\mu(U) U^{\otimes k} \Lambda (U^\dagger)^{\otimes k}$ is called the twirl over G .

Recall that

- the moments of $\mathcal{C} = (C_j(\rho_i), \dots, C_j(\rho_{i'}), \dots)$ are of the form
$$\mathbb{E}_{\mathbb{U}(d)} \left[\prod_{\lambda=1}^k C_j(\rho_{i_\lambda}) \right] = \mathbb{E}_{\mathbb{U}(d)} \left[\prod_{\lambda=1}^k \text{Tr}[U \rho_{i_\lambda} U^\dagger O_j] \right]$$
- $\text{Tr}[A \otimes B] = \text{Tr}[A] \text{Tr}[B]$

The twirl belongs to the k -th order commutant.

The k -th order commutant of the tensor representation of $U(d)$ is the set of permutations, i.e. the symmetric group S_k (Schur-Weyl duality)



We can expand the twirl as

$$\mathcal{T}_G^{(k)}[\Lambda] = \sum_{\mu=1}^D c_{\mu}(\Lambda) P_{\mu}, \quad \text{with } P_{\mu} \text{ a basis of } \mathcal{C}^{(k)}(G)$$

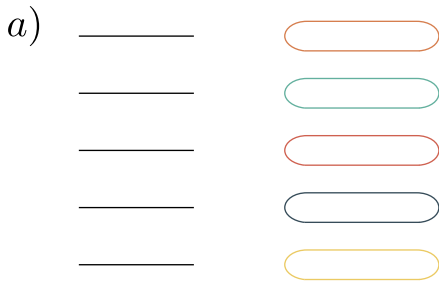
The coefficients $c_{\mu}(\Lambda)$ can be calculated using **Weingarten calculus**, i.e.

$$\vec{c}(\Lambda) = A^{-1} \cdot \vec{b}(\Lambda),$$

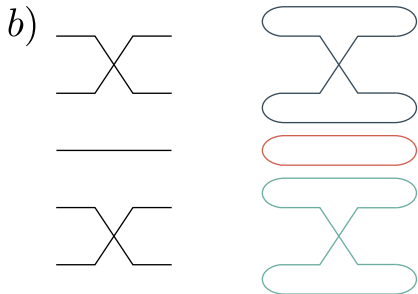
where $\vec{b}(\Lambda) = (\text{Tr}[\Lambda P_1], \dots, \text{Tr}[\Lambda P_D])$ and $(A)_{\nu\mu} = \text{Tr}[P_{\nu} P_{\mu}]$ is the Gram matrix.

$$\text{Tr}[AB] = \text{---} \boxed{A} \text{---} \boxed{B} \text{---}$$

$$\text{Tr}[\mathbf{1}] = \text{---} \text{---} = d$$



$$e = (1)(2)(3)(4)(5) \quad \chi(e) = d^5$$



$$\sigma = (12)(3)(45) \quad \chi(\sigma) = d^3$$

Let the representation of the symmetric group S_k be

$$P_d(\sigma) = \sum_{i_1, \dots, i_k=0}^{d-1} |i_{\sigma^{-1}(1)}, \dots, i_{\sigma^{-1}(k)}\rangle \langle i_1, \dots, i_k|$$

We can prove that

$$\begin{aligned} \mathbb{E}_{\mathbb{U}(d)}[U^{\otimes k} \Lambda(U^\dagger)^{\otimes k}] &= \frac{1}{d^k} \sum_{\sigma \in S_k} \text{Tr}[\Lambda P_d(\sigma)] P_d(\sigma^{-1}) \\ &\quad + \frac{1}{d^k} \sum_{\sigma, \pi \in S_k} c_{\sigma, \pi} \text{Tr}[\Lambda P_d(\sigma)] P_d(\pi), \end{aligned}$$

with $c_{\sigma, \pi} \in \mathcal{O}(\frac{1}{d})$.

The previous implies

$$\begin{aligned}\mathbb{E}_{\mathbf{U}(d)} \left[\text{Tr} \left[\mathbf{U}^{\otimes k} \Lambda(\mathbf{U}^\dagger)^{\otimes k} \mathbf{O}^{\otimes k} \right] \right] &= \frac{1}{d^k} \sum_{\sigma \in \mathcal{S}_k} \text{Tr}[\Lambda P_d(\sigma)] \text{Tr}[P_d(\sigma^{-1}) \mathbf{O}^{\otimes k}] \\ &+ \frac{1}{d^k} \sum_{\sigma, \pi \in \mathcal{S}_k} c_{\sigma, \pi} \text{Tr}[\Lambda P_d(\sigma)] \text{Tr}[P_d(\pi) \mathbf{O}^{\otimes k}]\end{aligned}$$

- $\text{Tr}[P_d(\sigma) \mathbf{O}^{\otimes k}] = 0$ if σ contains a cycle of odd length.
- $\text{Tr}[P_d(\sigma) \mathbf{O}^{\otimes k}] = d^r$ if k is even and σ is a product of r disjoint cycles of even length.
- In particular, $\text{Tr}[P_d(\sigma) \mathbf{O}^{\otimes k}] = d^{k/2}$ if k is even and σ is the product of $k/2$ disjoint transpositions.

We find that the moments of \mathcal{C} are

$$\mathbb{E}_{\mathbf{U}(d)} \left[\text{Tr} \left[U^{\otimes k} \Lambda(U^\dagger)^{\otimes k} O^{\otimes k} \right] \right] = \frac{1}{d^{k/2}} \sum_{\sigma \in \mathcal{T}_k} \prod_{\{t, t'\} \in \sigma} \text{Tr}[\rho_t \rho_{t'}] \quad (1)$$

According to Wick's theorem, the moments of a multivariate Gaussian are

$$\mathbb{E}[C_j(\rho_i), \dots, C_j(\rho_{i'}) \dots] = \sum_{\sigma \in \mathcal{T}_k} \prod_{t, t' \in \sigma} \text{Cov}[C_j(\rho_t) C_j(\rho_{t'})]$$

(Recall that asymptotically $\sum_{i, i'}^{\mathbb{U}} = \frac{\text{Tr}[\rho_i \rho_{i'}]}{d}$)

Conclusions

Conclusions

- Under certain conditions, the output distribution of deep Haar random QNNs converges to a Gaussian process in the limit of large Hilbert space dimension.
- The situation is more nuanced than in classical NNs. We need to make assumptions on the states processed by the QNN, as well as on the measurement operator.
- Our results may be useful in more general settings where Haar random unitaries / t -designs are considered, such as quantum information scramblers and black holes, or many-body physics.



DGM, Martin Larocca and M. Cerezo (2023)

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arXiv preprint: [arXiv.2305.09957](https://arxiv.org/abs/2305.09957)