

# Gibbs Sampling of Periodic Potentials on a Quantum Computer

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## Reference

Gibbs Sampling of Periodic Potentials on a Quantum Computer. arXiv:2210.08104 (2022),  
Joint with Arsalan Motamedi.

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# Energy-based models

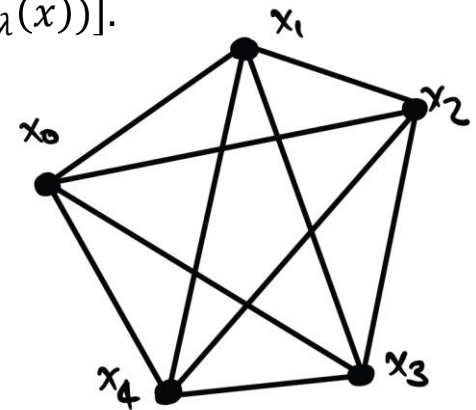
- In an energy-based model the model distribution is a Gibbs distribution  $p_\lambda(x) = \frac{1}{Z_{\lambda,\beta}} e^{-\beta E_\lambda(x)}$ . The model provides a representation of the **energy potential**  $E_\lambda: \Omega \rightarrow \mathbb{R}$ .



- Training:

$$-\nabla_\lambda \mathcal{L}(\lambda) = \int_x \nabla_\lambda \log p_\lambda(x) p(x) dx = \dots = -\beta \mathbb{E}_p[\nabla_\lambda(E_\lambda(x))] + \beta \mathbb{E}_{p_\lambda}[\nabla_\lambda(E_\lambda(x))].$$

- The first term is called the **positive phase** (easy to approximate from data) and the second term is called the **negative phase** (requires samples from the model/Gibbs distribution).
- Old EBMs (e.g., Boltzmann machines) used the Ising potential  $E_\lambda = -\sum_{i,j=1,\dots,m} \lambda_{i,j} \sigma_i \sigma_j$  which didn't work out so great!

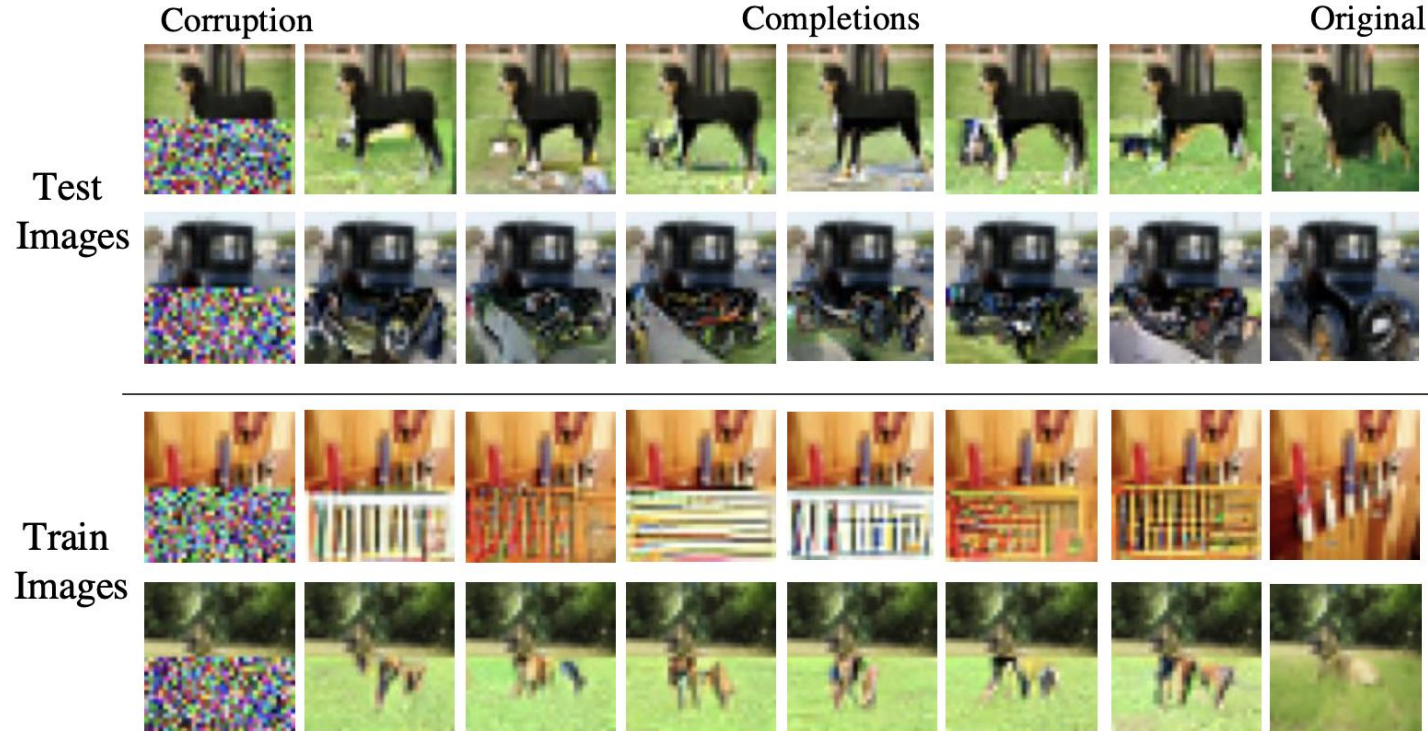


# Modern (deep) EBMs

- Du and Mordatch (2019) showed the first large-scale EBM trainings.
- Langevin dynamics:

$$dX_t = -\frac{1}{2}\beta\sigma^2\nabla_X E_\theta(X) dt + \sigma dW_t.$$

Models	Parameters	Training Time	Sampling Time
EBM	5M	48	3 Hour (Variable)
PixelCNN++	160M	1300	72 Hour
Glow	115M	1300	0.5 Hour
SNGAN	5M	9	0.02 Hour



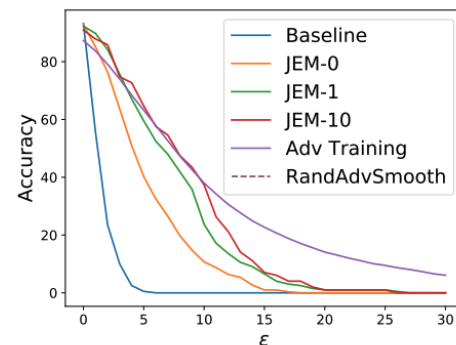
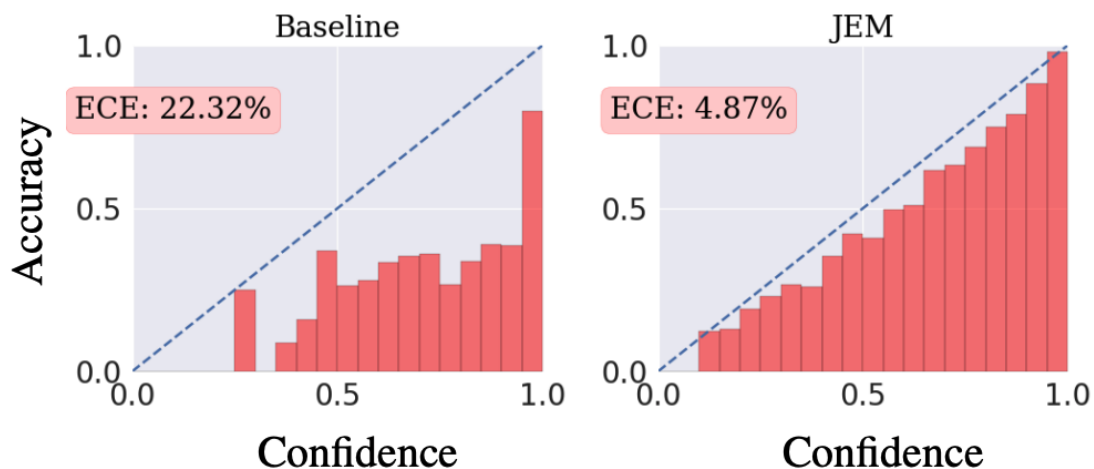
Model	Inception*	FID
<b>CIFAR-10 Unconditional</b>		
PixelCNN [Van Oord et al., 2016]	4.60	65.93
PixelIQN [Ostrovski et al., 2018]	5.29	49.46
EBM (single)	6.02	40.58
DCGAN [Radford et al., 2016]	6.40	37.11
WGAN + GP [Gulrajani et al., 2017]	6.50	36.4
EBM (10 historical ensemble)	6.78	38.2
SNGAN [Miyato et al., 2018]	<b>8.22</b>	21.7
<b>CIFAR-10 Conditional</b>		
Improved GAN	8.09	-
EBM (single)	8.30	37.9
Spectral Normalization GAN	<b>8.59</b>	25.5
<b>ImageNet 32x32 Conditional</b>		
PixelCNN	8.33	33.27
PixelIQN	10.18	22.99
EBM (single)	<b>18.22</b>	<b>14.31</b>
<b>ImageNet 128x128 Conditional</b>		
ACGAN [Odena et al., 2017]	28.5	-
EBM* (single)	28.6	43.7
SNGAN	<b>36.8</b>	<b>27.62</b>

# Benefits of EBMs

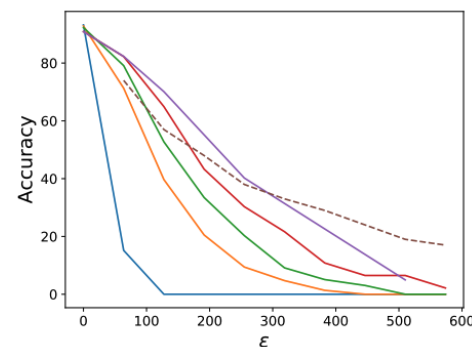
- Guo et al. (2017): Over the years conventional classifiers have become less and less calibrated.
- Grathwohl et al. (2019): Joint EBM improve the calibration and adversarial robustness of the representation.
- Du and Mordatch (2019): Good at OOD detection by EBMs.

Model	PixelCNN++	Glow	EBM (ours)
SVHN	0.32	0.24	<b>0.63</b>
Textures	0.33	0.27	<b>0.48</b>
Constant Uniform	0.0	0.0	<b>0.30</b>
Uniform	1.0	1.0	<b>1.0</b>
CIFAR10 Interpolation	<b>0.71</b>	0.59	0.70
Average	0.47	0.42	<b>0.62</b>

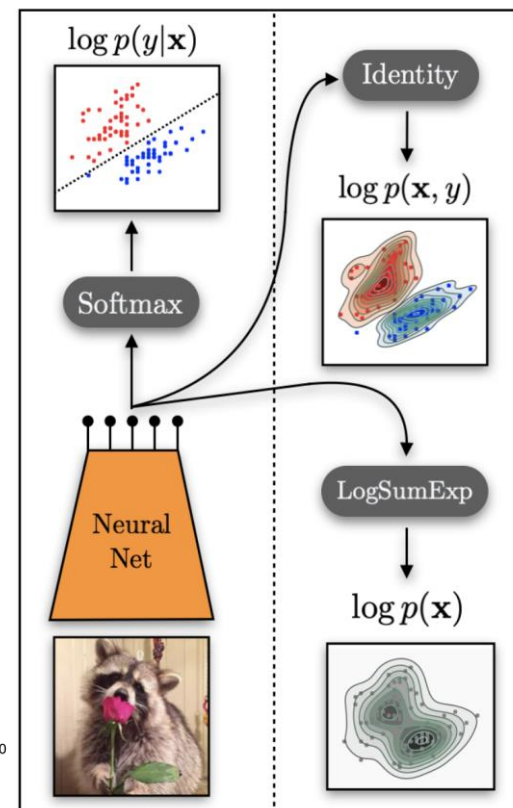
Figure 10: AUROC scores of out of distribution classification on different datasets. Only our model gets better than chance classification.



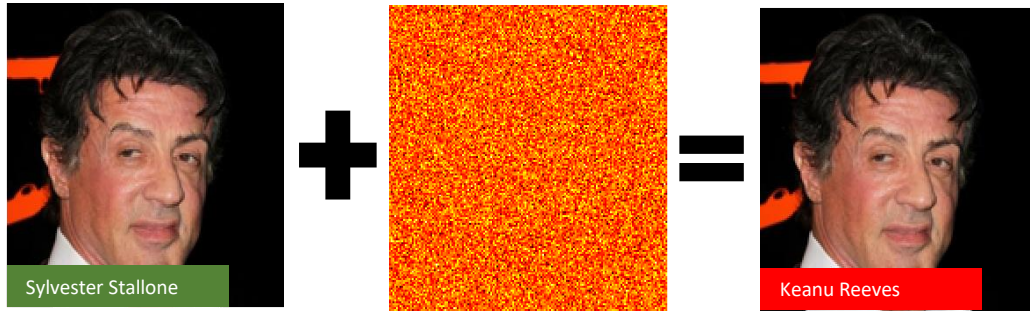
(a)  $L_\infty$  Robustness



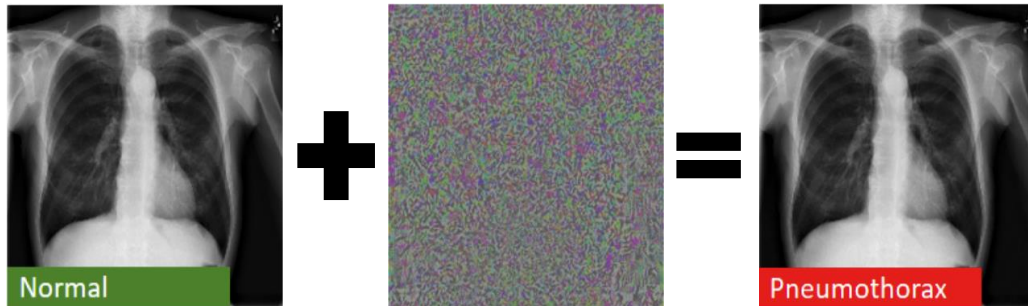
(b)  $L_2$  Robustness



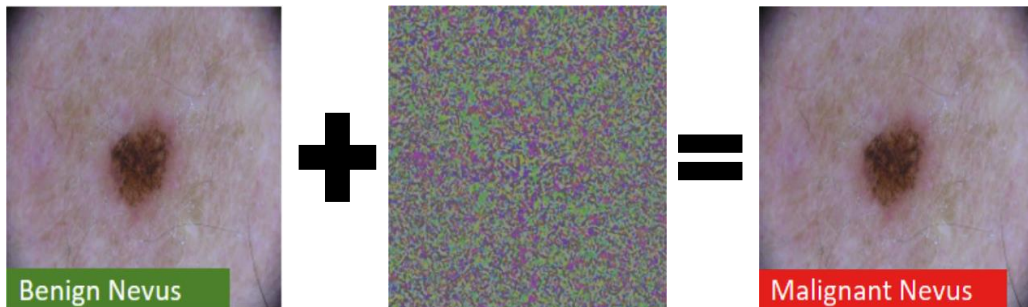
# Stochastic security



Mikhailov, Trusov 2017, blog.ycombinator.com

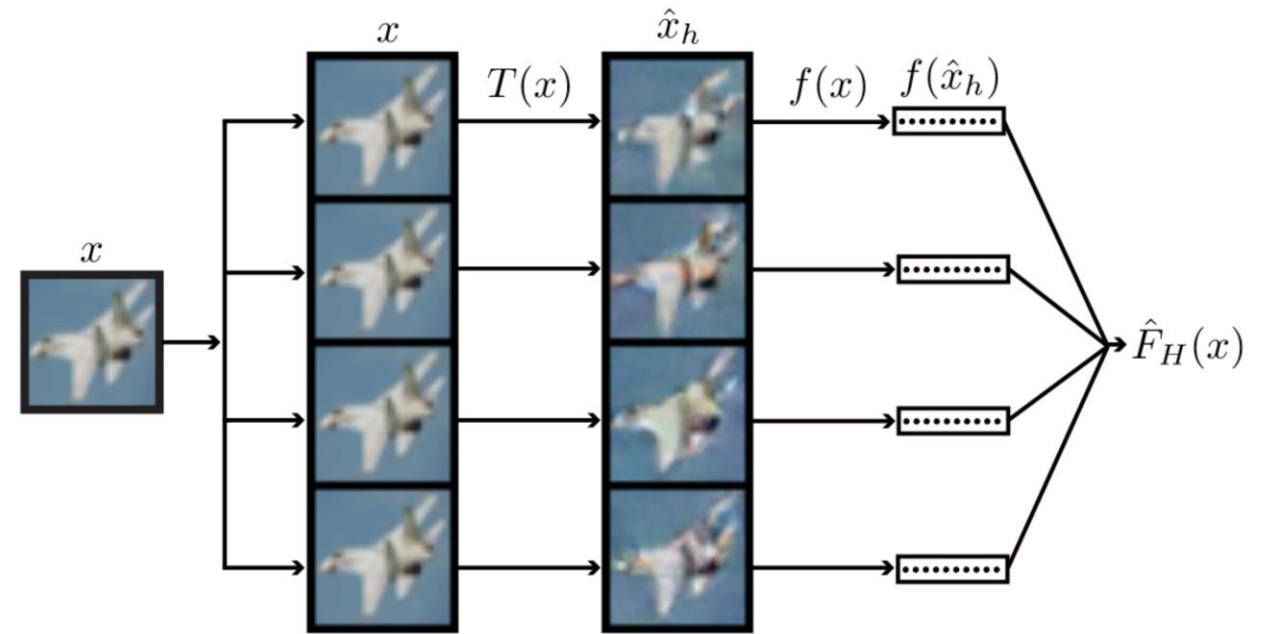


Ma et al. 2020, arXiv:1907.10456v2



Ma et al. 2020, arXiv:1907.10456v2

- Elflein et al., (2021): EBMs outperform normalizing flows in OOD detection.
- Hill-Mitchel-Zho (2021): Convergent EBMs purify an image from adversarial attacks.



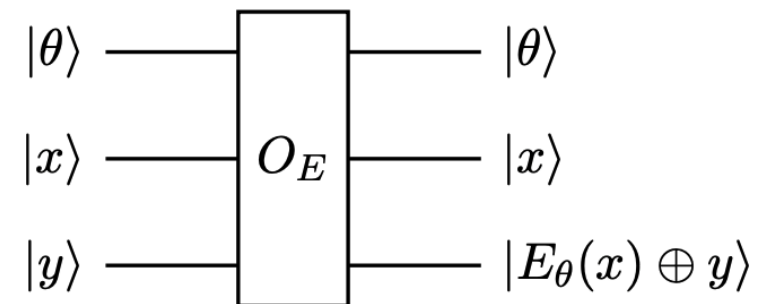
# Our agenda

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- Instead of Monte-Carlo simulation of Langevin dynamics  $dX_t = -\nabla E(X_t)dt + \sqrt{2\beta^{-1}}dW_t \dots$
- ... solve the **Fokker-Planck equation**

$$\frac{\partial \rho}{\partial t} = \nabla \cdot (\nabla E \rho) + \beta^{-1} \Delta \rho.$$

- A second order PDE; a.k.a. the diffusion equation; a.k.a. Kolmogorov's forward equation.
- Our approach: Linearize the generator  $\mathcal{L} = -\nabla E \cdot \nabla + \beta^{-1} \Delta$  and solve using quantum linear system solvers.
- Theorem [Motamedi-R]. A quantum computer can sample from the Gibbs state of an analytic periodic potential using  $O(d^3 e^\beta \text{polylog}(1/\epsilon))$  queries to the energy oracle.
- The training scheme does not require QRAM.
- Non convexity is now a non-issue and the complexity of algorithm depends rather on the rate of decay of Fourier coefficients of  $E$ .



# High precision sampling using Fourier interpolation

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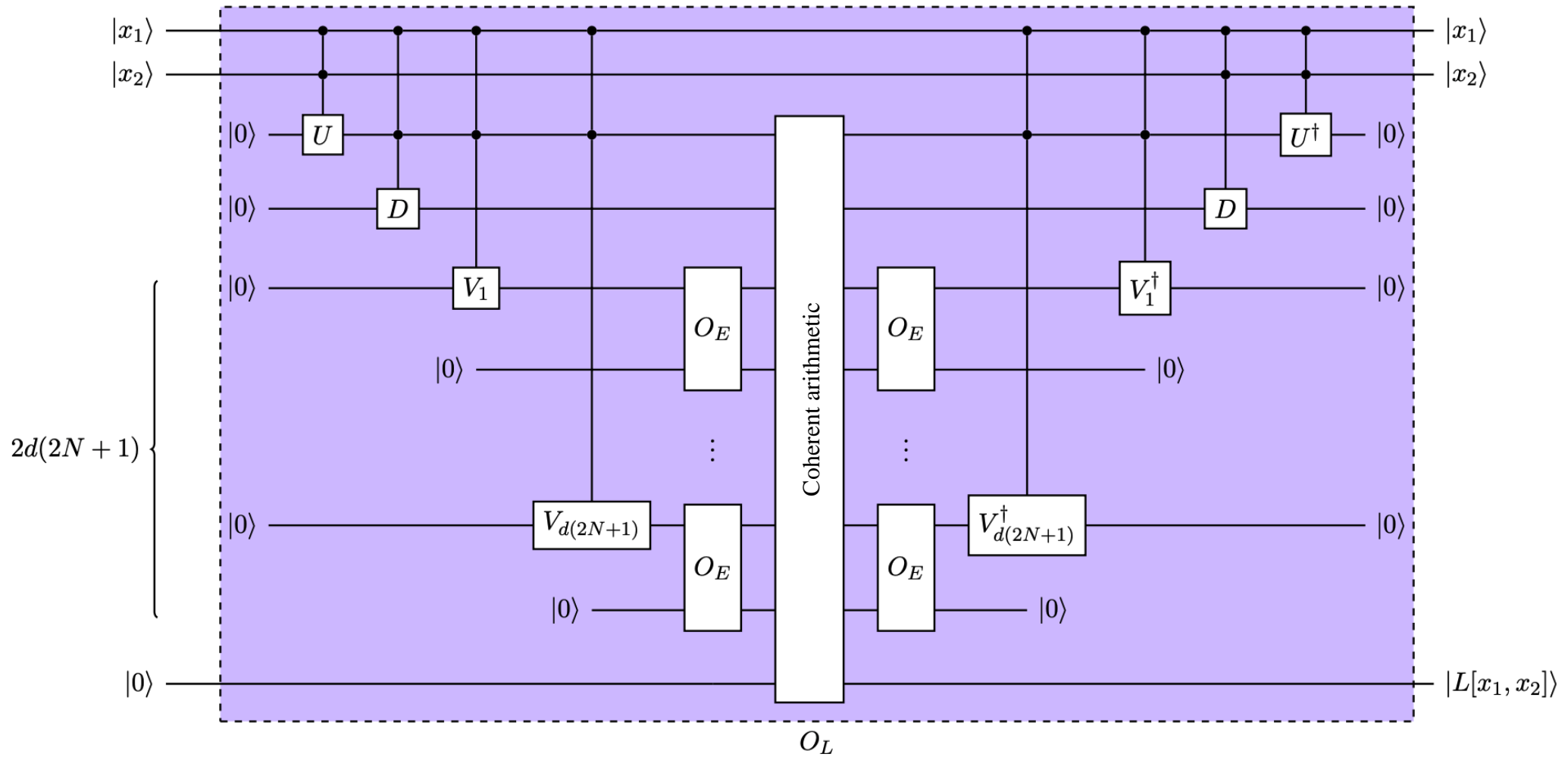
**Input:** Energy function oracle  $O_E$ , lattice parameters  $N, M \in \mathbb{N}$ , solution time  $T > 0$

- 1: Construct an oracle for the discretization  $\mathbb{L}$  of the generator of the Fokker–Planck equation (see Fig. 4 in the appendix).
- 2: Deploy the algorithm of [BCOW17] to prepare a quantum state approximating  $|u(T)\rangle$  pertaining to the solution of  $\frac{d}{dt} \vec{u} = \mathbb{L} \vec{u}$ , with  $\vec{u}(0) = \mathbf{1}$ , at time  $t = T$ .
- 3: Apply the upsampling isometry  $F_M^{-1} \iota F_N$  involving quantum Fourier transformations on the prepared state (Theorem 3.1).
- 4: Measure the resulting state in the computational basis to obtain a lattice point  $x \in [-\ell/2, \ell/2]^d$ .
- 5: Draw a sample  $\tilde{x}$  uniformly at random from the box  $\prod_{i=1}^d \left[ x_i - \frac{\ell}{4M+2}, x_i + \frac{\ell}{4M+2} \right]$  around  $x$ .

**Output:** Sample point  $\tilde{x}$ .

- By writing FPE in the Fourier domain we boost our precision with a technique we call **upsampling**:
- Given a  $(C, a)$ -semi-analytic periodic function  $u$ , an integer  $N \geq 2ad$ , and a quantum state  $|\psi_N\rangle \in H_N$  that is  $\epsilon$ -close to  $|u_N\rangle$  we can sample from  $u^2$  with an  $O(\epsilon, C)$  precision.

# The oracle of the ODE solver





# Long mixing time in equilibrium dynamics

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- Let  $\rho_s$  be the corresponding steady distribution. Further, assume that for a constant  $\lambda > 0$  any differentiable function  $f \in L^2(\rho_s)$  satisfies

$$\text{Var}_{\rho_s}[f] \leq \kappa \mathbb{E}_{\rho_s}[|\nabla f|^2]$$

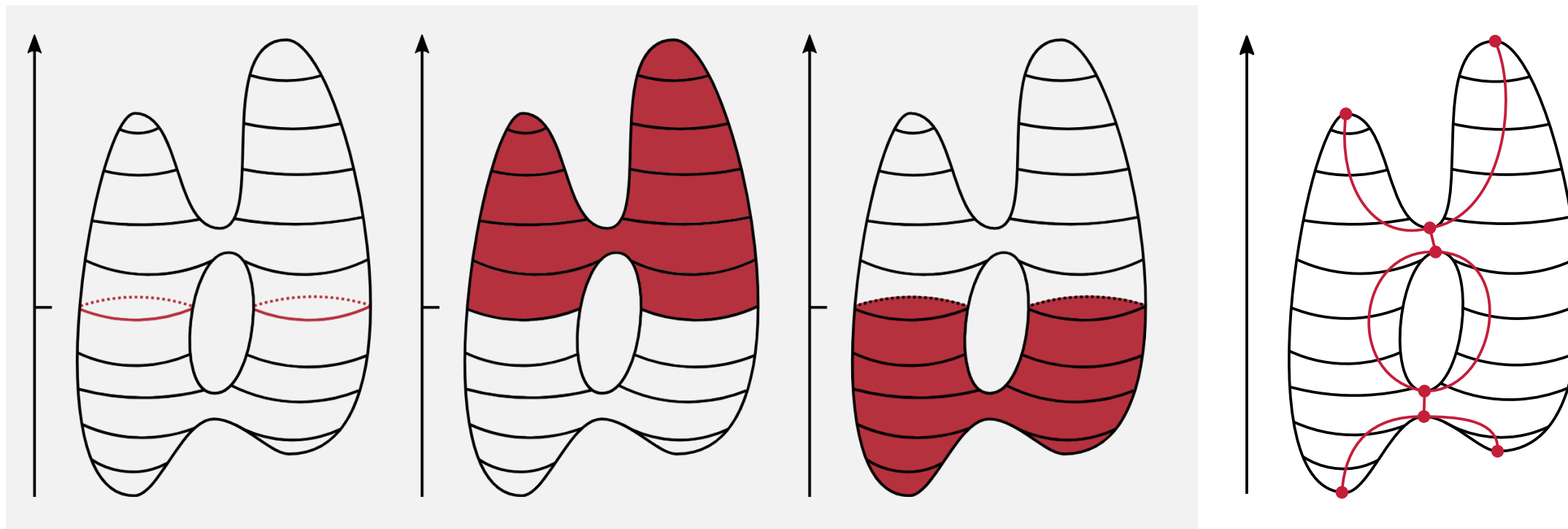
- Then

$$\left\| \frac{\rho_t}{\rho_s} - 1 \right\|_{L^2(\rho_s)} \leq e^{-\frac{t}{\kappa}} \left\| \frac{\rho_0}{\rho_s} - 1 \right\|_{L^2(\rho_s)} .$$

- Note:  $TV(\rho_t, \rho_s) \leq \frac{1}{2} \sqrt{\text{Var}_{\rho_s}(\rho_t/\rho_s)}$ .
- Proposition [Motamedi-R]: Analytic period functions (on a torus) admit a universal Poincare constant  $\kappa = O(e^\Delta)$ .
- So, no speed up in terms of the mixing time. However, in practice  $\Delta$  is bounded.

# Exponential quantum speedups for Morse potentials

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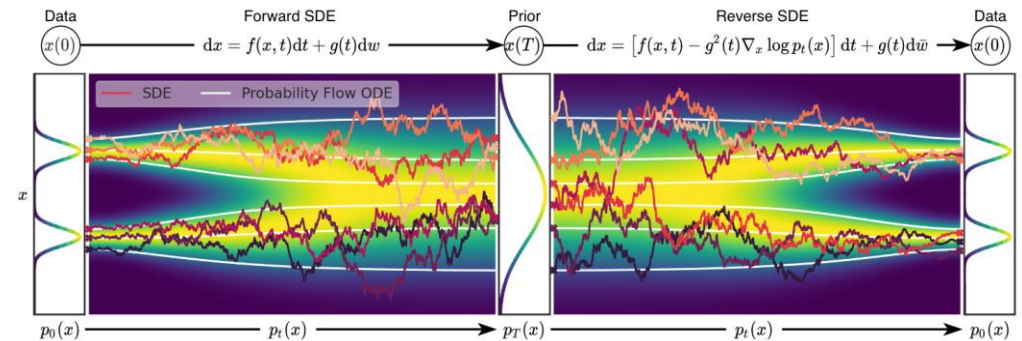
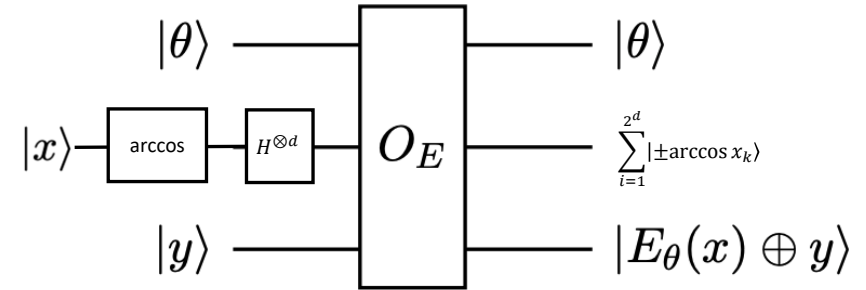
[https://bastian.riek.me/blog/posts/2019/morse\\_theory/](https://bastian.riek.me/blog/posts/2019/morse_theory/)

# Comparison with other algorithms

Method	Potential type	Query order	Sampling complexity	Norm	Mean estimation complexity
This paper	non-convex periodic	zeroeth	$\tilde{O}\left(\kappa_{E/2}e^{\Delta/2}d^7\right)$	TV	$\tilde{O}\left(\kappa_{E/2}e^{\Delta/2}d^7\Delta_f\varepsilon^{-1}\right)$
Rejection sampling	non-convex	zeroeth	$O\left(e^\Delta\right)$	TV	$O\left(e^\Delta\Delta_f^2\varepsilon^{-2}\right)$
This paper	Morse and periodic	zeroeth	$\tilde{O}\left(\lambda^{-2}e^{\Delta/2}d^7\right)$	TV	$\tilde{O}\left(\lambda^{-2}e^{\Delta/2}d^7\Delta_f\varepsilon^{-1}\right)$
Cl. RLA [LE20]	Morse and periodic	first	$\tilde{O}\left(\lambda^{-4}L^4d^3\varepsilon^{-2}\right)$	TV	$\tilde{O}\left(\lambda^{-4}L^4d^3\Delta_f^2\varepsilon^{-4}\right)$
Cl. MRW [DCWY18]	convex	first	$\tilde{O}\left(L^2d^3\varepsilon^{-2}\right)$	TV	$\tilde{O}\left(L^2d^3\Delta_f^2\varepsilon^{-4}\right)$
Q. ULD [CLL <sup>+</sup> 22]	strongly convex	zeroeth	$\tilde{O}\left(\mu^{-2}L^2d^{1/2}\varepsilon^{-1}\right)$	$W_2$	–
Cl. ULD [CCBJ17]	strongly convex	first	$\tilde{O}\left(\mu^{-2}L^2d^{1/2}\varepsilon^{-1}\right)$	$W_2$	–
Q. MALA [CLL <sup>+</sup> 22]	strongly convex	first	$\tilde{O}\left(\mu^{-1/2}L^{1/2}d\right)$	TV	$\tilde{O}\left(\mu^{-1/2}L^{1/2}d\Delta_f\varepsilon^{-1}\right)$
Cl. MALA [LST20]	strongly convex	first	$\tilde{O}\left(\mu^{-1}Ld\right)$	TV	$\tilde{O}\left(\mu^{-1}Ld\Delta_f^2\varepsilon^{-2}\right)$

# Outlook

- Outlook 1: Beyond periodic functions
  - Periodicity does not create a constrain in data-driven training.
  - Direct sampling from  $[-1, 1]^d$  can be achieved by moving to Chebychev spectral methods (WIP).
- Outlook 2:
  - In diffusion models the diffusion equation is solved for a finite time avoiding the long mixing time.
  - This requires solving FPEs with time-dependent drift and diffusion terms (WIP).
- Outlook 3:
  - Digital computation is very expensive for FTQC. Perhaps all this machinery should rather be developed in the bosonic setting (Please talk to me if interested...).



## Quantum Physics

[Submitted on 26 Feb 2000 (v1), last revised 7 Jun 2000 (this version, v2)]

## Quantum searching with continuous variables

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