Hadronic cross sections measurement for the muon *g-2* **calculation**

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22nd Conference on Flavor Physics and CP Violation **(FPCP 2024)** **Bangkok, Thailand - May 27-31, 2024**

The anomalous magnetic moment of the leptons The anomalous magnetic moment of the leptons

$$
\vec{\mu} = g \frac{e}{2m} \vec{s}, \quad a = \frac{g-2}{2}
$$

- **(1928) pointlike Dirac particles:** $g=2$ **,** $a=f$
- (1948) **anomaly discovered for the electro** – *ae* **exp= (1.19** ± **0.05) 10**−**3** (Kusch-Foley) ϵ (1048) enemaly discovered for the ϵ $e^{i\theta} = (1,10,0.05), 10^{-3}$ (Kucob I
	- (1948) **explained by** $O(\alpha)$ **QED correction** e^{ϵ} (1048) explained by $O(\epsilon)$ OFD corre
- $a_e^{\text{th}} = \alpha/2\pi = 1.16 \ 10^{-3}$ (Schwinger) (1948) explained by $O(n)$ QED corrections u_e $\omega = 1.16$ 10 (SCHW)
- First triumph of QED! $\overline{}$ • First triumph of Q

 a_l sensitive to quantum fluctuations, not only from QED. ==> Must include all contributions for a precise calculation

Leading Order HVP calculations

- Main contribution to a_u from very low energy region
	- $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ provides ~73% of the total
- Uncertainty from multi-hadronic channels significant
- Strategy to calculate the total cross section:
	- Sum up exclusive cross sections up to \sqrt{s} \sim 2 GeV
	- Estimate the remaining (*e.g.* by using isospin relations)
	- Use inclusive cross sections measurements and/or pQCD above $\sqrt{s} \sim 2$ GeV

The impact of CMD-3 on SM prediction of $\mathbf{D_{r}}$ of a hadred **Present status**

FNAL (Run2+3): <u>[Phys.Rev.Lett. 131 \(2023\) 161802](https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.131.161802)</u> • **Preliminary results on** $e^+e^ \overline{P}$ BNL: <u>[Phys. Rev. D. 73 \(2006\) 072003](https://journals.aps.org/prd/abstract/10.1103/PhysRevD.73.072003)</u> FNAL (Run1): [Phys. RevL ett. 126 \(2021\) 141801](https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.126.141801) BMW20 (lattice): [Nature 593 \(2021\) 51](https://www.nature.com/articles/s41586-021-03418-1) CMD-3: [arXiv:2303.08834 \(2023\)](https://arxiv.org/abs/2302.08834)

Muon g-2 Theory Initiative White Paper: T. Toyama *et al.***, [Phys.Rep. 887 \(2020\) 1](https://www.sciencedirect.com/science/article/pii/S0370157320302556)** Using only 2π from CMD-3 million construction of the CMD-3 million construction of the CMD-3 million construction *a_µSM* = **(116 591 810** ± 43) \cdot **10**⁻¹¹

The CMD-3 is one now one now over many or now over many or now over many or now over \sim

Experiment (BNL + FNAL Run 1-3) $a_{\mu}^{\text{exp}} = (116592059 \pm 22) \times 10^{-11} (0.19 \text{ ppm})$

Theory-experiment deviation at 5 sigma! But... After WP2020 publication:

- **• First lattice calculation with small uncertainties (BMW Collaboration)**
- **Preliminary results on** $e^+e^- \rightarrow \pi^+\pi^-$ by CMD-3 **at VEPP2000**
- Both in tension with previous SM results from *e*⁺*e*[−] data

Hadronic cross sections measurement

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How to measure σ*had*

"Conventional" method: Energy Scan

- σ*had* measured varying the beam energies within the accessible range ==> $\sqrt{s} = 2E_{beam}$
- Well-defined center-of-mass energy
- Low background
- Very good energy resolution
	- precise study of narrow resonances
- Systematics from point-to-point normalization
- Limited energy range covered
- $\cdot \varepsilon \rightarrow 0$ approaching processes thresholds (zeromomentum tracks)

hadrons $e^ \longrightarrow$ m_{max} *m2=s*′=*s*(*1-x*) f *e*+*e*[−] ➝ γISR *+ hadrons* $x = \frac{E_{\gamma}}{Z}$ *s*

"Novel method": Radiative Return

- Developed by KLOE and *BABAR*
- Effective c.m. energy: $\sqrt{s'} = \sqrt{s(1-x)}$
- Continuous and wide spectrum of energies √s*′* below the nominal \sqrt{s}
- Uniform data quality all-over energy range
- Boost of hadronic system: ε ≠ *0* at threshold
- Higher (and different) background sources \triangleright main backgrounds from different ISR processes and $e^+e^- \rightarrow q\bar{q}$ production
- Limited mass resolution (\sim few MeV)

SND and CMD-2 @VEPP-2M:

- √s <1.4 GeV
- ⁻ $π$ ⁺ $π$ [−] and several multi-hadron channels

SND and CMD-3 @VEPP-2000:

- $-0.32 \leq \sqrt{s} \leq 2.0$ GeV
- $-L_{int}$ = > 1 fb⁻¹; ~65 pb⁻¹ in the ρ region
- ⁻ $π⁺π[−]$ and many multi-hadron channels

KEDR @VEPP-4M:

- 2 $\leq \sqrt{s}$ < 4 GeV
- inclusive cross section

Older experiments:

- DM1, DM2, FENICE, BES, $BESII,...$ **BESII,... BESII,... BESII,...**

The players Energy scan
Energy scan ISR

BABAR **@PEP-II:**

- $\sqrt{s} = 10.6$ GeV (Y(4S) peak)
- $-$ *L*_{int} = 470 fb⁻¹; 232 fb⁻¹ used for e^+e^- → $\pi^+\pi^-$
- Covered energy range: 0.3 4.5 GeV
	- >50 final states measured

KLOE @DAΦNE:

- $\sim \sqrt{s} = 1.02 \text{ GeV} (\phi(1020) \text{ peak})$
- $-$ *L*_{int} = ∼2 fb⁻¹; 240 pb⁻¹ used for e^+e^- → $\pi^+\pi^-$
- Covered energy range: 0.3 0.9 GeV

BES III @BEPC-II:

- \sqrt{s} = 2 4.7 GeV
- main ISR results from data at ψ (3770)
- $-$ *L*_{int}[ψ(3770)] = ~20 fb⁻¹; 2.93 fb⁻¹ used for e^+e^- → $\pi^+\pi^-$
- *−* ISR: 0.6 *−* 0.9 GeV (π ⁺π[−])
- Scan: 2 3.8 GeV (spectroscopy and inclusive cross section)

CLEO-c data

- \sqrt{s} = 3.77 & 4.17GeV
- $L_{int} = \sim 1.4$ fb⁻¹ used for $e^+e^- \rightarrow \pi^+\pi^-$

- $\sqrt{s} = 10.6$ GeV (Y(4S) peak)
- ISR program in progress

e+*e*[−] ➝ π+π−

- Channel dominated by the huge "ρ" peak.
- Cross section must be measured at few per mil level in the ρ region
- Accuracy of all measurements limited by systematic uncertainties
- Several measurements with sub percent error available, but tension observed among some of them

$e^+e^- \rightarrow \pi^+\pi^-$: *BABAR*

- Data taken at Y(4S) peak. Tagged analysis (require $E_{\gamma ISR} > 3$ GeV).
- **Measure the ratio** $\frac{\partial (e^+e^- \rightarrow \mu^+ \mu^- (\gamma))}{\partial (e^+e^- \rightarrow \mu^+ \mu^- (\gamma))}$ = Results rather insensitive to details of MC generator $\sigma(e^+e^- \to \pi^+\pi^-(\gamma))$ $\sigma(e^+e^- \to \mu^+\mu^-(\gamma))$
- Kinematic fit to data allows for one extra photon => measurement at NLO:
- Systematic uncertainty at the ρ peak at 0.5% level, dominated by Particle-ID effects

- data set (232 fb-1).
	- \sim 500M selected π + π pairs
- It is the only experiment to cover the cross section from threshold up to \sim 3 GeV
- New measurement with different analysis strategy in progress.

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CMD-3: new *e*+*e*[−] ➝ π+π−

Three VEPP-2000 data taking periods used:

- **RHO2013** mainly *ρ*-peak $L = 17.8$ pb⁻¹
- **RHO2018** mainly *ρ*-peak $L = 45.4$ pb-1
- **LOW2020** \sqrt{s} < 0.6 GeV $L = 1$ pb⁻¹
- additional data taken after 2021 at $\sqrt{s} > 1.2$ GeV
- Significant accelerator and detector (new drift chamber) upgrade after RHO2013

Total uncertainties in the ρ -peak region:

- 0.7% (RHO2018) / 0.9% (RHO2013)

	uncertainty
Radiative corrections	0.3%
ee/ $\mu\mu/\pi\pi$ separation	0.2%
Fiducial volume	$\rightarrow 0.5\% / 0.8\% (2013)$
Correlated inefficiency	0.1%
Trigger	0.05%
Beam energy	0.1%
Bremsstrahlung loss	0.05%
Pion specific loss	0.2% nuclear interaction
	0.1% pion decay
	0.7% / 0.9% (2013)

Measured Pion FF

 $N_{\pi\pi} = 34 \times 10^6$ **}** $N_{\pi\pi} = 34 \times 10^6$
==> x30 CMD-2

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CMD-3 *e*+*e*[−] ➝ π+π− **: Results**

CMD-3 *vs* CMD-3 *vs* energy • $|F_\pi|^2$ (and $a_\mu^{\pi\pi}$) significantly scan measurements **ISR** measurements
CMD-3 syst. uncert. higher than all previous results! CMD-3 syst. uncert. **0.2** ្_{រ្} 0.2
្ត្⁸0.15 **0.2** $\frac{1}{11}$ 50 **2-1 BABAR 2-1** $-CMD-2$ $-$ CMD-2 $\mathbf{e}^{\mathbf{\frac{5}{3}}}_{\mathbf{e}}$ 0.15 **|CMD3 fit BESIII BABAR** CLEO₁₈ **0.15 SND** -45 **0.15** K CLOE10 SND **←** SND2k **|F| |F|** K LOE12 SND_{2k} **0.1 0.1 0.1** $+40$ **BES** KLOE10 **0.05 0.05** KLOE12 35 CMD3 **0 0** +− ((1020)): PLB 779 −**0.05** −**0.05** −**0.05** −**0.1** −**0.1** −**0.1** $\frac{2}{2}$ −**0.15** −**0.15** −**0.15** 0.82 0.84 0.7 0.72 0.74 0.76 0.78 0.8 −**0.2** s. GeV −**0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1 1.1 1.2 11 December 2023 Roma Tre g-2 Seminar 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1 1.1 1.2 s, GeV** \sim \sim $a_{\mu}^{\pi\pi,LO}$ [$0.6 < \sqrt{s} < 0.88$ GeV] Figure 34: The relative di↵erence of the ISR mea- (1000) $F = \frac{1}{2}$ • **CMD-3:** $a_{\mu}^{\pi\pi,LO} = (379.35 \pm 0.30 \pm 2.95)10^{-10}$ $35 + 0.30 + 2.9511$ surements to the CMD-3 fit. **e** CMD-3: $a^{\pi\pi, LO} = (3$ **before CMD2** 3+3−0: PLB 792 (2019), **CMD2** before CMD2 368*.*8 *±* 10*.*3 generally shows larger pion form factor in the whole energy range under discussion. The whole energy range under \mathbb{R} **SND** CMD2 366*.*5 *±* 3*.*4 t would explain the difference **KLOE comb** $\frac{1}{2}$ = $\frac{1}{2}$ and $\frac{1}{2}$ including previous CMD-2 including • No clear smoking gun that would explain the difference measurement, is observed at the left side of **at the left side of ps** $\frac{1}{2}$, where it reaches i **BABAR** KLOE 360*.*6 *±* 2*.*1 **• A** new scan in the *ρ* region is started with the aim to **BES** +: submitted to PLB $\mathbb{R} \rightarrow \mathbb{R}$ is the new statistical errors of the new scan in the pregion is started with the aim to **CLEO** $r = \frac{1}{2}$ $r = \frac$ **SND2k** $\frac{1}{2}$ One of the main applications for measured *e*⁺*e* ! *hadrons* cross sections is the evalu-One of the main applications for measured *e*⁺*e* ! *hadrons* cross sections is the evalu- $\begin{array}{c} \begin{array}{c} \begin{array}{c} \end{array} \\ \end{array}$ $\begin{array}{c} \end{array}$ $\begin{array}{c} \end{array}$ $\begin{array}{c} \end{array}$ $\begin{array}{c} \end{array}$ \end{array} $\begin{array}{c} \end{array}$ $\begin{array}{c} \end{array}$ $\begin{array}{c} \end{array}$ huge amount of data for better understanding of **CMD3** systematic effects **is calculated via the dispersive integral calculated via the dispersive integral** $a_{\mu}^{\pi^+\pi^-}$ **(** $0.6 <$ $\sqrt{s} < 0.88$ **GeV), 10⁻¹⁰**

s, GeV

arXiv:2302.08834 (2023)

Exclusive cross sections other than $π$ ⁺

- Main contributions from **3π, 2K and 4π final states**
- *BABAR*: almost complete set of measurements of final states with 2 to 7 hadrons $(\pi, K \text{ and } \eta's)$
- VEPP-2M: \sqrt{s} < 1.4 GeV and relatively low statistics
- VEPP-2000: \sqrt{s} < 2 GeV with much higher luminosity. Many new results coming
- BES-III measured a few channels included in the WP2020. More results on the way[†]

Contribution to $a_{\mu}^{HVP,LO}$ (\sqrt{s} < 1.8 GeV) from DHMZ19 [Phys. J. C 80 \(2020\) 410](http://dx.doi.org/10.1140/epjc/s10052-020-7792-2)

 $\mathbf{F} = \mathbf{F} \cdot \mathbf{V} + \mathbf{V}$, SND \mathbf{F} $e^+e^- \rightarrow K^+K^-$

Recent $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ results

- Second largest contribution to $a_{\mu}^{had,LO}$ and its uncertainty
- Dominated at low energies by the ω and φ resonances, and then by ω recurrences
- Old data:
	- Precise SND and CMD-2 @VEPP-2M below 1.1 GeV
	- *BABAR* (80 fb⁻¹) above the φ
- Five new analyses, not in the WP2020, recently produced:
- *BABAR*: ISR@Y(4S) $0.62 < m_{3\pi} < 3.5$ GeV [PRD 104 (2021) 112003]
- **BESIII**: ISR@ψ(3770) 0.7 < $m_{3π}$ < 3 GeV [arXiv:1912.11208]
- **SND**: scan $1.2 < \sqrt{s} < 2$ GeV [EPJC80, 993 (2020)]
- **BESIII**: scan $2 < \sqrt{s} < 3.08$ GeV [arXiv2401.14711]
- **CMD-3**: scan $0.66 \leq \sqrt{s} \leq 0.97$ GeV [preliminary]

BABAR analysis \sim 469 fb⁻¹

- tagged ISR: 2.5 MeV bin-width at the resonances
- Syst. uncert. at the ω and ϕ peaks ~1.3%.
- Precision on $a_{\mu}^{3\pi}$ improved by a factor \sim 2

(for $m_{3\pi} < 2 \text{ GeV}$):

 $a_{\mu}^{3\pi} = (45.86 \pm 0.14 \pm 0.58) 10^{-10}$

BESIII: Inclusive measurement DEDITII: INCRUSTVC INCRUSH CHICHC accuracy of the SM predictions of α \mathbf{m}

Inclusive cross sections above 2 GeV

- Previous results:
- KEDR
	- $1.84 < \sqrt{s} < 3.72$ [PLB 770 (2017) 174] [PLB 788 (2019) 42]
	- tot. uncert. $2.6\% > 4\%$
	- Agree well with pQCD in the region between $\frac{1}{2}$ (by 1.9 σ) data and nOCD pred 2 and 3.8 GeV $\left(0, 1.50\right)$ data and $p \ge 25$ from
- ections above 2 GeV New results from BESIII (not in the WP)
	- 14 points: $2.23 < \sqrt{s} < 3.67$ GeV
	- Precision slightly better than KEDR
- t_0 , t_2 [PLB $t/0$ (2017) 174]
• dominated by systematic uncertainties n
terrtainties
- $r_{\text{e}} > 4\%$ \rightarrow 2% -> 3% from 2.2 to 3.7 GeV GeV
	- Measured cross section somewhat above KEDR (by 1.9 σ) data and pQCD predictions (by 2.7 σ)

Measurement of additional radiation in ISR processes

Study of $e^+e^- \to \mu^+\mu^-\gamma_{ISR}$ and $e^+e^- \to \pi^+\pi^-\gamma_{ISR}$ with 1 (NLO) or 2 (NNLO) additional hard photons • The full BaBar data sample $(\sim 468 \text{ fb}^{-1})$ is used. Data compared with MC generators:

- **AfkQed** : up to NNLO additional ISR collinear to beams, FSR from PHOTOS
- **Phokara 9.1**: full matrix elements at NLO (no NNLO)
- **Study ISR Processes e+e−→**!**+**!**−/**"**+**"**−**#**ISR with 1 or 2 more Hard Photons** NLO event topologies se concidential several fits consistent with the event topology ● All events are subjected to $\frac{1}{2}$ ^{NSR} \mathcal{S}
- r_{th} α notages determine the category of that • The fit with the best χ^2 event
- gie (LA) γ s detected $\Omega_{\rm P}$ $(S \Delta)$ ν 's ass • Small-angle (SA) $γ$'s assumed collinear with the beams. Large-angle (LA) γ's detected

Two NLO fits:

- $\gamma_{ISR}\gamma_{SA} (E^*_{\gamma_{SA}} > 200 \,\text{MeV})$
- $\gamma_{ISR} \gamma_{LA} (E_{\gamma_{LA}} > 200 \text{ MeV})$

Three categories:

- NLO LA
- NLO SA
- LO: events with no γ 's above threshold

Three NNLO fits and three categories:

- *γ_{ISR}γ*_{SA}*γ*_{SA} -> NNLO 2SA
- *γ_{ISR}γ_{SA}γ*_{LA} -> NNLO SA+LA
- *γ_{ISR}γ*_{LA} γ _{LA} -> NNLO 2LA

Executive 15 μ ➣ ISR photon candidate: . D108 (202. ➣ Largest E*, with E* > 4 GeV **BABAR: Phys. Rev. D108 (2023) 111103** ➣ ISR photon candidate:

 M_{\odot} - Hadronic cross sections M_{\odot} \rightarrow FFCF 2024 Fabio Anulli - Hadronic cross sections — FPCP 2024

➣ ": 0.4−2.45 rad

NLO fits results NLO fits rest \sim \sim \sim

- Similar results for $\mu\mu$ and $\pi\pi$ samples
	- suppressed by a BDT-based selection 2 other multihadronic processes, largely • π *π* affected by larger background due to $\frac{1}{\pi}$

 $\begin{array}{l} \delta \mathrm{R} = \mathrm{SA}\text{-}\mathrm{ISR} + \ \delta \mathrm{A} = 2\mathrm{LA}\text{-}\mathrm{ISR} \end{array}$ $\frac{1}{\sqrt{2}}$ (a), $\frac{1}{\sqrt{2}}$ (b) of the 2 (a), (b) of the ISR2SA kinematic fit for the muon (right) samples. The mu

gies. The same-beam NNLO component in the measured in the measured in the measured in the measured in the measured

 \overline{Q} , (d) \overline{Q} of the corresponding of the more (see Full dots) \overline{Q} ● AfkQed reproduces well the additional SA SA *<* 2*.*3 GeV. photons energy distributions

Consequences for the $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ measurements

- *BABAR* analyses essentially unaffected:
	- loose selection, $\pi\pi/\mu\mu$ ratio, efficiencies from data \Rightarrow no particular dependence from any MC
	- Phokara used for acceptance corrections \Rightarrow Estimated effect $(0.03\pm0.01)\%$ well below the quoted syst. uncertainty of $\sim 0.5\%$
	- NLO already included in 2009 analysis (one extra photon, SA or LA) allowed
	- New analysis in progress will allow for two extra photons (NNLO)
- Other ISR results (**KLOE**, **BESIII**), relying on Phokara, might be affected by:
	- Missing NNLO and higher order contributions
	- Too large MC SA hard-photon rate
	- Recent DHMZ study with fast simulation to estimate such effects [arXiv:2312.02053]
	- A precise estimate of such effects could/should be performed by the experiment themselves

What's in the future

New π *π* analysis on full data

20 fb-1 total collected data at \sqrt{s} =3770 MeV (\sim x6 previous

Expected improved ISR analysis

sample

BESTIT

analyses).

on many final states

- The $\pi\pi$ channel is still the major source of uncertainty
	- CMD-3 result increases the tension
	- Progress in the theoretical description of radiative corrections is necessary to reach few per mil precision

1fb⁻¹ collected above 1 GeV **For analysis on IUII** data New $\pi \pi$
set in progress => ~x10 data with diff = > many analysis in progress on multi hadron production $\overline{110}$ co $200 - 10$ 1 fb⁻¹ collected above 1 GeV

1 fb⁻¹ planned in the ρ region $+1$ fb⁻¹ planned in the ρ region

 \equiv Future $\pi \pi$ analyses with much improved control of systematic effects rovea co

• Impact of other final states becoming very important.

- Goal: measure the main channels at 1-3% total uncertainty
- Sophisticated amplitude analyses needed in most cases
- KEDR @VEPP-4M, BESIII and BelleII may provide new inclusive measurement at $\sqrt{s} > 2$ GeV

New $\pi\pi$ analysis on full data set, with different technique, in progress \Rightarrow Effectively \sim x7 data sample

ISR program, analogous to BABAR, already started

New *BABAR* $e^+e^- \rightarrow \gamma_{ISR}\pi^+\pi^-(\gamma\gamma)$

 n^2 1.25-1.30

- Run 1-4 (232 fb^{-1})
- Main systematics from Particle Identification (PID) related effects $\frac{1}{100}$ and $\frac{1}{100}$ momentum tracking (MC).
- Track momentum selection: p>1 GeV/*c* �/� separation using particle identification (PID), one of dominant systematics,
- \bullet Total uncert. $[0.6 0.9 \text{GeV}] = 0.5\%$ m \mathbb{M}_{q_0} \mathbb{M}_{q_1} \cdots

 New analysis (in progress) $\left[\begin{array}{c|c} \setminus \end{array}\right]$ $\left[\begin{array}{c|c} \setminus \end{array}\right]$

- \bullet Run -6 (469 fb \bullet 1)
- \bullet New muon detector in additional data
- \bullet No PID selection and no momentum cut (detector limit $\stackrel{1}{\bullet}$ pT>0.1 GeV/*c* **µµ**g
- \bullet π /μ/K separation from angular distribution $P1 \sim 1$ OC V/C $\frac{1}{2}$ or $\frac{1}{2}$ • New method to separate all processes: fit of angular distributions in 2-particle CM frame

 $=$ effectively x7 statistics increase = ϵ enectively x/statistics increase ℓ

 \bullet Up to two extra photon admitted in the fit = \Rightarrow NNLO $\frac{1}{2}$ $\frac{1}{2}$ measurement m., 0.70-0.75 GeV → independent method allowing to check the previous BaBar result + improve the precision. M. Davier ISR BABAR g-2

Challenge: understand data/MC differences for tracking&trigger efficiencies below 1 GeV

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 $\pi\pi$

Data

Importance of low-momentum tracking (MC)

eev

 $\mu\mu\gamma$

 $\cos\theta^*_\pi$

mponent fit 0.600-0.605 Ge

µµg

p>0.2 GeV

 $KK\gamma$

 \mathbb{R}^n and \mathbb{R}^n $\pi\pi$ and $\qquad \qquad \text{for all}\qquad \mu\mu$

 \blacksquare

2000

 \mathbb{Z}_2 workshop, FermiLab 4, \mathcal{A}_1 , \mathcal{A}_2 , \mathcal{A}_3 , \mathcal{A}_4 , \mathcal{A}_5 , \mathcal{A}_6 , \mathcal{A}_7 , \mathcal{A}_8 , \mathcal{A}_7 , \mathcal{A}_8 , \mathcal{A}_9 , \mathcal{A}_9 , \mathcal{A}_9 , \mathcal{A}_9 , \mathcal{A}_9 , \mathcal{A}_9

Importance of low-momentum tracking (MC)

ppg

no p cut

p>0.2 GeV

p>1.0 GeV

10000

 $p>1.0$ GeV

μμγ MC

m., 0.70-0.75 GeV

Conclusions

- The muon *g*-2 is a long standing and among the most solid and significant discrepancies between experiment and SM theory
- Impressive recent improvements in every direction:
	- high quality e^+e^- data from new experiments/analyses
	- new direct measurement at Fermilab $(g-2)$ _µ experiment
	- accuracy of lattice calculations becoming competitive

But the puzzle now is not at all simpler w.r.t. WP!!

 $\mathcal{L}(\mathcal{A})=\mathcal{L}(\mathcal{A})$, outside of $\mathcal{L}(\mathcal{A})$, and Many open questions:

- KLOE vs BABAR vs CMD-3 Where the differences come from?
- Lattice *vs e*+*e* data?
- Will the final FNAL and JPARC results confirm the current central value?

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BACKUP slides

Fabio Anulli - Measurement of $e^+e^- \to \pi^+\pi^-\pi^0$ *-* ICHEP 2022 24 + *a*⁺ *n*⁺*n*[−]*n*⁰ + *a*⁺ *n*⁺*n*[−]*n*⁰ + *a*⁺ *a*⁺*n*[−]*n*⁰ + *a*⁺ *a*⁺*n*[−]*n*⁰ + *a*⁺ *a*⁺*n*^{−*n*} + *a*⁺ + *a*⁺

Why studying e^+e^- annihilation

Bernardini,Corazza,Ghigo,Touschek Il Nuovo Cimento 18, 1293 (1960)

- Many more *e+e-* collider since ADA:
	- c.m. energy up to 200 GeV (LEP-2)
	- Luminosity above

 1034 cm-2s-1 (*B*-factories) and 1035 cm-2s-1 (SuperKEKB)

All had begin with ADA in Frascati…. e^+e^- annihilation at low energies is a very effective tool to study the structure of the hadrons

- initial state well defined (momentum and quantum numbers)
- the final state can be fully reconstructed
- in general very simple event topology

Perturbative QCD approaches do not apply at low energies

- QCD-based models need experimental data as input
- Nature of many observed exotic states not yet understood
- Precise knowledge of **Hadronic Vacuum Polarization (HVP)** needed for the calculation of:
	- $\alpha_{\text{OED}}(M_Z^2)$: running fine structure constant at the Z^0 mass
	- **• (g-2)**µ **: anomalous magnetic moment of the muon**

HVP calculations

- Quark loops not computable from QCD (low mass scale)
- Can use dispersion relations, with analyticity and optical theorem to relate the vertex corrections to the *e*+*e*[−] ➝ *hadrons* cross section

$$
\boxed{\text{Im}\big[\text{Hom}\big]\longleftrightarrow\text{Im}\big[\text{Hom}\big]\times\text{Im}\big[\text{Hom}\big]\times\text{Im}\big[\text{Hom}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im}\big]\times\text{Im}\big[\text{Im
$$

$$
\text{Disperson integral}
$$
\n
$$
a_{\mu}^{HVP,LO} = \frac{\alpha^2}{3\pi^2} \int_{s_{th}}^{\infty} ds \frac{K(s)}{s} R(s)
$$
\n
$$
K(s)/s \approx 1/s^2
$$
 emphasizes the role

 $K(s)/s \simeq 1/5^2$ comphasizes the role of the processes at low energies Fabio Anulli - Hadronic cross sections — FPCP 2024

 ${\boldsymbol{\mathsf{u}}}$

γ

γ

 μ

hadrons

Muon *g***-2 Theory²**

Data combination for calculation of is quite a delicate procedure: *aHVP*−*LO μ*

 $>$ 300 data sets from $>$ 50 channels by dozens of experiments

combination and integration. Among them: Various approaches used to perform the data

- DHMZ, [Eur. Phys. J. C 80 \(3\) \(2020\) 241](https://link.springer.com/article/10.1140/epjc/s10052-020-7792-2)
- KNT, [Phys. Rev. D 101 \(2020\) 014029](https://doi.org/10.1103/PhysRevD.101.014029)
- F. Jegerlehner, [Springer Tracts Mod. Phys. 274 \(2017\) 1](https://link.springer.com/book/10.1007/978-3-319-63577-4)
- WP 2020 recommended a_u ^{HVP,LO} value is based on a combination of DHMZ and KNT evaluations
- They agree within roughly one sigma.
- Results on $a_{\mu} \pi \pi$ (for \sqrt{s} < 1.8 GeV):
	- KNT: *aµHVP,LO***[***ππ* **]***=* **(503.2 ± 1.9) ·10-10**
	- **DHMZ:** $a_{\mu}HVP,LO[\pi\pi] = (507.9 \pm 0.8 \pm 3.2)$

 $0₄$

split regions are assumed to have uncorrelated systematics. Data from Refs. $\frac{37}{4}$ Recommended WP2020 value from conservative merging of the various evaluations (dominated by the results from DHMZ19 and KNT19): $a_{\mu}HVP,LO = (693.1 \pm 2.8_{\text{exp}} \pm 2.8_{\text{sys}} \pm 0.7_{\text{DV-QCD}}) \cdot 10^{-10}$ KLOE-BABAR ππ experimental data *vs* QCD

 $\text{experiments} \quad \text{KLUE-BABAK }\pi\pi \quad \text{data vs QCD}$ is applicable to 9.2 GeV and 5.2 GeV and 5.2 GeV as above 11.6 GeV as above 11.5 GeV, see 11.6 GeV as above 11.5 GeV as above 11.5 GeV as above 11.5 GeV as above 11.5 GeV, see 11.5 GeV, se Fig. 29, using the contramners and $\frac{1}{2}$ differences discrepancy uncertainties differences

The other contributions to *aµ*

Higher order HVP:

various calculations, all in agreement

- NLO: a_u HVP,NLO = (-98.3 \pm 0.7) ·10⁻¹¹
- NNLO: a_u HVP, NNLO = $(12.4 \pm 0.1) \cdot 10^{-11}$

Hadronic Light-By-Light:

estimates with a dispersive relations approach and with lattice calculations compatible. Average value:

• a_{μ} HLBL = (92 ± 18) ·10⁻¹¹

QED contributions:

full perturbative calculation (α = 1/137.035) 10th order corrections calculated!

• $a_u QED = (116\,584\,718.931 \pm 0.104) \cdot 10^{-11}$

Electroweak contributions:

calculated up to two loops and an estimate of leading log contribution beyond 2-loop level \mathcal{O} between considered as \mathcal{O}

value for LO HVP given in Eq. (2.33), we adopt the central value from Ref. α

is of a similar size as the final accuracy goal of the Fermilab *g* 2 experiment. We will adopt this value for the NNLO

• $a_{\mu} E^{W} = (153.6 \pm 1.0) \cdot 10^{-11}$

contribution, which agrees well with the subsequent evaluation from Ref. [27], *a*HVP, NNLO

Recommended WP20 theory (SM) value: a_{μ} SM = (116 591 810 \pm 43) \cdot 10⁻¹¹ *a***HV** (SM) value: \blacksquare which, in particular, ensures that the uncertainty can still be considered anticorrelated with the one assigned *^µ* = ¹*.*24(1) ⇥ ¹⁰¹⁰ (2.35)

Overall estimate of a_{μ} ^{HVP,LO}

The Initial State Radiation method

$$
\frac{d\sigma_{e^+e^-\to f\gamma}(s,m_f)}{dm_f\,d\cos\theta_\gamma^*} = \frac{2m_f}{s}W(s,x,\theta_\gamma^*)\cdot\sigma_{e^+e^-\to f}(m_f)
$$

The hadronic cross section $e^+e^- \rightarrow f$ can be extracted from the ISR cross section $e^+e^- \rightarrow \gamma f$.

 Γ The radiator function *W(s,x)* is calculated in QED with accuracy better than 1% level where $\frac{10}{v}$ is calculated than $\frac{10}{v}$ level.

Common ISR analysis strategy

- **Tagged analysis** $(E_{\gamma}^* > 3 \text{ GeV})$ $U_{\gamma} > 0$ at \log analysis $(E_{\gamma} > 0eV)$
- Back-to-back topology btw ISR γ and the rest of the event
	- π /K/p discrimination based on dE/dx e Cherenkov angle
- **External acts •** Kinematic fit for 4-momentum conservation
- **Fitted** χ^2 **used for signal selection and background subtraction**
	- Detector acceptances and selection efficiencies estimated with MC simulation

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ISR method in a nutshell ISR method in a nutshell ISR method in a nutshell ISR method in a nutshell

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To tag or not to tag

Tagged approach:

☺ fully reconstructed events \rightarrow great background reduction ☹ ~90% signal loss

Untagged approach:

 $$\odot$ typically higher efficiency$ [®] higher background reduced by requiring the missing mass consistent with zero

So, what is the more convenient approach?

==> It depends on experimental situation

- At \sqrt{s} =10.58 GeV and for low m_f , (i.e. large *x*) the hadronic system has a large boost opposite to the photon direction \equiv \rightarrow the efficiency is almost insensitive to tagging
	- This is why, at BABAR:
	- Light Quarks final states \iff Tagged analyses
	- Heavy Quarks final states \iff Untagged analyses
- At $\sqrt{s} \approx 1$ GeV (**KLOE**) **untagged analyses** are more efficient

e+**e**[−] → π+π[−]: BESIII and "CLEO-c" *Source:* Reprinted from Ref. [60].

BESIII:

- Tagged ISR. $600 < \sqrt{s} < 900$ MeV. Analysis of **CLEO-c** data:
	- with $e^+e^- \rightarrow \mu^+\mu^-$. • Normalisation from MC, cross checked
- *2.2.3. The missing channels* • Syst. uncertainties at the *ρ*-peak: $\sim 0.9\%$ • Statistical uncertainty: 0.7%

Analysis of **CLEO-c** data:

- Tagged ISR $300 \leq s \leq 1000$ MeV.
- Normalisation from MC
- Statistical uncertainty: 0.7%
- Thanks to the BABAR systematic program of measurements of exclusive cross sections, very few channels are now few • Syst. uncertainties on a_{μ} : 1.5%

SND: new $e^+e^- \rightarrow \pi^+\pi^-$ (DSND: JHEP 01 (2021) 113

Energy scan at **VEPP-2000**: $0.32 < \sqrt{s} < 2$ GeV

- peak luminosity 7x1031 cm-2s-1
- higher than VEPP-2M (CMD-2)
- \sim 70% of data in the p-peak region
- This analysis:
	- 10% of collected data, with $525 < \sqrt{s} < 883$ MeV
	- Syst. uncertainties $\sim 0.8\%$
- $a_{\mu}^{\pi\pi}$ [525 883 MeV] = (409.79 ± 1.44 ± 3.87) × 10⁻¹⁰

CMD-3 *e*+*e*[−] ➝ π+π− **: Results**

 $Deta$ is α between than all other experiments Interference between Details of ρ-region: CMD-3 significantly

 $F \leftarrow C$ $\frac{1}{\sqrt{1000}}$ Interference between φ (1020) and non-resonant $\pi\pi$ **Production** interference between $\phi(1020)$

Fabio Anulli - Hadronic cross sections — FPCP 2024 ranges 0*.*3*÷*0*.*6/0*.*6*÷*0*.*9/ 0*.*9*÷*1*.*1 GeV the deviations are (+0*.*62*±*0*.*22)% / (0*.*06*±*0*.*03)%

Recent $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ results

χ 68.73 / 75 $1.094 + 0.042$

- and its unctreality 1.4 • Second largest contribution to $a_{\mu}^{had,LO}$
- Dominated at low energies by the ω and φ resonances, and then by ω recurrences $Q11, 14$.
	- Old data:

 $.04$

- \bullet Cross cross sections the tagged and until the ratios of the two. • Precise SND and CMD-2 @VEPP-2M below 1.1 GeV
	- *BABAR* (80 fb⁻¹) above the φ

Five new analyses, not in the WP, recently produced:

- *BABAR*: ISR@Y(4S) 0.62 < $m_{3\pi}$ < 3.5 GeV [PRD 104 (2021) 112003]
- **BESIII**: ISR@ψ(3770) 0.7 < $m_{3π}$ < 3 GeV [arXiv:1912.11208]
- **SND**: scan $1.2 < \sqrt{s} < 2$ GeV [EPJ C80, 993 (2020)]
- **BESIII**: scan $2 < \sqrt{s} < 3.08$ GeV [arXiv2401.14711]
- **CMD-3**: scan $0.66 < \sqrt{s} < 0.97$ GeV [preliminary]

Clear observation of $\omega(1420)$ and $\omega(1650)$ All recent data inconsistent with old DM2 data

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Measurement of additional radiation in ISR processes 1 **EIVI eas!**

● Study of and with 1 (NLO) or 2 (NNLO) additional hard photons *e*+*e*[−] → *μ*+*μ*−*γISR e*+*e*[−] → *π*+*π*−*γISR* $H - 468$ fb⁻¹) is used. Data compared with MC generators: *u***p** to NNLO additional ISR collinear to beams, FSR from PHOTOS **• Phokara 9.1:** full matrix elements at NLO (no NNLO) **Study ISR Processes e+e−→**!**+**!**−/**"**+**"**−**#**ISR with 1 or 2 more Hard Photons** 1 10 10 ² 1 $\sum_{n=10}^{\infty} \frac{10^{n} \text{ FWHM}}{10^{2}}$ $\frac{1}{10}$ ³

YISR YSA

 γ_{SA} e^-

- consistent with the event topology \sim Small Angle (SA) fitted **Study ISR Processes** to the series of the Sevent Europe is to the Europe Hard Photons in the Europe of the Europe Hard Photons in the Europe of the Eu • All events are subjected to several fits $\text{TL} \times \text{C}$ and the measured one C from the ISRLA
- best χ^2 determine ategory of that event • The fit with the best χ^2 determine the category of that event
- $\frac{1}{\sqrt{2}}$ e -angle (LA) .
ion from anou $\frac{1}{4}$ • Small-angle (SA) γ 's assumed collinear with ➣ Three NNLO fits $\frac{1}{1}$ *γ*_{*LA*} to the closest track $\overline{\mathbf{L}}$ the data of $\overline{\mathbf{L}}$ of \mathbf{L} and $\overline{\mathbf{L}}$ ($\overline{\mathbf{L}}$ \mathbf{L}) \mathbf{L} ² d the beams. Large-angle (LA) $γ$'s detected FCD/ICD conoration from anoular • FSR/ISR separation from angular distance of

NNLO event topologies

NLO event topologies

 γ _{ISR} γ _{LA}

➣ ": 0.35−2.4 rad

Two NLO fits:

Three categories:

- NLO LA
- NLO SA
- LO: events with no γ 's above threshold

Three NNLO fits and three categories:

- *γ*_{*ISR}γ_{SA}γ_{SA}* -> NNLO 2SA</sub>
- *γ*_{*ISR}γ_{SA}γ_{LA}* -> NNLO SA+LA</sub>

```
• γ<sub>ISR</sub>γLAγ<sub>LA</sub> -> NNLO 2LA
```
BABAR: Phys. Rev. D108 (2023) 111103 ➣ ": 0.35−2.4 rad

 $e^+e^- \rightarrow K^+K^-$

- Cross section dominated by the $\varphi(1020)$ peak
- Precise measurements from BABAR, CMD-2/3 and SND
- Some tension at the peak with new CMD-3 data significantly above the older Novosibirsk results and BABAR data in the middle
- Note: $\varphi(1020)$ just above KK threshold. Very low momentum tracks in energy scan experiments
- Green combination from DHMZ

$e^+e^- \rightarrow \text{KK}\pi$ and KK $\pi\pi$

The E821 and E989 direct *a* μ The E821 and E989 direct a_n measurements positron trajectories to the FQ20 directories to the trajectories to the TQ20 directories to the TQ20 director $\sum_{i=1}^{\infty}$ with $\sum_{i=1}^{\infty}$. The storage aperture $\sum_{i=1}^{\infty}$

- Intense proton beam on target to produce the pion beam. $\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$ \bullet Intense proton beam on target to produce the pion beam
- \bullet Pions of 3.1 GeV are selected.
- $\pi \mapsto \mu + \nu$ violates $P = \nu > 95\%$ polarisation for forward muons T_{tot} \sim 54–56 for a containing and containing a containing
	- μ ⁺ stored in a cyclotron: constant *B* field
- cyclotron frequency ω_c ; spin precessing with freq. ω_s \cdot μ÷ stored in a cyclotron; constant *B* frequency compared in a cyclotron frequency ω ; and precessing with freq. ω • cyclotron frequency ω_c ; spin precessing with freq. ω_s
- $\bullet \ \omega_a = \omega_s \omega_c = a_\mu eB/m_\mu$
- $\mu \rightarrow e \nu \nu$ violates $P = \rightarrow e^+$ direction (energy in LAB) "remembers" the μ polarization μ and σ and σ (i.e., σ) and σ) are σ . μ poiarization
- fraction of detected e^+ with $E>E_{threshold}$ modulated with E98 frequency ω_a ϵ from of detected at with $E\setminus E$ modulated w rection of detected ϵ

$$
N_{ideal}(t) = N_0 \exp(-t/\gamma \tau_\mu) \left[1 - A \cos(\omega_a t + \phi)\right], \frac{1}{\sqrt{\frac{a}{\pi}}} \frac{1}{\sqrt{\frac{a}{\pi}}} \frac{1}{\sqrt{\frac{a}{\pi}}}.
$$

- where N0, A and φ are all implicitly dependent on Eth. For a threshold energy of 1.8 GeV Precise knowledge of B field critical frequency !⁰ *p*(*Tr*) of a proton shielded in a spherical sam-• Precise knowledge of B field critical
- (y ≈ 0.58 in Fig. 1b), the asymmetry is ≈ 0.4 and the average FOM is maximized. A use of pulsed proton NMR The magnetic field *B* is determined from the precession
- B determined from proton precession frequency *p* determined from proton precession frequent $\frac{1}{2}$ and $\frac{1}{2}$
	- a_{μ} extracted from:

$$
a_\mu = \frac{\omega_a}{\tilde \omega_p'(T_r)} \frac{\mu_p'(T_r)}{\mu_e(H)} \frac{\mu_e(H)}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}
$$

- measure ω_a and ω_p
- Other quantities from external inputs

Perspectives in direct measurements $\frac{1}{2}$ A $\frac{1}{2}$ Experiment of $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2}$

- 1. The "Muon g-2" experiment at Fermilab is taking new data
- Final goal: **uncertainty of 0.14 ppm** (0.1 ppm stat., 0.07 *ppm* syst.) its current status is described.
- (now they reached 0.35 *ppm*)
- very ambitious program, but it looks under reach

2. J-PARC E34 Collaboration

- **five.** The projected sensitivity of the first phase of magnitude $\frac{1}{2}$ $\mathbf S$ denomination than $\mathbf S$
- ultra-cold muons produced from pion decay at rest
- accelerated by a linac at $p = 300$ MeV
- beam with negligible transverse momentum
	- no need for strong electromagnetic focusing => no need of the 3.1 GeV *magic momentum*
- small storage ring (33 cm radius); cyclotron period only 7.4 ns; 3T dipole magnet
- Predicted statistics 10x FNAL
- Final goal: **uncertainty of 0.1** *ppm*

Alternative approach for a_μ ^{HLO}: MUonE experiment where *^R*had(*s*) is the ratio of the total *^e*⁺*e* ! hadrons and the Born *^e*⁺*e* ! *^µ*⁺*µ* cross sections, *^K*^ˆ (*s*) is a smooth for a *MLO* (*MI* lonG experiment \sum is due thus \sum approach for α_{μ} and \sum integral \sum experiment $\overline{13}$ app μ_{W} thus and the form and the HIO. MI Long and the muon to the muon μ ialive approach for a_{μ} . IVFOOILE ex

- $u_1 = 1, 1, \ldots, u_n = 1, 1, \ldots, u_n = 1$ • Standard approach using e^+e^- data $a_\mu^{\text{HLO}} = \left(\frac{\alpha m_\mu}{3\pi}\right)^2$ $\int ds \frac{K(s)R_{\text{had}}(s)}{s^2}$, $(s > 0)$ integral in the time-like region) $\int_a^b \int_a^\infty ds \frac{\hat{K}(s)R_{\text{had}}(s)}{2}$ $g(R_{\text{had}}(s))$ *x*²*m*² *x* $\frac{1}{2}$ **1** $\frac{1}{$ $a_{\mu}^{\mathrm{HLO}} =$ $\sqrt{\alpha m_\mu}$ 3π λ^2 \int^∞ $4m_\pi^2$ $ds \frac{\hat{K}(s)R_{\text{had}}(s)}{2}$ $a_\mu^{\text{HLO}} = \left(\frac{\omega m_\mu}{3\pi}\right) \int_{4m_\pi^2} ds \frac{1 + (b)2v_{\text{Had}}(b)}{s^2},$ Measuring the running of ↵(*t*) in the space-like region, using a muon beam with *Eⁱ* • Standard approach using *e+e-* data
	- $a_{\mu}^{\text{HLO}} = \frac{\alpha}{\pi}$ \int_0^1 $\int\limits_{0}^{t} dx \left(1-x\right) \Delta \alpha_{\text{had}}[t(x)] \, , \qquad t(x) = \frac{x-\mu}{x-1} < 0 \, .$ $\frac{x^2m_{\mu}^2}{x-1} < 0$ is a space-like (negative) squared four-momentum and r^2m^2 fine-structure constant. In contrast with the integrand function of Eq. (1), the integrand in Eq. (2) is smooth and free of $\frac{1}{2}$ **R**_d is the ratio of \int_{0}^{1} **e** \int_{0}^{1} **e** \int_{0}^{1} cross section is and \int_{0}^{1} \int_{0}^{1} cross section is a section is a section of \int_{0}^{1} **x**² m_{μ}^{2} cross sections, **x**² m_{μ}^{2} and **x**² s approach: $a_{\mu}^{\text{mLO}} = \frac{1}{\pi} \int dx (1-x) \Delta a_{\text{had}}[t(x)]$, $t(x) = \frac{R}{\pi} < 0$ $i₀$ fluctuating at low energy due to responsive integral integral in $x-1$ **µ ' 150 GeV on a fixed electron** th a technique similar to the one described in $x^2m_\mu^2$ • Alternative approach:

is a space-like region \rightarrow integrand smooth and Iree of resonances $t = q^2 < 0$ integral in the space-like region \Rightarrow integrand smooth and free of resonances perturbative QCD (pQCD) [19] in the high-energy tail. For the calculation of *a*HLO

• Measure the running of $\alpha(t)$ from the elastic scattering $\mu^+e^- \to \mu^+e^-$ to extract $\Delta \alpha(t)$ \overline{dt}

g of
$$
\alpha(t)
$$
 from the elastic
\n μ^+e to extract $\Delta\alpha(t)$
$$
\frac{d\sigma}{dt} = \frac{d\sigma_0}{dt} \left| \frac{\alpha(t)}{\alpha(0)} \right|^2, \quad \alpha(t) = \frac{\alpha(0)}{1 - \Delta\alpha(t)},
$$

• Subtract the leptonic contribution (calculable in perturbation theory) to get: $\Delta a_{bad}(t) = \Delta a(t)$ - $\Delta a_{can}(t)$ • Subtract the leptonic contribution (calculable in perturbation theory) to get: $\Delta \alpha_{had}(t) = \Delta \alpha(t) - \Delta \alpha_{lep}(t)$ *dx* (1 *x*) ↵had[*t*(*x*)] *,* (2)

MUonE at CERN North-Area

where

-
-
- Experimental coverage:
	-
	- \sim 87% of all a_μ ^{HVP,LO}
- Remaining part estimated using time-
• EM calorimeter and muon tracking downstream to measure like data + pQCD, or Lattice+QCD all components of the moments of the momenta of the mo • Remaining part estimated using time-

II. EXPERIMENTAL PROPOSALE

 $\left\{\begin{array}{c} \lambda^{n} \text{ and } -257.57154 \end{array}\right\}$

44

 $\mathbf{y} = \begin{bmatrix} 1.02 \end{bmatrix}$

• \sim 87% of all a_{μ} HVP,LO

• 40 Be targets, 1.5cm-thick, equipped with Si trackers

Fig. 2 (left) shows ↵lep and ↵had as functions of the scale *t*(*x*) defined in Eq. (3). The integrand function of

- for **the peak of the integrand of the integrand of the final state** $\frac{1}{2}$ and *t* $\frac{1}{2}$ and *t***_p and** *t***_p an** neter and muon
ents of the final all components of the final state *,* (4) • EM calorimeter and muon tracking downstream to measure
- $\frac{d}{dx}$ (fig. 2). The scale $\frac{d}{dx}$ defined in Eq. (3). The integrand functions of the scale $\frac{d}{dx}$ • With three years of running at planned conditions, expect a statical undertainty of $\sim 0.3\%$ (which uses a ratio of $\sim 0.3\%$) • With three years of running at planned conditions, expect a statical undertainty of $\sim 0.3\%$

We propose to use Eq. (2) to determine *a*HLO

Eq. (2) involves ↵had(*q*²) evaluated at negative space-like momenta *t <* 0.

- Systematic and theoretical uncertainties to be carefully evaluated $\begin{bmatrix} 1 & x^2 \end{bmatrix}$ $\begin{bmatrix} x^2 \end{bmatrix}$ $\begin{bmatrix} x^2 \end{bmatrix}$ $\begin{bmatrix} 1.08 \\ 1.06 \end{bmatrix}$
- Results should be competitive and fully independent from e^+e^- data 10^3 $\left[\right]$ $\left[\right]$