# Prospects for New Physics in Rare Kaon Decays Zachary Polonsky (University of Zurich) University of Zurich<sup>UZH</sup> May 29, 2024 22nd FPCP 2024, Bangkok

#### Kaons historically vital for new physics discoveries!



	<u>Kaons</u>	<u>B-mesons</u>
GIM Suppression:	$ V_{ts}^*V_{td} \sim\lambda^5$	$ V_{tb}^*V_{td(s)} \sim\lambda^{3(2)}$
Decay Suppression:	$\Gamma \sim M_K^5/M_W^4$	$\Gamma \sim M_B^5/M_W^4$
Light NP:	$\mathcal{B} \sim (M_W/M_K)^n$	$\mathcal{B} \sim (M_W/M_B)^n$

# Promising Observables

(See talks from Xu Feng, Silvia Martellotti, Yu-Chen Tung, and Rainer Wanke for more details!)

$$\mathcal{H}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} \frac{\alpha}{2\pi \sin^2 \Theta_W} \sum_{\ell=e,\mu,\tau} \left( V_{cs}^* V_{cd} X_c^\ell + V_{ts}^* V_{td} X_t \right) (\bar{s}_L \gamma^\mu d_L) (\bar{\nu}_L^\ell \gamma_\mu \nu_L^\ell) + \text{h.c.}$$

 $\blacktriangleright \text{ Re } V_{ts}^* V_{td} \sim \text{Im } V_{ts}^* V_{td} \sim \lambda^5, \qquad \text{Re } V_{cs}^* V_{cd} \sim \lambda, \qquad \text{Im } V_{cs}^* V_{cd} \sim \lambda^5$ 

#### $K_L \rightarrow \pi^0 \bar{\nu} \nu$ : A Theorist's Dream Decay

$$|K_L\rangle = p \left| K^0 \right\rangle - q \left| \bar{K}^0 \right\rangle \quad \Rightarrow \quad \left\langle \pi^0 \bar{\nu} \nu \right| \mathcal{H}_{\text{eff}} \left| K_L \right\rangle = \frac{4G_F}{\sqrt{2}} \frac{\alpha}{2\pi \sin^2 \Theta_W} \operatorname{Im} \left( V_{ts}^* V_{td} X_t \right) \left\langle Q_\nu \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts}^* V_{td} X_t \right\rangle \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts}^* V_{td} X_t \right\rangle \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{ts} \right\rangle + \mathcal{O} \left( \lambda^5 \frac{m_c^2}{M_W^2} \right) \left\langle W_{$$

▶ Nearly pure *CP*-violating ⇒ Top-quark dominated (tiny long-distance)

$$\mathcal{B}(\mathcal{K}_L o \pi^0 ar{
u} 
u)_{\mathsf{SM}} = ig( 2.59(6)_{\mathsf{SD}}(2)_{\mathsf{LD}}(28)_{\mathsf{param}} ig) imes 10^{-11}$$

[Brod, Gorbahn, Stamou; 2105.02868], [Mescia, Smith; 0705.2025], [Buchalla, Buras; 9607447]

Very challenging experimentally (fully neutral final state)

 $K_I \rightarrow \pi^0 \bar{\nu} \nu$ : KOTO at J-PARC

2015 dataset results:

$$\mathcal{B}(\mathcal{K}_L o \pi^0 ar{
u} 
u) < 3.0 imes 10^{-9} \ (90\% \ ext{CL})$$

• Charged-K veto counter and reduction of halo  $K_L \rightarrow 2\pi^0$ 

Updated analysis:

$${\cal B}({\cal K}_L o \pi^0 ar 
u 
u) < 2.0 imes 10^{-9} \ (90\% \, {
m CL})^{-9}$$



 $K_I \rightarrow \pi^0 \bar{\nu} \nu$ : KOTO-II at J-PARC



#### $\blacktriangleright$ Expected $\sim$ 25% sensitivity for SM BR

$$K^+ \to \pi^+ \bar{\nu} \nu$$

► Real part of  $\mathcal{H}_{eff}$  plays non-negligible role ⇒ charm contributions compete  $(\lambda m_c^2/M_W^2 \text{ vs. } \lambda^5)$ 

- ► More dependent on long-distance than K<sub>L</sub> case; O(5%) [Isidori, Mescia, Smith; 0503107]
- Future lattice calculations [Bai et al.; 1806.11520]



$${\cal B}({\cal K}^+ o \pi^+ ar 
u 
u)_{\sf SM} = ig( 7.73(16)_{\sf SD}(25)_{\sf LD}(54)_{\sf param}ig) imes 10^{-11}$$

[Brod, Gorbahn, Stamou; 2105.02868], [Mescia, Smith; 0705.2025]

#### Run I results:

$$\mathcal{B}(\mathcal{K}^+ o \pi^+ ar{
u} 
u) = ig(10.6^{+4.0}_{-3.4}|_{\mathsf{stat}} {\pm} 0.9|_{\mathsf{syst}}ig) { imes} 10^{-11}$$

#### Run II:

2022 signals ~ all Run I [2023 NA62 Status Report]
 15% precision (similar data taking up to LS3)



## $\epsilon_{\mathcal{K}}$ : Experiment Ahead of Theory

- Measure of indirect *CP*-violation in  $K^0 \bar{K}^0$  mixing
- ▶ Theoretical uncertainty  $\sim \mathcal{O}(1\%)$

Perturbative [Brod, Gorbahn, Stamou: 1911.06822], [Brod, Kvedaraité, ZP; 2108.00017]. [Brod, Kvedaraité, ZP, Youssef; 2207.07666]

Hadronic ME [FLAG; 2111.09849]

► m<sub>c</sub> power corrections [Ciuchini, et al.; 2111.05153]

$$\epsilon_{\mathcal{K}}|_{\mathsf{th}} = (2.170(65)_{\mathsf{pert}}(76)_{\mathsf{nonpert}}(153)_{\mathsf{param}}) \times 10^{-3}$$

Experimentally measured to per-mil accuracy [PDG 2022]

$$|\epsilon_{K}|_{\mathsf{ex}} = (2.228 {\pm} 0.011) { imes} 10^{-3}$$

Theory calculation will be improved with 3loop QCD top, NLO RI/SMOM-MS
matching and improved lattice calculations!



$$K \rightarrow \mu^+ \mu^-$$

- Large LD contaminations from CP-even parts of decays: hard to control
- Interference effects dominated by SD  $\Rightarrow$  theoretically clean [D'Ambrosic, Kitahara; 1707.06999]
- ▶ Mostly *CP*-violating  $\mathcal{B}(K_S \to (\mu \mu)_{\ell=0})$  can be determined from\*

[Dery, Ghosh, Grossman; 2104.06427]

$$\mathcal{B}(\mathcal{K}_{\mathcal{S}} 
ightarrow (\mu\mu)_{\ell=0}) = \mathcal{D}_{\mathcal{F}}\mathcal{B}(\mathcal{K}_{L} 
ightarrow \mu\mu) rac{ au_{\mathcal{S}}}{ au_{L}} \Big(rac{\mathcal{C}_{\mathsf{int}}}{\mathcal{C}_{L}}\Big)^{2}$$

LD effects from mixing can be enhanced by  $|\mathcal{A}_{L}^{0}|/|\mathcal{A}_{S}^{0}| \sim 10 \ (\lesssim \text{few \%})$ 

Experimentally challenging:  $C_{int}$  only seen in  $K^0 - \bar{K}^0$ -asymmetric beam

- Need to extract ChPT FFs from data
- ►  $K_{L/S} \rightarrow \pi^0 \bar{\ell} \ell$ : both depend on same FF ⇒ Measuring  $K_S$  decay can help prediction of  $K_L$  (LHCb)

► 
$$K^+ \rightarrow \pi^+ \bar{\ell} \ell$$
: FF difference b/n  $\ell = e, \mu$  gives LFUV test

$${\sf LFUV}(a_+^{\mu\mu}{-}a_+^{ee})=-0.014{\pm}0.016$$

[D'Ambrosio, Mahmoudi, Neshatpour; 2209.07445], [E865; 9907045], [NA62; 2209.05076]



## Heavy New Physics

- ► Treated in model-independent\* way with corrections to WCs:  $C \rightarrow C_{SM} + \delta C$
- ► Heavy GIM suppressions ⇒ high-scale/weakly coupled heavy NP

$$\frac{4G_{F}}{\sqrt{2}}\frac{\alpha}{2\pi\sin^{2}\Theta_{W}}\lambda_{t}\sim-(130 \text{ TeV})^{-2}+i(200 \text{ TeV})^{-2}$$

#### HNP in $K \to \pi \bar{\nu} \nu$ : Current Status

• Compatible with SM at  $1\sigma$ 

Charged Decay (90% CL):

 $|\operatorname{\mathsf{Re}} \delta C_L| \lesssim (120 \, \operatorname{\mathsf{TeV}})^{-2}, \quad |\operatorname{\mathsf{Im}} \delta C_L| \lesssim (70 \, \operatorname{\mathsf{TeV}})^{-2}$ 

▶ Neutral Decay (90% CL):

 $|\operatorname{Im} \delta C_L| \lesssim (50 \text{ TeV})^{-2}$ 



- Projected ~ 15% from NA62
   SM-like Central Value (90% CL):
   | Re δC<sub>L</sub>| ≲ (225 TeV)<sup>-2</sup>, | Im δC<sub>L</sub>| ≲ (100 TeV)<sup>-2</sup>
- Unchanged Central Value: Compatible at 2σ





SM-like Central Value (90% CL):

$$\operatorname{Re} \delta C_L | \lesssim (240 \text{ TeV})^{-2}, \quad |\operatorname{Im} \delta C_L| \lesssim (280 \text{ TeV})^{-2}$$

Unchanged NA62 Central Value: Compatible at 2σ



- Extra NA62 data-taking:  $\sim 10\%$  (optimistic)
- SM-like Central Value (90% CL):

$$\operatorname{Re} \delta C_L | \lesssim (290 \text{ TeV})^{-2}, \quad |\operatorname{Im} \delta C_L| \lesssim (280 \text{ TeV})^{-2}$$

Unchanged NA62 Central Value:  $\sim 3\sigma$  tension



#### Other BSM Operators

• Currently  $C_i \lesssim (\sim \text{few 10s-100 TeV})^{-2}$ : will be improved in future



[Gorbahn, et al.; 2312.06494]

### Other (Semi-)Leptonic Decays

•  $SU(2)_L$  symmetry:  $(\bar{s}_L \gamma_\mu d_L)(\bar{\ell}_L \gamma^\mu \ell_L) \Rightarrow K \rightarrow \pi \nu \bar{\nu}, K \rightarrow \pi \ell \bar{\ell}, K \rightarrow \ell \bar{\ell}$  complimentary



[D'Ambrosio, Mahmoudi, Neshatpour; 2311.04878]

Light New Physics

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Searches less model-independent than heavy case (can "see" light dynamics)

- Plethora of possible NP models: see review [Goudzovski, et al.; 2201.07805]
- ▶ Light flavored NP can appear in many solutions to other SM problems:
  - Inflation
  - Strong CP
  - Baryogenesis
  - Dark Matter (via dark photons/Higgs portals)
  - Flavor puzzle
  - Neutrino oscillations/masses
  - And combinations thereof!

### Light New Physics

Light NP benchmarks (Physics Beyond Colliders BSM study group):

- 1. Minimal dark photon
- 2. Light DM-coupled dark photon
- 3. Millicharged particles
- 4. Higgs-mixed scalar
- 5. Higgs-mixed scalar + pair production
- 6. Single HNL  $(U_{eN})$
- 7. Single HNL  $(U_{\mu N})$
- 8. Single HNL  $(U_{\tau N})$
- 9. Photon-coupled ALP
- 10. Fermion-coupled ALP
- 11. Gluon-coupled ALP

<sup>[</sup>Beacham, et al.; 1901.09966]

$$K^+ 
ightarrow \pi^+ X_{
m inv}$$

Generalized Grossman-Nir bound:

$$\mathcal{B}(\mathcal{K}_L o \pi^0 X) \leq 4.3 imes \mathcal{B}(\mathcal{K}^+ o \pi^+ X)$$
 [Grossman, Nir; 9701313]

 $\Rightarrow$  Stronger bounds from charged decays

- NA62 sensitive in  $K^+ \rightarrow \pi^+ \bar{\nu} \nu$  signal regions:  $m_X = 0 - 110$  MeV, 160 - 260 MeV
- ► Can also probe  $K^+ \to \pi^+(\pi^0 \to X_{inv})(\gamma)$  in intermediate region [NA62; 2010.07644]





[NA62; 2103.15389]

#### Dark Photon



[NA48/2; 1504.00607]



$$K^+ 
ightarrow \pi^+ (\pi^0 
ightarrow \gamma A')$$

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 $K^+ 
ightarrow \pi^+ (\pi^0 
ightarrow \gamma (A' 
ightarrow e^+ e^-))$ 

#### Heavy Neutral Leptons



$$K^+ 
ightarrow \ell^+ N \ (\ell = e, \mu)$$

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### Light New Physics

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### Light New Physics

Light NP benchmarks (Physics Beyond Colliders BSM study group):

- 1. Minimal dark photon ←
- 2. Light DM-coupled dark photon ←
- 3. Millicharged particles
- 4. Higgs-mixed scalar ←
- 5. Higgs-mixed scalar + pair production
- 6. Single HNL  $(U_{eN}) \leftarrow$
- 7. Single HNL  $(U_{\mu N}) \leftarrow$
- 8. Single HNL  $(U_{\tau N})$
- 9. Photon-coupled ALP
- 10. Fermion-coupled ALP ←
- 11. Gluon-coupled ALP  $\leftarrow$

<sup>[</sup>Beacham, et al.; 1901.09966]

# $K^+$



#### Generalized GN-violating NP

Violations of generalized GN bound possible (almost universally light):

- 1. Additional CPV in Decay (challenging to make work)
- 2. New states enhance  $\Delta I = 3/2$  transition [He, et al.; 2005.02942]
- 3. Charge conservation:  $K_L \rightarrow X^2$  vs  $K^+ \rightarrow X^2 \pi^+$  [Gori, Perez, Tobioka; 2005.05170]. [Hostert, Kaneta, Pospelov; 2005.07102]
- 4. Charged vs neutral mass difference:  $m_{K^+} m_{\pi^+} < 2m_X < m_{\mathcal{K}_L} m_{\pi^0}$  [Fabbrichesi, Gabrielli; 1911.03755]
- 5. Variety of experimental loopholes (blind spots, unstable  $X \rightarrow SM$ , etc.)
- $\blacktriangleright$  K<sub>L</sub> still very important for LNP searches!

- We don't control where or what NP is!
- ► Kaon decays cast wide BSM net (GIM suppression/*B* enhancement)
- Simultaneously test theoretically well-understood hadronic decays
- ▶  $K^{\pm}$  and  $K_{L/S}$  provide complimentary information for both HNP and LNP

Thank you!