

# Searching new physics in $c \rightarrow ul^+l^-$ decays with non-universal $Z'$ model

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**Introduction**

**Theoretical framework**

**Estimation of NP couplings in  $Z'$  model**

**Results**

**Conclusion**

**References**

**Acknowledgement**

**OUTLINE**

## Motivation

- Recent results from the LHCb have provided the limit of the branching ratio of lepton flavour violating decays  $D^0 \rightarrow \mu^\mp e^\pm$  as  $\mathcal{B}(D^0 \rightarrow \mu^\mp e^\pm) < 1.3 \times 10^{-8}$  at 90% confidence level [1].
- The SM calculation including the short distance effects for the decays  $D \rightarrow l^+ l^-$  predicts the branching fractions to be of the order of  $10^{-18}$  [2].
- For the leptonic decays  $D^0 \rightarrow e^+ e^-$  and  $D^0 \rightarrow \mu^+ \mu^-$ , the experimental bounds for the branching fractions are constrained at  $7.9 \times 10^{-8}$  [3] and at  $3.1 \times 10^{-9}$  [4], respectively.
- Effective GIM (Glashow-Iliopoulos-Maiani) suppression in  $c \rightarrow u$  decays make these channels sensitive to new physics effects [5].

## Motivation

- The rare charm decays  $D_s^+ \rightarrow K^+ e^- e^+$ ,  $D_s^+ \rightarrow K^+ e^- \mu^+$ ,  $D^0 \rightarrow e^- e^+$  and  $D^0 \rightarrow e^- \mu^+$  have been studied in  $U(1)'$  model,  $2HDM$  model, and unparticle model where only  $Z'$  model and  $2HDM$  model significantly increased the branching ratios [6].
- The decays  $D^{+(0)} \rightarrow \pi^{+(0)} \mu^+ \mu^-$ ,  $D^0 \rightarrow \mu e, \tau e$  and  $D^{+(0)} \rightarrow \pi^{+(0)} \mu^- e^+$  have been investigated in leptoquark model and  $Z'$  model [7].
- Along with  $D \rightarrow \pi l^+ l^-$ ,  $D_s \rightarrow K l^+ l^-$  has also been studied in leptoquark model, SUSY models and flavourful  $Z'$  model [8].

[6] Xing-Dao Guo, Xi-Qing Hao, Hong-Wei Ke, Ming-Gang Zhao and Xue-Qian Li, *Chinese Phys. C* **41** 093107 (2017)

## Non-universal $Z'$ model

- The gauge group  $SU(3) \times SU(2) \times U(1)$  of SM is extended to  $SU(3) \times SU(2) \times U(1) \times U'(1)$ .
- Extending the gauge group introduces a new particle, the  **$Z'$  boson**.
- $Z'$  boson is a massive, charge neutral, colourless, hypothetical particle.
- $Z'$  boson couples to quarks as well as leptons through FCNC (flavour changing neutral current) transitions.

## Search for $Z'$ boson

### Experimental searches

- ATLAS collaboration [9]
  1.  $M_{Z'} \sim 0.5 - 2.5$  TeV
  2.  $M_{Z'_{SSM}} > 1.90$  TeV
  3.  $M_{Z'_{SFM}} > 1.82 - 2.17$  TeV at  $\sqrt{s} = 13$  TeV
- CMS collaboration [10]
  1.  $M_{Z'_{SSM}} > 4.50$  TeV
  2.  $M_{Z'_\psi} > 3.90$  TeV
- EW data of weak neutral current processes and mixing between  $Z - Z'$ :  $M_{Z'} > \mathcal{O}(500)$  GeV [11]
- Tevatron:  $M_{Z'} > \mathcal{O}(800)$  GeV [12]

### Theoretical searches

- Sahoo et al.:  $M_{Z'} \sim 1352 - 1665$  GeV [13]
- Allanach:  $M_{Z'} \sim 0.8 - 1.8$  TeV [14]
- Luzio et al.:  $M_{Z'} \leq 9$  TeV [15]

[14] B. Allanach, LHC di-lepton searches for  $Z'$  bosons which explain measurements of  $b \rightarrow sl^+ l^-$  transitions, arXiv: 2404.14748 [hep-ph] (2024)

# Theoretical Framework

We have studied the decays  $D^+ \rightarrow \pi^+ \mu^+ \mu^-$ ,  $D^0 \rightarrow \pi^0 \mu^+ \mu^-$ ,  $D_s^+ \rightarrow K^+ \mu^+ \mu^-$ ,  $D^+ \rightarrow \rho^+ \mu^+ \mu^-$  and  $D^0 \rightarrow \rho^0 \mu^+ \mu^-$  in non-universal  $Z'$  model

Effective Hamiltonian for  $c \rightarrow ul^+l^-$  transition [8, 16]

$$\mathcal{H}_{eff} = -\frac{4G_F}{\sqrt{2}} \frac{\alpha}{4\pi} \left[ \sum_{i=7,9,10,S,P} (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i) + \sum_{i=T,T_5} C_i \mathcal{O}_i + \sum_{q=d,s} V_{cq}^* V_{uq} \sum_{i=1}^2 C_i \mathcal{O}_i^q \right]$$

The differential decay distribution of  $D \rightarrow Pl^+l^-$  in standard model is given by [8, 16]

$$\begin{aligned} \frac{d\Gamma}{dq^2} &= \frac{G_F^2 \alpha^2 \beta_l}{1024 \pi^5 M_D^3} \left\{ \frac{2}{3} \left| C_9 + C_7 \frac{2m_c}{M_D + M_P} \frac{f_T}{f_+} \right|^2 \left( 1 + \frac{2m_l^2}{q^2} \right) \lambda f_+^2 \right. \\ &\left. + |C_{10}|^2 \left[ \frac{2}{3} \left( 1 - \frac{4m_l^2}{q^2} \right) \lambda f_+^2 + \frac{4m_l^2}{q^2} (M_D^2 - M_P^2)^2 f_0^2 \right] \right\} \end{aligned}$$

The decay rate for  $D \rightarrow \rho l^+ l^-$  is given by [17]

$$\begin{aligned} \frac{d\Gamma}{d\hat{s}} &= \frac{G_F^2 \alpha^2 M_D^2}{2^{10} \pi^5} \left\{ \frac{|A|^2}{3} \hat{s} \lambda \left( 1 + \frac{2\hat{m}_l^2}{\hat{s}} \right) + |E|^2 \hat{s} \frac{\hat{u}(\hat{s})^2}{3} \right. \\ &+ \frac{1}{4\hat{M}_\rho^2} \left[ |B|^2 \left( \lambda - \frac{\hat{u}(\hat{s})^2}{3} + 8\hat{M}_\rho^2 (\hat{s} + 2\hat{m}_l^2) \right) \right. \\ &+ |F|^2 \left( \lambda - \frac{\hat{u}(\hat{s})^2}{3} + 8\hat{M}_\rho^2 (\hat{s} - 4\hat{m}_l^2) \right) \left. \right] \\ &+ \frac{\lambda}{4\hat{M}_\rho^2} |C|^2 \left( \lambda - \frac{\hat{u}(\hat{s})^2}{3} \right) \\ &+ |G|^2 \left( \lambda - \frac{\hat{u}(\hat{s})^2}{3} + 4\hat{m}_l^2 (2 + 2\hat{m}_l^2 - \hat{s}) \right) \left. \right] \\ &- \frac{1}{2\hat{M}_\rho^2} \left[ \text{Re}(BC^*) \left( \lambda - \frac{\hat{u}(\hat{s})^2}{3} \right) (1 - \hat{M}_\rho^2 - \hat{s}) \right. \\ &+ \text{Re}(FG^*) \left( \left( \lambda - \frac{\hat{u}(\hat{s})^2}{3} \right) (1 - \hat{M}_\rho^2 - \hat{s}) + 4\hat{m}_l^2 \lambda \right) \left. \right] \\ &\left. - \frac{2\hat{m}_l^2 \lambda}{\hat{M}_\rho^2} \left[ \text{Re}(FH^*) - \text{Re}(GH^*) (1 - \hat{M}_\rho^2) \right] + \frac{\hat{m}_l^2}{\hat{M}_\rho^2} \hat{s} \lambda |H|^2 \right\} \end{aligned}$$

# Theoretical framework

Effective Hamiltonian for  $c \rightarrow ul^+l^-$  transition in  $Z'$  model [7]

$$H_{Z'} = H_{Z'}^q + H_{Z'}^\ell$$

$$H_{Z'}^q = g_{Z'1} \bar{u}_L \gamma_\mu c_L Z'^\mu + g_{Z'2} \bar{u}_R \gamma_\mu c_R Z'^\mu + h.c.$$

$$H_{Z'}^\ell = g'_{Z'1} \bar{\ell}_L \gamma_\mu \ell_L Z'^\mu + g'_{Z'2} \bar{\ell}_R \gamma_\mu \ell_R Z'^\mu + h.c.$$

New Physics (NP)  
Wilson coefficients [8]

$$C_{9(10)}^{NP} = -\frac{\pi}{\sqrt{2}G_F\alpha} \frac{g_{Z'1}(g'_{Z'1} \pm g'_{Z'2})}{M_{Z'}^2}$$

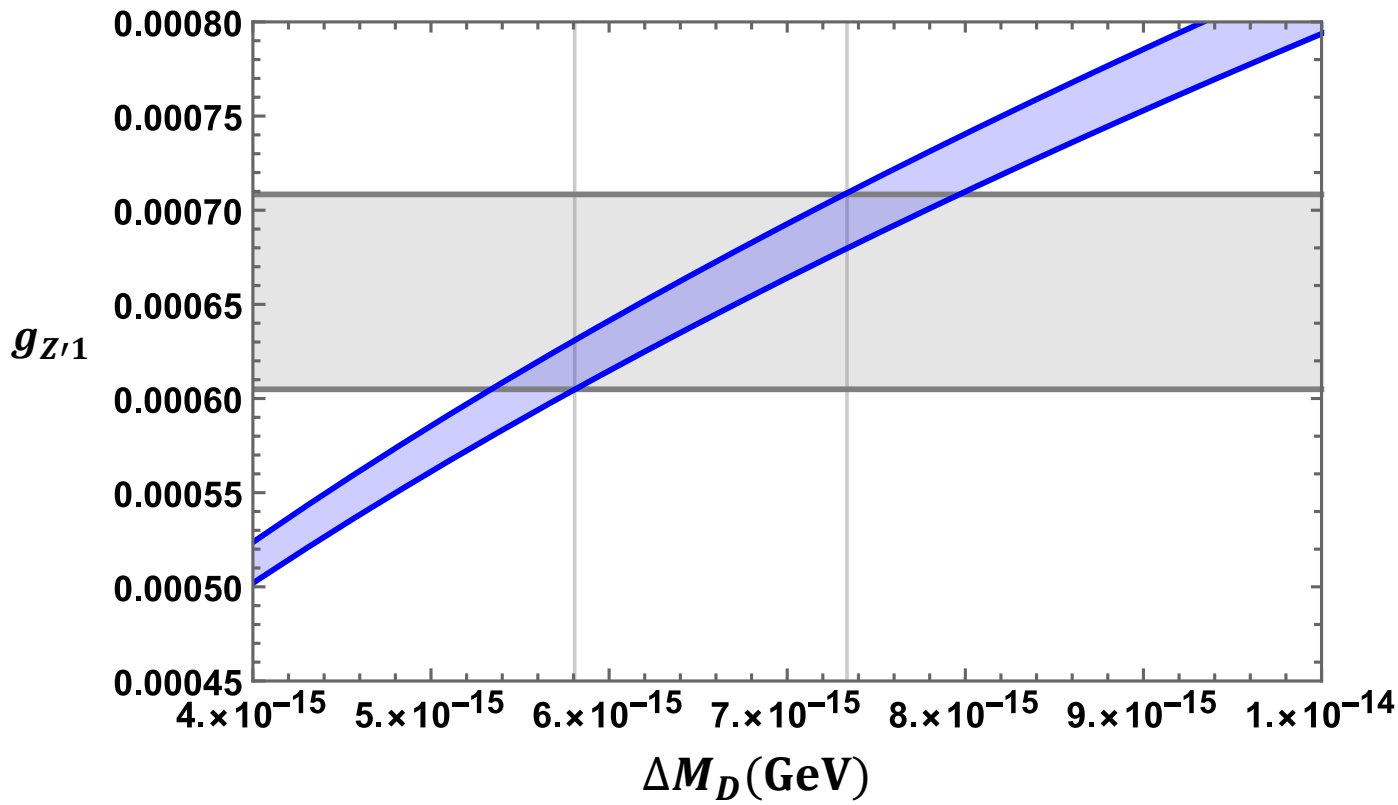
$$C'_{9(10)}{}^{NP} = -\frac{\pi}{\sqrt{2}G_F\alpha} \frac{g_{Z'2}(g'_{Z'1} \pm g'_{Z'2})}{M_{Z'}^2}$$



# Estimation of NP couplings in $Z'$ model

Quark coupling

$D^0 - \bar{D}^0$  mixing



$$\Delta M_D^{Z'} = \frac{f_D^2 M_D B_D r(m_c, M_{Z'})}{3} \frac{g_{Z'1}^2}{M_{Z'}^2} [7, 18]$$

$$\Delta M_D = (99.7 \pm 11.6) \times 10^8 \text{ } \hbar s^{-1} [19]$$

$$g_{Z'1} = 6.47 \times 10^{-4}$$

Fig 1. Determination of quark coupling

# Estimation of NP couplings in $Z'$ model

Leptonic coupling

$D^0 \rightarrow \mu^+ \mu^-$  process

$$\mathcal{B}(D^0 \rightarrow \mu^+ \mu^-)|_{Z'} = \tau_D \frac{f_D^2 m_\mu^2 M_D}{32\pi M_{Z'}^4} \sqrt{1 - \frac{4m_l^2}{M_D^2}} (g_{Z'1} - g_{Z'2})^2 (g'_{Z'2} - g'_{Z'2})^2 [7, 18]$$

$$g'_{Z'1} \leq 35.31$$

# Estimation of NP couplings in $Z'$ model

Decay mode	Experimental upper limit [19]
$D^+ \rightarrow \pi^+ \mu^+ \mu^-$	$6.7 \times 10^{-8}$
$D^0 \rightarrow \pi^0 \mu^+ \mu^-$	$1.8 \times 10^{-4}$
$D_s^+ \rightarrow K^+ \mu^+ \mu^-$	$1.4 \times 10^{-7}$
$D^+ \rightarrow \rho^+ \mu^+ \mu^-$	$5.6 \times 10^{-4}$
$D^0 \rightarrow \rho^0 \mu^+ \mu^-$	$2.2 \times 10^{-5}$

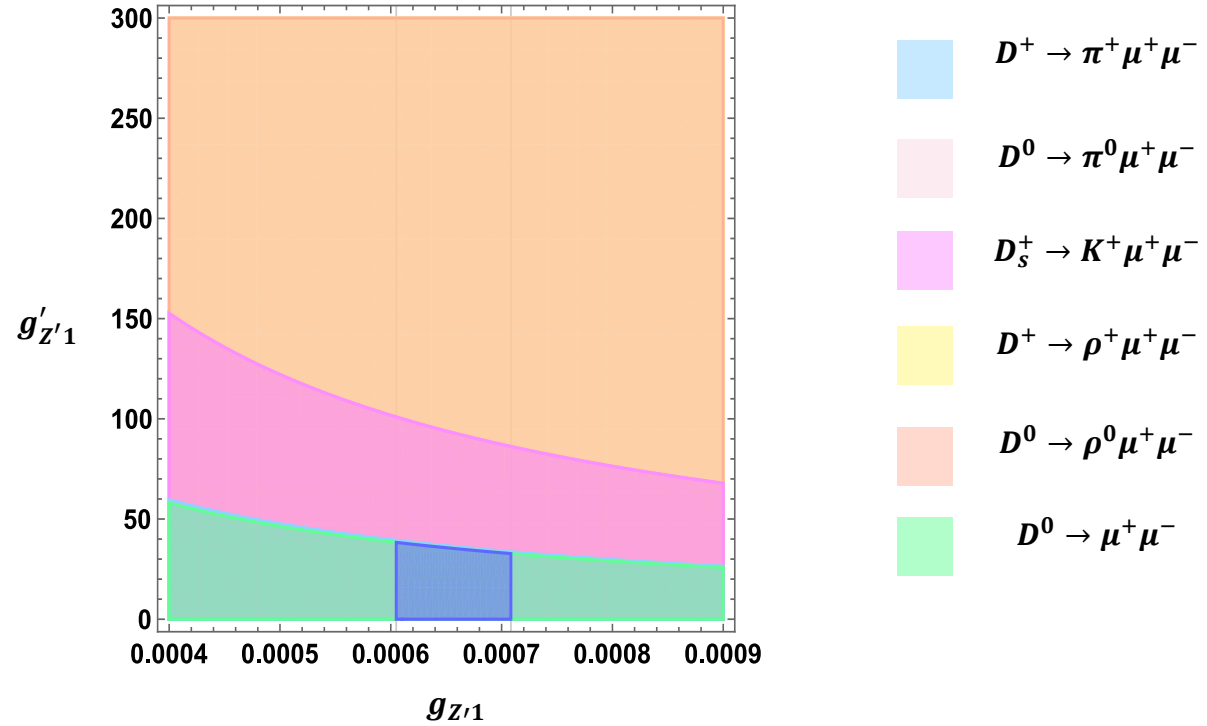


Fig. 2: Plot of leptonic coupling  $g'_{Z'1}$  vs quark coupling  $g_{Z'1}$ . The blue region shows the allowed region of the couplings, constrained from  $D^+ \rightarrow \pi^+ \mu^+ \mu^-$ ,  $D^0 \rightarrow \pi^0 \mu^+ \mu^-$ ,  $D_s^+ \rightarrow K^+ \mu^+ \mu^-$ ,  $D^+ \rightarrow \rho^+ \mu^+ \mu^-$ ,  $D^0 \rightarrow \rho^0 \mu^+ \mu^-$  and  $D^0 \rightarrow \mu^+ \mu^-$  decays

## Estimation of branching fraction

Upper limit ———  
 Z' model ———  
 SM ———

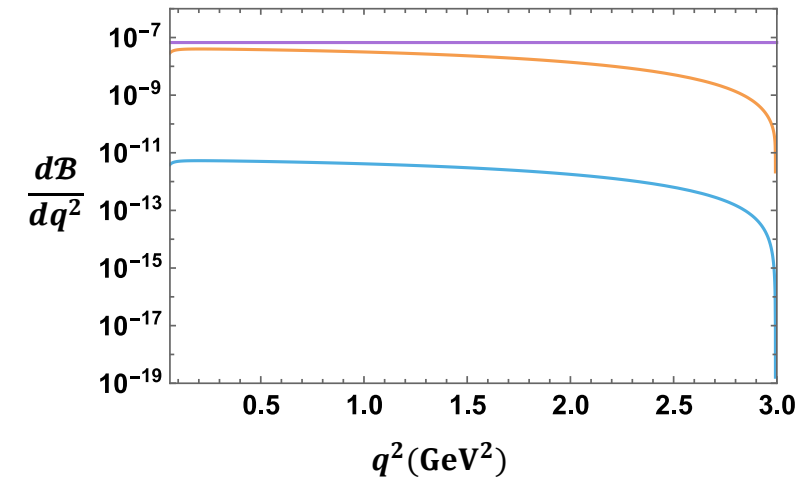


Fig. 3a

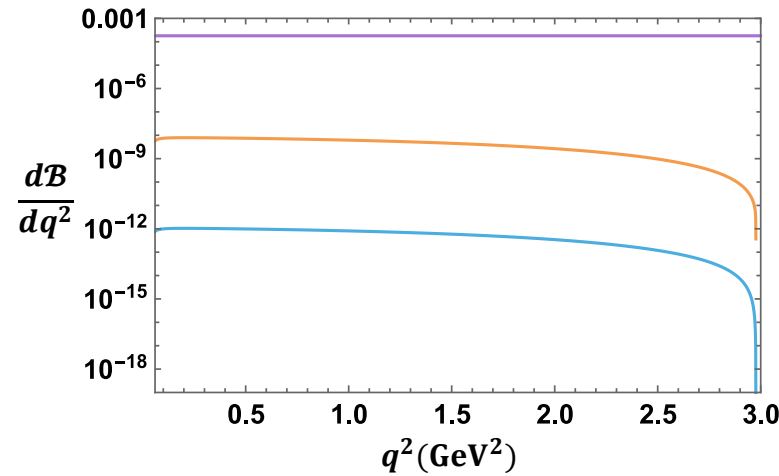


Fig. 3b

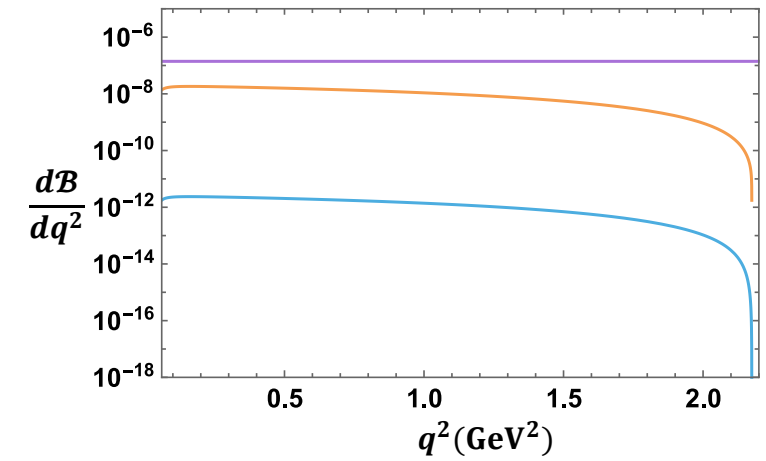


Fig. 3c

Fig. 3: Variation of differential branching fractions within allowed kinematic region for (a)  $D^+ \rightarrow \pi^+ \mu^+ \mu^-$  and (b)  $D^0 \rightarrow \pi^0 \mu^+ \mu^-$  (c)  $D_s^+ \rightarrow K^+ \mu^+ \mu^-$

## Estimation of branching fraction

Upper limit —  
 Z' model —  
 SM —

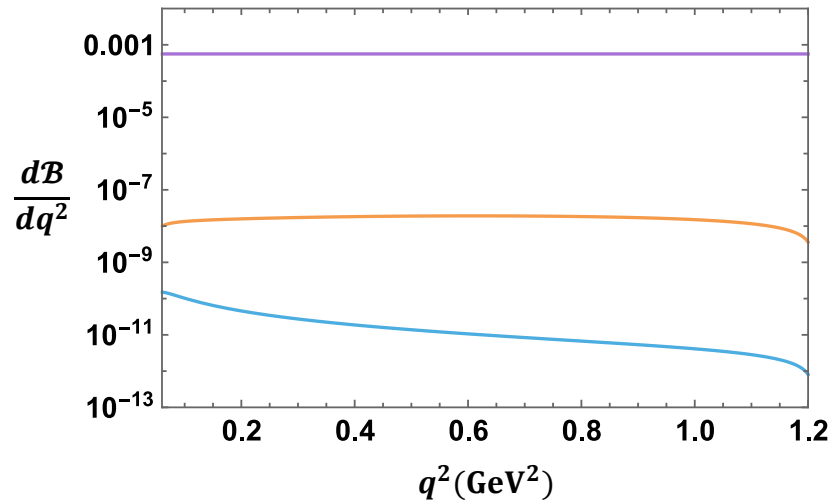


Fig. 4 (a)

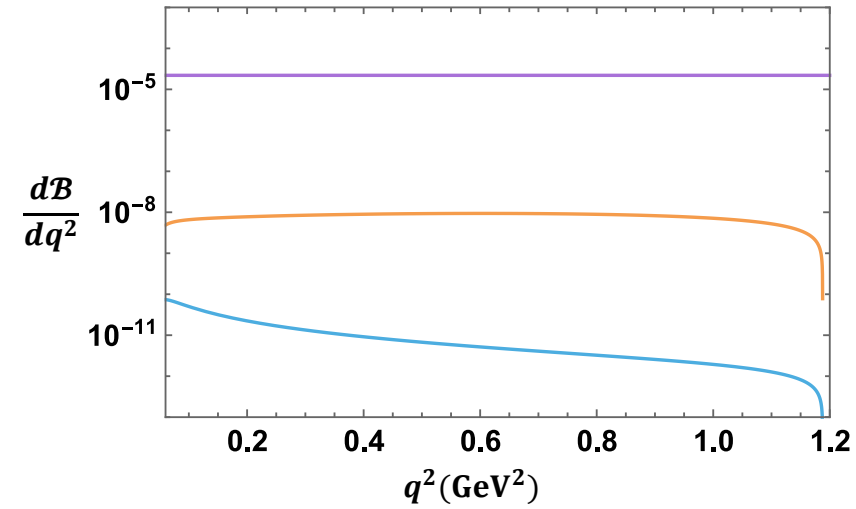


Fig. 4 (b)

Fig. 4: Variation of differential branching fractions within allowed kinematic region for  
 (a)  $D^+ \rightarrow \rho^+ \mu^+ \mu^-$  and (b)  $D^0 \rightarrow \rho^0 \mu^+ \mu^-$

## Estimation of forward backward asymmetry

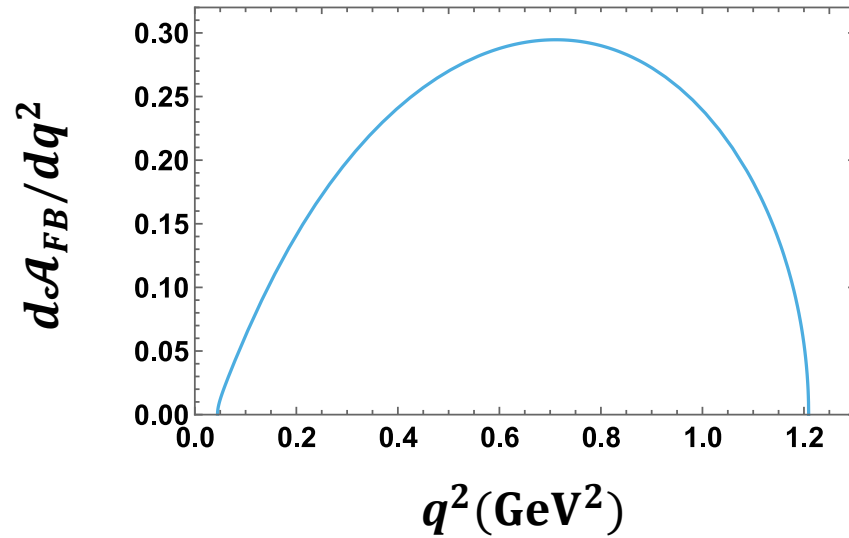


Fig. 5 (a)

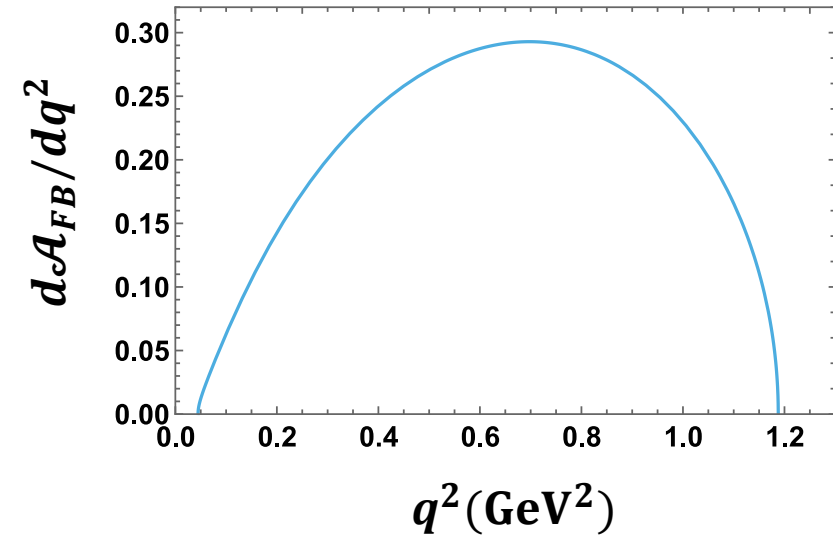


Fig. 5 (b)

Fig. 5: Variation of normalized forward backward asymmetry for (a)  $D^+ \rightarrow \rho^+ \mu^+ \mu^-$  and (b)  $D^0 \rightarrow \rho^0 \mu^+ \mu^-$  in non-universal  $Z'$  model.

# Results

**Table 1: Branching ratio values integrated in the allowed  $q^2$  region**

Decay mode	SM	$Z'$ model	Experimental upper limit [19]
$D^+ \rightarrow \pi^+ \mu^+ \mu^-$	$8.29 \times 10^{-12}$	$6.33 \times 10^{-8}$	$6.7 \times 10^{-8}$
$D^0 \rightarrow \pi^0 \mu^+ \mu^-$	$1.62 \times 10^{-12}$	$1.24 \times 10^{-8}$	$1.8 \times 10^{-4}$
$D_s^+ \rightarrow K^+ \mu^+ \mu^-$	$2.56 \times 10^{-12}$	$2.00 \times 10^{-8}$	$1.4 \times 10^{-7}$

**Table 2: Branching ratio values integrated in the allowed  $q^2$  region**

Decay mode	SM	$Z'$ model	Experimental upper limit [19]
$D^+ \rightarrow \rho^+ \mu^+ \mu^-$	$2.39 \times 10^{-11}$	$1.85 \times 10^{-8}$	$5.6 \times 10^{-4}$
$D^0 \rightarrow \rho^0 \mu^+ \mu^-$	$1.16 \times 10^{-11}$	$8.92 \times 10^{-8}$	$2.2 \times 10^{-5}$

**Table 3: Normalized forward backward asymmetry values integrated in the allowed  $q^2$  region**

Decay mode	Normalized forward backward asymmetry in $Z'$ model
$D^+ \rightarrow \rho^+ \mu^+ \mu^-$	0.252
$D^0 \rightarrow \rho^0 \mu^+ \mu^-$	0.246

## Conclusions

- We have calculated the differential branching ratio and forward backward asymmetry for the decays  $D^+ \rightarrow \pi^+ \mu^+ \mu^-$ ,  $D^0 \rightarrow \pi^0 \mu^+ \mu^-$ ,  $D_s^+ \rightarrow K^+ \mu^+ \mu^-$ ,  $D^+ \rightarrow \rho^+ \mu^+ \mu^-$  and  $D^0 \rightarrow \rho^0 \mu^+ \mu^-$  in non-universal  $Z'$  model.
- The branching fractions for all the decay channels studied here show a significant enhancement in the  $Z'$  model.
- Although forward backward asymmetry is zero for  $D \rightarrow P \mu^+ \mu^-$ , the  $D \rightarrow \rho \mu^+ \mu^-$  decays show a variation of the forward backward asymmetry in the  $Z'$  model.
- The results in this study indicate that these decays are sensitive to new physics contributions in the  $Z'$  model. We expect that our study will contribute more to the present understanding of charm decays.



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