

Possible New Physics on $D_s^+ \rightarrow \eta^{(\prime)} l^+ \nu_l$ Decays in Scalar Leptoquark Model

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Motivation

B Anomalies for the $b \rightarrow cl\nu_l$ decays

Observable	Exp. Result	SM Prediction	Deviation
R_D	$0.339 \pm 0.026 \pm 0.014$ HFLAG [1]	0.299 ± 0.003 [1]	1.4σ
R_{D^*}	$0.295 \pm 0.010 \pm 0.010$ HFLAG [1]	0.254 ± 0.005 [1]	2.8σ
$R_{J/\psi}$	$0.71 \pm 0.17 \pm 0.18$ [2]	0.283 ± 0.048 [3]	2σ
$P_{\tau D^*}$	$-0.38 \pm 0.51_{-0.16}^{+0.21}$ [7, 8]	-0.497 ± 0.013 [4]	–
F_{LD^*}	$0.60 \pm 0.08 \pm 0.035$ [9]	0.46 ± 0.04 [5]	1.7σ
R_{X_c}	0.223 ± 0.030 [6]	0.216 ± 0.003 [6]	–

Is there any possibility to find such anomalous results in charm decays?

Therefore, there are some possibilities to exist some NP. Several theoretical efforts [10, 13, 14] have been done recently to find the NP contribution in D meson decays.

Decay	SM result [10]	Experiment [11]	P_{ul}_{SM}
$\mathcal{B}(D_s^+ \rightarrow \eta\mu^+\nu_\mu) \times 10^{-2}$	1.52 ± 0.31	2.4 ± 0.5	1.5σ
		$2.215 \pm 0.051 \pm 0.052$ [12]	2.2σ
$\mathcal{B}(D_s^+ \rightarrow \eta'\mu^+\nu_\mu) \times 10^{-3}$	5.64 ± 1.10	11.0 ± 5.0	1.05σ
		$8.01 \pm 0.55 \pm 0.31$ [12]	1.9σ
$\mathcal{B}(D_s^+ \rightarrow \eta e^+\nu_e) \times 10^{-2}$	1.55 ± 0.33	2.29 ± 0.19	2σ
$\mathcal{B}(D_s^+ \rightarrow \eta' e^+\nu_e) \times 10^{-3}$	5.91 ± 1.26	7.4 ± 1.4	0.8σ

[11] R. L. Workman et al. (Particle Data Group), *Prog. Theor. Exp. Phys.* **2022**, 083C01 (2022).

[12] M. Ablikim et al., (BESIII Collaboration), [arXiv:2307.12852 [hep-ex]] (2023).



Theoretical Framework

Considering all possible Lorentz structures and assuming only left-handed neutrinos, the general effective Lagrangian for the $c \rightarrow q \bar{l} \nu_l$ transitions can be written as [10, 13]

$$\mathcal{L}_{eff} = -\frac{4G_F}{\sqrt{2}} V_{cq}^* \left[(1 + C_{V_L}^l) O_{V_L}^l + C_{V_R}^l O_{V_R}^l + C_{S_L}^l O_{S_L}^l + C_{S_R}^l O_{S_R}^l + C_T^l O_T^l \right] + h.c, \quad (1)$$

The two-fold differential angular decay distribution of $D \rightarrow Pl^+ \nu_l$ ($P = \eta, \eta'$) decay can be expressed as

$$\frac{d^2\Gamma(D \rightarrow Pl^+ \nu_l)}{dq^2 d\cos\theta_l} = \frac{G_F^2 |V_{cq}|^2 \sqrt{Q_+ Q_-}}{256\pi^3 m_D^3} \left(1 - \frac{m_l^2}{q^2}\right)^2 \left[q^2 A_1^P + \sqrt{q^2} m_l A_2^P + m_l^2 A_3^P \right], \quad (2)$$

with

$$A_1^P = |C_{S_L} + C_{S_R}|^2 |H^{PS}|^2 + \text{Re}[(C_{S_L} + C_{S_R})C_T^*] H^{PS} (H_{0,t}^{PT} + H_{0,-1}^{PT}) \cos\theta_l + 4|C_T|^2 |H_{0,t}^{PT} + H_{1,-1}^{PT}|^2 \cos^2 \theta_l + |1 + C_{V_L} + C_{V_R}|^2 |H_0^{PV}|^2 \sin^2 \theta_l, \quad (3)$$

Theoretical Framework

$$A_2^P = 2\{Re[(C_{S_L} + C_{S_R})(1 + C_{V_L} + C_{V_R})^*]H^{PS}H_t^{PV} - 2Re[C_T(1 + C_{V_L} + C_{V_R})^*]H_0^{PV}(H_{0,t}^{PT} + H_{1,-1}^{PT})\} - 2\{Re[(C_{S_L} + C_{S_R})(1 + C_{V_L} + C_{V_R})^*]H^{PS}H_0^{PV} - 2Re[C_T(1 + C_{V_L} + C_{V_R})^*]H_t^{PV}(H_{0,t}^{PT} + H_{1,-1}^{PT})\}cos\theta_l, \quad (4)$$

$$A_3^P = 4|C_T|^2 |H_{0,t}^{PT} + H_{1,-1}^{PT}|^2 \sin^2 \theta_l + |1 + C_{V_L} + C_{V_R}|^2 (|H_0^{PV}|^2 \cos^2 \theta_l - 2H_0^{PV}H_t^{PV} \cos\theta_l + |H_t^{PV}|^2), \quad (5)$$

where θ_l is the angle between the charged lepton and opposite direction of the motion of the final meson in the virtual W^{*+} rest frame and $M_{\pm} = m_D \pm m_P$ and $Q_{\pm} = M_{\pm}^2 - q^2$.

Theoretical Framework

The $D \rightarrow Pl^+ \nu_l$ decays have contributions through only five helicity amplitudes, which are given by

$$H_0^{PV} = \frac{f_+ \sqrt{Q_+ Q_-}}{\sqrt{q^2}}, \quad H_t^{PV} = \frac{f_0 M_+ M_-}{\sqrt{q^2}}, \quad H^{PS} = \frac{f_0 M_+ M_-}{m_c - m_q},$$

$$H_{0,t}^{PT} = -\frac{f_T \sqrt{Q_+ Q_-}}{M_+}, \quad H_{1,-1}^{PT} = -\frac{f_T \sqrt{Q_+ Q_-}}{M_+}. \quad (6)$$

In this work, we have used the form factors obtained from light cone sum rules (LCSR) [15], which are parameterized as

$$F^i(q^2) = \frac{F^i(0)}{1 - a \frac{q^2}{m_D^2} + b \left(\frac{q^2}{m_D^2} \right)^2}. \quad (7)$$

Table: Form factors for $D \rightarrow \eta_q$ and $D_s \rightarrow \eta_s$ [15].

Decay	$D \rightarrow \eta_q$		$D_s \rightarrow \eta_s$	
	f_+	f_0	f_+	f_0
$F(0)$	$0.56_{-0.05}^{+0.06}$	$0.56_{-0.05}^{+0.06}$	$0.61_{-0.05}^{+0.06}$	$0.61_{-0.05}^{+0.06}$
a	$1.25_{+0.05}^{-0.04}$	$0.65_{+0.02}^{-0.01}$	$1.20_{+0.03}^{-0.02}$	$0.64_{+0.02}^{-0.01}$
b	$0.42_{+0.05}^{-0.06}$	$-0.22_{+0.02}^{-0.03}$	$0.38_{+0.01}^{-0.01}$	$-0.18_{-0.03}^{+0.04}$

Theoretical Framework

Scalar Leptoquark Model

To describe the $c \rightarrow q \bar{l} \nu_l$ transitions through LQ as mediator, the Lagrangian can be defined as [16]

$$\mathcal{L}_{int}^{\phi} \supset \bar{Q}_L^c \lambda^L i \tau_2 L \phi^* + \bar{u}_R^c \lambda^R l_R \phi^* + h.c.. \quad (8)$$

In terms of Yukawa coupling the NP Wilson coefficients are defined as

$$\begin{aligned} C_{V_L}(M_{\phi}) &= \frac{\lambda_{c\nu_l}^L \lambda_{ql}^{L*}}{4\sqrt{2} G_F V_{cq} M_{\phi}^2}, \\ C_{S_L}(M_{\phi}) &= -\frac{\lambda_{c\nu_l}^L \lambda_{ql}^{R*}}{4\sqrt{2} G_F V_{cq} M_{\phi}^2}, \\ C_T(M_{\phi}) &= -\frac{1}{4} S_L(M_{\phi}). \end{aligned} \quad (9)$$

CMS Collaboration [17] has predicted the mass of Leptoquark to be in the range of 0.98 – 1.73 TeV.

[16] X.-L. Mu, Y. Li, Z.-T. Zou and B. Zhu, *Phys. Rev. D* **100**, 113004 (2019) [arXiv:1909.10769 [hep-ph]].

[17] A. M Sirunyan *et al.* (CMS Collaboration), *Phys. Lett. B* 819, 136446 (2021) [arXiv:2012.04178 [hep-ex]].

Results

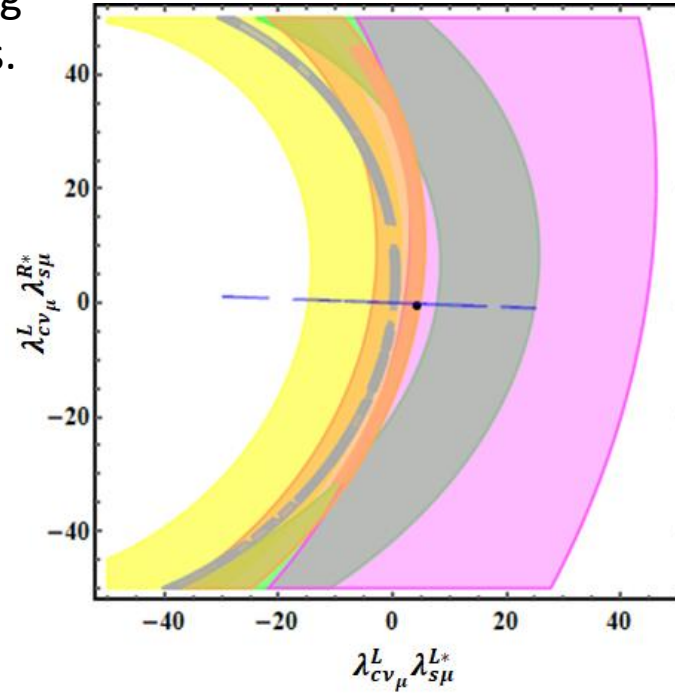
Table: the SM and experimental results of branching fraction of semileptonic and leptonic D meson decays.

$$c \rightarrow s\mu^+\nu_\mu$$

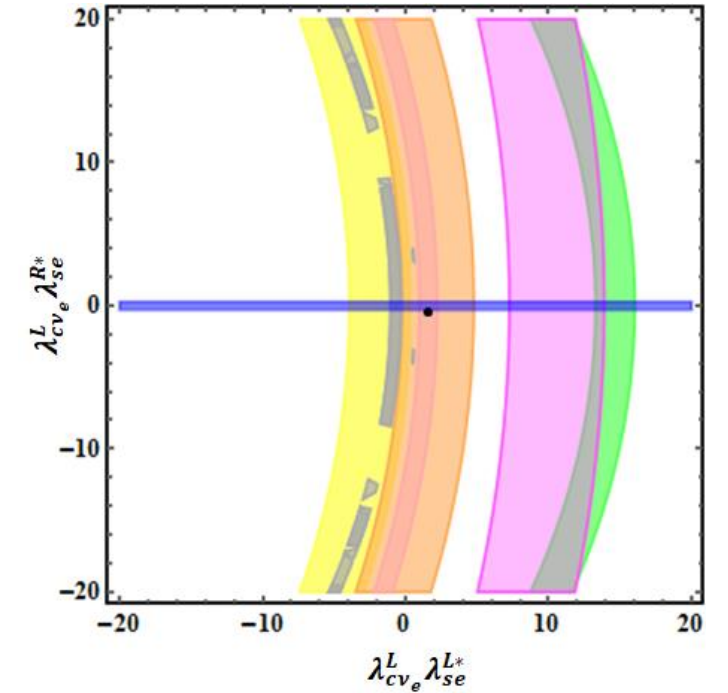
Decay	SM Result [10]	Exp. Result [11]
$D^0 \rightarrow K^-\mu^+\nu_\mu$	$(3.40 \pm 0.22) \times 10^{-2}$	$(3.41 \pm 0.04) \times 10^{-2}$
$D^+ \rightarrow \bar{K}^0\mu^+\nu_\mu$	$(8.69 \pm 0.57) \times 10^{-2}$	$(8.76 \pm 0.19) \times 10^{-2}$
$D^0 \rightarrow K^{*-}\mu^+\nu_\mu$	$(1.81 \pm 0.16) \times 10^{-2}$	$(1.89 \pm 0.24) \times 10^{-2}$
$D^+ \rightarrow \bar{K}^{*0}\mu^+\nu_\mu$	$(4.71 \pm 0.42) \times 10^{-2}$	$(5.27 \pm 0.15) \times 10^{-2}$
$D_s^+ \rightarrow \phi\mu^+\nu_\mu$	$(2.33 \pm 0.40) \times 10^{-2}$	$(1.90 \pm 0.50) \times 10^{-2}$
$D_s^+ \rightarrow \eta\mu^+\nu_\mu$	$(1.52 \pm 0.31) \times 10^{-2}$	$(2.4 \pm 0.5) \times 10^{-2}$
$D_s^+ \rightarrow \eta'\mu^+\nu_\mu$	$(5.64 \pm 1.10) \times 10^{-3}$	$(11.0 \pm 5.0) \times 10^{-3}$
$D_s^+ \rightarrow l^+\nu_\mu$	$(5.28 \pm 0.08) \times 10^{-3}$	$(5.43 \pm 0.15) \times 10^{-3}$

$$c \rightarrow se^+\nu_e$$

Decay	SM Result [10]	Exp. Result [11]
$D^0 \rightarrow K^-e^+\nu_e$	$(3.49 \pm 0.23) \times 10^{-2}$	$(3.542 \pm 0.0035) \times 10^{-2}$
$D^+ \rightarrow \bar{K}^0e^+\nu_e$	$(8.92 \pm 0.59) \times 10^{-2}$	$(8.73 \pm 0.10) \times 10^{-2}$
$D^0 \rightarrow K^{*-}e^+\nu_e$	$(1.92 \pm 0.17) \times 10^{-2}$	$(2.15 \pm 0.16) \times 10^{-2}$
$D^+ \rightarrow \bar{K}^{*0}e^+\nu_e$	$(4.98 \pm 0.45) \times 10^{-2}$	$(5.40 \pm 0.10) \times 10^{-2}$
$D_s^+ \rightarrow \phi e^+\nu_e$	$(2.46 \pm 0.42) \times 10^{-2}$	$(2.39 \pm 0.16) \times 10^{-2}$
$D_s^+ \rightarrow \eta e^+\nu_e$	$(1.55 \pm 0.33) \times 10^{-2}$	$(2.29 \pm 0.19) \times 10^{-2}$
$D_s^+ \rightarrow \eta' e^+\nu_e$	$(5.91 \pm 1.26) \times 10^{-3}$	$(7.4 \pm 1.4) \times 10^{-3}$
$D_s^+ \rightarrow e^+\nu_e$	$(1.24 \pm 0.02) \times 10^{-7}$	$< 8.3 \times 10^{-5}$



(a)



(b)

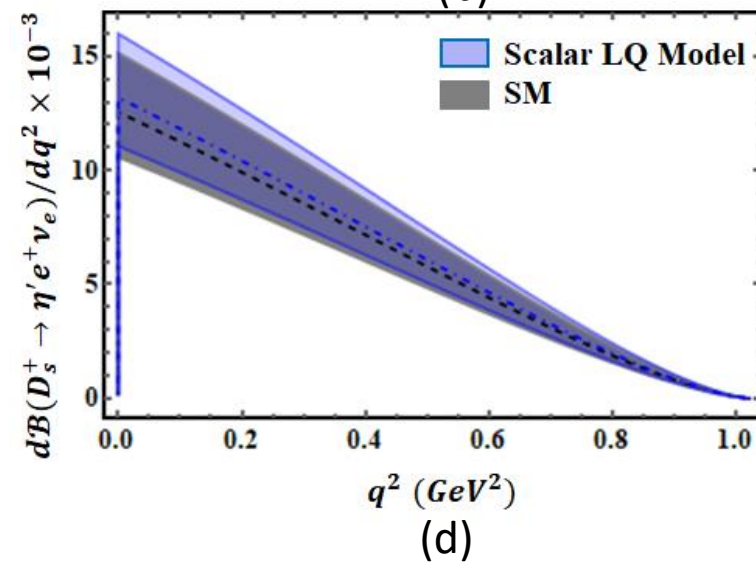
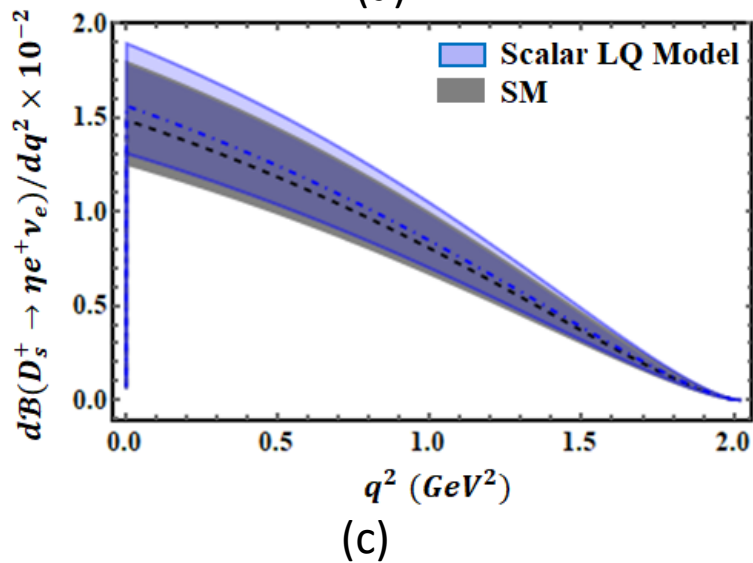
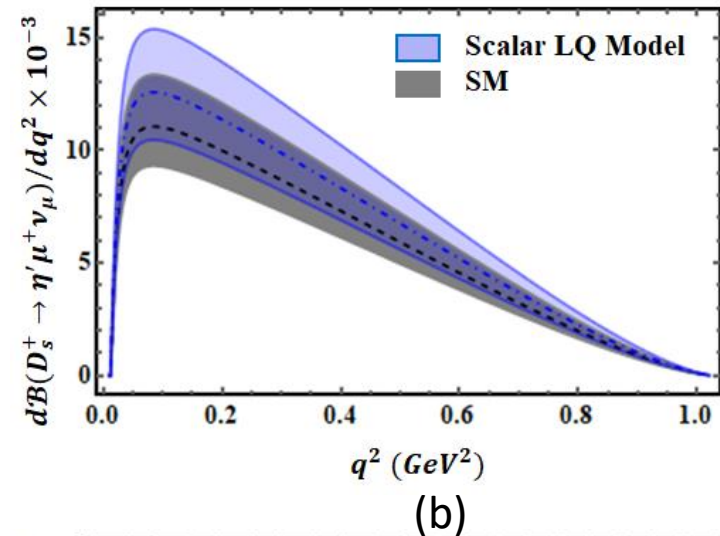
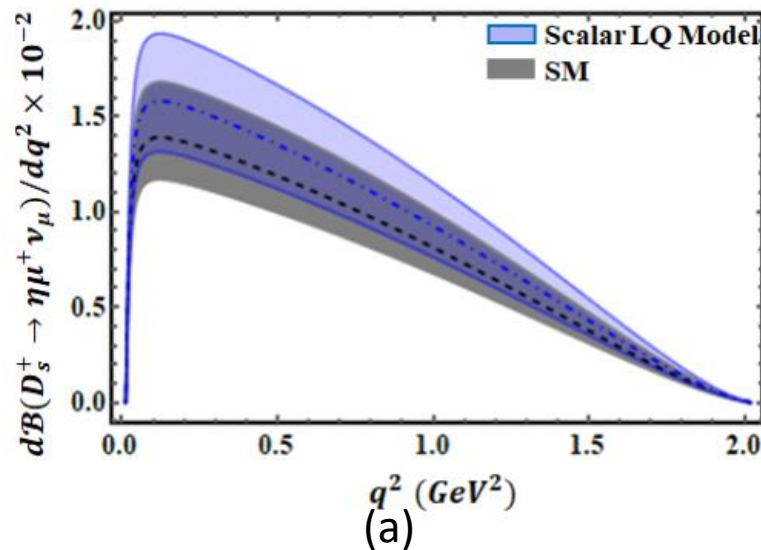
Fig: Allowed parameter space in the $(\lambda_{cv_l}^L, \lambda_{sl}^{L*}, \lambda_{cv_l}^L, \lambda_{sl}^{R*})$ plane for (a) $c \rightarrow s\mu^+\nu_\mu$ and (b) $c \rightarrow se^+\nu_e$ transitions. Here, colour notations are: Green ($D_s^+ \rightarrow \eta$), Magenta ($D_s^+ \rightarrow \eta'$), Yellow ($D_s^+ \rightarrow \phi$), Pink ($D^+ \rightarrow \bar{K}^{*0}$), Orange ($D^0 \rightarrow K^{*-}$), Grey ($D^0 \rightarrow K^-$ and $D^+ \rightarrow \bar{K}^0$) and Blue ($D_s^+ \rightarrow l^+\nu_l$). Black dot represents the fitted point of coupling parameters.

Results

Table: Values of the scalar LQ coupling parameters.

Current	Coupling Parameter	Mid-value
$c \rightarrow s\mu^+\nu_\mu$	$\lambda_{c\nu_\mu}^L \lambda_{s\mu}^{L*}$	4.30
	$\lambda_{c\nu_\mu}^L \lambda_{s\mu}^{R*}$	-0.47
$c \rightarrow se^+\nu_e$	$\lambda_{c\nu_e}^L \lambda_{se}^{L*}$	1.63
	$\lambda_{c\nu_e}^L \lambda_{se}^{R*}$	-0.44

Fig: Distribution of branching fraction.



Results

Table: Values of the branching fractions in scalar LQ model and their corresponding pull from the experimental results.

Decay	Scalar LQ Model	Experiment [11]	$pull_{SM}$	$Pull_{LQ}$
$\mathcal{B}(D_s^+ \rightarrow \eta \mu^+ \nu_\mu) \times 10^{-2}$	1.73 ± 0.41	2.4 ± 0.5	1.5σ	1.0σ
		$2.215 \pm 0.051 \pm 0.052$ [12]	2.2σ	1.1σ
$\mathcal{B}(D_s^+ \rightarrow \eta' \mu^+ \nu_\mu) \times 10^{-3}$	6.42 ± 1.47	11.0 ± 5.0	1.05σ	0.9σ
		$8.01 \pm 0.55 \pm 0.31$ [12]	1.9σ	0.8σ
$\mathcal{B}(D_s^+ \rightarrow \eta e^+ \nu_e) \times 10^{-2}$	1.62 ± 0.38	2.29 ± 0.19	2σ	1.5σ
$\mathcal{B}(D_s^+ \rightarrow \eta' e^+ \nu_e) \times 10^{-3}$	6.21 ± 1.36	7.4 ± 1.4	0.8σ	0.6σ

[11] R. L. Workman et al. (Particle Data Group), *Prog. Theor. Exp. Phys.* **2022**, 083C01 (2022).

[12] M. Ablikim et. al., (BESIII Collaboration), [arXiv:2307.12852 [hep-ex]] (2023).

Results

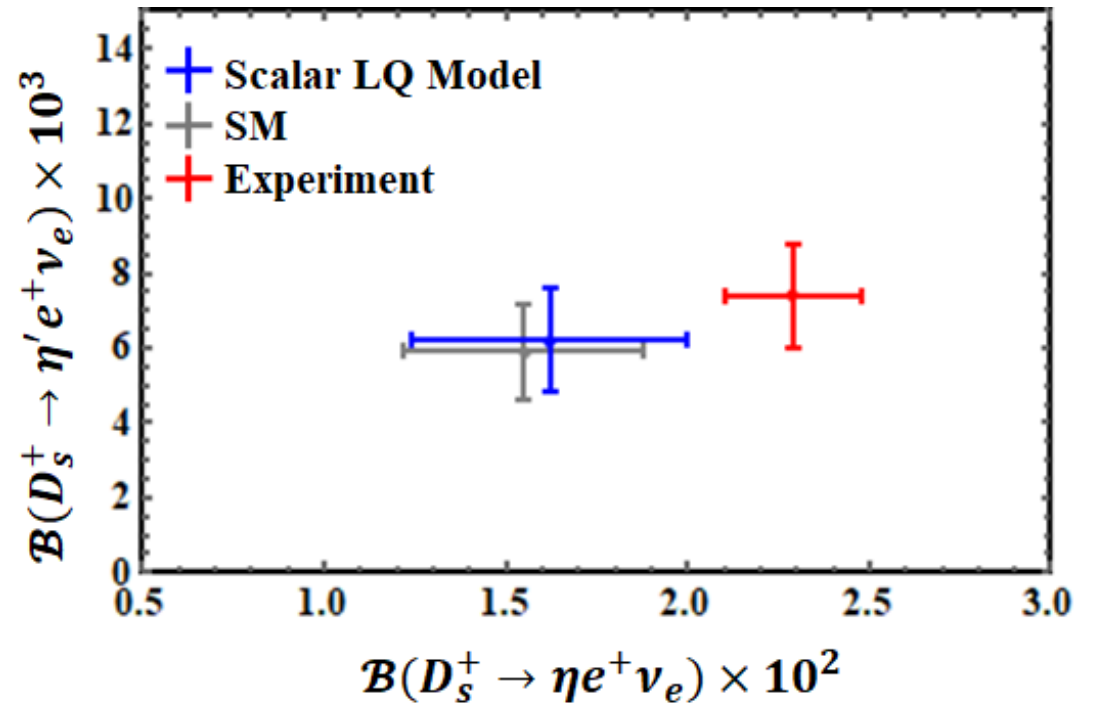
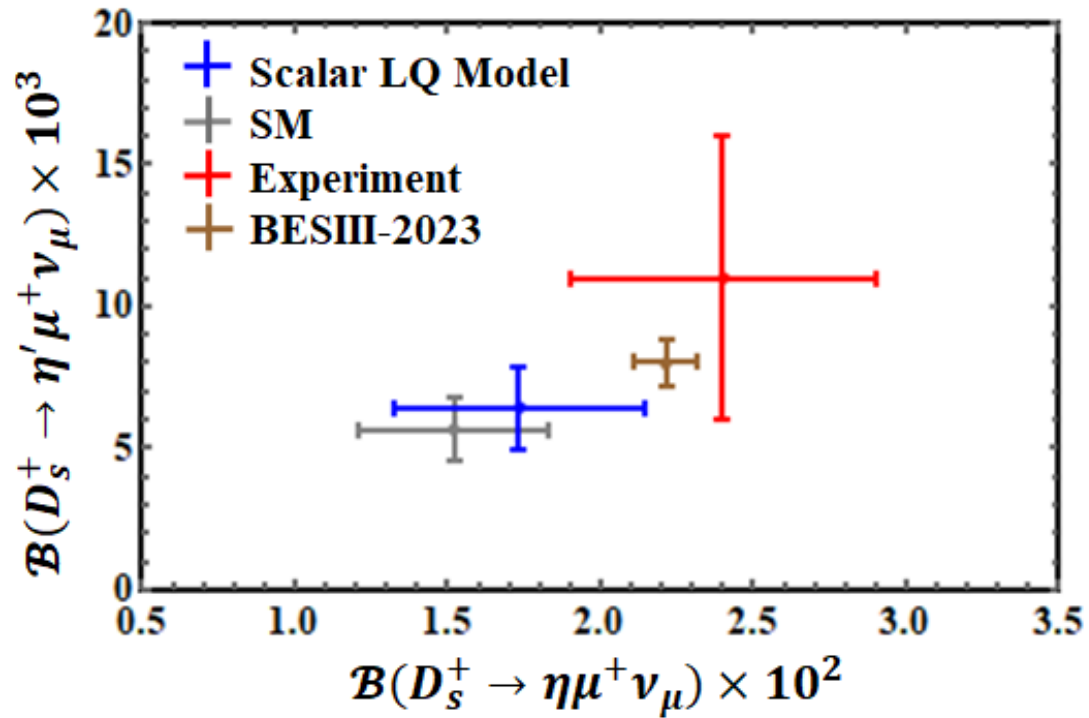


Fig: Comparison of the scalar LQ model results with the experiment and SM in 2D ($\mathcal{B}(D_s^+ \rightarrow \eta l^+ \nu_l), \mathcal{B}(D_s^+ \rightarrow \eta' l^+ \nu_l)$) plane.

Conclusion

- Our predicted branching fraction results are found round 1σ range of the experimental results excluding $D_s^+ \rightarrow \eta e^+ \nu_e$ decay.
- We have also found good consistency with the BESIII 2023 results.
- The effectiveness of our models is clearly observed through the comparison of our predicted results with the experimental results and SM predictions in 2D branching fraction plane.
- The scalar LQ model is reliable to study $c \rightarrow s \bar{l} \nu_l$ transitions in the framework of NP.
- We will be able to better comprehend the existence of NP with the help of the impending measurement of D meson decays in the BESIII and Belle II experiments. We expect that our predicted values in scalar LQ model will be closer to the outcomes of upcoming experiments.

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Thank You