



1



FPCP2024

Chulalongkorn University, Bangkok

28th May 2024

Asymmetric Self-interacting Dark Matter via Dirac Leptogenesis

Dr. Manoranjan Dutta
Department Of Physics
North Lakhimpur University



COLD & COLLISIONLESS DARK MATTER: Λ CDM

2



- ▶ CMB Anisotropy
- ▶ Bullet Cluster
- ▶ Gravitational Lensing
- ▶ Galaxy Rotation Curves
- ▶ Example: WIMPs

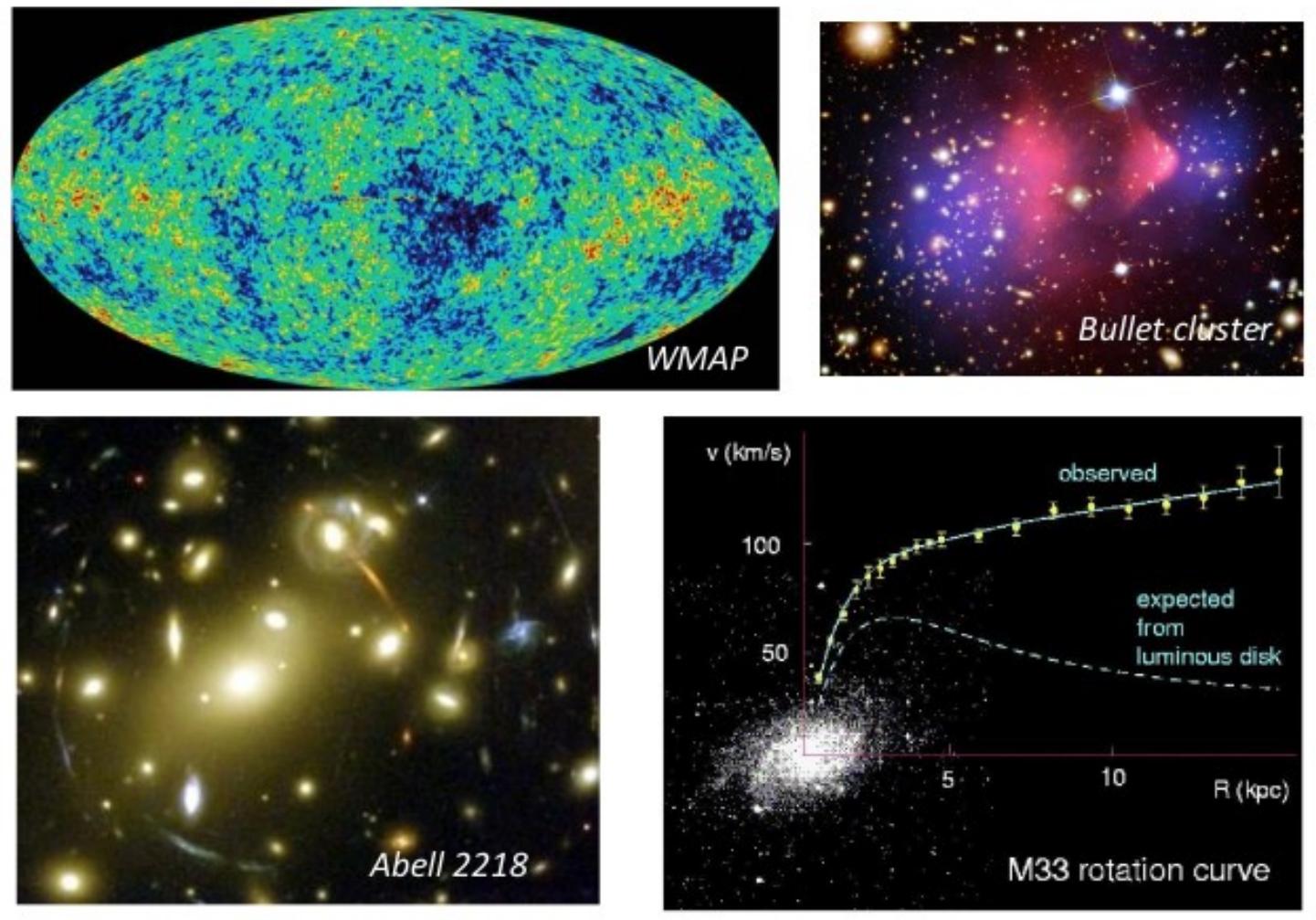


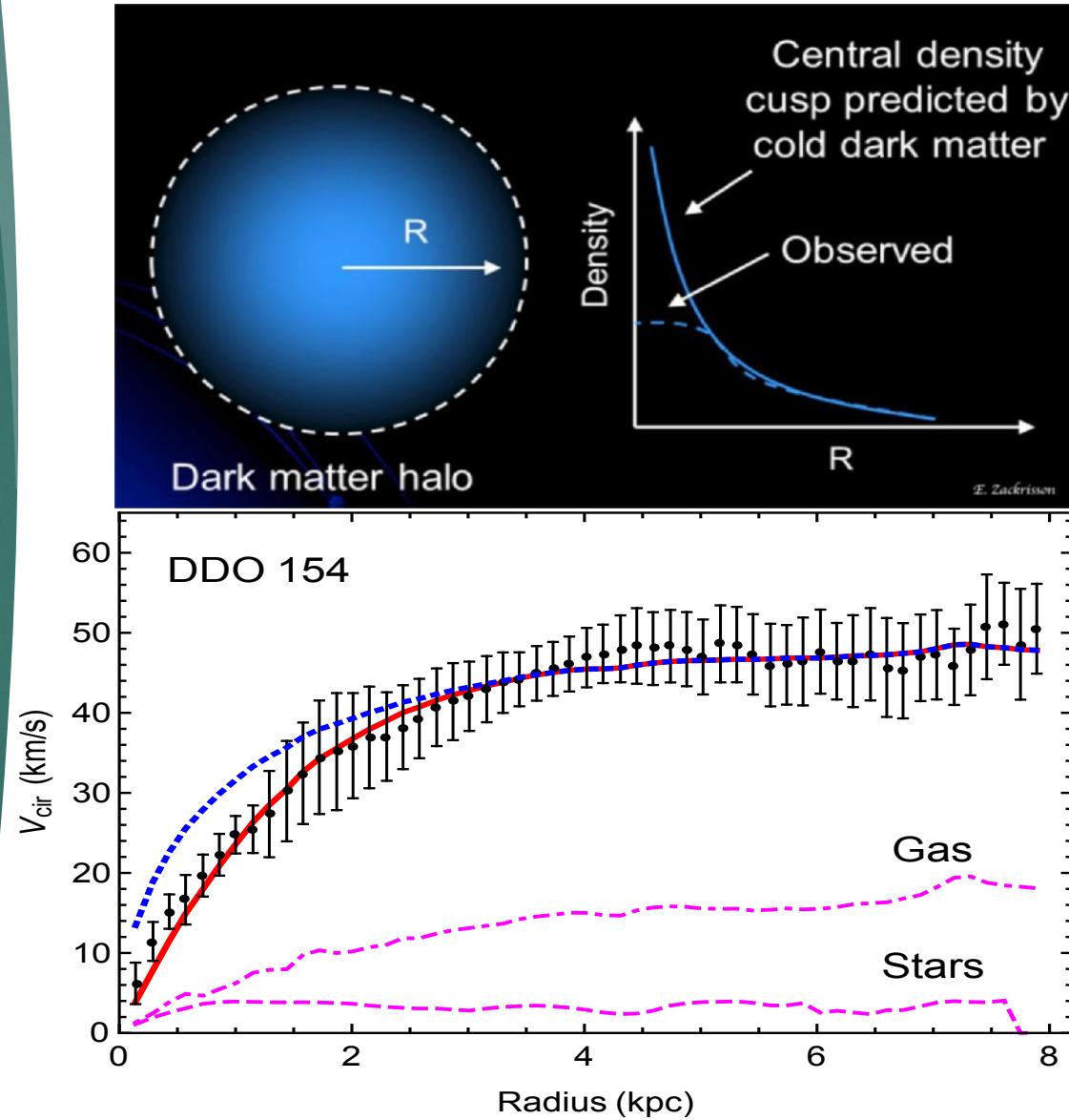
Image Courtesy: Sean Tulin, Frascati'15



COLD DARK MATTER: DISADVANTAGES

- ▶ Cusp-Core problem
- ▶ Missing Satellite Problem
- ▶ Too-Big-to-Fail Problem
- ▶ Prominent at Dwarf Scale
- ▶ Less prominent at Cluster scale

Image Courtesy: S Tulin, Frascati'15 & arXiv:1705.02358

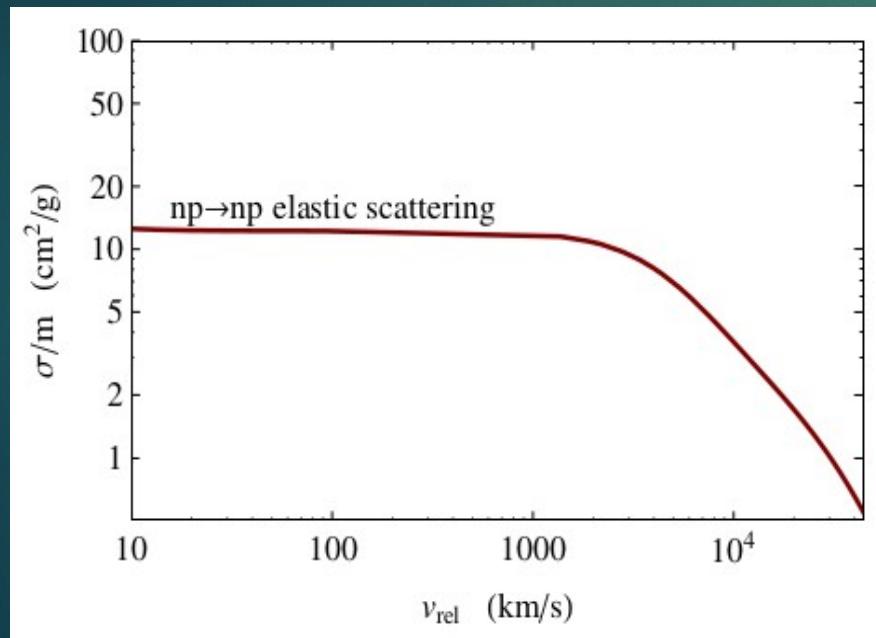




Self-Interacting Dark Matter

$$\text{SIDM: } \frac{\sigma}{m} = 1 \text{ cm}^2/\text{g} = 10^{-24} \text{ cm}^2/\text{GeV}$$

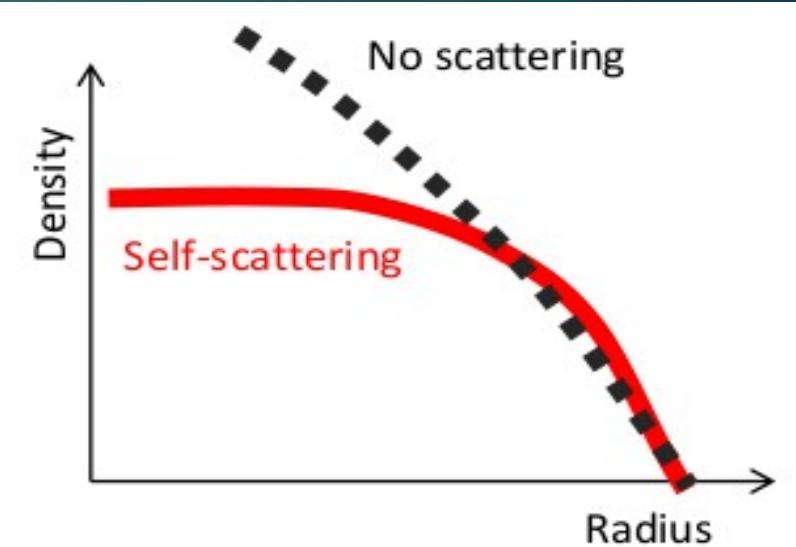
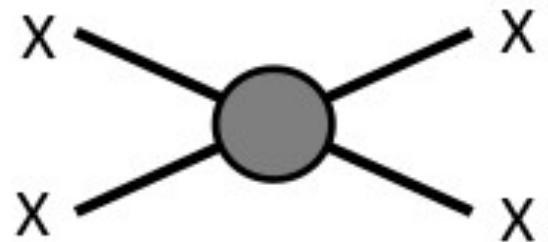
$$\text{WIMP: } \frac{\sigma}{m} = 10^{-38} \text{ cm}^2/\text{GeV}$$



GeV Scale DM
With
MeV scale Mediator

Velocity-dependent
Cross-section

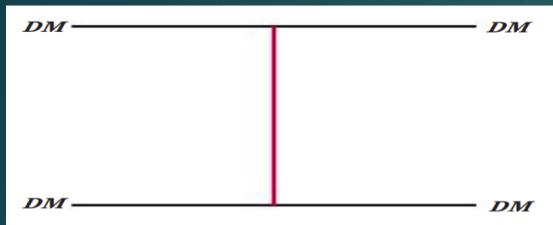
Dark matter self-scattering



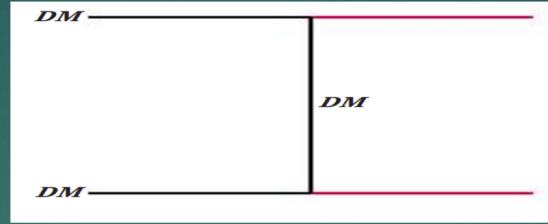


SIDM with Light Mediators

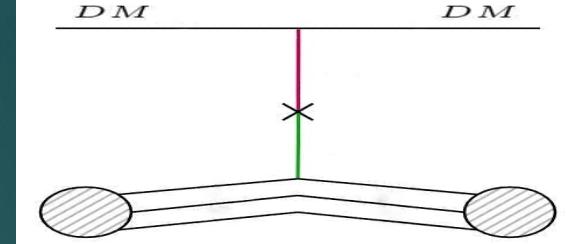
5



Self-scattering

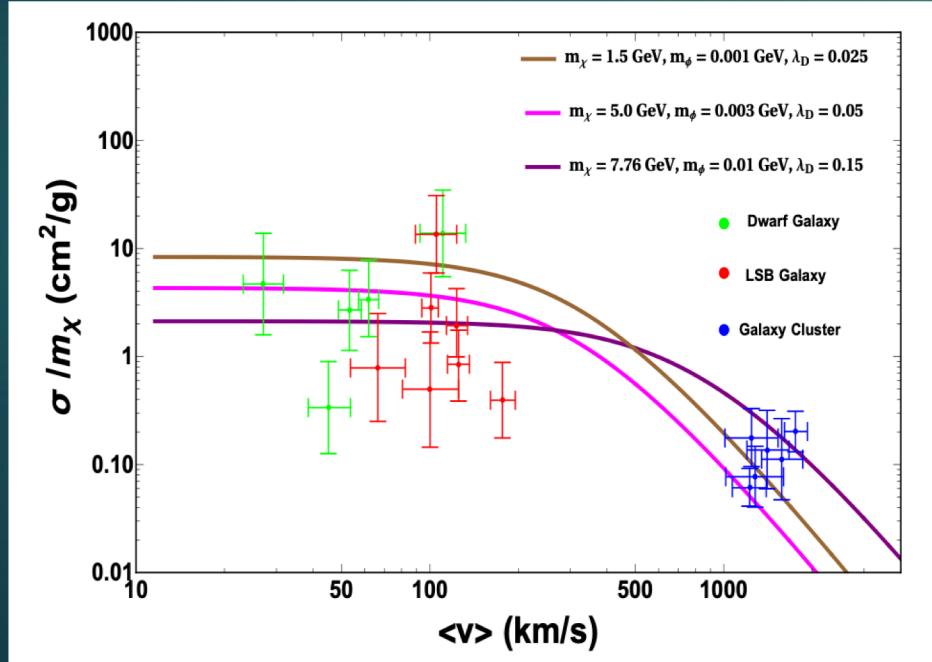


Relic Density

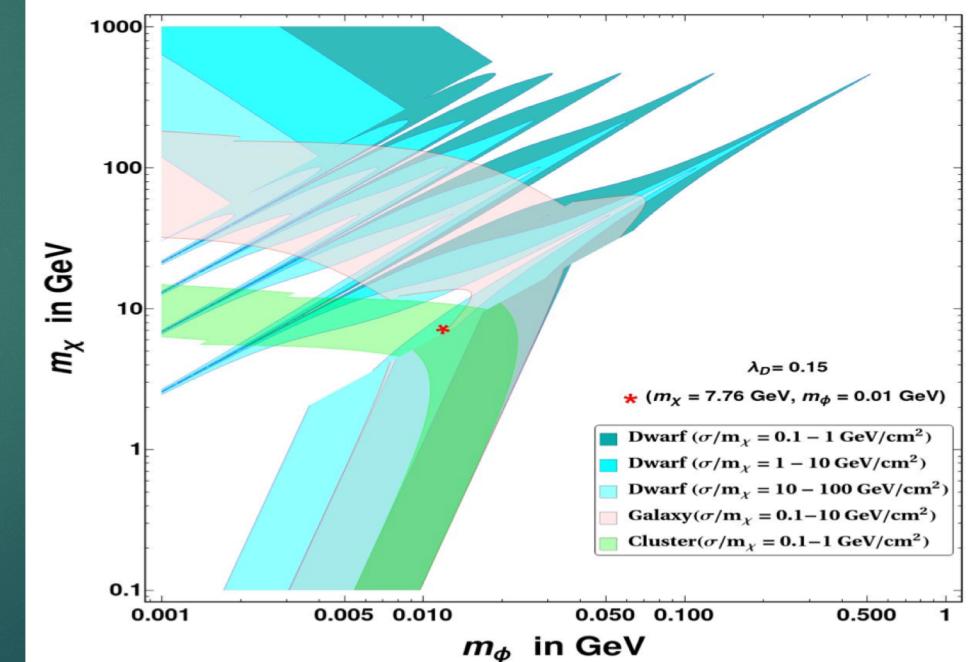


Direct Search

$$\mathcal{L}_{\text{int}} = \begin{cases} g_D \bar{\chi} \gamma^\mu \chi \phi_\mu & (\text{vector mediator}) \\ g_D \bar{\chi} \chi \phi & (\text{scalar mediator}) \end{cases}$$



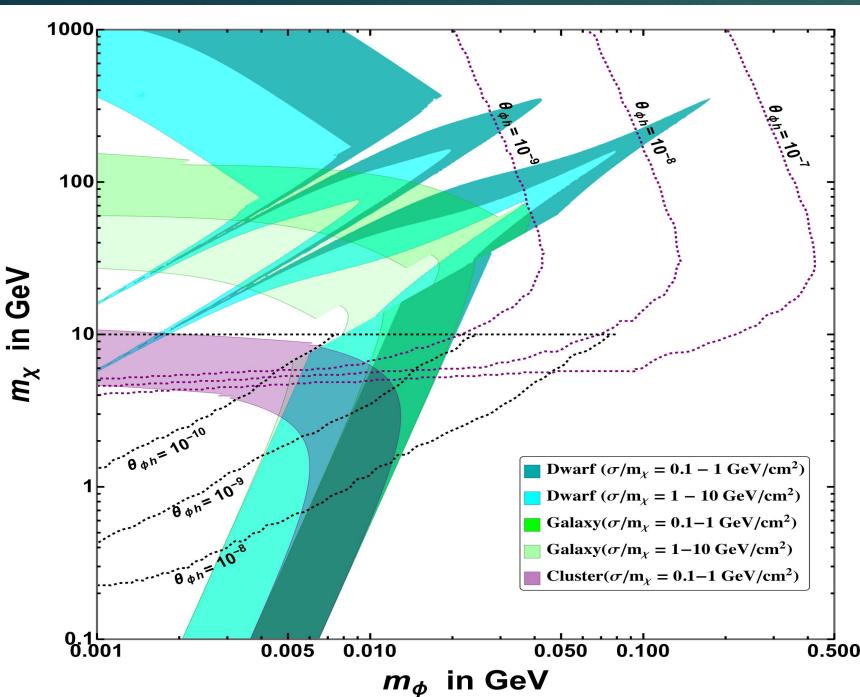
$$V(r) = \pm \frac{\alpha_D}{r} e^{-M_\phi r}$$





Issues with Symmetric SIDM

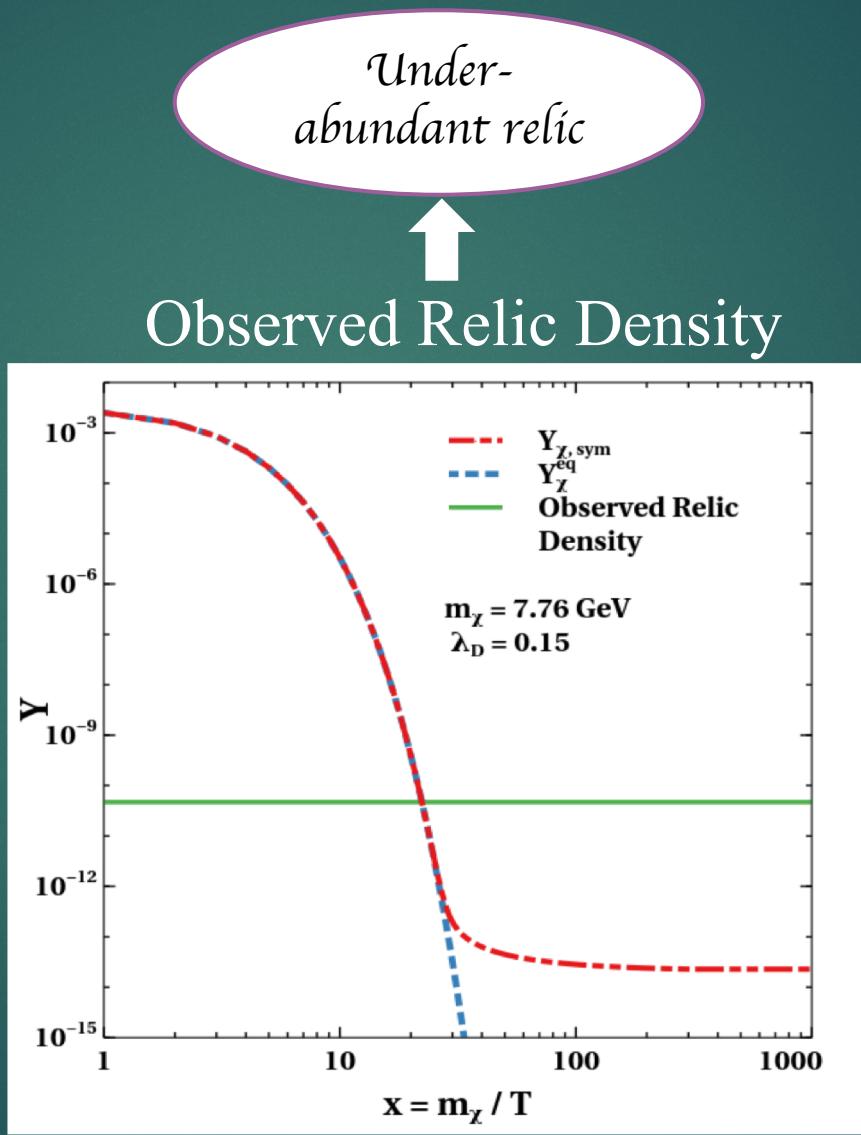
6



Direct Search

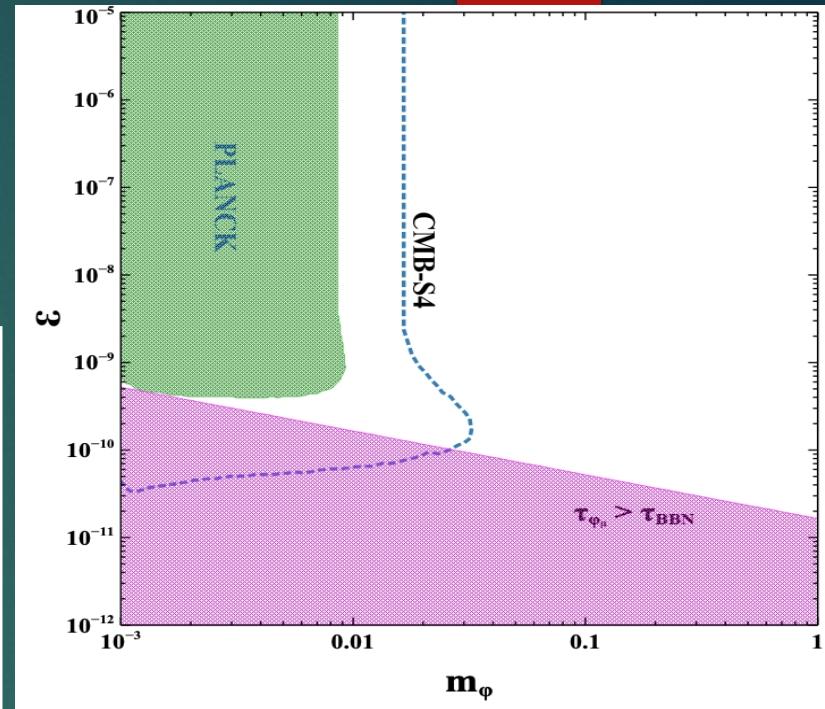


Tight upper bound on portal mixing



Under-abundant relic

Observed Relic Density



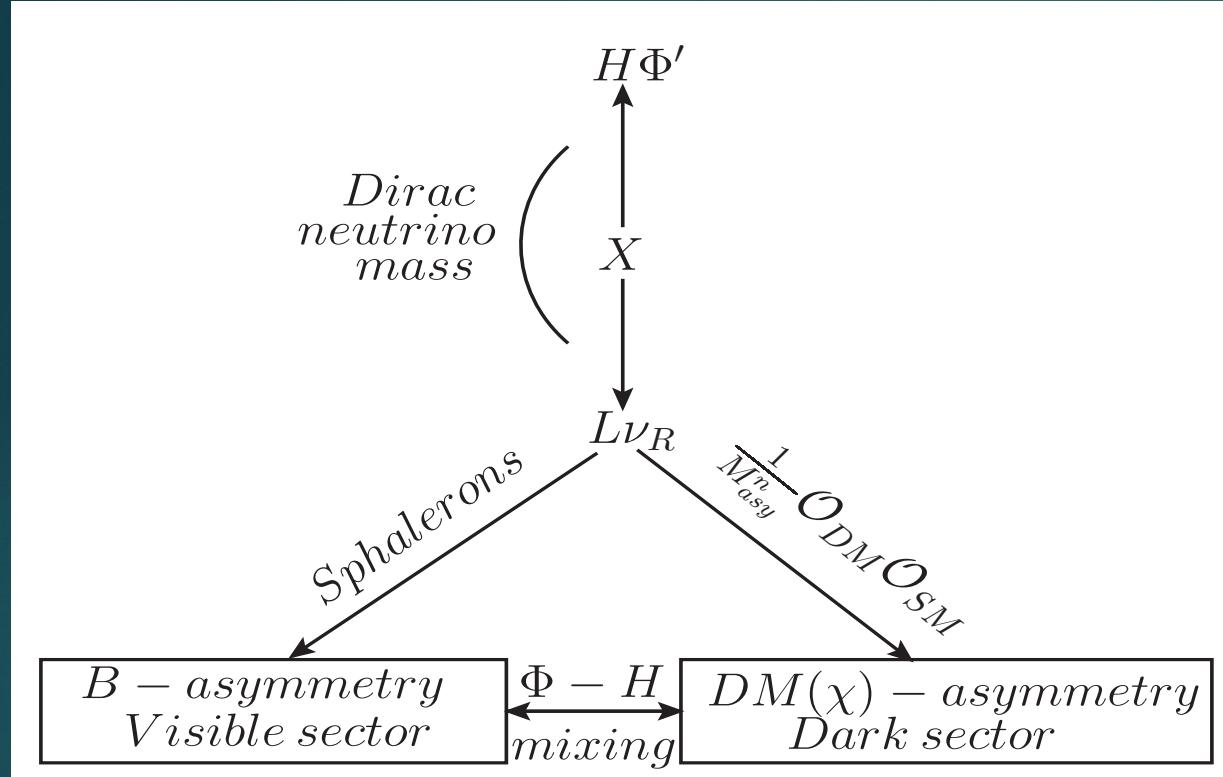
Indirect Search



Tight lower bound on portal mixing

ASIDM via Dirac Leptogenesis

7



$$\begin{aligned} -\mathcal{L} \supset & m_\chi \bar{\chi} \chi + \lambda_D \bar{\chi} \chi \Phi + y \bar{L} \tilde{X} \nu_R + \rho \Phi'^* X^\dagger H \\ & + \lambda \bar{\chi} L \eta + \text{h.c.} + V(X, \eta, H, \Phi, \Phi'), \end{aligned}$$

(M. Dutta, N. Narendra, N. Sahu, S. Shil,
Phys. Rev. D 106 (2022) 9, 095017)

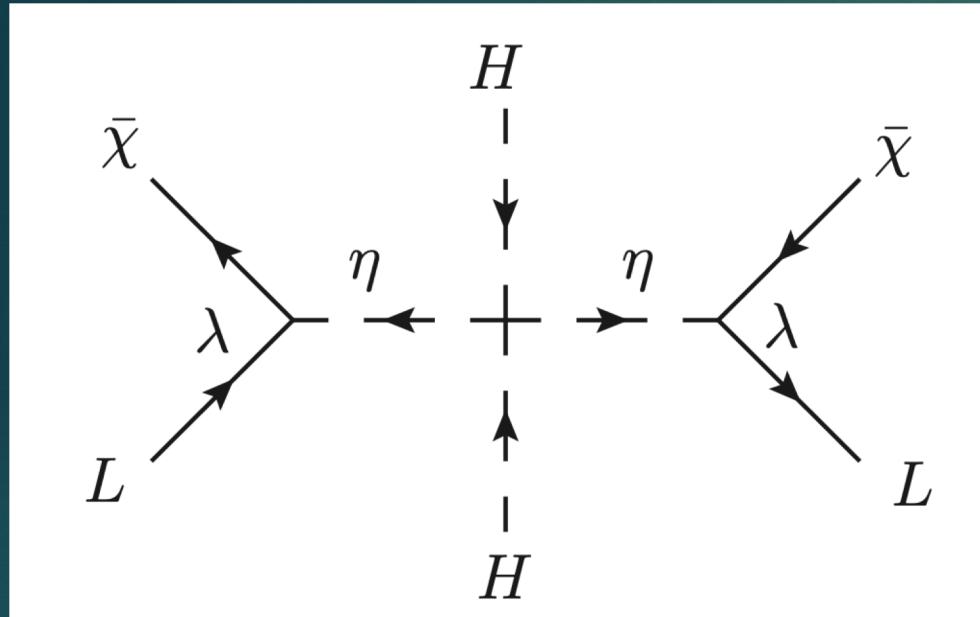
| Fields | SU(2) _L | U(1) _Y | U(1) _D | U(1) _{B-L} |
|---------|--------------------|-------------------|-------------------|---------------------|
| X_i | 2 | +1 | -1 | 0 |
| η | 2 | +1 | 1/2 | 0 |
| ν_R | 1 | 0 | -1 | -1 |
| Φ' | 1 | 0 | +1 | 0 |
| Φ | 1 | 0 | 0 | 0 |
| χ | 1 | 0 | 1/2 | -1 |

$$\begin{aligned} V(\eta, H, \Phi, \Phi') = & M_\eta^2 (\eta^\dagger \eta) + \lambda_\eta (\eta^\dagger \eta)^2 + \lambda'_{\eta H} (\eta^\dagger \eta)(H^\dagger H) \\ & + [\lambda_{\eta H} (\eta^\dagger H)^2 + \text{h.c.}] \\ & - \mu_H^2 H^\dagger H + \lambda_H (H^\dagger H)^2 + \frac{1}{2} m_\phi^2 \Phi^2 + \frac{1}{3} \mu_\Phi \Phi^3 \\ & + \frac{1}{4} \lambda_\Phi \Phi^4 - \mu_{\Phi'}^2 (\Phi'^\dagger \Phi') + \lambda_{\Phi'} (\Phi'^\dagger \Phi')^2 \\ & + \frac{\mu_1}{\sqrt{2}} \Phi H^\dagger H + \frac{\mu_2}{\sqrt{2}} \Phi (\Phi'^\dagger \Phi') + \frac{\lambda_{H\Phi}}{2} H^\dagger H \Phi^2 \\ & + \lambda_{H\Phi'} H^\dagger H (\Phi'^\dagger \Phi') + \frac{\lambda_{\Phi\Phi'}}{2} \Phi^2 (\Phi'^\dagger \Phi') \end{aligned}$$



ASIDM via Dirac Leptogenesis

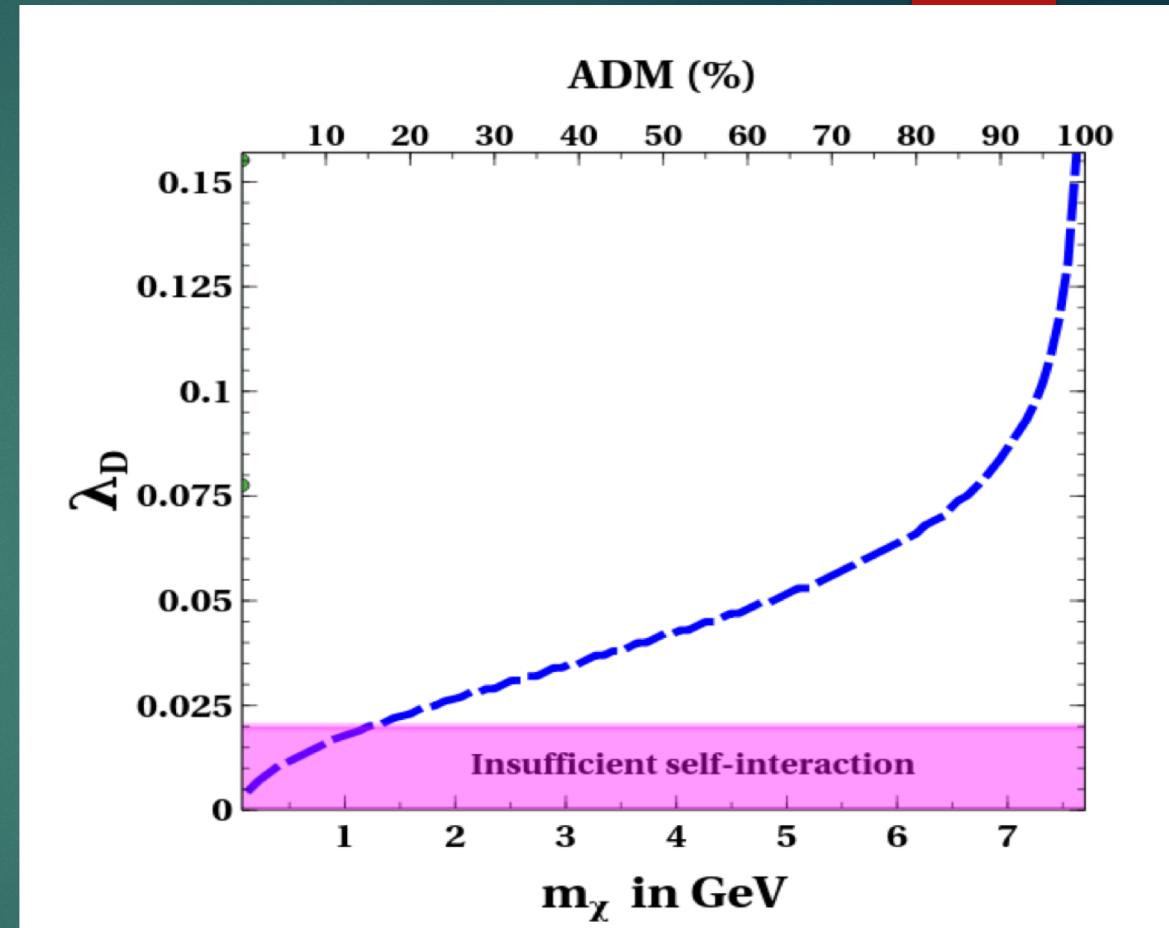
8



$$\mathcal{O}_8 = \frac{\lambda^2 \lambda_{\eta_H} \bar{\chi} L H H \bar{\chi} L}{M_\eta^4} \equiv \frac{\bar{\chi}^2 (L H)^2}{M_{asy}^4}$$

$$n_\chi = (n_{B-L})_{\text{dark}} = \frac{58}{291} (n_{B-L})_{\text{vis}}$$

(M. Dutta, N. Narendra, N. Sahu, S. Shil, Phys. Rev. D 106 (2022) 9, 095017)

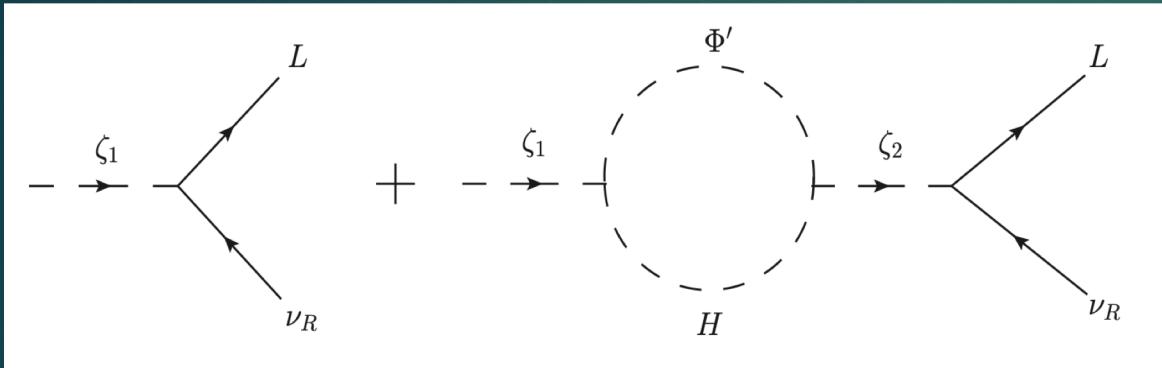


$$m_\chi = \frac{90}{58} \times 5 \times \frac{\text{ADM}(\%)}{100}.$$



ASIDM via Dirac Leptogenesis

9

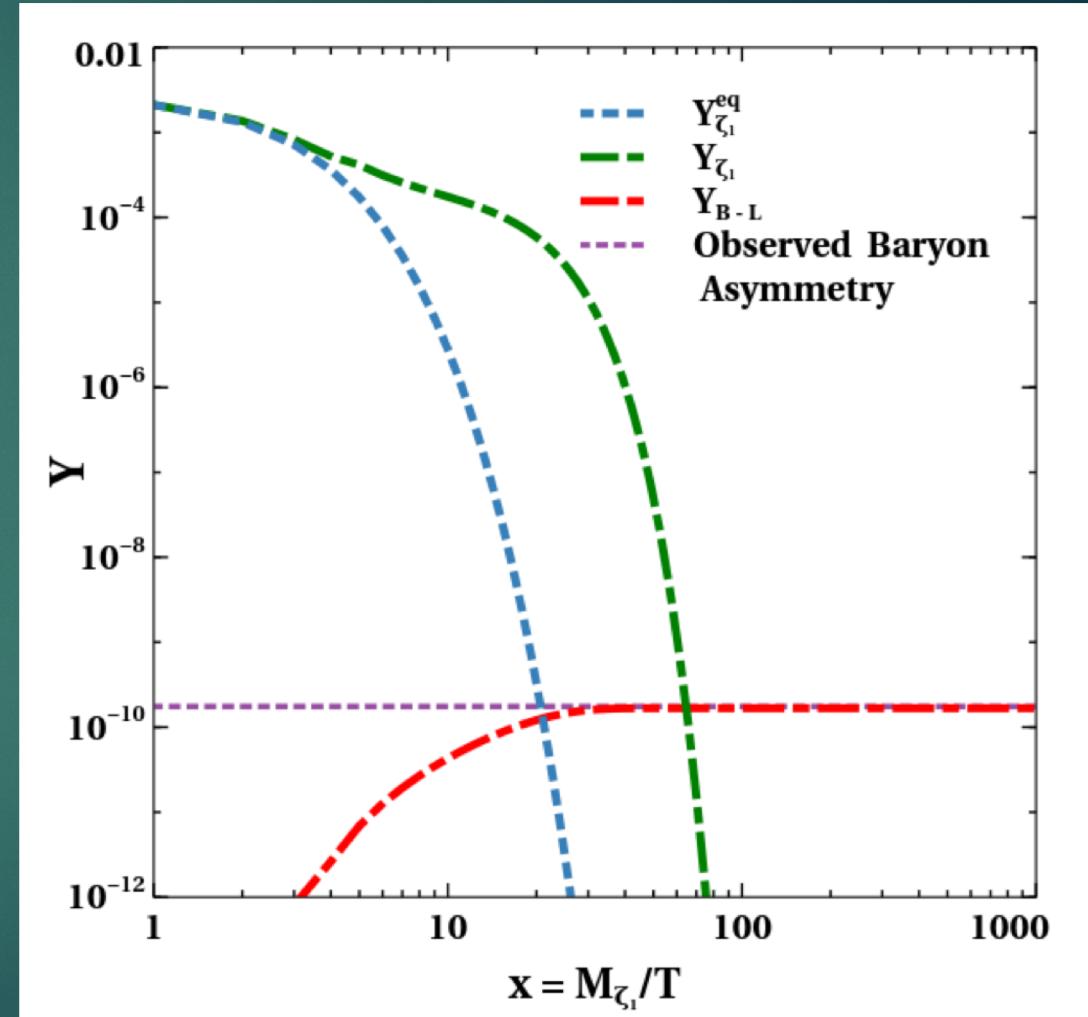


$$\frac{dY_{\zeta_1}}{dx} = -\frac{x}{H(M_{\zeta_1})} s \langle \sigma |v|_{(\zeta_1 \zeta_1 \rightarrow All)} \rangle [Y_{\zeta_1}^2 - (Y_{\zeta_1}^{eq})^2]$$

$$-\frac{x}{H(M_{\zeta_1})} \Gamma_{(\zeta_1 \rightarrow All)} [Y_{\zeta_1} - Y_{\zeta_1}^{eq}]$$

$$\frac{dY_{B-L}}{dx} = \frac{x}{H(M_{\zeta_1})} \left[\epsilon_L \Gamma_{(\zeta_1 \rightarrow All)} B_L \left(Y_{\zeta_1} - Y_{\zeta_1}^{eq} \right) - \Gamma_W Y_{B-L} \right]$$

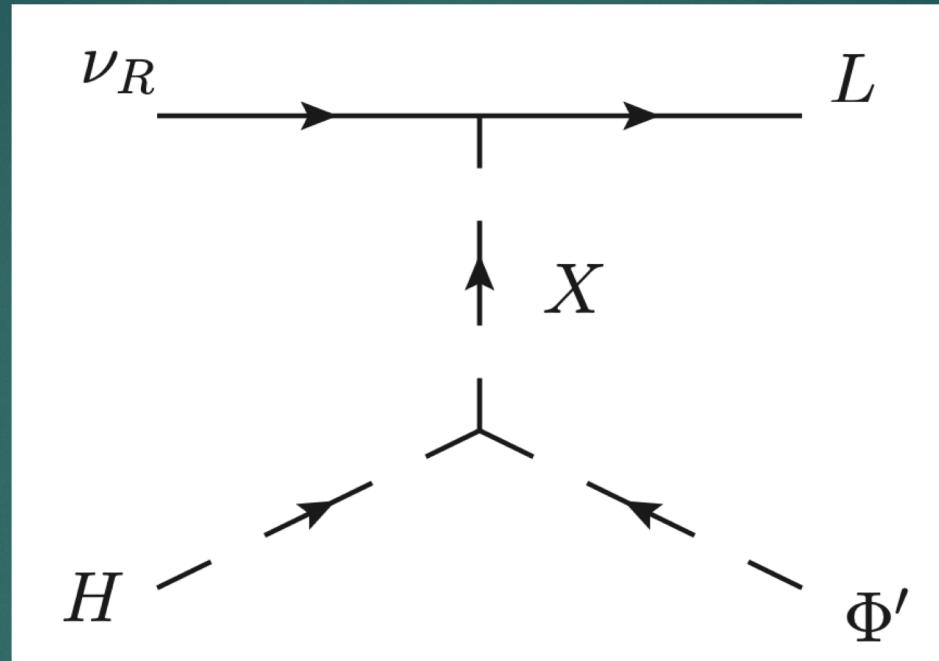
$$\begin{aligned} \epsilon_L &= [B_L(\zeta_1^- \rightarrow l^- \nu_R) - B_L(\zeta_1^+ \rightarrow (l^-)^c \nu_R^c)] \\ &= -\frac{\text{Im} \left(\rho_1^* \rho_2 \sum_{k,l} y_{1kl}^* y_{2kl} \right)}{8\pi^2 (M_{\zeta_2}^2 - M_{\zeta_1}^2)} \left[\frac{M_{\zeta_1}}{\Gamma_{\zeta_1}} \right], \end{aligned}$$





ASIDM via Dirac Leptogenesis: Neutrino Mass

10



$$\mathcal{L}_{Dirac} = -y_1 \frac{\rho_1}{M_{X_1}^2} \bar{L} H \Phi' \nu_R - y_2 \frac{\rho_2}{M_{X_2}^2} \bar{L} H \Phi' \nu_R$$

$$M_\nu \simeq y_1 \frac{\rho_1 v w}{M_{X_1}^2} + y_2 \frac{\rho_2 v w}{M_{X_2}^2}.$$



Summary

- ▶ The inadequacies of symmetric SIDM (relic density, indirect detection) can be alleviated in Asymmetric SIDM Scenario.
- ▶ It can naturally explain the cosmic coincidence “ $\Omega_{\text{DM}} \simeq 5\Omega_{\text{B}}$ ”, while simultaneously alleviating the small scale anomalies of Λ CDM.
- ▶ Provides an avenue to explain SIDM, neutrino mass and baryon asymmetry in a common frame-work.
- ▶ $0\nu\beta\beta$ experiment will play the crucial role in deciding nature of neutrinos, and hence the type of leptogenesis.
- ▶ Next decade will be a testing time for Self-interacting Dark Matter.



FPCP2024

Chulalongkorn University, Bangkok

28th May 2024

Thank You

Dr. Manoranjan Dutta
Department Of Physics
North Lakhimpur University

