

# Quantum decoherence and CP violation at Protvino to ORCA

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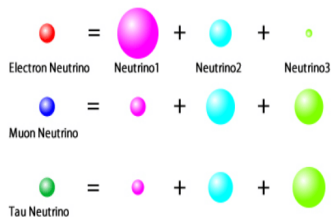
- 1 Introduction and motivation
- 2 Decoherence
- 3 Current bounds on decoherence parameters
- 4 Simulation details
- 5 Analysis and results
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# Neutrino oscillation

- Neutrino flavor ( $\nu_\alpha$ ) oscillations are generated by the quantum interference of neutrino mass states ( $\nu_j$ ),

$$|\nu_\alpha\rangle = \sum_j U_{\alpha j} |\nu_j\rangle .$$

- Coherence between  $\nu_j$  is essential for neutrino oscillations.
- Neutrino system behaves as a closed system.



- Time evolution of  $\rho$  represented by Liouville-Von Neumann equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] .$$

- $P_{\alpha\beta}(t) = Tr[\rho_\alpha(t)\rho_\beta(0)] .$

# Motivation

- $P_{\alpha\beta}(t \approx L) = P_{\alpha\beta}(\Delta m_{21}^2, |\Delta m_{31}^2|, \theta_{12}, \theta_{13}, \theta_{23}, \delta_{\text{CP}}; E, L, V(x))$  .
- Significant opportunity to probe new physics (NP) phenomenon with upcoming high-precision neutrino oscillation experiments.
- One such interesting phenomenon is the **environmentally induced decoherence**.
- This effect arises in the oscillation probabilities through the damping term  $e^{-\Gamma L}$ .

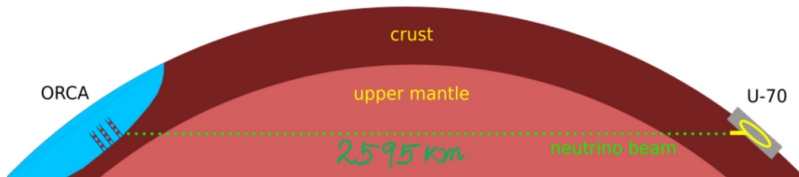


Fig 1. Protvino to ORCA.

- Neutrino system interacts with the stochastic environment.
- $$\frac{d\tilde{\rho}_m(t)}{dt} = -i [H, \tilde{\rho}_m(t)] + \mathcal{D} [\tilde{\rho}_m(t)] .$$
- **Assumptions:**
  - (a) complete positivity,
  - (b) trace preserving conditions,
  - (c) increasing von Neumann entropy,
  - (d) average energy conservation of the system.

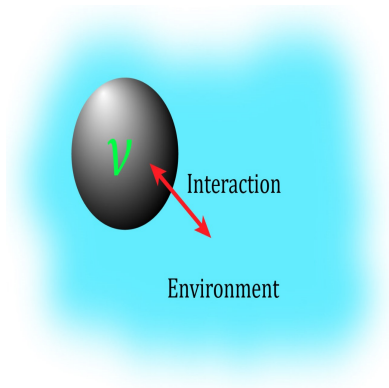


Fig 2. Neutrino system as an open quantum system.

# Oscillation probability in presence of decoherence

- $P_{\alpha\beta}(t) = \text{Tr}[\tilde{\rho}_\alpha(t)\tilde{\rho}_\beta(0)]$  .
- $P_{\alpha\beta}(L) = \delta_{\alpha\beta} - 2 \sum_{j>k} \text{Re} \left( \tilde{U}_{\beta j} \tilde{U}_{\alpha j}^* \tilde{U}_{\alpha k} \tilde{V}_{\beta k}^* \right)$   
 $+ 2 \sum_{j>k} \text{Re} \left( \tilde{U}_{\beta j} \tilde{U}_{\alpha j}^* \tilde{U}_{\alpha k} \tilde{U}_{\beta k}^* \right) \exp(-\Gamma_{jk}L) \cos \left( \frac{\tilde{\Delta}m_{jk}^2 L}{2E} \right)$   
 $+ 2 \sum_{j>k} \text{Im} \left( \tilde{U}_{\beta j} \tilde{U}_{\alpha j}^* \tilde{U}_{\alpha k} \tilde{U}_{\beta k}^* \right) \exp(-\Gamma_{jk}L) \sin \left( \frac{\tilde{\Delta}m_{jk}^2 L}{2E} \right)$  .
- Damping of interference terms by a factor  $e^{-\Gamma L}$  in the oscillation probability.
- Energy dependency on  $\Gamma$  :

$$\Gamma_{jk}(E_\nu) = \Gamma_0 \left( \frac{E_\nu}{\text{GeV}} \right)^n ; n = 0, \pm 1, \pm 2 .$$

# Current upper bounds on decoherence parameters

## Search for decoherence from quantum gravity with atmospheric neutrinos

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The IceCube Collaboration\* 

Neutrino oscillations at the highest energies and longest baselines can be used to study the structure of spacetime and test the fundamental principles of quantum mechanics. If the metric of spacetime has a quantum mechanical description, its fluctuations at the Planck scale are expected to introduce non-unitary effects that are inconsistent with the standard unitary time evolution of quantum mechanics. Neutrinos interacting with such fluctuations would lose their quantum coherence, deviating from the expected oscillatory flavour composition at long distances and high energies. Here we use atmospheric neutrinos detected by the IceCube South Pole Neutrino Observatory in the energy range of 0.5–10.0 TeV to search for coherence loss in neutrino propagation. We find no evidence of anomalous neutrino decoherence and determine limits on neutrino–quantum gravity interactions. The constraint on the effective decoherence strength parameter within an energy-independent decoherence model improves on previous limits by a factor of 30. For decoherence effects scaling as  $E^2$ , our limits are advanced by more than six orders of magnitude beyond past measurements compared with the state of the art.



## Search for decoherence from quantum gravity with atmospheric neutrinos

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**Table 1 | Summary of constraints on decoherence models obtained in this analysis**

$n$	Phase perturbation $\Gamma_{90}$	State selection $\Gamma_{90}$
0	$1.18 \times 10^{-15}$ eV	$1.17 \times 10^{-15}$ eV
1	$6.89 \times 10^{-16}$ eV	$6.67 \times 10^{-16}$ eV
2	$9.80 \times 10^{-18}$ eV	$9.48 \times 10^{-18}$ eV
3	$1.58 \times 10^{-19}$ eV	$1.77 \times 10^{-19}$ eV

The 90% CL upper limits on decoherence strength parameter  $\Gamma_0$  (which we name  $\Gamma_{90}$ ) are reported for each power-law index  $n$  with power-law pivot energy  $E_0=1$  TeV in the state selection and phase perturbation models.

## Neutrino oscillation bounds on quantum decoherence

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ABSTRACT: We consider quantum-decoherence effects in neutrino oscillation data. Working in the open quantum system framework we adopt a phenomenological approach that allows to parameterize the energy dependence of the decoherence effects. We consider several phenomenological models. We analyze data from the reactor experiments **RENO**, **Daya Bay** and **KamLAND** and from the accelerator experiments **NOvA**, **MINOS/MINOS+** and **T2K**. We obtain updated constraints on the decoherence parameters quantifying the strength of damping effects, which can be as low as  $\Gamma_{ij} \lesssim 8 \times 10^{-27}$  GeV at 90% confidence level in some cases. We also present sensitivities for the future facilities **DUNE** and **JUNO**.

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Decoherence Model	$n = -2$	$n = -1$	$n = 0$	$n = +1$	$n = +2$
A: $\Gamma_{21} = \Gamma_{31} = \Gamma_{32}$	$7.8 \times 10^{-27}$ (KL)	$1.8 \times 10^{-24}$ (KL)	$5.1 \times 10^{-24}$ (M)	$3.0 \times 10^{-25}$ (M)	$1.3 \times 10^{-26}$ (M)
B: $\Gamma_{21} = \Gamma_{31}, \Gamma_{32} = 0$	$7.9 \times 10^{-27}$ (KL)	$1.8 \times 10^{-24}$ (KL)	$2.4 \times 10^{-23}$ (N)	$2.3 \times 10^{-24}$ (M)	$1.0 \times 10^{-25}$ (M)
C: $\Gamma_{21} = \Gamma_{32}, \Gamma_{31} = 0$	$7.9 \times 10^{-27}$ (KL)	$1.8 \times 10^{-24}$ (KL)	$9.4 \times 10^{-24}$ (M)	$5.7 \times 10^{-25}$ (M)	$2.5 \times 10^{-26}$ (M)
D: $\Gamma_{31} = \Gamma_{32}, \Gamma_{21} = 0$	$6.9 \times 10^{-25}$ (R)	$2.1 \times 10^{-23}$ (T2K)	$5.6 \times 10^{-24}$ (M)	$3.3 \times 10^{-25}$ (M)	$1.5 \times 10^{-26}$ (M)
E: $\Gamma_{21}, \Gamma_{31} = \Gamma_{32} = 0$	$7.9 \times 10^{-27}$ (KL)	$1.8 \times 10^{-24}$ (KL)	$3.2 \times 10^{-23}$ (M)	$2.2 \times 10^{-24}$ (M)	$1.0 \times 10^{-25}$ (M)
F: $\Gamma_{31}, \Gamma_{21} = \Gamma_{32} = 0$	$1.0 \times 10^{-24}$ (R)	$1.9 \times 10^{-23}$ (T2K)	$2.3 \times 10^{-23}$ (N)	$2.2 \times 10^{-24}$ (M)	$1.0 \times 10^{-25}$ (M)
G: $\Gamma_{32}, \Gamma_{21} = \Gamma_{31} = 0$	$4.0 \times 10^{-23}$ (T2K)	$6.5 \times 10^{-23}$ (T2K)	$1.1 \times 10^{-23}$ (M)	$6.6 \times 10^{-25}$ (M)	$3.0 \times 10^{-26}$ (M)

TABLE III: Summary of results: each column shows the most constraining upper limit on  $\Gamma_{ij}$ , in GeV, for each model (A to G) and value of  $n$ . We also clarify, within parenthesis, which experiment sets the bound (KL = KamLAND, R = RENO, M = MINOS/MINOS+, N = NOvA).

## Quantum decoherence and relaxation in long-baseline neutrino data

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**ABSTRACT:** We investigate the effect of quantum decoherence and relaxation in neutrino oscillations using MINOS and T2K data. The formalism of open quantum systems is used to describe the interaction of a neutrino system with the environment, where the strength of the interaction is regulated by a decoherence parameter  $\Gamma$ . We assume an energy dependence parameterized by  $\Gamma = \gamma_0(E/\text{GeV})^n$ , with  $n = -2, 0, +2$ , and consider three different scenarios, allowing the investigation of the effect of relaxation and of constraining the solar and atmospheric sectors to the same decoherence parameter. The MINOS and T2K data present a complementary behavior, with regard to our theoretical model, resulting in a better sensitivity for  $n = +2$  and  $n = -2$ , respectively. We perform a combined analyses of both experimental data, which also include a reactor constraint on  $\sin^2 \theta_{13}$ , and observe an independence of the results to the scenarios we investigate. Our analyses obtain limits on  $\gamma_0$  based on long-baseline data for scenarios allowing or not relaxation. We improve some previous bounds on  $\gamma_0$  and outline which data (solar, reactor, atmospheric, long-baseline) determine the more stringent constraints for different scenarios and energy dependencies.

## Quantum decoherence and relaxation in long-baseline neutrino data

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	$n = -2$	$n = 0$	$n = 2$
<b>MINOS (this work)</b>			
Case 1 ( $\Gamma_{31} = \Gamma_{32} = \Gamma_{21}$ , with relaxation)	$(0.33 - 37.0) \times 10^{-23}$	$6.8 \times 10^{-23}$	$1.7 \times 10^{-25}$
Case 2 ( $\Gamma_{31} = \Gamma_{32} = \Gamma_{21}$ , no relaxation)	$30.0 \times 10^{-23}$	$6.5 \times 10^{-23}$	$2.4 \times 10^{-25}$
Case 3 ( $\Gamma_{31} = \Gamma_{32}$ , $\Gamma_{21} = 0$ , no relaxation)	$19.0 \times 10^{-23}$	$5.9 \times 10^{-23}$	$2.5 \times 10^{-25}$
<b>T2K (this work)</b>			
Case 1	$2.8 \times 10^{-23}$	$6.2 \times 10^{-23}$	$3.1 \times 10^{-23}$
Case 2	$2.9 \times 10^{-23}$	$5.2 \times 10^{-23}$	$3.3 \times 10^{-23}$
Case 3	$1.7 \times 10^{-23}$	$3.9 \times 10^{-23}$	$4.1 \times 10^{-23}$
<b>MINOS+T2K (this work)</b>			
Case 1	$2.9 \times 10^{-23}$	$6.6 \times 10^{-23}$	$2.3 \times 10^{-25}$
Case 2	$3.4 \times 10^{-23}$	$6.1 \times 10^{-23}$	$2.9 \times 10^{-25}$
Case 3	$2.0 \times 10^{-23}$	$5.0 \times 10^{-23}$	$3.3 \times 10^{-25}$

Our work:  
[arXiv:2405.03286](https://arxiv.org/abs/2405.03286)

- Case 1:  $\Gamma_{21} = \Gamma_{31} = \Gamma_{32} \neq 0$
- Case 2:  $\Gamma_{21} = \Gamma_{31}$  ,  $\Gamma_{32} = 0$
- Case 3:  $\Gamma_{21} = \Gamma_{32}$  ,  $\Gamma_{31} = 0$
- Case 4:  $\Gamma_{31} = \Gamma_{32}$  ,  $\Gamma_{21} = 0$

$$\Gamma_{jk}(E_\nu) = \Gamma_0 \left( \frac{E_\nu}{\text{GeV}} \right)^n ; n = 0, \pm 1, \pm 2 .$$

**Objectives:** To study

- the effects of  $\Gamma$  on different oscillation channels.
- the bounds on decoherence parameters in P2O considering  $\Gamma \propto E_\nu^n$ .
- MH and CPV sensitivity considering  $3\sigma$  value of  $\Gamma$ .

# Simulation details

Baseline	2595 km
Exposure	$0.8 \times 10^{20}$ POT/year
Beam power	90 kW (modest) and 450 kW (upgradable)
Run time	6 years ( $3(\nu) + 3(\bar{\nu})$ )
Detector mass	8 Mton (Water Cherenkov)
Matter density	$3.25 \text{ gm/cm}^3$

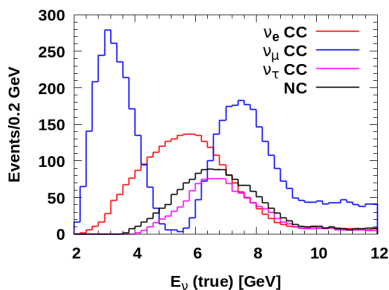


Fig 3. Events rate considering standard interaction in the presence of earth matter.  
[1] *Eur. Phys. J. C* (2019) 79: 758



# Oscillation probabilities

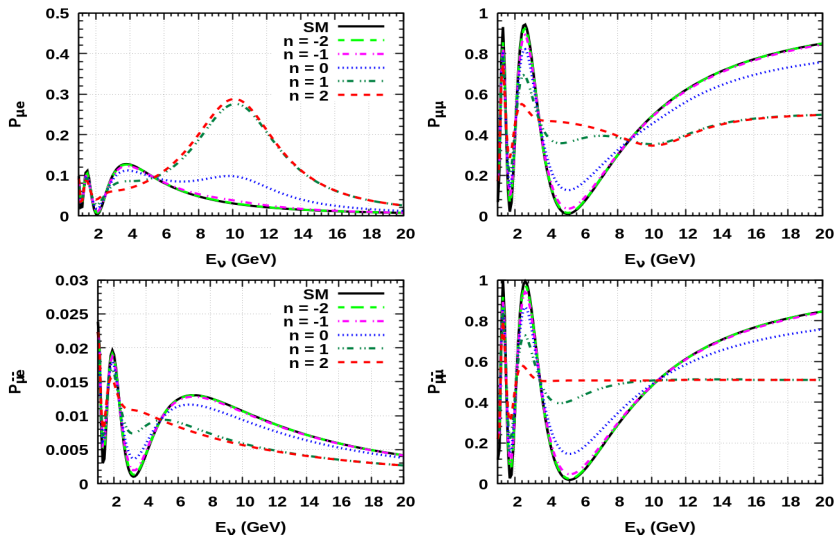


Fig 4. The left (right) plot in the upper panel represents  $\nu_e$  appearance ( $\nu_\mu$  disappearance) probability and in the lower panel corresponding anti-neutrino probability. We consider a representative case  $\Gamma_{21} = \Gamma_{31} = \Gamma_{32} = 2.3 \times 10^{-23}$  GeV.

- We incorporated a new oscillation probability engine into GLoBES [2,3] by taking into account the effect of decoherence on neutrino propagation.
- $\Delta\chi_{\Gamma}^2 = \chi^2(\Gamma(\text{true}) = 0, \Gamma(\text{test}) \neq 0)$  .  
Marginalized over  $\theta_{23}$  and  $\delta_{CP}$ ,  $(\Delta m_{31}^2)_{NH}$ .
- $\Delta\chi_{MH}^2 = \chi_{\text{true}}^2(\Gamma \neq 0, \Delta m_{31}^2 > 0) - \chi_{\text{test}}^2(\Gamma = 0, \Delta m_{31}^2 < 0)$  .  
Marginalized over  $\theta_{23}$  and  $\delta_{CP}$ .
- $\Delta\chi_0^2 = \chi_{\text{true}}^2(\delta_{CP}(\text{true}), \Gamma \neq 0) - \chi_{\text{test}}^2(\delta_{CP} = 0, \Gamma = 0)$  ,  
 $\Delta\chi_{\pi}^2 = \chi_{\text{true}}^2(\delta_{CP}(\text{true}), \Gamma \neq 0) - \chi_{\text{test}}^2(\delta_{CP} = \pi, \Gamma = 0)$  ,  
 $\Delta\chi_{CPV}^2 = \min[\Delta\chi_0^2, \Delta\chi_{\pi}^2]$  .  
Marginalized over  $\theta_{23}$ ,  $(\Delta m_{31}^2)_{NH}$ .

[2] Comput. Phys. Commun. 167 (2005) 195 , [3] Comput. Phys. Commun. 177 (2007) 432–438 .

# Results

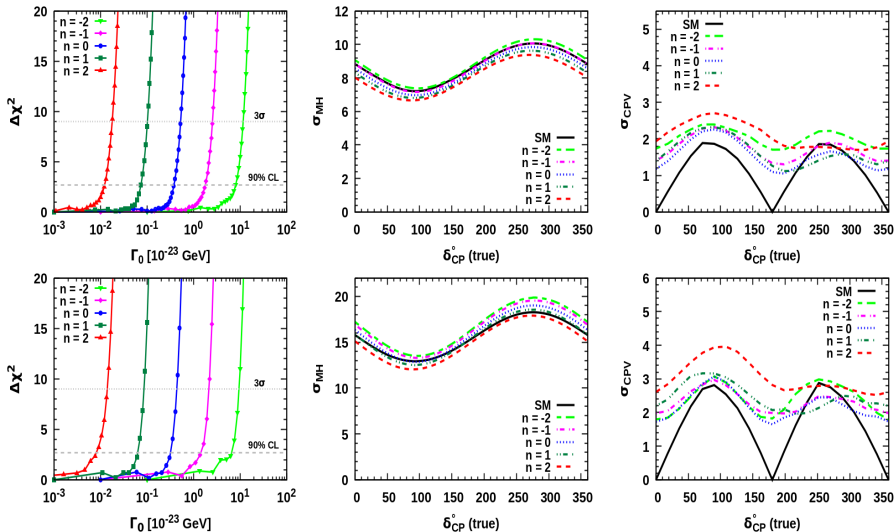


Fig 5. Case 1:  $\Gamma_{21} = \Gamma_{31} = \Gamma_{32} \neq 0$ . Upper and lower panel are the results of 90 kW and 450 kW beam respectively.

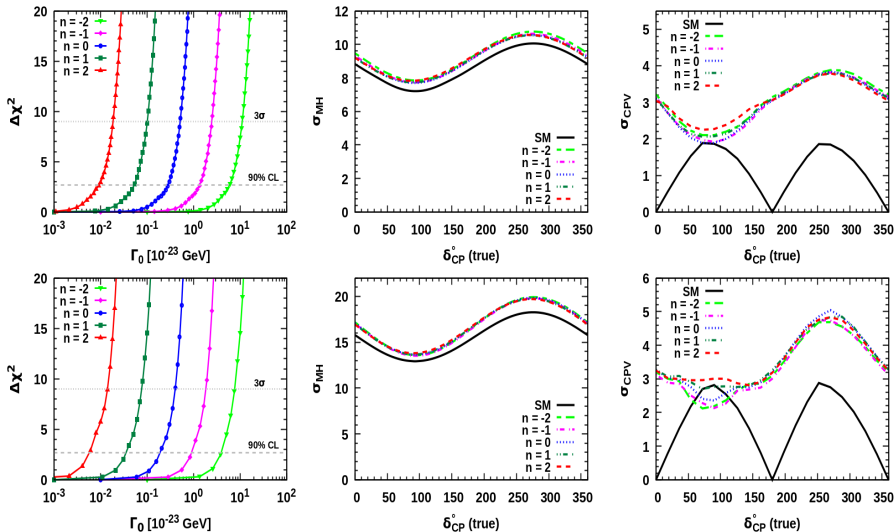


Fig 6. Case 2:  $\Gamma_{21} = \Gamma_{31}$ ,  $\Gamma_{32} = 0$ . Upper and lower panel are the results of 90 kW and 450 kW beam respectively.

# Summary table

$n$	CL	90 kW		450 kW	
		Case 1	Case 2	Case1	Case 2
$n = -2$	90%	$8.06 \times 10^{-23}$	$6.0 \times 10^{-23}$	$6.4 \times 10^{-23}$	$3.9 \times 10^{-23}$
	$3\sigma$	$1.15 \times 10^{-22}$	$1.11 \times 10^{-22}$	$9.0 \times 10^{-23}$	$7.6 \times 10^{-23}$
$n = -1$	90%	$1.8 \times 10^{-23}$	$1.36 \times 10^{-23}$	$1.4 \times 10^{-23}$	$8.5 \times 10^{-24}$
	$3\sigma$	$2.55 \times 10^{-23}$	$2.46 \times 10^{-23}$	$2.1 \times 10^{-23}$	$1.75 \times 10^{-23}$
$n = 0$	90%	$3.76 \times 10^{-24}$	$2.76 \times 10^{-24}$	$3.1 \times 10^{-24}$	$1.89 \times 10^{-24}$
	$3\sigma$	$5.3 \times 10^{-24}$	$5.18 \times 10^{-24}$	$4.2 \times 10^{-24}$	$4.0 \times 10^{-24}$
$n = 1$	90%	$7.4 \times 10^{-25}$	$5.4 \times 10^{-25}$	$6.2 \times 10^{-25}$	$3.6 \times 10^{-25}$
	$3\sigma$	$1.04 \times 10^{-24}$	$1.02 \times 10^{-24}$	$8.5 \times 10^{-25}$	$7.7 \times 10^{-25}$
$n = 2$	90%	$1.24 \times 10^{-25}$	$9.07 \times 10^{-26}$	$7.8 \times 10^{-26}$	$5.87 \times 10^{-26}$
	$3\sigma$	$1.79 \times 10^{-25}$	$1.87 \times 10^{-25}$	$1.5 \times 10^{-25}$	$1.4 \times 10^{-25}$

Table 1: Upper bounds on  $\Gamma_0$  (GeV) obtained for different power law dependencies.

# Conclusions

- 1 The upper bounds on  $\Gamma$  are comparatively stronger for  $n \geq 0$ .
- 2 Similar conclusion can be drawn for upgraded 450 kW beam.
- 3 Case-2 ( $\Gamma_{21} = \Gamma_{31}$ ,  $\Gamma_{32} = 0$ ) poses strong constraints for  $n = 0$  on  $\Gamma \leq 2.76 \times 10^{-24}$  GeV for P2O and  $\Gamma \leq 1.89 \times 10^{-24}$  GeV for P2O-upgrade.
- 4 MH sensitivity does not vary significantly w.r.t SM.
- 5 In the presence of decoherence, CP conserving values ( $\delta_{CP} = 0, \pm\pi$ ) also show nonzero  $\delta_{CP}$

# Thank you

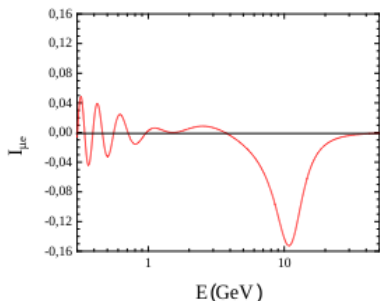
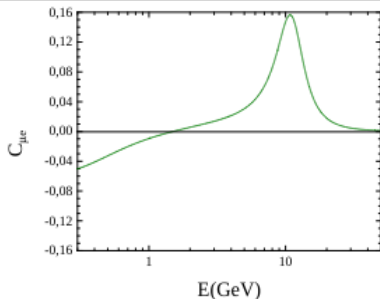
chinmay20pphy014@mahindrauniversity.edu.in

$$\begin{aligned} P_{\nu_{\alpha}\nu_{\alpha'}} &= \delta_{\alpha\alpha'} - 2 \sum_{j>k} \operatorname{Re}(\tilde{U}_{\alpha'j}\tilde{U}_{\alpha j}^*\tilde{U}_{\alpha k}\tilde{U}_{\alpha'k}^*) \\ &\quad + 2 \sum_{j>k} \operatorname{Re}(\tilde{U}_{\alpha'j}\tilde{U}_{\alpha j}^*\tilde{U}_{\alpha k}\tilde{U}_{\alpha'k}^*)e^{-\Gamma_{jk}x} \cos\left(\frac{\tilde{\Delta}_{jk}}{2E}x\right) \\ &\quad + 2 \sum_{j>k} \operatorname{Im}(\tilde{U}_{\alpha'j}\tilde{U}_{\alpha j}^*\tilde{U}_{\alpha k}\tilde{U}_{\alpha'k}^*)e^{-\Gamma_{jk}x} \sin\left(\frac{\tilde{\Delta}_{jk}}{2E}x\right) \end{aligned}$$



$$C_{\mu e} = -2 \sum_{i < j} \text{Re} \left[ \tilde{U}_{\mu i}^* \tilde{U}_{ei} \tilde{U}_{\mu j} \tilde{U}_{ej}^* \right]$$

$$I_{\mu e} = 2 \sum_{i < j} \text{Re} \left[ \tilde{U}_{\mu i}^* \tilde{U}_{ei} \tilde{U}_{\mu j} \tilde{U}_{ej}^* \right] \cos \tilde{\Delta}_{ij} \\ - 2 \sum_{i < j} \text{Im} \left[ \tilde{U}_{\mu i}^* \tilde{U}_{ei} \tilde{U}_{\mu j} \tilde{U}_{ej}^* \right] \sin \tilde{\Delta}_{ij}$$



- NH,  $\delta_{CP} = \pi/2$ ; 90 kW beam power.

Signal events	$\nu_e$	$\nu_\mu$	$\bar{\nu}_e$	$\bar{\nu}_\mu$
SM ( $\Gamma = 0$ )	2115	4746	160	1473
$\Gamma = 1.79 \times 10^{-25}$ GeV ( $n = 2$ )	2390	5160	154	1623
$\Gamma = 1.15 \times 10^{-22}$ GeV ( $n = -2$ )	2139	5039	159	1567

Bkg in $\nu_e$ app channel	$\nu_\mu$	NC	$\nu_\tau$
SM ( $\Gamma = 0$ )	2362	1190	876
$\Gamma = 1.79 \times 10^{-25}$ GeV ( $n = 2$ )	2596	1190	852
$\Gamma = 1.15 \times 10^{-22}$ GeV ( $n = -2$ )	2370	1190	868