

Determination of α_s from azimuthal correlations among jets at CMS

Paris Gianneios

on behalf of the CMS Collaboration



Parallel session talk - Strong Interactions and Hadron Physics

July 18, 2024

- 1 Introduction - Motivation
- 2 $R_{\Delta\phi}$ observable definition and event selection
- 3 $R_{\Delta\phi}$ measurement with CMS
- 4 Fixed Order pQCD predictions
- 5 Determination of $\alpha_s(m_z)$ and $\alpha_s(Q)$ evolution

Talk based on [arXiv: 2404.16082](#) *

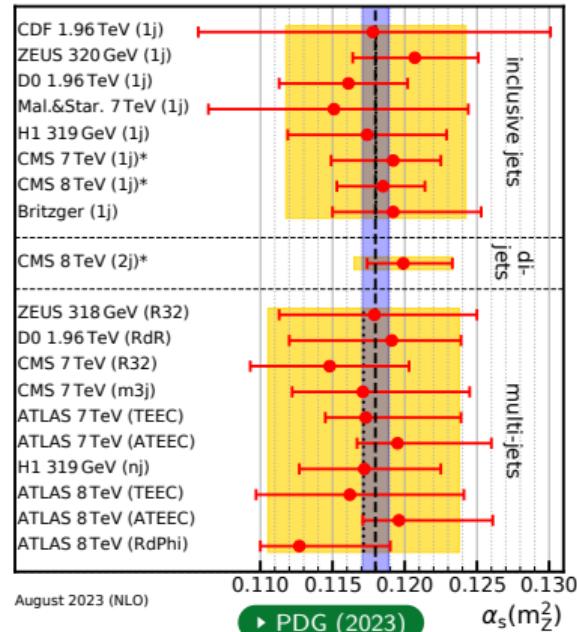
*Accepted by EPJC

Introduction - Motivation

Multijet cross section ratios

Goals

- ① Determination of the strong coupling constant $\alpha_s(m_z)$.
- ② Investigation of the energy scale (Q) dependence $\alpha_s(Q)$.



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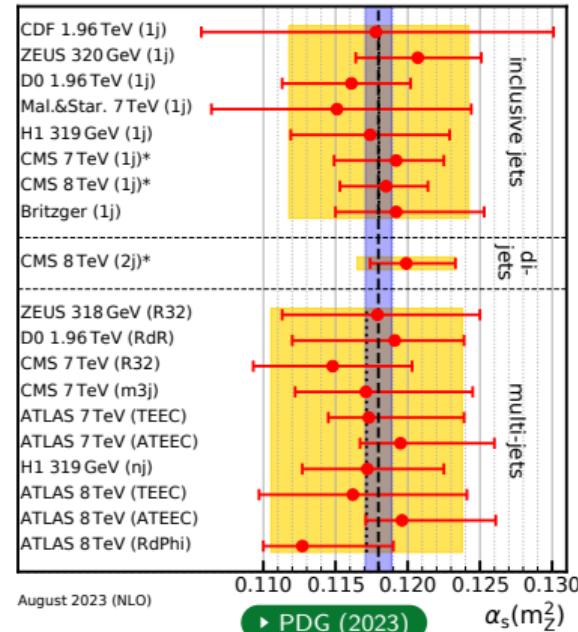
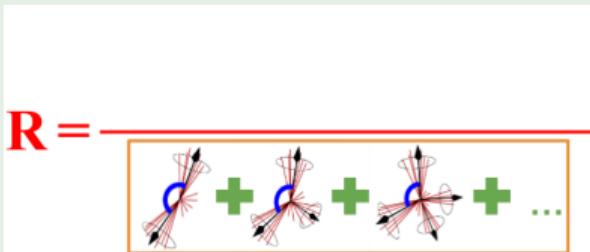
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- Denominator: topologies with at least 2-jets ($\sim \alpha_s^2$ @LO).



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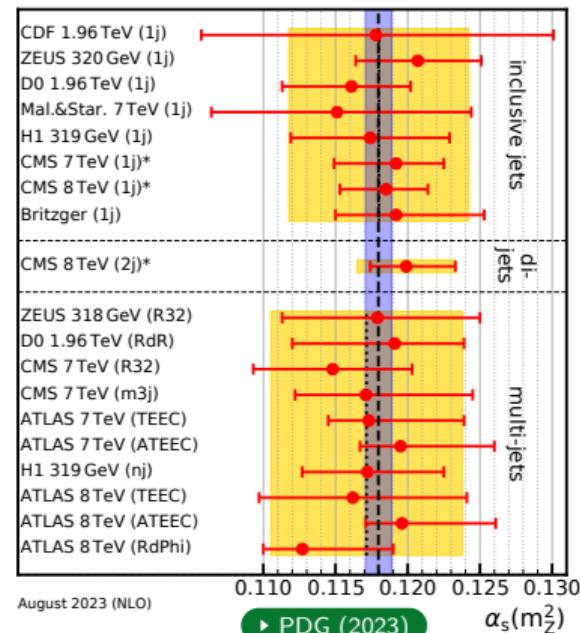
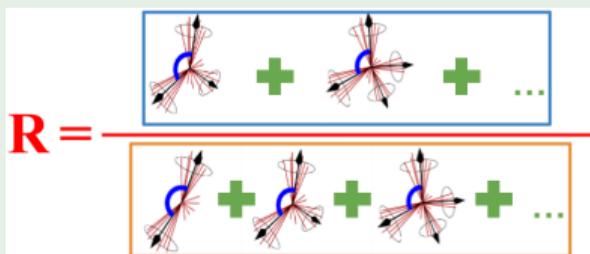
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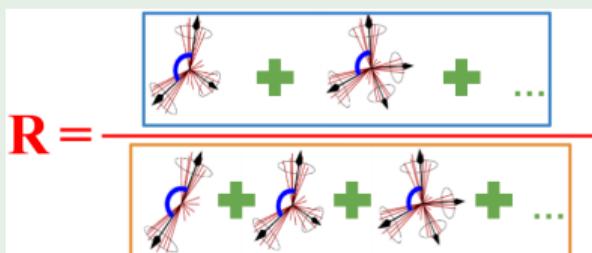
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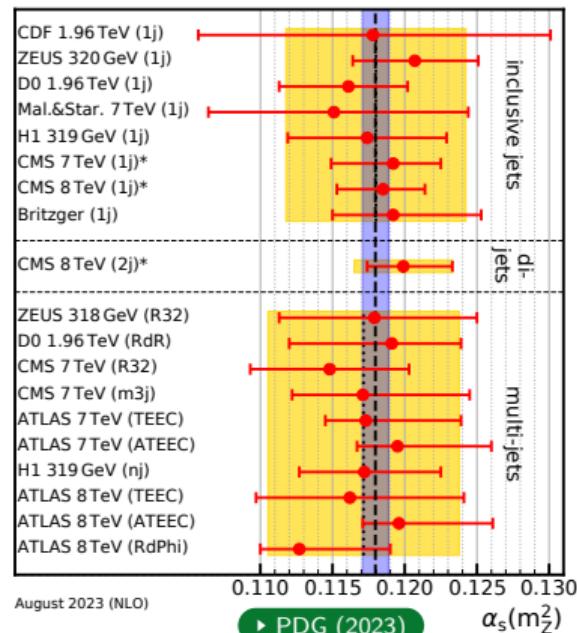
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Benefits

- ✓ Reduction/cancellation of systematic effects e.g. luminosity.
- ✓ Reduction of theoretical uncertainties e.g. non-perturbative.



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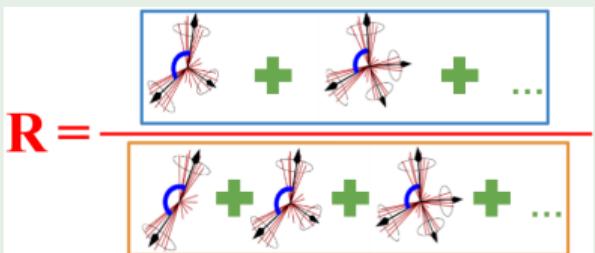
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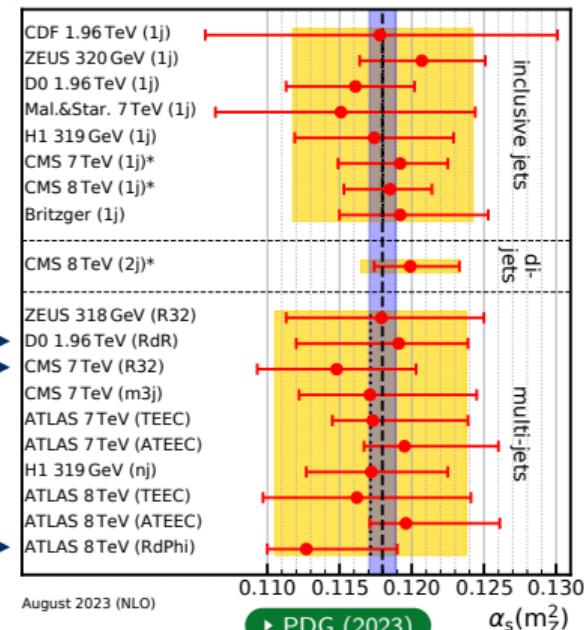


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Examples

- $R_{\Delta R}$ (D0, 1.96 TeV) ▶ PLB 718:1 (2012)
- R_{32} (CMS, 7 TeV) ▶ EPJC 73:2604 (2013)
- $R_{\Delta\phi}$ (ATLAS, 8 TeV) ▶ PRD 98:092004 (2018)



$R_{\Delta\phi}$ observable definition

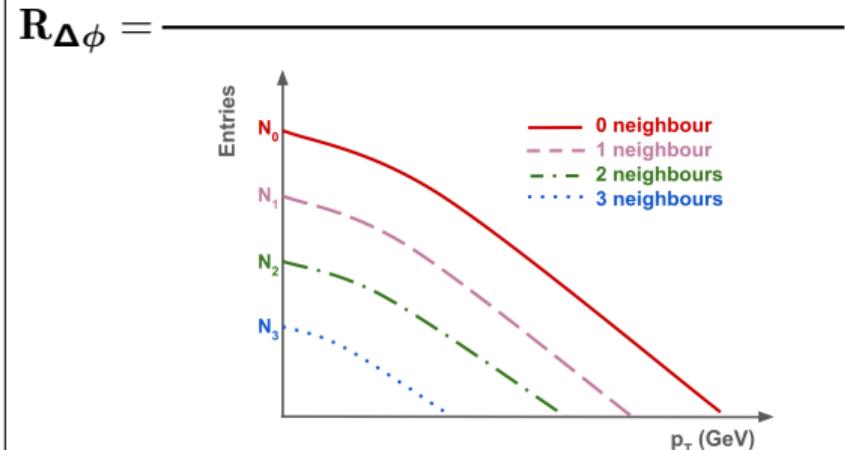
$$R_{\Delta\phi} = \frac{N_{\text{jet}}(p_T)}{\text{Number of jets in a jet } p_T \text{ bin } (\sim \alpha_s^2 \text{ @LO})}$$

Number of jets in a jet p_T bin ($\sim \alpha_s^2$ @LO)

2-jet topology



$R_{\Delta\phi}(p_T)$ entries
 $\Delta\phi \approx \pi$
 Numerator: 0
 Denominator: 2



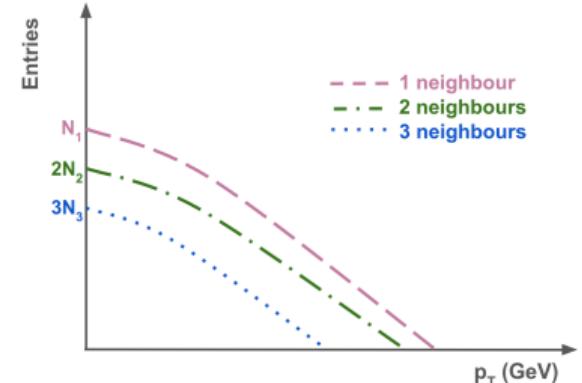
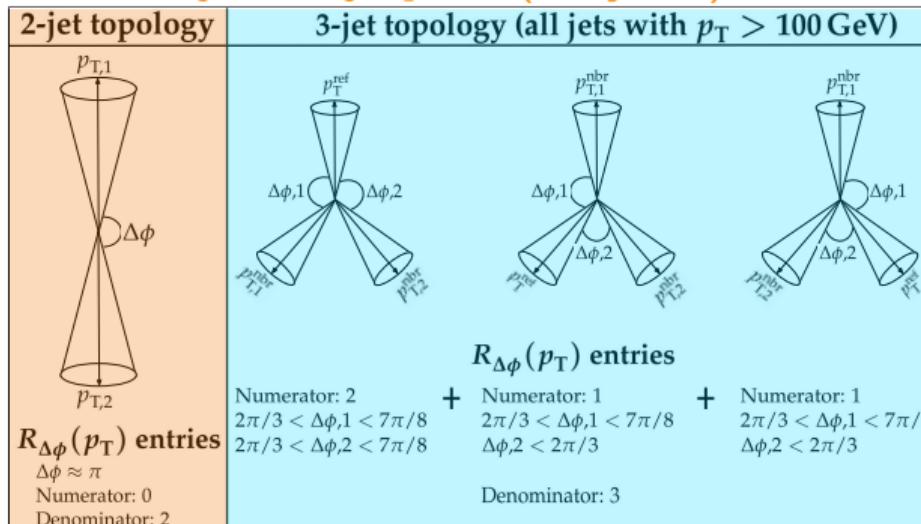
$R_{\Delta\phi}$ observable definition

Jets with neighbours within azimuthal separation:

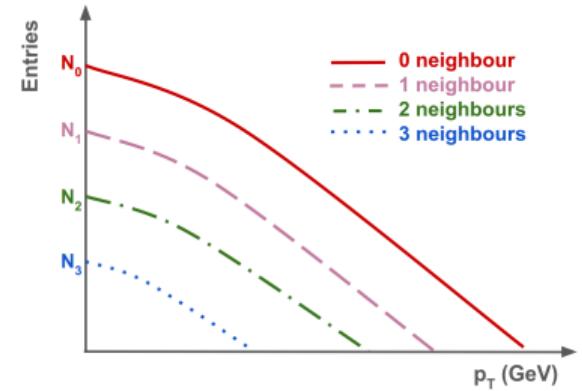
$2\pi/3 < \Delta\phi < 7\pi/8$ and $p_T > 100$ GeV ($\sim \alpha_s^3$ @LO)

$$R_{\Delta\phi} = \frac{\sum_{i=1}^{N_{\text{jet}}(p_T)} N_{\text{nbr}}^{(i)}(\Delta\phi, p_{T,\min}^{\text{nbr}})}{N_{\text{jet}}(p_T)}$$

Number of jets in a jet p_T bin ($\sim \alpha_s^2$ @LO)



$$R_{\Delta\phi} = \frac{\dots}{\dots}$$



Event selection - α_s sensitivity

Experimental data

- Full Run 2: $\mathcal{L}_{\text{int}} = 134 \text{ fb}^{-1}$
- Event sample with jets:
 - anti- k_T with $R = 0.7$
 - $p_T > 50 \text{ GeV}, |y| < 2.5$
- Numerator selection criteria:
 - $(\Delta\phi_{\min}, \Delta\phi_{\max}) = (2\pi/3, 7\pi/8)$
 - $p_{T\min}^{\text{nbr}} = 100 \text{ GeV}$

Results from: ▶ arXiv: 2404.16082

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Theoretical predictions

- Fixed-Order predictions pQCD @NLO.
- NLOJet++ (up to 3 jets @NLO)
- Using the **fastNLO** framework.
- $\mu_R = \mu_F = \hat{H}_T/2$, with $\hat{H}_T = \sum_{i \in \text{partons}} p_{T,i}$

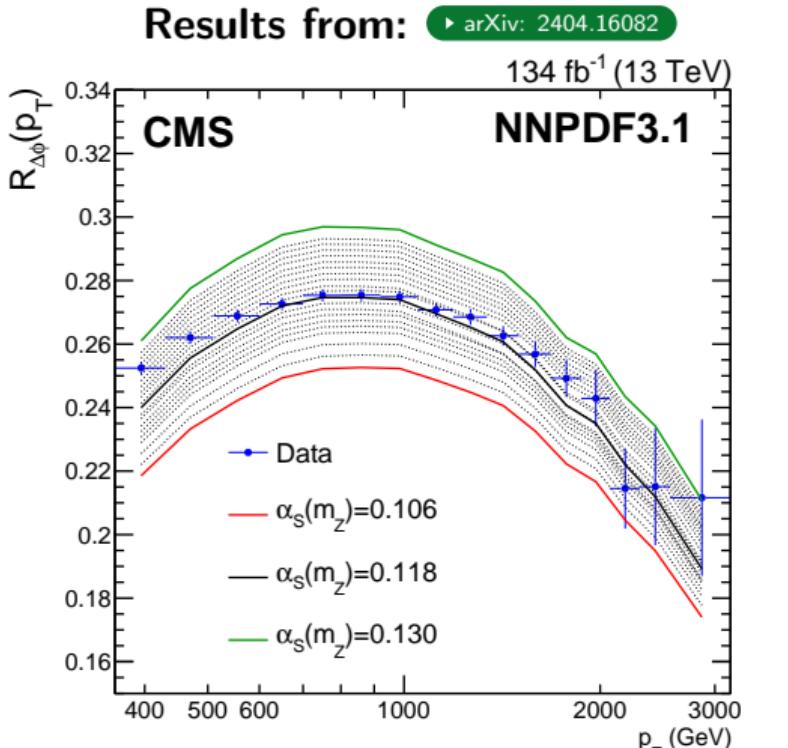
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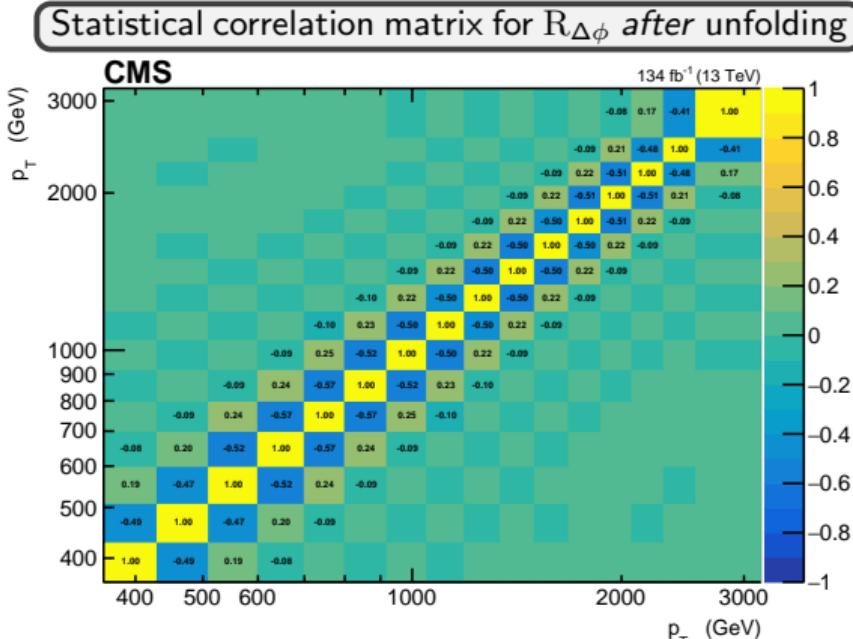
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$R_{\Delta\phi}$ observable has large sensitivity to α_s .

Measurement: uncertainties



- **Statistical** (0.18 – 10.49 %): from the covariance matrix *after unfolding*.
- **JES** (0.65 – 5.00 %): Jet Energy Scale uncertainty sources $\rightarrow p_T = p_T(1 \pm \text{unc. source})$.
- **JER** (0.04 – 0.77 %): Jet Energy Resolution smearing process applied to MC samples.
- **Other** (< 1%): Prefiring corrections, PU profile reweighting, MC modeling.

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- **Fixed Order predictions up to NLO in pQCD**
 - 4 different NLO PDF sets.
- **PDF uncertainties**
 - 68% CL Hessian/MC methods.
- **Scale uncertainties**
 - $\frac{1}{2} \leq \mu_R/\mu_F \leq 2$

PDFs available via LHAPDF

PDF set	Default $\alpha_s(m_z)$	Alternative $\alpha_s(m_z)$
ABMP16	0.1191	0.114 - 0.123
CT18	0.1180	0.110 - 0.124
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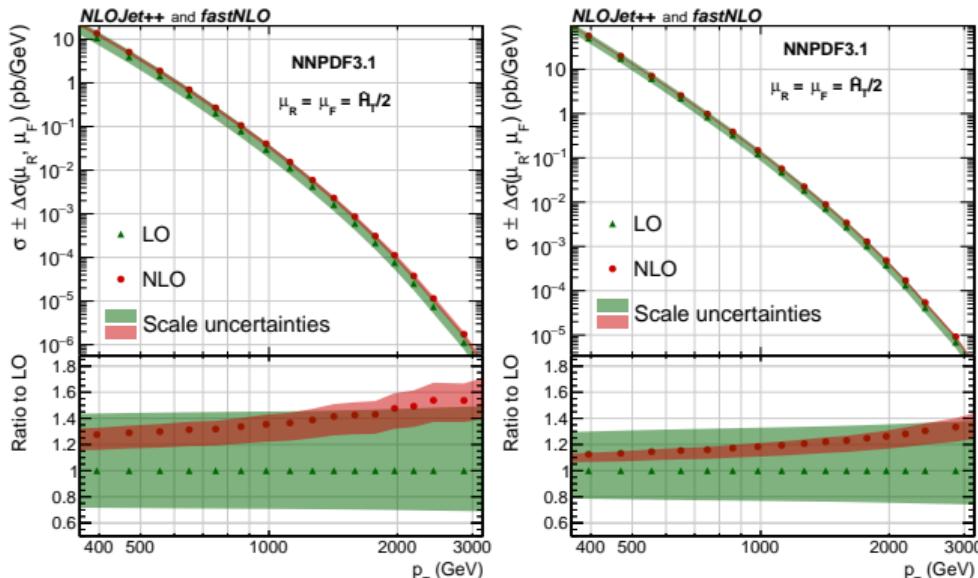
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NLO scale uncertainties

(3-jet) Numerator : 9 – 17 %
 (2-jet) Denominator : 5 – 10 %

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- Fixed Order pQCD predictions are available at parton level only: non-perturbative (NP) corrections account for multiple parton interactions (MPI) and hadronization (HAD) effects:

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- NP correction factors:

$$C^{\text{NP}} = \frac{\sigma^{\text{PS+MPI+HAD}}}{\sigma^{\text{PS}}}$$

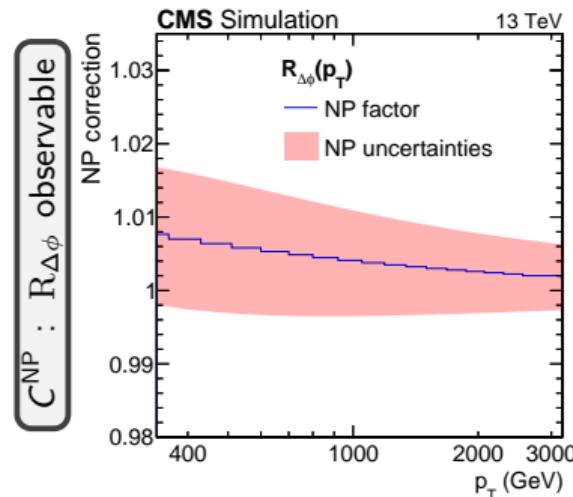
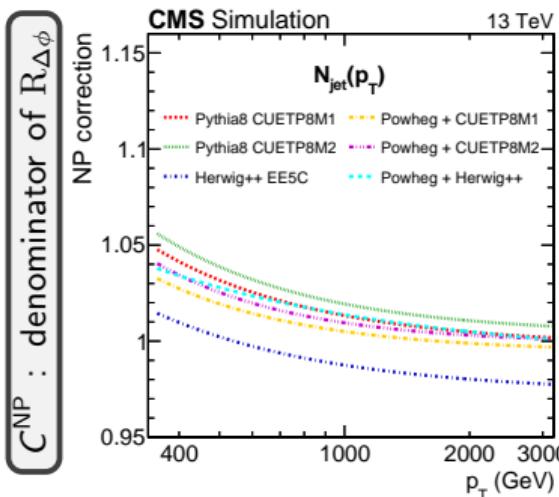
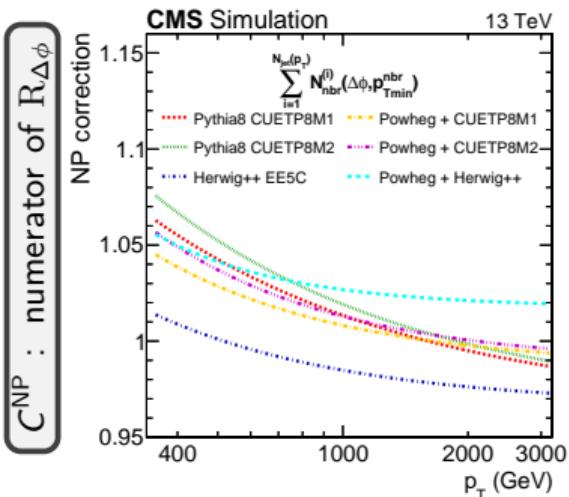
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ElectroWeak corrections

- Fixed Order pQCD predictions are also corrected for the **ElectroWeak (EW)** effects.

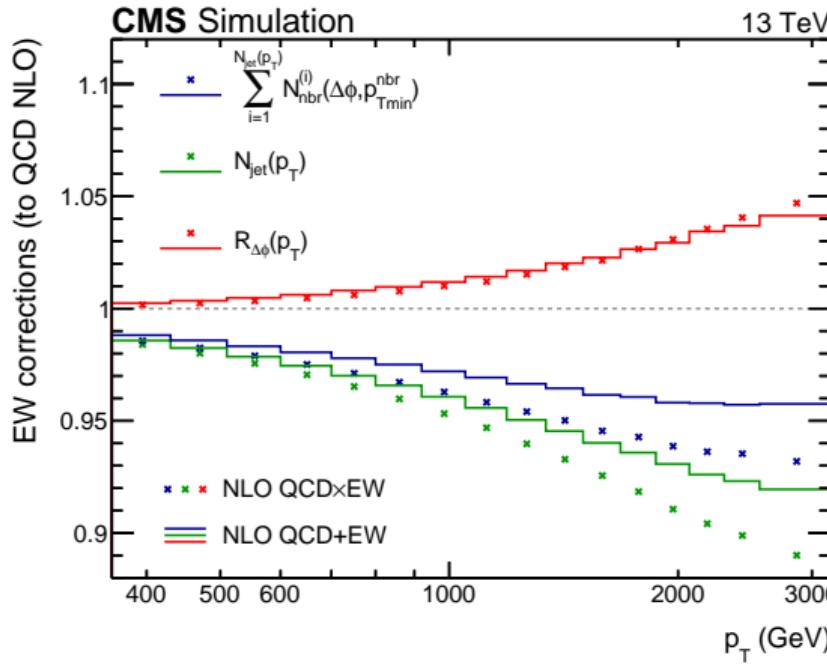
- *Full NLO corrections to 3-jet production and R₃₂ at the LHC*

M. Reyer, M. Schönherr, S. Schumann [► arXiv:1902.01763](#) → $\mathcal{O}(\alpha_s^n \alpha^m)$, with $n + m = 2$ and $n + m = 4$.

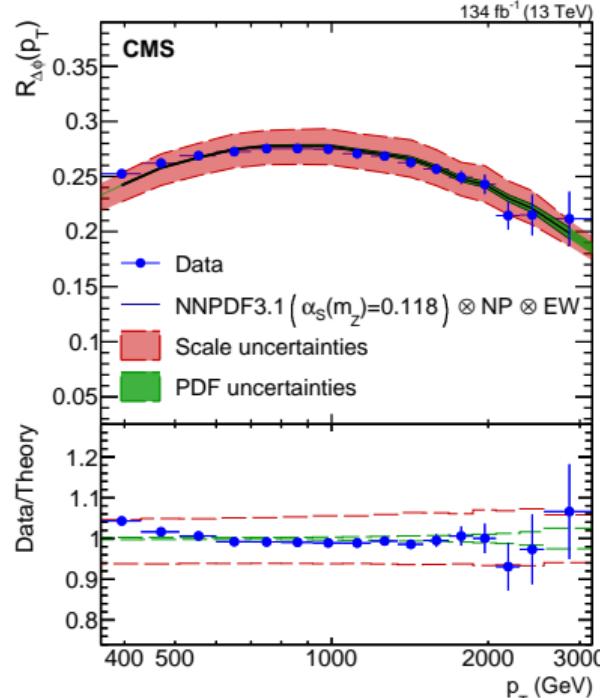
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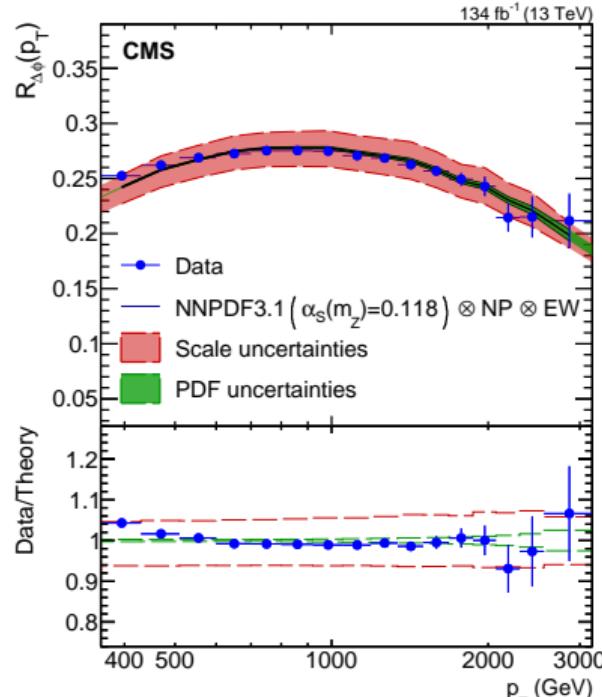
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EW corrections for $R_{\Delta\phi} < 5\%$ and EW correction uncertainties $< 0.6\%$.

Determination of $\alpha_s(m_z)$ 

- Agreement between data and theory (within the uncertainties) for all the PDF sets.
- Uncertainties: **PDF: 1 – 2 % , Scale: 2 – 8 %**

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Determination of $\alpha_s(m_z)$

- Minimization of:

$$\chi^2 = \sum_{ij} (D_i - T_i) C_{ij}^{-1} (D_j - T_j)$$

N : number of measurements

D_i : experimental data

T_i : theoretical predictions

C_{ij} : covariance matrix

- Covariance matrix composition:

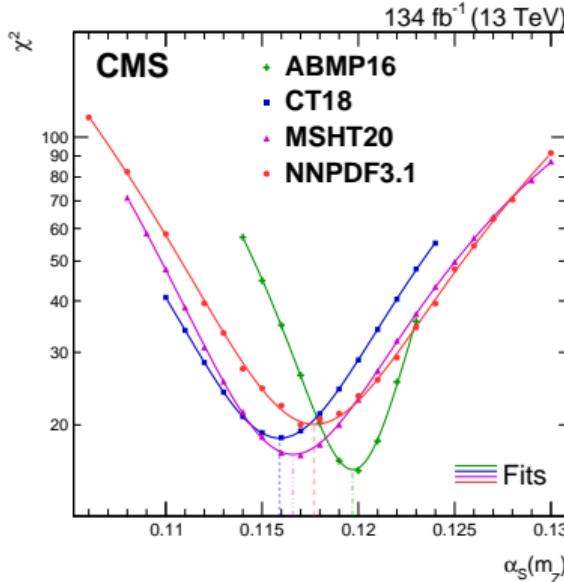
$$C_{ij} = C_{\text{uncor}} + C_{\text{exp}} + C_{\text{theo}}$$

C_{uncor} : numerical precision of FO predictions

C_{exp} : all the experimental uncertainties

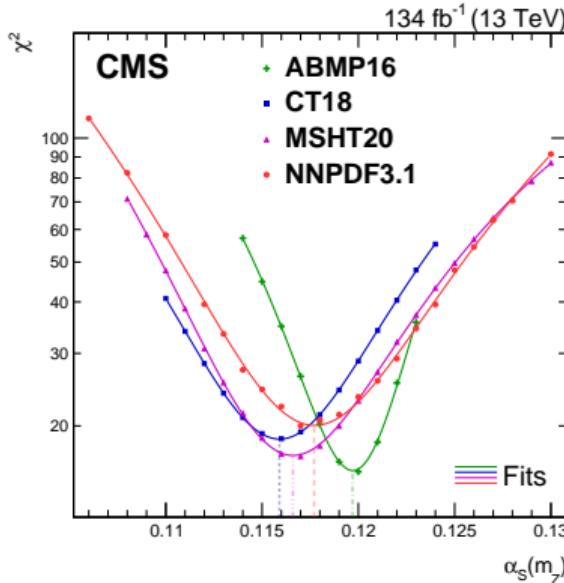
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ABMP16	0.1197	0.0008	0.0007	0.0007	0.0002	+0.0043 -0.0042	+0.0045 -0.0044	16/16
CT18	0.1159	0.0013	0.0009	0.0014	0.0002	+0.0099 -0.0067	+0.0101 -0.0070	19/16
MSHT20	0.1166	0.0013	0.0008	0.0010	0.0003	+0.0112 -0.0063	+0.0114 -0.0066	17/16
NNPDF3.1	0.1177	0.0013	0.0011	0.0010	0.0003	+0.0114 -0.0068	+0.0116 -0.0071	20/16

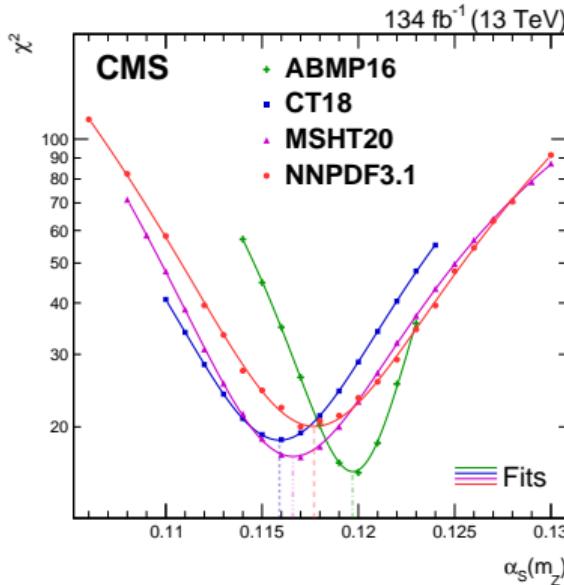
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- Extracted $\alpha_s(m_z)$ are compatible among each other within the uncertainties.

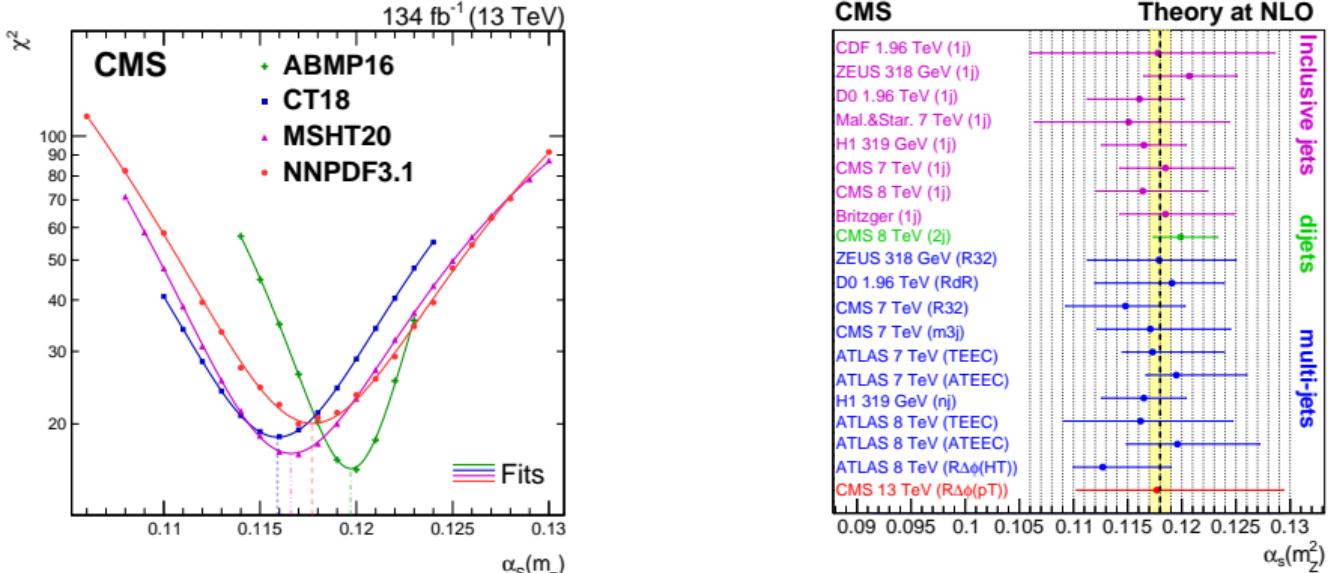
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- Scale uncertainties (theoretical) by far the dominant: 4 – 10 %.

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- Results also compatible with the world average: $\alpha_s(m_z) = 0.1180 \pm 0.0009$.

Evolution of α_s as a function of the energy scale (Q).

$\alpha_s(Q)$ evolutionEvolution of α_s as a function of the energy scale (Q).

- ① $\alpha_s(m_z)$ determination in 4 p_T ranges.
- ② $\langle Q \rangle$: cross-section-weighted average.
- ③ $\alpha_s(m_z)$ evolution to $\alpha_s(Q)$ via RGE.

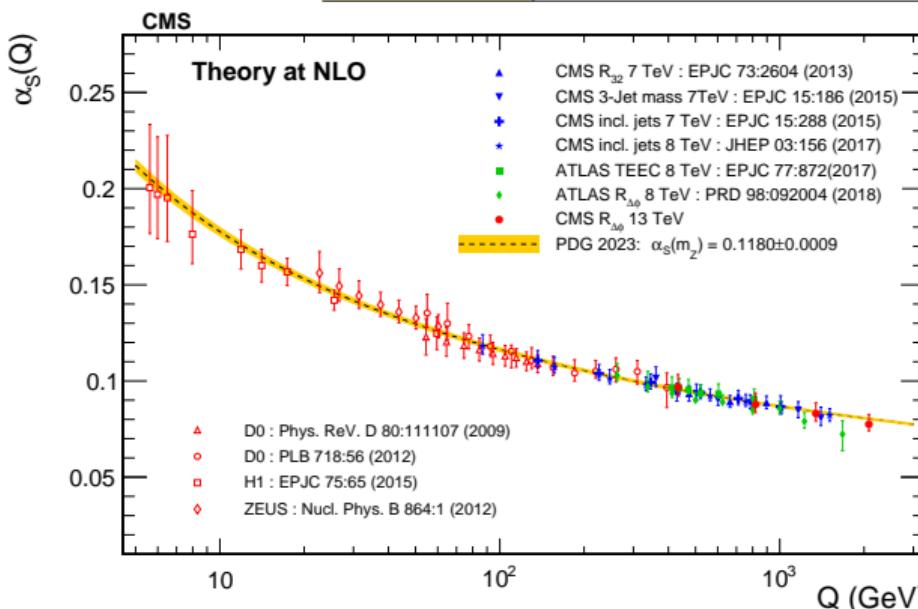
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360 – 700	$0.1177^{+0.0104}_{-0.0067}$	433.0	$0.0967^{+0.0066}_{-0.0044}$
700 – 1190	$0.1162^{+0.0108}_{-0.0073}$	819.0	$0.0878^{+0.0060}_{-0.0042}$
1190 – 1870	$0.1159^{+0.0112}_{-0.0077}$	1346.0	$0.0830^{+0.0055}_{-0.0040}$
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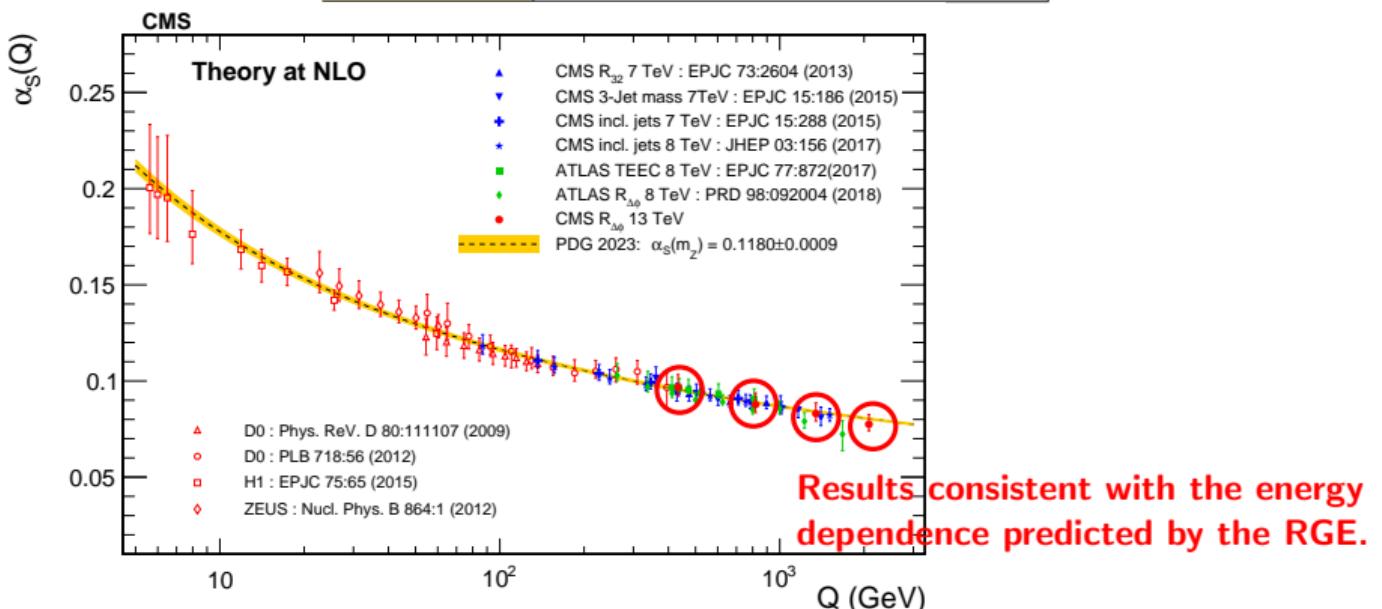


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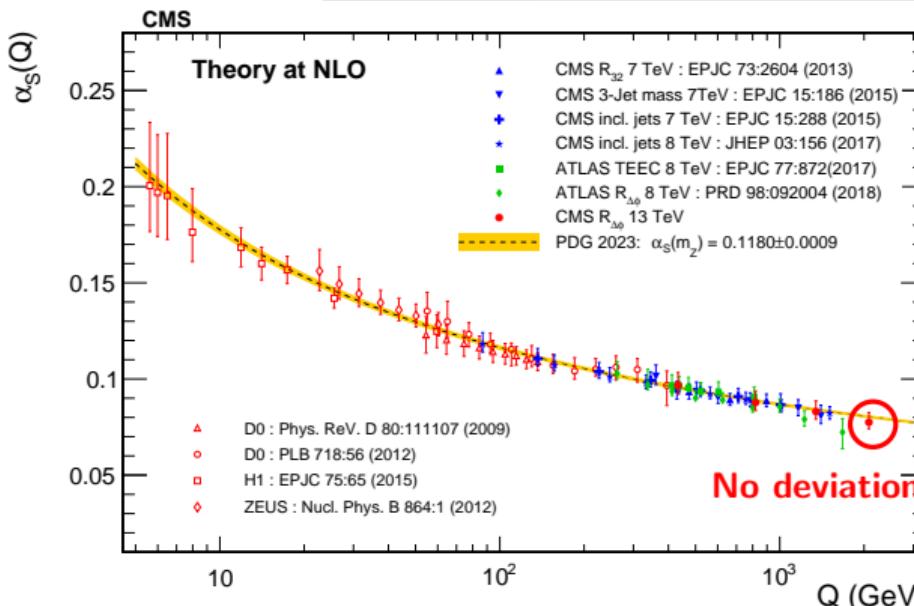


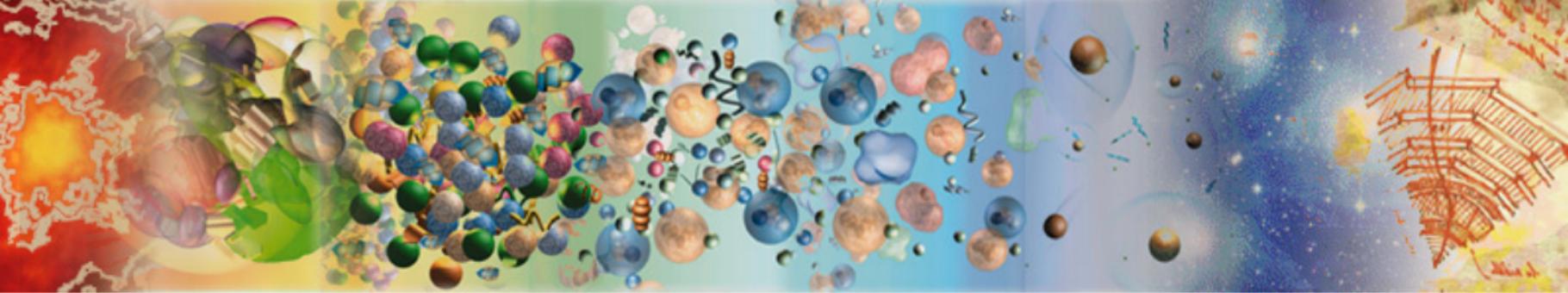
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THANK YOU FOR YOUR ATTENTION

BACK UP

$R_{\Delta R}$ and $R_{\Delta\phi}$ observables

(D0) $R_{\Delta R}$ definition

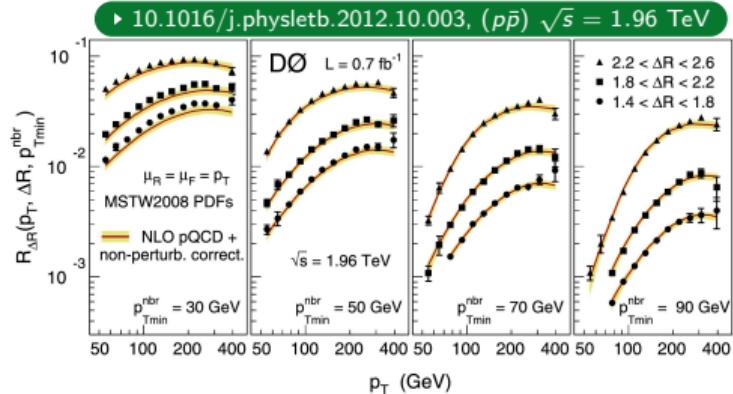
$$R_{\Delta R} = \frac{\sum_{i=1}^{N_{\text{jet}}(p_T)} N_{\text{nbr}}^{(i)}(\Delta R, p_{T\text{min}}^{\text{nbr}})}{N_{\text{jet}}(p_T)}$$

- N_{jet} : inclusive number of jets.
- N_{nbr} : jets with neighbours within angular separation interval ΔR and $p_T > p_{T\text{min}}^{\text{nbr}}$.

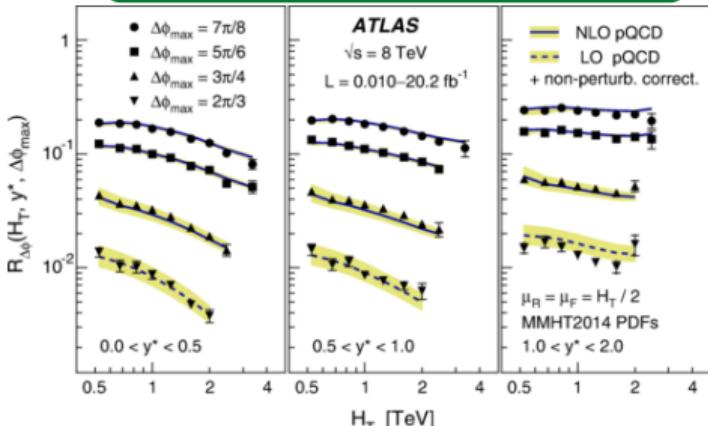
(ATLAS) $R_{\Delta\phi}$ definition

$$R_{\Delta\phi} = \frac{\frac{d^2\sigma_{\text{dijet}}(\Delta\phi_{\text{dijet}} < \Delta\phi_{\text{max}})}{dH_T dy^*}}{\frac{d^2\sigma_{\text{dijet (inclusive)}}}{dH_T dy^*}}$$

- σ_{dijet} : inclusive dijet sample.
- H_T : $\sum_{i \in \text{jets}} p_{T,i}$, y^* : $|y_1 - y_2|/2$.



► 10.1016/j.physletb.2012.10.003, ($p\bar{p}$) $\sqrt{s} = 1.96$ TeV



Measurement: data unfolding

Equivalent observable definition

$$R_{\Delta\phi} = \frac{\sum_{i=1}^{N_{\text{jet}}(p_T)} N_{\text{nbr}}^{(i)}(\Delta\phi, p_{T\min}^{\text{nbr}})}{N_{\text{jet}}(p_T)} = \frac{\sum_n n N(p_T, n)}{\sum_n N(p_T, n)}$$

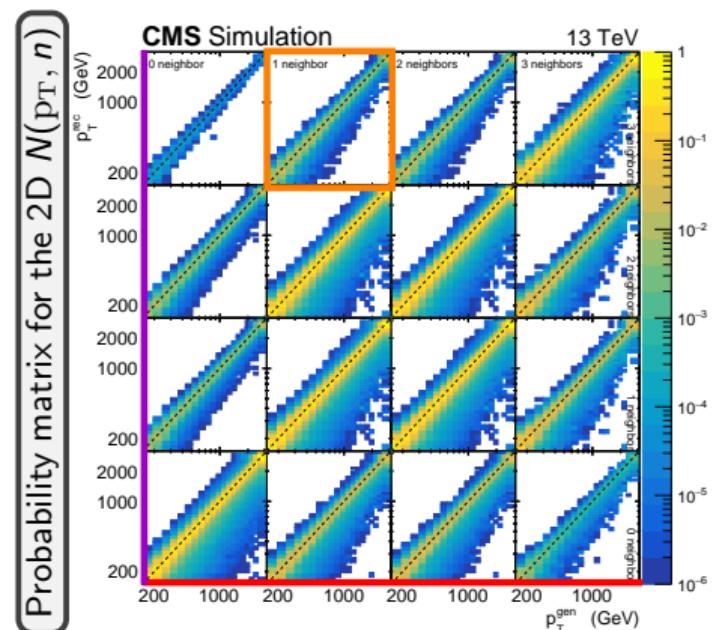
where n is the number of neighbours and p_T is jet's transverse momentum.

Motivation

- ✓ 2D unfolding of $N(p_T, n)$ distribution.
- ✓ Migrations among p_T **and** among n bins.
- ✓ Account for non-trivial numerator-denominator correlations.

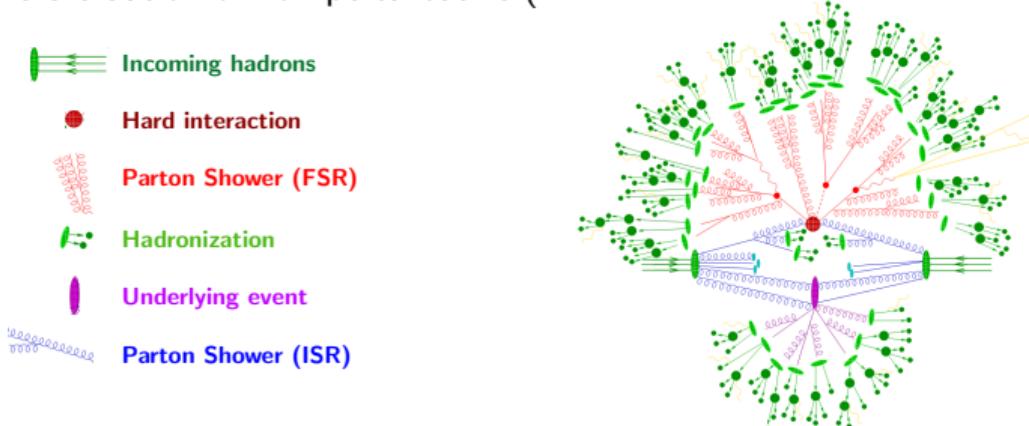
Matrix structure

- **x axis: generator-level p_T .**
- **y axis: reconstructed-level p_T .**
- **inner cells: neighbouring jet bins.**



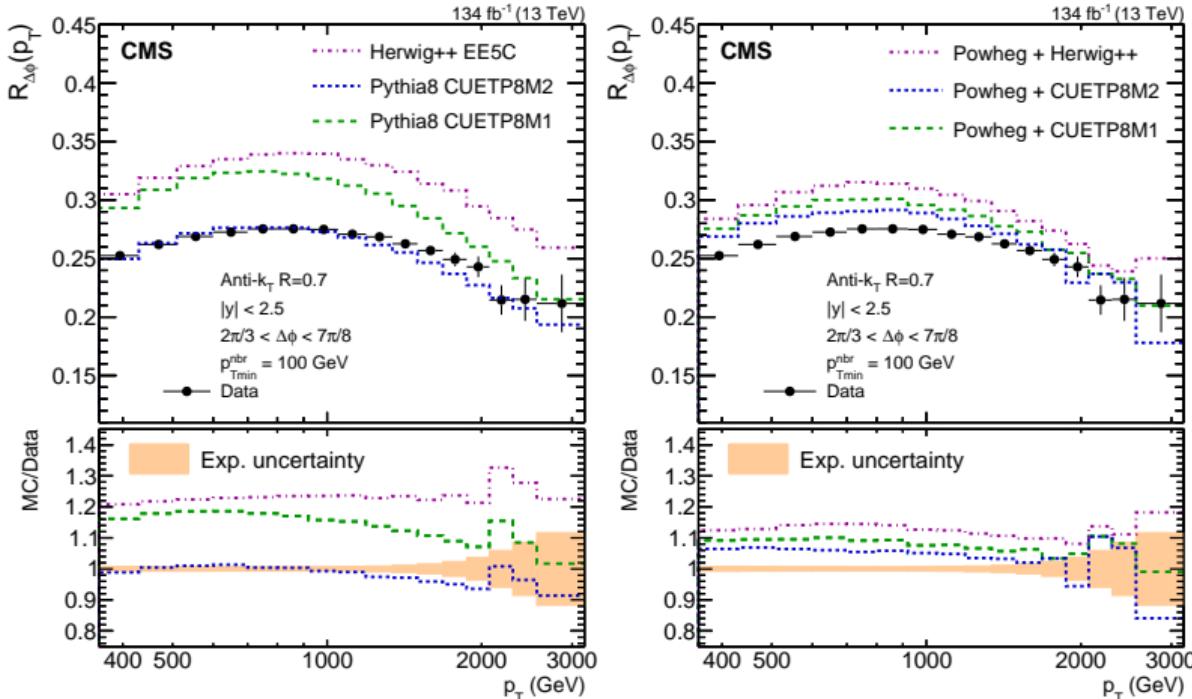
Monte Carlo predictions

- Predictions from MC event generators at particle level using RIVET toolkit.
 - For the comparison with experimental data.
 - For the evaluation of non-perturbative (NP) effects.



MC	Matrix Element	Parton Shower	Hadronization	Tune	PDF set
PYTHIA8	$2 \rightarrow 2$ (LO)	p_T ordered	Lund string	CUETP8M1	NNPDF2.3
				CUETP8M2	NNPDF3.0
HERWIG++	$2 \rightarrow 2$ (LO)	Angular ordered	Cluster model	EE5C	CTEQ6.1M
POWHEG	$2 \rightarrow 2$ (NLO), $2 \rightarrow 3$ (LO)	PYTHIA8 PYTHIA8 HERWIG++	PYTHIA8 PYTHIA8 HERWIG++	CUETP8M1 CUETP8M2 EE5C	NNPDF3.0 NNPDF3.0 NNPDF3.0

Monte Carlo predictions



- Predictions from Powheg overestimate the measurement by $\sim 5\text{-}12\%$.
- Herwig++ EE5C (Pythia8 CUETP8M1) overestimate $R_{\Delta\phi}$ by $\sim 20\%$ ($\sim 12\text{-}18\%$).
- Nice description from (LO) Pythia8 tune CUETP8M2T4.

Electroweak corrections

Combination of QCD and EW corrections

Pure NLO EW corrections for n-jet:

$$\sigma_{nj}^{\text{NLO EW}} = \sigma_{nj}^{\text{LO}} + \sigma_{nj}^{\Delta\text{NLO}_1}$$

ΔNLO_1 : virtual and real EW corrections.

Combination to QCD process:

- ① **Additive:** $\sigma_{nj}^{\text{NLO QCD+EW}}$

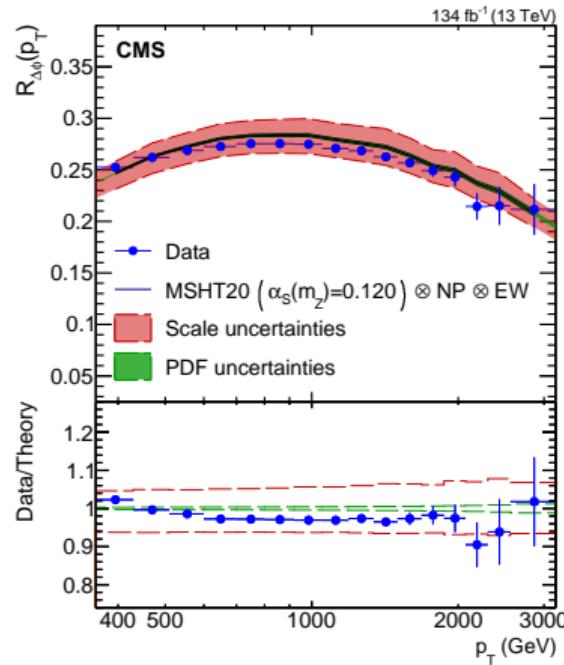
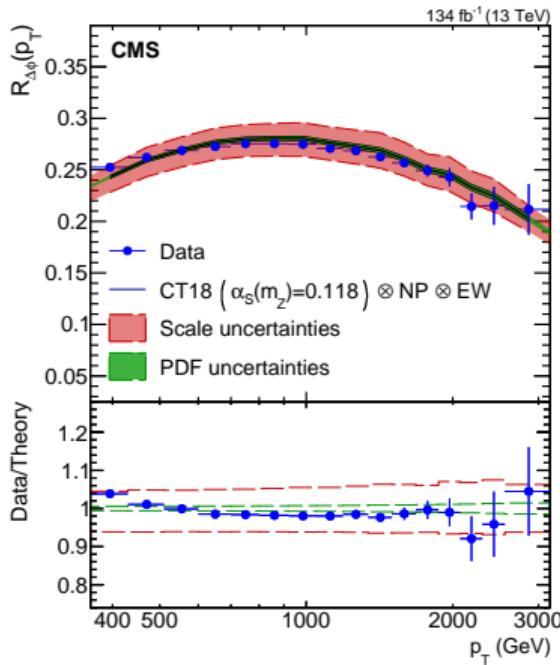
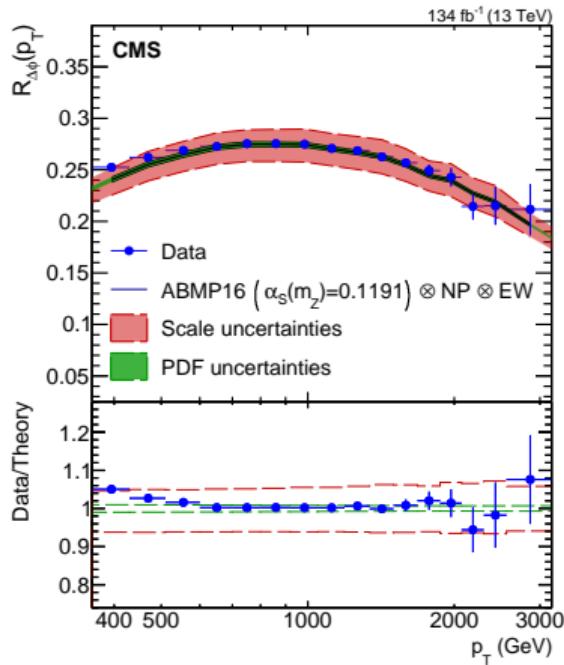
$$\sigma_{nj}^{\text{LO}} + \sigma_{nj}^{\Delta\text{NLO}_0} + \sigma_{nj}^{\Delta\text{NLO}_1}$$

ΔNLO_0 : virtual and real QCD corrections.

- ② **Multiplicative:** $\sigma_{nj}^{\text{NLO QCD}\times\text{EW}}$

$$\sigma_{nj}^{\text{LO}} \left(1 + \frac{\sigma_{nj}^{\Delta\text{NLO}_0}}{\sigma_{nj}^{\text{LO}}} \right) \left(1 + \frac{\sigma_{nj}^{\Delta\text{NLO}_1}}{\sigma_{nj}^{\text{LO}}} \right)$$

Data vs Theory



Sensitivity to α_s

