

ICHEP, Prague, 17-24 Jul. 2024

***bbH* as a background for *HH*:**

taming a leading theoretical uncertainty in *HH*  
measurements via accurate simulation of *bbH* production

JHEP 09 (2023) 179

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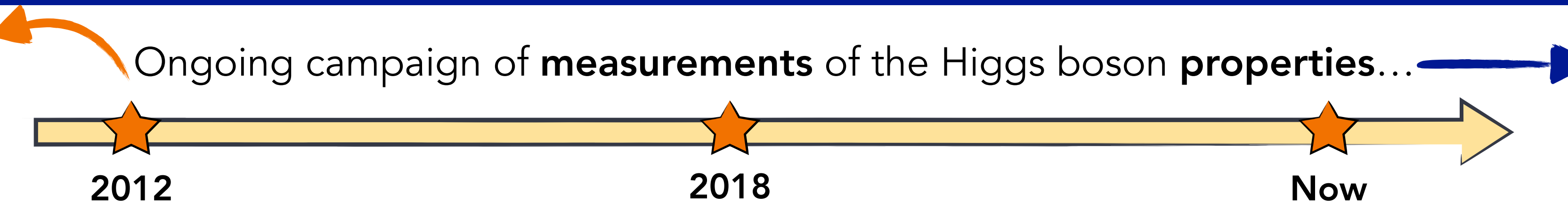


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# Outline

Observation of the Higgs boson by ATLAS and CMS  
( $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ^* \rightarrow 4l$  channels)

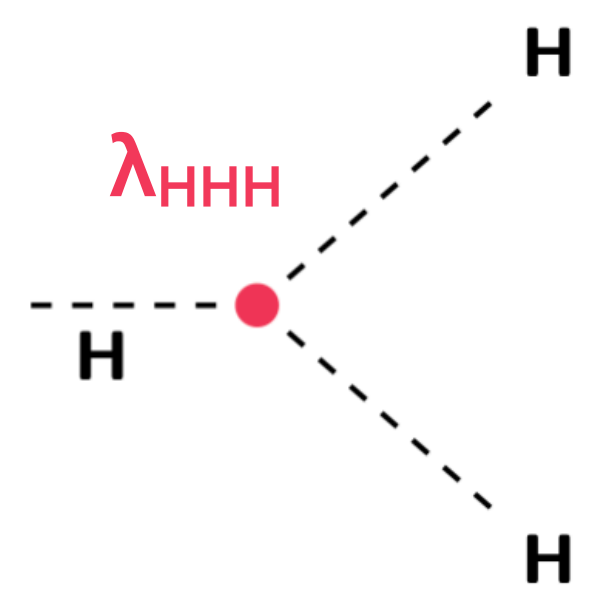
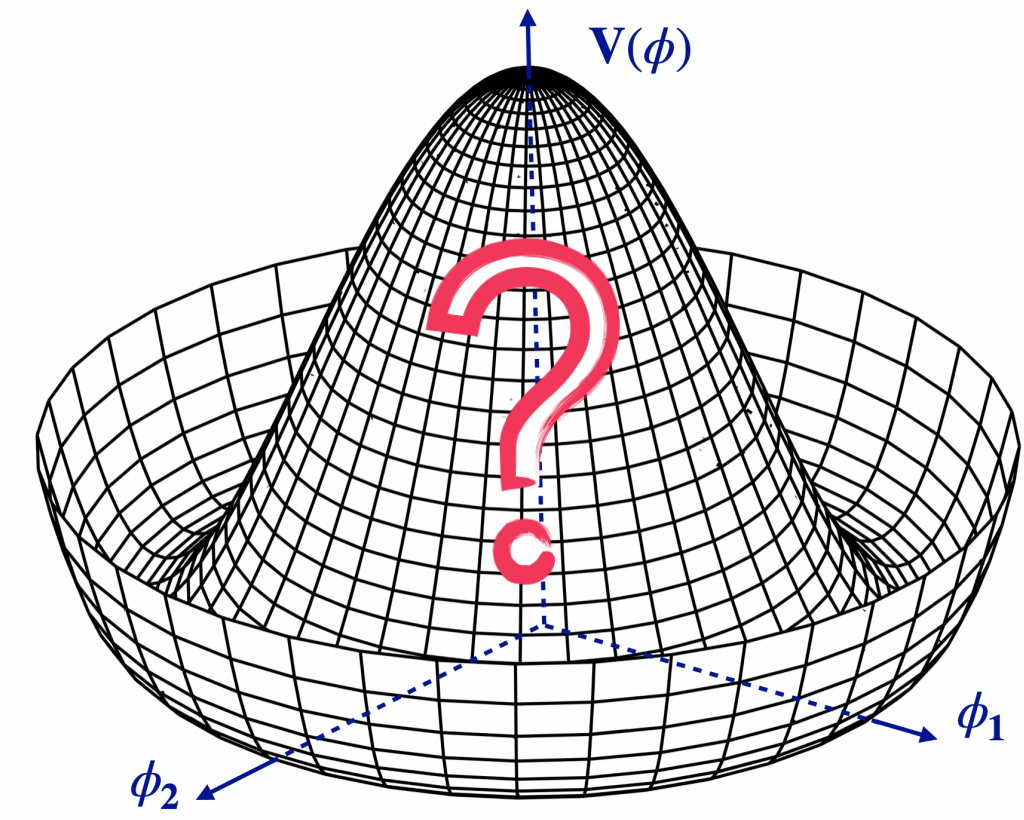


Excellent agreement with the **Standard Model (SM)!**

Observation of the  $H \rightarrow b\bar{b}$  decay

More and more **pp collision data** delivered by the LHC during **Run 3!**

- Precision of measurements will increase.
- New couplings will become accessible! The **Higgs boson trilinear self-coupling  $\lambda_{HHH}$ !**



- Measuring  $\lambda_{HHH}$  allows to shed light on the **shape** of the **Higgs boson potential**.
- The **only direct way** to access  $\lambda_{HHH}$  is via **Higgs boson pair production (HH)!**

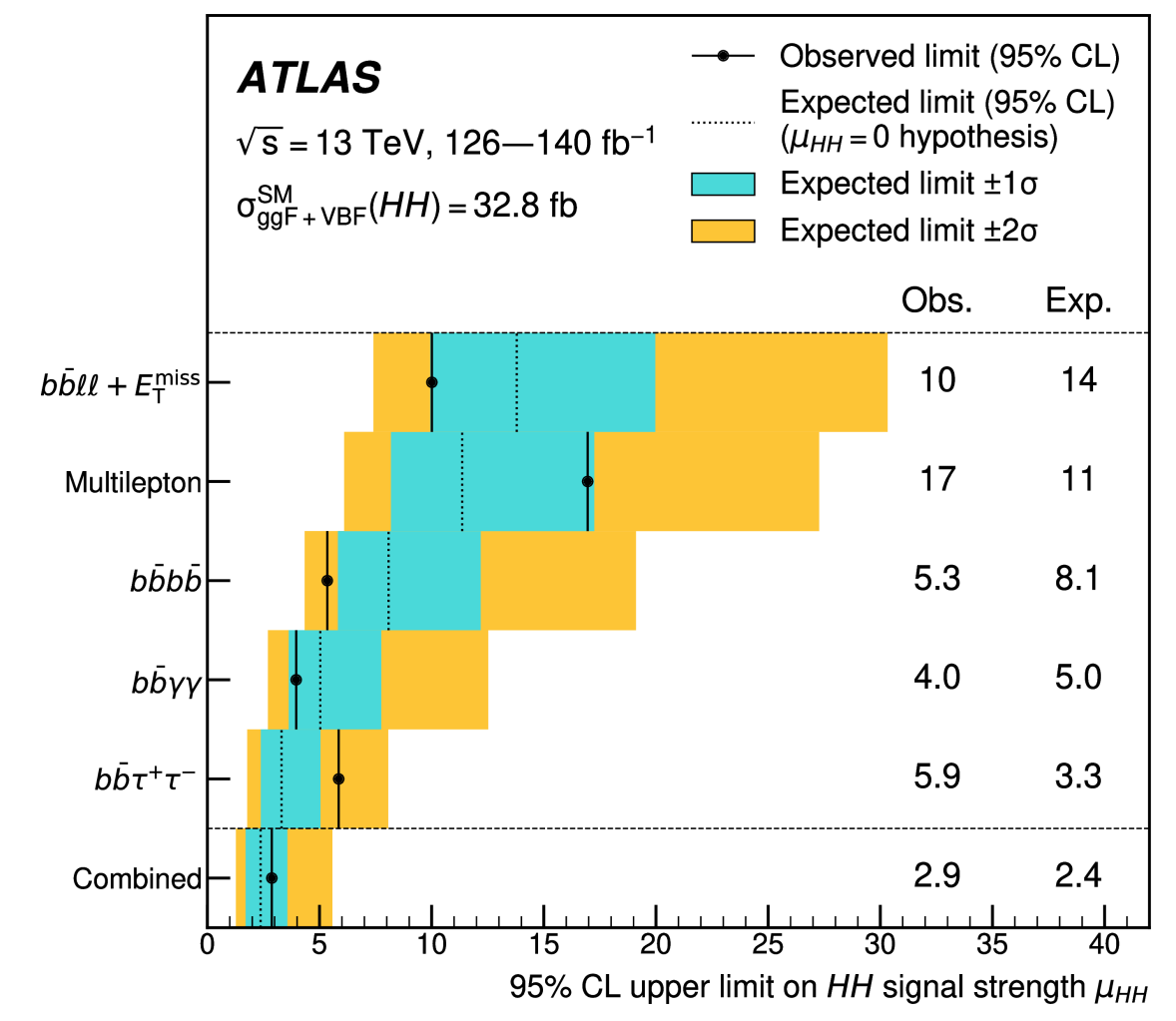
**Extremely rare process!** In the SM,  $\sigma(HH) \approx 10^{-3} \times \sigma(H)$ !

- **No hint of HH** with **Run 2** data. But we are getting **close to SM-like sensitivity...**

Current observed (expected) **constraints** on **HH** production:

- **ATLAS** =  $\sigma(HH) < 2.9$  (2.4)  $\times \sigma^{SM}(HH)$ ;
- **CMS** =  $\sigma(HH) < 3.4$  (2.5)  $\times \sigma^{SM}(HH)$ .

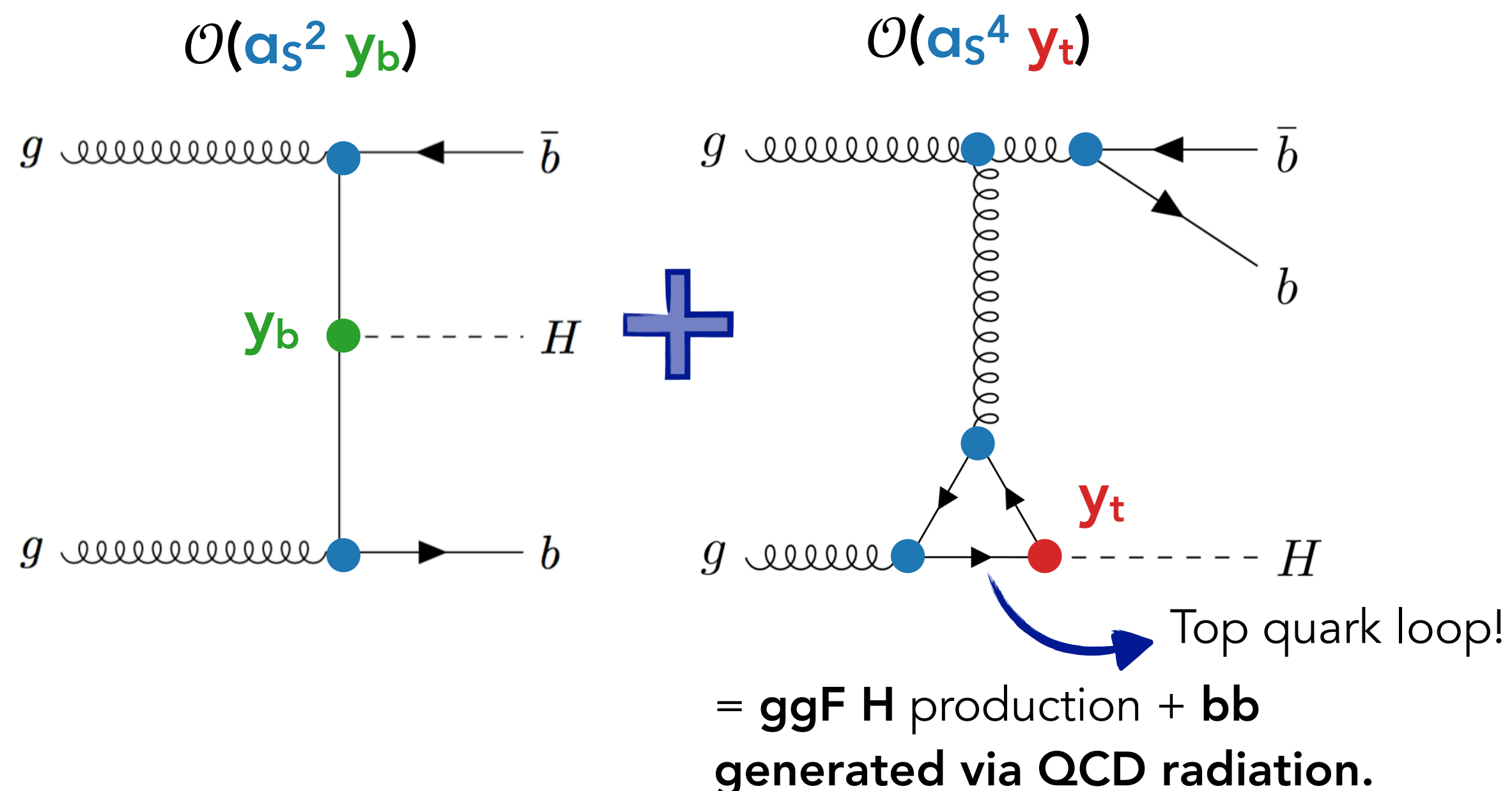
- According to current projections, an **observation** of **SM HH** production could be possible only at the end of **HL-LHC**.



# The $bbH$ background for $HH$ searches

The **most sensitive**  $HH$  channels all involve the  $H \rightarrow bb$  decay.

➔ **Single Higgs boson production +  $bb$**  (=  $bbH$ ) is an irreducible bkg. for all these searches!



➔

$$d\sigma(bbH) \propto |\mathcal{O}(\alpha_s^2 y_b) + \mathcal{O}(\alpha_s^4 y_t)|^2$$

$$\propto \mathcal{O}(y_b^2) + \mathcal{O}(y_t^2) + \mathcal{O}(y_b y_t).$$

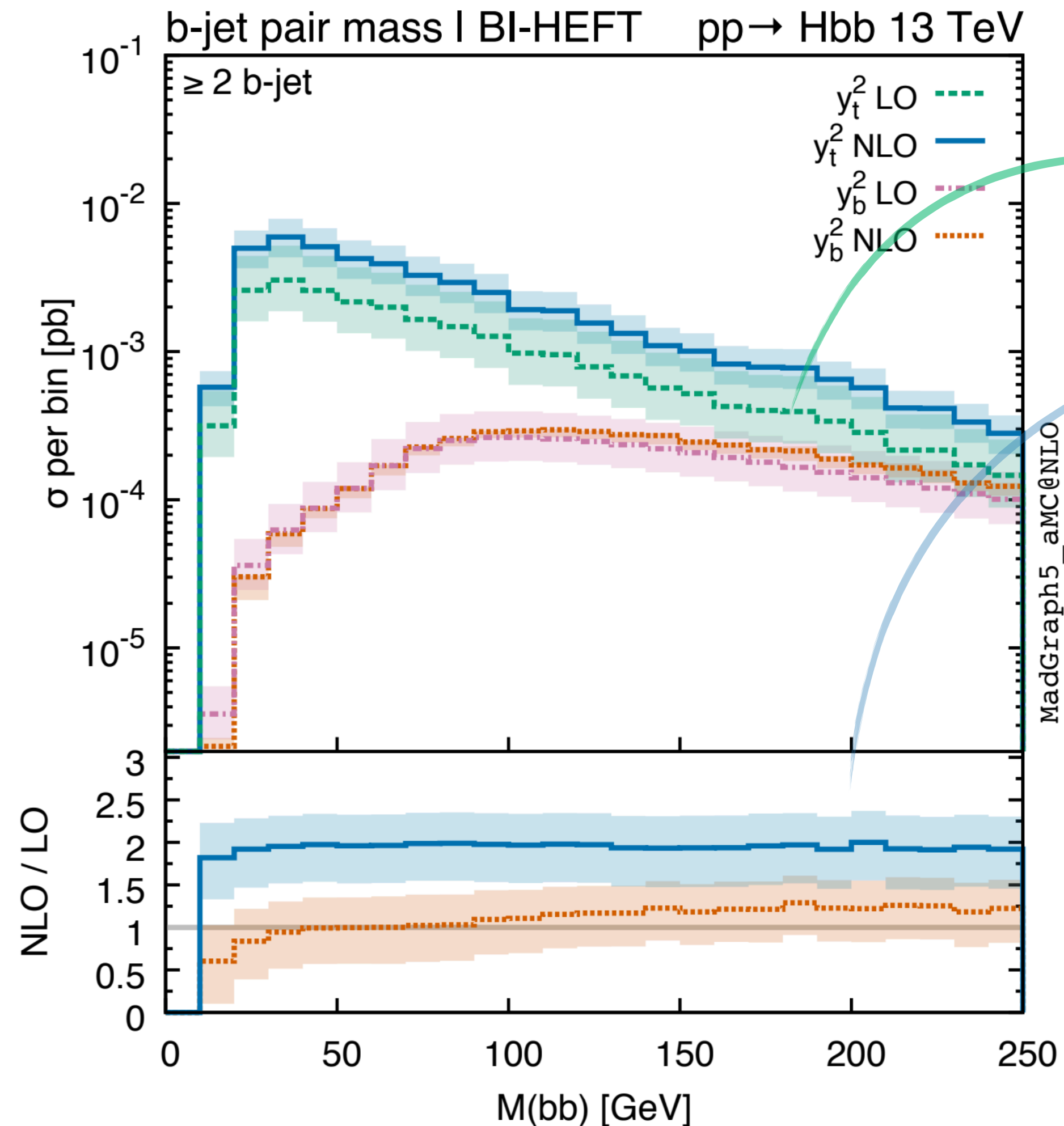
- $bbH$  rate comparable to **SM  $HH$  signal** in  $HH$  phase space.
- **$y_t^2$  contribution** currently simulated by the experiments using **ggF NNLOPS (ATLAS)** or **MinLO (CMS)**:
  - ➔ Within 5-flavor scheme (5FS) = Massless bottom quarks.
  - Only at most LO-accurate for **ggF Higgs + 2 jets**.
- ATLAS  $HH$  searches assign a **100% uncertainty** to the **ggF H +  $bb$  background**.
  - ➔ Also motivated by data - MC disagreement in other analyses (see [JHEP 08 \(2022\) 027](#) and [Eur. Phys. J. C 80 \(2020\) 942](#)).
  - Now the **primary systematic uncertainty** for  $HH$  searches, affecting the [current constraints on  \$HH\$  production from ATLAS](#) by **25%**!

=

Crucial to have **better description** of the  **$bbH$  background** for the **next generation of  $HH$  searches!**

The (fiducial) rates for the  $bbH$  process was computed @ NLO in the **4-flavor scheme (4FS)**, using MadGraph5\_aMC@NLO.

➔ Including the  $\propto y_b^2$  and the  $\propto y_t^2$  contributions and the interference ( $\propto y_b y_t$ ). ➔ = Massive bottom quarks.



- The  $y_t^2$  contribution to the  $bbH$  cross-section is **dominant already @ LO**.  
➔ Becomes even **larger @ NLO**, with **K-factors = 2 - 3**.
- **Large NLO / LO K-factors** for both  $y_t^2$  and  $y_b^2$  components, with **strong dependence** on the **differential distributions**.  
➔ Critical to have **reliable** predictions for the  $bbH$  process @ NLO.
- The **interference** term is **subdominant** (= **5-10%** w.r.t. the total).  
➔ Subleading w.r.t. the relatively large scale uncertainties ( $\approx$  **50%**).

# The $bbH$ process @ NLO + PS (4FS) in the phase space of $HH \rightarrow bb\gamma\gamma$

We present a simulation of  $bbH$  process @ NLO + matching to parton shower (PS) in the 4FS + [JHEP 09 \(2023\) 179](#)  
dedicated analysis targeting the  $HH$  phase space!  $\longrightarrow$  Using the  $HH \rightarrow bb\gamma\gamma$  search as a representative case.

## Setup

- Same approach & settings as [JHEP 07 \(2019\) 054](#).

- $\longrightarrow$  -  $m_b = 4.92$  GeV,  $m_t = 172.5$  GeV,  $m_H = 125$  GeV.  
- PDF set NNPDF31\_nlo\_as\_0118\_nf\_4.  
- Central scales  $\mu_R = \mu_F = \mu_{Sh} = H_T / 4$ .

- We simulate only the  $y_b^2$  and  $y_t^2$  contributions (= the  $y_b y_t$  interference is ignored).

- Fiducial cuts inspired by  $HH \rightarrow bb\gamma\gamma$  analysis ([JHEP 01 \(2024\) 066](#)).

- $\longrightarrow$  - Consider **di-photon decay** of the **Higgs boson**.  
- **Anti- $k_T$  jets** with  $R = 0.4$ , with  $p_T > 25$  GeV and  $|\eta| < 2.5$ .  
- At least **2 photons** and exactly **2 b-jets** (= jets containing at least one B-hadron).  
- **Photon cuts**:  $p_T(\gamma_{1(2)}) > 0.35$  (0.25)  $\times m_{\gamma\gamma}$ ,  $|\eta(\gamma)| < 2.37$ ,  $105 < m_{\gamma\gamma} < 160$  GeV.  $\longrightarrow$  Targeting the  $H \rightarrow \gamma\gamma$  kinematics.  
- **b-jet cuts**:  $80 < m_{bb} < 140$  GeV.  $\longrightarrow$  Targeting the  $H \rightarrow bb$  kinematics.

# The $bbH$ process @ NLO + PS (4FS) in the phase space of $HH \rightarrow bby\gamma$

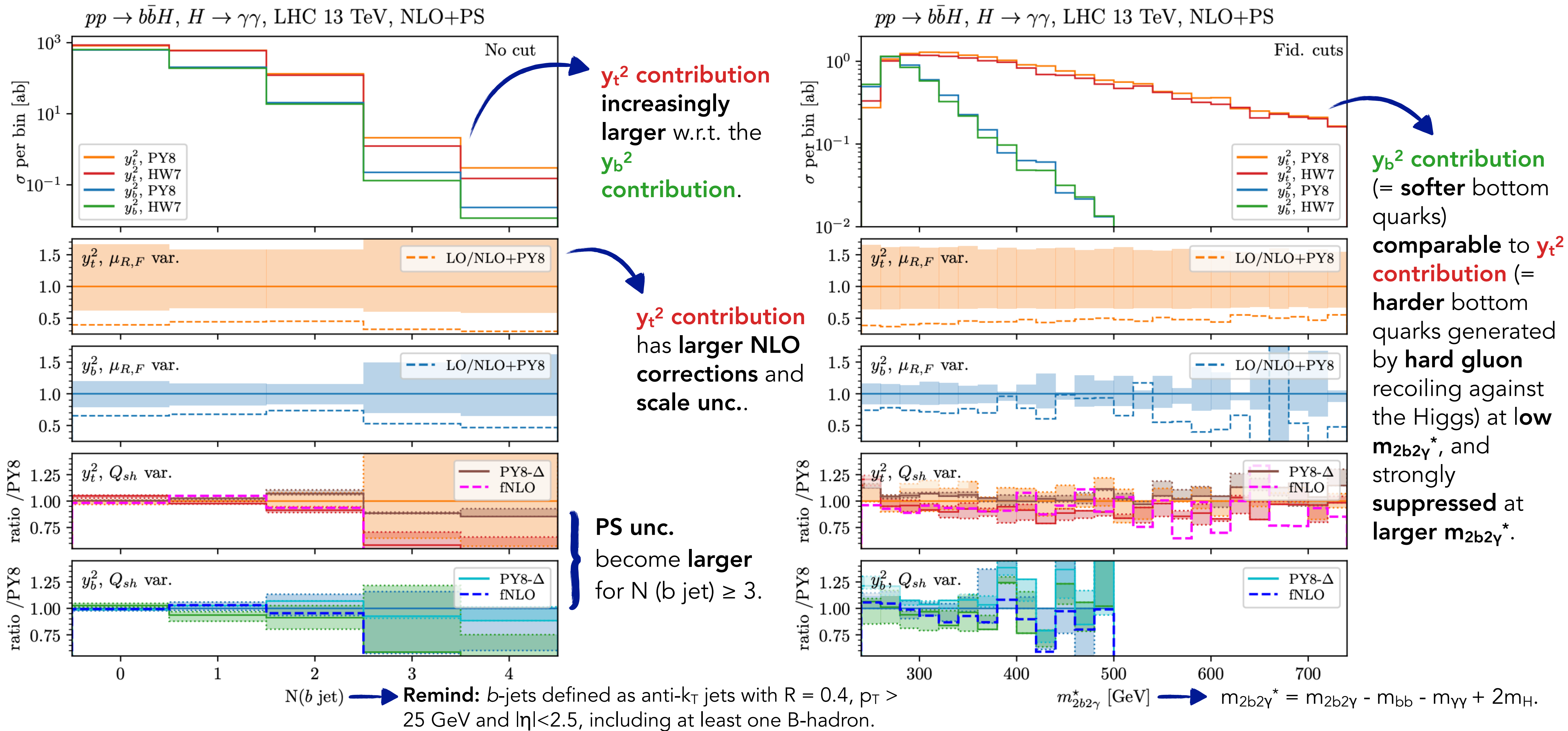
We present a simulation of  $bbH$  process @ NLO + matching to parton shower (PS) in the 4FS + dedicated analysis targeting the  $HH$  phase space!  $\longrightarrow$  Using the  $HH \rightarrow bby\gamma$  search as a representative case.

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Fiducial cross-sections									
Cut	Contr.	Run	LO	NLO	$\delta\mu_{R,F}$	$\delta Q_{sh}$	NNLOPS ( $y_t^2$ LO)	HH signal	
Fid. cuts	$y_b^2$	PY8	3.15	4.22	+15%	+10%	29.9	22.7	
		HW7	2.59	4.08	-15%	-4%			
	$y_t^2$	PY8	8.24	18.1	+58%	+10%			$g \rightarrow b\bar{b}$ : 17.2
		HW7	6.83	16.6	-34%	-7%			
	sum	PY8	11.4	22.3	+50%	+10%			
		HW7	9.42	20.7	-30%	-6%			

- The **NLO** corrections to the  $bbH$  rates are **very large!**  $\longrightarrow$ 
  - $y_b^2 = +50\%$ .
  - $y_t^2 = +150\%$ .
- The **scale uncertainties** @ NLO are still **sizeable** (especially for  $y_t^2$ , where they are within  $\sim 60\%$ ).
- The **impact of the PS** (Pythia8 or Herwig7) on the  $bbH$  NLO rates is within  $\sim 10\%$  (= **subdominant** w.r.t. the **scale uncertainties!**).
- The  $bbH$  background rate is **comparable** with the **SM HH signal**.
- The **ggF 5FS NNLOPS** simulation provides  $bbH$  cross-sections **larger** by a **factor  $\sim 2$**  w.r.t. our 4FS  $bbH$  @ NLO prediction.
  - $\longrightarrow$  - Could be traced back to the  **$g \rightarrow b\bar{b}$  splittings** in the PS.
  - **When turning these off**, the **ggF 5FS NNLOPS rates drops to half** its nominal prediction.  $\longrightarrow$  More details in the next slides.

# Differential distributions

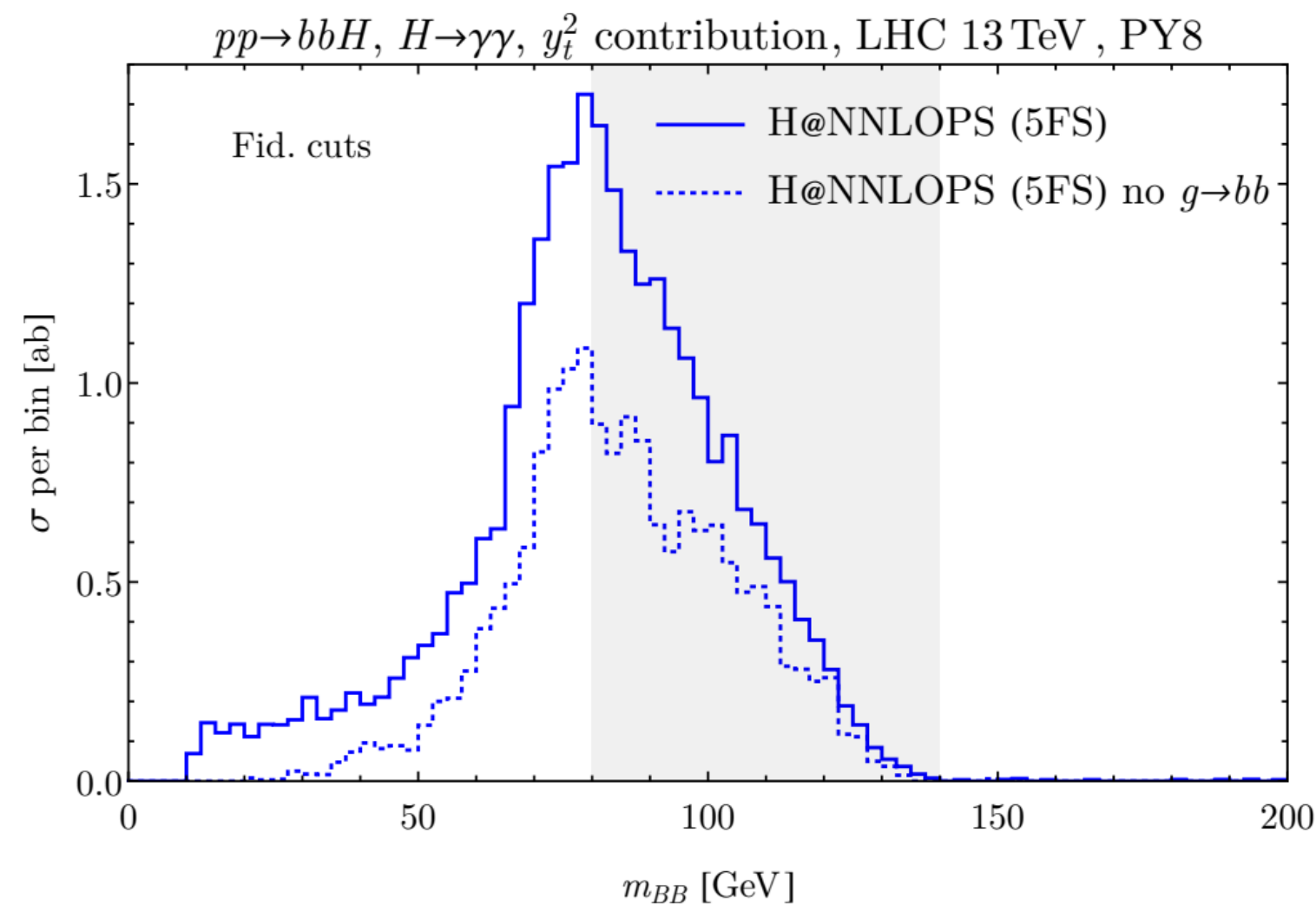
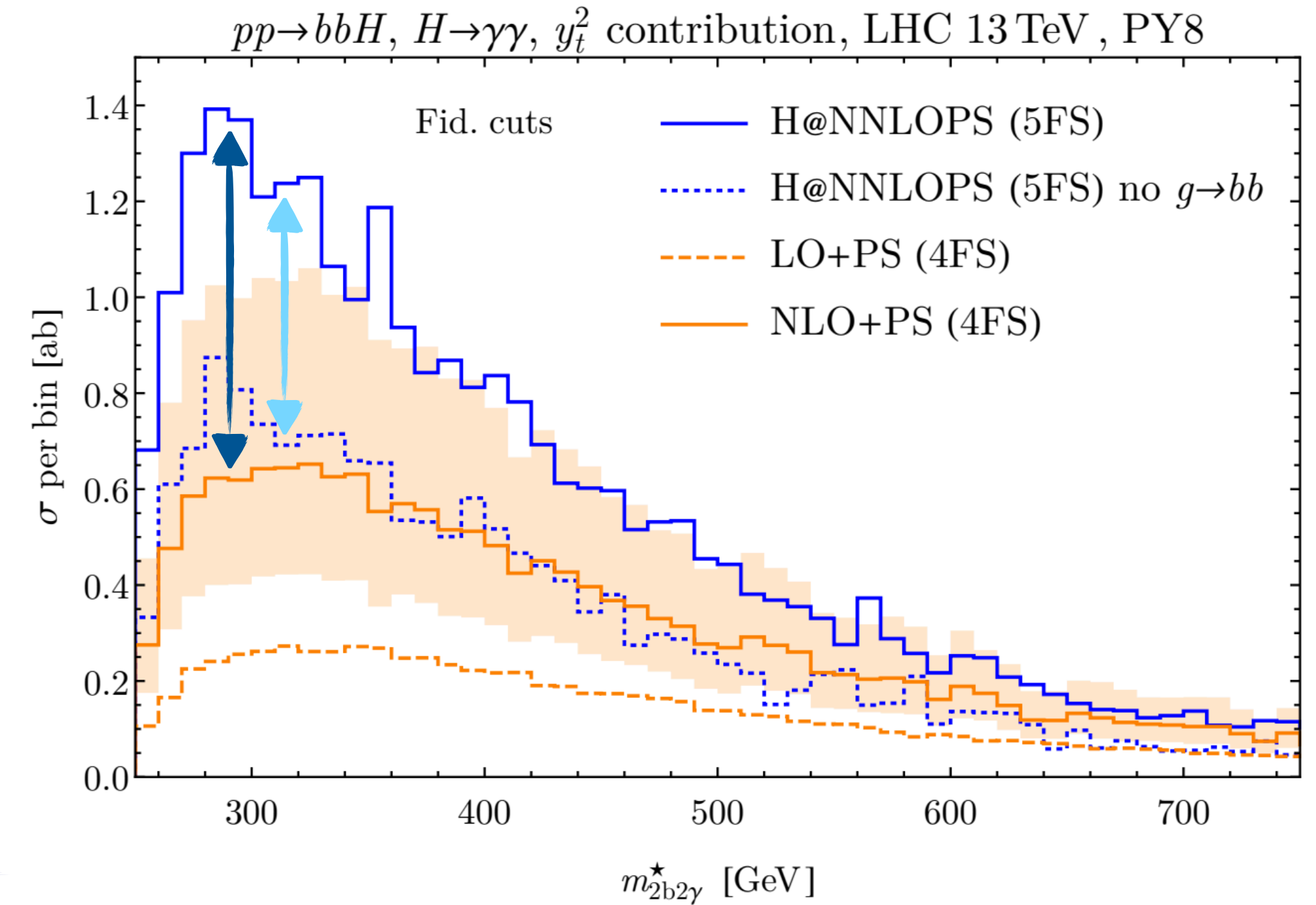


# Comparison and combination with inclusive NNLOPS prediction

- The  $bbH$  cross-sections provided by the **ggF 5FS NNLOPS** simulation are **larger** by a factor  $\sim 2$  w.r.t. our **4FS  $bbH$  @ NLO** prediction.
- This **discrepancy** is **recovered** if we **turn off** the  **$g \rightarrow bb$  splittings** in the PS.

➔ **Question:** why does the  **$g \rightarrow bb$  splittings** generate a large amount of events with **2 hard and central** ( $p_T > 25$  GeV,  $|\eta| < 2.5$ )  **$b$ -jets**?

- ➔
1. Maybe **soft wide-angle bottom quarks** from  **$g \rightarrow bb$  splittings** are **clustered** with **other hard partons** to form **two hard  $b$ -jets** (=  $pp \rightarrow H + bb + jj$  events)?
  2. Or maybe  **$g \rightarrow bb$  splittings** are generating **hard, wide-angle bottom quarks**?



**$g \rightarrow bb$  splittings** populate region with **high B-hadron invariant mass!**

- ➔
- **$g \rightarrow bb$  splittings** are generating events with **hard B-hadrons!**
  - Probably bottom quarks from  **$g \rightarrow bb$  splittings** are hard and well separated.
- ➔
- Kinematics poorly described in soft / collinear approximation of the PS.
  - The PS is acting outside of its validity range.

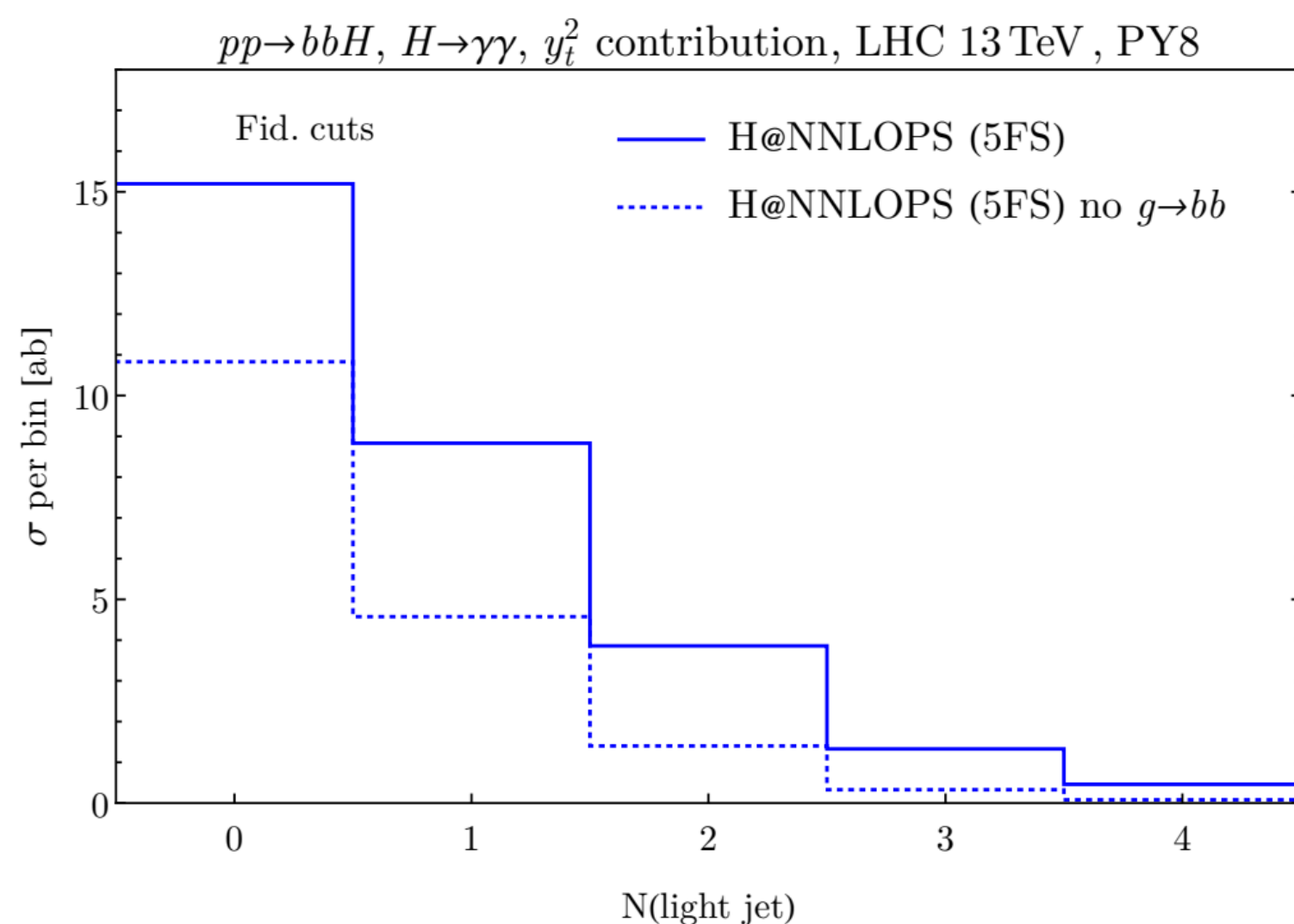
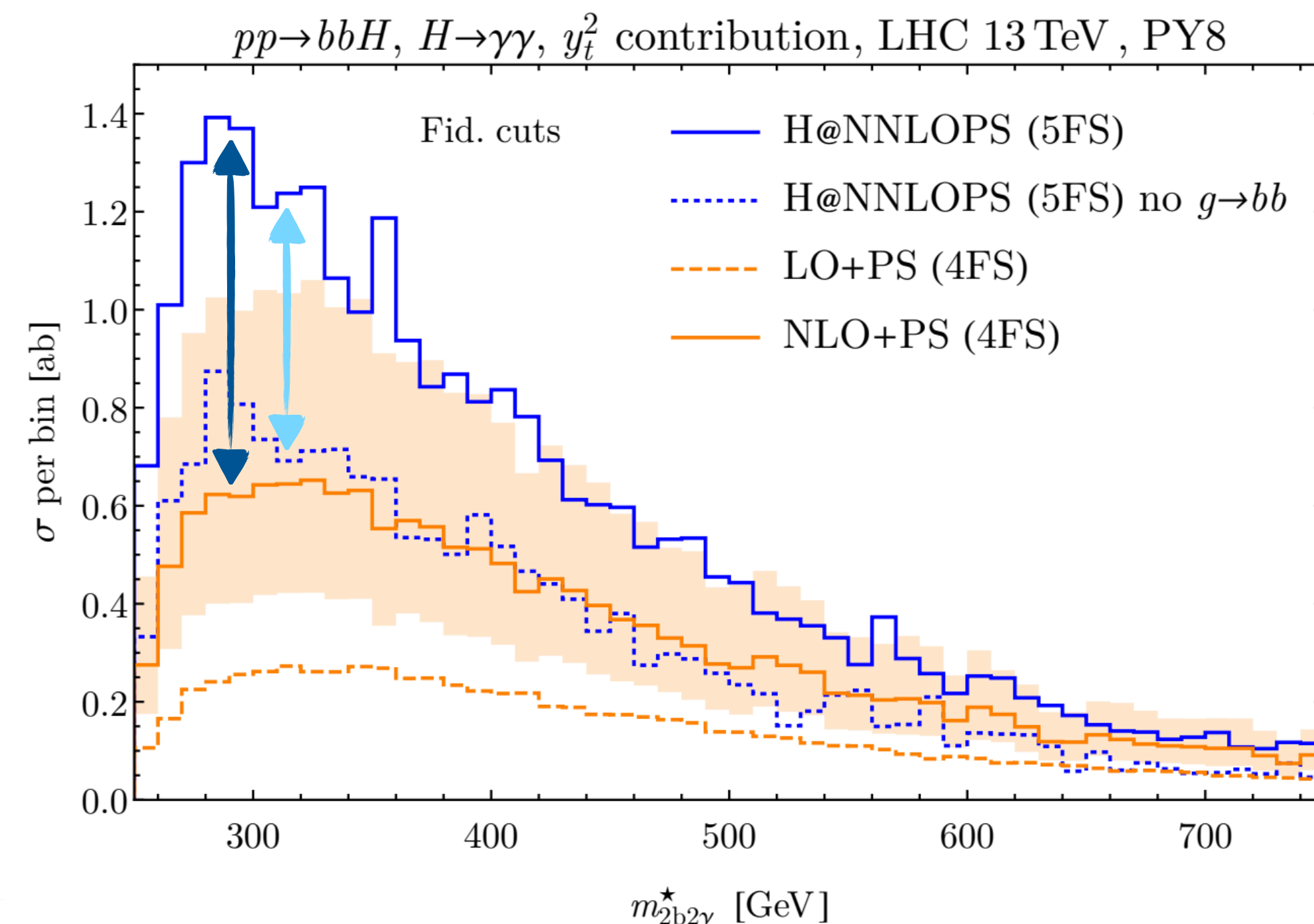


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**$g \rightarrow bb$  splittings** populate region with **no light jets** (in addition to the two  $b$ -jets from the fiducial cuts).

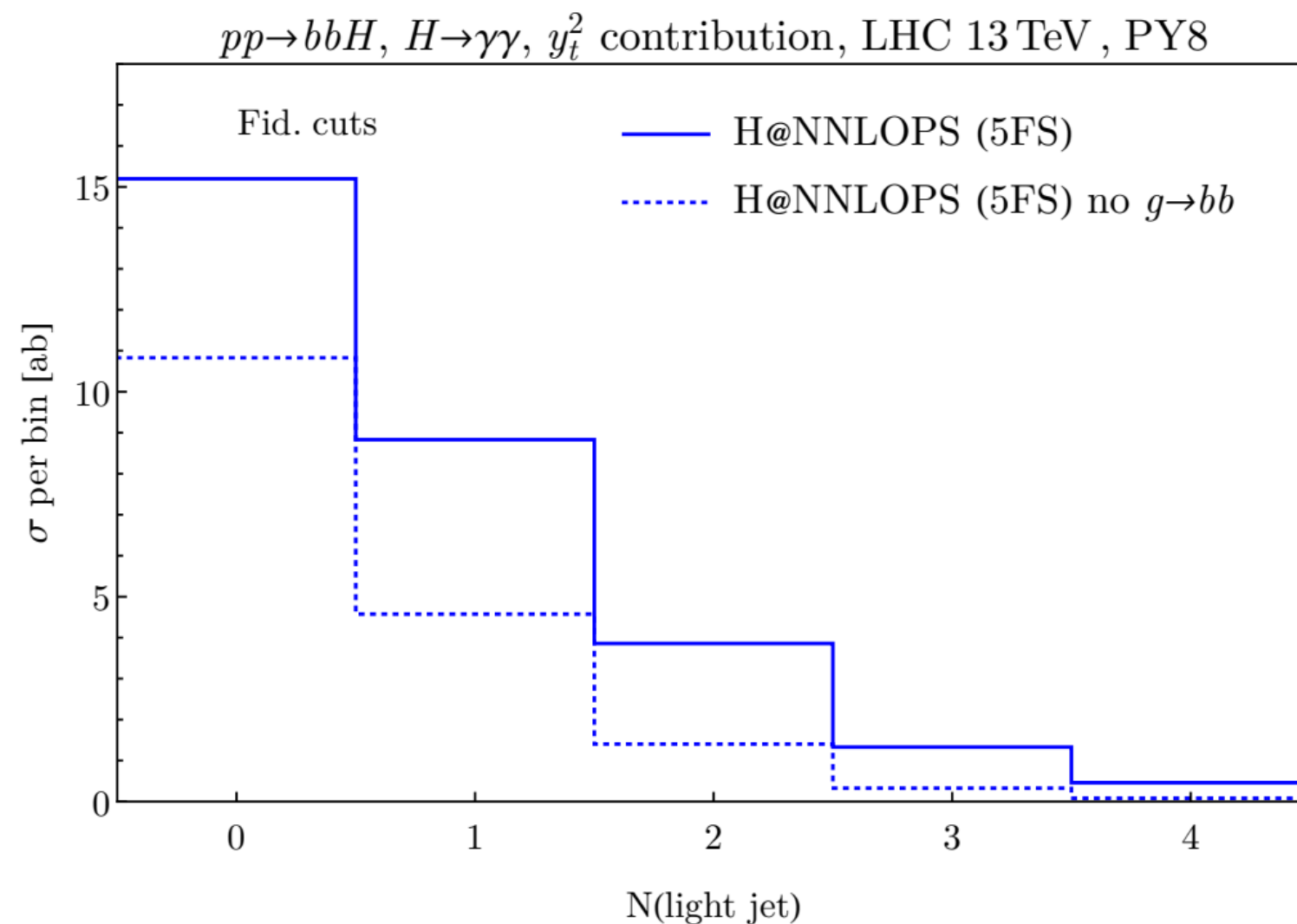
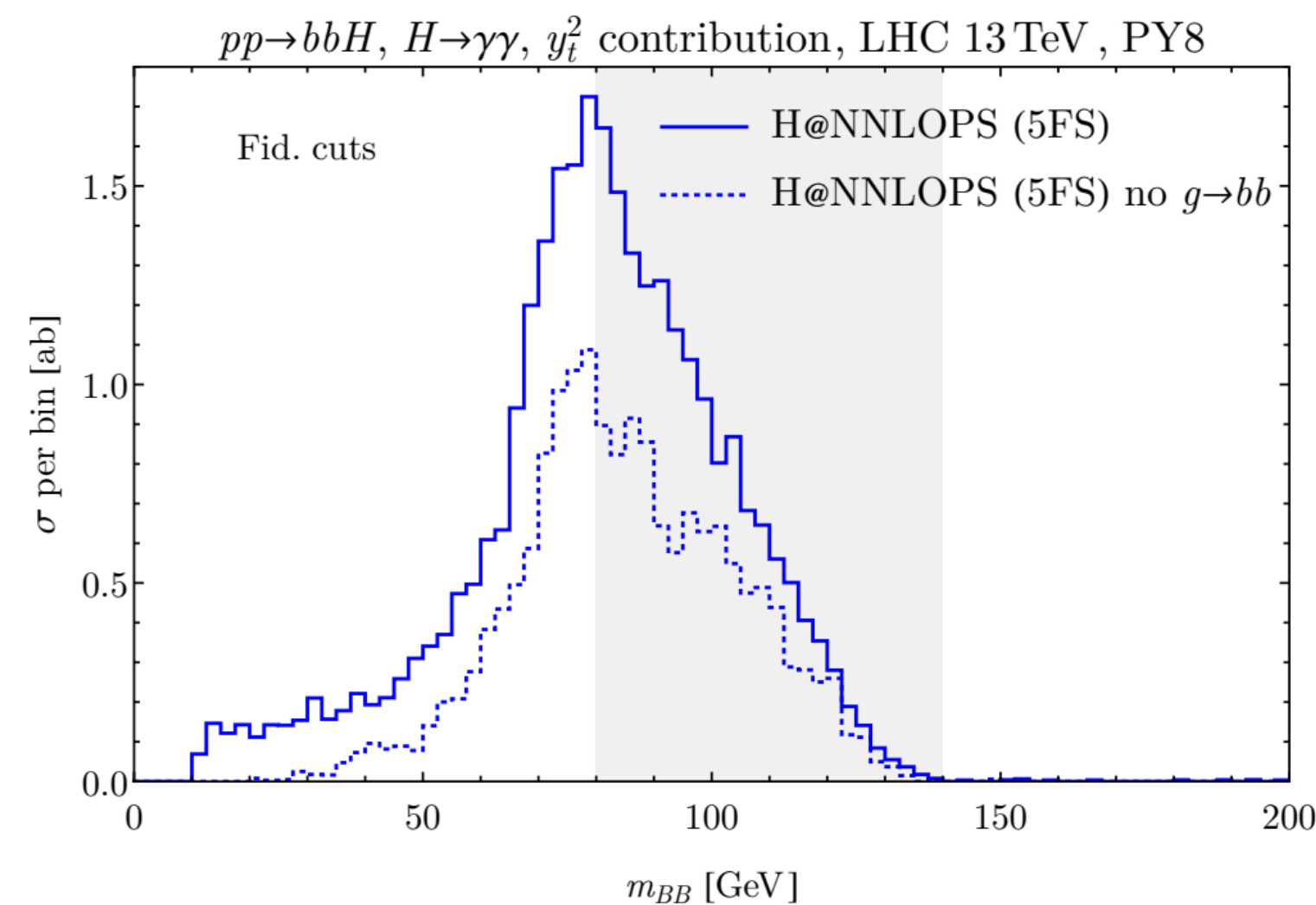
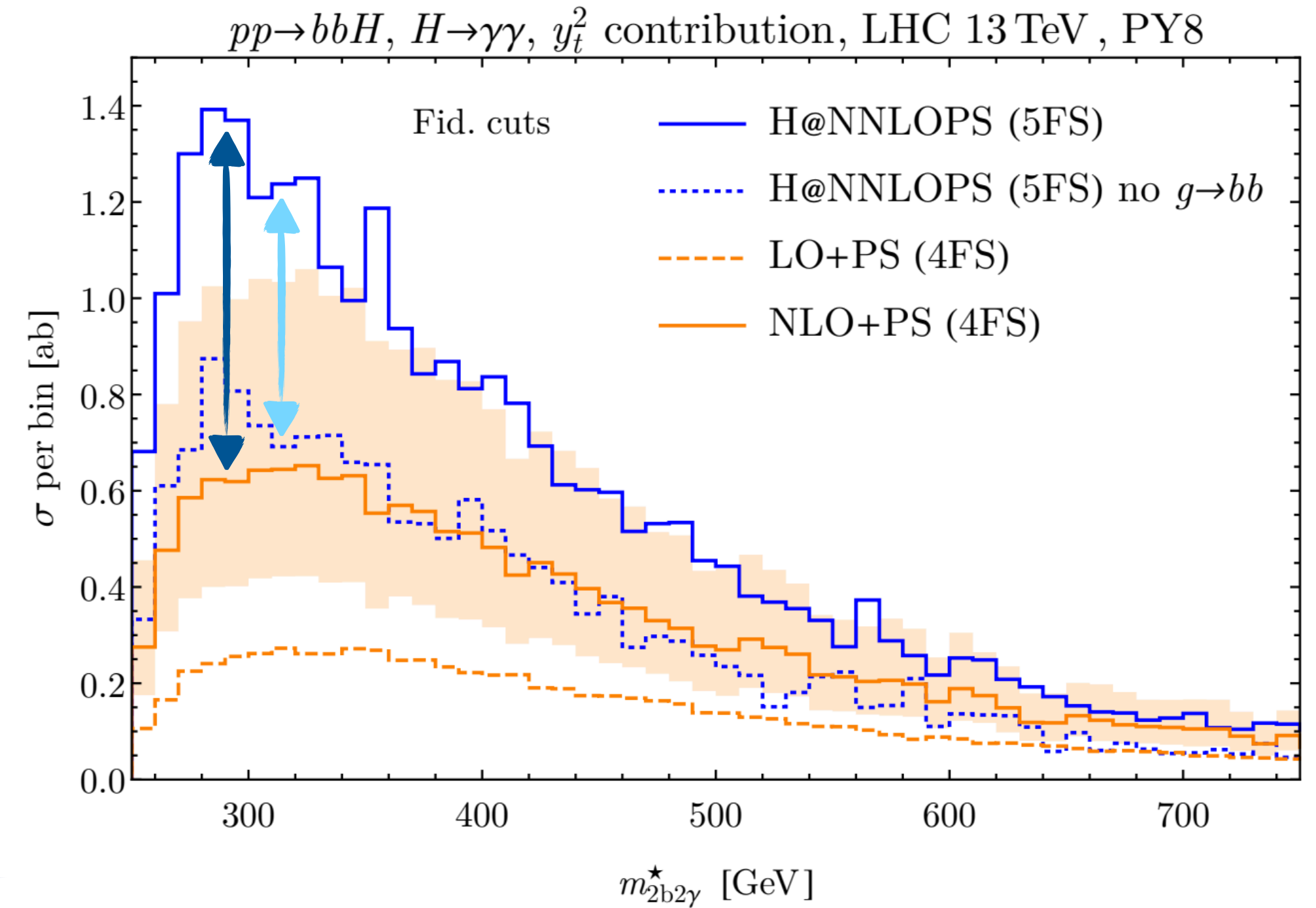
- ➡
- **$g \rightarrow bb$  splittings** are **not contributing** to  **$pp \rightarrow H + bb + jj$  events** (i.e.  $N(\text{light jets}) = 2$ ).
  - **Potential double-counting** of  $pp \rightarrow H + bb$  events (already covered by the Matrix-Element calculations in ggF NNLOPS 5FS simulation) from the **PS**.

# Comparison and combination with inclusive NNLOPS prediction

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  2. Or maybe  **$g \rightarrow bb$  splittings** are generating **hard, wide-angle bottom quarks**?



The **100% uncertainty** assigned to the ggF  $H + bb$  background by the HH searches is **probably adequate**.

# Impact of new $bbH$ modeling for $HH$ searches

**Question:** what is the **impact** of the **new  $bbH$  modeling** from the  $bbH @ NLO (4FS) + PS$  for the  $HH$  searches?

➔ We have propagated the **new  $bbH$  rates** and **uncertainties** to **two  $HH$  searches** in the  $bb\gamma\gamma$  and  $bb\tau\tau$  channels.

➔ Using full Run 2 ATLAS analyses as representative cases (= [Phys. Rev. D 106 \(2022\) 052001](#) and [JHEP 07 \(2023\) 040](#)).

1.  **$bbH @ NLO (4FS)$  rates**  $\approx 0.5 \times$   **$ggF NNLOPS (5FS)$  rates.**

➔ **Rates** of the  $bbH$  background from the two analyses  $\approx$  **halved.**

2. **100% uncertainty** on  $ggF + bb$  bkg. **replaced** with **scale uncertainties** for  **$bbH @ NLO (4FS)$**  in fid. region ( $\approx 50\%$ )

➔ **Uncertainty** on the  $bbH$  background  $\approx$  **halved!**

**Positive impact** on upper limits on  $HH$  production &  $HH$  discovery significance.

**Impact of rates and uncertainties** from  $bbH @ NLO + PS (4FS)$  on the upper limits on  $HH$  production &  $HH$  discovery significance.

	$HH \rightarrow bb\gamma\gamma$	$HH \rightarrow bb\tau\tau$
Run 2	2%	5%
HL-LHC projection	10%	20%

➔ **Larger improvement for  $HH \rightarrow bb\tau\tau$** , because the analysis is less statistically limited and because of the larger contribution of the  $bbH$  background.

## Subtlety:

- The  **$ggF NNLOPS (5FS)$**  sample is also used to estimate the **bkg.** from single **Higgs + jets**, where light or c-jets are **mistagged as b-jets.**
- The new  $bbH @ NLO$  does not cover this!
- For this exercise, we **only rescaled** the **true b-jet contribution** ( $\approx 80\%$  of the  $ggF NNLOPS$  estimation).

# Summary

- **Single Higgs boson production + bb ( $bbH$ )** is an **irreducible background** for the most sensitive  **$HH$  searches** (= all involving at least one  **$H \rightarrow bb$  decay**).
  - ➔ **Current predictions** adopted by ATLAS rely on the inclusive **ggF NNLOPS (5FS)** sample, and assign a **100% uncertainty**.
- We studied the  **$bbH$  process @ NLO (within the 4FS) + matching to PS** in a fiducial region **targeting the  $HH$  phase space**.
  - ➔ - Simulated both the  $y_b^2$  and the  $y_t^2$  contributions using MadGraph5\_aMC@NLO.
  - **Large NLO corrections** (especially for the  $y_t^2$  case).
  - Still **sizable scale uncertainties** @ NLO, especially for the  $y_t^2$  contribution (= +58%<sub>-34%</sub>).
- We **compared** the **new  $bbH$  @ NLO simulation (4FS)** with the **current ggF NNLOPS (5FS)** sample.
  - ➔ The **rates** from the **ggF NNLOPS (5FS)** prediction appear to be **largely influenced** by the  **$g \rightarrow bb$  splittings in the PS** (= probably acting **outside** of its **validity range!**).
- We estimated the **impact** of the **new  $bbH$  @ NLO (4FS) predictions** to  **$HH$  searches** (using  $HH \rightarrow bb\gamma\gamma$  and  $HH \rightarrow bb\tau\tau$  analyses as representative cases).
  - ➔ - Propagating **lower rates** and **smaller uncertainties**.
  - **2% - 20% improvement** on upper limits on  $HH$  cross-section /  $HH$  discovery significance, depending on the  $HH$  channel and the luminosity.

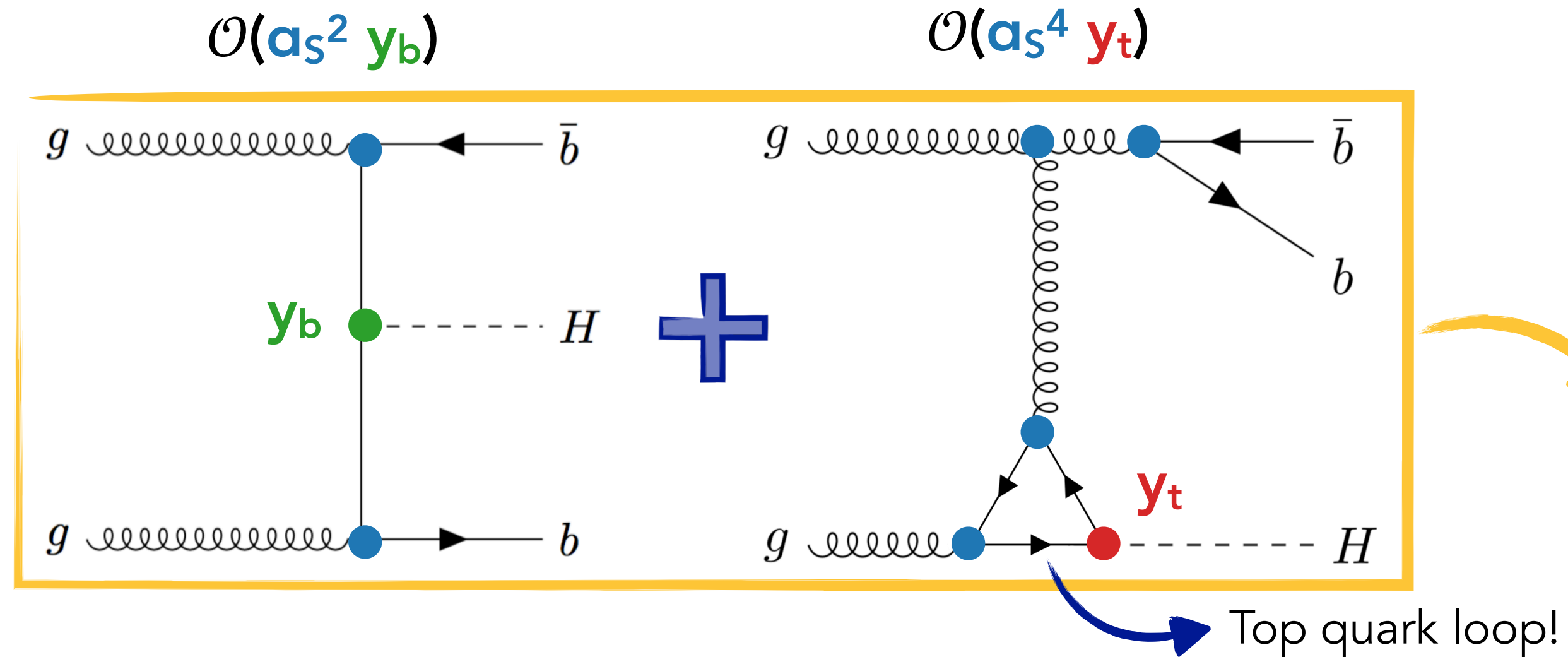
**Thank you for your attention!**

**Backup**

# The $bbH$ background for $HH$ searches

The **most sensitive**  $HH$  channels all involve the  $H \rightarrow bb$  decay.

➔ **Single Higgs boson production +  $bb$  (=  $bbH$ ) is an irreducible bkg. for all these searches!**



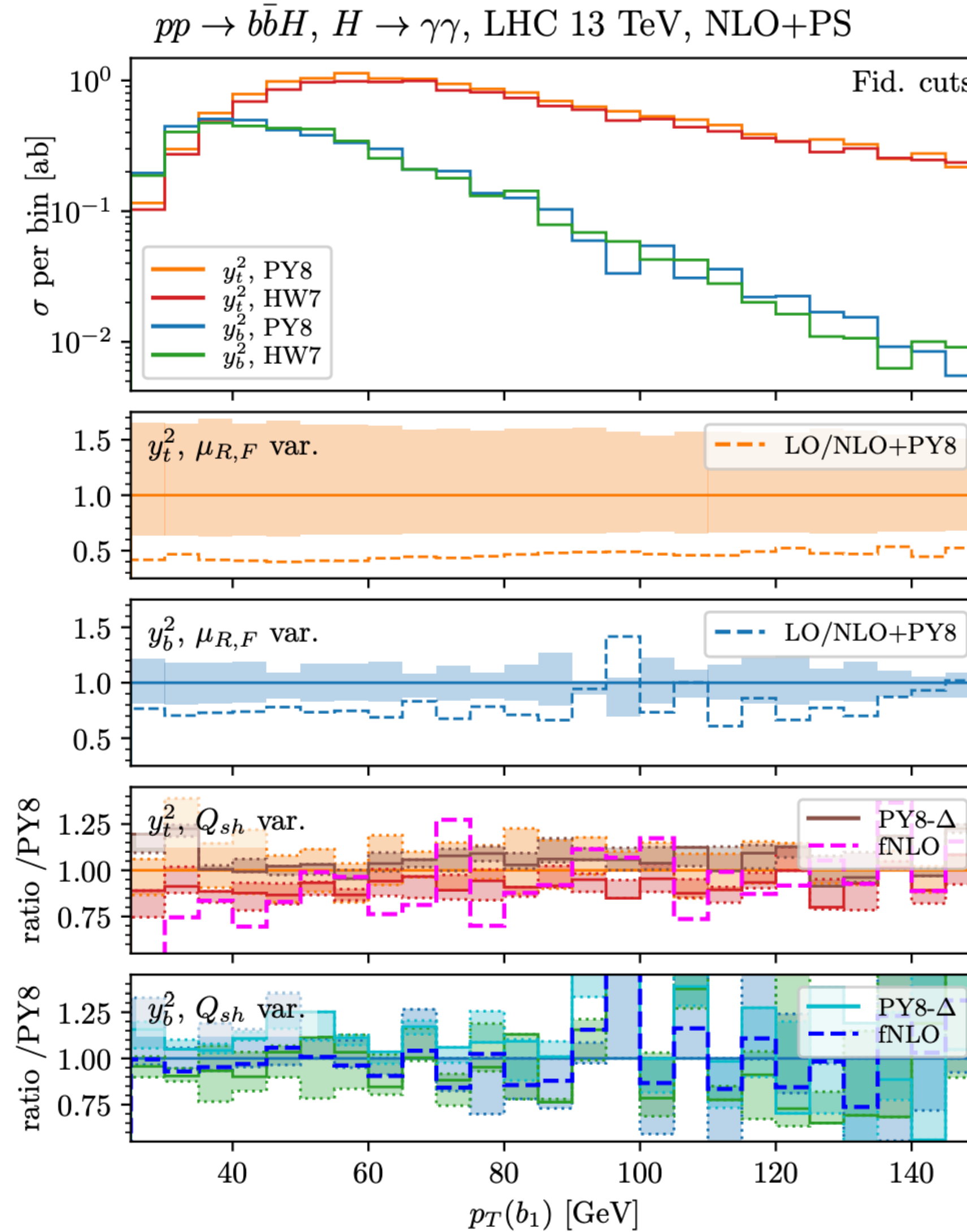
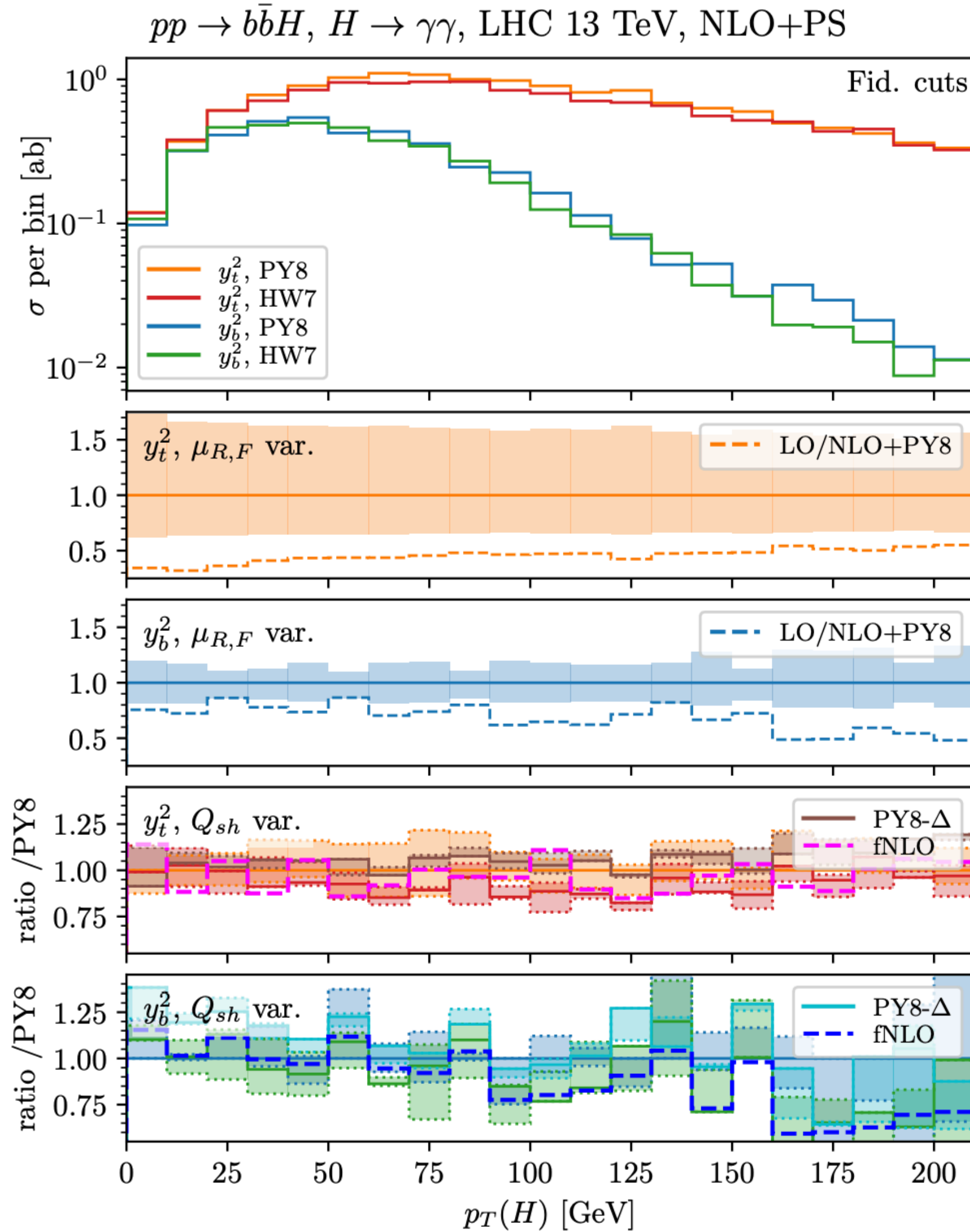
## ATLAS Run 2 $HH \rightarrow bby\gamma$ search ([JHEP 01 \(2024\) 066](#))

	High Mass 1	High Mass 2	High Mass 3	Low Mass 1	Low Mass 2	Low Mass 3	Low Mass 4
SM $HH(\kappa_\lambda = 1)$ signal	$0.26^{+0.03}_{-0.04}$	$0.194^{+0.021}_{-0.032}$	$0.84^{+0.10}_{-0.14}$	$0.048^{+0.007}_{-0.008}$	$0.038^{+0.004}_{-0.006}$	$0.039^{+0.004}_{-0.006}$	$0.032^{+0.004}_{-0.004}$
ggF	$0.25^{+0.03}_{-0.04}$	$0.188^{+0.021}_{-0.032}$	$0.81^{+0.10}_{-0.14}$	$0.046^{+0.007}_{-0.008}$	$0.036^{+0.004}_{-0.006}$	$0.037^{+0.004}_{-0.006}$	$0.025^{+0.004}_{-0.004}$
VBF [ $10^{-3}$ ]	$7.9^{+0.6}_{-0.5}$	$5.3^{+0.5}_{-0.4}$	$29^{+4}_{-3}$	$1.98^{+0.28}_{-0.24}$	$1.71^{+0.16}_{-0.14}$	$1.96^{+0.21}_{-0.19}$	$7.4^{+0.6}_{-0.5}$
Alternative $HH(\kappa_\lambda = 10)$ signal	$2.5^{+0.4}_{-0.3}$	$1.81^{+0.25}_{-0.20}$	$6.2^{+0.8}_{-0.6}$	$5.0^{+1.2}_{-0.9}$	$3.8^{+0.7}_{-0.5}$	$3.7^{+0.7}_{-0.6}$	$3.6^{+0.4}_{-0.4}$
ggF	$2.3^{+0.4}_{-0.3}$	$1.64^{+0.25}_{-0.19}$	$4.9^{+0.8}_{-0.6}$	$4.7^{+1.0}_{-0.8}$	$3.6^{+0.7}_{-0.6}$	$3.3^{+0.7}_{-0.5}$	$2.04^{+0.34}_{-0.27}$
VBF	$0.231^{+0.019}_{-0.017}$	$0.170^{+0.019}_{-0.017}$	$1.29^{+0.15}_{-0.14}$	$0.28^{+0.20}_{-0.11}$	$0.23^{+0.23}_{-0.12}$	$0.36^{+0.10}_{-0.08}$	$1.57^{+0.17}_{-0.16}$
Alternative VBF $HH(\kappa_{2V} = 3)$ signal	$0.23^{+0.04}_{-0.04}$	$0.20^{+0.05}_{-0.04}$	$3.8^{+0.7}_{-0.6}$	$0.03^{+0.04}_{-0.02}$	$0.03^{+0.06}_{-0.02}$	$0.048^{+0.023}_{-0.015}$	$0.17^{+0.04}_{-0.03}$
Single Higgs boson background	$1.5^{+0.5}_{-0.3}$	$0.48^{+0.21}_{-0.10}$	$0.57^{+0.25}_{-0.14}$	$1.72^{+0.31}_{-0.19}$	$0.53^{+0.08}_{-0.06}$	$0.29^{+0.14}_{-0.07}$	$0.16^{+0.06}_{-0.03}$
ggF	$0.5^{+0.5}_{-0.2}$	$0.14^{+0.21}_{-0.09}$	$0.25^{+0.25}_{-0.12}$	$0.29^{+0.31}_{-0.15}$	$0.08^{+0.08}_{-0.04}$	$0.07^{+0.13}_{-0.06}$	$0.04^{+0.06}_{-0.03}$
$t\bar{t}H$	$0.302^{+0.034}_{-0.032}$	$0.069^{+0.009}_{-0.008}$	$0.063^{+0.008}_{-0.007}$	$0.77^{+0.09}_{-0.08}$	$0.214^{+0.029}_{-0.026}$	$0.100^{+0.012}_{-0.012}$	$0.048^{+0.005}_{-0.005}$
$ZH$	$0.61^{+0.06}_{-0.05}$	$0.174^{+0.020}_{-0.016}$	$0.188^{+0.035}_{-0.029}$	$0.49^{+0.05}_{-0.04}$	$0.149^{+0.028}_{-0.025}$	$0.069^{+0.033}_{-0.023}$	$0.028^{+0.010}_{-0.007}$
Rest	$0.17^{+0.08}_{-0.04}$	$0.089^{+0.030}_{-0.016}$	$0.07^{+0.04}_{-0.02}$	$0.181^{+0.030}_{-0.019}$	$0.089^{+0.016}_{-0.009}$	$0.046^{+0.007}_{-0.004}$	$0.039^{+0.008}_{-0.004}$
Continuum background	$11.3^{+1.5}_{-1.6}$	$3.2^{+0.8}_{-0.8}$	$2.8^{+0.8}_{-0.8}$	$37.2^{+2.9}_{-2.9}$	$10.8^{+1.5}_{-1.5}$	$4.4^{+0.9}_{-1.0}$	$1.1^{+0.5}_{-0.5}$
Total background	$12.8^{+1.6}_{-1.6}$	$3.7^{+0.9}_{-0.8}$	$3.4^{+0.8}_{-0.8}$	$38.9^{+2.9}_{-2.9}$	$11.3^{+1.5}_{-1.5}$	$4.7^{+0.9}_{-1.0}$	$1.3^{+0.5}_{-0.5}$
Data	12	4	1	29	8	5	4

➔ The  **$bbH$  background** is not negligible w.r.t. the **SM  $HH$  signal!**

➔ Example of ATLAS Run 2  $HH \rightarrow bby\gamma$  search:  **$bbH$  background** comparable with **SM  $HH$  signal** in most sensitive analysis categories (= High Mass 3 and Low Mass 4).

# Differential distributions

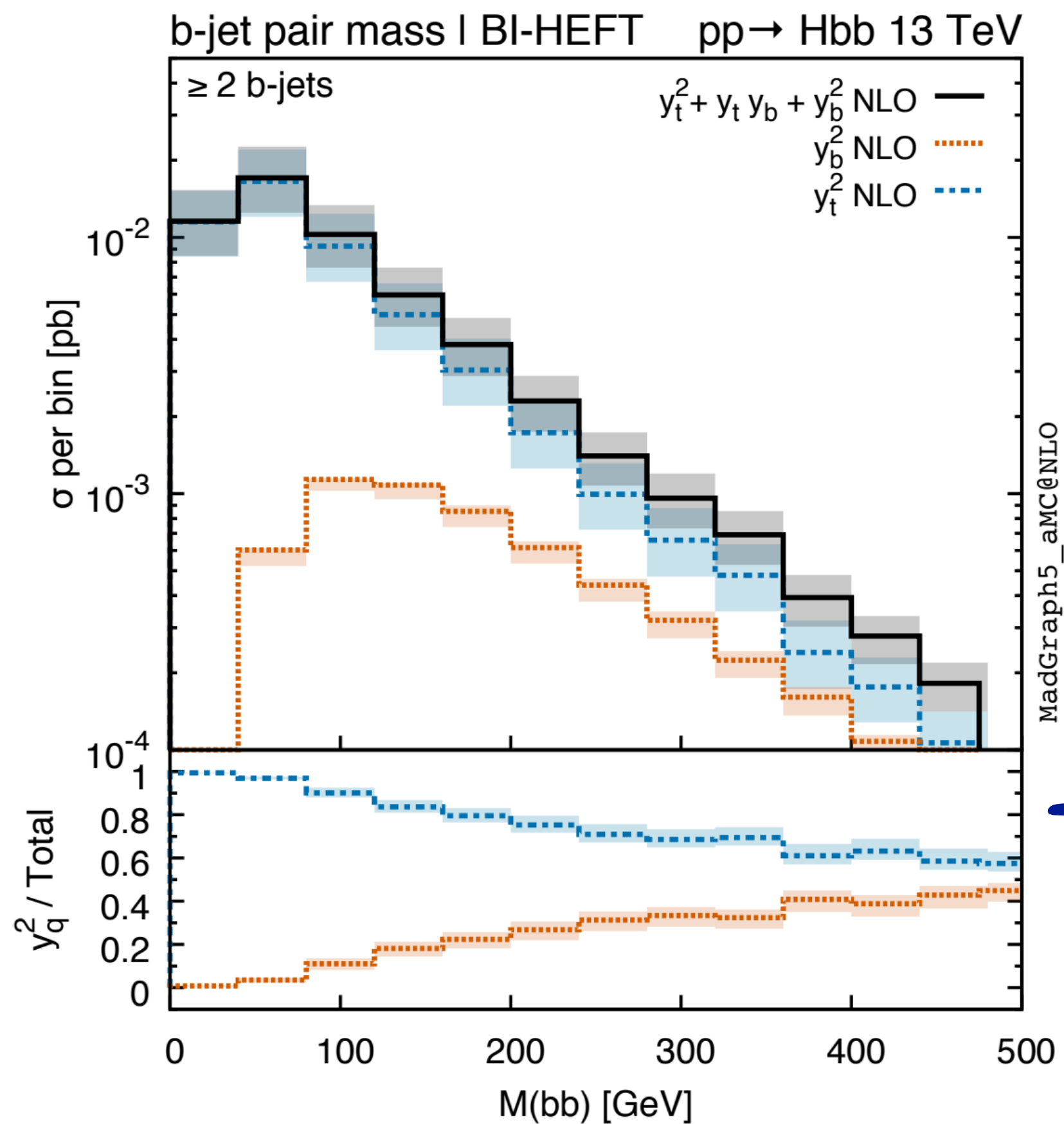


The shape difference between the  $y_b^2$  contribution and the  $y_t^2$  contribution for  $p_T(H)$  and  $p_T(b_1)$  confirms that the  $y_t^2$  contribution favors harder spectra for the Higgs boson and bottom quarks.



The (fiducial) rates for the  $bbH$  process was computed @ NLO in the **4-flavor scheme (4FS)**, using MadGraph5\_aMC@NLO.

- ➔ - Including the  $\propto y_b^2$  and the  $\propto y_t^2$  **contributions** and the **interference** ( $\propto y_b y_t$ ). ➔ = Massive bottom quarks.
- **Heavy top approximation (HTL)** adopted for the **NLO corrections** to the  $y_t^2$  component. ➔ Verified to be reliable from a comparison @ LO.



➔ The relative contribution of the **interference** ( $\propto y_b y_t$ ) corresponds to  $1 - y_b^2 - y_t^2$ .

# The $bbH$ process @ NLO + PS (4FS) in the phase space of $HH \rightarrow b\bar{b}\gamma\gamma$

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We present a simulation of  $bbH$  process @ NLO + matching to parton shower (PS) + dedicated analysis targeting the  $HH$  phase space!  $\longrightarrow$  Using the  $HH \rightarrow b\bar{b}\gamma\gamma$  search as a representative case.

Cut	Contr.	Run	LO	NLO	$\delta\mu_{R,F}$	$\delta Q_{sh}$	NNLOPS ( $y_t^2$ LO)	HH signal	
No cut	$y_b^2$	PY8	561	849	+50%	+8%	4867	82.1	
		PY8- $\Delta$		848	-20%	+0%			
		HW7	561	851	+0%	+0%			
	$y_t^2$	PY8	655	1565	+61%	+0%			2140
		PY8- $\Delta$		1595	-35%	+0%			
		HW7	655	1578	+0%	+0%			
	sum	PY8	1217	2414	+46%	+0%			
		PY8- $\Delta$		2443	-29%	+0%			
		HW7	1216	2429	+0%	+0%			
Fid. cuts	$y_b^2$	PY8	3.15	4.22	+35%	+10%	29.9	22.7	
		PY8- $\Delta$		4.75	-15%	-4%			
		HW7	2.59	4.08	+8%	-2%			
	$y_t^2$	PY8	8.24	18.1	+58%	+10%			17.2
		PY8- $\Delta$		19.2	-34%	-7%			
		HW7	6.83	16.6	+3%	+4%			
	sum	PY8	11.4	22.3	+50%	+10%			
		PY8- $\Delta$		23.9	-30%	-6%			
		HW7	9.42	20.7	+2%	+4%			
Fid. cuts + $m_{2b2\gamma}^* < 500$ GeV	$y_b^2$	PY8	3.11	4.15	+33%	+11%	22.3	15.7	
		PY8- $\Delta$		4.69	-15%	-4%			
		HW7	2.56	4.02	+8%	-2%			
	$y_t^2$	PY8	5.33	12.3	+60%	+12%			13.3
		PY8- $\Delta$		12.8	-34%	-8%			
		HW7	4.31	11.3	+8%	+5%			
	sum	PY8	8.44	16.5	+49%	+12%			
		PY8- $\Delta$		17.5	-29%	-7%			
		HW7	6.86	15.3	+1%	+6%			
Fid. cuts + $m_{2b2\gamma}^* < 350$ GeV	$y_b^2$	PY8	2.71	3.65	+33%	+9%	11.5	2.84	
		PY8- $\Delta$		4.11	-16%	-4%			
		HW7	2.22	3.54	+8%	-2%			
	$y_t^2$	PY8	2.32	5.78	+61%	+13%			6.82
		PY8- $\Delta$		6.05	-34%	-9%			
		HW7	1.88	5.43	+9%	+5%			
	sum	PY8	5.03	9.43	+44%	+12%			
		PY8- $\Delta$		10.2	-27%	-7%			
		HW7	4.10	8.97	+0%	+6%			

## Fiducial cross-sections (before cuts and in $m_{2b2\gamma}^*$ categories)

- The  $bbH$  rates were evaluated also **before the fiducial cuts**, and in **three** categories, based on cuts on the  $m_{2b2\gamma}^*$  variable (on top of the fiducial selection).

$\longrightarrow$   $m_{2b2\gamma}^* < \infty$ ,  $m_{2b2\gamma}^* < 500$  GeV, and  $m_{2b2\gamma}^* < 350$  GeV.

- The  $bbH$  cross-section changes substantially depending on the cuts!

$\longrightarrow$

- After applying the **fiducial cuts**, the  $bbH$  cross-section drops of a factor  $\sim 100!$
- In the fiducial region, the  $bbH$  background rate is **comparable** with the **SM HH signal**, and becomes **dominant** in the  $m_{2b2\gamma}^* < 350$  GeV category.

- The relative contributions of the  $y_b^2$  and  $y_t^2$  components changes with the cuts.

$\longrightarrow$  The  $y_b^2$  contribution is **subleading** w.r.t  $y_t^2$  in all categories, **except** in the  $m_{2b2\gamma}^* < 350$  GeV category, where the **two contributions are similar**.

# The $bbH$ process: the state of the art

## References ([TWiki](#) and ongoing work)

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