

#### Nº INTERNATIONAL CONFERENCE ZN **HIGH ENERGY PHYSICS Julu 2024**

# In search of new Physics with Lepton Flavor Violation in  $\Upsilon(3S) \rightarrow e^{\pm} \mu^{\mp}$  at BABAR

Nafisa Tasneem On behalf of the BaBar Collaboration July 18 - 24, 2024.

## Outline of the Talk

- Motivation: Charged Lepton Flavor Violation
- Theoretical Expectations and Experimental Limits
- Asymmetric PEP-II Collider and BaBar Detector
- Data, MC and Data Driven Background
- Analysis Strategy
	- ➢ Signal and Background Characteristic
	- ➢ A Selection Criterion
	- ➢ Validation of Using the Data Driven Background
- Results and Indication for New Physics
- Conclusion

### Motivation

- Lepton Flavor Violation (LFV) in the neutral lepton sector has already indicated by the observation of neutrino oscillations by the Super-Kamiokande Observatory and the Canadian Sudbury Neutrino Observatories (SNO).
- Such an oscillation mechanism cannot induce observable LFV in the charged lepton sector.
- Charged Lepton flavor violating (CLFV) processes are strongly suppressed in the standard model by powers of (small) neutrino masses, e.g.  $\left(\frac{\Delta m_V^2}{M^2}\right)$  $M_W^2$ 2  $≤ 10^{-48}$ [1]
- Observation of CLFV is, therefore, a clear sign of new physics (NP) beyond the SM.

[1] Phys. Rev. Lett. 104, 151802

#### **Motivation**

Table: CLEO and BABAR results on different decay modes of  $\Upsilon$ .



Phys. Rev. Lett. 128 (2022) 9, 091804

 $\blacktriangledown$ 

- BABAR reports the first search for electron-muon LFV in the decay of a  $\bullet$ Y (3S) resonance.
- With the results of  $Y(3S)$  we placed constraints on New Physics (NP)  $\bullet$ processes that include LFV.

### Theoretical Background

• Let us assume that a vector boson  $V_i$  (here  $V_i$  could be either a fundamental state, such as the  $Z^0$ , or a quark-antiquark bound state such as the  $\phi$ ,  $J/\psi$ , or Y) couples to  $e^{\pm}\mu^{\mp}$ . The effective coupling between the vector boson  $V_i$  and  $e^{\pm} \mu^{\mp}$  can be written as

$$
[\mathcal{B}(\mu \to 3e)]_{V-\text{exch.}} \approx \frac{\Gamma(V \to e^+e^-)\Gamma(V \to e^{\pm}\mu^{\mp})}{\Gamma^2(W \to e\nu)} \left(\frac{M_W}{M_V}\right)^6.
$$



Using  $\mathcal{B}$  ( $\mu \rightarrow eee$ ) <1.0 × 10<sup>-12</sup> [2] and other data pertaining to the  $e^+e^-$  widths of the various vector mesons  $V_i$ , we find

$$
\mathcal{B}\left(\Upsilon(3S) \to e^{\pm} \mu^{\mp}\right) \leq 2.5 \times 10^{-8}
$$

- $\mu \rightarrow eee$  decay amplitude can be significantly reduced if there are kinematical suppressions. Such suppressions are possible when the effective vector boson couplings involve momentum factors. According to Nussinov [3], the size of the vector boson exchange contribution reduce by a factor of  $\frac{M_{\mu}^2}{2M_{\tau}^2}$  $2M_{\Upsilon(3S)}^2$  $\frac{\sqrt{n\mu}}{2}$  which is **3** × **10**<sup>-5</sup>.
- Thus, **new modified bound on**  $\mathcal{B}(\Upsilon(3S) \to e^{\pm} \mu^{\mp})$  **reduced by**  $\leq 1 \times 10^{-3}$

#### **ntasneem@stfx.ca**

<sup>[2]</sup> Bellgardt, et al., Nucl.Phys. B299 (1988) [3] S. Nussinov, et. al. PRD 63 (2001)

# **Charged Lepton Flavor Violation** in Upsilon Decays



## Data, MC Sample



# Signal and Background

- $\Upsilon(3S) \rightarrow e^{\pm} \mu^{\mp}$ : required two primary track signal of  $e^\pm$  and  $\mu^\mp$
- Centre of Mass (CM) momentum:  $p_{e^{\pm}}{\sim}\frac{\sqrt{s}}{2}$  $\frac{\sqrt{s}}{2}$  =  $E_B$  and  $p_{\mu^{\pm}} \sim \frac{\sqrt{s}}{2}$  $\frac{\sqrt{3}}{2} = E_B$ where  $\mathbf{E}_B$  = beam energy in CM
- Angle between the two lepton tracks must satisfy  $\theta_{12}^{CM} > 179^{\circ}$  to emerge as back-to-back.
- Energy deposit by  $\mu^{\pm}$  track on the Electromagnetic Calorimeter (EMC) > 50 MeV
- EMC acceptance  $24^{\circ} < \theta_{Lab} < 130^{\circ}$  etc.
- Optimized PID selection for  $e^{\pm} \mu^{\mp}$  track





### Selection criteria on the lepton momentum plane



The lepton momenta must satisfy the condition which is defining a circle of radius  $p_e$  $\bm{E}_{\bm{B}}$  $-1$  $\mathbf{z}$  $+$  $p_\mu$  $\bm{E}_{\bm{B}}$  $-1$  $\mathbf{2}$  $= (0.1)^2 = 0.01$  Where,  $p_{e^{\pm}, \mu^{\pm}} \sim \frac{\sqrt{s}}{2}$  $\frac{\sigma}{2} = E_B$ 

#### **ntasneem@stfx.ca**

# Systematic Study



- We reversed two major cuts to check the  $Y(4S)$ data/MC agreement.
- $e^+e^- \rightarrow \Upsilon(4S)$  of 78.31/fb pre-selected as  $e^{\pm} \mu^{\mp}$ and  $\mu^{\pm} \mu^{\mp}$
- Disagreement arises due to uncertainties in PID, Tracking, kinematics, trigger etc.
- The difference in continuum background and MC background (in the energy band 6-8 GeV) was 1.2%



Impact of each component of the selection on the signal efficiency, background and data

## **Final Invariant Mass Distribution of**  $e^{\pm}\mu^{\mp}$



# Summary and Result  $\ \Upsilon(3S) \rightarrow e^{\pm} \mu^{\mp}$



Table: Summery of systematic uncertainties. The values of efficiency, background, and number of  $Y(3S)$  decays are presented in the first column and their uncertainties in the second column.

- $N_{\rm{Candidate}}~-N_{BG}$ Calculating the branching fraction from  $\varepsilon_{sig}\times N_{\Upsilon}$  $\mathcal{B}[\Upsilon(3S) \to e^{\pm} \mu^{\mp}] = [1.0 \pm 1.4 \text{(stat)} \pm 0.8 \text{(syst)}] \times 10^{-7}$
- We set an upper limit at 90% confidence level  $(C.L.)$  on the branching fraction by using the "CLs" method.

$$
\mathcal{B}[\Upsilon(3S) \to e^{\pm} \mu^{\mp}] < 3.6 \times 10^{-7} \text{ @ } 90\% \text{ C.L.}
$$

• This result is the first reported experimental upper limit on  $\Upsilon$  (3S)  $\rightarrow e^{\pm}\mu^{\mp}$ .

# Implication For New Physics (NP)



Using the relationship,

$$
\frac{\left(\frac{1}{A_{NP}}\right)}{\frac{\pi\alpha_{QED}Q_b}{M_{\Upsilon(3S)}}} = \frac{\mathcal{B}(\Upsilon(3S))}{\mathcal{B}(\Upsilon(3S))}
$$

we can set an upper limit on NP  $\frac{24N}{\epsilon_2} \geq 80 \text{ TeV}$  @90% C.L.

 $\left(\frac{g_{NP}^2}{2}\right)^2$ 

 $g^2$ NP  $\Lambda_{NP}$ = Effective coupling of the new physics Energy scale of the NP, given by the mass of the NP propagator. where,

- 
$$
Q_b = -\frac{1}{3}
$$
 is the b-quark charge

- $\alpha_{3S}$  is the fine structure constant at the  $M_{\gamma(3S)}$  energy scale.
- Using the world average  $\mathcal{B}(Y(3S) \to \mu^+ \mu^-) = 2.18 \pm 0.21$  [4]

[4] P. A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. 2020, 083C01 (2020).

#### **ntasneem@stfx.ca**

### Conclusion

• Upper limit at 90% confidence level (C.L.) on the branching fraction

#### $\mathcal{B}[\Upsilon(3S) \to e^{\pm} \mu^{\mp}] < 3.6 \times 10^{-7}$  @ 90% C.L.

• **A limit on NP is**  $\frac{P_1P_2}{q_{NP}^2} \ge 80 \,\text{TeV}$  (*Q*) 90% CL [Phys. Rev. Lett. 128 (2022) 9, 091804]

- This is the first reported experimental upper limits on the branching fraction of  $\varUpsilon(3S)\to e^\pm\mu^\mp$  .
- The measurement we report here is several orders of magnitude more sensitive than the indirect limit  $\mathcal{B}(\Upsilon(3S) \to e^\pm \mu^\mp) \leq 1 \times 10^{-3}$  .

#### Thanks

#### **Back up: Theoretical Upper limit (Indirect)**

**Nussinov, Peccei, Zhang [1]**

- Assume coupling of Y to eµ looks like:  $L_{eff} = gV_{e\mu}\bar{\mu}\gamma_{\alpha}eV^{\alpha}$
- Through Fig 1. this coupling contributes to  $A(\mu \rightarrow 3e)$

$$
A(\mu \to 3e) = (\bar{u}_{\mu}(p)\gamma^{\alpha}u_{e}(k_{3}))(\bar{v}_{e}(k_{1})\gamma_{\alpha}u_{e}(k_{2}))\frac{g_{V_{e\mu}}g_{V_{ee}}}{M_{V}^{2} - S} \dots (1)
$$

$$
\frac{\left[\Gamma(\mu \to 3e)\right]_{V-exch}}{\left[\Gamma(\mu \to e\nu\bar{\nu})\right]} \approx \frac{g^2 V_{e\mu} g^2 V_{ee}}{M_V^4} / \frac{g_W^4}{M_W^4} \qquad \qquad \text{---(2)}
$$

 $\sim$ 



(Left) A vector exchange diagram contributing to  $\mu \rightarrow 3e$ (Right) Ordinary muon decay,  $\mu \to e\nu\bar{\nu}$ , which proceeds via W exchange.

• 
$$
\text{BF}(\mu \to \text{eee}) \leq 1.0 \times 10^{-12}
$$

- BF $(\mu \to e\nu\bar{\nu}) \simeq 100\%$
- BF(W  $\rightarrow$  e<sup>+</sup>v)  $\simeq$  (10.71  $\pm$  0.09) %

• 
$$
\text{BF}(\Upsilon(3S) \to l^+l^-) \simeq (2.18 \pm 0.21)
$$
 %

• 
$$
\Gamma(\Upsilon(3S) = (20.32 \pm 1.85) \; keV
$$

•  $\Gamma(W) = (2.046 \pm 0.049) \; GeV$ 



Since 
$$
[\Gamma(V \to e^+e^-)] \sim g^2 V_{ee} M_V
$$
 and  $[\Gamma(V \to e^{\pm} \mu^{\mp})] \sim g^2 V_{e\mu} M_V$ , while  $[\Gamma(W \to e\nu)] \sim g_W^2 M_W$ 

$$
[BR(\mu \to 3e)]_{V-exch} \approx \frac{[\Gamma(V \to e^+e^-)][\Gamma(V \to e^{\pm} \mu^{\mp})]}{[\Gamma^2(W \to e\nu)]} (\frac{M_W}{M_V})^6 \quad \text{---} (3)
$$

$$
BR(\Upsilon \to e\mu) = BR(\mu \to eee) \frac{\Gamma(W \to e\nu)^2}{\Gamma(\Upsilon)\Gamma \to ee} (\frac{M_{\Upsilon}}{M_W})^6
$$
---(4)  
BR(\Upsilon(3S) \to e^{\pm} \mu^{\mp} \le 2.5 \times 10^{-8}.

S.Nussinov, et. al. estimate that the contribution of the virtual  $\Upsilon(3S) \rightarrow e^{\pm} \mu^{\mp}$  to the  $\mu \rightarrow eee$  rate would be reduced by approximately  $M_{\mu}^2$  (2  $M_{\gamma}^2$ ) leading to a re-calculated indirect bound:  $BF(Y (3S) \rightarrow e^{\pm} \mu^{\mp}) < 1 \times 10^{-3}$  $BF(T(3S) \rightarrow e^{\pm} \mu^{\prime}) < 1 \times 10^{-5}$ <br>
[1] Nussinov, et. al. PRD 63, 016003 (2001) 15

#### **Back up: Analysis Scheme**

- **Blind Analysis:** To eliminate experimenter's bias.
- **Pre-Selection:** Needs a special background filter to collect  $e^{\pm} \mu^{\mp}$  events efficiently.
- **Final Selection by the analyst:** Applied on the pre-selected events
- **PID Selection:** Multivariate Technique applied, tested 16 different PID selectors.
- **Optimized Electron and Muon selectors:**  $\frac{\varepsilon_{e\mu}}{\varepsilon_{e\mu}}$  $(1+N_{BG})$ where

 $\varepsilon_{e\mu}$  is the final efficiency as determined by signal MC and

 $N_{BG}$  is the number of expected background events



#### Back up: Impact of each component of the selection on the signal efficiency, background and data



- The first row provides information on the pre-selection of the events as  $\Upsilon$  (3*S*)  $\rightarrow e^{\pm} \mu^{\mp}$ .
- The last row provides information after applying all selection criteria.
- Rows 2-7 provides information when all requirements are applied except the criterion associated with the particular row.
- Third column represents the background events from  $e^+e^- \to \gamma(3S)$  defining as generic MC Background.
- Fourth column represents the background events from  $e^+e^- \to \Upsilon(4S)$  defining as data-driven Continuum **Background.**
- Event numbers in the third and forth columns are luminosity-normalized.
- The last column gives the number of events in the  $27 \text{ fb}^{-1}$  data sample after unblinding.
- We have added an additional uncertainty of  $\pm$  0.9 any potential peaking background from the  $\gamma$ (3S) decays which estimate the background events  $12.2 \pm 2.3$