

Measurements of single and double D^0 meson production in pp collisions at $\sqrt{s} = 13.6$ TeV with ALICE

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1. Introduction

We measured the D^0 non-prompt fraction in pp collisions using Run 3 data from ALICE. This measurement can be used to **test pQCD** calculations in the beauty sector. It was already performed in Run 2 [1], but in Run 3 we can:

- **Improve the precision** with respect to Run 2, thanks to much larger data samples.
- Get **more granular results**, even down to $p_T = 0$.
- Provide tighter constraints for **distinguishing different hadronisation implementations in models** (EPOS [2], PYTHIA [3] with different tunes).

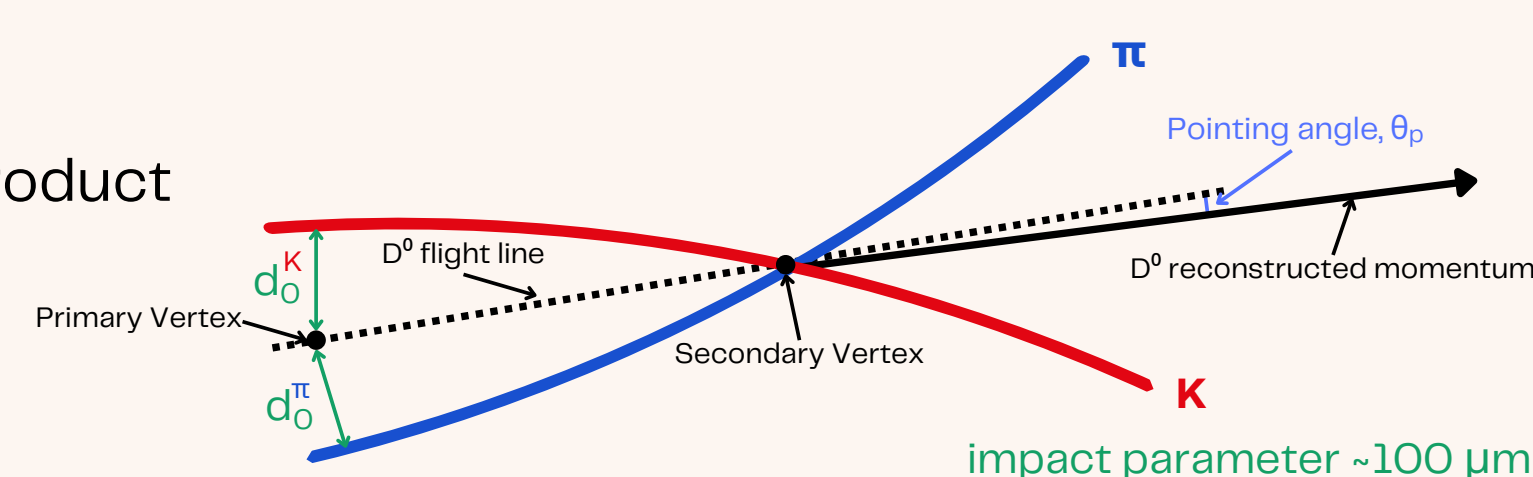
We can also measure double D^0 production in Run 3 to investigate and **distinguish single and double parton scattering** better.

3. Methodology

To select D^0 mesons, we trained a set of Boosted Decision Trees (BDTs) using candidates with $0 < p_T < 24$ GeV/c with **three different classes**: prompt D^0 , non-prompt D^0 and combinatorial background.

To perform the model training, we used **variables that exploit the differences in the topology and the particle identification (PID)** between our classes, in particular:

- Daughter impact parameter product
- Cosine of pointing angle
- Decay length
- TPC + TOF PID information



The distributions of the variables for prompt and non-prompt D^0 mesons were taken from an ideal MC, and those for the background were obtained from sidebands of the data invariant-mass distribution. To separate between prompt and non-prompt candidates accepted by the model, we used a data-driven approach called **cut-variation**:

1. Define a **set of cuts on the non-prompt BDT score** to get different prompt and non-prompt D^0 contributions. For each cut, we can define an equation of the type:

$$(\text{Acc} \times \epsilon)_i^{\text{prompt}} \cdot N_{\text{prompt}} + (\text{Acc} \times \epsilon)_i^{\text{non-prompt}} \cdot N_{\text{non-prompt}} = Y_i$$

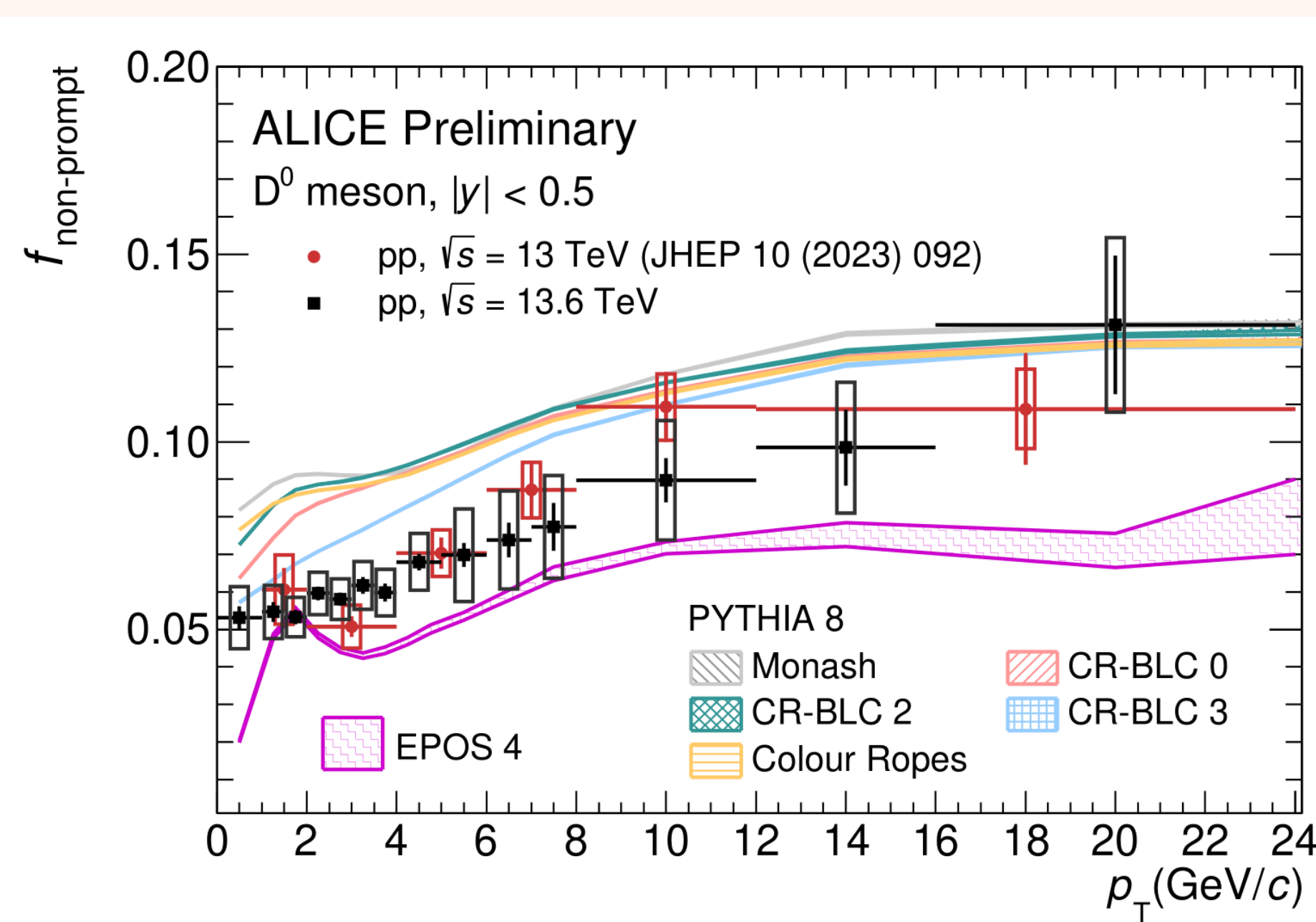
where Y is the raw-yield, and $(\text{Acc} \times \epsilon)$ is the acceptance multiplied by the efficiency.

2. From the full set of cuts, we obtain a system of equations that can be solved iteratively by **minimising the χ^2** .

3. Once we have the number of prompt (N_p) and non-prompt (N_{np}) candidates, **extract the fraction** per cut as

$$f_{\text{non-prompt}} = \frac{(\text{Acc} \times \epsilon)^{\text{non-prompt}} \cdot N_{\text{non-prompt}}}{(\text{Acc} \times \epsilon)^{\text{non-prompt}} \cdot N_{\text{non-prompt}} + (\text{Acc} \times \epsilon)^{\text{prompt}} \cdot N_{\text{prompt}}}$$

5. Results and outlook



The D^0 non-prompt fraction obtained with Run 3 data (in **black**), taken at $\sqrt{s} = 13.6$ TeV, is compared with the results from Run 2 (in **red**), taken at $\sqrt{s} = 13$ TeV, and several theoretical models:

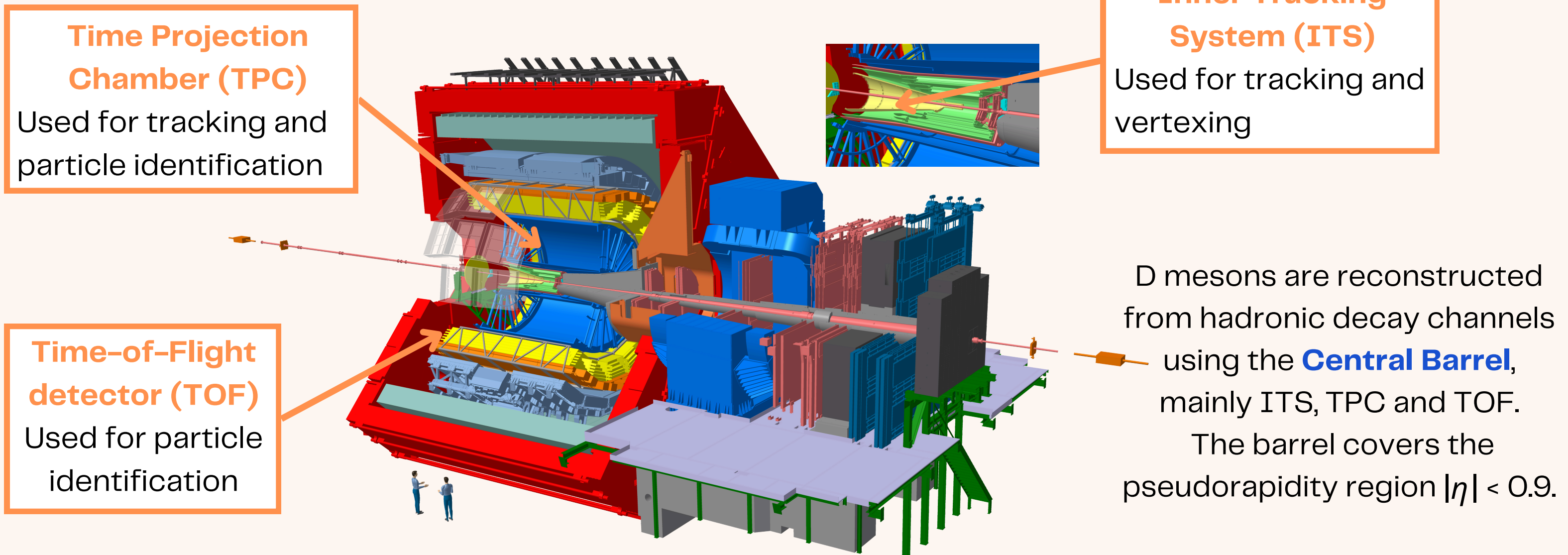
- **Run 2 and Run 3** results are **compatible** with each other.
- **PYTHIA**-based models tend to slightly overestimate the data, independently of the tune employed.
- **EPOS 4** tends to underestimate the data.

In Run 3, we are **dominated by systematic uncertainties**, which are being reduced in view of the results publication.

A **higher granularity**, especially at low p_T , is achieved thanks to access to higher statistics.

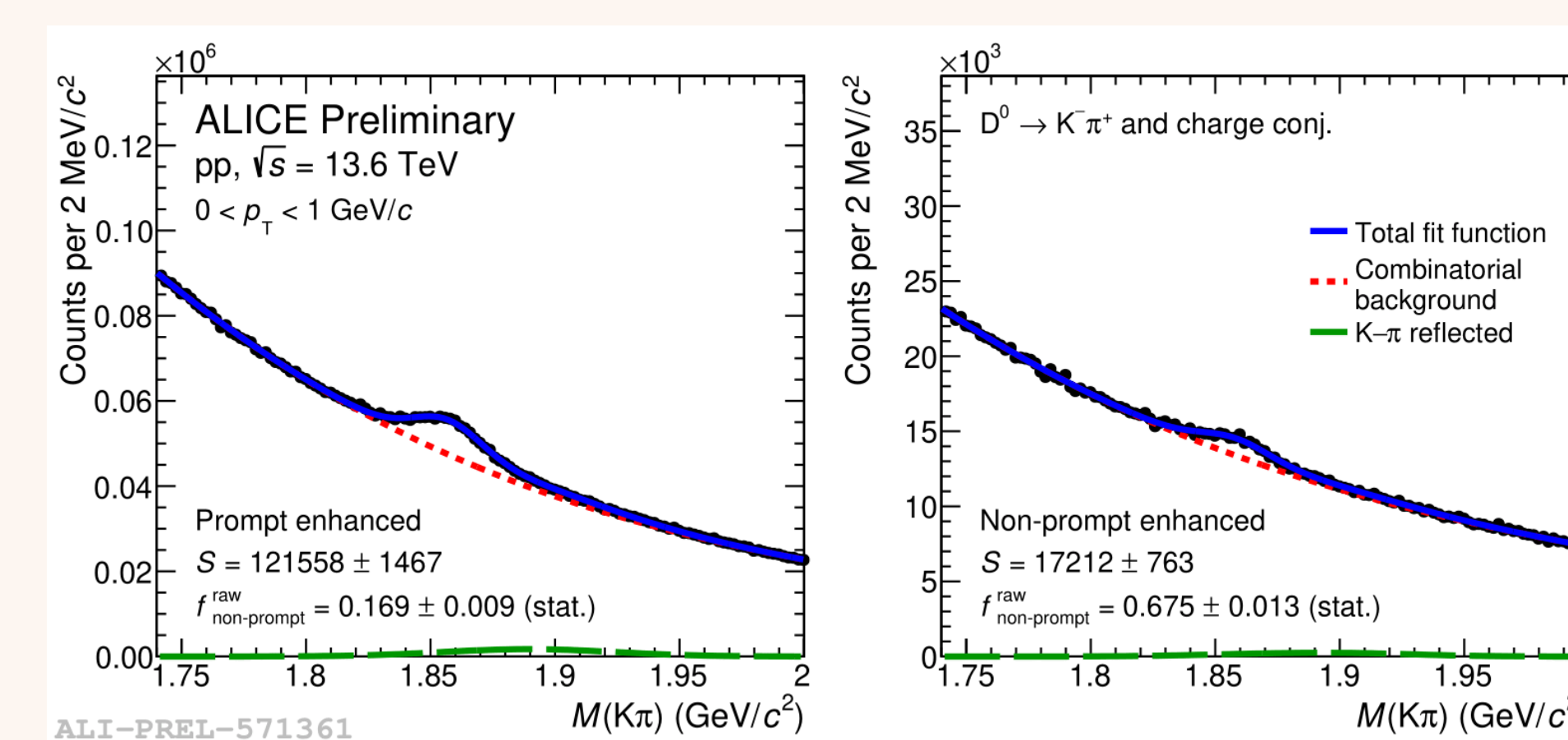
The non-prompt fraction at $0 < p_T < 1$ GeV/c is successfully measured.

2. Experimental apparatus: ALICE in Run 3



4. Analysis

We started by selecting a working point that maximised the significance of the signal by applying some cuts to the background BDT score. Then, we defined a set of cuts on the non-prompt BDT score, and we **extracted the raw yields and the efficiencies** for each cut:



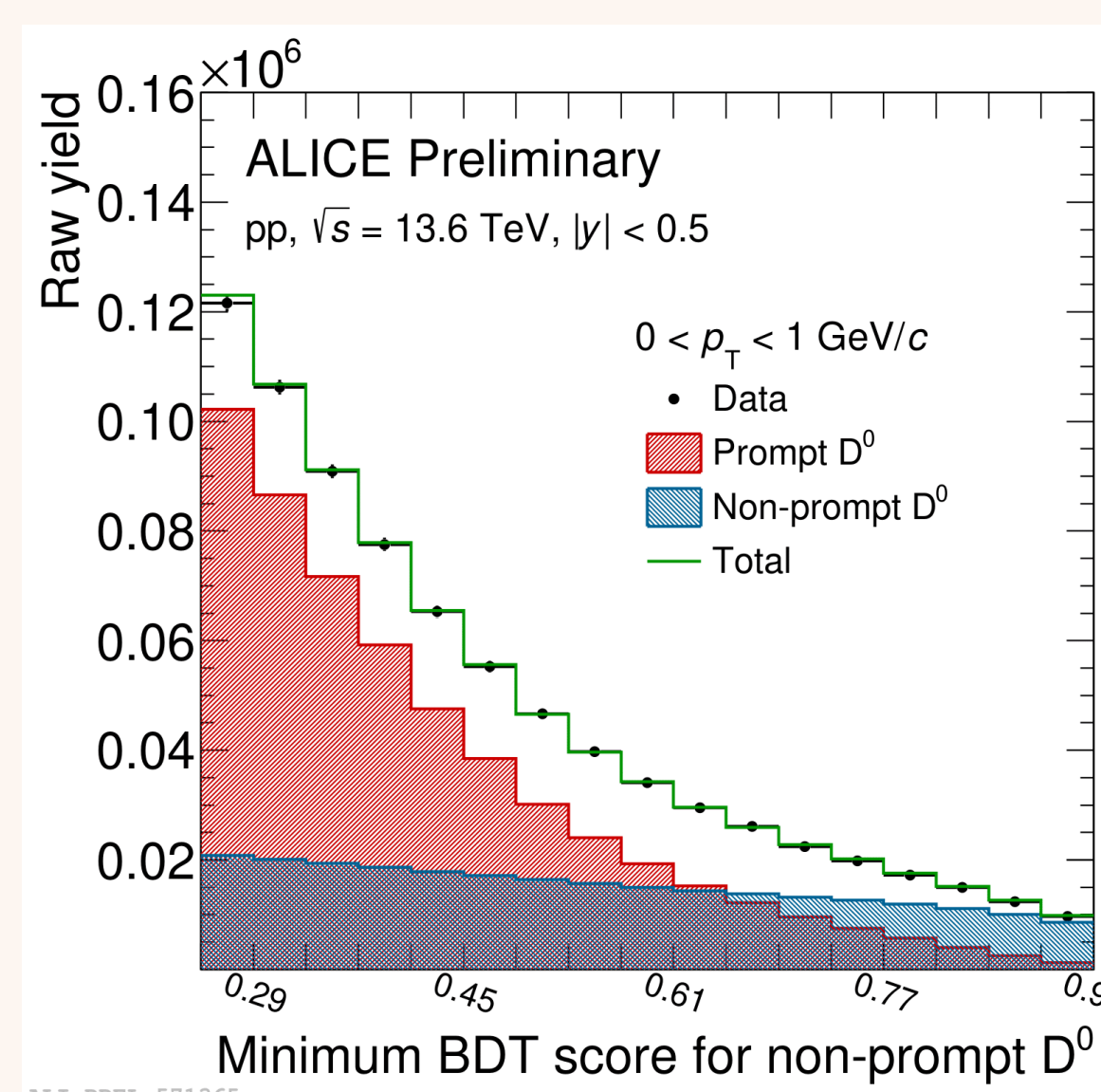
Invariant-mass fit:

- **Signal function**: single Gaussian
- **Background function**: third-degree polynomial
- **Reflection function**: double Gaussian

Efficiency extraction:

- Calculated using a MC anchored to our data sample

Solving the system of equations, we get:

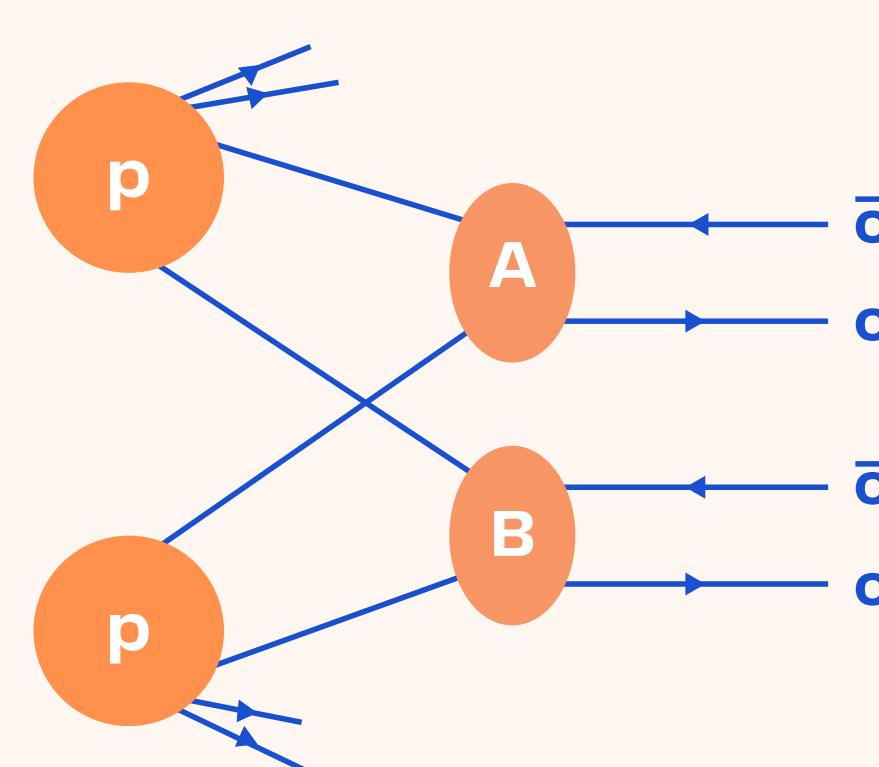


- The **red** and **blue** templates contain the corrected prompt and non-prompt yields, respectively.
- The **green line** represents the result of fitting the sum of the two corrected yields to the data.
- The **black points** represent the measured raw yields.

Systematic uncertainties considered

- **Raw-yield extraction**: between 4% and 7%.
- **Data vs MC discrepancies for topological distributions**: around 7% in all p_T ranges.
- **ML selections**: largest uncertainty at high p_T , up to 15%. At low p_T , its values are around 5%.
- **p_T weighting**: around 2% in all ranges.

6. Perspectives: study of double D^0 production



We can study the double production of D^0 mesons to investigate **double parton scattering (DPS)** [4]. In general, we measure the DPS cross section for the production of particles A and B as:

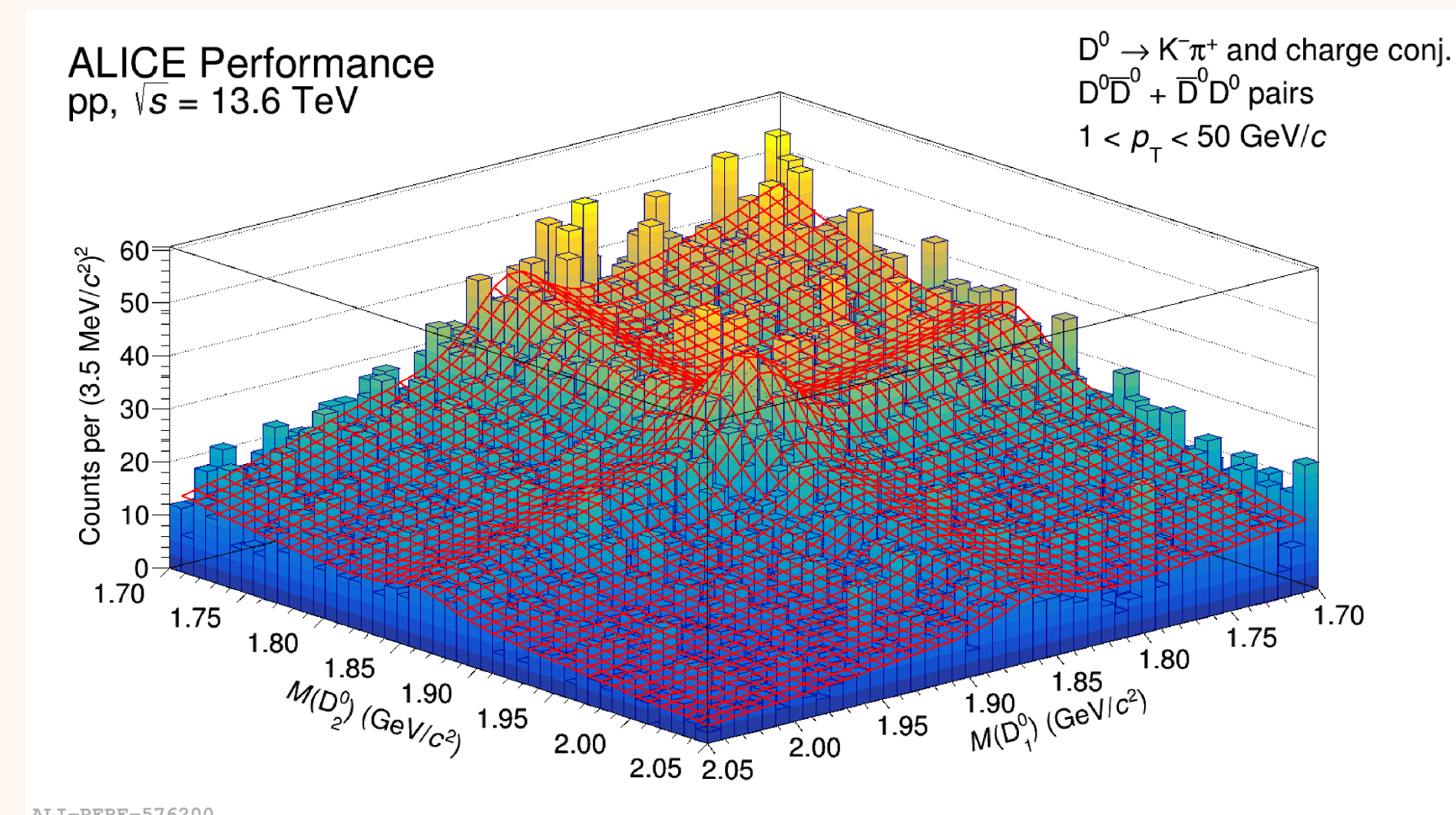
$$\sigma_{\text{DPS}}^{AB} = \frac{N}{\mathcal{L} \times (\text{Acc} \times \epsilon) \times BR_A \times BR_B}$$

where N is the raw yield, which can be obtained via a 2D invariant mass fit; \mathcal{L} is the integrated luminosity; $(\text{Acc} \times \epsilon)$ is the acceptance multiplied by the efficiency, calculated using MC simulations; and $BR_{A,B}$ is the branching ratio of the decays from which particles A and B are reconstructed, respectively.

For the double D^0 production case, to distinguish DPS from single parton scattering (SPS), we can look at different types of D^0 meson pairs:

- **Opposite-sign pairs**, like $D^0\bar{D}^0$, will be primarily produced by SPS.
- **Like-sign pairs**, such as D^0D^0 , will mostly come from DPS.

The image shows the 2D invariant-mass fit of opposite-sign pairs ($D^0\bar{D}^0 + \bar{D}^0D^0$), which will be used to extract the SPS raw yield and will be compared with the DPS one.



Related literature

[1] ALICE Collaboration, *J. High Energ. Phys.* **2023**, *92* (2023).

[2] T. Sjöstrand et al., *Comput. Phys. Commun.* **191** (2015) 159–177.

[3] K. Werner, *Phys.Rev.C* **108** (2023) *6*, 064903.

[4] LHCb Collaboration, *J. High Energ. Phys.* **2012**, *141* (2012).