

Test-beam Measurements of Instrumented Sensor Planes for a Highly Compact and Granular Electromagnetic Calorimeter

Grzegorz Grzelak
(on behalf of the LUXE ECAL group)

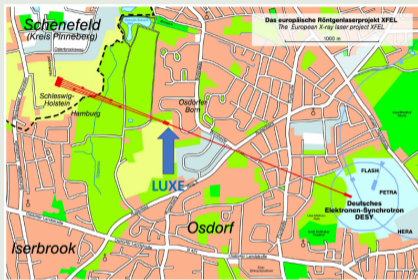
Faculty of Physics
University of Warsaw

LUXE



42nd International Conference on High Energy Physics, ICHEP-2024, 17 - 24 July 2024, Prague

LUXE: Laser Und XFEL Experiment



LUXE milestones documents:

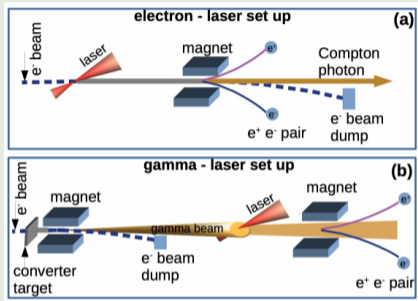
- **LOI (2019)** 1909.00860
- **CDR (2021)** EPJ ST 230, 2445 - 2560
- **TDR (2023)** 2308.00515 (EPJ ST Accepted)

LUXE Collaboration



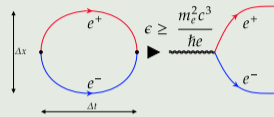
- at DESY as the host laboratory
- at Eu.XFEL 16.5 GeV electron beam
- over 20 participating institutes
- about 130 active scientists

LUXE: Two modes of operation



- e-laser: using 16.5 GeV XFEL e^- beam
- γ -laser: using bremsstrahlung γ photons
- collide them with **High Power** (40 or 350 TW) optical **Laser (HPL)** [phase-0 / phase-1]

non-linear and non-perturbative QED

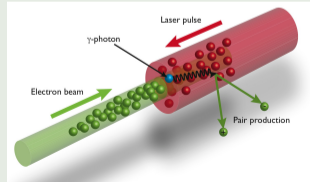


- physics at and above Schwinger limit:

$$\mathcal{E}_{cr} = \frac{m_e^2 c^3}{\hbar e} = 1.3 \times 10^{18} \text{ V/m,}$$

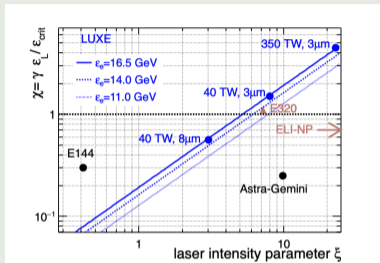
- boosted frame: $\chi = \gamma \frac{\mathcal{E}_{HPL}}{\mathcal{E}_{cr}}$

$$RMS(\mathcal{E}_{HPL}) \sim 10^{14} \text{ V/m } (\times 10^4 e^- \text{ boost})$$

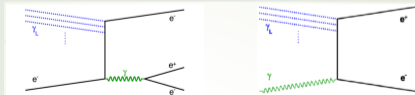


LUXE: the detectors challenge: very high rate of particles

Parameters space and e^+ rate

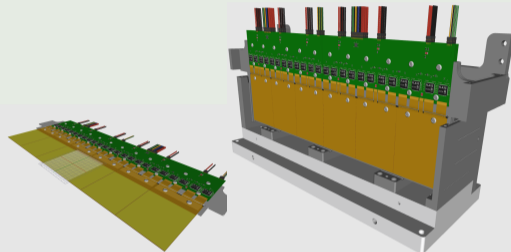


- laser intensity, dimensionless amplitude of \mathcal{E} field):
$$\xi = \frac{m_e \mathcal{E}_L}{\omega_L \mathcal{E}_{cr}}, \quad \omega_L - \text{laser frequency}$$



- **expected positron rate:** $10^{-5} - 10^6$ per BX, EM showers overlap at high multiplicity

Solution for e^+ calorimetry: ECAL-p



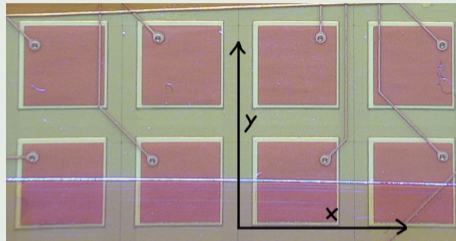
- compact, high density sampling calorimeter
- small Molière radius: ~ 9.3 mm
- high granularity
- 21 layers of 3.5 mm ($1X_0$) tungsten absorber
- 1 mm gaps instrumented with active sensors

Gallium Arsenide sensor:



- National Research Tomsk State University
- GaAs crystals compensated with chromium
- $4.7 \times 4.7 \text{ mm}^2$ pads, 0.3 mm gap between pads
- pads are made of $0.05 \text{ }\mu\text{m}$ vanadium layer
- sensor thickness $500 \text{ }\mu\text{m}$
- total wafer area: $51.9 \times 75.6 \text{ mm}^2$
- $1 \text{ }\mu\text{m}$ thick Al traces in the gaps between pads
- better radiation tolerance than silicon

Silicon sensor:



- produced by Hamamatsu (CALICE design)
- Si crystals: p+ on n substrate diodes
- $5.5 \times 5.5 \text{ mm}^2$ pads, 0.01 mm gap between pads
- few nm pads Al metalization
- sensor thickness $500 \text{ }\mu\text{m}$ ($320 \text{ }\mu\text{m}$)
- total wafer area: $89.7 \times 89.7 \text{ mm}^2$
- external kapton fan-outs with copper traces connected to the sensor pads with conductive glue

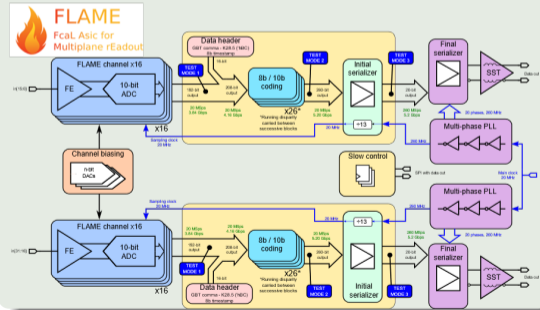
FLAME/FLAXE front-end ASIC designed by AGH University of Krakow

- FLAME (FcaL Asic for Multiplane rEadout) is a 32-channel ASIC in CMOS 130 nm
- 10-bit ADC in each channel, two fast (5.2 Gbps) serializers and data transmitters
- FLAME has been already used in several test-beams of FCAL and LUXE-ECAL collaborations
- final DAQ version will use a new front-end ASIC FLAXE, which is based on FLAME (in progress)

FLAME ASIC specification

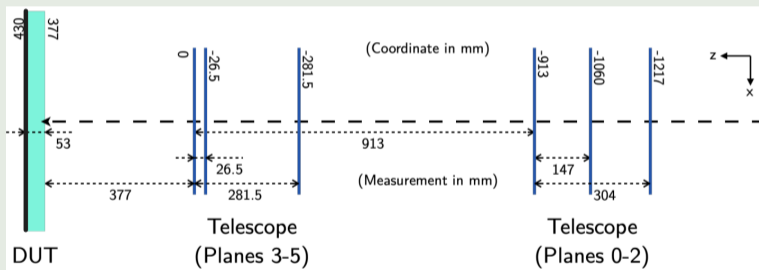
- Analog front-end in each channel:
 - CR-RC shaping ($T_{peak} \sim 50$ ns)
 - two switched gains (high gain for MIPs, low gain for showers)
 - $C_{in} \sim 20 - 40$ pF
- 10-bit ADC per channel:
 - $f_{sample} = 20$ MHz
 - ENOB > 9.5 (effective resolution)
 - FPGA to extract amplitude and time
 - Power < 350 μ W @ 20 MHz

32-channel FLAME ASIC



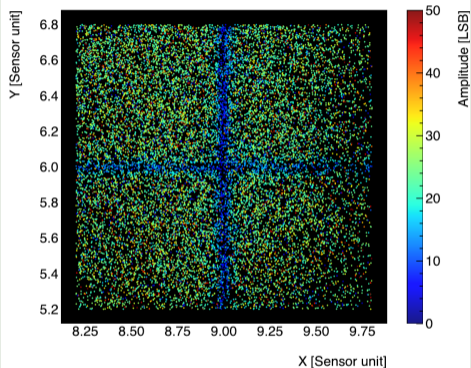
Beam telescope (EUNET collaboration), scintillators and Detector Under Test (DUT)

- Electrons arrive from the right, pass the first scintillator, then six ALPIDE pixel sensors, the second scintillator, and hit the sensor, denoted as DUT (Detector under Test)
~ 35 μm resolution of the track extrapolated from the TB telescope

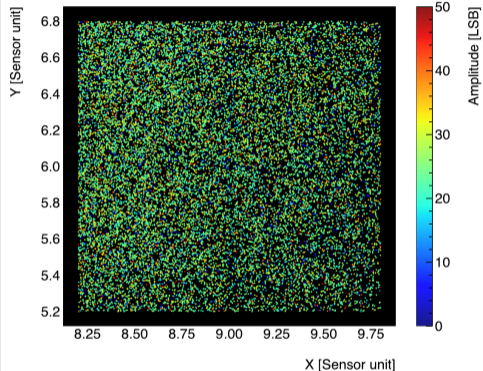


- Two 16×8 pad arrays of Silicon sensors and two 15×10 pad arrays of GaAs sensors were tested on 5 GeV electron beam at the DESY-II facility
- investigated were homogeneity of the sensor response, edge effects and signal sharing between pads
- in addition: test of the FPGA based data on-line preprocessing (amplitude and time reconstruction)

GaAs sensor (2×2 pads)

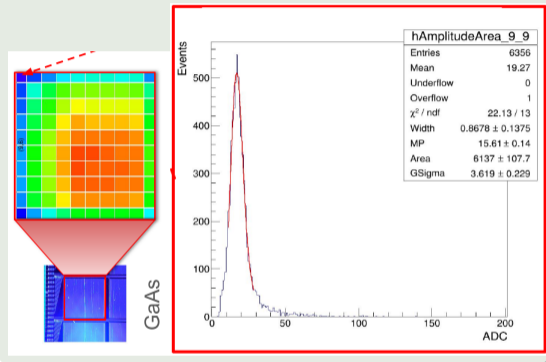


Si sensor (2×2 pads)

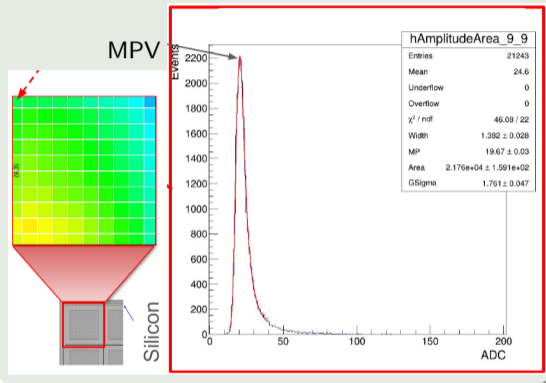


- after alignment with beam telescope ($\sim 35 \mu\text{m}$ resolution on DUT XY)
- color (Z scale) indicates the size of the signal
- loss of signal for GaAs sensor in the region between pads

GaAs sensor

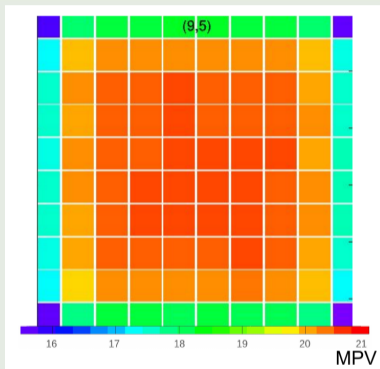


Si sensor



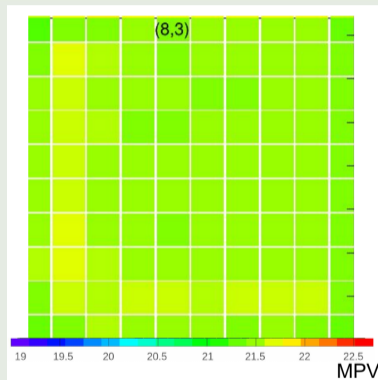
- Pads were subdivided into 10×10 XY sections and plotted was amplitude distribution in each section
- Fits of Landau distribution convoluted with Gaussian \rightarrow Most Probable Value (MPV) on next page
- color (Z scale) encodes the statistic of hits

GaAs sensor (single pad)



- GaAs: Drop in amplitude around edges and esp. in corners
- color (Z scale) encodes the MIP value

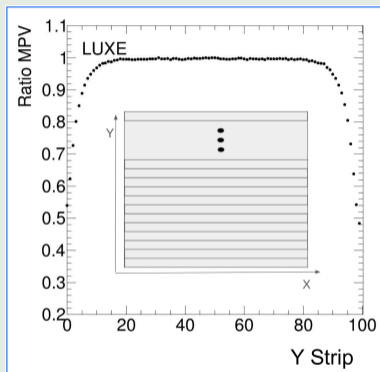
Si sensor (single pad)



- Si: more uniform response, but...
- ... L-shaped area of a bit higher amplitude

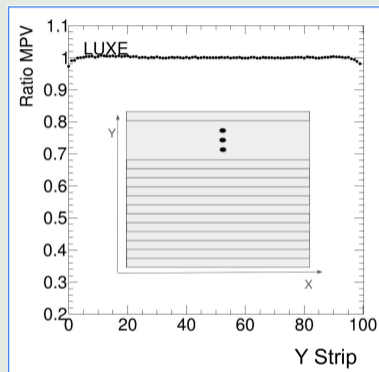
Sensor homogeneity study (cont.) : 100 vertical (Y) strips on pad

GaAs sensor (single pad)



- GaAs Y scan: **MPV drop : 50%** wrt pad center

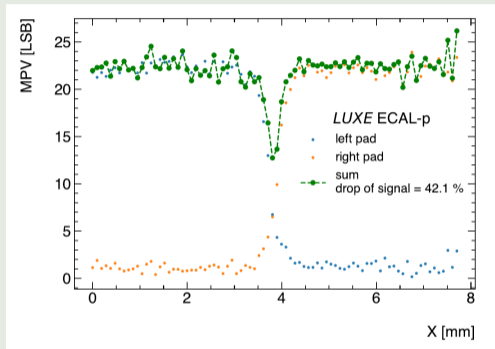
Si sensor (single pad)



- Si Y scan: **MPV drop : 2-3%** wrt pad center

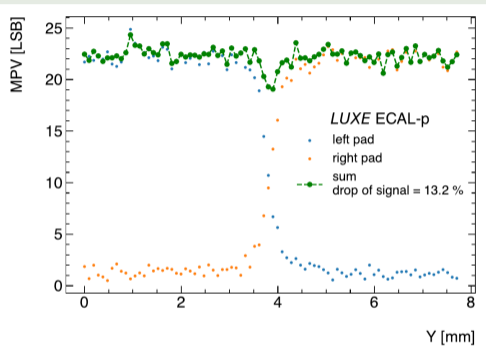
- Normalized to MPV of central strip. Similar response along X direction

GaAs sensor: 2×1 neighbor pads: X scan



- X scan - traces between pads
- sum of MPV drops of 40%

GaAs sensor: 2×1 neighbor pads: Y scan

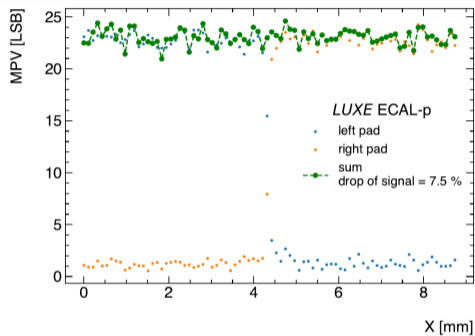


- Y scan - without traces between pads
- sum of MPV drops of 15%

- MPV measured as a function of x and y , crossing the area between two pads
- gap between GaAs pads $300 \mu\text{m}$

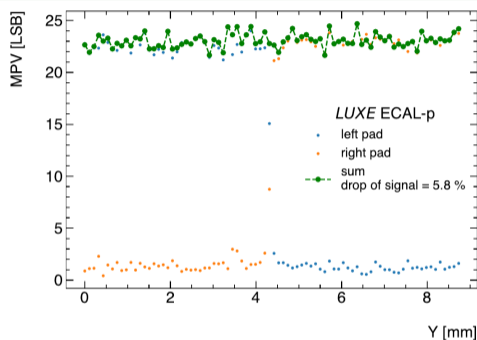
Si sensor: signal sharing between pads

Si sensor: 2×1 neighbor pads: X scan



- no signal drop is observed

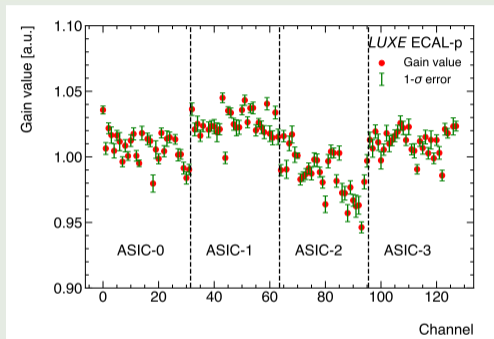
Si sensor: 2×1 neighbor pads: Y scan



- no signal drop is observed

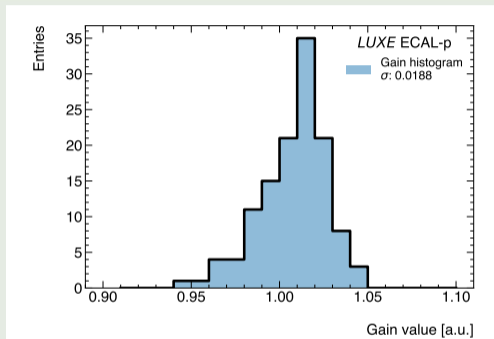
- MPV measured as a function of x and y, crossing the area between two pads
- gap between Si pads $10 \mu\text{m}$

Relative gain [a.u.]



- relative gain vs. channel number for 4 ASICs
- using calibrated charge injector (2 – 12 fC for MIPs gain)

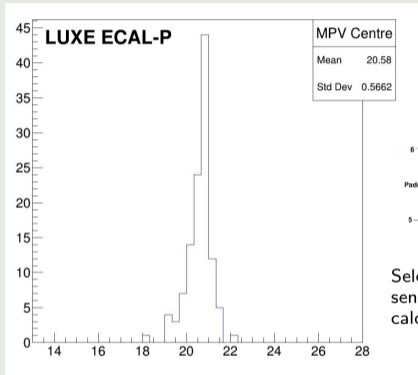
Relative gain [a.u.]



- distribution of relative gains
- RMS $\sim 2\%$

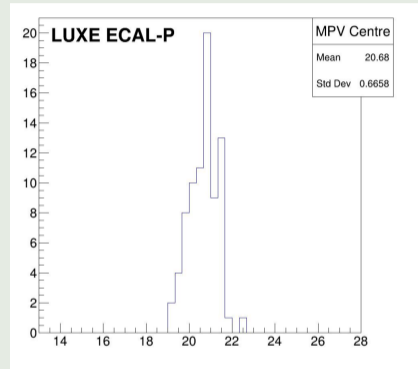
- good homogeneity of front-end preamplifiers, some dependence on ASIC fabrication

MPV: GaAs sensors



- RMS: 2.8%

MPV: Si sensors



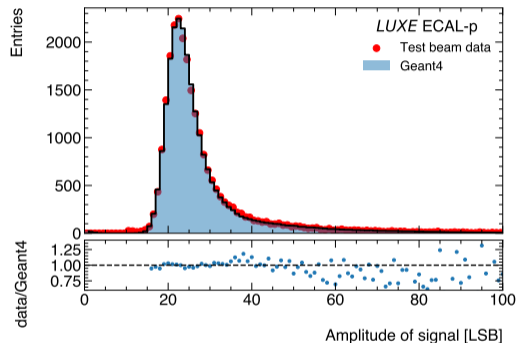
- RMS: 3.2%

- MPV distribution after gain correction, excluding edge effect (20% margin)

response of a Si sensors to 5 GeV electrons

- energy loss dE/dx [GeV] in 500 μm Si sensor from Geant4
- energy loss converted into number of charge carriers using 3.6 eV per electron-hole pair
- gain of the read-out chain determined from charge injection: 3.45 LSB/fC
- as a cross-check 3.46 LSB/fC was obtained fitting the gain as a free parameter

response of a Si sensors to 5 GeV electrons



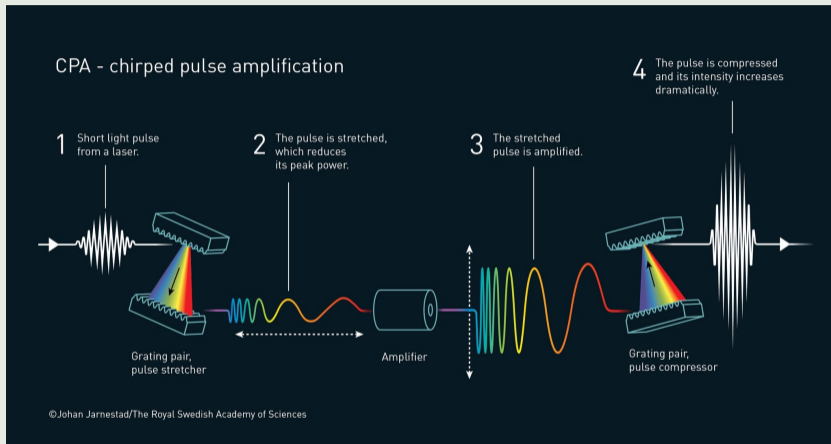
- good agreement DATA/MC for MIP

- **two types of semiconductor sensors (GaAs and Si)** for high density EM calorimeters **were tested** at 5 GeV electron beam at DESY
- energy losses for MIPs are well described by Landau distribution convoluted with Gaussian function
- homogeneity and signal sharing study were performed using hit position from the beam telescope
- **for GaAs sensors edge effect are observed** related to aluminum tracers and bigger gap between pads (up to 40-50% signal drop)
- **for silicon sensor edge effects are barely visible**
- after gain correction, in the central region of pads the **homogeneity** of the sensors amounts to **2.8 and 3.2 %** for the GaAs and Si sensors, respectively
- collected data are in **good agreement with Geant4 based MC**
- readout electronics absolute gain agrees between MC simulations and independent lab measurement (converting the energy loss into charge and using the gain of the readout chain)

- **two types of semiconductor sensors (GaAs and Si) for high density EM calorimeters were tested** at 5 GeV electron beam at DESY
- energy losses for MIPs are well described by Landau distribution convoluted with Gaussian function
- homogeneity and signal sharing study were performed using hit position from the beam telescope
- **for GaAs sensors edge effect are observed** related to aluminum tracers and bigger gap between pads (up to 40-50% signal drop)
- **for silicon sensor edge effects are barely visible**
- after gain correction, in the central region of pads the **homogeneity** of the sensors amounts to **2.8 and 3.2 %** for the GaAs and Si sensors, respectively
- collected data are in **good agreement with Geant4 based MC**
- readout electronics absolute gain agrees between MC simulations and independent lab measurement (converting the energy loss into charge and using the gain of the readout chain)

● **Thank You Very Much for Your Attention !**

BACKUP PLOTS FOLLOWS...

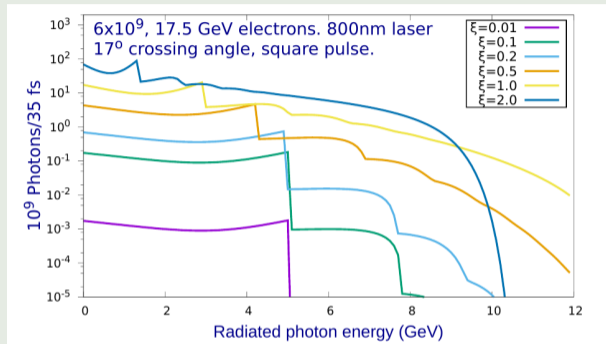


2018 Nobel Prize Donna Strickland and Gerard Mourou
“for method of generating **high-intensity, ultra-short optical pulses**”

dimensionless intensity parameter (field energy density) ξ^2

- $\xi^2 = 4\pi\alpha\left(\frac{\mathcal{E}_L}{m_e\omega_L}\right)^2 = \left(\frac{m_e\mathcal{E}_L}{\omega_L\mathcal{E}_{cr}}\right)^2$, ← “classical picture”
 ω_L - laser frequency, ξ - “dimensionless amplitude” of \mathcal{E} field
- $\xi^2 = 4\pi\alpha\tilde{\lambda}_L\tilde{\lambda}_C^2 n_L$, ← “quantum picture”
 $\tilde{\lambda}_L$ and $\tilde{\lambda}_C$ - reduced laser and Compton wavelengths,
 $\tilde{\lambda}_L \sim 1 \mu\text{m}$
 $\tilde{\lambda}_C \sim 10^{-6} \mu\text{m}$
 n_L - number density of laser photons
- for low and moderate $\xi \lesssim 1$ the probability of net absorption of n laser photons $\propto (\xi^2)^n \sim \alpha^n$
(consistent with perturbative QED vertex counting)

Non-linear Compton γ spectrum

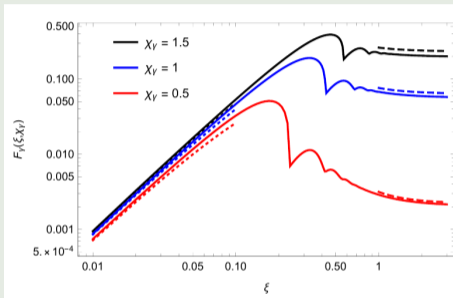


- low laser intensity (ξ) \rightarrow KleinNishina process
- $\xi \nearrow$: increasing flux of Compton photons
- $\xi \nearrow$: **shift of Compton edge** with laser intensity (\rightarrow next page)
- **additional structure** due to multi-photon absorption

$$e^- + n\gamma_L \rightarrow e^- + \gamma$$

- for monochromatic, circularly polarized laser pulse: $|\vec{\mathcal{E}}| = \text{const}$
- in transverse plane circular motion of electron with frequency ω_L
- **energy accumulated in this transverse degree of freedom can be treated as extra, effective mass of the electron**
- electron transverse momentum: $P_{\perp} \sim \xi m$
- $E^2 = m^2 + P_{\perp}^2 + P_{\parallel}^2 \sim (1 + \xi^2)m^2 + P_{\parallel}^2$
- **electron effective mass: $\bar{m} = m\sqrt{1 + \xi^2}$**
- \rightarrow **shift of the lowest order Compton edge**
(scaling as $1/\sqrt{1 + \xi^2}$)
- \rightarrow can be used to monitor the intensity parameter ξ

full calculation and asymptotic behavior (dotted-dashed)



- in a constant static field: $\propto \exp\left(-\pi \frac{\mathcal{E}_{cr}}{\mathcal{E}}\right)$ (Schwinger process)
- in plane wave laser (asymptotic): $\propto \exp\left(-\frac{8}{3} \frac{1}{1+\cos\theta} \frac{m_e}{\omega_L} \frac{\mathcal{E}_{cr}}{\mathcal{E}}\right)$
- good agreement for $\xi \ll 1$ and $\xi > 1$
- initial growth with ξ then drop due to the Compton edge shift