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Flavourful global fit to LHCb data with the general 2HDM: An update

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ICHEP 2024

2024年07月19日

Peter Athron, Andreas Crivellin, Tomas Gonzalo, Syuhei Iguro, **CS [WIP]**

- **SM**
 - Electroweak interactions
- **Second Higgs: Yukawa Lagrangian**
 - Flavour changing transitions
- **Flavour Anomalies**
 - Charged anomalies
 - Neutral anomalies
 - New diagrams from G2HDM
 - Wilson coefficients at LO
- **Scans and Results**
- **Summary**

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Fermions

Flavour states

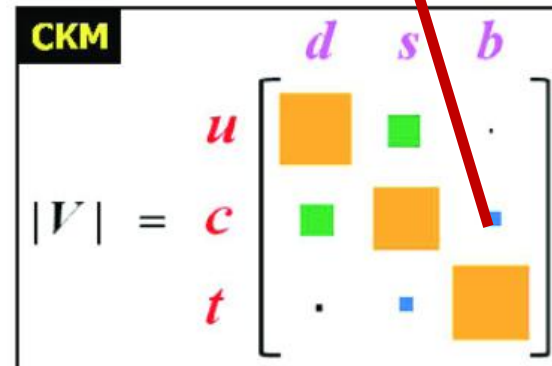
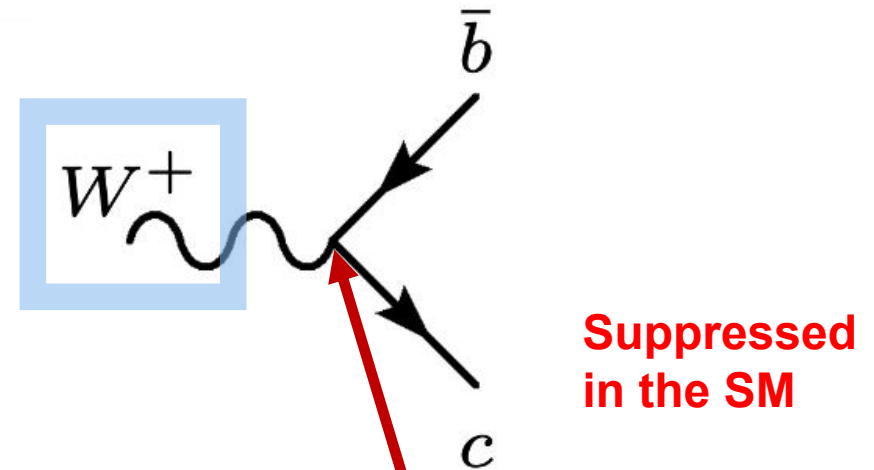
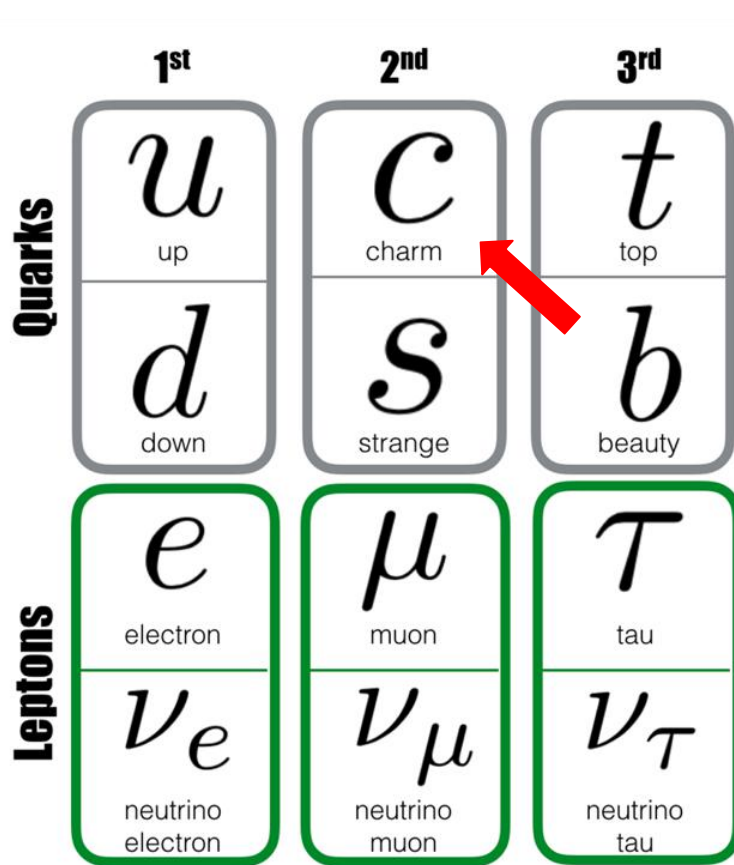
	1 st	2 nd	3 rd
Quarks	u up	C charm	t top
	d down	S strange	b beauty
Leptons	e electron	μ muon	τ tau
	ν_e neutrino electron	ν_μ neutrino muon	ν_τ neutrino tau

12 flavours



Image credit: Physik-Institut - UZH

Flavour changing transitions



Cabibbo-Kobayashi-Maskawa (CKM)

Image credit: Physik-Institut - UZH

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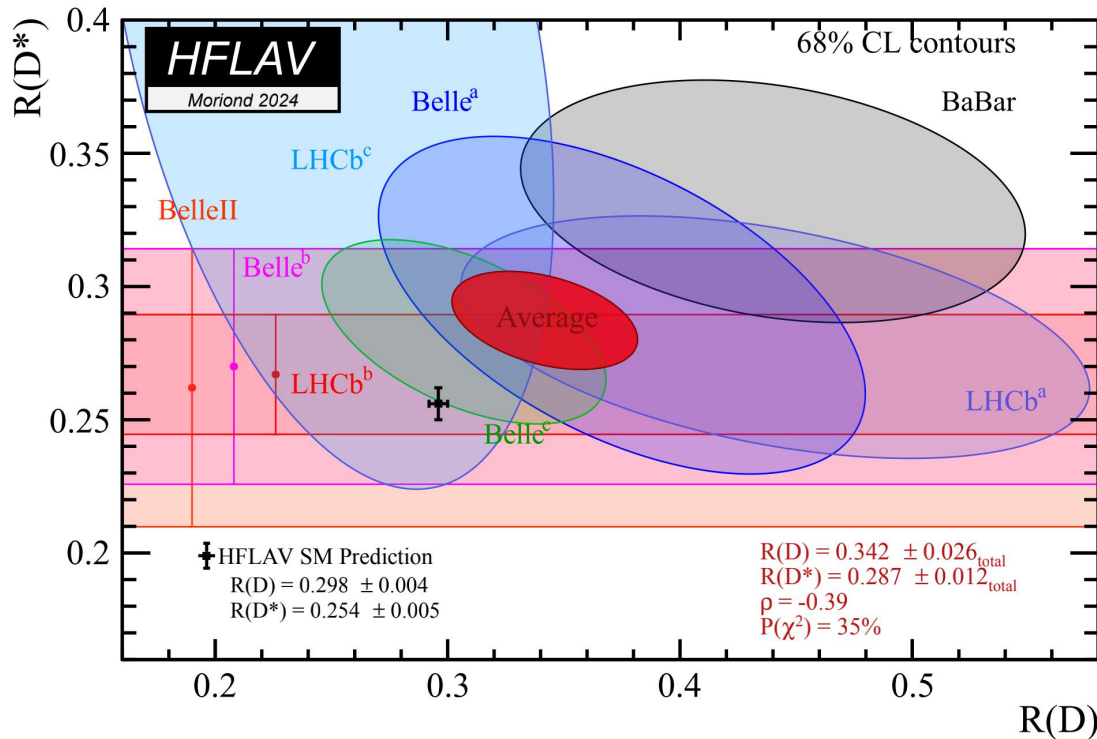
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Charged anomalies

$$R_D = \frac{\Gamma(\bar{B} \rightarrow D\tau\bar{\nu})}{\Gamma(\bar{B} \rightarrow Dl\bar{\nu})} \quad R_{D^*} = \frac{\Gamma(\bar{B} \rightarrow D^*\tau\bar{\nu})}{\Gamma(\bar{B} \rightarrow D^*l\bar{\nu})}$$

$$R_{D^{(*)}}^{\text{exp}} > R_{D^{(*)}}^{\text{SM}}$$



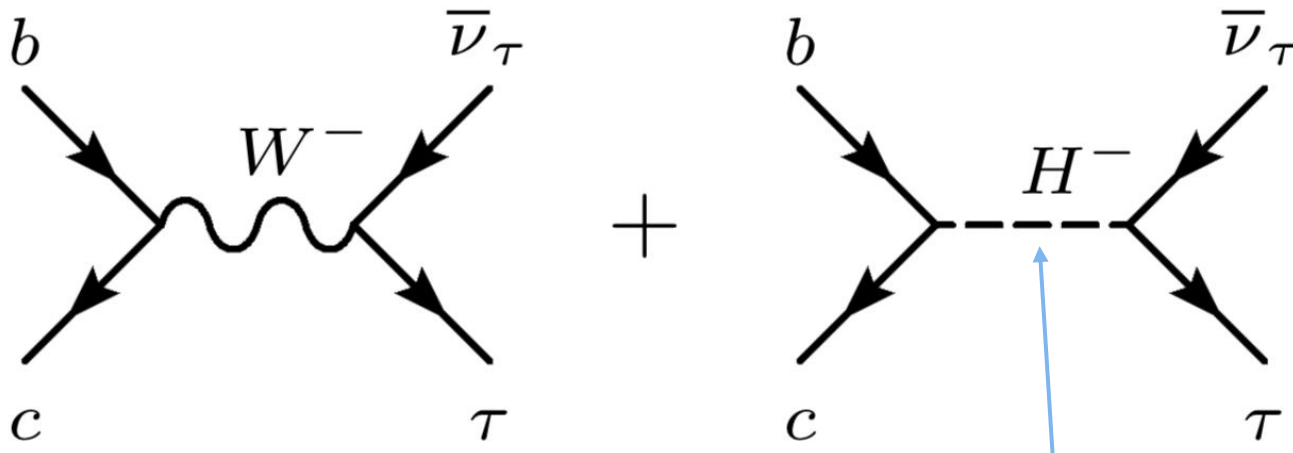
At 3.2σ

**Interference
with NP?**

From **Moriond 2024**

Charged anomalies

$$R_{D^{(*)}}^{\text{exp}} > R_{D^{(*)}}^{\text{SM}}$$



A second Higgs doublet can provide answers: Charged Higgs!



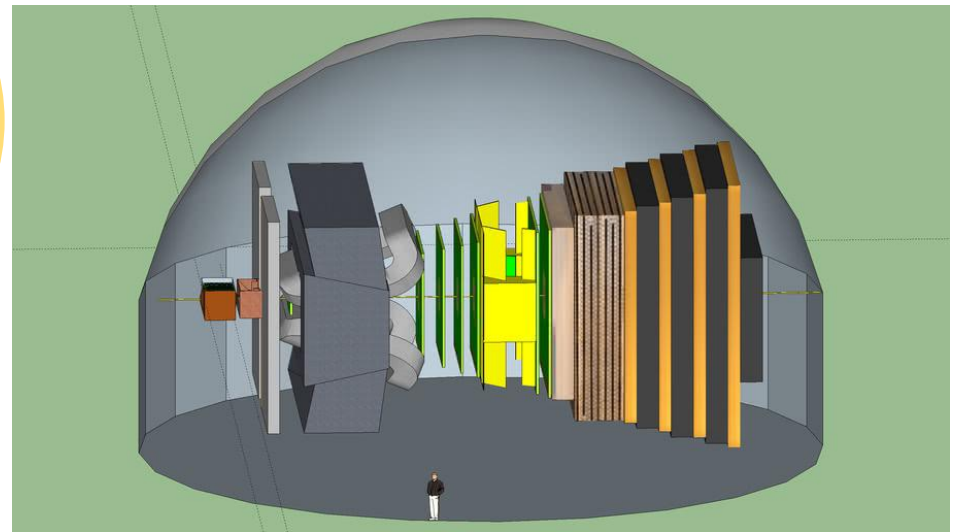
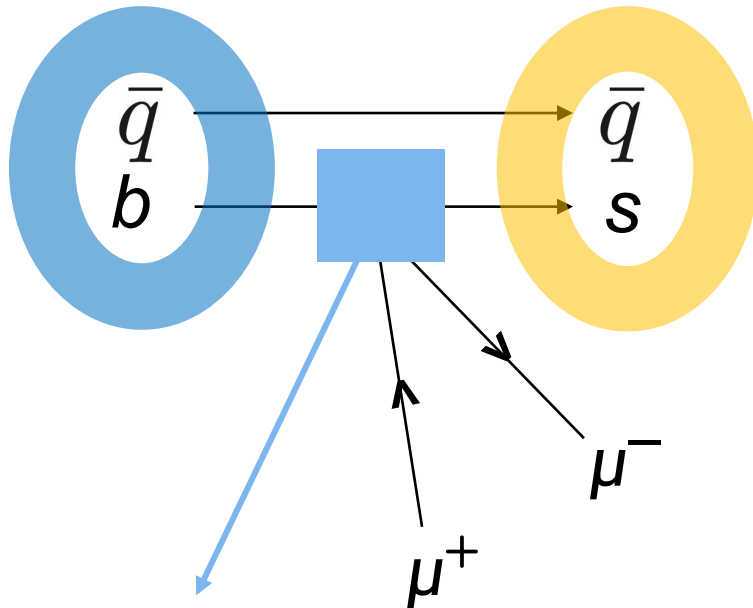
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Neutral anomalies

Most precise measurements
come from the LHCb detector

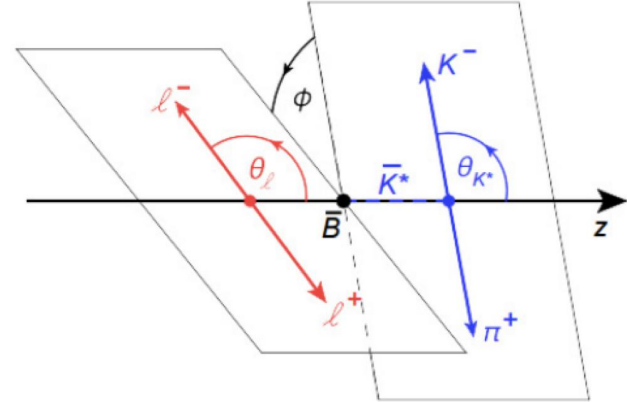
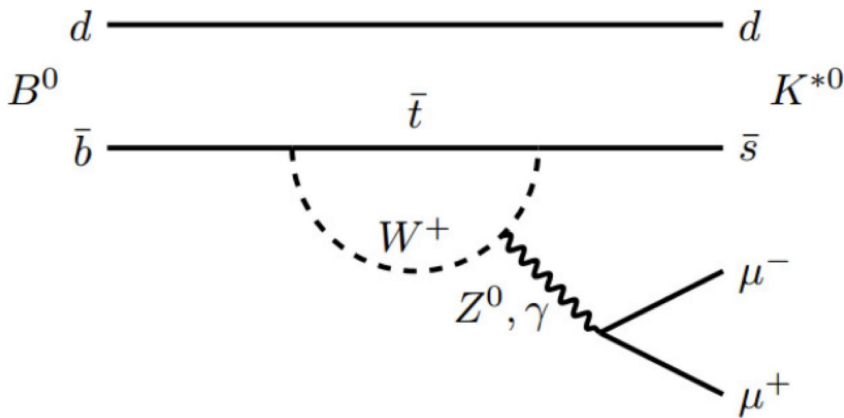


$$B \rightarrow K^{(*)} \mu^+ \mu^-$$



Explanations from the NP point of view of any
anomaly will be via effective **Wilson coefficients**

B meson semilep. decays: Angular observables



Decay fully described by three helicity angles $\vec{\Omega} = (\theta_\ell, \theta_K, \phi)$ and $q^2 = m_{\mu\mu}^2$

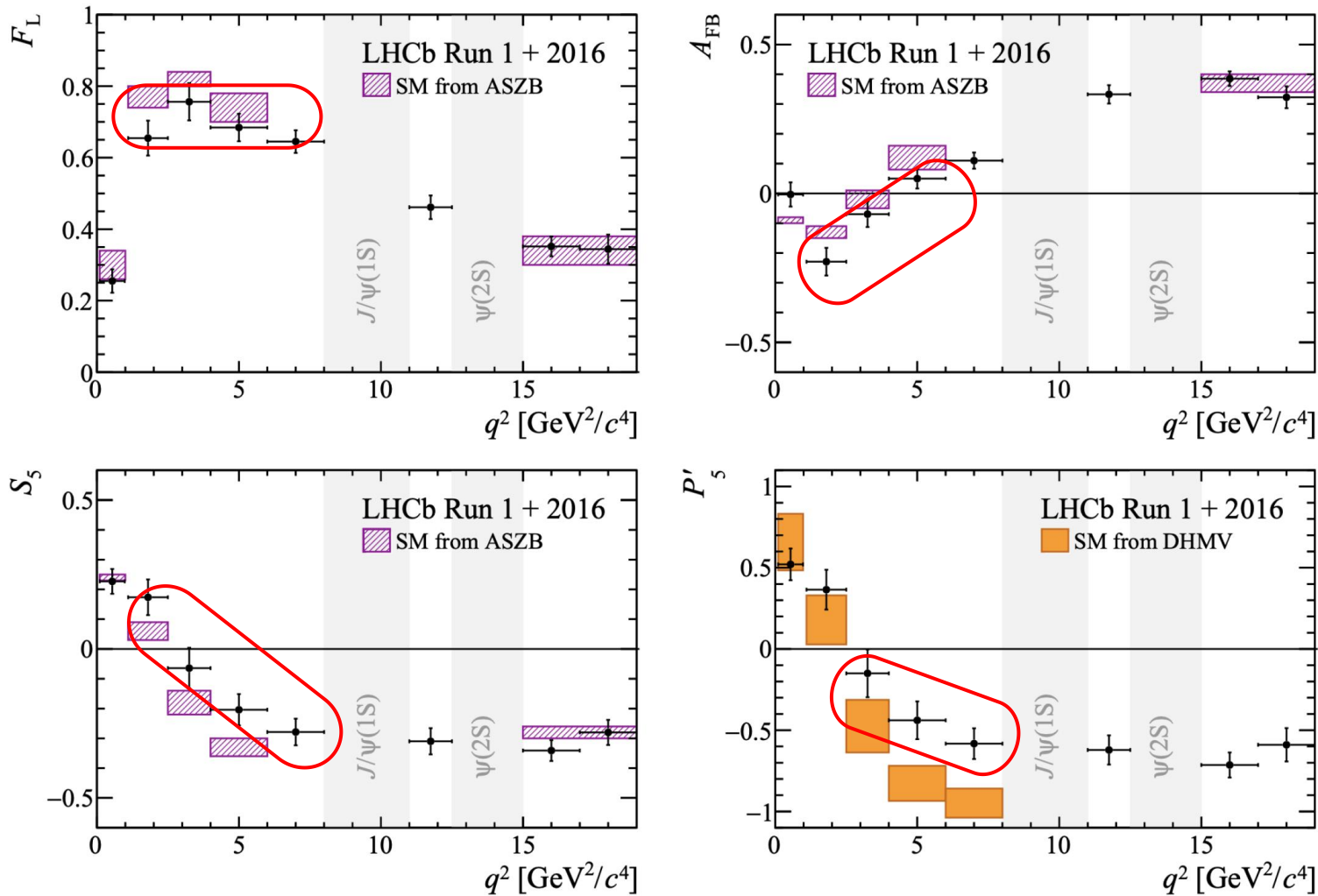
$$\frac{1}{d(\Gamma + \bar{\Gamma})/dq^2} \frac{d^3(\Gamma + \bar{\Gamma})}{d\vec{\Omega}} = \frac{9}{32\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\ \left. - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi \right. \\ \left. + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi \right. \\ \left. + \frac{4}{3} A_{\text{FB}} \sin^2 \theta_K \cos \theta_\ell + S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi \right. \\ \left. + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

$$P'_5 = \frac{S_5}{\sqrt{F_L(1 - F_L)}}$$



Most famous

Anomalies in B meson semileptonic decays



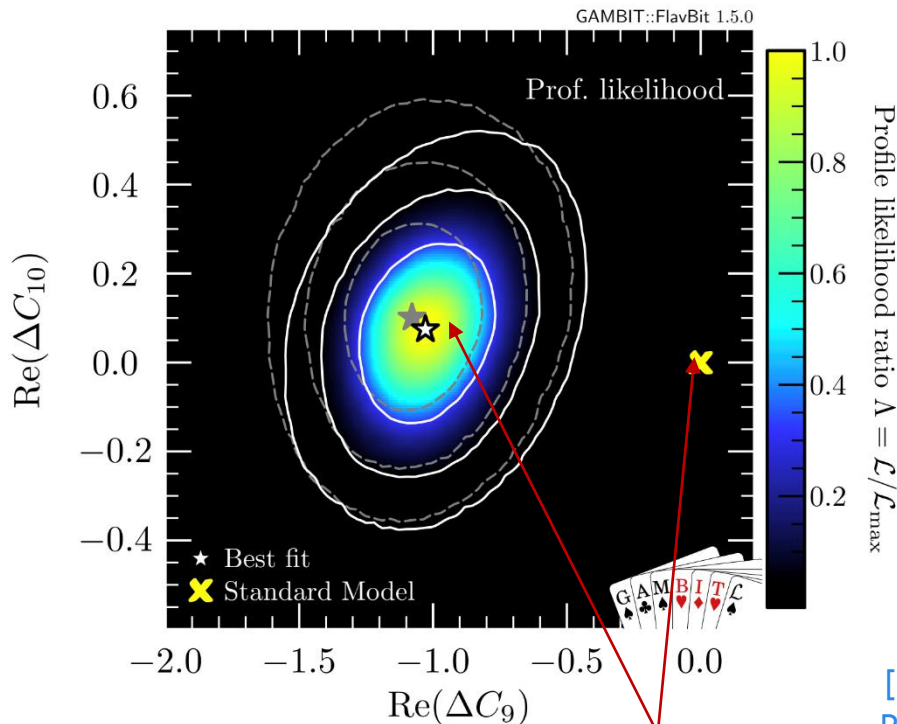
**$\sim 3\sigma$
deviation**

LHCb collaboration, Measurement of CP-averaged observables in the $B^0 \rightarrow K^* \mu^+ \mu^-$ decay, Phys. Rev. Lett. 125 (2020) 011802 [arXiv:2003.04831]

Fit all Wilson Coefficients (WCs)

- $b \rightarrow s\mu^+\mu^-$ observables included in the model independent analysis from [J. Bhom et al, arXiv: 2006.03489] get modified by the new Wilson coefficients.

>200 observables on $b \rightarrow s$ transitions
(most of them are angular observables)



~6σ deviation

All observables with $\chi_{\text{SM}}^2 = 157.28$ ($\chi_{\text{min}}^2 = 100.34$; Pull _{SM} = 4.3σ)			
δC_7 0.05 ± 0.03		δC_8 -0.71 ± 0.43	
$\delta C_7'$ -0.01 ± 0.02		$\delta C_8'$ -0.09 ± 0.86	
δC_9^μ -1.11 ± 0.19	δC_9^e -6.69 ± 1.37	δC_{10}^μ 0.08 ± 0.25	δC_{10}^e 3.97 ± 4.99
$\delta C_9'^\mu$ 0.18 ± 0.35	$\delta C_9'^e$ 1.84 ± 1.75	$\delta C_{10}'^\mu$ -0.13 ± 0.21	$\delta C_{10}'^e$ 0.05 ± 5.01
$C_{Q_1}^\mu$ -0.07 ± 0.12	$C_{Q_1}^e$ -1.52 ± 0.98	$C_{Q_2}^\mu$ -0.10 ± 0.14	$C_{Q_2}^e$ -4.36 ± 1.46
$C_{Q_1}'^\mu$ 0.05 ± 0.12	$C_{Q_1}'^e$ -1.40 ± 1.56	$C_{Q_2}'^\mu$ -0.17 ± 0.15	$C_{Q_2}'^e$ -4.33 ± 2.33

[T. Hurtha, F. Mahmoudi, S. Neshatpour
Phys.Rev.D 102 (2020) 5, 055001]

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Fermions + second Higgs doublet

	1 st	2 nd	3 rd
Quarks	u up	C charm	t top
	d down	S strange	b beauty
Leptons	e electron	μ muon	τ tau
	ν_e neutrino electron	ν_μ neutrino muon	ν_τ neutrino tau

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + H_1 + i\eta_1) \end{pmatrix}$$

$$\Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + H_2 + i\eta_2) \end{pmatrix}$$

Adding a second Higgs doublet is one of the simplest extensions of the SM

Flavour changing transitions

	1 st	2 nd	3 rd
Quarks	u up	C charm	t top
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Leptons	e electron	μ muon	τ tau
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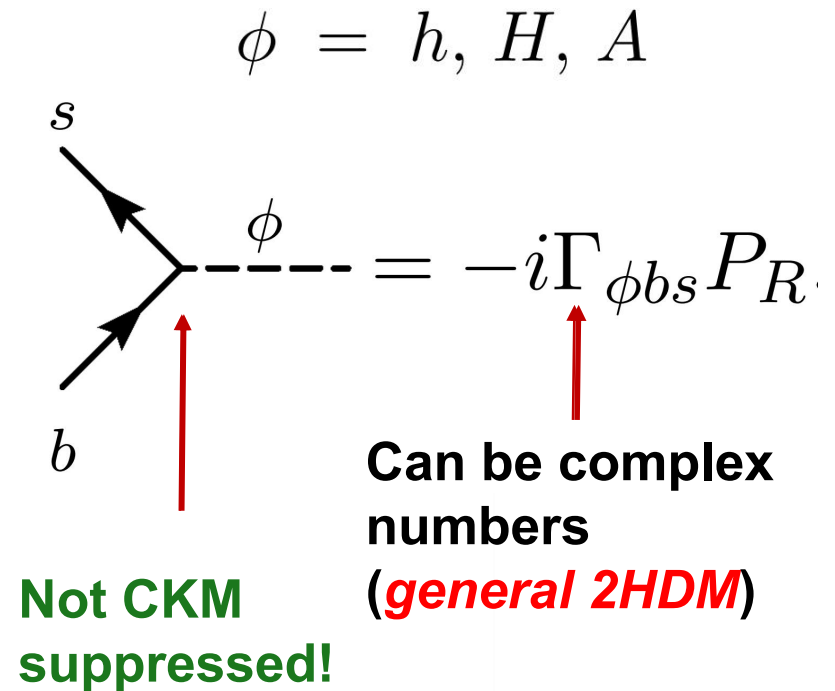


Image credit: Physik-Institut - UZH

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Yukawa Lagrangian

$$-\mathcal{L}_{Yukawa} = \bar{u}_b \left(V_{bc} \rho_d^{ca} P_R - V_{ca} \rho_u^{cb*} P_L \right) d_a H^+ + \bar{\nu}_b \rho_\ell^{ba} P_R l_a H^+ + \text{h.c.} \\ + \sum_{f=u,d,\ell} \sum_{\phi=h,H,A} \bar{f}_b \Gamma_f^{\phi ba} P_R f_a \phi + \text{h.c.},$$

general 2HDM (G2HDM)

New couplings constrained by

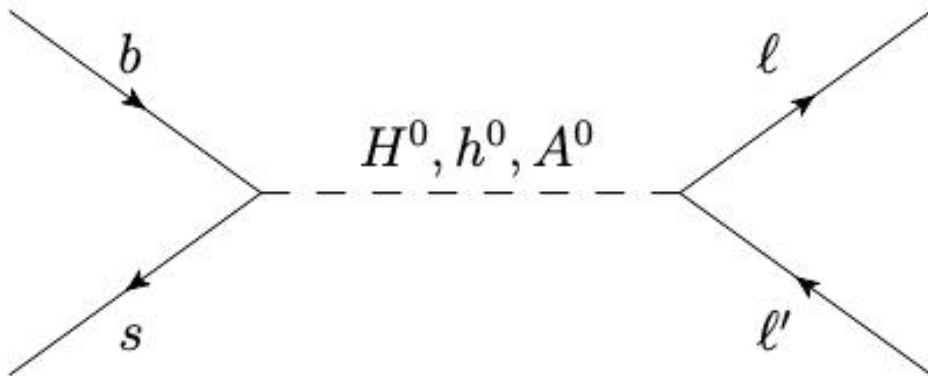
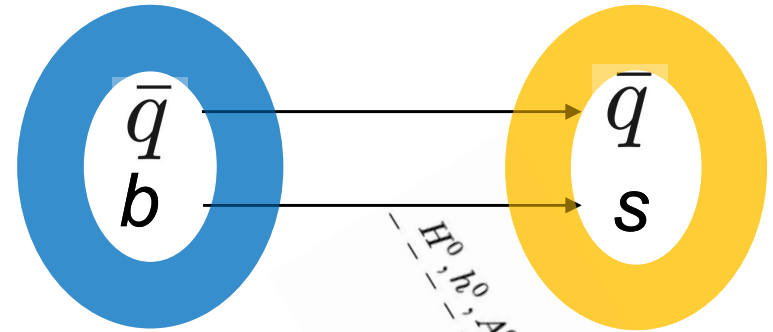
- *Stability, perturbativity and unitarity*
- *Strong flavour constraints*

WCs with in the G2HDM

Tree level diagrams



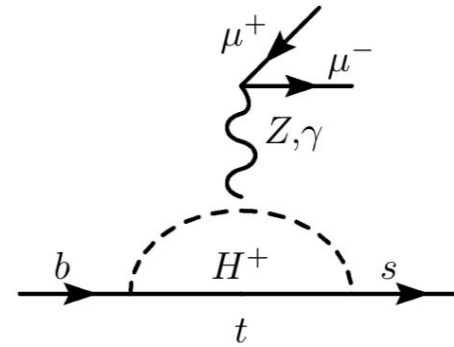
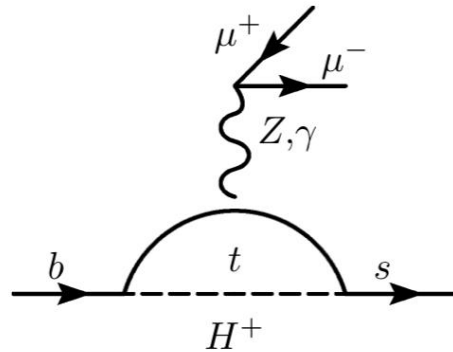
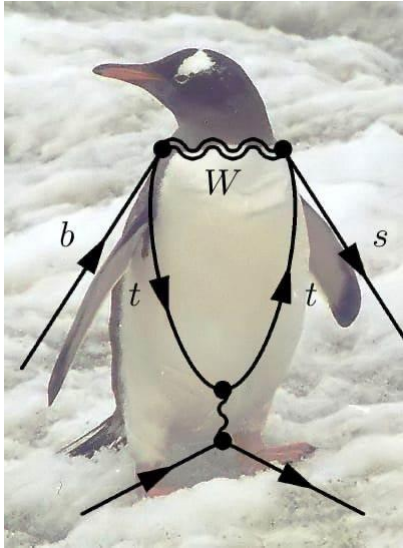
树
Shù



Neutral currents

Pseudo-scalar and
scalar NP encoded
in $C_{P,S}$

Penguin diagrams



企鵝 - Qì'é

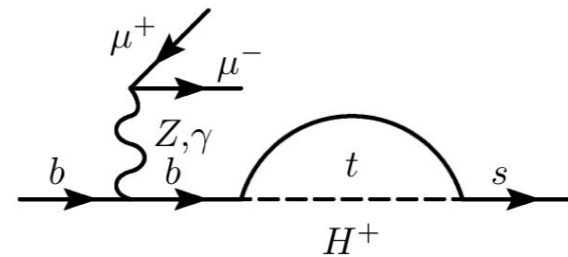
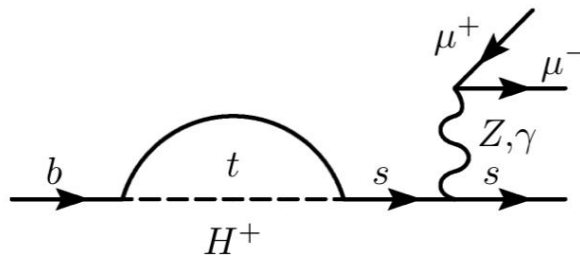


Figure 2: Penguin diagrams at one loop level for $b \rightarrow s\mu^+\mu^-$ transitions.

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The code: GAMBIT

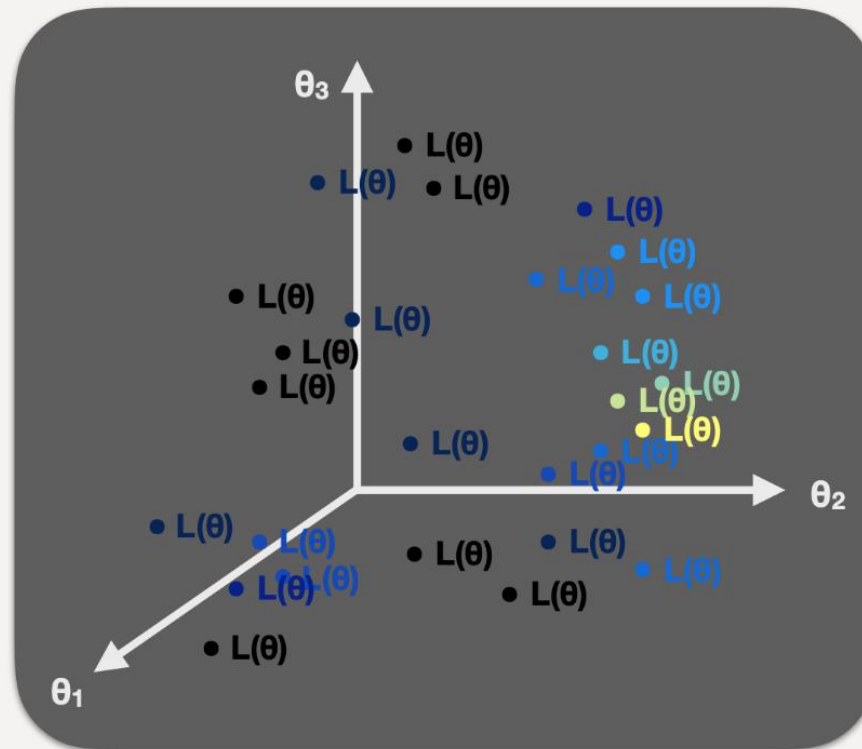
The **G**lobal **A**nd **M**odular **B**SM Inference **T**ool

- Open-source code in **C++** to calculate observables and *likelihoods* for generic Beyond the Standard Model(s) theories.
- Modular: *modules* provide **GAMBIT** with a range of functions (capabilities) to calculate a certain quantity.
- **GAMBIT** samples the parameter space by calling the necessary modules and backend functions for each parameter point, e.g., performing a global fit.



Likelihood functions and Global fits

- Explore the model parameter space $(\theta_1, \theta_2, \theta_3, \dots)$
- At every point θ : calculate $\text{predictions}(\theta) \rightarrow$ evaluate joint likelihood $L(\theta)$



- Region of highest $L(\theta)$ or $\ln L(\theta)$: **model's best simultaneous fit to all data**
(but not necessarily a *good* fit, or the most probable θ ...)

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Scans

CP conserving potential close to the alignment limit

**LUMI supercomputer (Large Unified Modern Infrastructure)
in Kajaani (LUMI also means snow in Finish.). 5th most powerful
in the world.**

Scanner:

```
use_scanner: de
```

scanners:

```
de:
```

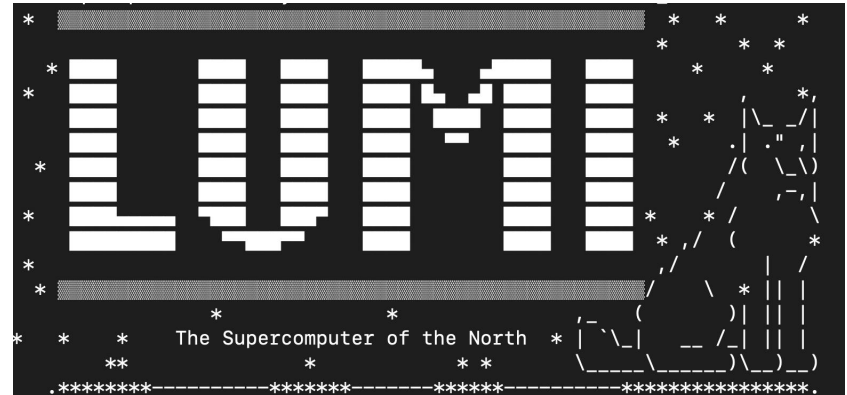
```
plugin: diver
```

```
like: LogLike
```

```
NP: 20000
```

```
convthresh: 1e-6
```

```
verbosity: 1
```

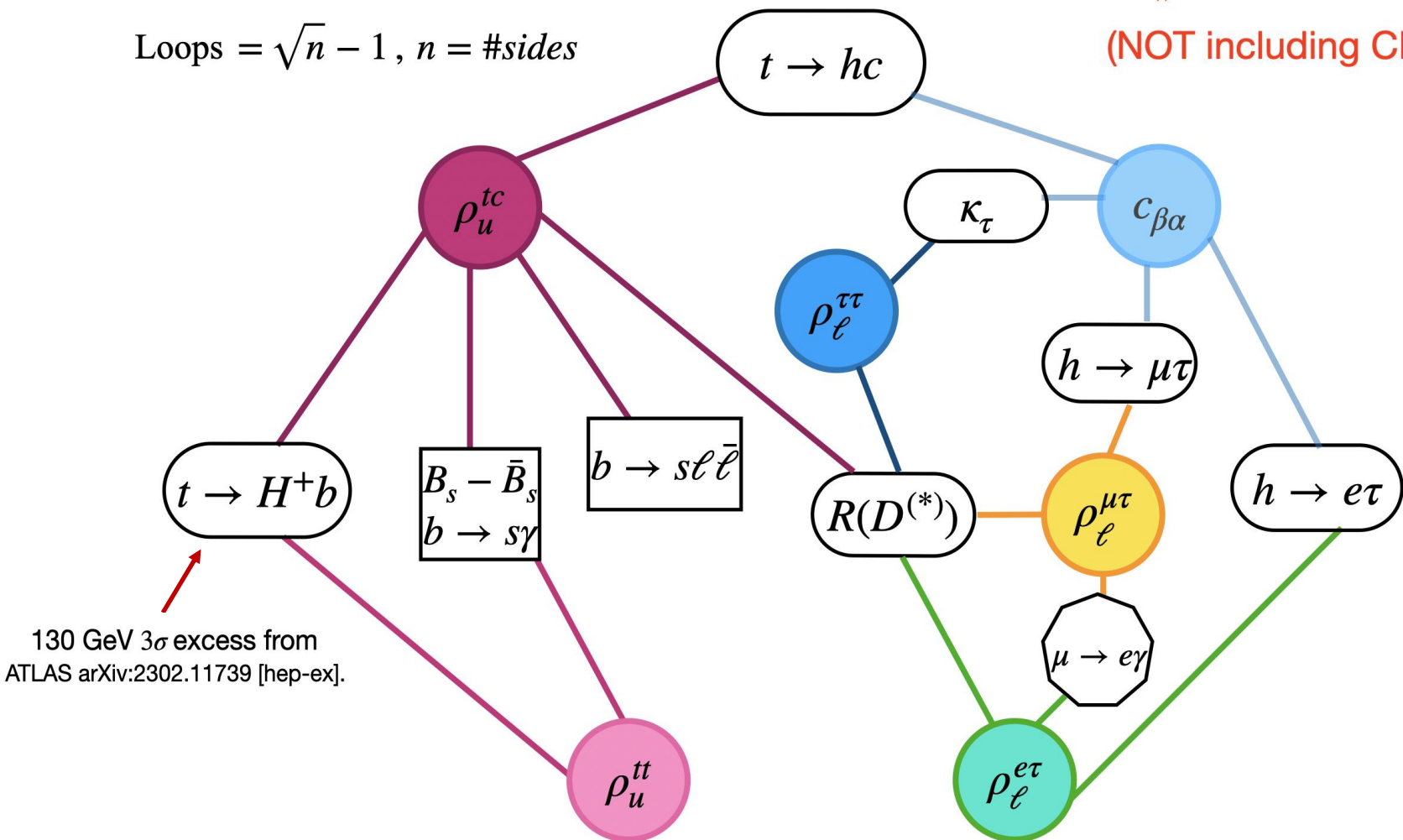


(~22 hours/scan with 512 cores on
the small partition)
So far I have used more than 300 hours.

Observables and couplings

m_W from PDG 2024
(NOT including CDF-II)

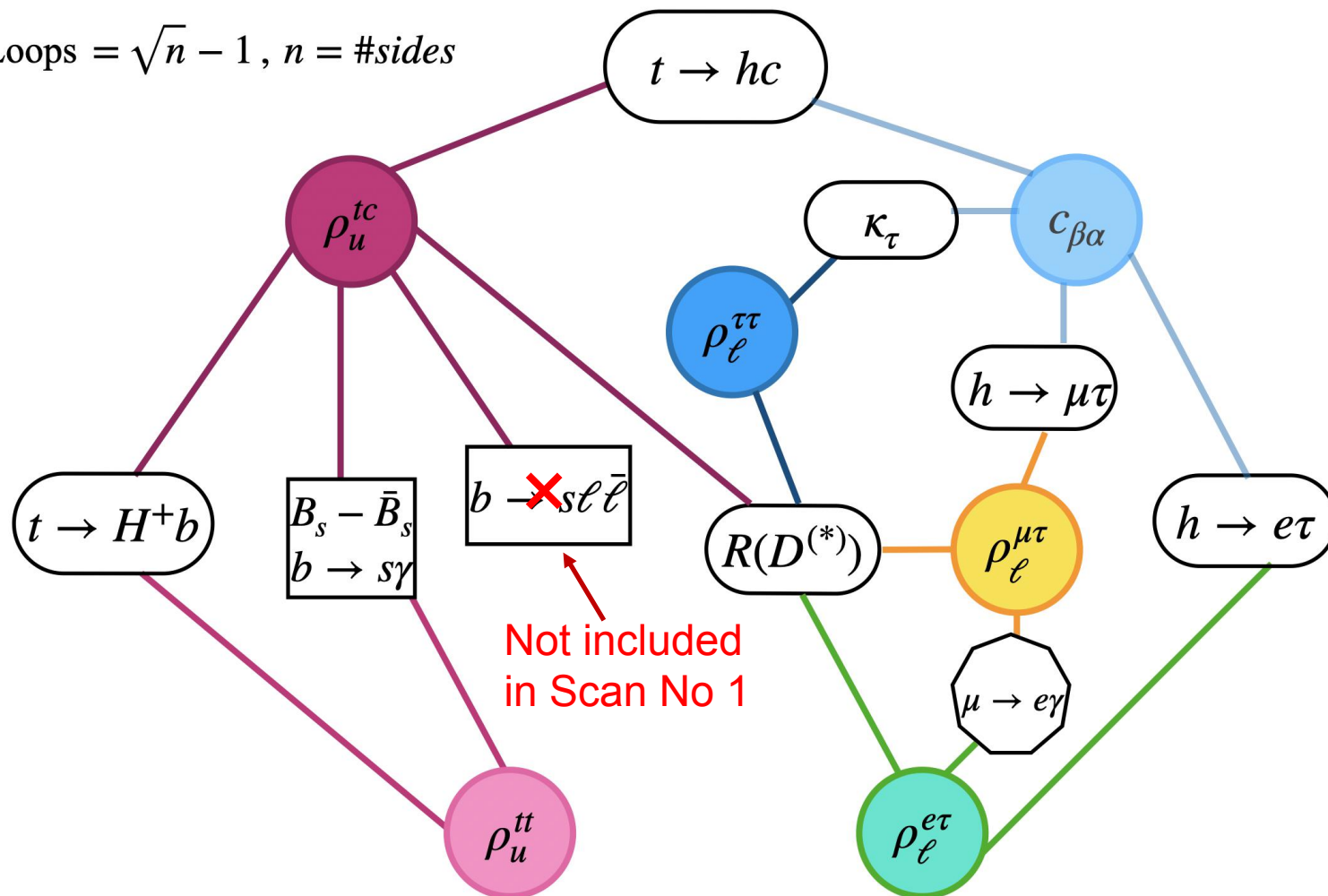
Loops = $\sqrt{n} - 1$, $n = \#sides$



Adapted from Phys.Rev.D 110 (2024) 1, 015014 • e-Print: 2311.03430 [hep-ph]

Scan No 1

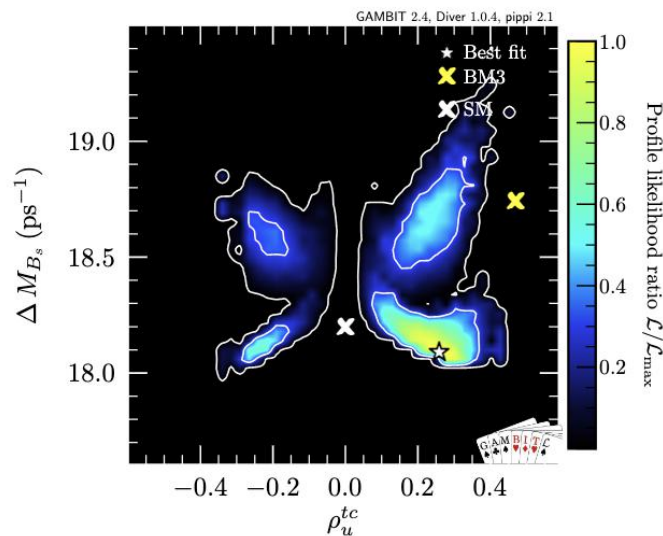
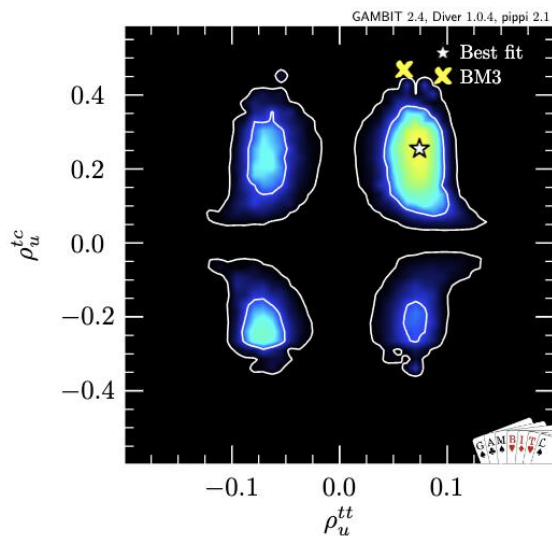
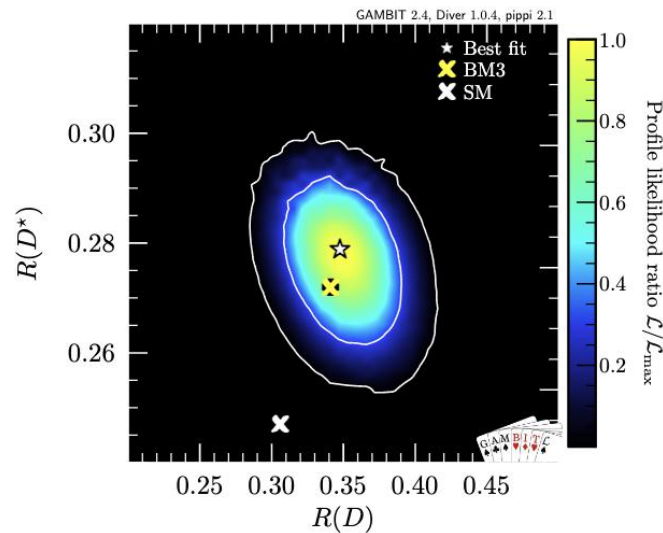
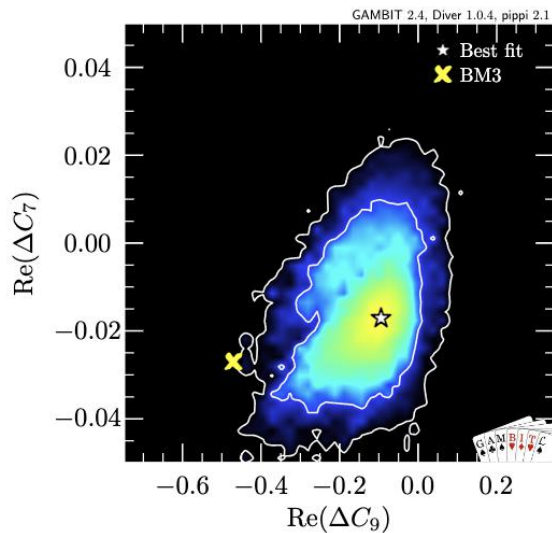
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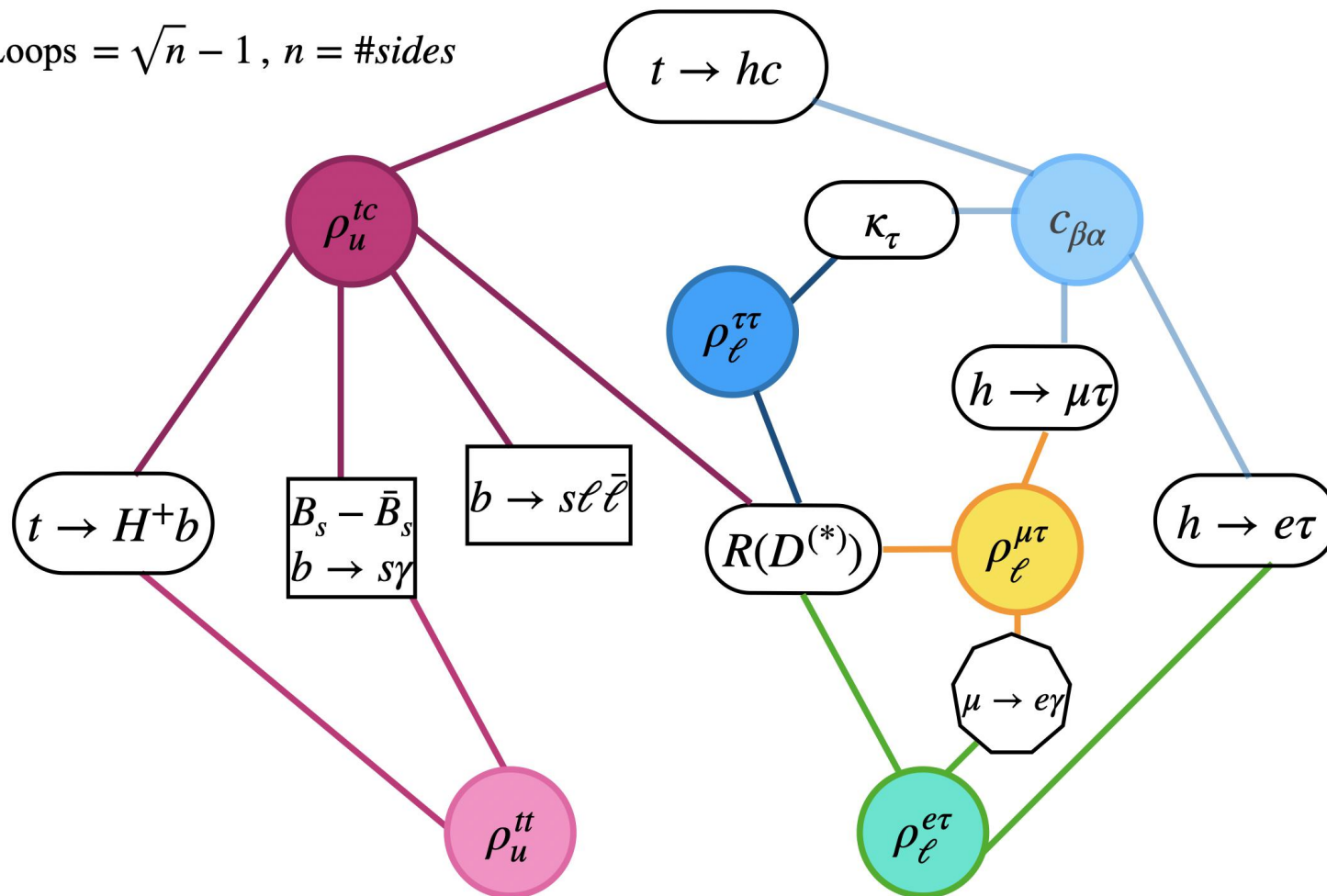
Scan No 1

Parameter space and best fit values



Scan No 2

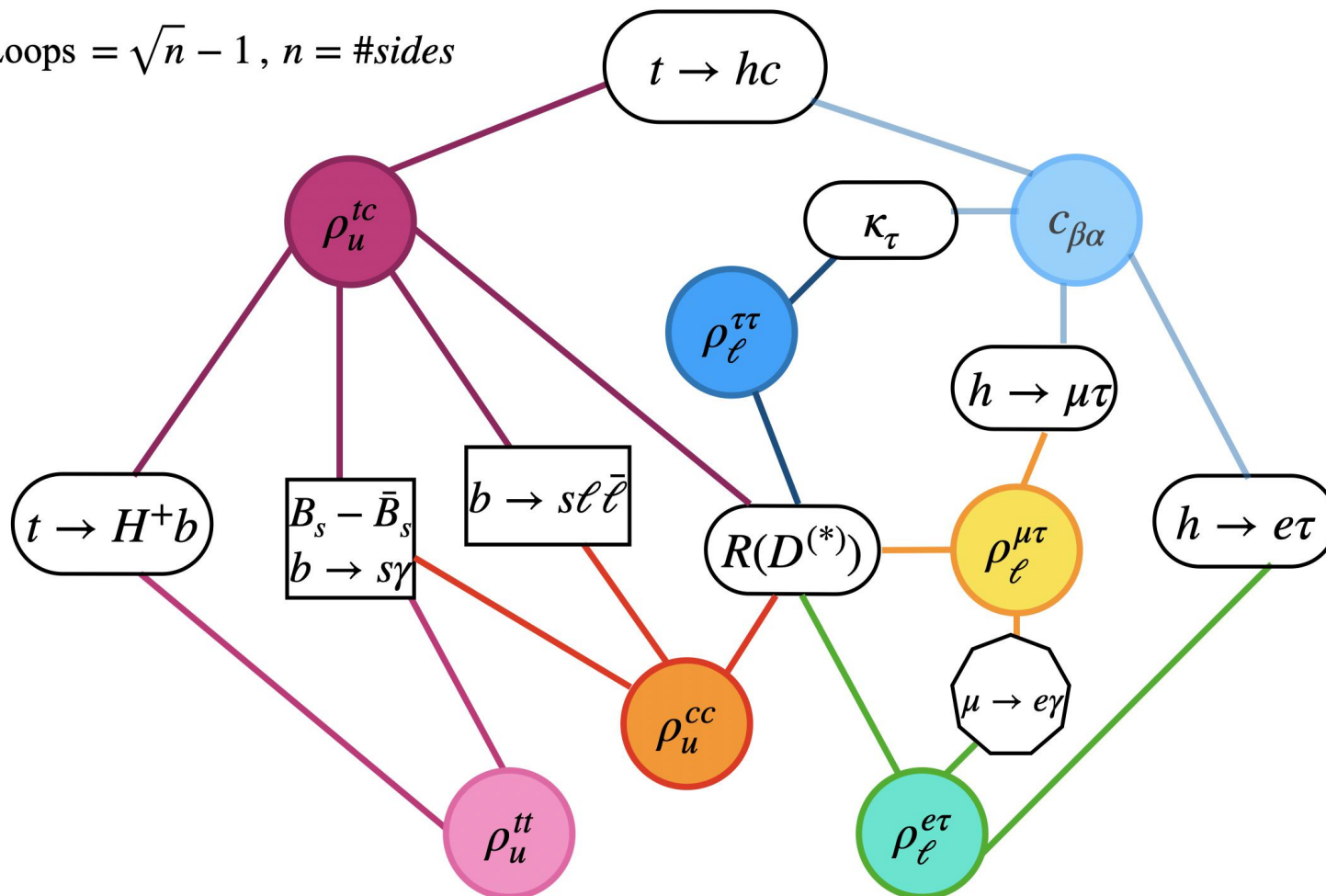
Loops = $\sqrt{n} - 1$, $n = \#sides$



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Scan No 2

Loops = $\sqrt{n} - 1$, $n = \#sides$

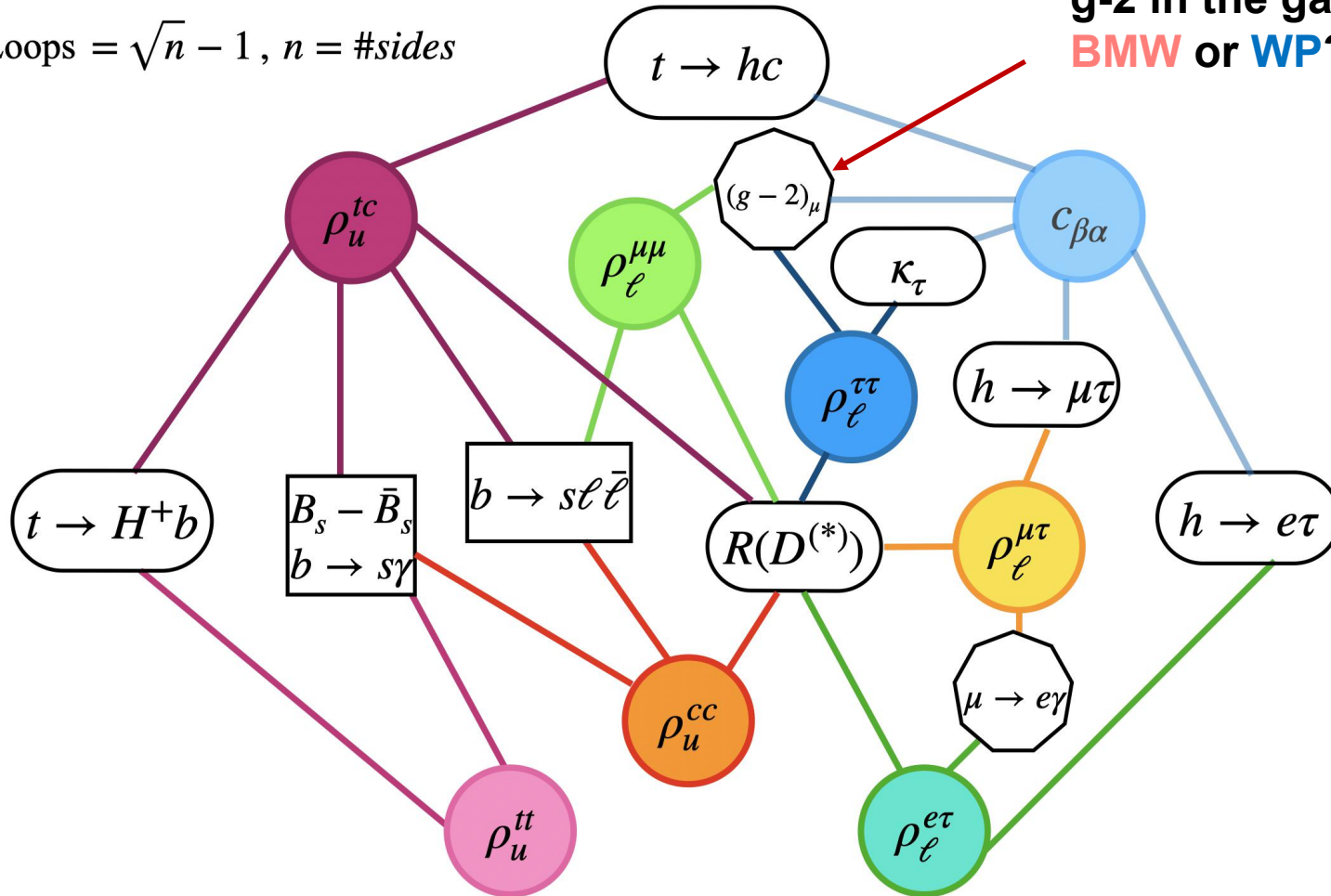


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Loops = $\sqrt{n} - 1$, $n = \#sides$

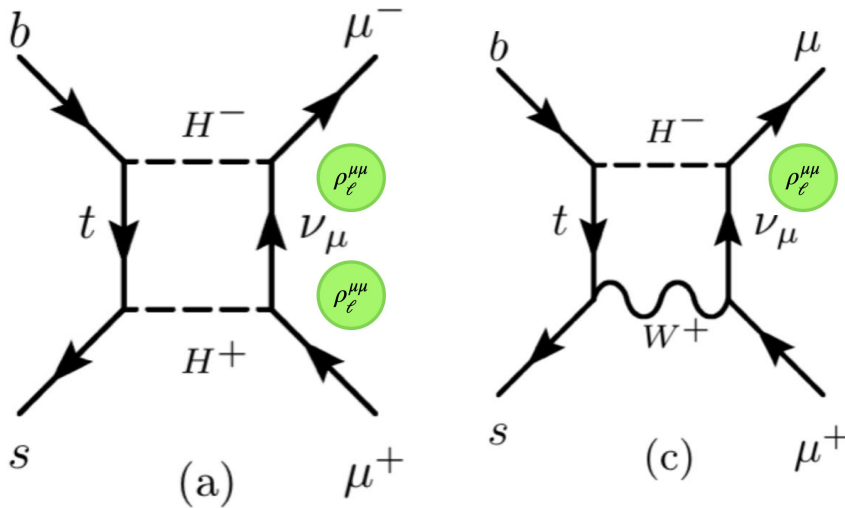
g-2 in the game,
BMW or WP?



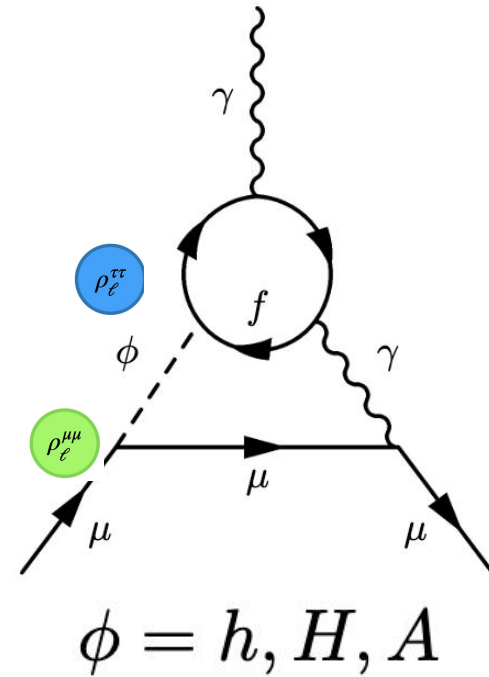
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New diagrams!

Box diagrams



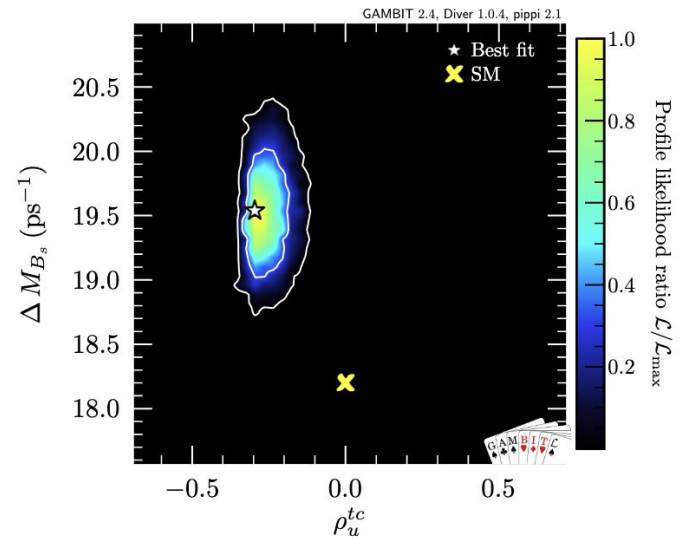
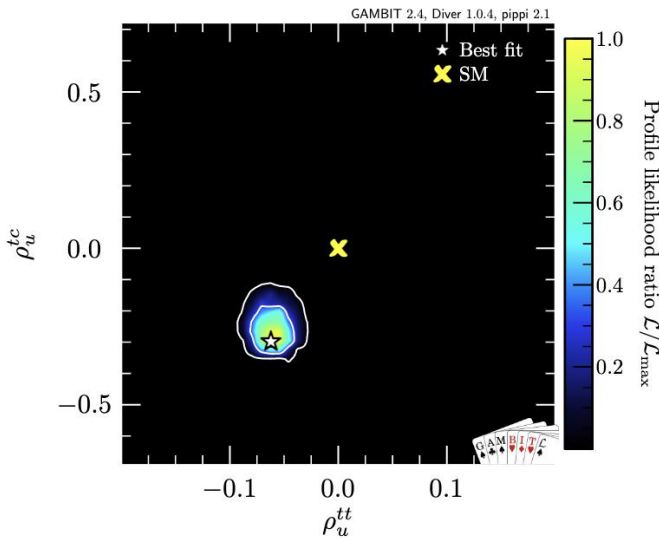
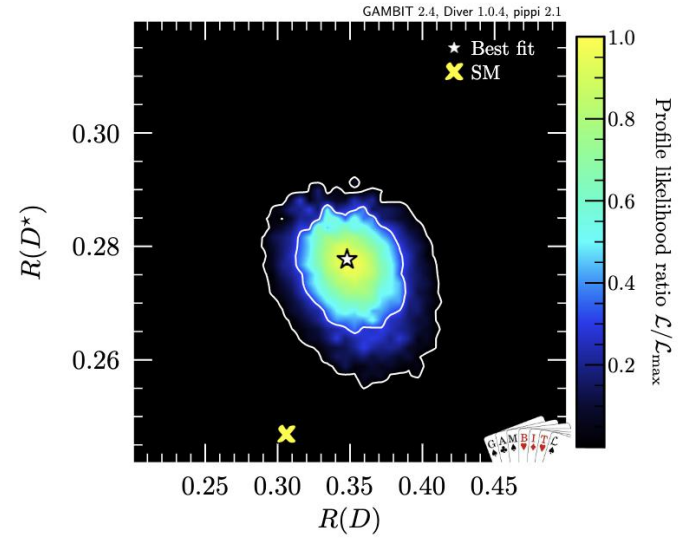
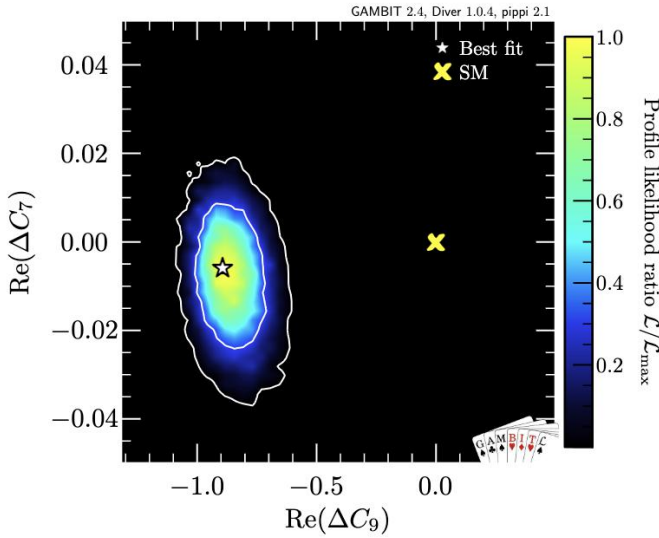
Barr-Zee diagrams
in muon g-2



Scan No 2

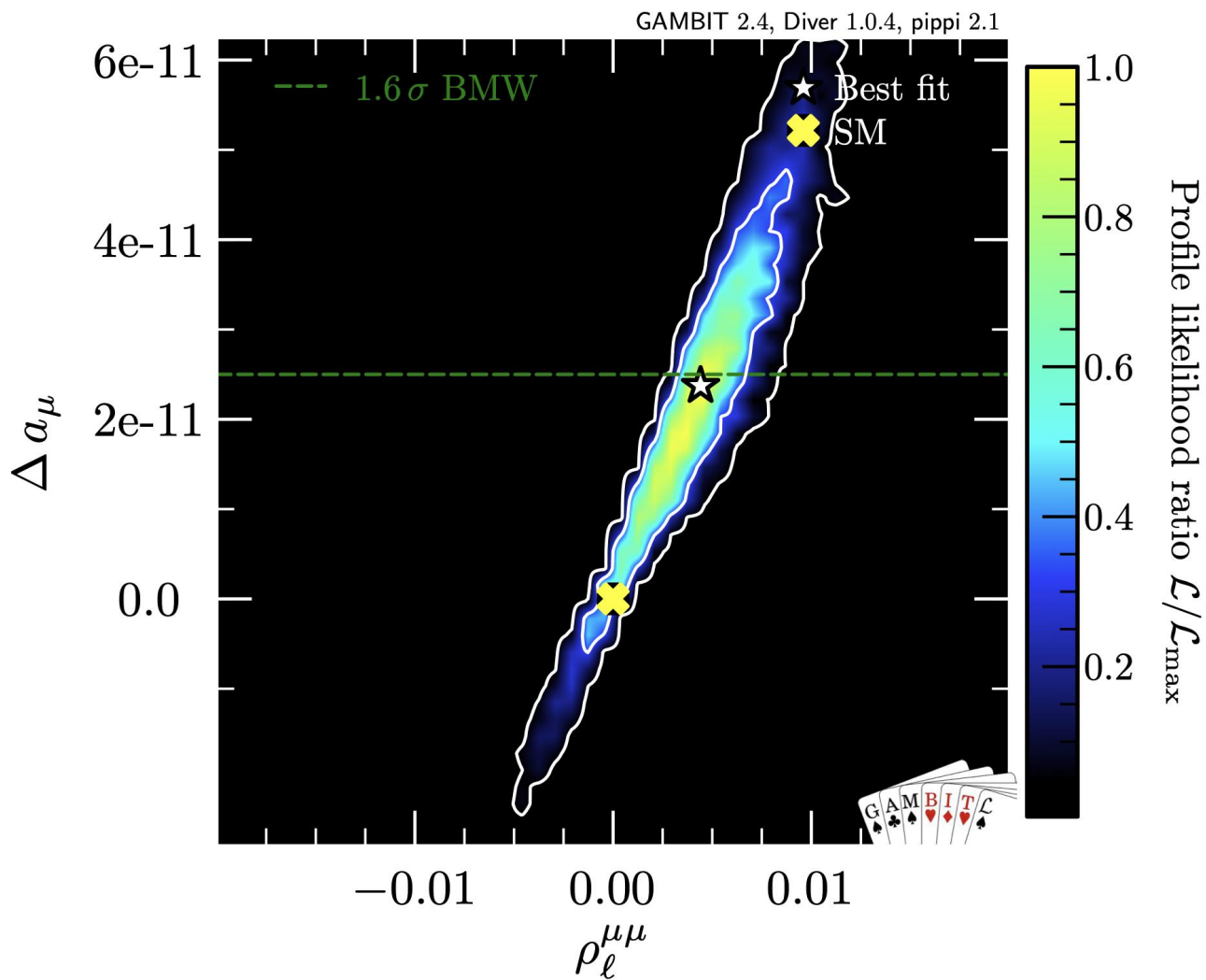
g-2 in the game,
BMW

Parameter space and best fit values



Scan No 2

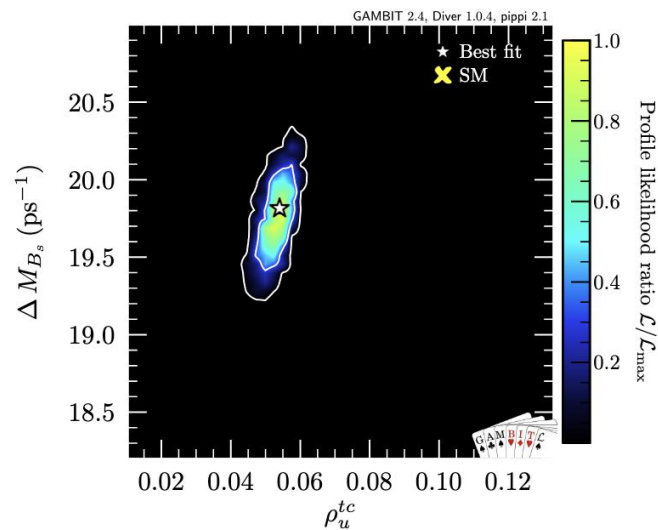
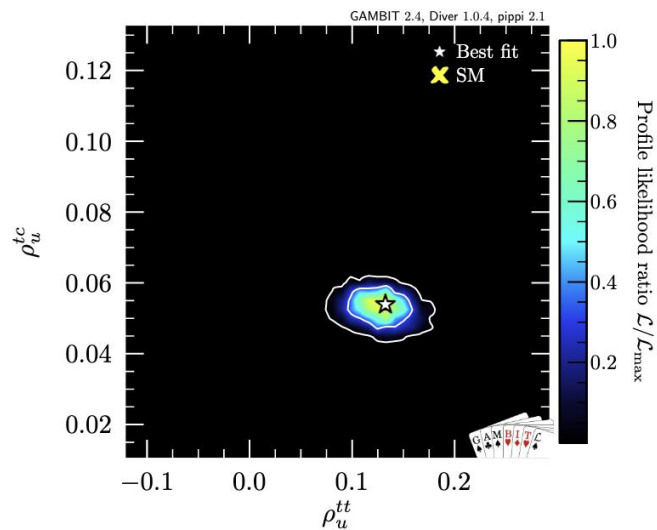
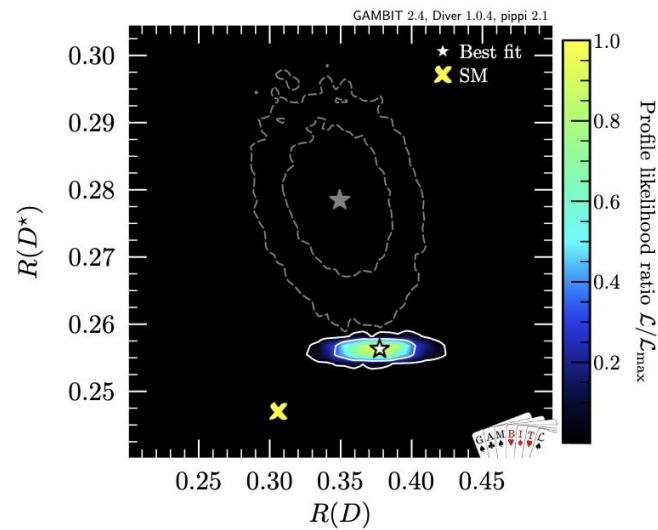
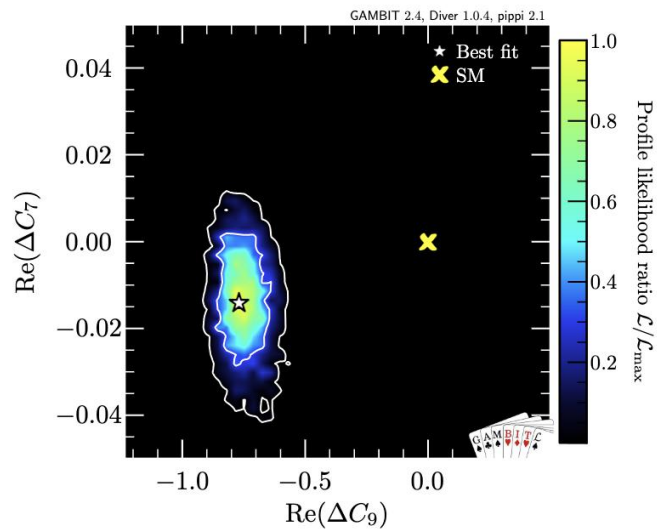
g-2 in the game,
BMW



Scan No 3

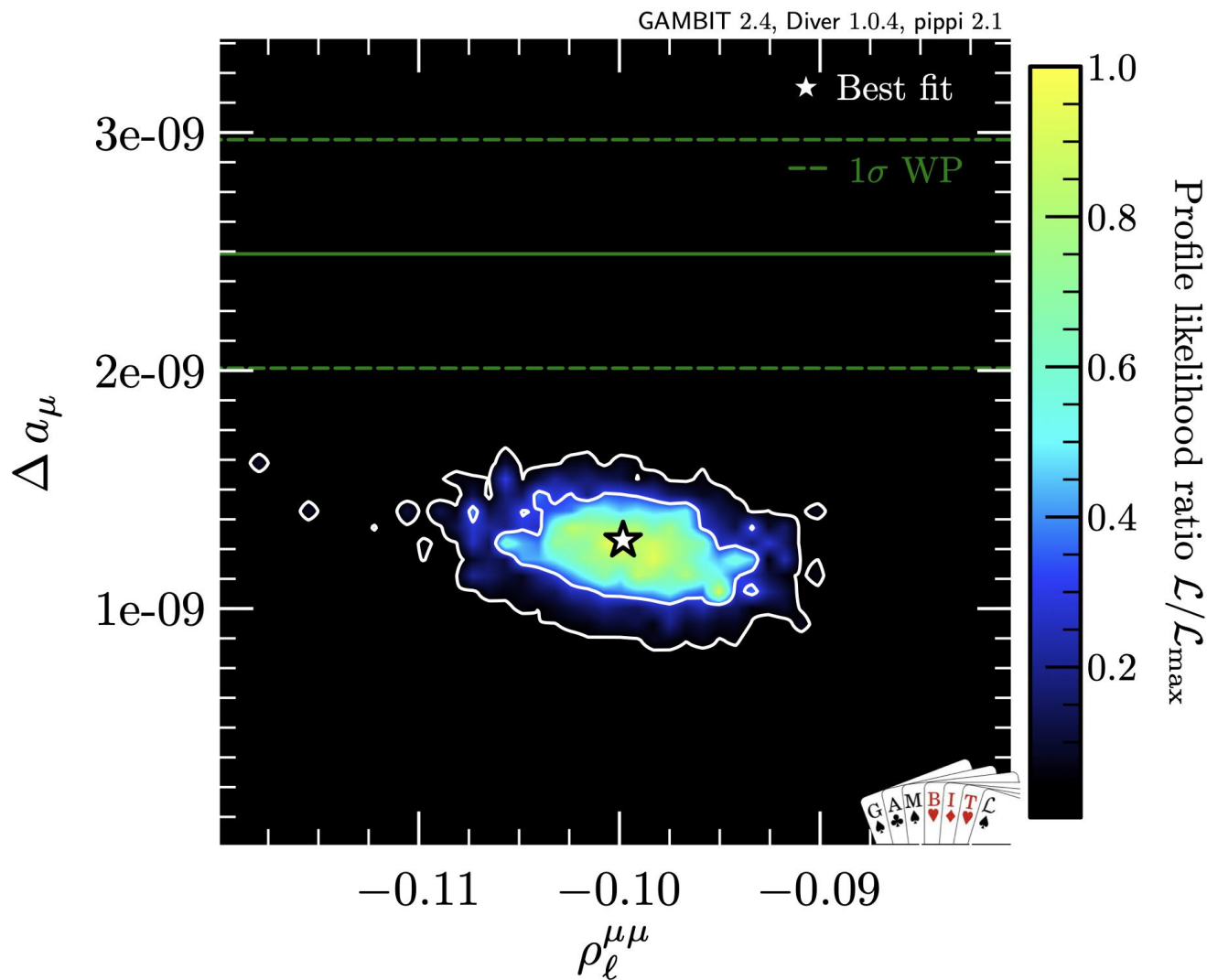
g-2 in the game,
WP

Parameter space and best fit values



Scan No 2

g-2 in the game,
WP



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Summary

- I presented a likelihood analysis for the G2HDM including both the charged and neutral anomalies along other flavour observables.
- We found that the model can explain the neutral anomalies at the 1 sigma level at the same time that the BMW muon $g-2$ value and the PDG 2024 data for the m_W mass (not including CDF-II data).
- The model will require small $b-s$ flavour violation at tree level in order to explain B_s - B_s mixing.
- When using PDF 2024 data, the model can explain W_P value at the 2 sigma level although large charm-charm extra Yukawas are needed and the charged anomalies can not be explained.
- Next step is to study what happens when using CDF-II data.

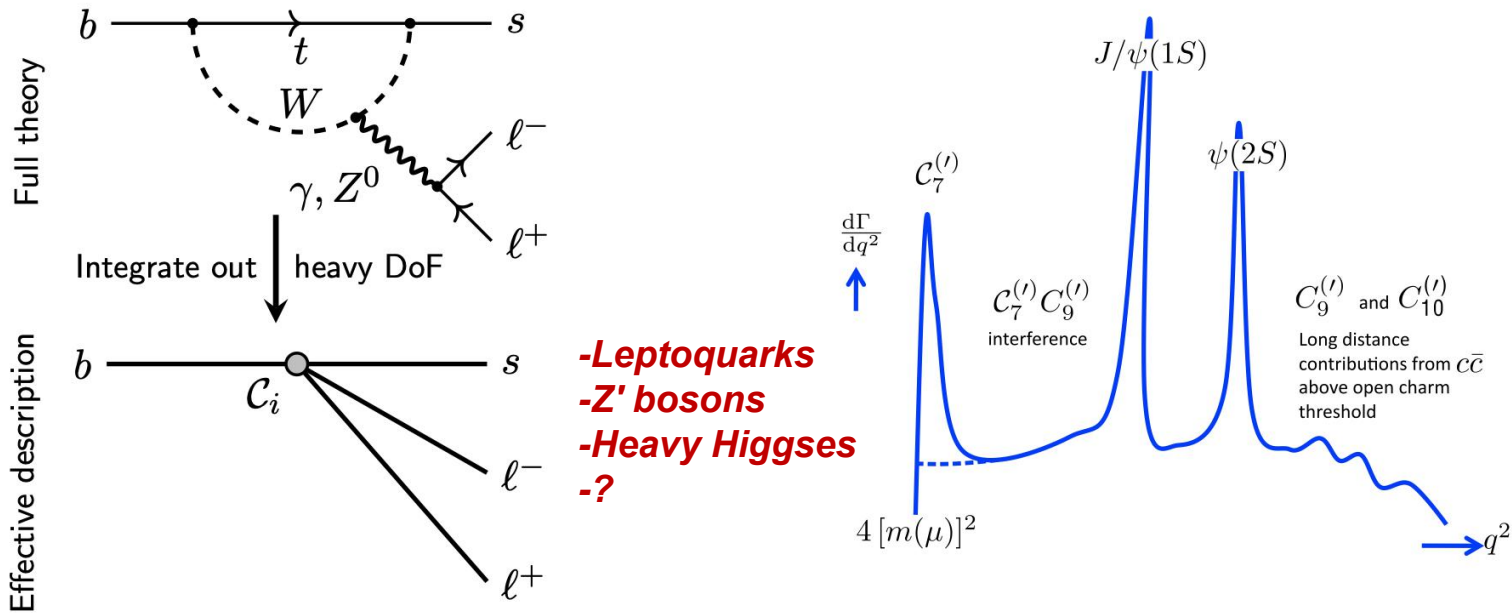
Thanks!

(谢谢您!)



Backup slides

How to explain them?: *New Physics Story*



- $b \rightarrow sll$ transitions described model-independently in effective theory

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i \underbrace{C_i}_{\text{Wilson coefficient ("effective coupling")}} \underbrace{\mathcal{O}_i}_{\text{Local operator}}$$

Effective couplings in $b \rightarrow sll$ transitions		
Wilson coefficient		Operator
γ -penguin	$C_7^{(l)}$	$\frac{e}{g^2} m_b (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu}$
ew. penguin	$C_9^{(l)}$	$\frac{e^2}{g^2} (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\mu} \gamma^\mu \mu)$
	$C_{10}^{(l)}$	$\frac{e^2}{g^2} (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\mu} \gamma^\mu \gamma_5 \mu)$
scalar	$C_S^{(l)}$	$\frac{e^2}{16\pi^2} m_b (\bar{s} P_{R(L)} b) (\bar{\mu} \mu)$
pseudoscalar	$C_P^{(l)}$	$\frac{e^2}{16\pi^2} m_b (\bar{s} P_{R(L)} b) (\bar{\mu} \gamma_5 \mu)$

For completeness

- Different $q^2 = m^2(\ell^+ \ell^-)$ regions probe different operator combinations

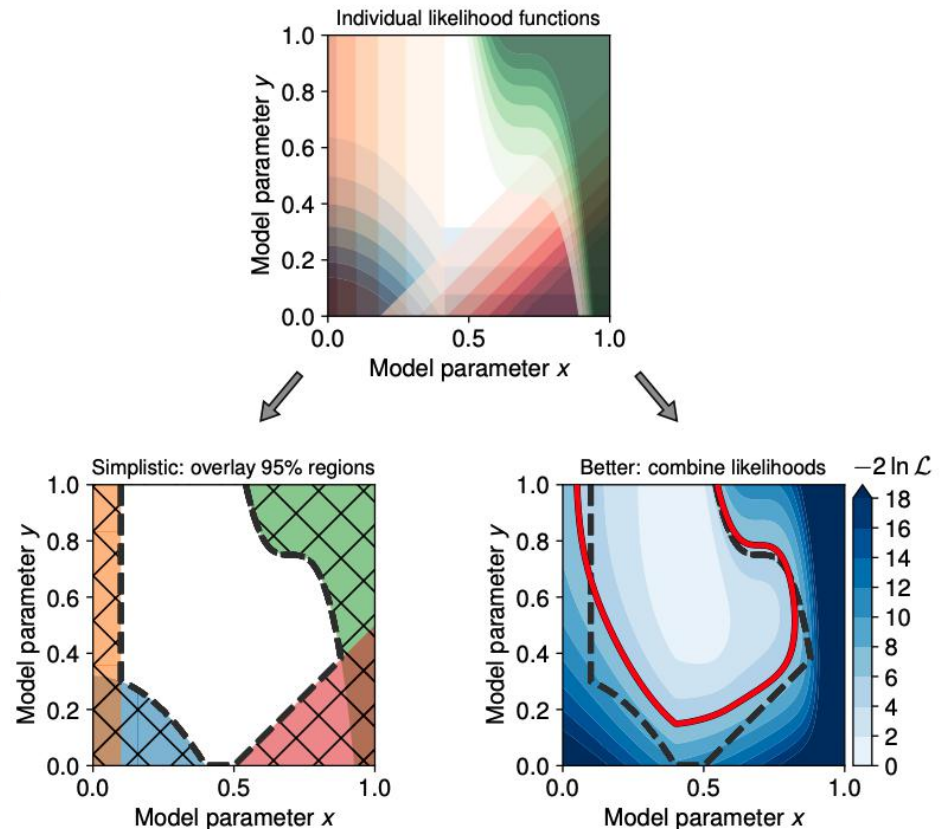
Likelihood functions and Global fits

- In general, we have several likelihood functions from different observables: Combine all constraints into a composite likelihood,

$$\mathcal{L} = \mathcal{L}_{Flavour} \mathcal{L}_{Higgs} \mathcal{L}_{Collider} \dots$$

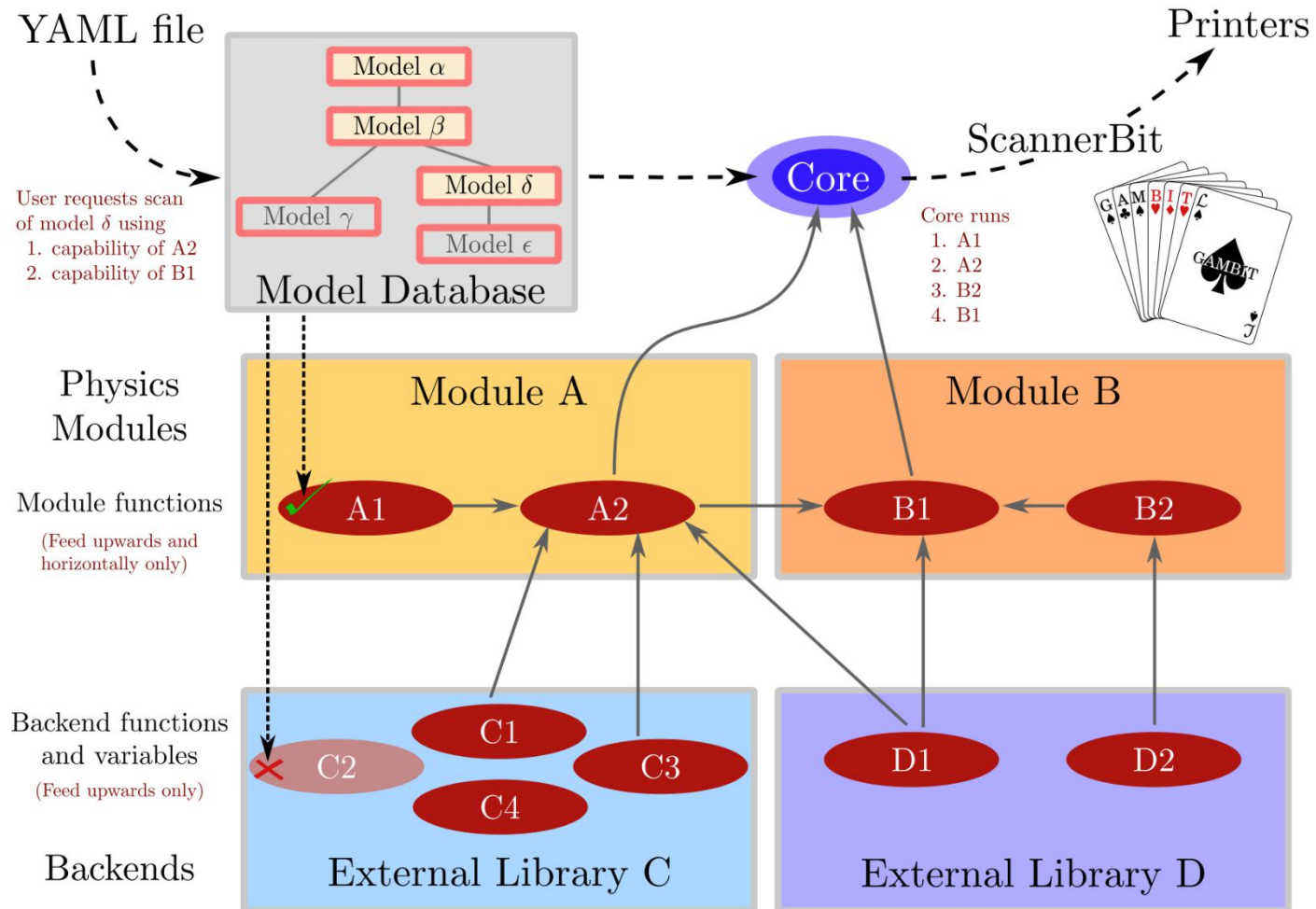
- Perform an extensive parameter scan with rigorous statistical interpretation (frequentist/Bayesian):
 - Parameter estimation.
 - Model comparison.

[GAMBIT Community, [arXiv:2012.09874 \[hep-ph\]](https://arxiv.org/abs/2012.09874)]



GAMBIT: The Global And Modular BSM Inference Tool

gambit.hepforge.org



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7/19/2024

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Likelihood functions

Probability Distribution Functions (PDFs)

Statistical Model: PDFs for obtaining observations x given a set of params θ . $\rightarrow p(x|\theta)$

Experiments provide observations of x which are used for inferences about components of θ .

Likelihood: We can compute theory predictions $x^{th}(\theta)$ so that

$$p(x^{exp} | \theta) = p(x^{exp}, x^{th}(\theta)).$$

Evaluate the PDF only for the specific x^{exp} that was observed, and examine how it varies with θ

$$\mathcal{L}(\theta) = p(x^{exp}, x^{th}(\theta))$$

Yukawa textures

- The number of free parameters from the Yukawas will be 54 (3x3x3x2).

$$\xi^u = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \xi_{cc}^u & \xi_{ct}^u \\ 0 & \xi_{tc}^u & \xi_{tt}^u \end{pmatrix}, \quad \xi^d = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \xi_{ss}^d & \xi_{sb}^d \\ 0 & \xi_{bs}^d & \xi_{bb}^d \end{pmatrix}, \quad \xi^l = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \xi_{\mu\mu}^l & \xi_{\mu\tau}^l \\ 0 & \xi_{\tau\mu}^l & \xi_{\tau\tau}^l \end{pmatrix},$$

Effective Hamiltonian for $b \rightarrow s\mu^+\mu^-$

$$\mathcal{H}_{\text{eff}}^{NP} = -\frac{4G_F}{\sqrt{2}}V_{tb}V_{ts}^* \left[\sum_{i=S,P} C_i^{(0)} \mathcal{O}_i + C_i^{\prime(0)} \mathcal{O}'_i + \sum_{i=7,9,10} C_i^{(1)} \mathcal{O}_i + C_i^{\prime(1)} \mathcal{O}'_i \right]$$

$$\mathcal{O}_9 = \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \ell),$$

$$\mathcal{O}_{10} = \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell),$$

$$\mathcal{O}_S = \frac{e^2}{16\pi^2} m_b (\bar{s} P_R b) (\bar{\ell} \ell),$$

$$\mathcal{O}_P = \frac{e^2}{16\pi^2} m_b (\bar{s} P_R b) (\bar{\ell} \gamma_5 \ell),$$

$$\mathcal{O}_7 = \frac{e}{16\pi^2} m_b (\bar{s} \sigma^{\mu\nu} P_R b) F_{\mu\nu},$$

and prime operators from $P_R \rightarrow P_L$.

Fermions + second Higgs doublet

NEW PARTICLES!

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + H_1 + i\eta_1) \end{pmatrix}$$

$$\Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + H_2 + i\eta_2) \end{pmatrix}$$

h	A	H	H^\pm
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$$\tan \beta = \frac{v_2}{v_1}$$

Mixing parameters

$$\sin(\beta - \alpha)$$

Results

Branching ratios, arXiv: 2111.10464

