

ICHEP 2024



Search for LFV with the Z'
model in future lepton colliders

arXiv:2302.02203

JHEP 03,190 (2023)

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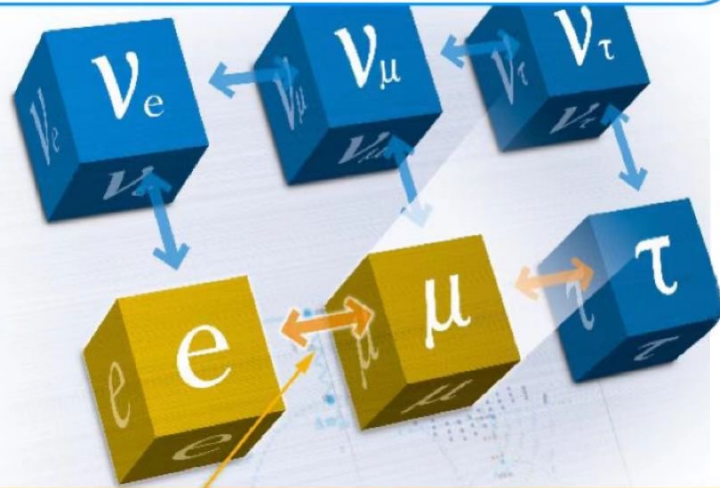
2024-7-19



- ◆ Motivation
- ◆ Z' model
- ◆ Future collider
- ◆ Event selection
- ◆ Analysis framework
- ◆ Upper limits and coupling
- ◆ Summary

Motivation

Neutrino Flavor Violation is observed !



charged Lepton Flavor Violation !? (cLFV)

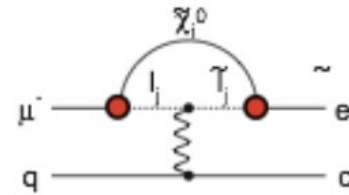
- ◆ Since LFV decay is forbidden in the SM, the observation of any LFV decay would be a signal of new physics beyond SM.
- ◆ In SM, Lepton Flavour is conserved for zero degenerate ν masses and now we have clear indication that ν s have finite mass.

◆ Models may enhance LFV effects up to a detectable level, such as leptoquark, Compositeness, Supersymmetry, Heavy Z' and Anomalous boson Coupling model.

Eur. Phys. J. C57:13-182, 2008

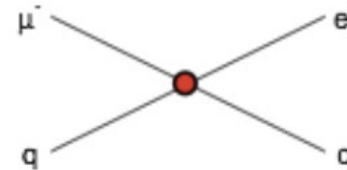
Supersymmetry

rate $\sim 10^{-15}$



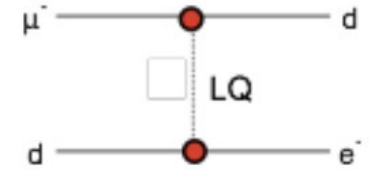
Compositeness

$\Lambda_c \sim 3000$ TeV



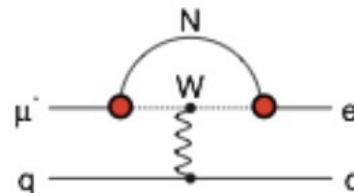
Leptoquark

$M_{LQ} = 3000 (\lambda_{\mu d} \lambda_{e d})^{1/2} \text{ TeV}/c^2$



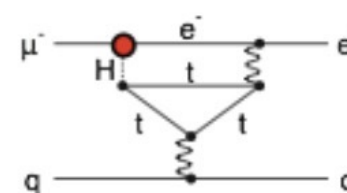
Heavy Neutrinos

$|U_{\mu N} U_{e N}|^2 \sim 8 \times 10^{-13}$



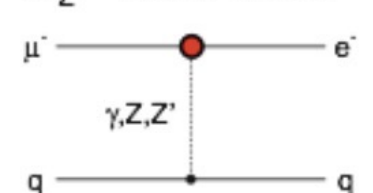
Second Higgs Doublet

$g(H_{\mu e}) \sim 10^{-4} g(H_{\mu\mu})$



Heavy Z' Anomal. Z Coupling

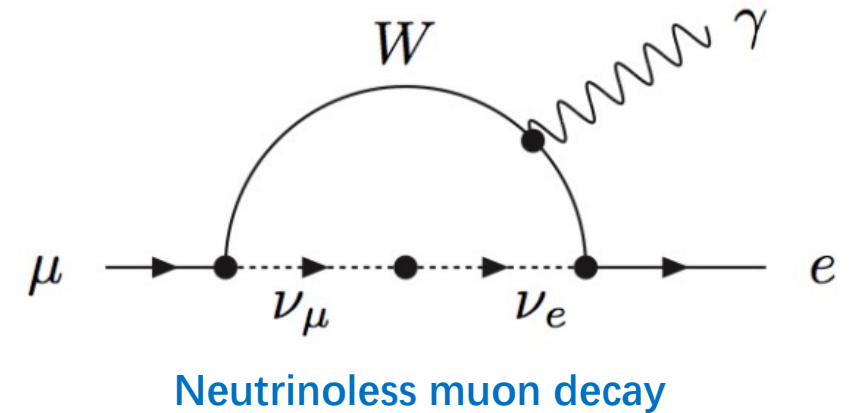
$M_{Z'} = 3000 \text{ TeV}/c^2$



- ◆ In the charged lepton sector, LFV is heavily suppressed in the Standard Model.

$$BR(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_{i=2,3} U_{\mu i}^* U_{ei} \frac{\Delta m_{1i}^2}{M_W^2} \right|^2 < 10^{-54}$$

Nucl. Phys. 25, 340 (1977)

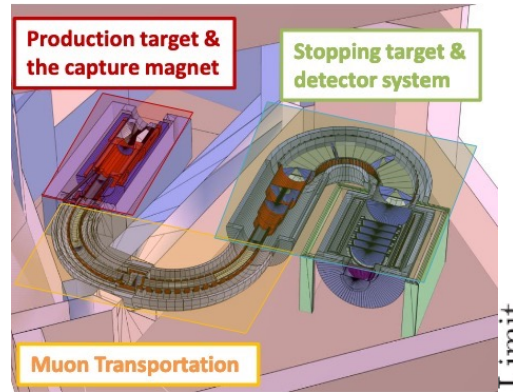


- ◆ Both experimental searches and upper-limit predictions, including μ , τ LFV decays, π , K LFV decays and ϕ , J/ψ two-body LFV decays, etc.

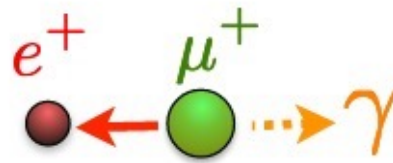
CLFV

- ◆ $B(\mu^+ \rightarrow e^+ \gamma) < 3.1 \times 10^{-13}$ @ 90% C.L. **MEGII**
- ◆ $B(\tau^+ \rightarrow e^+ \gamma) < 3.3 \times 10^{-8}$ @ 90% C.L. **BABAR**
- ◆ $B(\mu \rightarrow 3e) < 1.0 \times 10^{-12}$ @ 90% C.L. **SINDRUM**
- ◆ $B(Z \rightarrow e^\pm \mu^\mp) < 7.5 \times 10^{-7}$ @ 95% C.L. **ATLAS**
- ◆ $B(\phi \rightarrow e^\pm \mu^\mp) < 2 \times 10^{-6}$ @ 90% C.L. **SND**
- ◆ $B(J/\psi \rightarrow e^\pm \tau^\mp) < 7.1 \times 10^{-8}$ @ 90% C.L. **BESIII**
- ◆ $B(J/\psi \rightarrow e^\pm \mu^\mp) < 4.5 \times 10^{-9}$ @ 90% C.L. **BESIII**

Current best limit



COMET



MEGII process

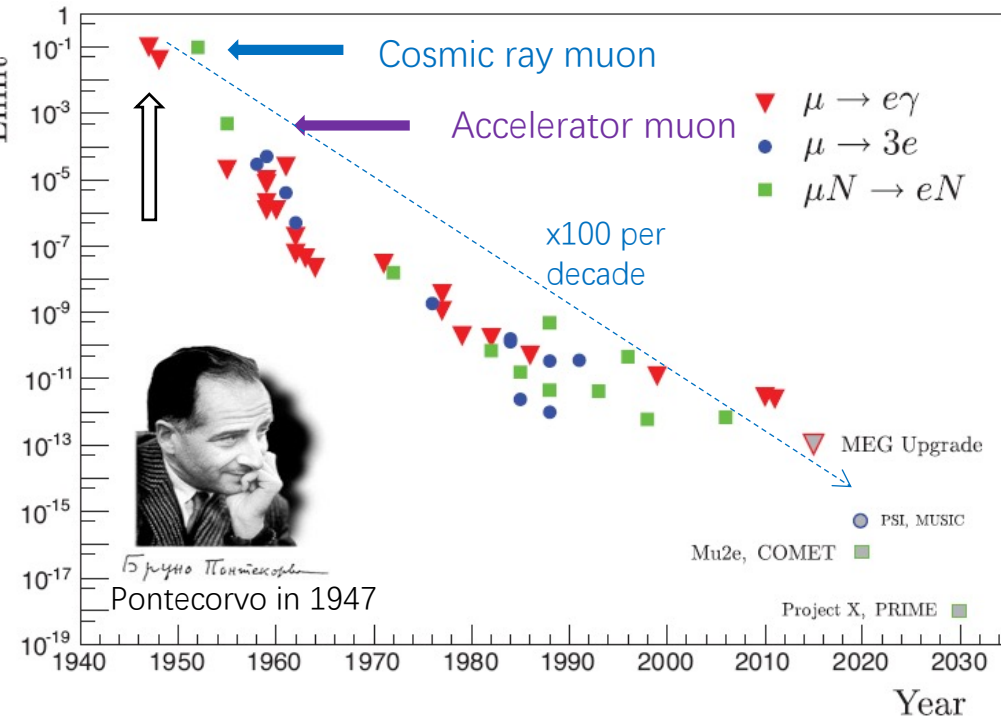
- ◆ **Mu2e** and **COMET** will search for CLFV with $\mu N \rightarrow e N$

Improve the current limit by a factor of 10^4

Next goal $< 6 \times 10^{-17}$ (90% C.L.)

Search for New Physics with energy scale up to 10^4 TeV

- ◆ **MEGII** and **Mu3e** has similar beam requirements. Intensity $O(10^8 \text{ muon/s})$, low momentum $p = 28 \text{ MeV/c}$. MEGII was started in 2021 and will continue to run until 2026 aiming at a sensitivity down to 6×10^{-14} (90% C.L.)



Eur. Phys. J. C 84, 216 (2024)
 Phys. Rev. Lett. 104, 021802 (2010)
 Nucl. Phys. B 299, 1 (1988).
 Phys. Rev. D 90, 072010 (2014)
 Phys. Rev. D 81, 057102 (2010)
 Phys. Lett. B 598, 172 (2004)
 Phys. Rev. D 103, 112007 (2021)
 Sci. Chin. Mech. Astron. 66 2 (2023)



- ◆ A new $U(1)$ gauge symmetry \longrightarrow Z'
- ◆ Z' , a **neutral vector boson** with the same couplings to fermion-antifermion as the Z , but with a larger mass.
- ◆ May interact with different particles and produce different decay modes \longrightarrow New physics
- ◆ May benefit from the development of new technologies, such as higher-energy particle accelerators and more sensitive detectors.
- ◆ Searching for Z' \longleftarrow LHC, FNAL...

f	$\Gamma_{f\bar{f}}$
ℓ	$\frac{\alpha M_{Z'}}{24s_W^2 c_W^2} (1 - 4s_W^2 + 8s_W^4)$
ν	$\frac{\alpha M_{Z'}}{24s_W^2 c_W^2}$
u	$\frac{3\alpha M_{Z'}}{24s_W^2 c_W^2} (1 - \frac{8}{3}s_W^2 + \frac{32}{9}s_W^4)$
d	$\frac{3\alpha M_{Z'}}{24s_W^2 c_W^2} (1 - \frac{4}{3}s_W^2 + \frac{8}{9}s_W^4)$

α : the fine structure constant

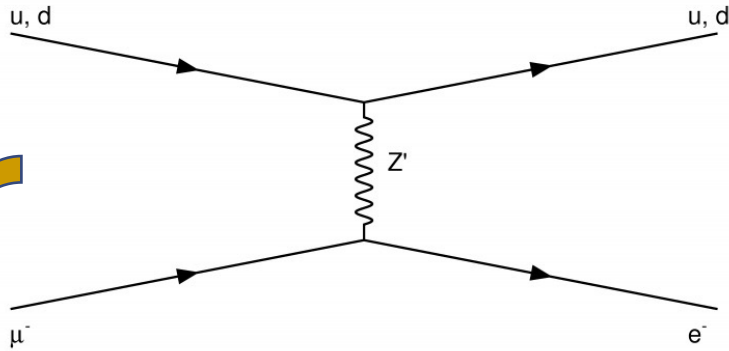
$M_{Z'}$: the mass of Z'

c_W : the cosine of the weak mixing angle

s_W : the sine of the weak mixing angle



◆ μ to e conversion



Z' - e - μ coupling

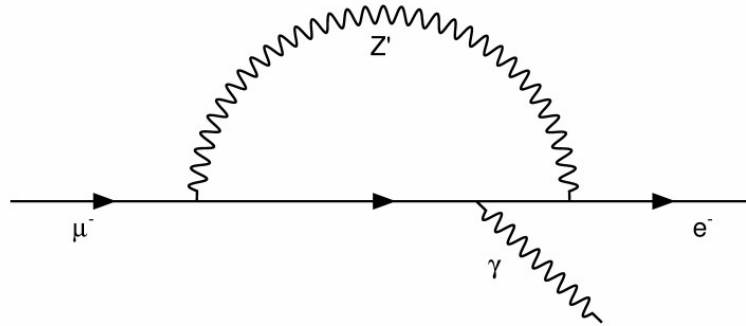
$$\lambda_{e\mu}^2 = \frac{2\pi^2 \Gamma_{capture} Z R}{G_F^2 \alpha^3 m_\mu^5 Z_{eff}^4 |F(q)|^2} \cdot \frac{M_{Z'}^4}{M_Z^4} \times \frac{1}{S_W^4 + \left(S_W^2 - \frac{1}{2}\right)^2} \times \frac{1}{\left[(2Z+N)\left(\frac{1}{2} - \frac{4}{3}S_W^2\right) + (Z+2N)\left(-\frac{1}{2} + \frac{2}{3}S_W^2\right)\right]^2}$$

G_F : the Fermi constant, α : the fine structure constant, $\Gamma_{capture}$: nuclear muon capture rate,

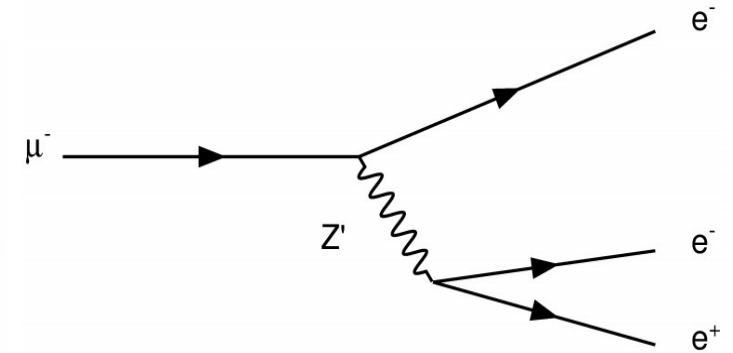
Z_{eff}, F_p : nuclear parameters, Z : the atomic number, N : the number of neutrons in the nucleus,

S_W : the sine of the weak mixing angle, $M_{Z'}$: the mass of Z' , m_μ : the muon mass

◆ $\mu^- \rightarrow e^- \gamma$

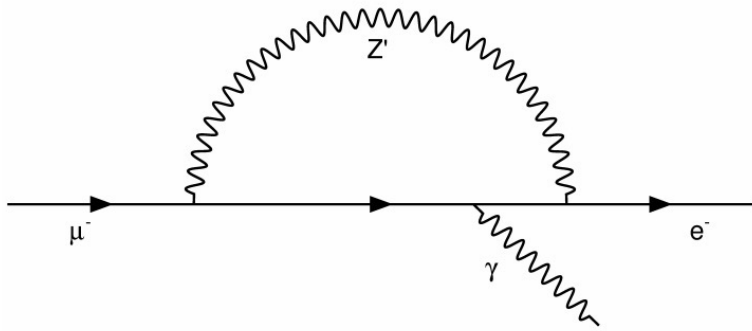


◆ $\mu^- \rightarrow e^- e^- e^+$





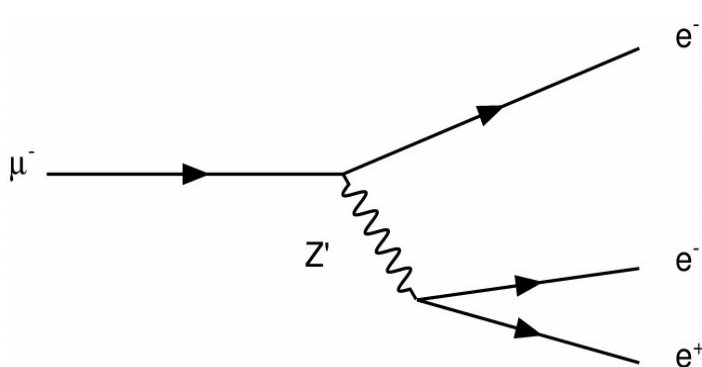
◆ $\mu^- \rightarrow e^- \gamma$



$$BR(\mu \rightarrow e\gamma) = \frac{\Gamma_{\mu \rightarrow e\gamma}}{\Gamma_{\mu \rightarrow e\nu\nu}} = \frac{48\alpha}{\pi} S_W^4 \left(S_W^2 - \frac{1}{2}\right)^2 \lambda_{e\mu}^2 \cdot \frac{M_Z^4}{M_{Z'}^4}$$

Z' - e - μ coupling

◆ $\mu^- \rightarrow e^- e^- e^+$



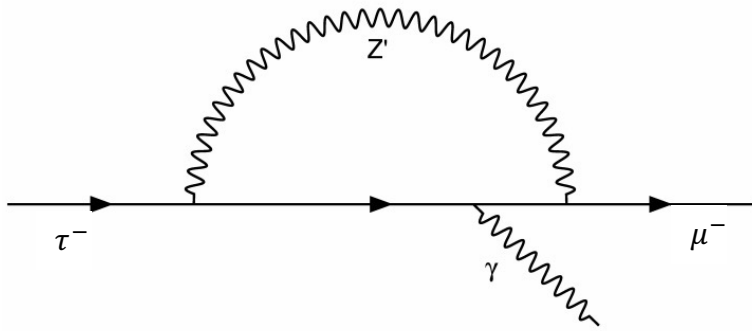
$$BR(\mu \rightarrow eee) = \frac{\Gamma_{\mu \rightarrow eee}}{\Gamma_{\mu \rightarrow e\nu\nu}} = 4 \cdot \lambda_{e\mu}^2 \cdot \frac{M_Z^4}{M_{Z'}^4} \left[S_W^4 + \left(S_W^2 - \frac{1}{2}\right)^2 \right]^2$$

α : the fine structure constant, M_Z : the Z boson mass,

$M_{Z'}$: the mass of Z', S_W : the sine of the weak mixing angle

Phys. Rev. D 62, 013006 (2000)
Phys. Lett. B 723, 15, (2013)

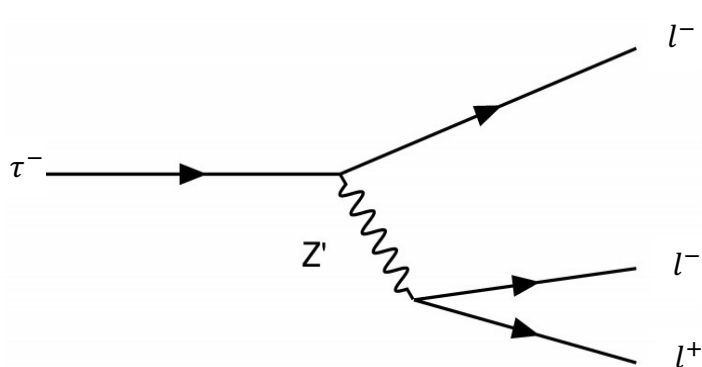
◆ $\tau^- \rightarrow \mu^- \gamma$



$$BR(\tau \rightarrow \mu \gamma) = \frac{48\alpha}{\pi} S_W^4 \left(S_W^2 - \frac{1}{2}\right)^2 \lambda_{\mu\tau}^2 \cdot \frac{M_Z^4}{M_{Z'}^4} BR(\tau \rightarrow \mu \nu \nu)$$

$Z' - \mu - \tau$ coupling

◆ $\tau^- \rightarrow l^- l^- l^+$



$$BR(\tau \rightarrow lll) = 4 \cdot \lambda_{\mu\tau}^2 \cdot \frac{M_{Z'}^4}{M_Z^4} \left[S_W^4 + \left(S_W^2 - \frac{1}{2}\right)^2 \right]^2 BR(\tau \rightarrow \mu \nu \nu)$$

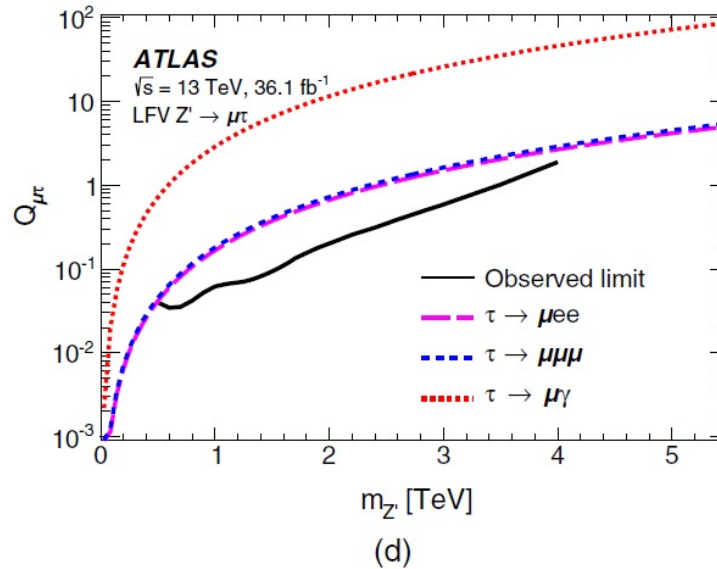
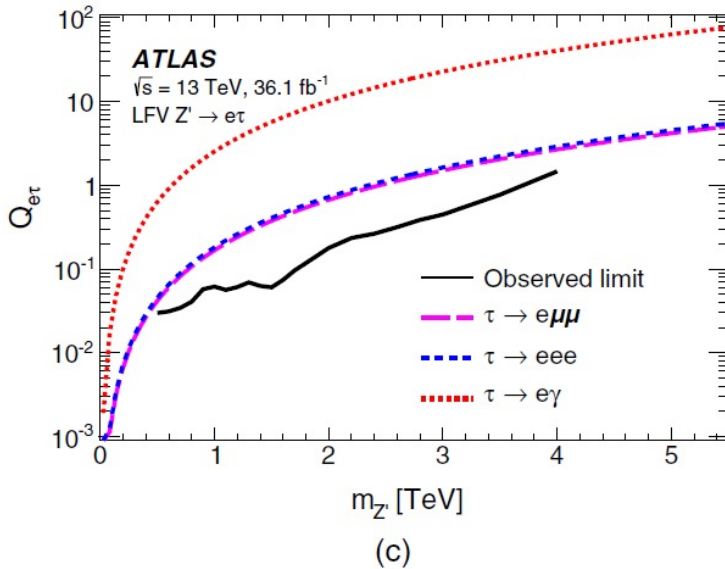
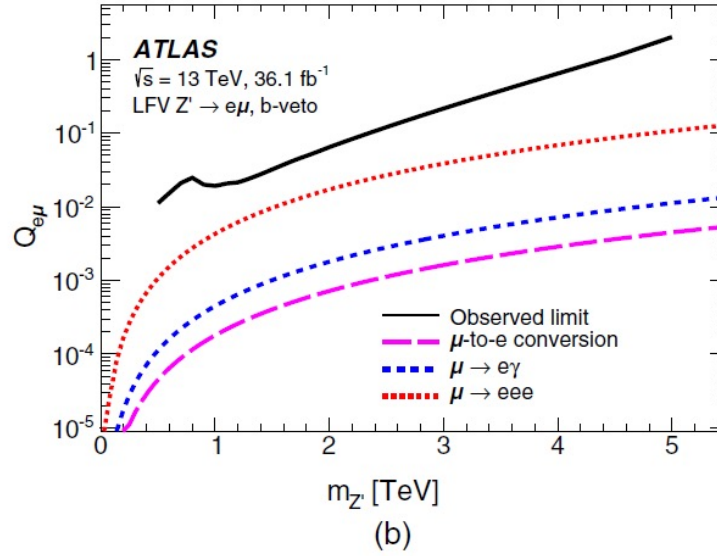
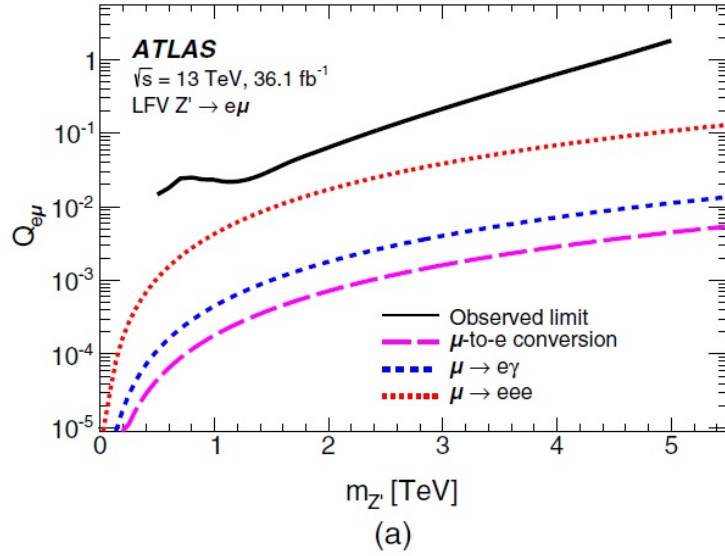
α : the fine structure constant, M_Z : the Z boson mass,

$M_{Z'}$: the mass of Z', S_W : the sine of the weak mixing angle

Phys. Rev. D 62, 013006 (2000)
Phys. Lett. B 723, 15, (2013)



ATLAS CLFV Z' decay result



◆ The ATLAS cross-section times branching ratio limits (solid lines) compared with similar limits from low-energy experiments

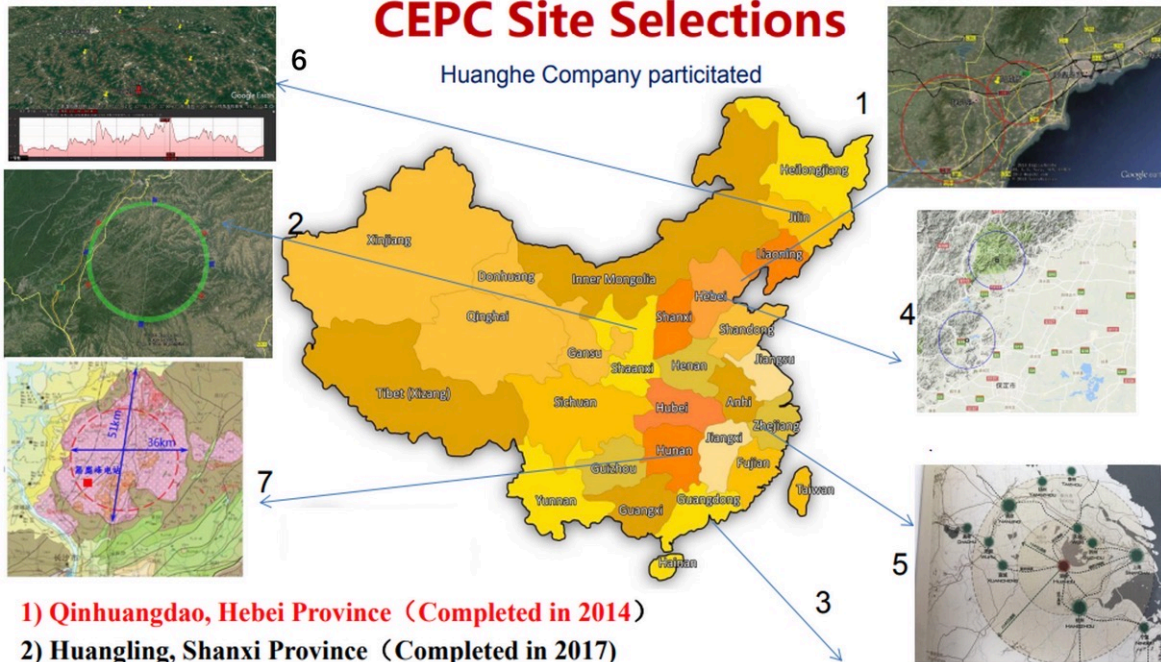
Phys. Rev. D 98, 092008 (2018)

Future collider

integrated luminosity $L \simeq 10 \text{ ab}^{-1} \times (\sqrt{s}/10 \text{ TeV})^2$



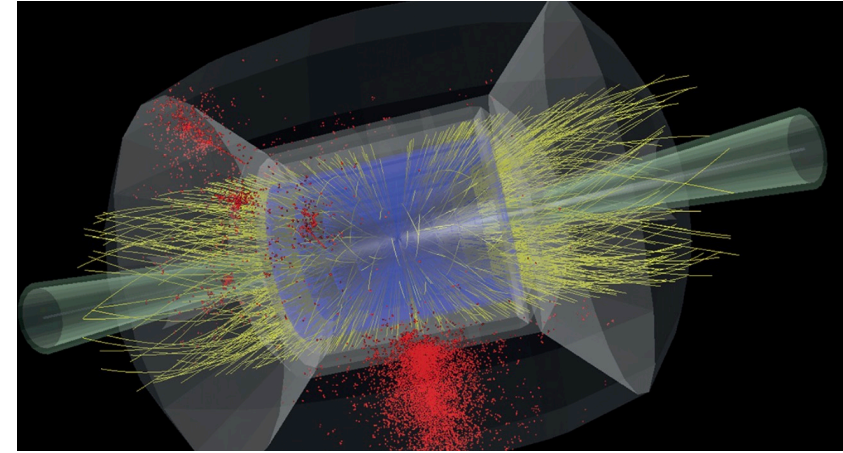
Circular Electron Positron Collider (CEPC)



- 1) Qinhuangdao, Hebei Province (Completed in 2014)
- 2) Huangling, Shanxi Province (Completed in 2017)
- 3) Shenshan, Guangdong Province (Completed in 2016)
- 4) Baoding (Xiong an), Hebei Province (Started in August 2017)
- 5) Huzhou, Zhejiang Province (Started in March 2018)
- 6) Chuangchun, Jilin Province (Started in May 2018)
- 7) Changsha, Hunan Province (Started in Dec. 2018)

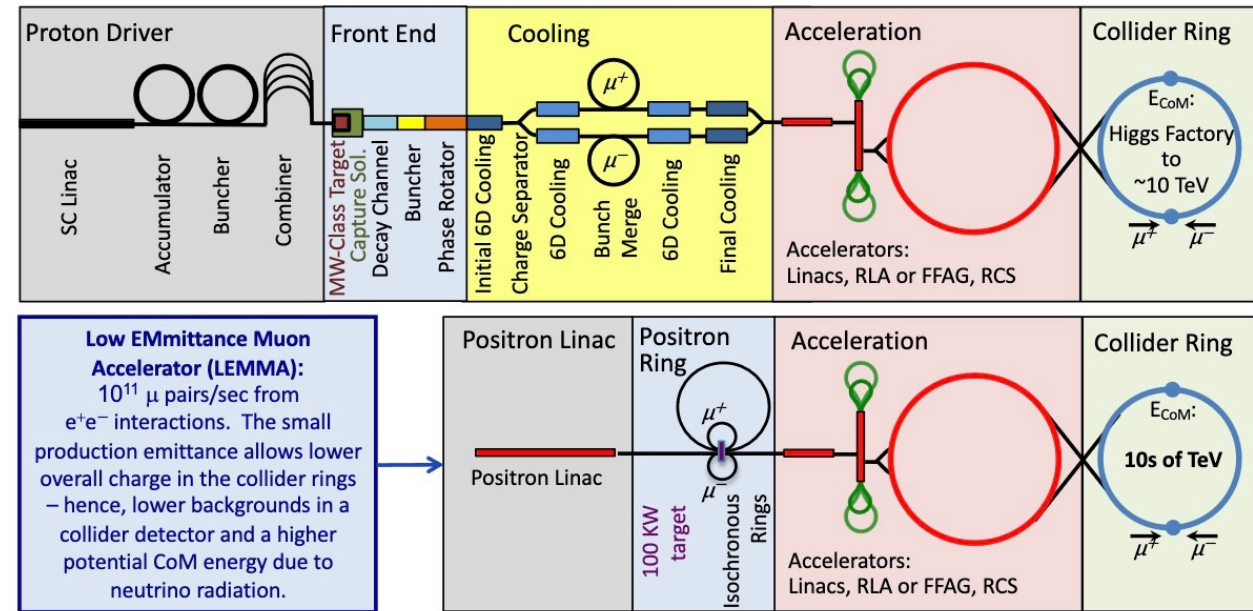
large-scale high-energy physics experimental facility perform high-precision detection of the Higgs boson

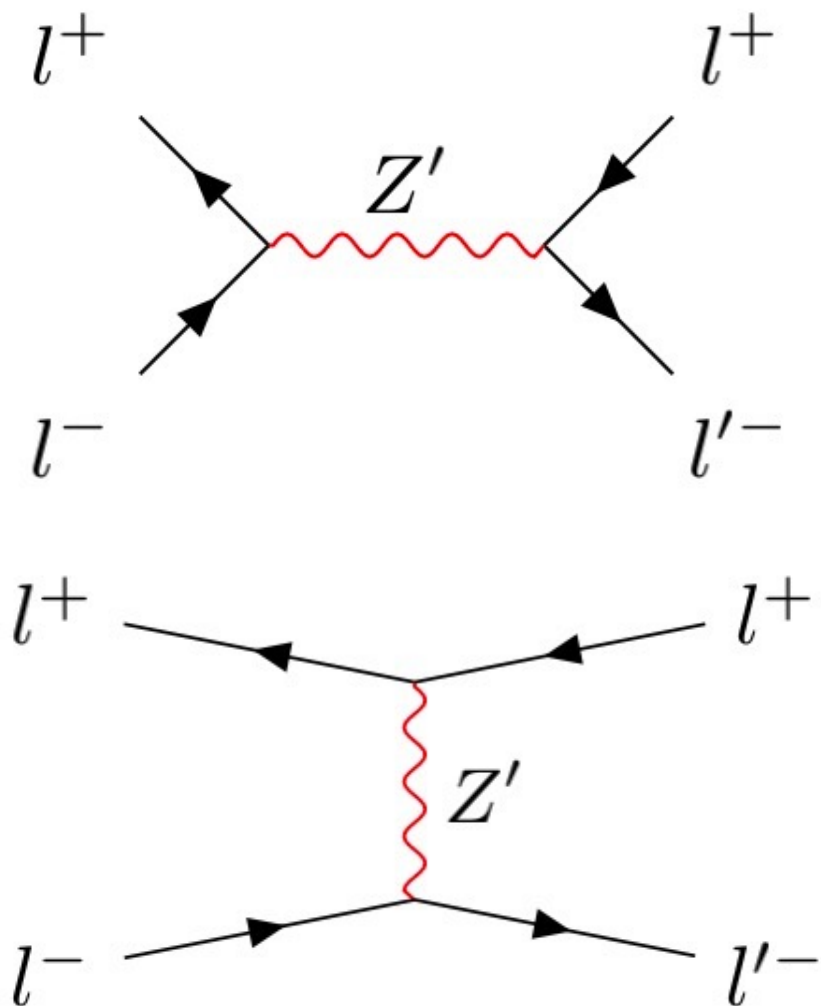
JACoW IPAC2023 (2023) MOPL051
Radiat. Detect. Technol. Methods 8 (2024)



arXiv:1901.06150
JINST 19 (2024)
02, T02015
Eur.Phys.J.C 84
(2024) 1, 36

Muon Collider





- ◆ CEPC: $ee \rightarrow e\mu, ee \rightarrow e\tau$
- ◆ Muon collider: $\mu\mu \rightarrow e\mu, \mu\mu \rightarrow \mu\tau$
- ◆ Only one CLFV coupling $\lambda_{ij} (i \neq j)$ is assumed to be non-zero while the diagonal couplings $\lambda_{ij} (i = j)$ are always set as 1.

$$\lambda_{ij} = \begin{pmatrix} \lambda_{ee} & \lambda_{e\mu} & \lambda_{e\tau} \\ \lambda_{\mu e} & \lambda_{\mu\mu} & \lambda_{\mu\tau} \\ \lambda_{\tau e} & \lambda_{\tau\mu} & \lambda_{\tau\tau} \end{pmatrix}$$

- ◆ Using @Madgraph, @Pythia8 and @Delphes to generate the processes



process	Cross section(pb)
$ee \rightarrow e\mu$	4.04×10^{-5}
$ee \rightarrow ww, w \rightarrow ev, w \rightarrow \mu\nu$	0.395
$ee \rightarrow \tau\tau, \tau \rightarrow ev\nu, \tau \rightarrow \mu\nu$	0.241
$ee \rightarrow hv\nu, h \rightarrow \tau\tau, \tau \rightarrow ev\nu/\mu\nu$	1.13×10^{-4}
$ee \rightarrow hv\nu, h \rightarrow ww, w \rightarrow ev/\mu\nu$	3.93×10^{-6}
$ee \rightarrow e\tau$	6.94×10^{-5}
$ee \rightarrow ww, w \rightarrow ev, w \rightarrow \tau\nu$	3.733
$ee \rightarrow \tau\tau, \tau \rightarrow \mu\nu$	0.658
$ee \rightarrow hv\nu, h \rightarrow \tau\tau, \tau \rightarrow \mu\nu$	4.28×10^{-3}

- ◆ Control τ decay to μ in MG5, and control another τ to hadrons in Pythia8.
- ◆ For the τ final state, only the hadronized τ is considered, the cross section needs to $\times 60\%$.

process	Cross section(pb)
$\mu\mu \rightarrow e\mu$ (14TeV collider)	3.38×10^{-4}
$\mu\mu \rightarrow ww\nu\nu, w \rightarrow ev, w \rightarrow \mu\nu$	0.013
$\mu\mu \rightarrow ww, w \rightarrow ev, w \rightarrow \mu\nu$	7.71×10^{-4}
$\mu\mu \rightarrow \tau\tau, \tau \rightarrow ev\nu, \tau \rightarrow \mu\nu$	3.20×10^{-5}
$\mu\mu \rightarrow hv\nu, h \rightarrow \tau\tau, \tau \rightarrow ev\nu/\mu\nu$	2.22×10^{-3}
$\mu\mu \rightarrow hv\nu, h \rightarrow ww, w \rightarrow ev/\mu\nu$	7.68×10^{-5}
$\mu\mu \rightarrow \mu\tau$ (6TeV collider)	0.042
$\mu\mu \rightarrow ww\nu\nu, w \rightarrow \tau\nu, w \rightarrow \mu\nu$	6.47×10^{-3}
$\mu\mu \rightarrow ww, w \rightarrow \tau\nu, w \rightarrow \mu\nu$	3.40×10^{-3}
$\mu\mu \rightarrow \tau\tau, \tau \rightarrow \mu\nu$	9.81×10^{-4}
$\mu\mu \rightarrow hv\nu, h \rightarrow \tau\tau, \tau \rightarrow \mu\nu$	2.28×10^{-3}
$\mu\mu \rightarrow hv\nu, h \rightarrow ww, w \rightarrow \tau\nu/\mu\nu$	2.92×10^{-5}



Preliminary selection and efficiency

◆ The events are required to satisfy the requirements of lepton flavor and charge conservation, i.e., $e^+e^- \rightarrow e^+\mu^-$, all **signal and background** events are required to have one e^+ and one μ^- .

◆ $e\mu$ final states:

$$p_T > 10 \text{ GeV}, |\eta| < 2.5$$

◆ Final state containing τ :

$$p_T > 20 \text{ GeV}, |\eta| < 5$$

p_T : the transverse momentum, $|\eta|$: the pseudo-rapidity

◆ μ tracking efficiency

Collider	Conditions	Efficiency
CEPC	$0.1 < \eta \leq 3$	100%
	$ \eta > 3$	0%
Muon Collider	$ \eta < 2.0, 0.5 < p_T < 1 \text{ GeV}$	95%
	$ \eta \leq 2.0, p_T > 1 \text{ GeV}$	99%
	$2.0 < \eta < 2.5, 0.5 < p_T \leq 1 \text{ GeV}$	90%
	$2.0 < \eta < 2.5, p_T > 1 \text{ GeV}$	95%
	$ \eta > 2.5$	0%

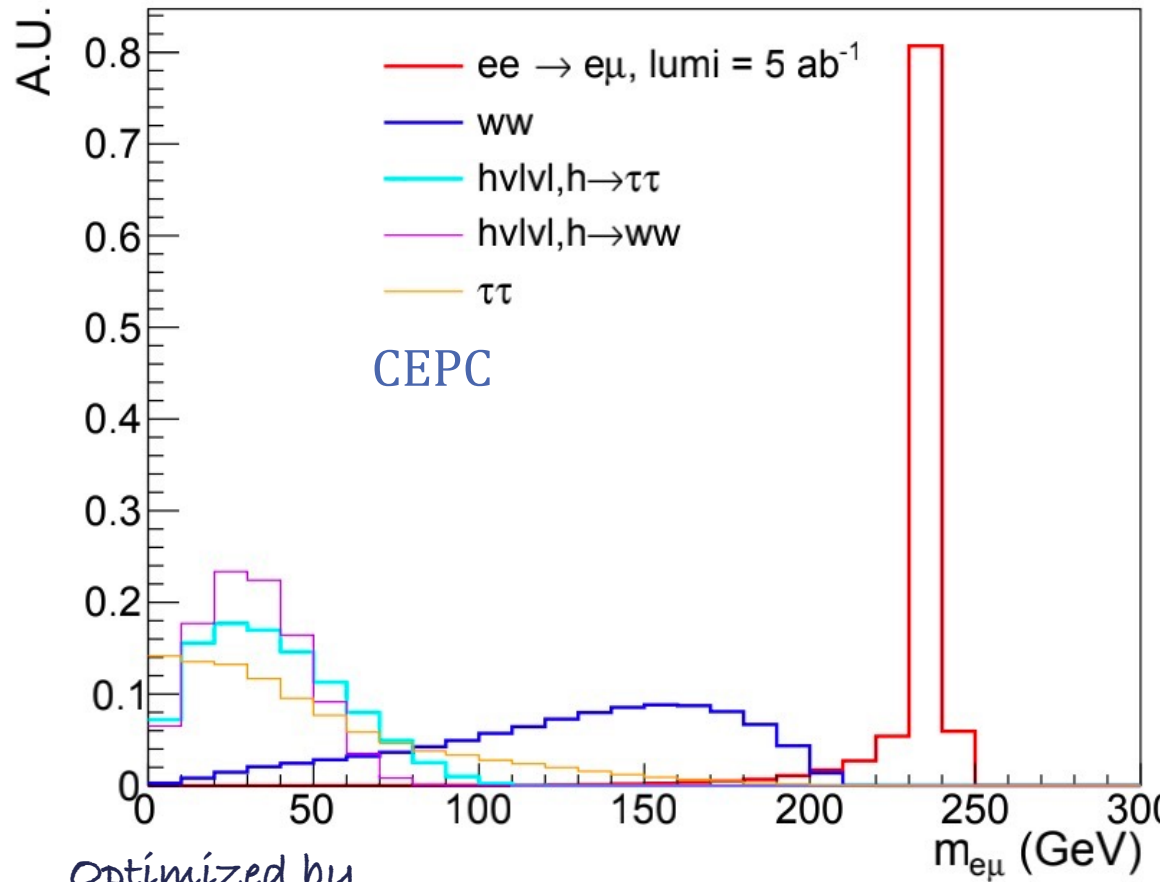
◆ τ tagging efficiency

Collider	Efficiency
CEPC	40%
Muon Collider	80%



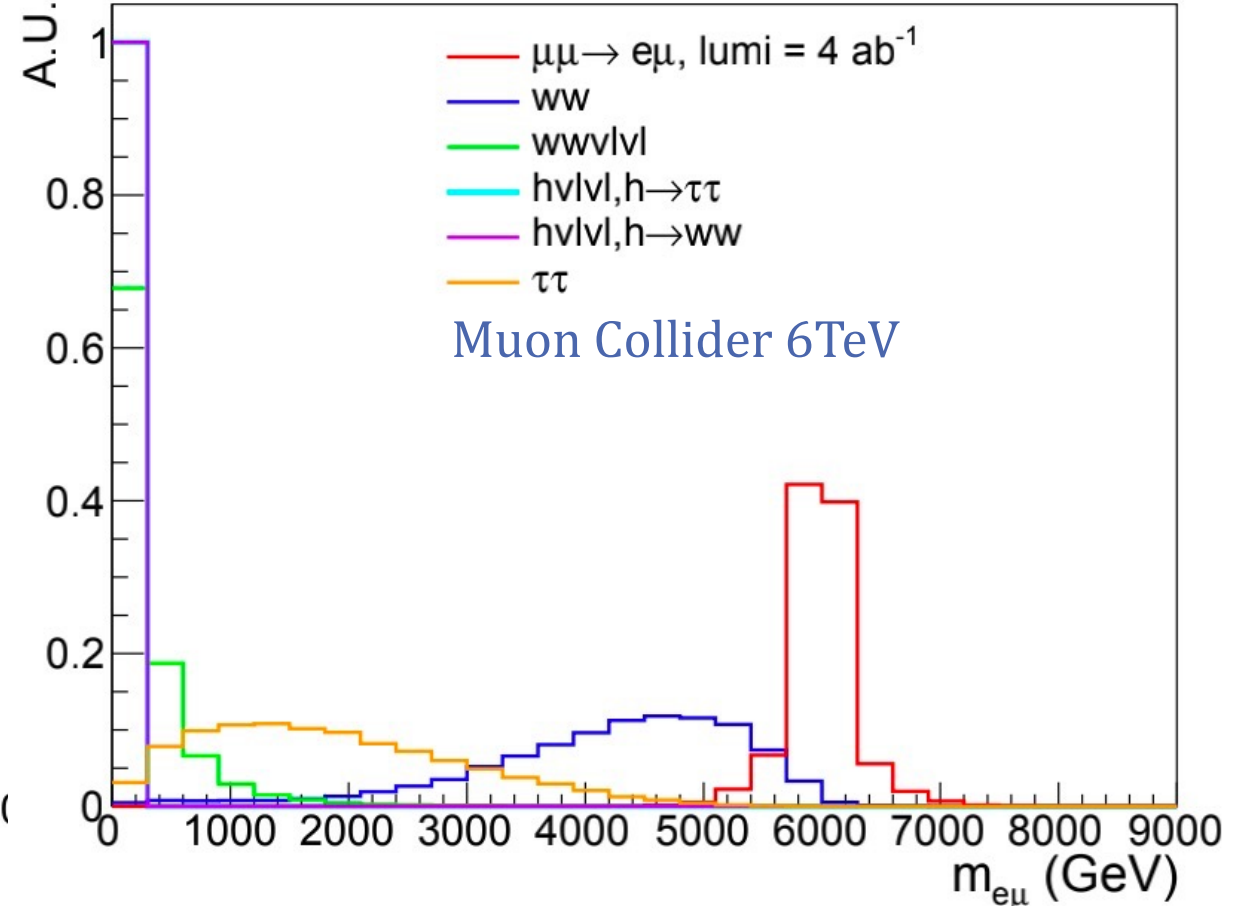
Further cuts

◆ Using $e\mu$ invariant mass to separate the signal and the backgrounds.



Optimized by maximizing the FOM $(\frac{s}{\sqrt{s+b}})$

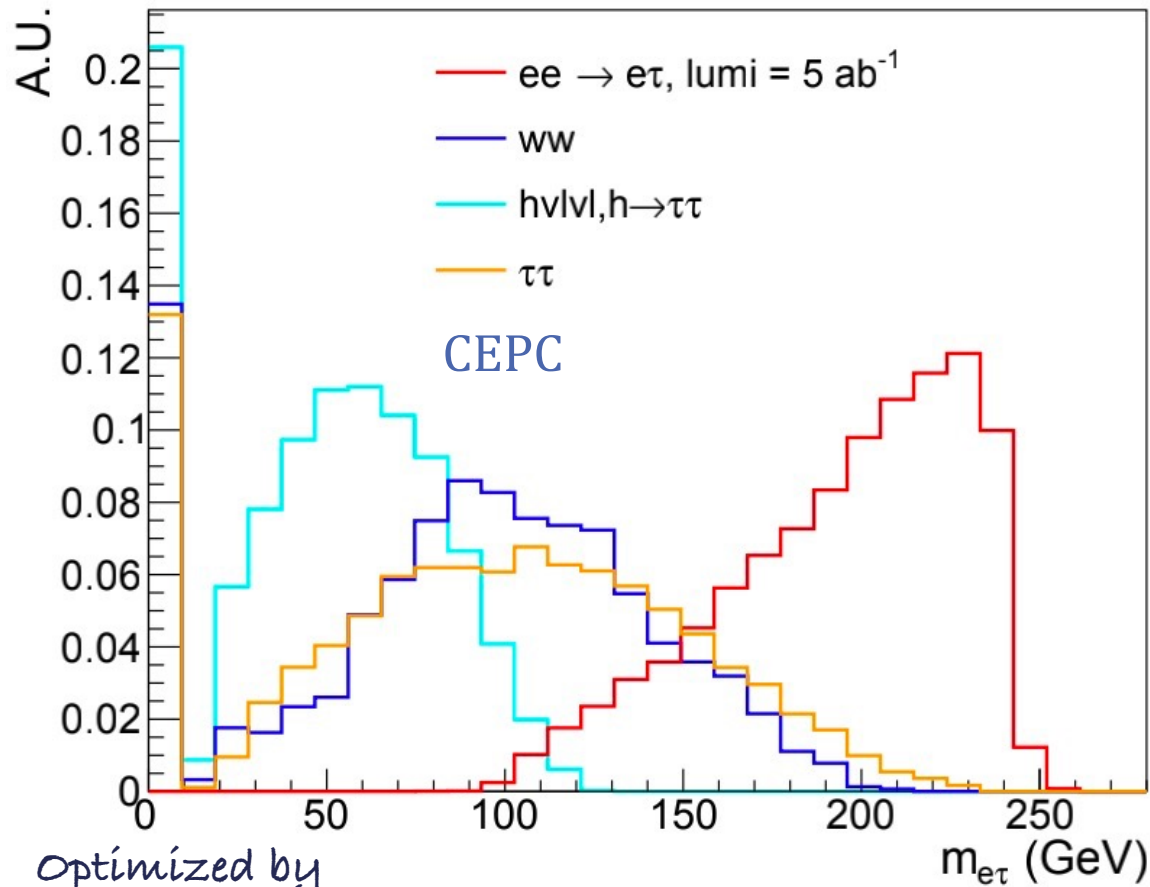
$m_{e\mu} > 220 \text{ GeV}$



$m_{e\mu} > 5.2 \text{ TeV}$ ($m_{e\mu} > 10 \text{ TeV}$ for 14TeV Muon collider)

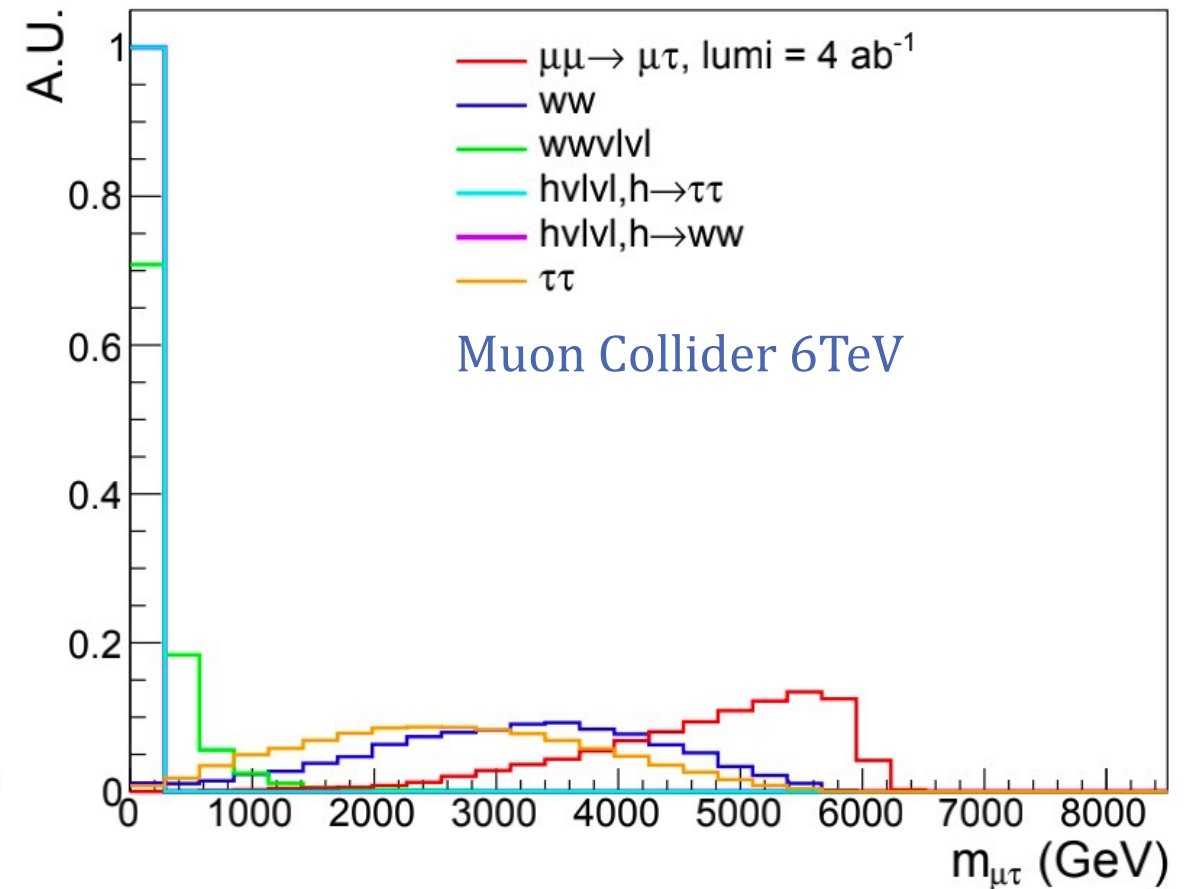


◆ Using $e\tau, \mu\tau$ invariant mass to separate the signal and the backgrounds.



Optimized by maximizing the FOM $(\frac{s}{\sqrt{s+b}})$

$m_{e\tau} > 160$ GeV



$m_{\mu\tau} > 4$ TeV ($m_{\mu\tau} > 9.5$ TeV for 14TeV Muon collider)



◆ After all selections, get binned histograms on the final state lepton p_T distributions

◆ Per-event weight to account for the cross-section difference : $n_{L_X} = \sigma_X L / N_X$

◆ Defined the negative log likelihood test statistics Z:

$$Z = \sum_{i=1}^{bins} Z_i$$

$$Z_i := 2[n_i - b_i + b_i \ln(b_i/n_i)] \quad 95\% \text{ C.L. Exclusion}$$

σ_X : cross section

L : luminosity

N_X : events generated

i : the bin number, s : the beyond SM signal, b : the SM background

$n = s + b$: the total yields containing both signal and background

Eur. Phys. J. C 73, 2501 (2013)

◆ The Z statistic subjects to a χ^2 distribution with degree of freedom equals to 1.



Using the Z' model formula, these upper limits can be converted into coupling upper limits

Eur. Phys. J. C 84, 216 (2024)
Eur. Phys.J. C 47, 337 (2006)
Nucl. Phys. B 299, 1 (1988).
Phys. Rev. Lett. 104, 021802 (2010)
High Energy. Phys. 2021, 19 (2021)
Phys.Lett.B 687, 139 (2010)

- ◆ $\mathcal{B}(\mu^- \rightarrow e^- \gamma) < 3.1 \times 10^{-13}$ @ 90% C.L. **MEGII**
- ◆ $\mathcal{B}(\mu^- N \rightarrow e^- N) < 7.0 \times 10^{-13}$ @ 90% C.L. **Mu2e**
- ◆ $\mathcal{B}(\mu^- \rightarrow e^- e^+ e^-) < 1.0 \times 10^{-12}$ @ 90% C.L. **Mu3e**
- ◆ $\mathcal{B}(\tau^- \rightarrow e^- \gamma) < 3.3 \times 10^{-8}$ @ 90% C.L. **BABAR**
- ◆ $\mathcal{B}(\tau^- \rightarrow \mu^- \gamma) < 4.2 \times 10^{-8}$ @ 90% C.L. **Belle**
- ◆ $\mathcal{B}(\tau^- \rightarrow e^- e^+ e^-) < 2.7 \times 10^{-8}$ @ 90% C.L. **Belle**
- ◆ $\mathcal{B}(\tau^- \rightarrow \mu^- \mu^+ \mu^-) < 2.1 \times 10^{-8}$ @ 90% C.L. **Belle**
- ◆ $\mathcal{B}(\tau^- \rightarrow \mu^- e^+ e^-) < 1.8 \times 10^{-8}$ @ 90% C.L. **Belle**
- ◆ $\mathcal{B}(\tau^- \rightarrow e^- \mu^+ \mu^-) < 2.7 \times 10^{-8}$ @ 90% C.L. **Belle**

Prospect limits



Symmetry 13 no.9, 1591 (2021)

PTEP 3, 033C01 (2020)

arXiv:1501.05241

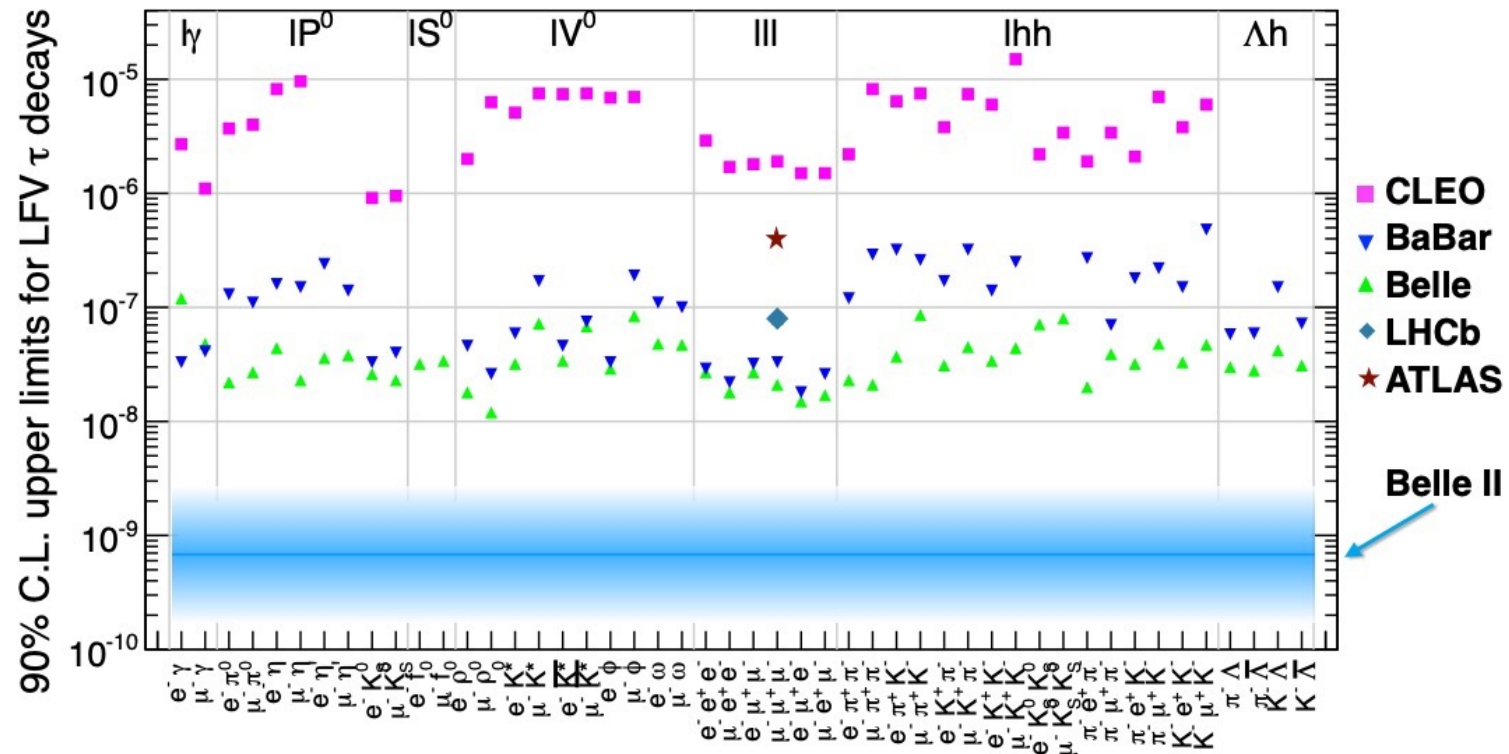
Nucl.Instrum.Meth.A 1014, 165679 (2021)

- ◆ $B(\tau^- \rightarrow e^- \gamma) < 9.0 \times 10^{-9}$
- ◆ $B(\tau^- \rightarrow \mu^- \gamma) < 6.9 \times 10^{-9}$
- ◆ $B(\tau^- \rightarrow e^- e^+ e^-) < 4.7 \times 10^{-10}$
- ◆ $B(\tau^- \rightarrow \mu^- \mu^+ \mu^-) < 3.6 \times 10^{-10}$
- ◆ $B(\tau^- \rightarrow \mu^- e^+ e^-) < 2.9 \times 10^{-10}$
- ◆ $B(\tau^- \rightarrow e^- \mu^+ \mu^-) < 4.5 \times 10^{-10}$

@ 90% C.L. **Belle II**

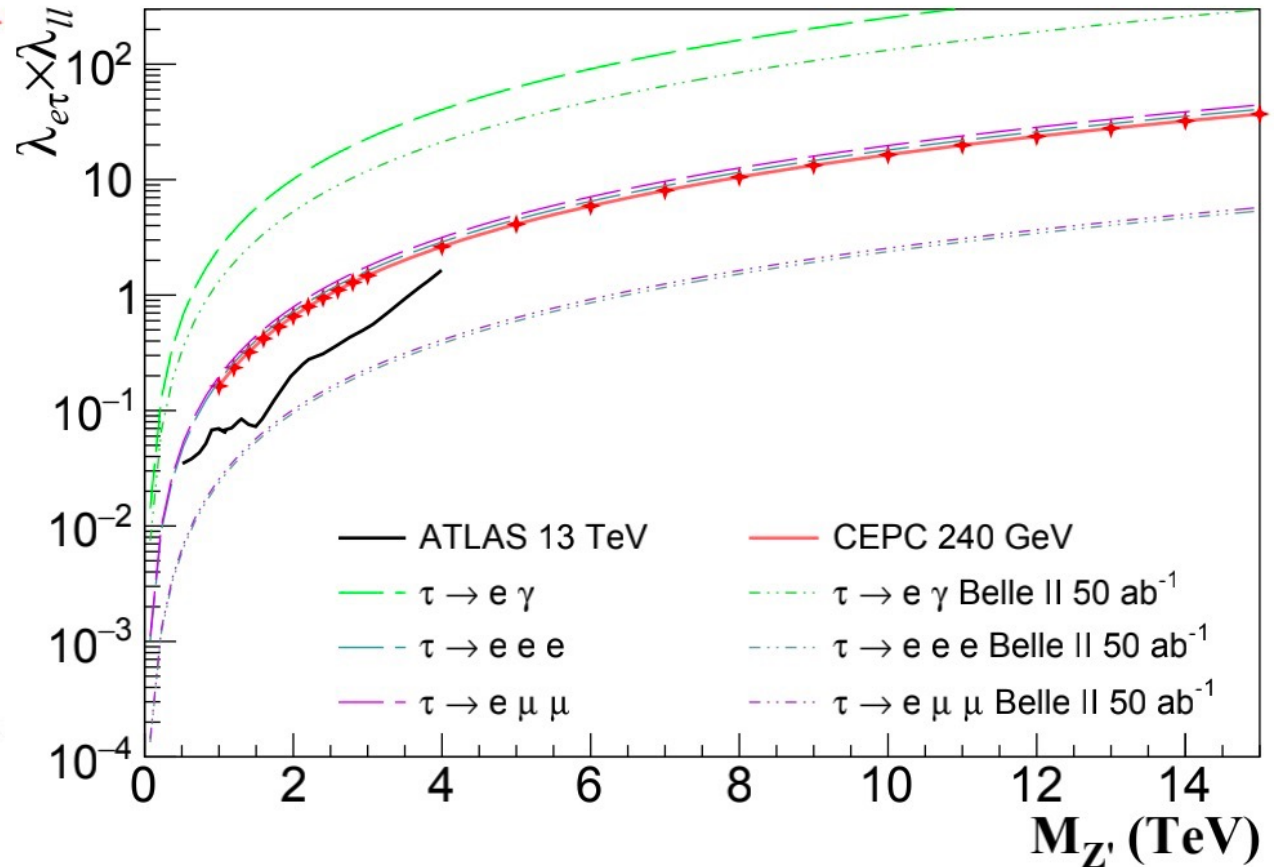
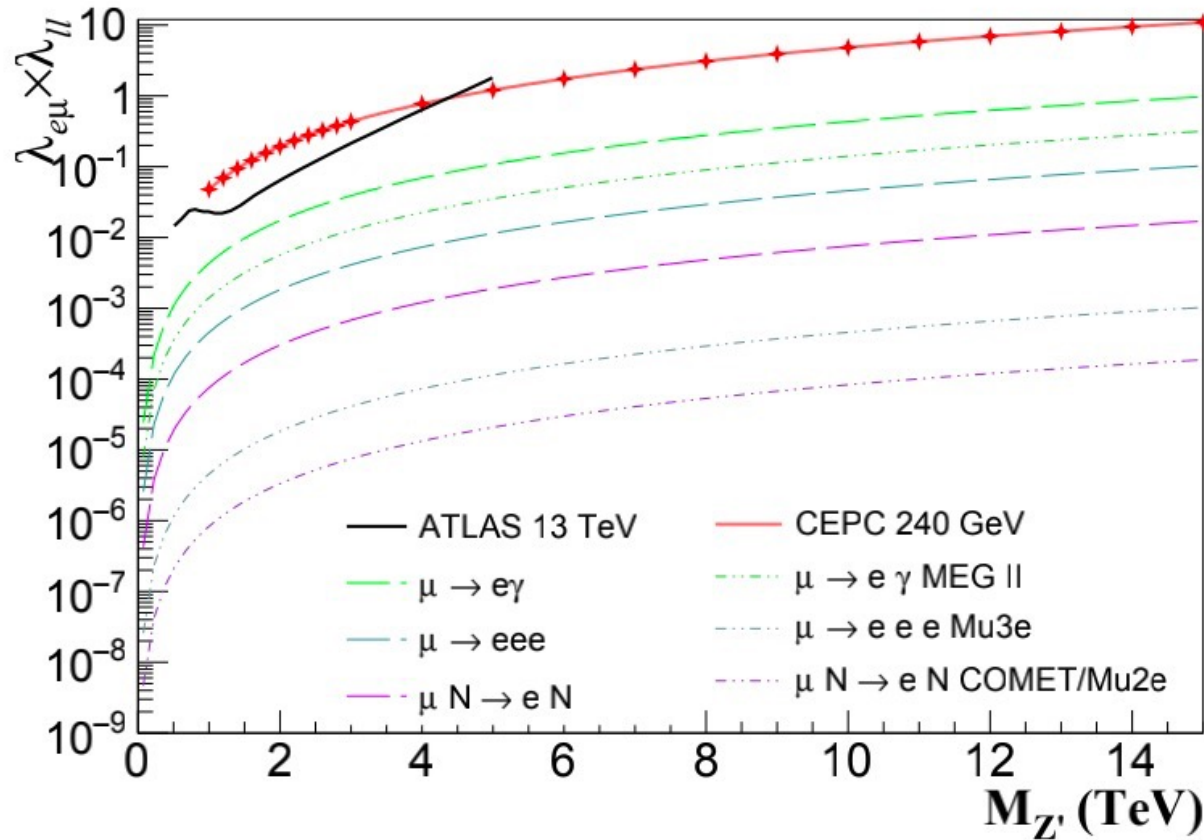
Universe 8 no.9, 480 (2022)

- ◆ $B(\mu^- \rightarrow e^- \gamma) < 6.0 \times 10^{-14}$ @ 90% C.L. **MEGII**
- ◆ $B(\mu^- N \rightarrow e^- N) < 3.0 \times 10^{-17}$ @ 90% C.L. **COMET**
- ◆ $B(\mu^- N \rightarrow e^- N) < 8.0 \times 10^{-17}$ @ 90% C.L. **Mu2e**
- ◆ $B(\mu^- \rightarrow e^- e^+ e^-) < 1.0 \times 10^{-16}$ @ 90% C.L. **Mu3e**





◆ The curves are plotted as functions of $M_{Z'}$, from the cross-section times branching ratio limits.

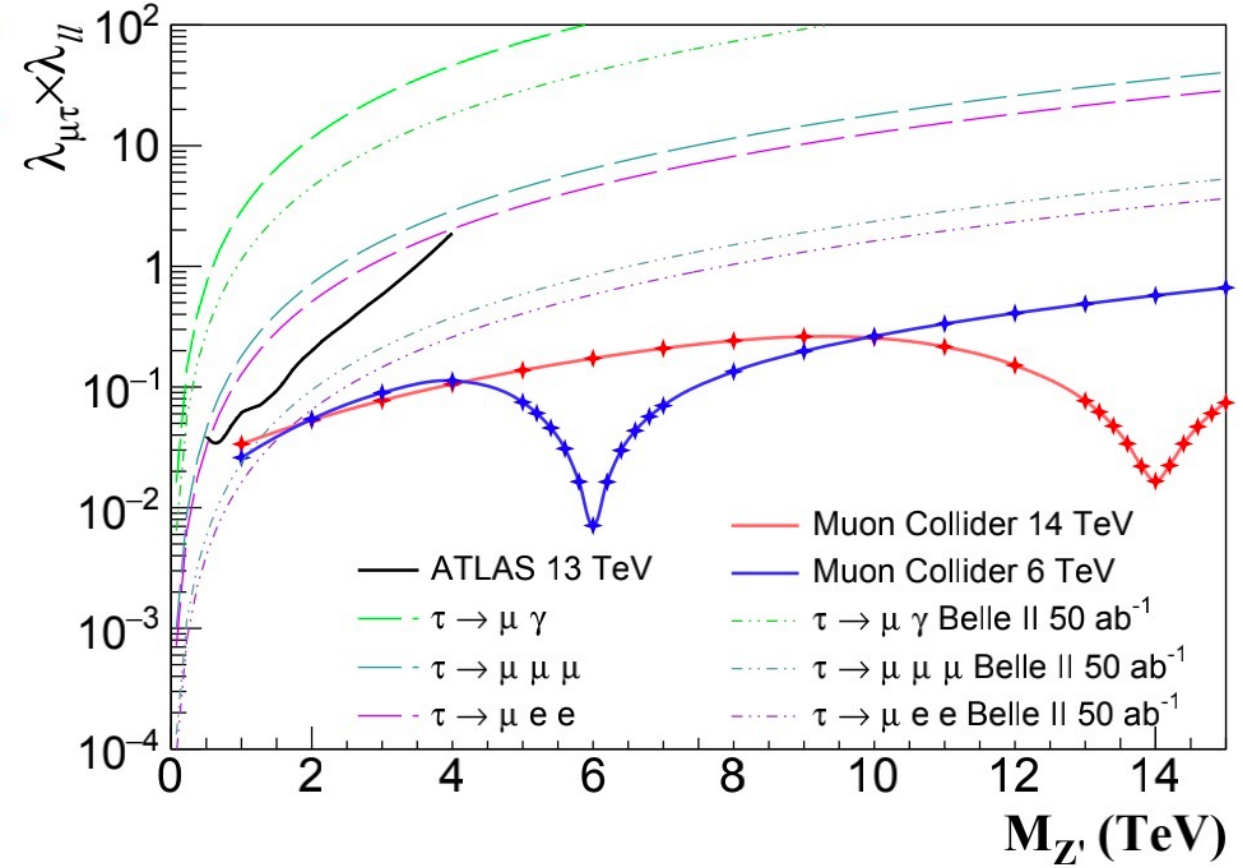
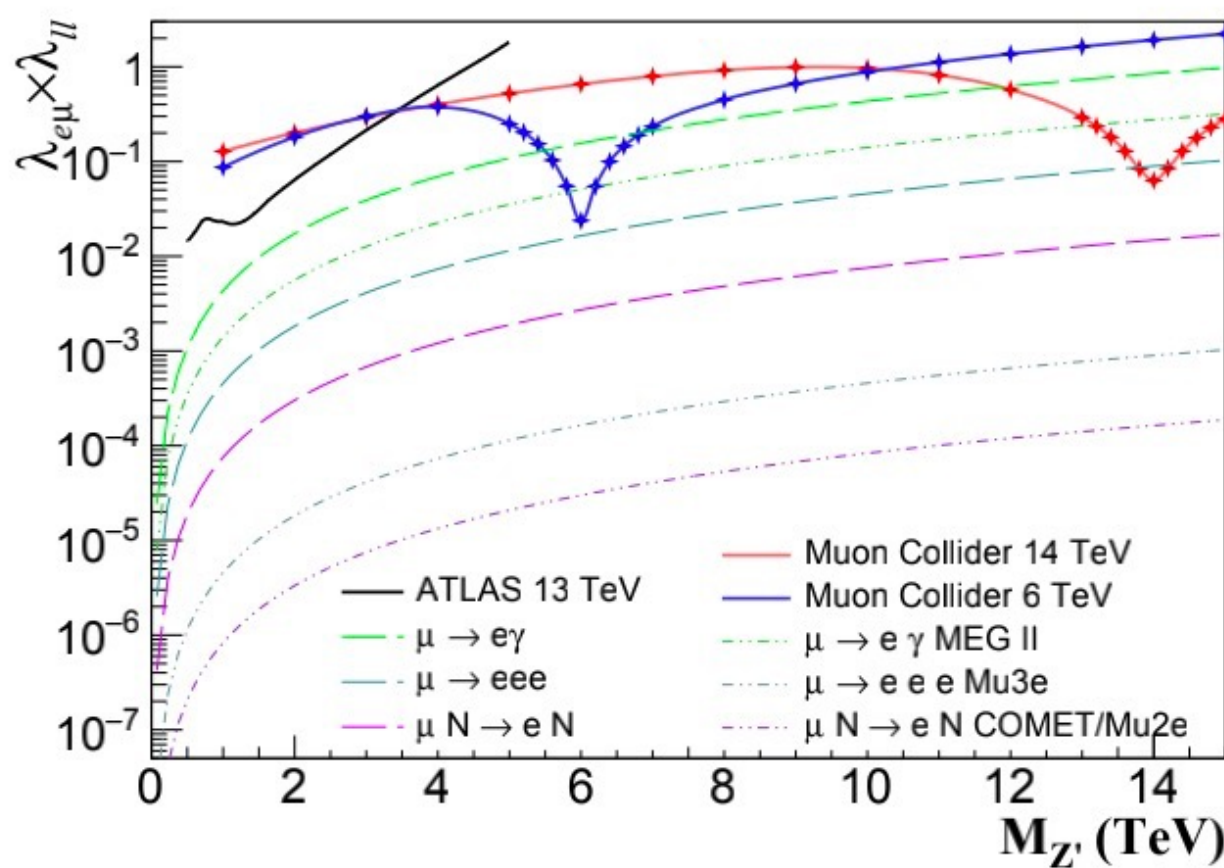


◆ Compared with the ATLAS experiment, current low-energy experiments (dashed lines) and future experiments (dash-dotted lines).



Upper limit on MuonC

◆ The curves are plotted as functions of $M_{Z'}$, from the cross-section times branching ratio limits.



◆ For $\mu\tau$ channel, the two coupling limits in this work are the most stringent when the mass of Z' is greater than 1.5 TeV.



- ◆ The observation of any CLFV process would be a clear signal of new physics beyond the SM.
- ◆ Perform a detailed comparative study on CLFV searches at a 6 (14) TeV scale muon collider and a 240 GeV electron-positron collider.
- ◆ The strongest constraint: $\mu\tau$ coupling at the 6TeV Muon collider, reaching 10^{-3} when Z' mass equals to 6 TeV, which is stronger than the current best limits on CLFV.
- ◆ The τ related CLFV coupling strength will be significantly improved.

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Thank you!

$$\Gamma(\mu \rightarrow e\gamma) \approx \frac{G_F^2 m_\mu^5}{192\pi^3} \left(\frac{\alpha}{2\pi}\right) \sin^2 2\theta \sin^2 \left(\frac{1.27\Delta m^2}{M_W^2}\right)$$

$\mu - \text{decay}$ $\gamma - \text{vertex}$ $\vartheta - \text{oscillation}$

$$\approx \frac{G_F^2 m_\mu^5}{192\pi^3} \left(\frac{3\alpha}{32\pi}\right) \left(\frac{\Delta m_{23}^2 s_{13} c_{13} s_{23}}{M_W^2}\right)^2$$

with $\Delta \sim 10^{-3} eV^2$, $M_W \sim O(10^{11}) eV \approx O(10^{-54})$

