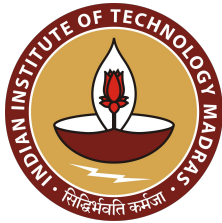


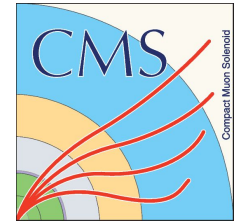


Probing Hydrodynamics in PbPb Collisions at 5.02 TeV using Higher-order Cumulants

18th July, 2024



Aryaa Dattamunsi
IIT Madras
On behalf of the CMS Collaboration





❖ Introduction

- Anisotropic flow
- Asymmetry of v_2 distribution

❖ Motivation

- Hydrodynamic probes

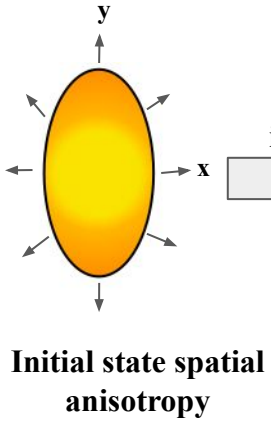
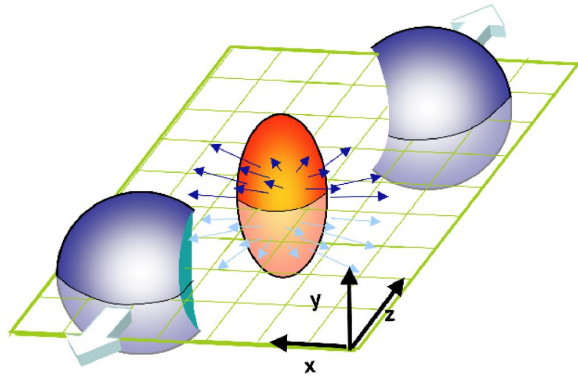
❖ Analysis Technique

- Collectivity - $v_2\{2k\}$ in terms of Q-cumulants in PbPb collisions

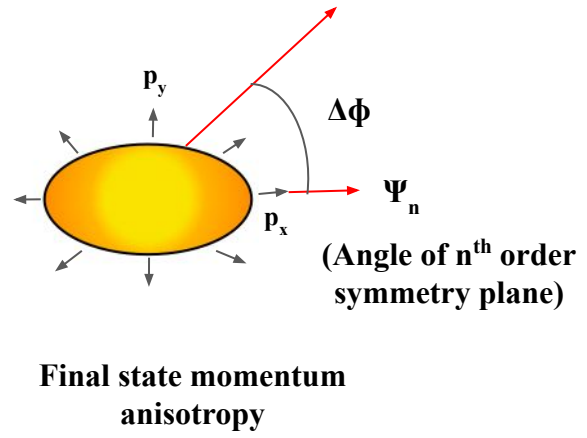
❖ Results

- Measurement of two hydrodynamic probes vs centrality
- High-precision measurement of central moments of v_2 distribution

❖ Summary



Collective interaction
Pressure gradient



Azimuthal Anisotropy

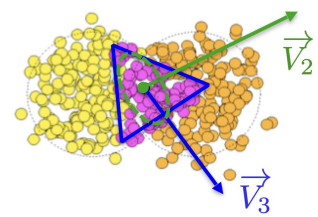
$$\frac{2\pi}{N} \frac{dN}{d\phi} = 1 + \sum_{n=1}^{\infty} 2v_n \cos n(\Delta\Phi)$$

$$\Delta\phi = \phi - \psi_n$$

$$v_n \equiv \langle \cos[n(\phi - \psi_n)] \rangle$$

v_n - **Fourier flow harmonics** depend on :

- initial state geometry
- initial state fluctuations
- medium transport properties (η/s)



$$\vec{V}_m = v_m e^{-im\Psi_m}$$

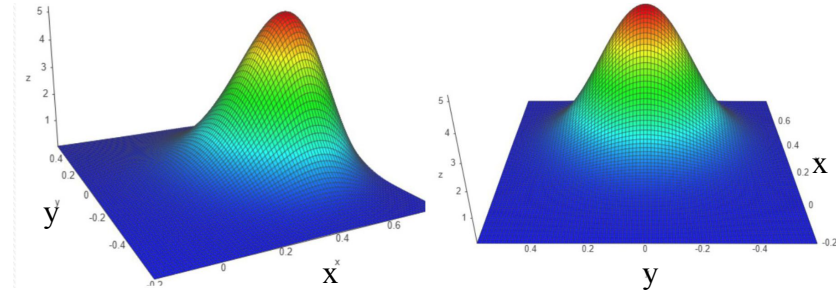
$$\vec{V}_n = v_n e^{-in\Psi_n}$$

- Hydrodynamics : azimuthally anisotropic expansion of QGP formed in AA collisions
- Event-by-event fluctuations : the early stage dynamics
- Non-Gaussianities in e-by-e v_2 distribution
- Hydrodynamic expansion - $v_2 \propto \epsilon_2$
- **Fluctuations** present in initial state = **non-gaussianities** in v_2 distribution in final state
- Precise measurements of the Non-Gaussian flow fluctuations : test hydrodynamics and constrain IS models

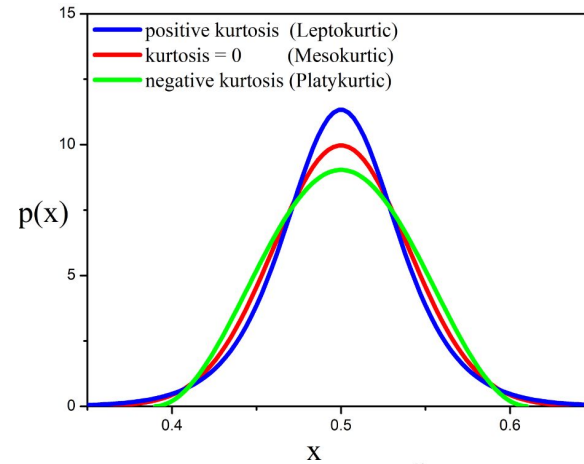
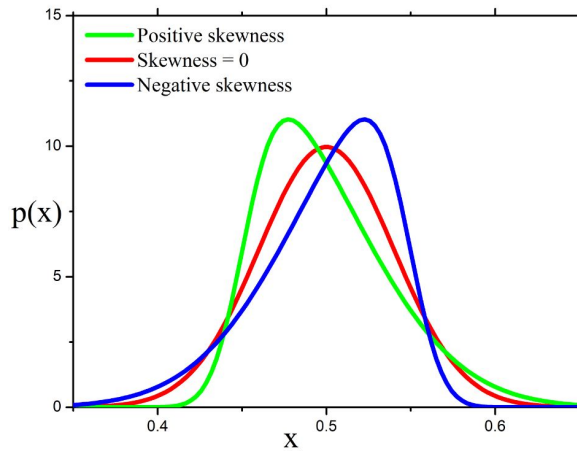
$$\mathbf{v}_2 = v_x \mathbf{e}_x + v_y \mathbf{e}_y$$

$$v_2 = \sqrt{v_x^2 + v_y^2}$$

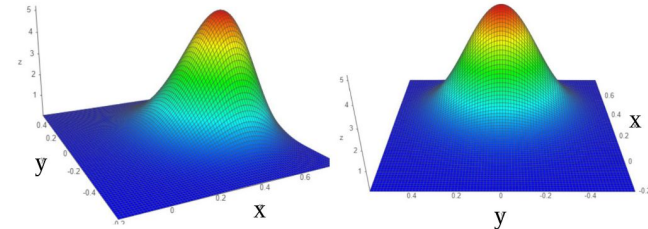
\mathbf{e}_x : along impact parameter
 \mathbf{e}_y : perpendicular to impact parameter



- Q-cumulants, $v_2\{2k\}$ ($k = 1, 2, \dots$) : good tool to study non-Gaussianities
- Non-Gaussianities : fine splitting between cumulants of different orders



- **Skewness** : degree of asymmetry of the distribution
- **Kurtosis** : degree of peakedness and flatness
- **Superskewness** : measure of asymmetry of tail



- Hydrodynamic probes : **observed azimuthal angle correlations** \longleftrightarrow **initial-state geometry**

Phys. Rev. C 95 (2017) 014913

$$\frac{v_2\{6\} - v_2\{8\}}{v_2\{4\} - v_2\{6\}} \approx \frac{1}{11} \approx 0.091$$

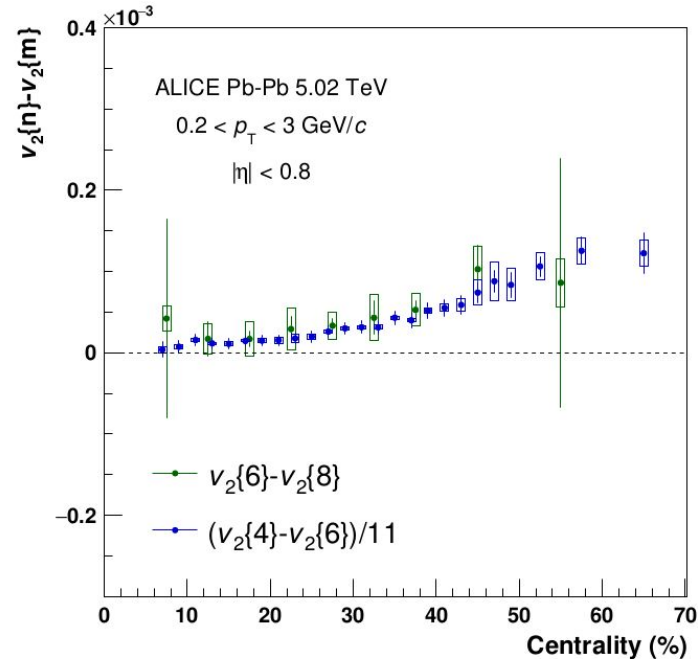
(limited to leading order term)

$0.143 \pm 0.008(stat) \pm 0.014(syst)$: 20-25% centrality
 $0.185 \pm 0.005(stat) \pm 0.012(syst)$: 55-60% centrality

Phys. Lett. B 789 (2019) 643 (CMS)

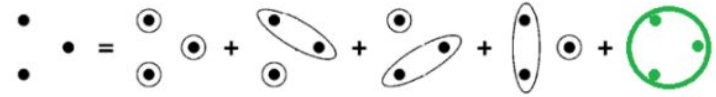
- Goal - to improve precision
- Possible Solution : Introducing higher-order terms in the cumulant expansion of the v_2 distribution**
 - large amount of statistics required

~~Centrality independent~~



JHEP 07 (2018) 103

- k-particle cumulant ($k > 1$) : **collective nature of flow**
 - suppresses non-flow



R. Kubo, J. Phys. Soc. Jpn. 17 (1962)

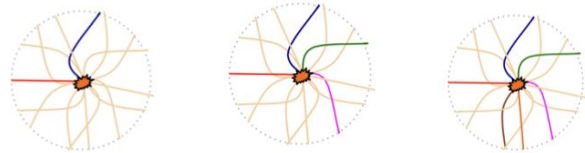
Q-Vector

$$Q_n = \sum_i e^{in\phi_i}$$

All-event average

$$\langle\langle 2 \rangle\rangle \equiv \langle\langle e^{in(\phi_1 - \phi_2)} \rangle\rangle$$

$$\langle\langle 4 \rangle\rangle \equiv \langle\langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle\rangle$$



k = 1

k = 2

k = 3

[Courtesy : Y. Zhou, 2019]

Single-event average over all particles

$$\langle 2 \rangle \equiv \langle e^{in(\phi_1 - \phi_2)} \rangle \implies \langle 2 \rangle \equiv \frac{|Q_n|^2 - M}{M(M-1)}$$

$$\langle 4 \rangle \equiv \langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle$$

$$\implies \langle 4 \rangle \equiv \frac{|Q_n|^4 + |Q_{2n}|^2 - 2\text{Re}[Q_{2n}Q_n^*Q_n^*]}{M(M-1)(M-2)(M-3)}$$

$$- 2 \frac{2(M-2)|Q_n|^2 - M(M-3)}{M(M-1)(M-2)(M-3)}$$

$v_n\{10\}$ from Q-cumulants

- **10-particle azimuthal correlator** $\langle\langle 10 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 + \phi_3 + \phi_4 + \phi_5 - \phi_6 - \phi_7 - \phi_8 - \phi_9 - \phi_{10})} \rangle\rangle$
- Recurrence relation (Phys. Rev. C 104 (2021) 034906) :

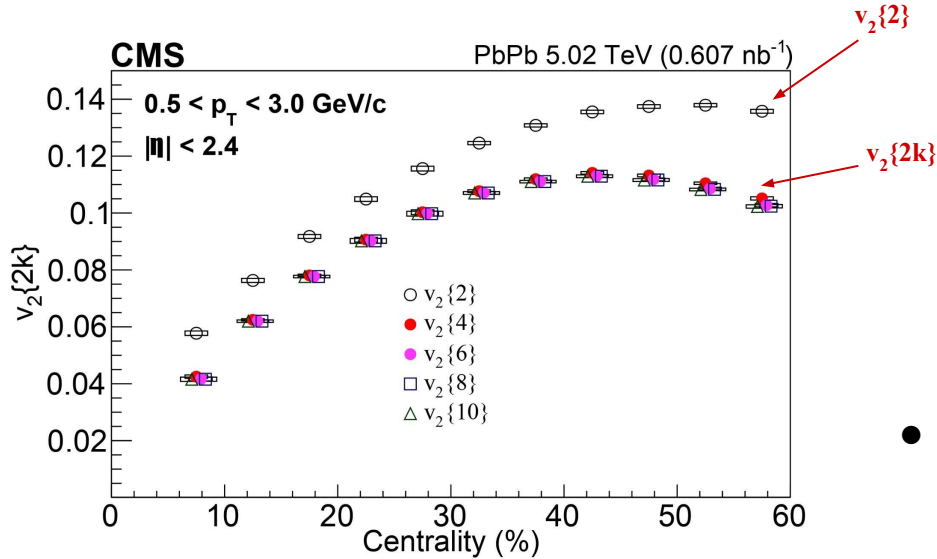
$$c_n\{2k\} = \langle\langle 2k \rangle\rangle - \sum_{m=1}^{k-1} \binom{k}{m} \binom{k-1}{m} \langle\langle 2m \rangle\rangle c_n\{2k-2m\}$$

$$c_n\{10\} = \langle\langle 10 \rangle\rangle - 25 \cdot \langle\langle 2 \rangle\rangle \langle\langle 8 \rangle\rangle - 100 \cdot \langle\langle 4 \rangle\rangle \langle\langle 6 \rangle\rangle + 400 \cdot \langle\langle 6 \rangle\rangle \langle\langle 2 \rangle\rangle^2 + 900 \cdot \langle\langle 2 \rangle\rangle \langle\langle 4 \rangle\rangle^2 - 3600 \cdot \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle^3 + 2800 \cdot \langle\langle 2 \rangle\rangle^5$$

$$v_n\{10\} = \sqrt[10]{\frac{1}{456} c_n\{10\}}$$

**First-time measurement by CMS -
huge amount of statistics!**

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Clear splitting between $v_2\{2\}$ and $v_2\{2k\}$ - larger towards more peripheral

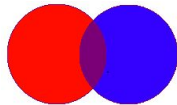
$$v_2\{2\} > v_2\{4\} \gtrsim v_2\{6\} \gtrsim v_2\{8\} \gtrsim v_2\{10\}$$

(v_2 variance)

Flow fluctuations : $v_2\{2\}^2 \approx v_2\{2k\}^2 + 2\sigma_v^2 (k > 1)$

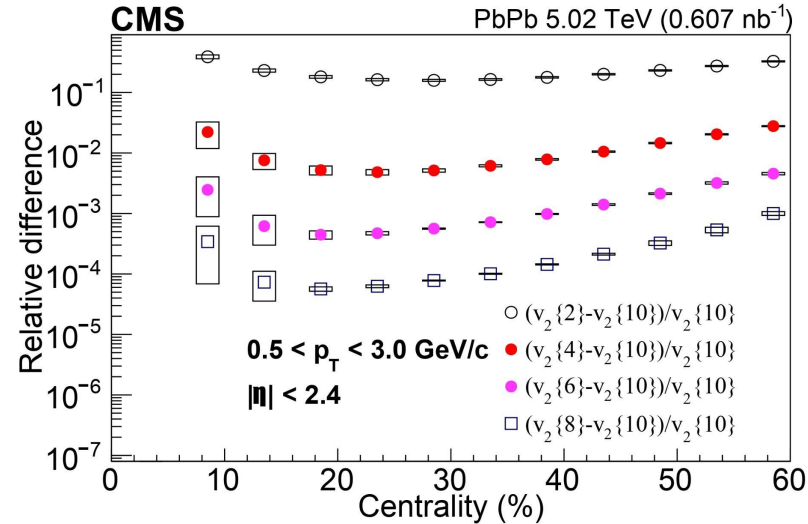
- systematic uncertainties ~ 2 orders of magnitude greater than statistical ones - **high precision measurement**

Decreasing overlap region



First-time measurement of $v_2\{10\}$ by CMS

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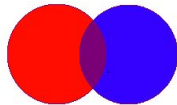
Fine splitting observed

$$v_2\{2\} > v_2\{4\} \gtrsim v_2\{6\} \gtrsim v_2\{8\} \gtrsim v_2\{10\}$$

Signature of non-Gaussian fluctuations

- The relative difference between adjacent $v_2\{2k\} - v_2\{10\}$ values decreases by about an order of magnitude for each increment in k

Decreasing overlap region



- $v_2\{2k\}$: Taylor expansion through **central moments** of the v_2 distribution (upto **5th moment**)

PRC 95 (2017) 014913, (mean, variance, skewness, kurtosis, superkurtosis)

PRC 99 (2019) 014907

+

Assuming $\sigma_x^2 = \sigma_y^2 \implies$ **hydrodynamic probes in terms of $v_2\{2k\}$**

$$1. \quad \frac{v_2\{6\} - v_2\{8\}}{v_2\{4\} - v_2\{6\}} \approx \frac{1}{11} - \frac{1}{11} \frac{v_2\{4\}^2 - 12v_2\{6\}^2 + 11v_2\{8\}^2}{v_2\{4\}^2 - v_2\{6\}^2 + \frac{(\sigma_y^2 - \sigma_x^2)s_{30}}{3v_2^3}}$$

neglecting this term

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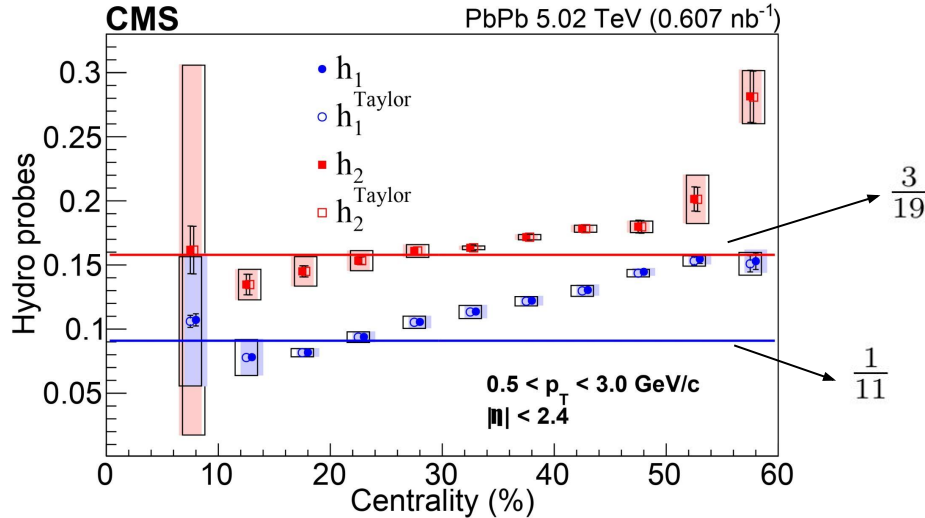
$$2. \quad \frac{v_2\{8\} - v_2\{10\}}{v_2\{6\} - v_2\{8\}} \approx \frac{3}{19} - \frac{1}{19} \frac{3v_2\{6\}^2 - 22v_2\{8\}^2 + 19v_2\{10\}^2}{v_2\{6\}^2 - v_2\{8\}^2 + \frac{(\sigma_y^2 - \sigma_x^2)s_{30}}{33v_2^3}}$$

introduction of higher order term

First-time measurement by CMS

neglecting this term

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$$h_1 = \frac{v_2\{6\} - v_2\{8\}}{v_2\{4\} - v_2\{6\}} \quad h_2 = \frac{v_2\{8\} - v_2\{10\}}{v_2\{6\} - v_2\{8\}}$$

Expansion using higher-order moments :

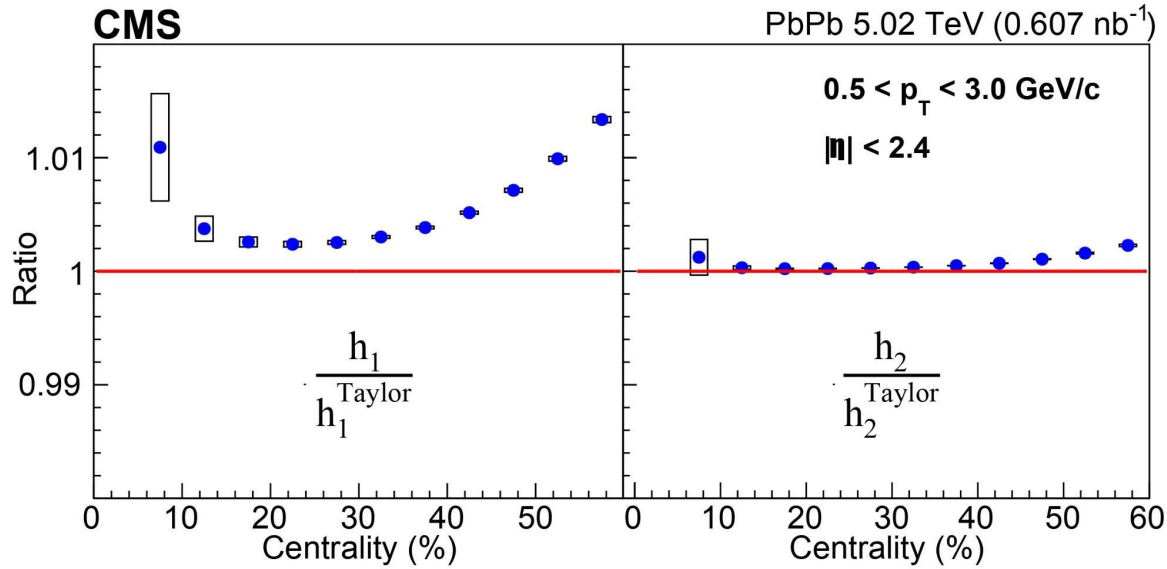
$$h_1^{Taylor} = \frac{1}{11} - \frac{1}{11} \frac{v_2\{4\}^2 - 12v_2\{6\}^2 + 11v_2\{8\}^2}{v_2\{4\}^2 - v_2\{6\}^2}$$

$$h_2^{Taylor} = \frac{3}{19} - \frac{1}{19} \frac{3v_2\{6\}^2 - 22v_2\{8\}^2 + 19v_2\{10\}^2}{v_2\{6\}^2 - v_2\{8\}^2}$$

**Clear dependence on centrality :
Higher-order moments necessary to
describe data**

First-time measurement by CMS

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$$\frac{h_1}{h_1^{Taylor}} \approx 1.000 \pm 0.013$$

$$\frac{h_2}{h_2^{Taylor}} \approx 1.000 \pm 0.003$$

small contribution from the term $(\sigma_x^2 - \sigma_y^2)$, but still negligible

- Systematic uncertainties \gg statistical uncertainties
- $\frac{h_2}{h_2^{Taylor}}$ is closer to unity

- Non-Gaussian fluctuations explain fine-splitting : $v_2\{2\} > v_2\{4\} \gtrsim v_2\{6\} \gtrsim v_2\{8\} \gtrsim v_2\{10\}$

- “Standardized” skewness : $\gamma_1^{\text{exp}} = -2^{3/2} \frac{v_2^3\{4\} - v_2^3\{6\}}{[v_2^2\{2\} - v_2^2\{4\}]^{3/2}} \approx -2^{3/2} \frac{-s_{30} - O_N}{[2\sigma_x^2 + O_D]^{3/2}} \approx \frac{s_{30}}{\sigma_x^3} = \gamma_1$

- Contributions from other moments :

Phys. Rev. C 99 (2019) 014907

$$O_N = \frac{3(\kappa_{40} + \kappa_{22})}{2\bar{v}_2} - \frac{3(p_{50} + 2p_{32} + p_{14})}{4\bar{v}_2^2} + \frac{3(\sigma_y^2 - \sigma_x^2)(s_{30} - 2s_{12})}{2\bar{v}_2^2} + \dots$$

non-negligible \swarrow

$$O_D = \frac{2}{\bar{v}_2} (s_{30} + s_{12}) + \frac{\kappa_{40} + 2\kappa_{22} + \kappa_{04}}{2\bar{v}_2^2} + \frac{(\sigma_y^2 - \sigma_x^2)^2}{\bar{v}_2^2} - 2 \frac{(\sigma_y^2 - \sigma_x^2)(s_{30} - s_{12})}{\bar{v}_2^3} + \dots$$

- “Corrected” skewness : free from contributions of other moments (eg. kurtosis, superskewness, ...)

$$\gamma_{1,corr}^{\text{exp}} = -2^{3/2} \frac{-s_{30} + 3 \frac{(\sigma_y^2 - \sigma_x^2)(s_{30} - 2s_{12})}{2\bar{v}_2^2} + o(>5)}{(2\sigma_x^2 + \frac{(\sigma_y^2 - \sigma_x^2)^2}{\bar{v}_2^2} + o(>5))^{3/2}} = -2^{3/2} \frac{187v_2^3\{8\} - 16v_2^3\{6\} - 171v_2^3\{10\} \checkmark}{[v_2^2\{2\} - 40v_2^2\{6\} + 495v_2^2\{8\} - 456v_2^2\{10\}]^{3/2}}$$

Free of other moments (up to 5th order)

- “Standardized” kurtosis :

$$\gamma_2^{exp} = -\frac{3}{2} \frac{v_2^4\{4\} - 12v_2^4\{6\} + 11v_2^4\{8\}}{[v_2^2\{2\} - v_2^2\{4\}]^2}$$

- “Corrected” kurtosis :

$$\gamma_{2,corr}^{exp} = -\frac{3}{2} \frac{v_2^4\{4\} + 24v_2^4\{6\} - 253v_2^4\{8\} + 228v_2^4\{10\}}{[v_2^2\{2\} - 40v_2^2\{6\} + 495v_2^2\{8\} - 456v_2^2\{10\}]^2}$$

- “Standardized” superskewness : $\gamma_3^{exp} = 6\sqrt{2} \frac{3v_2^5\{6\} - 22v_2^5\{8\} + 19v_2^5\{10\}}{[v_2^2\{2\} - v_2^2\{4\}]^{5/2}}$

- “Corrected” superskewness :

$$\gamma_{3,corr}^{exp} = 6\sqrt{2} \frac{3v_2^5\{6\} - 22v_2^5\{8\} + 19v_2^5\{10\}}{[v_2^2\{2\} - 40v_2^2\{6\} + 495v_2^2\{8\} - 456v_2^2\{10\}]^{5/2}}$$

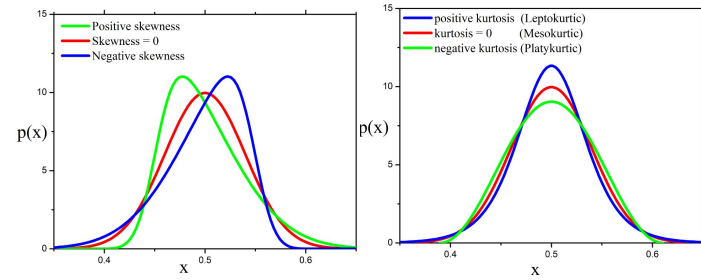
- **Free of other moments (up to 5th order)**

- **Additional constraints on initial-state geometry - “Cleaning” conditions require elliptic power distribution, with :**

$$\varepsilon_0 < 0.15 \text{ and}$$

$$V_n \propto \varepsilon_n$$

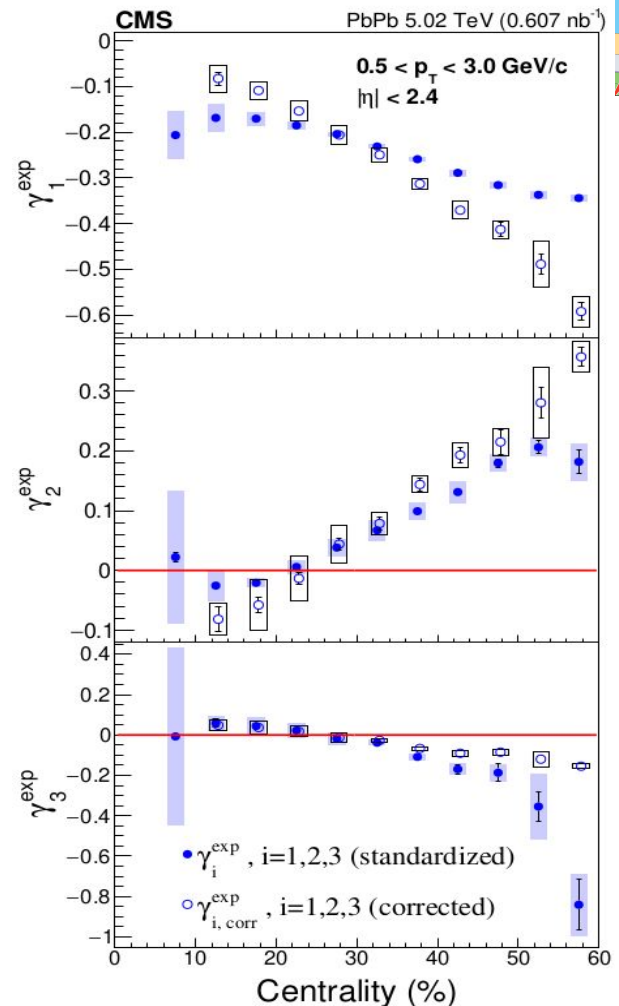
[Phys. Rev. C 90,
024903 (2014)]

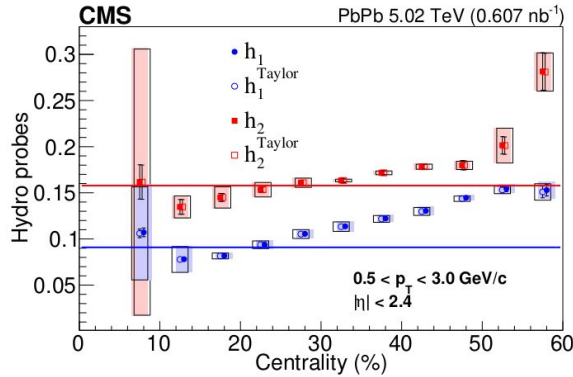


First-time measurement by CMS

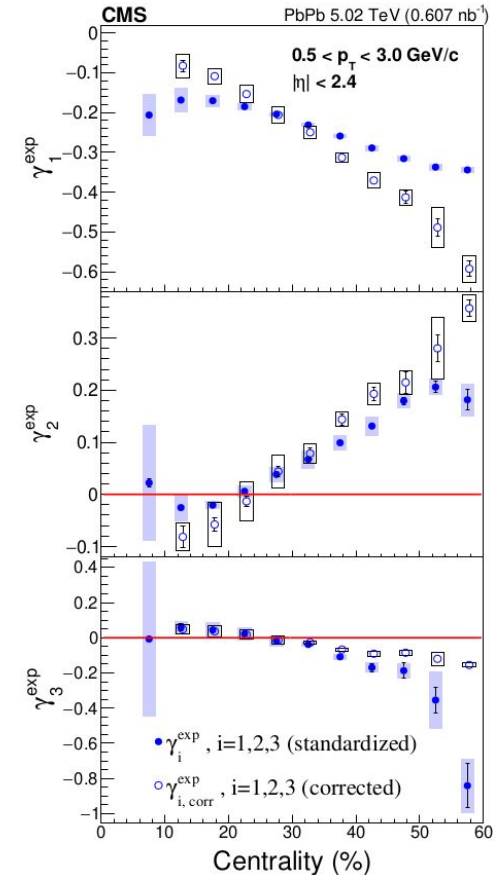
JHEP02 (2024) 106

- Negative values of skewness for all centralities - v_2 distribution has longer tail to the left
 - corrected skewness is steeper
- Kurtosis - negative for most central events, positive towards peripheral
 - qualitatively agrees with theory predictions [[Phys. Rev. C 99, 014907 \(2019\)](#)]
- Superskewness - measured for the first time
 - positive towards more central region, then becomes negative



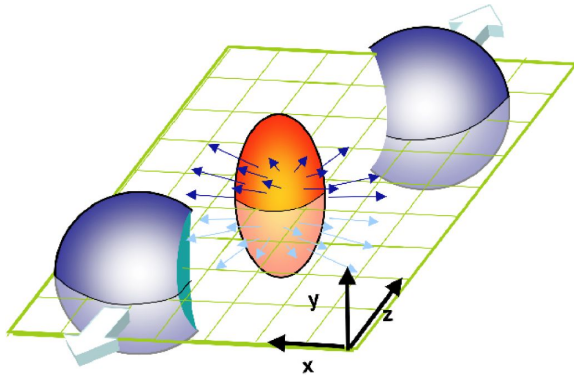


- Two hydrodynamics probes and first-time measurement of $v_2\{10\}$ performed with CMS PbPb data at 5.02 TeV energy
- High precision measurement of skewness, kurtosis and superskewness of the v_2 distribution
- Can provide novel constraints on the initial state geometry used in hydrodynamic calculations of the medium expansion in high energy nuclear collisions





Backup



$\langle\langle \dots \rangle\rangle$ Averaging over all events

Multi-particle correlations \implies **cumulants**

$$\langle\langle 2 \rangle\rangle = \langle\langle e^{in(\phi_1 - \phi_2)} \rangle\rangle \quad \langle\langle 4 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle\rangle$$

$$\langle\langle 6 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 + \phi_3 - \phi_4 - \phi_5 - \phi_6)} \rangle\rangle$$

$$\langle\langle 8 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 + \phi_3 + \phi_4 - \phi_5 - \phi_6 - \phi_7 - \phi_8)} \rangle\rangle$$

$$c_n\{2\} = \langle\langle 2 \rangle\rangle \quad c_n\{4\} = \langle\langle 4 \rangle\rangle - 2 \cdot \langle\langle 2 \rangle\rangle^2$$

$$c_n\{6\} = \langle\langle 6 \rangle\rangle - 9 \cdot \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle + 12 \cdot \langle\langle 2 \rangle\rangle^3$$

$$c_n\{8\} = \langle\langle 8 \rangle\rangle - 16 \cdot \langle\langle 6 \rangle\rangle \langle\langle 2 \rangle\rangle - 18 \cdot \langle\langle 4 \rangle\rangle^2 + 144 \cdot \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle^2 - 144 \cdot \langle\langle 2 \rangle\rangle^4$$

$$v_n\{2\} = \sqrt{c_n\{2\}} \quad v_n\{4\} = \sqrt[4]{-c_n\{4\}}$$

$$v_n\{6\} = \sqrt[6]{\frac{1}{4} c_n\{6\}} \quad v_n\{8\} = \sqrt[8]{-\frac{1}{33} c_n\{8\}}$$

$$\langle 2m \rangle = \frac{1}{P_{M,2m}} \sum_{i_1 \neq \dots \neq i_{2m} = 1}^M e^{in(\phi_{i_1} + \dots + \phi_{i_m} - \phi_{i_{m+1}} - \dots - \phi_{i_{2m}})}$$

$$P_{M,2m} = \frac{M!}{(M-2m)!}$$

Averaging over M particles in a single event

General formulae :

$$c_n\{2k\} = \langle\langle 2k \rangle\rangle - \sum_{m=1}^{k-1} \binom{k}{m} \binom{k-1}{m} \langle\langle 2m \rangle\rangle c_n\{2k-2m\}$$

$$v_n = \sqrt[2k]{a_{2k}^{-1} c_n\{2k\}}.$$

$$a_{2k} = 1 - \sum_{m=1}^{k-1} \binom{k}{m} \binom{k-1}{m} a_{2k-2m}, \quad \text{with } a_2 = 1.$$

Phys. Rev. C 104 (2021)
034906
Phys. Rev. C 95 (2017)
014913
Phys.Rev.C 64 (2001)
054901
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Example :

$$c_n\{2\} = \langle\langle 2 \rangle\rangle$$

$$c_n\{4\} = \langle\langle 4 \rangle\rangle - 2 \langle\langle 2 \rangle\rangle^2$$

$$v_n\{2\} = \sqrt[2]{c_n\{2\}}$$

$$v_n\{4\} = \sqrt[4]{-c_n\{4\}}$$

Can be extended till any
order

$$c_n\{10\} = \langle\langle 10 \rangle\rangle - 25 \cdot \langle\langle 2 \rangle\rangle \langle\langle 8 \rangle\rangle - 100 \cdot \langle\langle 4 \rangle\rangle \langle\langle 6 \rangle\rangle + 400 \cdot \langle\langle 6 \rangle\rangle \langle\langle 2 \rangle\rangle^2 + 900 \cdot \langle\langle 2 \rangle\rangle \langle\langle 4 \rangle\rangle^2 - 3600 \cdot \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle^3 + 2880 \cdot \langle\langle 2 \rangle\rangle^5$$

$$v_n\{10\} = \sqrt[10]{\frac{1}{456} c_n\{10\}}$$

**First-time measurement by CMS -
huge amount of statistics!**

- size of formula increases with the order k

$$c_n \{10\} = \langle\langle 10 \rangle\rangle - 25 \cdot \langle\langle 2 \rangle\rangle \langle\langle 8 \rangle\rangle - 100 \cdot \langle\langle 4 \rangle\rangle \langle\langle 6 \rangle\rangle + 400 \cdot \langle\langle 6 \rangle\rangle \langle\langle 2 \rangle\rangle^2 + 900 \cdot \langle\langle 2 \rangle\rangle \langle\langle 4 \rangle\rangle^2 - 3600 \cdot \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle^3 + 2880 \cdot \langle\langle 2 \rangle\rangle^5$$

$$v_n \{10\} = \sqrt[10]{\frac{1}{456} c_n \{10\}}$$

$$\begin{aligned} \langle 10 \rangle = & \frac{|Q_n|^{10} - 20\text{Re}[Q_{2n}|Q_n|^6 Q_n^4 Q_n^2] + 100|Q_{2n}|^2 |Q_n|^6 + 30\text{Re}[Q_{2n} Q_{2n} |Q_n|^2 Q_n^4 Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{225|Q_{2n}|^4 |Q_n|^2 - 300\text{Re}[Q_{2n}|Q_{2n}|^2 |Q_n|^2 Q_n^4 Q_n^2] + 40\text{Re}[Q_{3n}|Q_n|^4 Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{600\text{Re}[Q_{3n} Q_n |Q_n|^2 Q_{2n}^2] - 400\text{Re}[Q_{3n}|Q_n|^4 Q_n^2 Q_{2n}] + 40|Q_{3n}|^2 |Q_n|^4}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{400\text{Re}[Q_{3n}|Q_{2n}|^2 Q_n^2 Q_n^2] - 40\text{Re}[Q_{3n} Q_{2n} Q_n^2 Q_n^2 Q_n^2] - 600\text{Re}[Q_{3n}|Q_{2n}|^2 Q_n^2 Q_{2n}]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{400|Q_{3n}|^2 |Q_{2n}|^2 - 800\text{Re}[Q_{3n}|^2 Q_{2n} Q_n^2 Q_n^2] - 60\text{Re}[Q_{4n}|Q_n|^2 Q_n^2 Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{600\text{Re}[Q_{4n}|Q_n|^2 Q_n^2 Q_n^2 Q_n^2] - 900\text{Re}[Q_{4n}|Q_n|^2 Q_{2n} Q_{2n}] - 1200\text{Re}[Q_{4n}|Q_n|^2 Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{1200\text{Re}[Q_{4n} Q_n Q_{2n} Q_n^2] + 900|Q_{4n}|^2 |Q_n|^2 + 48\text{Re}[Q_{5n} Q_n^2 Q_n^2 Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{720\text{Re}[Q_{5n} Q_n^2 Q_{2n} Q_{2n}] - 480\text{Re}[Q_{5n} Q_n^2 Q_n^2 Q_n^2 Q_{2n}] + 960\text{Re}[Q_{5n} Q_n^2 Q_n^2 Q_{3n}]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{576|Q_{5n}|^2 - 960\text{Re}[Q_{5n} Q_n^2 Q_n^2] - 1440\text{Re}[Q_{5n} Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-8)(M-9)} + \\ & \frac{300\text{Re}[Q_{2n}|Q_n|^4 Q_n^2 Q_n^2] - 25|Q_n|^{18} - 900|Q_{2n}|^2 |Q_n|^4 - 150\text{Re}[Q_{2n} Q_{2n} Q_n^2 Q_n^2 Q_n^2]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-9)} + \\ & \frac{900\text{Re}[Q_{2n}|Q_{2n}|^2 Q_n^2 Q_n^2] - 225|Q_{2n}|^4 - 400\text{Re}[Q_{3n}|Q_n|^2 Q_n^2 Q_n^2] + 2400\text{Re}[Q_{3n}|Q_n|^2 Q_n^2 Q_{2n}]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-9)} + \\ & \frac{300\text{Re}[Q_{4n} Q_n^2 Q_n^2 Q_n^2] - 1200\text{Re}[Q_{3n} Q_n Q_{2n} Q_{2n}] - 1600|Q_{3n}|^2 |Q_n|^2 - 1800\text{Re}[Q_{4n} Q_n^2 Q_n^2 Q_{2n}]}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-9)} + \\ & \frac{900\text{Re}[Q_{4n} Q_{2n} Q_{2n}] + 2400\text{Re}[Q_{4n} Q_n^2 Q_n^2] - 900|Q_{4n}|^2}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-6)(M-7)(M-9)} + \\ & \frac{200|Q_n|^6 - 1200\text{Re}[Q_{2n}|Q_n|^2 Q_n^2 Q_n^2] + 1800|Q_{2n}|^2 |Q_n|^2}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-7)(M-8)} + \\ & \frac{800\text{Re}[Q_{3n} Q_n^2 Q_n^2 Q_n^2] - 2400\text{Re}[Q_{3n} Q_n^2 Q_n^2] + 800|Q_{3n}|^2}{M(M-1)(M-2)(M-3)(M-4)(M-5)(M-7)(M-8)} + \\ & \frac{1200\text{Re}[Q_{2n} Q_n^2 Q_n^2] - 600|Q_n|^4 - 600|Q_{2n}|^2}{M(M-1)(M-2)(M-3)(M-5)(M-6)(M-7)} + \\ & \frac{600|Q_n|^2}{M(M-1)(M-3)(M-4)(M-5)(M-6)} - \frac{120}{(M-1)(M-2)(M-3)(M-4)(M-5)} \end{aligned}$$

- Moments of v_2 distribution :

- Variance - 2nd moment $\sigma_x^2 = \langle (v_x - \langle v_x \rangle)^2 \rangle$, $\sigma_y^2 = \langle v_y^2 \rangle$ - (i)

- Skewness - 3rd moment $s_{30} = \langle (v_x - \langle v_x \rangle)^3 \rangle$, $s_{12} = \langle (v_x - \langle v_x \rangle) v_y^2 \rangle$ - (ii)

- Kurtosis - 4th moment $\kappa_{40} = \langle (v_x - \langle v_x \rangle)^4 \rangle$, $\kappa_{22} = \langle (v_x - \langle v_x \rangle)^2 v_y^2 \rangle - \sigma_x^2 \sigma_y^2$ - (iii)

Phys. Rev. C 95, 014913 (2017), Phys. Rev. C 64, 054901 (2001) :

$$v_2 \{4\} \approx \bar{v}_2 + \frac{\sigma_y^2 - \sigma_x^2}{2\bar{v}_2} - \frac{s_{30} + s_{12}}{\bar{v}_2^2} - \frac{\kappa_{40} + 2\kappa_{22} + \kappa_{04}}{4\bar{v}_2^3} - \frac{5(\sigma_y^2 - \sigma_x^2)^2}{8\bar{v}_2^3} + \frac{(\sigma_y^2 - \sigma_x^2)(3s_{30} + 3s_{12})}{2\bar{v}_2^4} \quad - \text{(iv)}$$

$$v_2 \{6\} \approx \bar{v}_2 + \frac{\sigma_y^2 - \sigma_x^2}{2\bar{v}_2} - \frac{2}{3} \frac{s_{30} + s_{12}}{\bar{v}_2^2} + \frac{\kappa_{40} - \kappa_{04}}{4\bar{v}_2^3} - \frac{5(\sigma_y^2 - \sigma_x^2)^2}{8\bar{v}_2^3} + \frac{p_{50} + 2p_{32} + p_{14}}{4\bar{v}_2^4} + \frac{(\sigma_y^2 - \sigma_x^2)(4s_{30} + 15s_{12})}{6\bar{v}_2^4} \quad - \text{(v)}$$

$$v_2 \{8\} \approx \bar{v}_2 + \frac{\sigma_y^2 - \sigma_x^2}{2\bar{v}_2} - \frac{7}{11} \frac{s_{30} + s_{12}}{\bar{v}_2^2} + \frac{31}{33} \kappa_{40} + \frac{2}{11} \kappa_{22} - \kappa_{04} - \frac{5(\sigma_y^2 - \sigma_x^2)^2}{8\bar{v}_2^3} + \frac{5}{3} p_{50} + \frac{14}{3} p_{32} + 3p_{14} + \frac{(\sigma_y^2 - \sigma_x^2)(13s_{30} + 57s_{12})}{22\bar{v}_2^4} \quad - \text{(vi)}$$

$$\frac{v_2\{6\} - v_2\{8\}}{v_2\{4\} - v_2\{6\}} \approx \frac{1}{11} \left(1 - \frac{4\kappa_{40} + \frac{8(p_{50} + p_{32})}{\bar{v}_2}}{2\bar{v}_2 s_{30} + 3(\kappa_{40} + \kappa_{22}) + \frac{3(p_{50} + 2p_{32} + p_{14}) - 2(\sigma_y^2 - \sigma_x^2)(5s_{30} - 6s_{12})}{2\bar{v}_2}} \right)$$

Using eqs. (iv), (v) and (vi) :

$$\frac{v_2\{6\} - v_2\{8\}}{v_2\{4\} - v_2\{6\}} \approx \frac{1}{11} - \frac{1}{11} \frac{v_2\{4\}^2 - 12v_2\{6\}^2 + 11v_2\{8\}^2}{v_2\{4\}^2 - v_2\{6\}^2 + \frac{(\sigma_y^2 - \sigma_x^2)s_{30}}{3v_2^3}}$$

\downarrow
 h_1
 \downarrow
 h_1^{Taylor}

negligible

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■ 5th moment -

$$p_{50} = \left\langle \left(v_x - \langle v_x \rangle \right)^5 \right\rangle - 10\sigma_x^2 s_{30}$$

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$$p_{32} = \left\langle \left(v_x - \langle v_x \rangle \right)^3 v_y^2 \right\rangle - \sigma_y^2 s_{30} - 3\sigma_x^2 s_{12}$$

$$p_{14} = \left\langle \left(v_x - \langle v_x \rangle \right) v_y^4 \right\rangle - 6\sigma_y^2 s_{12}$$

Phys. Rev. C 64,
054901 (2001)

$$v_2 \{10\} \approx \bar{v}_2 + \frac{\sigma_y^2 - \sigma_x^2}{2\bar{v}_2} - \frac{12}{19} \frac{s_{30} + s_{12}}{\bar{v}_2^2} + \frac{53}{57} \frac{\kappa_{40} + \frac{4}{19} \kappa_{22} - \kappa_{04}}{4\bar{v}_2^3} - \frac{5(\sigma_y^2 - \sigma_x^2)^2}{8\bar{v}_2^3} + \frac{163}{60} \frac{p_{50}}{19\bar{v}_2^4} + \frac{47}{6} \frac{p_{32}}{19\bar{v}_2^4} + \frac{21}{4} \frac{p_{14}}{19\bar{v}_2^4} + \frac{(\sigma_y^2 - \sigma_x^2) \left(11s_{30} + \frac{99}{2} s_{12} \right)}{19\bar{v}_2^4}$$

$$\frac{v_2 \{8\} - v_2 \{10\}}{v_2 \{6\} - v_2 \{8\}} \approx \frac{3}{19} - \frac{88p_{50}}{95 \left[4\bar{v}_2^2 s_{30} - 2\bar{v}_2 (\kappa_{40} - 3\kappa_{22}) - 13(p_{50} + 10p_{32} - 3p_{14}) - 2(\sigma^2 - \sigma_x^2)(5s_{30} - 6s_{32}) \right]}$$

**New
hydrodynamic
probe :**

$$\frac{v_2 \{8\} - v_2 \{10\}}{v_2 \{6\} - v_2 \{8\}} \approx \frac{3}{19} - \frac{1}{19} \frac{3v_2 \{6\}^2 - 22v_2 \{8\}^2 + 19v_2 \{10\}^2}{v_2 \{6\}^2 - v_2 \{8\}^2 + \frac{(\sigma_y^2 - \sigma_x^2) s_{30}}{33\bar{v}_2^3}}$$

h_2

h_2^{Taylor}

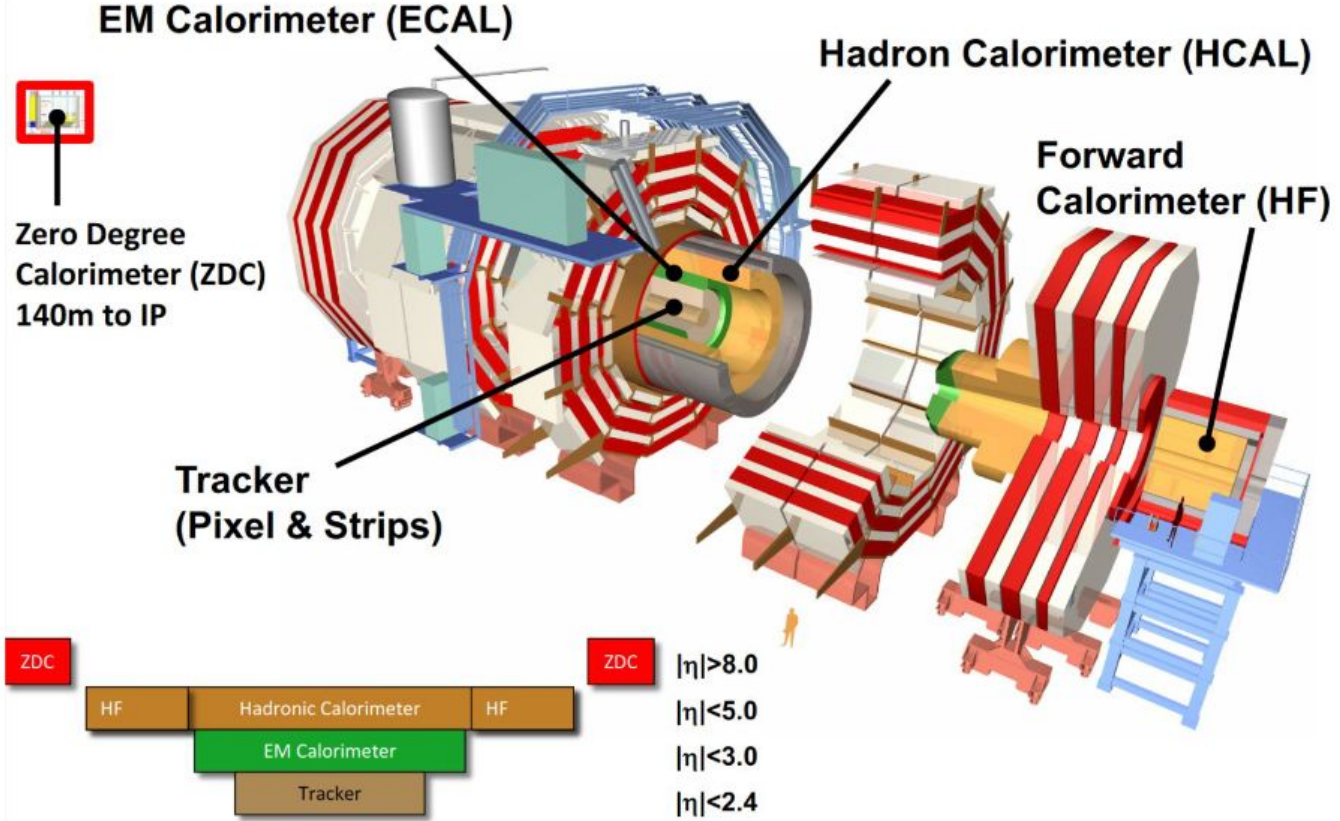
- Condition for “cleaning” : $s_{12} \approx \frac{s_{30}}{3} \quad \kappa_{22} \approx \frac{\kappa_{40}}{3} \quad p_{32} \approx p_{14} \approx \frac{p_{50}}{5}$

$$\varepsilon_0 \equiv \langle y_j^2 - x_j^2 \rangle / \langle y_j^2 + x_j^2 \rangle$$

Elliptic power distribution, ellipticity parameter : $\varepsilon_0 < 0.15$

Phys. Rev. C 90, 024903 (2014)

$$p(\varepsilon_x, \varepsilon_y) = \frac{\alpha}{\pi} (1 - \varepsilon_0^2)^{\alpha + \frac{1}{2}} \frac{(1 - \varepsilon_x^2 - \varepsilon_y^2)^{\alpha - 1}}{(1 - \varepsilon_0 \varepsilon_x)^{2\alpha + 1}}$$





- **Data**

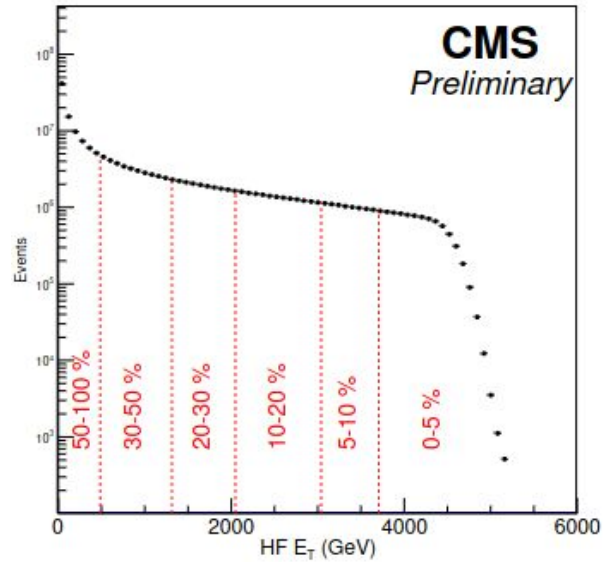
- 2018 PbPb Minimum Bias events

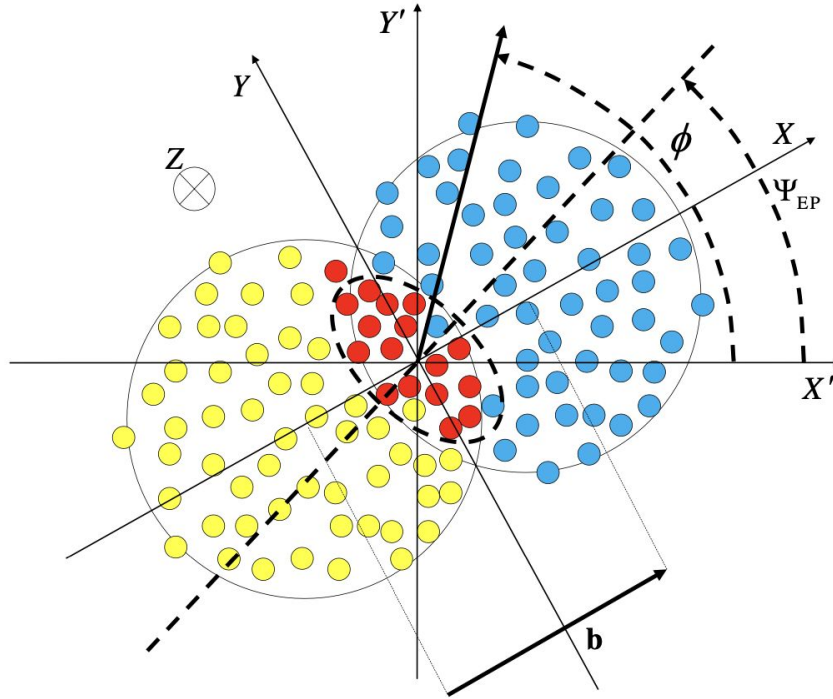
- **Event selections**

- primaryVertexFilter
- clusterCompatibilityFilter
- hfCoincFilter2Th4

- **Track selections**

- packedPFCandidates
- $0.5 < p_T < 3.0 \text{ GeV}/c$
- $|\eta| < 2.4$
- highPurity
- $DCA < 3.0$
- $N_{\text{hits}} \geq 11$
- $dp_T/p_T < 0.1$
- $\chi^2/\text{ndof}/N_{\text{layers}} < 0.18$





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