

Massive diphoton production at NNLO

Federico Coro - ICHEP 2024



VNIVERSITAT
DE VALÈNCIA



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Ongoing work

Outline of the talk

- ❖ Introduction
- ❖ Two-loop form factors
- ❖ MIs evaluation
- ❖ Phenomenological results
- ❖ Conclusions

Motivations

- ❖ Diphoton is an experimentally clean final state
- ❖ QCD background for Higgs
- ❖ Important to measure the fundamental parameters within the Standard Model
- ❖ Search for new physics

State of the art

❖ Massless NNLO QCD accuracy (5lf)

[S.Catani,L.Cieri,D.de Florian,G.Ferrera,M.Grazzini]

[J.M.Campbell,R.K.Ellis,Y.Li,C.Williams]

[R.Schuermann,X.Chen,T.Gehrmann,E.W.N.Glover,M.Höfer,A.Huss]

❖ Elements for N^3LO

[Z.Bern,A.De Freitas,L.J.Dixon]

[F.Caola,A.Chakraborty,G.Gambuti,A.Von Manteuffel,L.Tancredi]

[H.A.Chawdhry,M.Czakon,A.Mitov,R.Poncelet]

[B.Agarwal,F.Buccioni,A.Von Manteuffel,L.Tancredi]

[S.Badger,T.Gehrmann,M.Marcoli,R.Moodie]

❖ First order EW/QED

[M.Chiesa,N.Greiner,M.Schönherr,F.Tramontano]

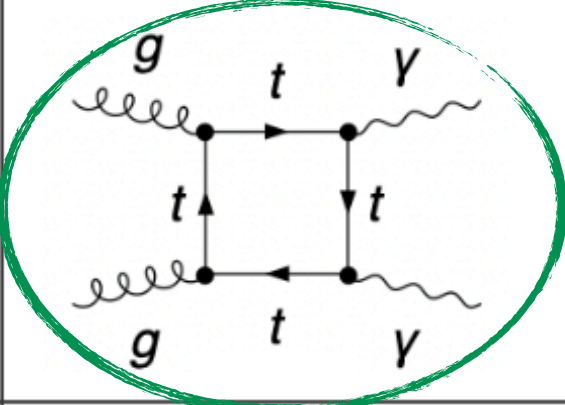

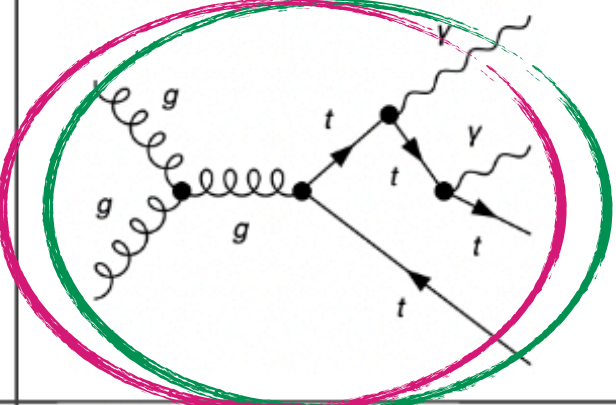
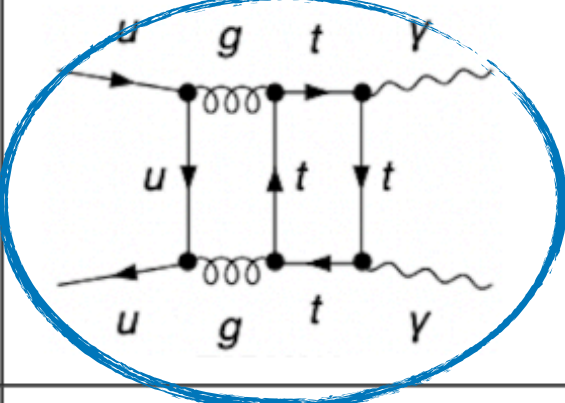
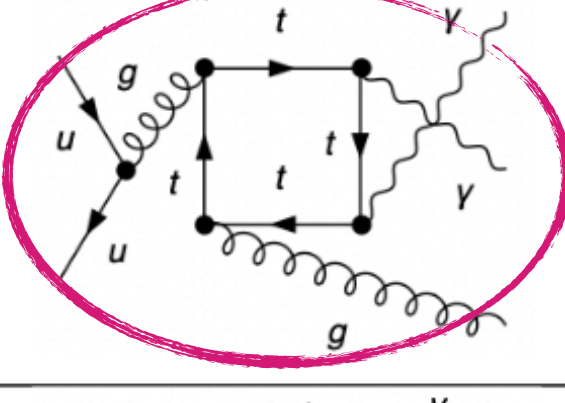
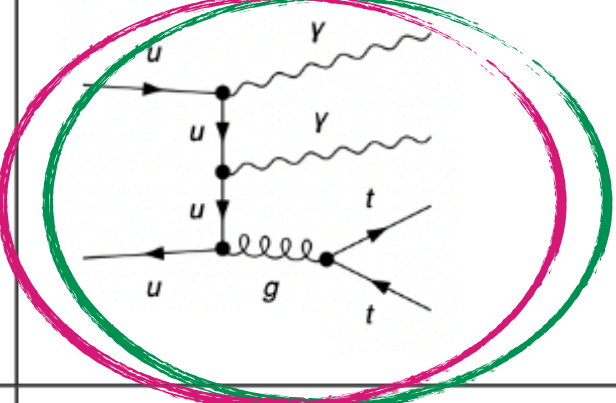

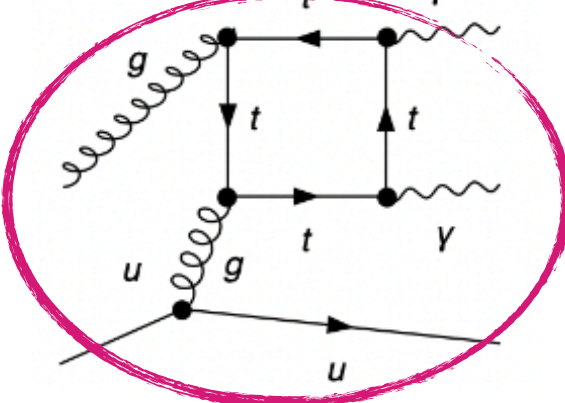

[L.Cieri,G.Sborlini]

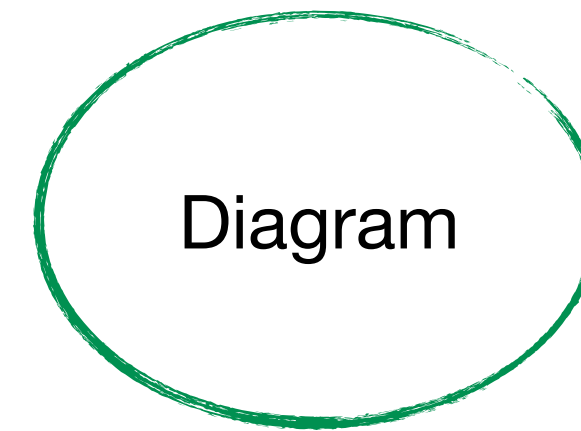
❖ Mass effect for $gg \rightarrow \gamma\gamma$

[F.Maltoni,M.K.Mandal,X.Zhao]

Massive Corrections

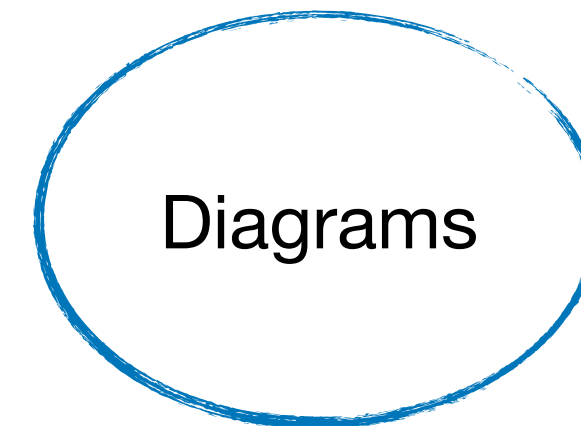
Massive corrections $\mathcal{O}(\alpha_s^2)$

Channels	$\gamma\gamma$	$\gamma\gamma j$	$\gamma\gamma jj$
gg			
$q\bar{q}$			
qg			

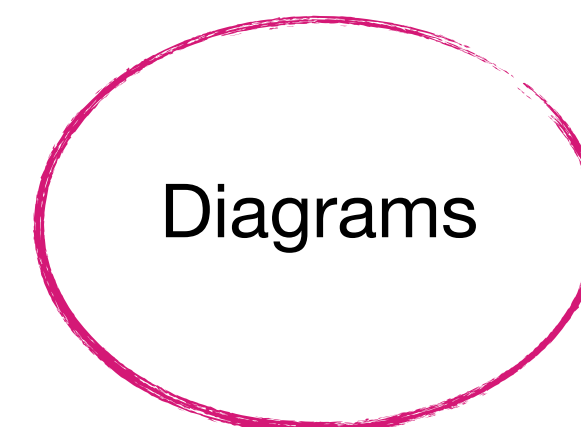


[J.M.Campbell,R.K.Ellis,Y.Li,C.Williams]

[F.Buccioni,J-N.Lang,J.M.Lindert,P.Maierhofer,
S.Pozzorini,H.Zhang,M.Zoller]

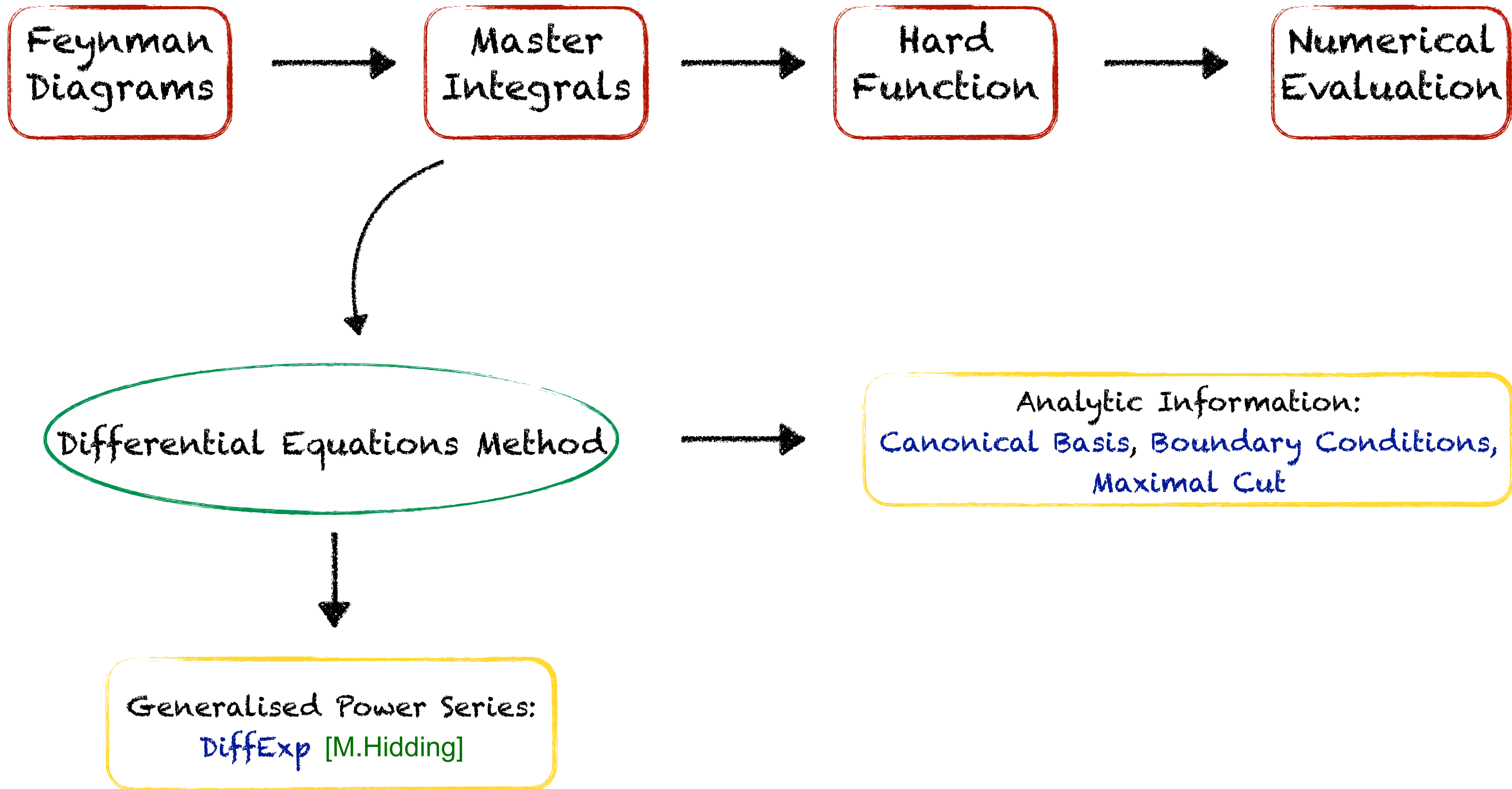


Original results and main focus of the talk



Evaluated for the final result

Computational pipeline



Form factors

The bare scattering amplitude $\mathcal{A}_{q\bar{q},\gamma\gamma} = \alpha_{em} \delta_{ij} \epsilon_{\lambda_3}^{*\mu}(p_3) \epsilon_{\lambda_4}^{*\nu}(p_4) \bar{v}_{s_2}(p_2) \mathcal{A}_{\mu\nu}(s, t, u, m_t^2) u_{s_1}(p_1)$

Can be decomposed in terms of a set of four independent tensors $\mathcal{A}_{q\bar{q},\gamma\gamma}(s, t, m_t^2) = \sum_{i=1}^4 \mathcal{F}_i(s, t, m_t^2) \bar{v}(p_2) \Gamma_i^{\mu\nu} u(p_1) \epsilon_{3,\mu} \epsilon_{4,\nu}$

$$\Gamma_1^{\mu\nu} = \gamma^\mu p_2^\nu, \Gamma_2^{\mu\nu} = \gamma^\nu p_1^\mu, \Gamma_3^{\mu\nu} = p_{3,\rho} \gamma^\rho p_1^\mu p_2^\nu, \Gamma_4^{\mu\nu} = p_{3,\rho} \gamma^\rho g^{\mu\nu}$$

[F.Caola,A.Von Manteuffel,L.Tancredi]

The form factors admits a perturbative expansion in α_s $\mathcal{F}_k = \mathcal{F}_k^{(0)} + \left(\frac{\alpha_s^B}{\pi}\right) \mathcal{F}_k^{(1)} + \left(\frac{\alpha_s^B}{\pi}\right)^2 \mathcal{F}_k^{(2)} + \dots$

Massive contribution appears at $\mathcal{O}(\alpha_s^2)$

$$\mathcal{F}_k^{(2)} = \delta_{ij} C_F (4\pi\alpha_{em}) \left[Q_q^2 \mathcal{F}_{k,top;0}^{(2)} + Q_t^2 \mathcal{F}_{k,top;2}^{(2)} \right]$$

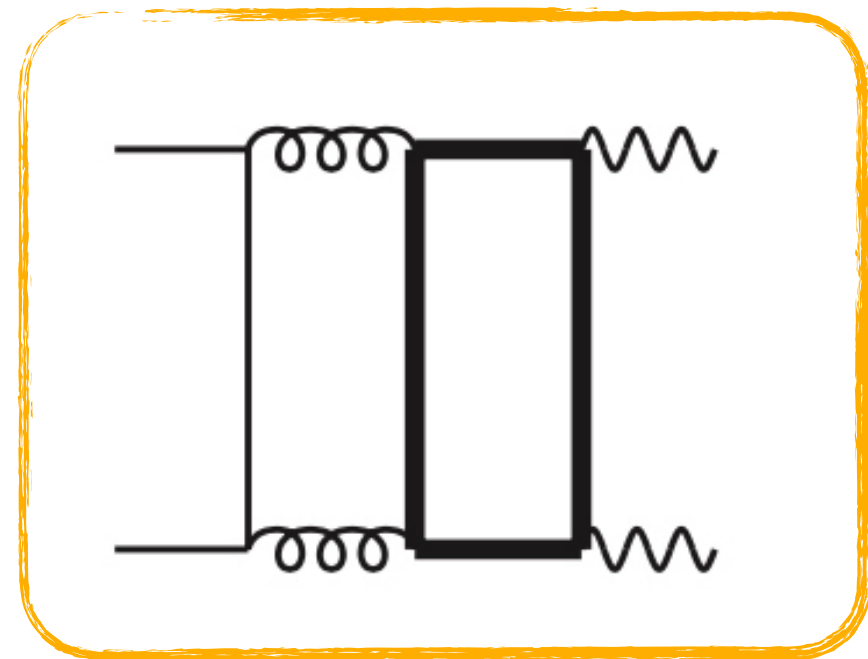
Q_q is the charge of light quark
 Q_t is the charge of heavy quark

Two-loop Feynman diagrams

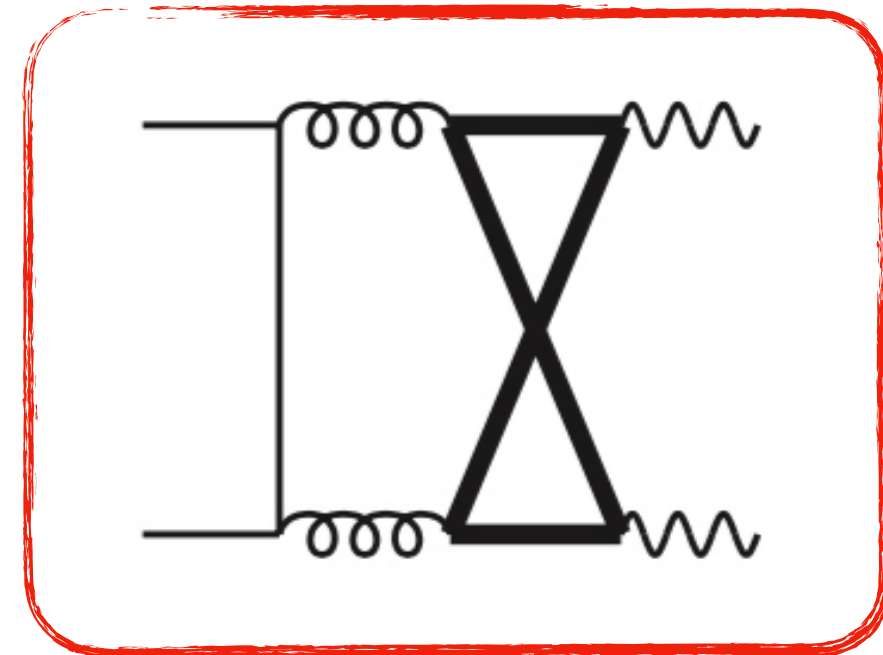
At partonic level the scattering process is: $q(p_1) + \bar{q}(p_2) \rightarrow \gamma(p_3) + \gamma(p_4)$

External particles on-shell and the top quark running in the loop

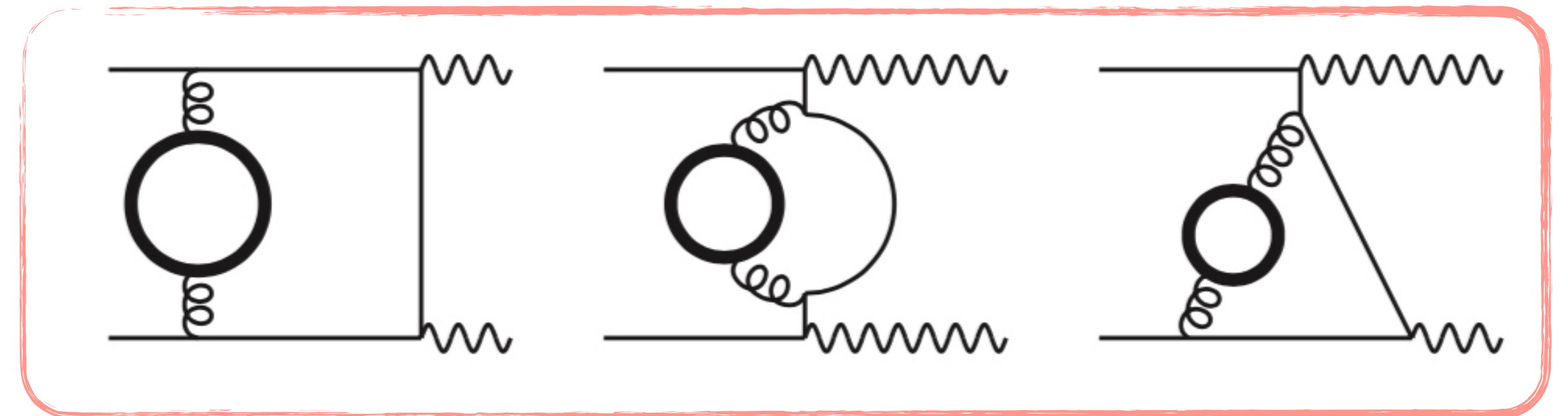
Feynman diagrams generated with **QGRAF** [P. Nogueira] and **FeynArts** [T.Hahn]



PLA



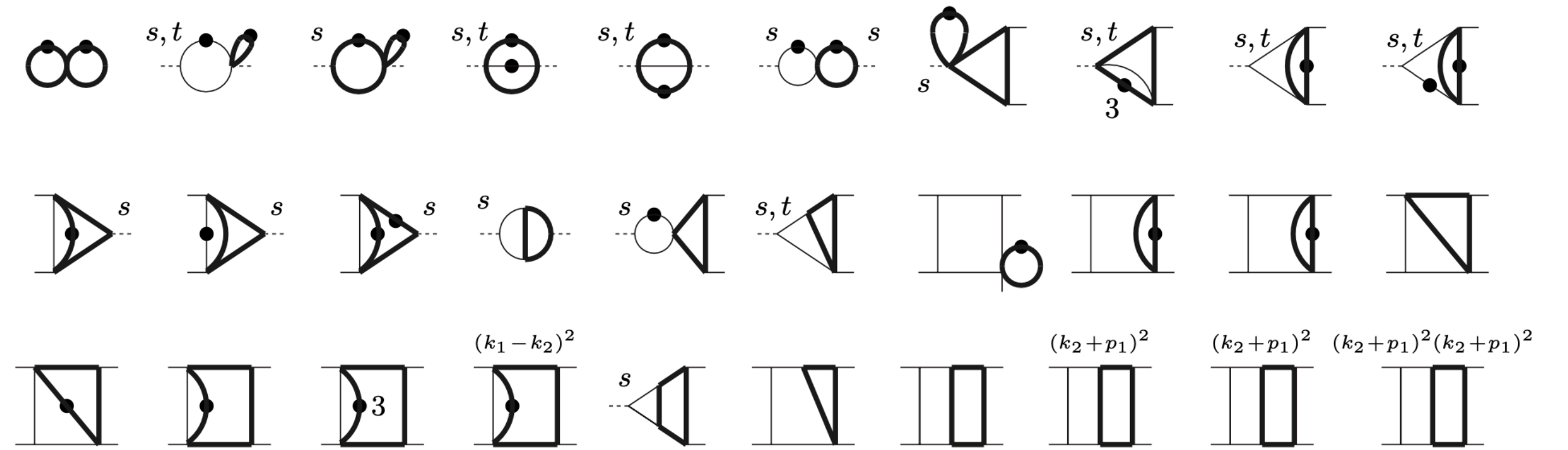
NPL



PLB

Master Integrals

PLA and PLB Master Integrals

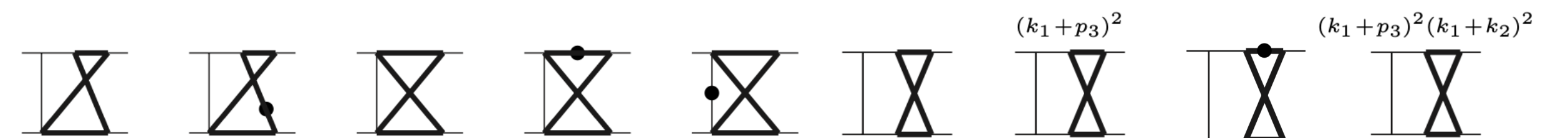


[M.Becchetti,R.Bonciani]



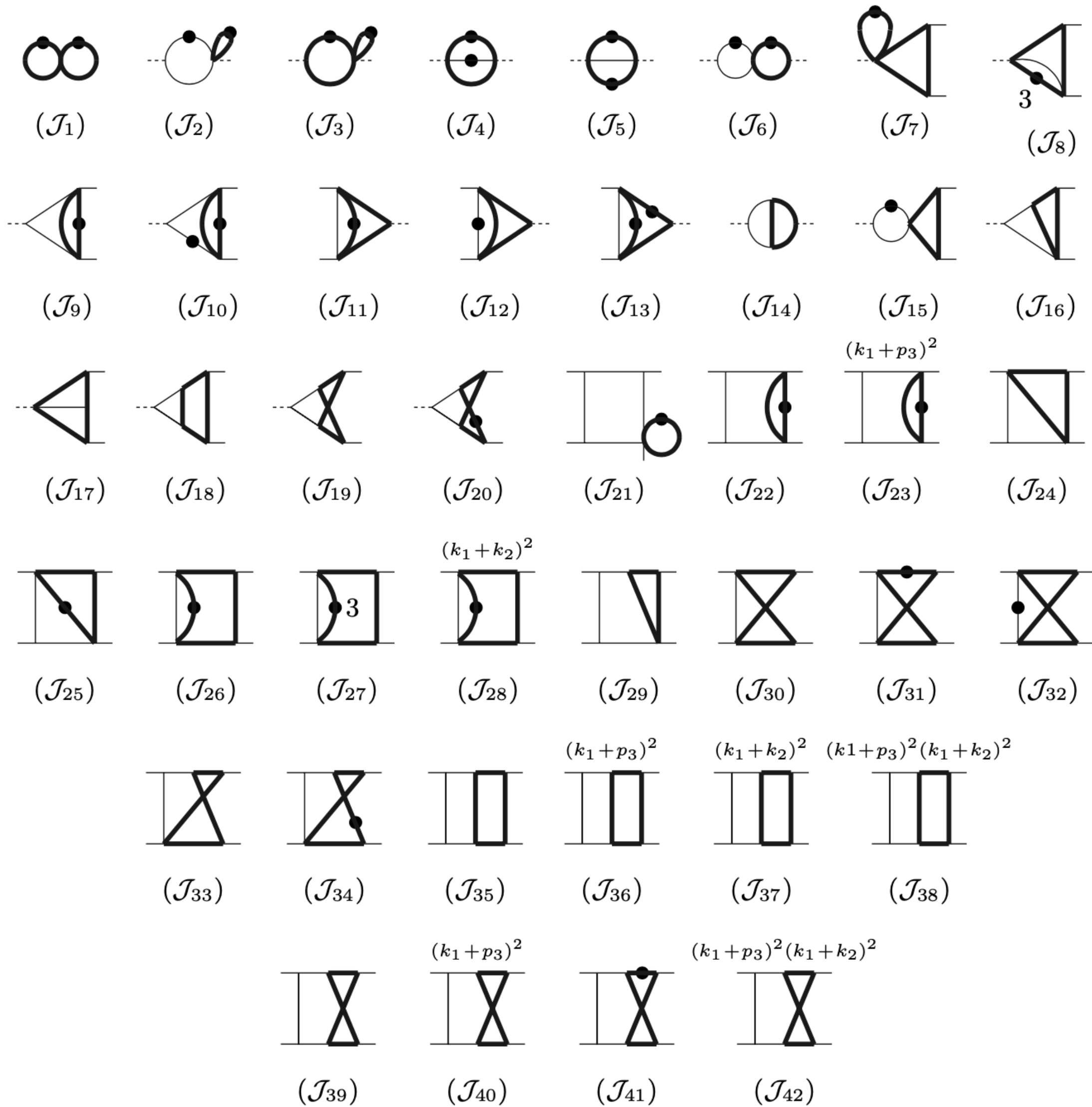
[A.Von Manteuffel,L.Tancredi]

NPL Master Integrals



Original MIs

Master Integrals



Now we have
42 MIs
for all the
process!

Evaluation of the MIs - PLA family:

The MIs are computed through the differential equations method:

$$df(\underline{x}, \epsilon) = \epsilon dA(\underline{x})f(\underline{x}, \epsilon)$$

Canonical logarithmic form [J.M.Henn]

with respect to the kinematic invariants: $\underline{x} = \{y, z\}$, $y = \frac{s}{m_t^2}$, $z = \frac{t}{m_t^2}$

Boundary Conditions: $\underline{x} = 0$

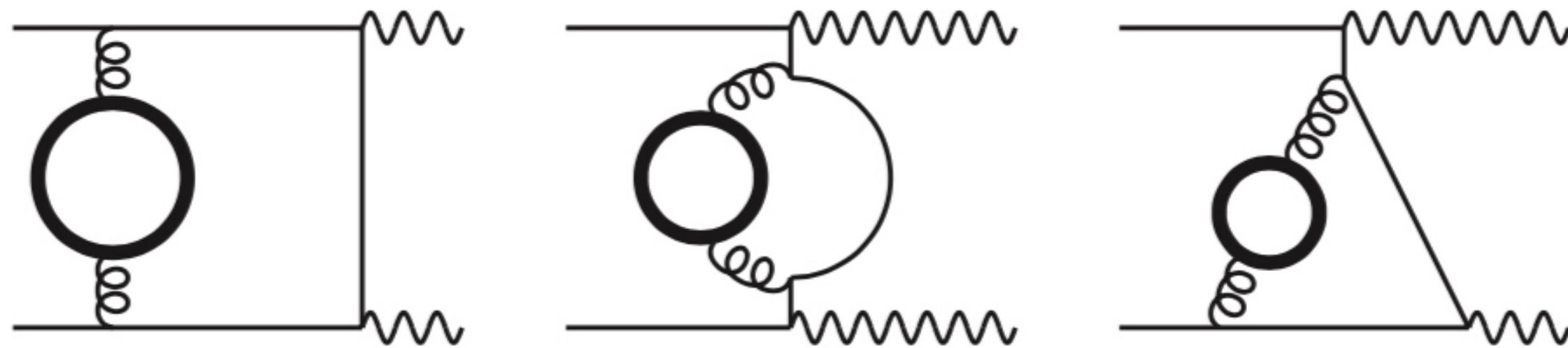
$$A(\underline{x}) = \sum c_i \log(w_i(\underline{x}))$$

$W_{PLA} = \{w_i(\underline{x})\} \rightarrow$ Set of 21 letters

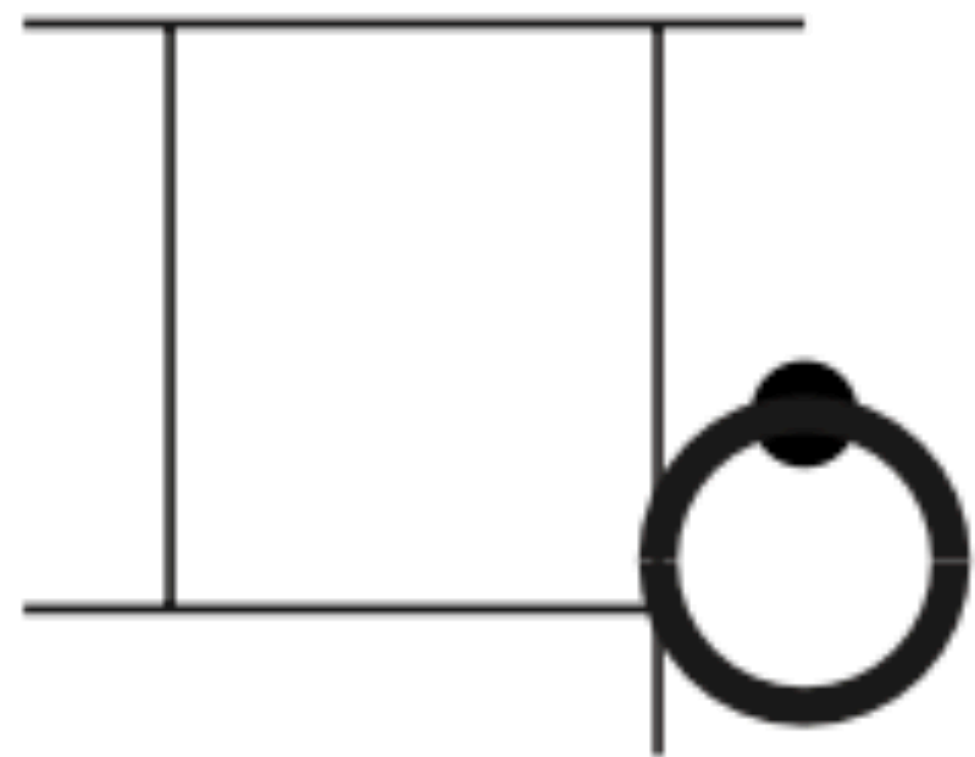
Cons of analytic evaluation :

- ❖ 5 non simultaneously linearizable square roots
- ❖ Non trivial solution!
- ❖ Big expressions!

Evaluation of the MIs - PLB family:



All the MIs coming from the PLB family, except J_{21} , are equal to one of the other two topologies PLA and NPL



Obtained by integrating analytically its differential equations

We don't need to set up a system of DEs for PLB

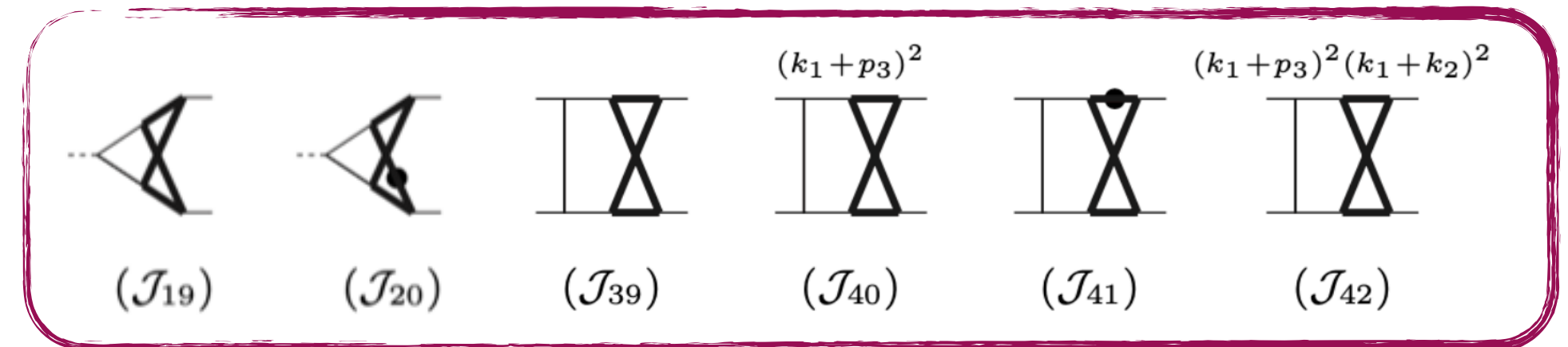
Evaluation of the MIs - NPL family

$$d\underline{g}(\underline{x}, \epsilon) = \epsilon dA(\underline{x})\underline{g}(\underline{x}, \epsilon) + d\tilde{A}(\underline{x}, \epsilon)\underline{g}(\underline{x}, \epsilon)$$

Two different subsets

Canonical
Logarithmic

ELLIPTIC
SECTORS

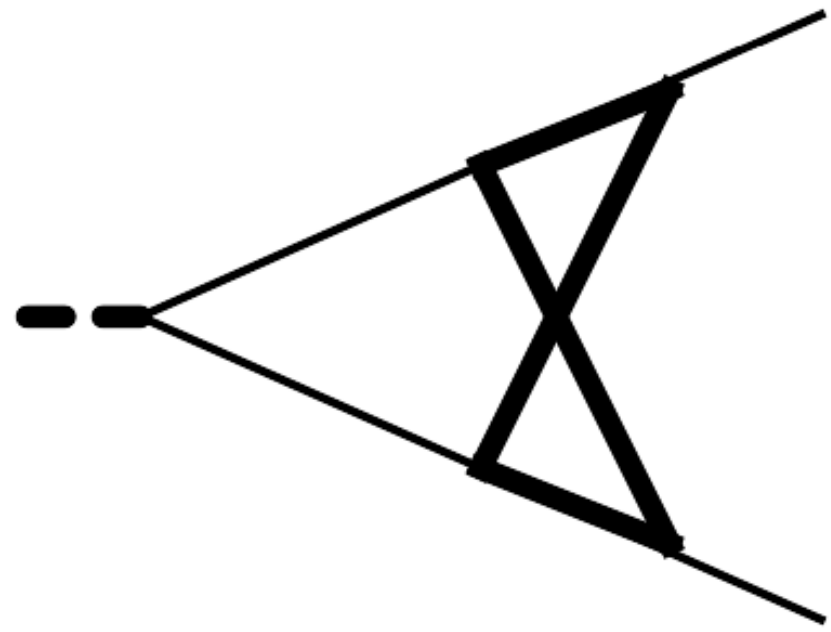


$W_{NPL} = \{w_i(\underline{x})\} \rightarrow$ Set of 30 letters

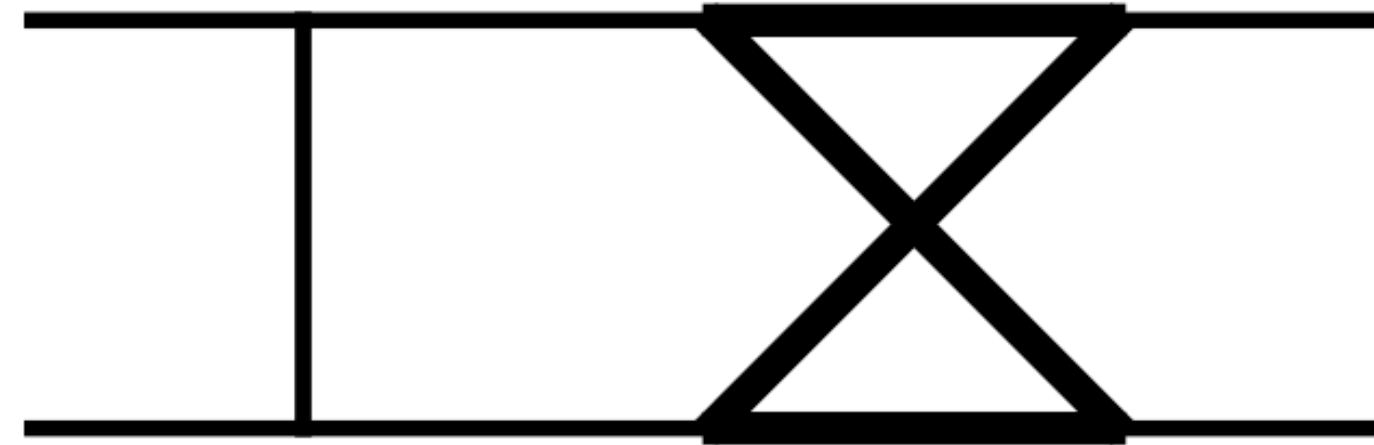
Cons of analytic
evaluation :

- ❖ Nine square roots in the alphabet
- ❖ Non trivial solution!
- ❖ Integrals involving eMPLs kernels

Elliptic sectors



Non-planar
triangle



Non-planar
double-box

The two sectors $(0,1,1,1,1,1,0,0)$ and $(1,1,1,1,1,1,1,0,0)$ for the NPL family are elliptic

$$MC(\text{triangle}) \propto \frac{1}{s} \int \frac{dz}{\sqrt{z(z+s)(z(s+z) - 4sm_t^2)}}$$

[J.Broedel,C.Duhr,F.Dulat,B.Penante,L.Tancredi]

Elliptic sectors

$$MC\left(\begin{array}{c} \text{---} \\ | \quad \times \\ \text{---} \end{array}\right) \propto \frac{1}{s} \int \frac{dz}{z \sqrt{(z+t)(s+t+z)((z+t)(s+t+z) - 4sm_t^2)}}$$

Möbius transformation: $z \rightarrow z - t$

The two elliptic curves are the same!

$$MC\left(\begin{array}{c} \text{---} \\ | \quad \times \\ \text{---} \end{array}\right) \propto \frac{1}{s} \int \frac{dz}{z \sqrt{z(z+s)(z(z+s) - 4sm_t^2)}}$$



[G.Fontana]

Generalised power series approach

$$\underline{f}(t, \epsilon) = \sum_{k=0}^{\infty} \epsilon^k \sum_{i=0}^{N-1} \rho_i(t) \underline{f}_i^{(k)}(t)$$

$$\rho_i(t) = \begin{cases} 1 & \text{if } t \in [t_i - r_i, t_i + r_i) \\ 0 & \text{if } t \notin [t_i - r_i, t_i + r_i) \end{cases}$$

$$\underline{f}_i^{(t)} = \sum_{l_1=0}^{\infty} \sum_{l_2=0}^{N_{i,k}} c_k^{(i,l_1,l_2)} (t - t_i)^{l_1/2} \log(t - t_i)^{l_2}$$

[R. N. Lee, A. V. Smirnov, V. A. Smirnov]

[M.K.Mandal,X.Zhao]

[F.Moriello]

Series solutions around DEs singular points

Pros:

- ❖ It doesn't depend on the function space, so it allows us to avoid elliptic integrals
- ❖ Values at arbitrary phase-space points
- ❖ Can be used to perform phenomenological studies

Numerical evaluation of the Master Integrals

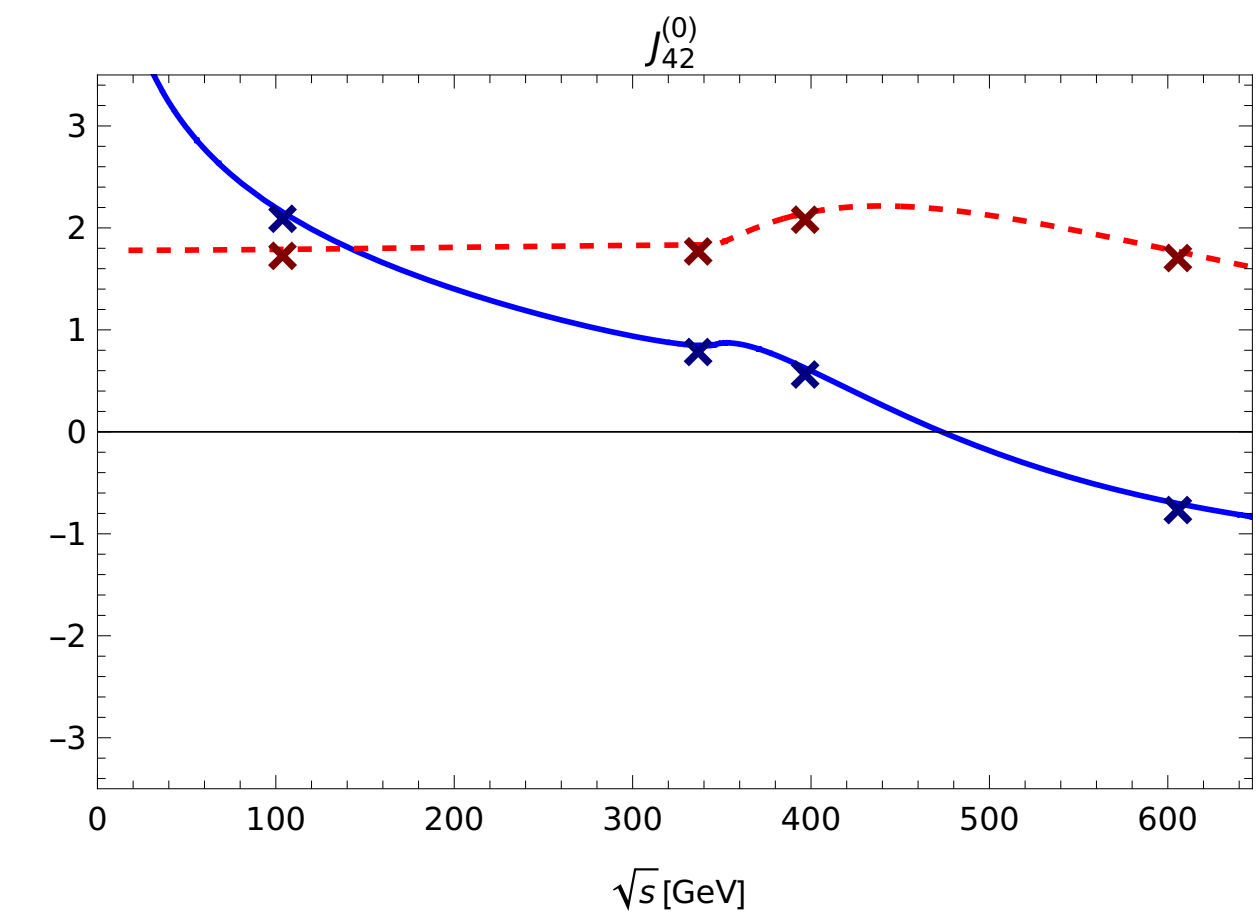
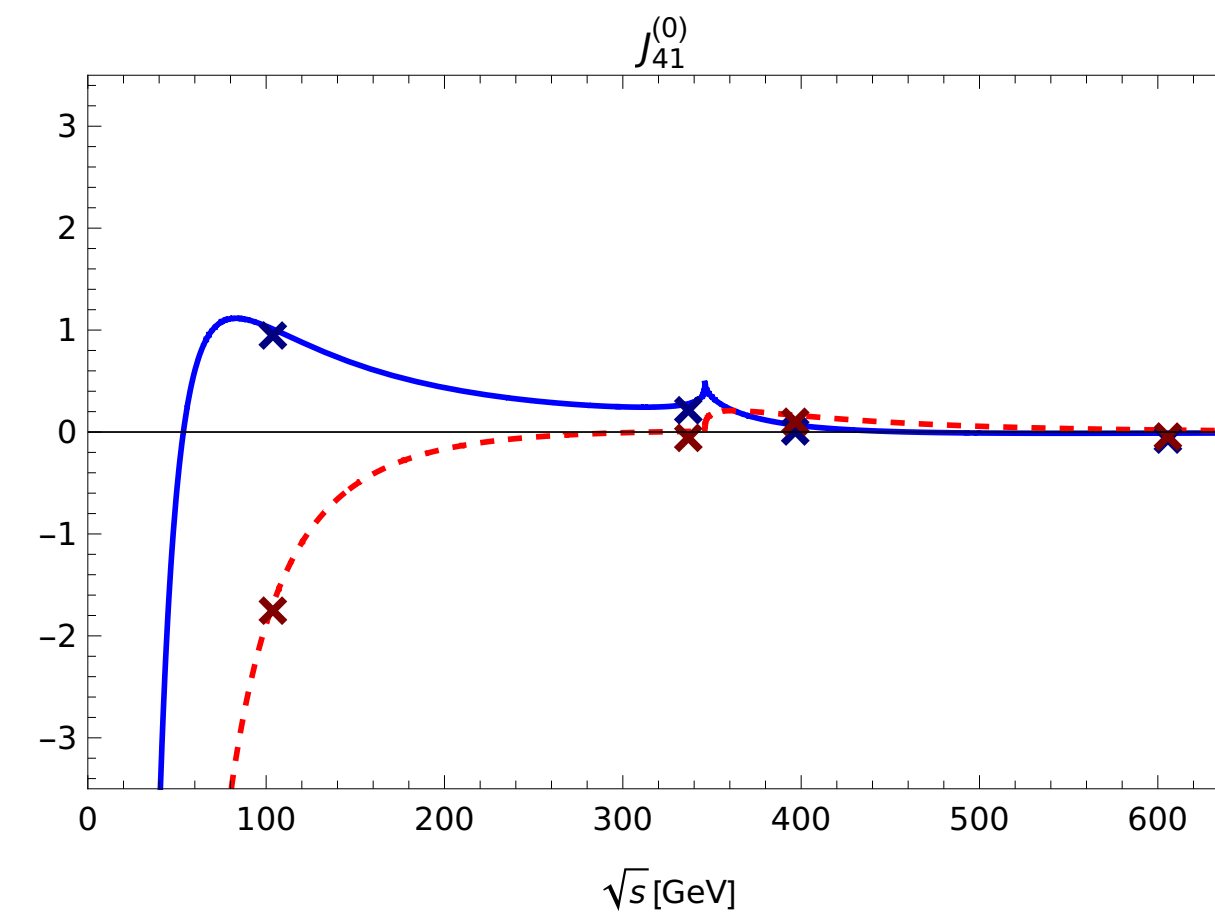
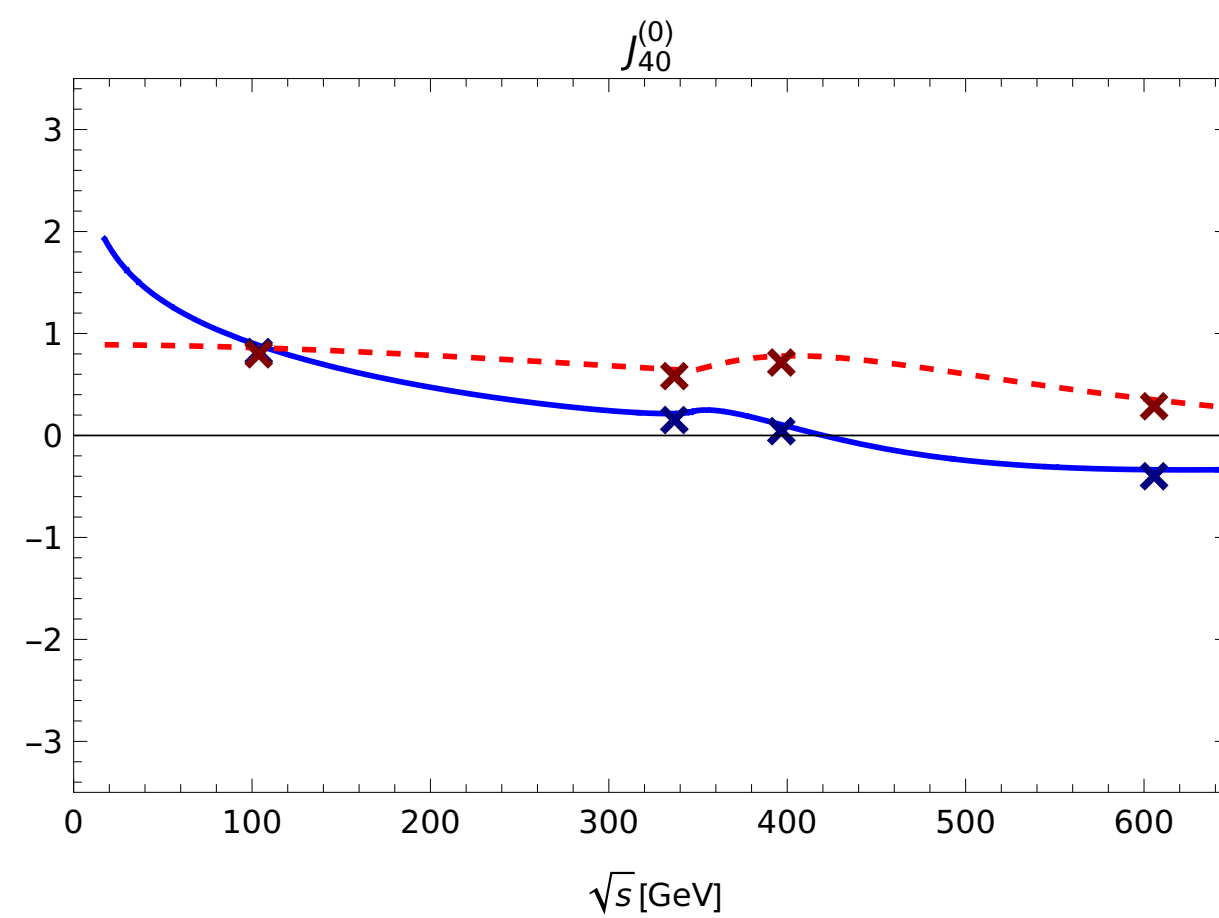
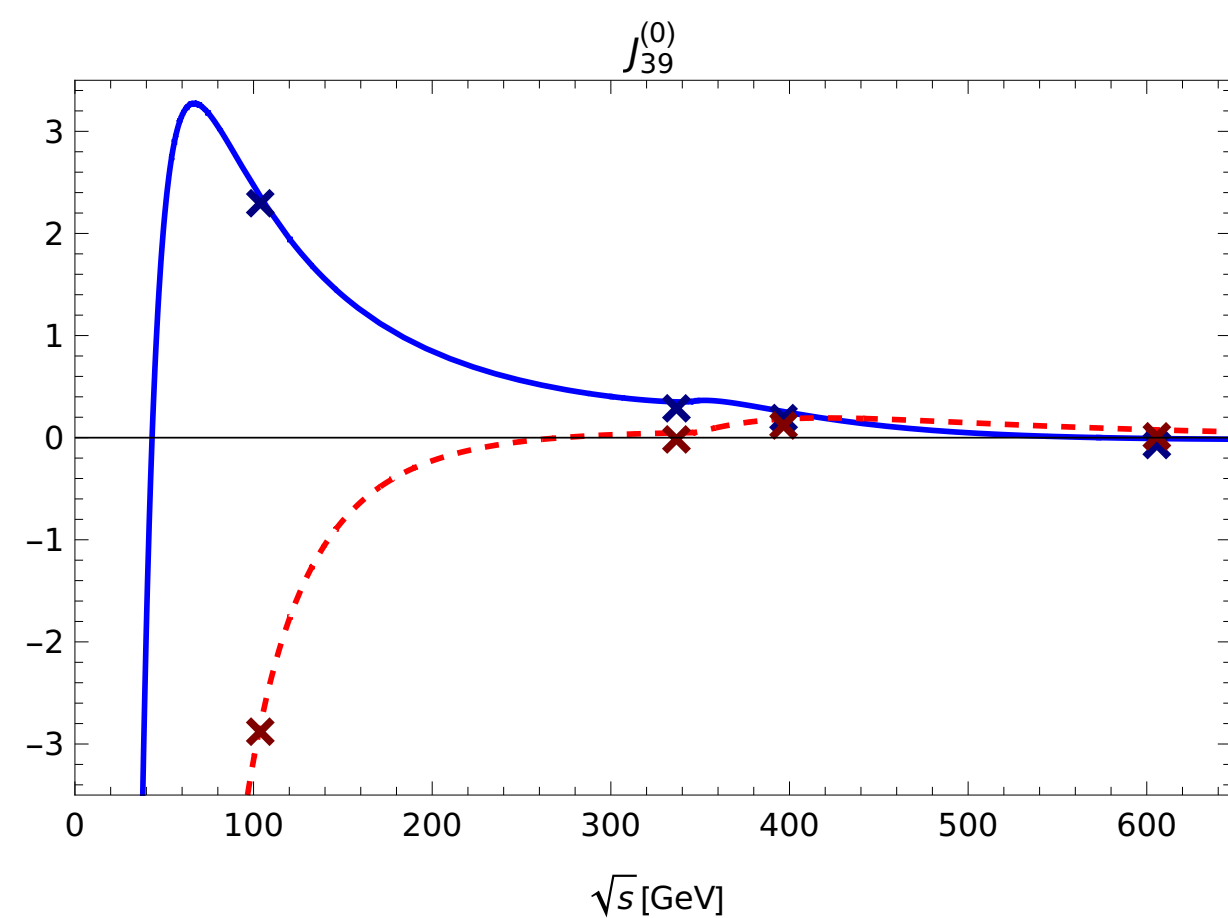
The numerical evaluation of the Master Integrals has been made with **DiffExp** [M.Hidding]

Several check for the numerical evaluation with **AMFLow** [X.Liu,Y.Ma]

Correspondance
up to 200 digits!

For the four elliptic boxes in NPL Topology, at a fixed angle:

----- Imaginary Part
————— Real Part



Numerical evaluation of the Hard Function

We prepared the numerical grid in the $2 \rightarrow 2$ physical phase-space region

$$s > 0, \quad t = -\frac{s}{2}(1 - \cos(\theta)), \quad -s < t < 0$$

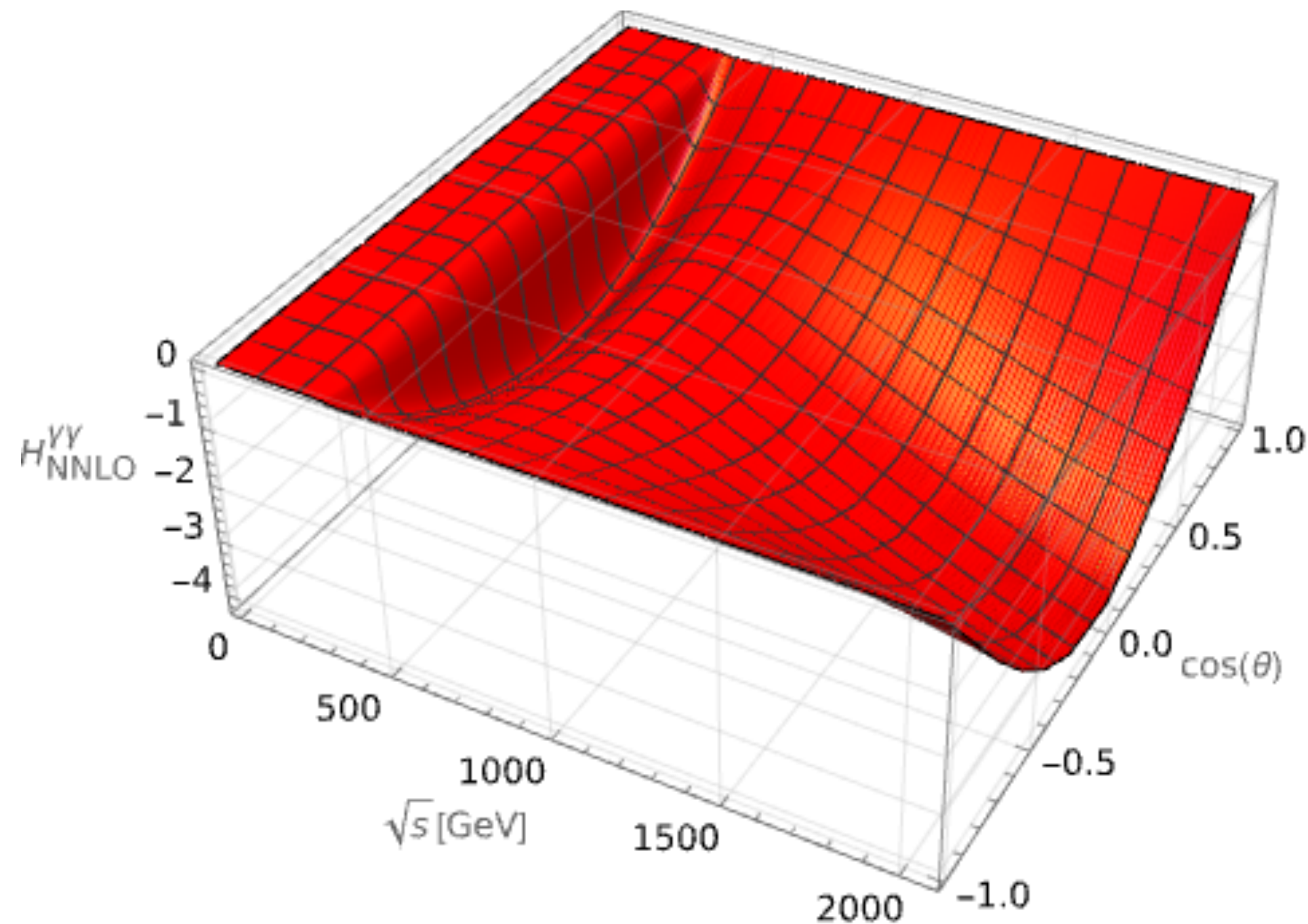
$$-0.99 < \cos(\theta) < +0.99$$

24 different values

$$8 \text{ GeV} < \sqrt{s} < 2.2 \text{ TeV}$$

573 different values

13752 points

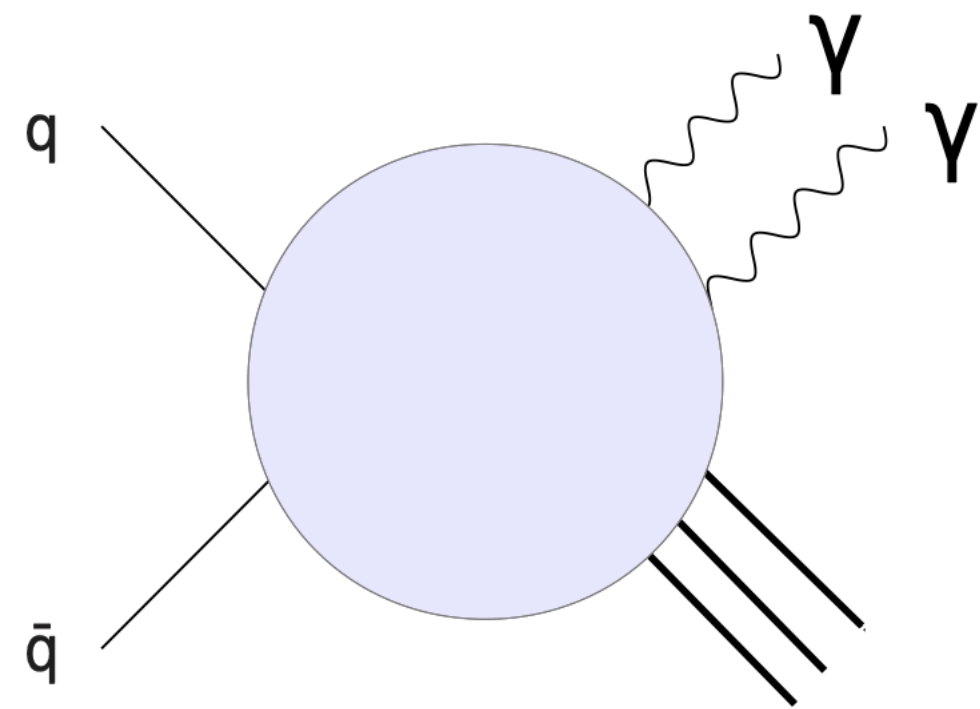


Planar Topology MIs
in $\mathcal{O}(2.5 \text{ h})$

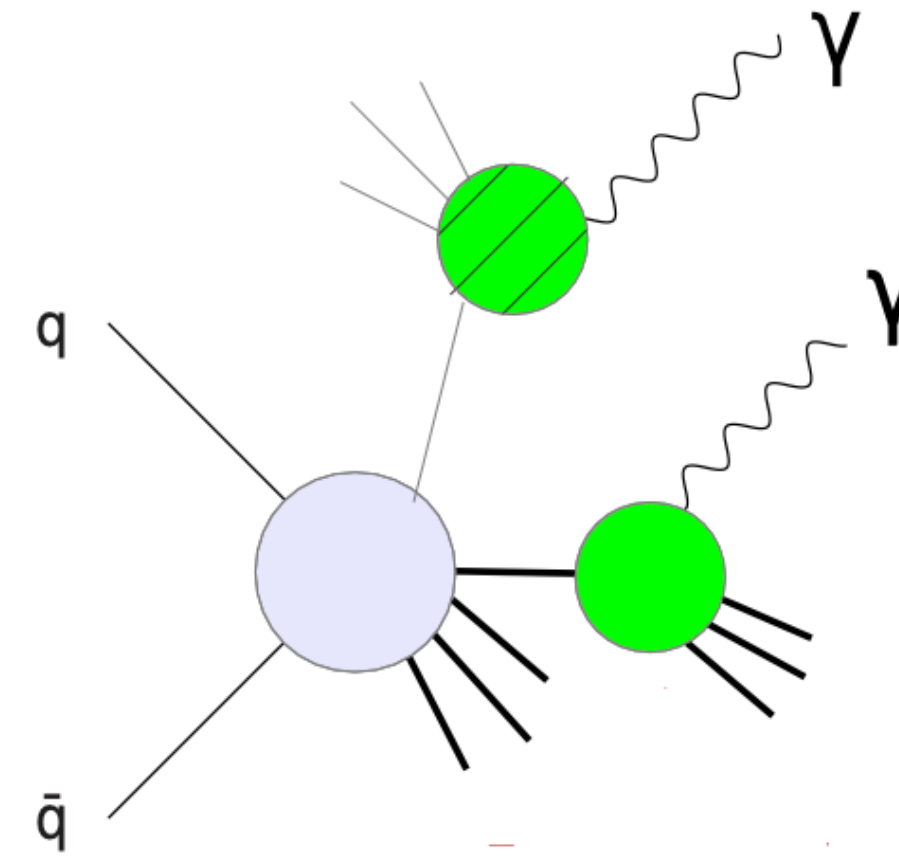
Non-planar Topology MIs
in $\mathcal{O}(10.5 \text{ h})$

On a single
core!

Photon production and isolation criteria



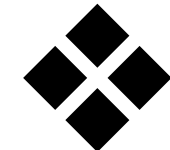
Direct component



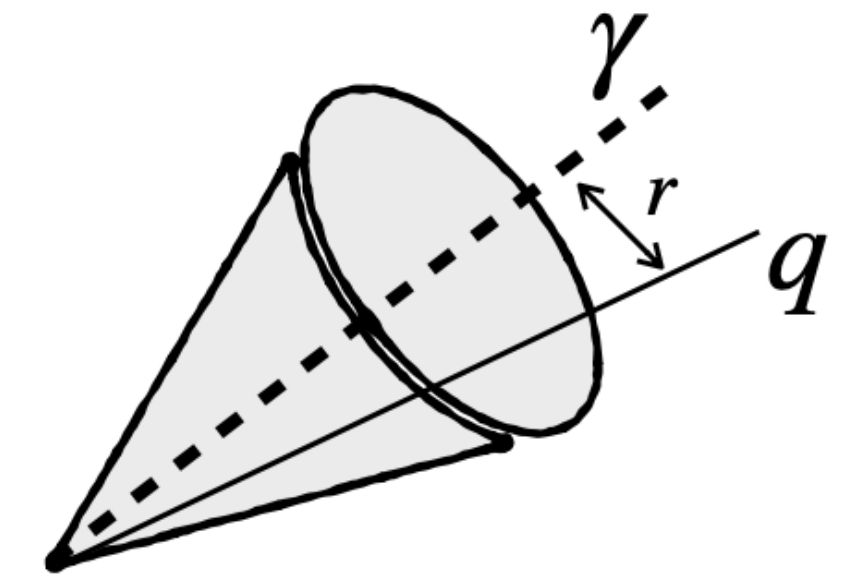
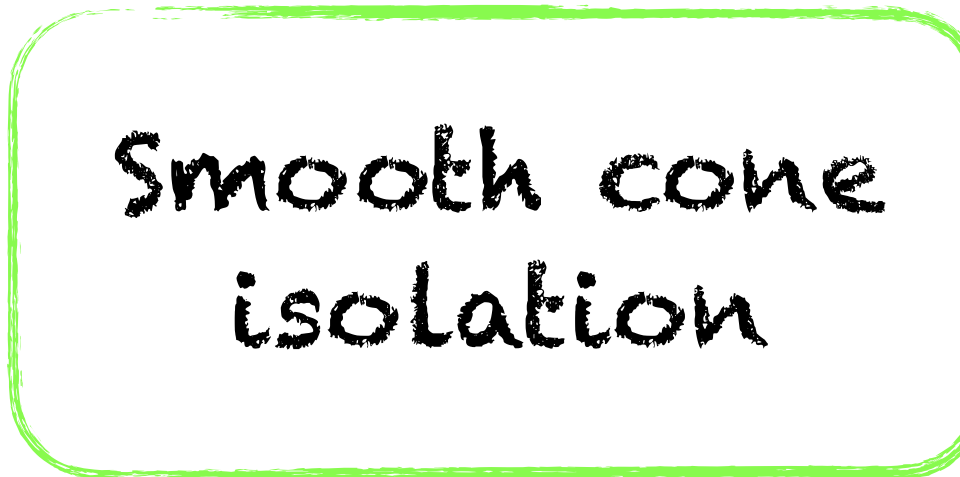
Fragmentation component



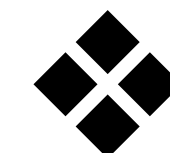
Experimentally photons must to be isolated



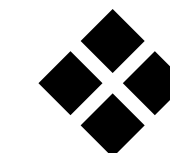
Isolation reduces fragmentation component



$$E_T^{had}(r) \leq \epsilon p_{T_\gamma} \chi(r; R)$$

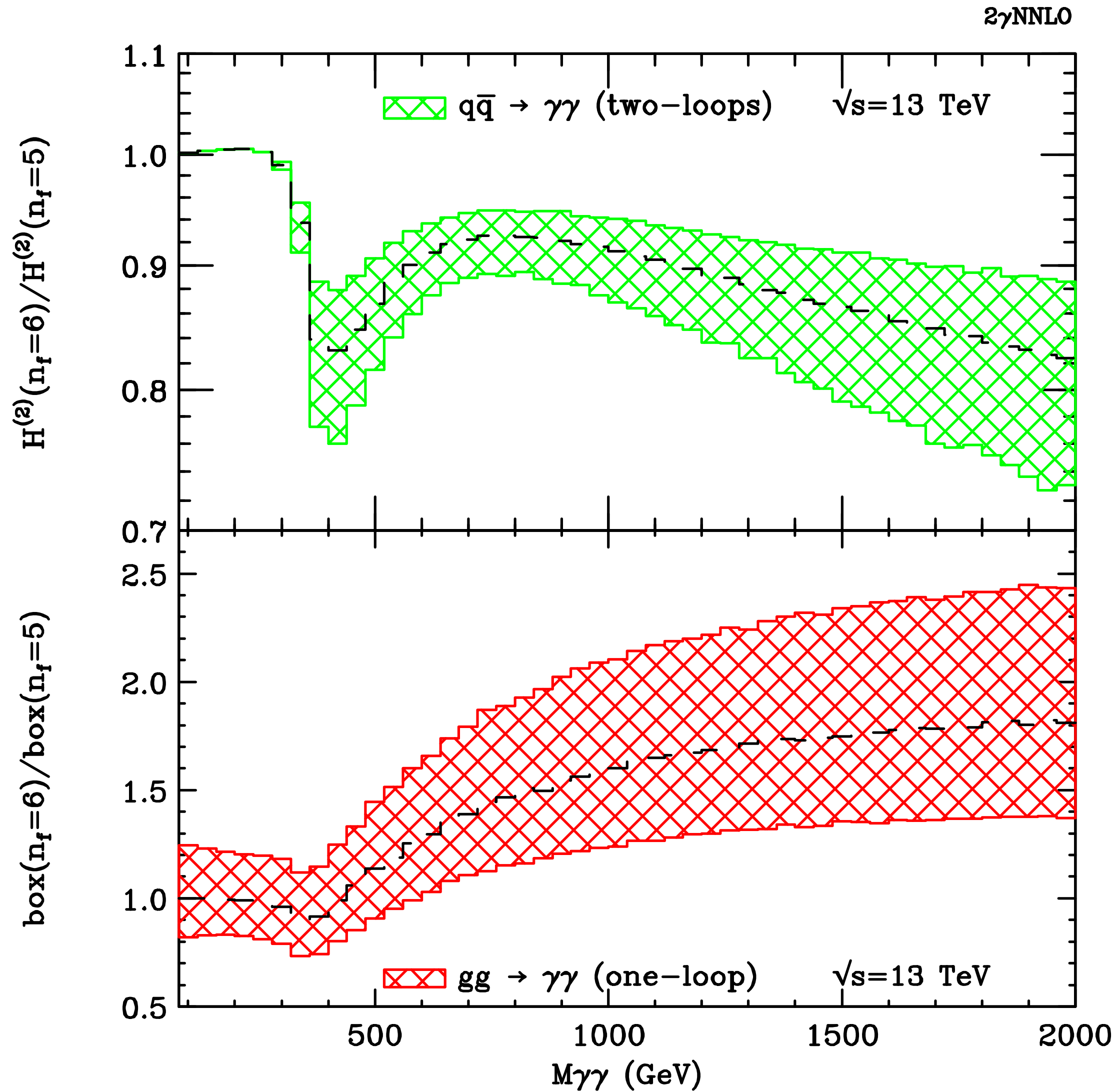


No quark-photon collinear divergences



No fragmentation component

Phenomenological Results



Upper panel: ratio between the fully massive and massless NNLO

$M_{\gamma\gamma} \sim 2m_t$ Negative peak \rightarrow Size of the ratio $\sim -15\%$

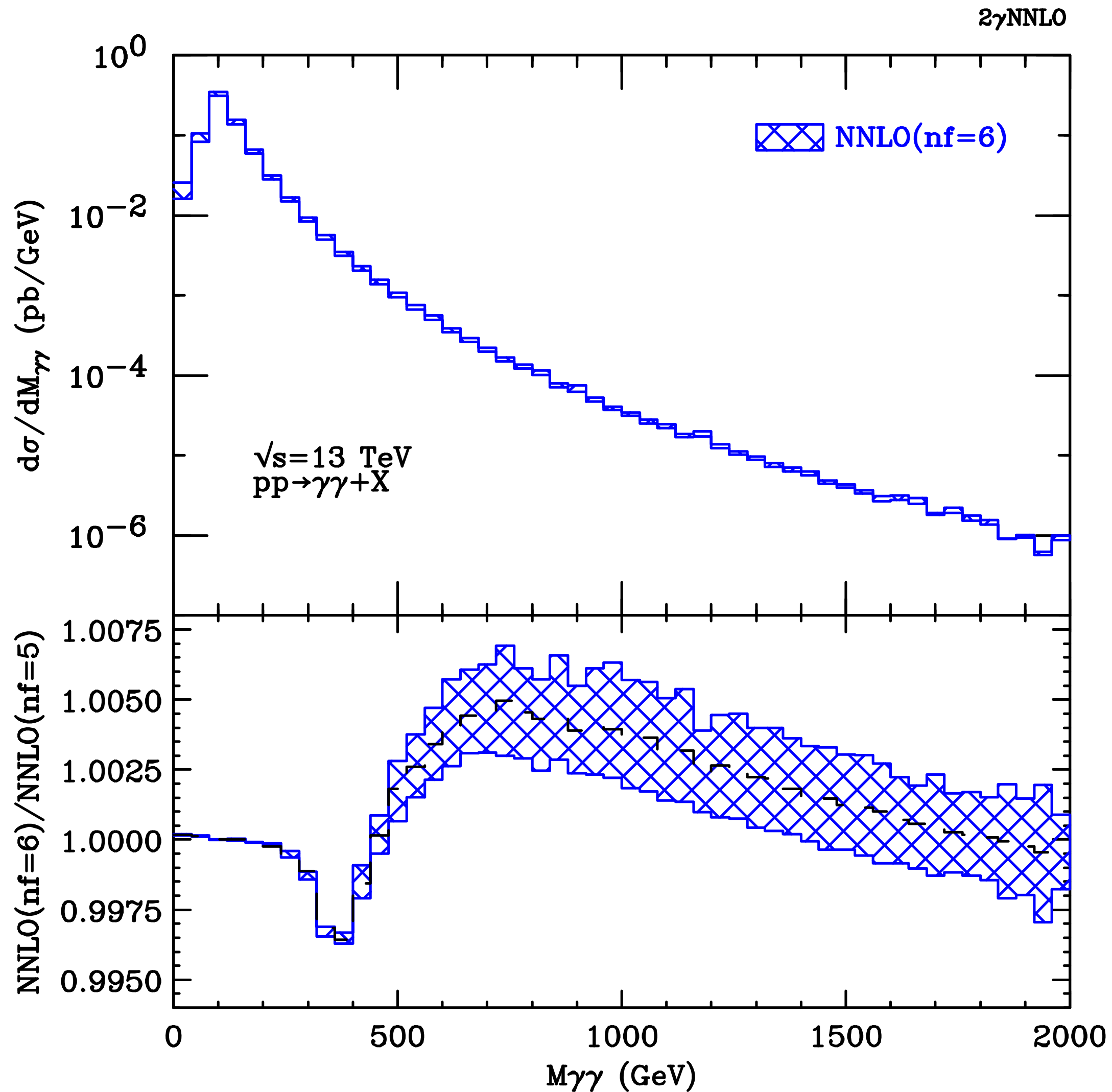
$M_{\gamma\gamma} > 2.3 \cdot 2m_t$ The tail decreases

Lower panel: ratio between the fully massive and massless one-loop box

For $M_{\gamma\gamma} \gg m_t$ the massive one-loop box contribution behaves as if it were composed by 6 light-quark flavors

$$\frac{(\sum_{n_f} e_q^2)^2}{(\sum_{n_f=5} e_q^2)^2} = \frac{225}{121}$$

Phenomenological Results



Ratio between the fully massive and massless invariant mass distribution at NNLO

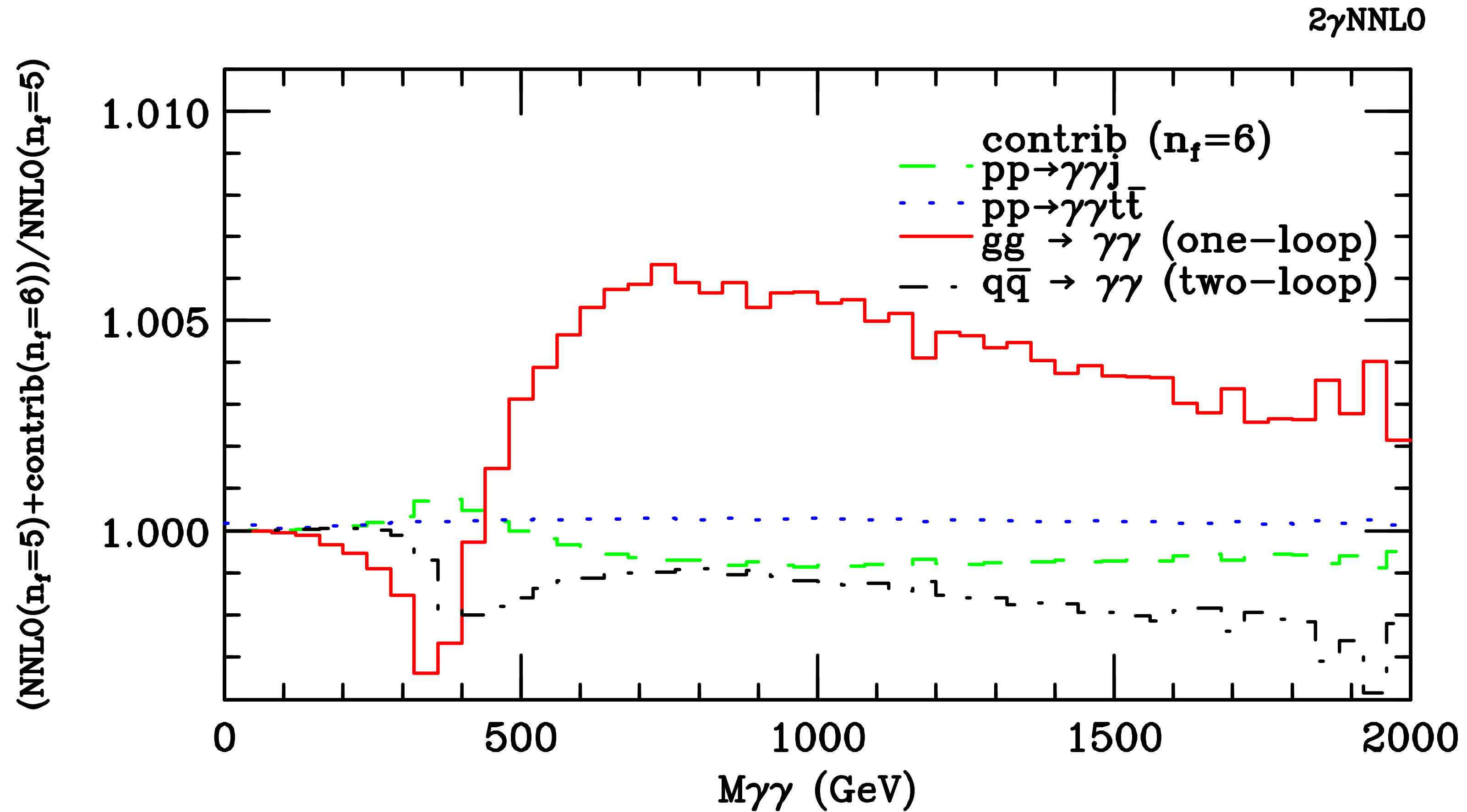
$M_{\gamma\gamma} \sim 2m_t$ Negative peak

$M_{\gamma\gamma} < 2m_t$ Massive corrections smaller than the massless one

$M_{\gamma\gamma} > 2m_t$ Massive corrections larger than the massless one

In the invariant mass region $1 \text{ GeV} < M_{\gamma\gamma} < 2 \text{ TeV}$
 deviation from the massless result in the range
 $[-0.4\%, 0.8\%]$

Phenomenological Results



Ratios of each one of the massive contributions with respect to the NNLO massless cross section as a functions of the invariant mass

Conclusions

- ❖ We computed the full NNLO QCD corrections to diphoton production
- ❖ The massive two loop $q\bar{q}$ -channel is one of the sizeable massive corrections
- ❖ Important corrections around the top quark threshold and along the distribution tail

Outlooks

- ❖ Analytic computation of the amplitude
- ❖ Inclusion of the partial massive N^3LO

Thank you for
your attention!