High Precision Results for Top-Quark Decay at N3LO in QCD

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Based on: [2309.01937 and work in progress]
In collaboration with Xiang Chen, Xin Guan and Yan-Qing Ma



Why Study the Top Quark

 Heaviest fundamental particle in SM $m_t = 172.69 \pm 0.30 \,\text{GeV}$

• Decay exclusively to b + W before

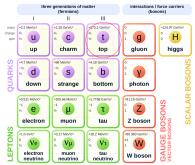
hadronization:

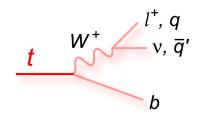
probs for possible BSM physics

$$\Gamma_t = 1.4 \, \text{GeV} \gg \Lambda_{\text{OCD}}$$

 Convergence of the perturbative QCD series (e.g. renormalon issue)

Standard Model of Elementary Particles





Why Study the Top Quark

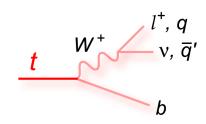
• Heaviest fundamental particle in SM

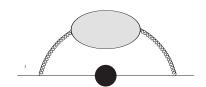
$$m_t = 172.69 \pm 0.30 \,\mathrm{GeV}$$

- Precision test of SM mechanism, and probs for possible BSM physics
- Decay exclusively to b + W before hadronization:

$$\Gamma_t = 1.4 \, \text{GeV} \gg \Lambda_{\text{QCD}}$$

• Convergence of the perturbative QCD series (e.g. renormalon issue)





Top Quark Mass m_t and **Decay Width** Γ_t

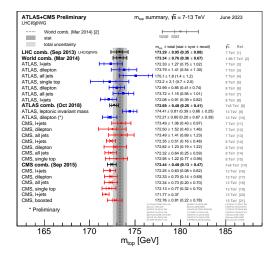
• PDG average for m_t :

$$172.69 \pm 0.30 \,\text{GeV}$$

• Current best measurement for Γ_t :

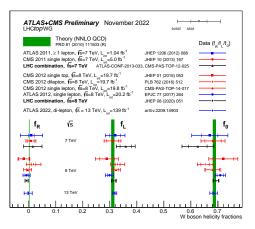
$$1.36 \pm 0.02 (\text{stat.})^{+0.14}_{-0.11} (\text{syst.}) \text{ GeV}.$$

 Experimental uncertainties anticipated at future colliders: 20 ~ 26 MeV



The W-helicity fractions in Top Decay

W from $t \rightarrow b + W^+ + X_{QCD}$ is polarized even if the *t*-quark is unpolarized



- The current best measurements: $f_0 = 0.684 \pm 0.005 \, (\mathrm{stat.}) \pm 0.014 \, (\mathrm{syst.})$, $f_\mathrm{L} = 0.318 \pm 0.003 \, (\mathrm{stat.}) \pm 0.008 \, (\mathrm{syst.})$ and $f_\mathrm{R} = -0.002 \pm 0.002 \, (\mathrm{stat.}) \pm 0.014 \, (\mathrm{syst.})$.
- Notoriously difficult to be predicted theoretically to high precision

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Much Theoretical Work Done So Far

Given the key role played by the top-quark both in SM precision test and searching for BSM, there have been vast amount of works done in literature regarding $t \to b + W^+ + X_{QCD}$.

• The inclusive Γ_t

▶ Up to NNLO in QCD: [Jezabek etc 88; Czarnecki etc 90; Li etc 90; Czarnecki etc 98; Chetyrkin etc 99; Fischer etc 01; Blokland etc 04'05;.....; Czarnecki etc 10; Meng etc 22; Chen etc 22]

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@NNNLO in QCD: [LC, Chen, Guan, Ma 23; Chen, Li, Li, Wang, Wang, Wu 23] [Datta, Rana, Ravindran, Sarkar 23 (only virtuals)]
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► NLO Electroweak: [Denner Sack 91; Eilam, Mendel Migneron Soni 91; Basso, Dittmaier, Huss, Oggero 15]

• W-helicities $f_{L,R,0}$

• @NNLO in QCD: [Czamecki, Korner, Piclum 10; Gao, Li, Zhu 12; Brucherseifer, Caola, Melnikov 13; Czarnecki, Groote Korner, Piclum 18]

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@NNNLO in QCD: [LC, Chen, Guan, Ma 23]
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► NLO Electroweak: [Do, Groote, Korner, Mauser 02]

Differential results

▶ QCD:

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    (a) NLO [Fischer, Groote, Korner, Mauser 01; Brandenburg, Si, Uwer 02; Bernreuther, Gonzalez, Mellei 14; Kniehl, Nejad 21]
    (a) NNLO [Gao, Li, Zhu 12; Brucherseifer, Caola, Melnikov 13; Campbell, Neumann, Sullivan 20]
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@NNNLO in QCD: [LC, Chen, Guan, Ma 23]

Cut Diagrams for Top Decay Width

 Γ_t in terms of the **semi-inclusive** $\mathcal{W}^{\mu
u}_{tb}$

$$\Gamma_t = \frac{1}{2 m_t} \int \frac{\mathrm{d}^{d-1} k}{(2\pi)^{d-1} 2E} \, \mathcal{W}_{tb}^{\mu\nu} \sum_{\lambda}^{L,R,0} \, \varepsilon_{\mu}^*(k,\lambda) \, \varepsilon_{\nu}(k,\lambda) \,,$$

$$\left|\mathcal{M}_{t\to b+W}\right|^2 \Rightarrow \frac{p}{\mu} \frac{k}{p_b} \frac{k}{\nu}$$

$$\mathcal{W}_{tb}^{\mu\nu}(p,k) = W_1 g^{\mu\nu} + W_2 p^{\mu} p^{\nu} + W_3 k^{\mu} k^{\nu} + W_4 (p^{\mu} k^{\nu} + k^{\mu} p^{\nu}) + W_5 i \epsilon^{\mu\nu\rho\sigma} p_{\rho} k_{\sigma},$$

Selection Criteria: the cut diagrams of *t*-quark self-energy function with exactly one (cut) *W* propagator interacting with the *external t*-quark plus (up to 3) QCD loops

Loop and Phase-space Integration

- Loop integrals are reduced using IBP [Chetyrkin 81] relations done with Blade [Guan, Liu, Ma 20], and the resulting masters are calculated using DE method [Kotikov 90; Remiddi 97] with AMFlow [Liu, Ma 22]
- The phase-space integrals, except for W-momentum *k*, are treated in the same manner as loop integrals by means of the reverse unitarity [Anastasiou, Melnikov 02]
- The phase-space integration of (divergent) $W_{tb}^{\mu\nu}$ over k are done "manually" using its power-log series representation (PSE) with ϵ assigned with non-zero numbers.
- ► Level of Complexity:

2988 master integrals, for which PSE about 200 orders in k_0 are derived with the above method. (c.f. only 185 MIs in the color-leading part)

Generalized Power-Log Series and the Integration Formula

According to Expansion-by-Region as well as Frobenius series solution for DE of dimensionally-regularized loop integrals:

$$f(\epsilon, x) = \sum_{a,b \in S} x^a \ln^b(x) T_{ab}(\epsilon, x) = \sum_{a,b \in S} x^a \ln^b(x) \left(\sum_{n=0}^{\infty} C_{abn}(\epsilon) x^n \right)$$

where $a = a_0 + a_1 \epsilon$ with rational a_0 , a_1 and non-negative integer b, belonging to a finite set S.

• Termwise integration formula we used:

$$\int_0^u x^a \ln^b(x) dx = \begin{cases} \frac{\ln^{1+b}(u)}{1+b} & \text{if } a = -1\\ u^{1+a} \frac{1}{(1+a)^{1+b}} \sum_{i=0}^b (-1)^{b-i} \frac{b!}{i!} (1+a)^i \ln^i(u) & \text{if } a \neq -1 \end{cases}$$

- The list of ϵ sampled: $10^{-3} + n \times 10^{-4}$ for $n = 0, 1, \dots, 15$.
- Fit in ε is done only at the very end for the final finite (physical) objects of interest. (e.g. we never reconstruct Laurent series results for individual masters)
- ► Consistency Check:

A perfect agreement in the result for the **inclusive** Γ_t with $\frac{\mathbf{d}^{d-1}k}{(2\pi)^{d-1}2E}$ calculated in this way and by directly applying the reverse unitarity [Anastasiou, Melnikov 02]

Results for the Inclusive Γ_t

The QCD effects on Γ_t in SM can be parameterized as

$$\Gamma_t = \Gamma_0 \left[\mathbf{c}_0 + \frac{\alpha_s}{\pi} \mathbf{c}_1 + \left(\frac{\alpha_s}{\pi} \right)^2 \mathbf{c}_2 + \left(\frac{\alpha_s}{\pi} \right)^3 \mathbf{c}_3 + \mathcal{O}(\alpha_s^4) \right],$$
 with $\Gamma_0 \equiv \frac{G_F \, m_W^2 \, m_t \, |V_{tb}|^2}{12 \sqrt{2}}$.

We choose $\mu = m_t/2$, motivated by the kinetic energy $m_t - m_W - m_b$ of the QCD radiations, at which our N3LO result reads: [LC, Chen, Guan, Ma 23]

$$\Gamma_t = 1.48642 - 0.140877 - 0.023306 - 0.007240 \text{ GeV}$$
= 1.31500 GeV

SM Inputs:

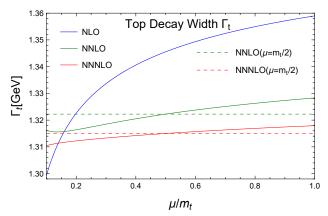
$$G_F = 1.166379 \times 10^{-5} \text{GeV}^{-2}$$
, $m_t = 172.69 \text{GeV}$, $m_W = 80.377 \text{GeV}$, $\alpha_s(m_t/2) \approx 0.1189$.

The leading-color part of Γ_t agrees with a parallel computation [Chen, Li, Li, Wang, Wang, Wu 23]

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The QCD Scale Uncertainty of Γ_t

The scale dependence of the fixed-order results for Γ_t in $\mu/m_t \in [0.1, 1]$



- NNLO scale-variation **never** cover the NNNLO result at any scales less than $\mu/m_t = 0.6$.
- Pure $\mathcal{O}(\alpha_s^3)$ correction decreases Γ_t by $\sim 0.8\%$ of the NNLO result at $\mu = m_t$ roughly 10 MeV(exceeding NNLO scale-hand)

The Offshell W and Finite m_b Effects

$$\bullet \ \ \text{With} \ \tfrac{1}{k^2 - m_W^2 + i\epsilon} \to \tfrac{1}{k^2 - m_W^2 + i m_W \Gamma_W},$$

$$\begin{split} \Gamma_t(m_W) \rightarrow \tilde{\Gamma}_t \; &= \; \int_0^{m_t^2} \frac{\mathrm{d}\,k^2}{2\pi} \frac{2m_W \Gamma_W}{(k^2 - m_W^2)^2 + (m_W \Gamma_W)^2} \, \Gamma_t(m_W^2 \rightarrow k^2) \\ &= \; \tilde{\Gamma}_0 \left[\tilde{\mathbf{c}}_0 + \frac{\alpha_s}{\pi} \, \mathbf{c}_1 + \left(\frac{\alpha_s}{\pi} \right)^2 \mathbf{c}_2 + \left(\frac{\alpha_s}{\pi} \right)^3 \mathbf{c}_3 + \mathcal{O}(\alpha_s^4) \right], \end{split}$$

we find: $\frac{\tilde{c}_i - c_i}{c_i}$ takes -1.54%, -1.53%, -1.39%, -1.23% for i = 0, 1, 2, 3.

- Similarly, keeping $m_b = 4.78 \, \text{GeV}$, we find: $\frac{c_1^{m_b} c_1}{c_1} \approx \frac{c_2^{m_b} c_2}{c_2} \approx -1.47\%$.
- The NLO electroweak K-factor is re-evaluated to be $K_{EW}^{NLO} = 1.0168$.

Taking these misc-effects into account, we finally obtain the **to-date most-precise high-precision theoretical prediction**:

$$\Gamma_t = 1.3148^{+0.003}_{-0.005} \times |V_{tb}|^2 + 0.027 (m_t - 172.69) \text{ GeV}$$

the error of which meets the request by future colliders.

Results for W-helicity Fractions

Top decay width with polarized *W*:

$$\Gamma_{\lambda} = \frac{1}{2 m_t} \int \frac{\mathrm{d}^{d-1} k}{(2\pi)^{d-1} 2E} \, \mathcal{W}_{tb}^{\mu\nu} \varepsilon_{\mu}^*(k,\lambda) \, \varepsilon_{\nu}(k,\lambda)$$

The *W*-helicity fractions $f_{\lambda}^{[n]} = \frac{\sum_{i=0}^{n} \Gamma_{\lambda}^{[n]}}{\sum_{i=0}^{n} \Gamma_{\lambda}^{[n]}}$ truncated to $\mathcal{O}(\alpha_s^3)$ in massless QCD:

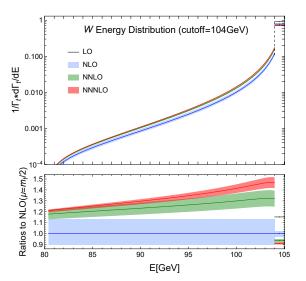
$$\begin{split} f_0^{[3]} &= 0.697706 - 0.008401 - 0.001954 - 0.000613, \\ &= 0.686737, \\ f_L^{[3]} &= 0.302294 + 0.007254 + 0.001799 + 0.000586, \\ &= 0.311933, \\ f_R^{[3]} &= 0. + 0.001147 + 0.000155 + 0.000027, \\ &= 0.001330. \end{split}$$

(The above results evaluated at $\mu=m_t$ agree with [Czarnecki, Korner, Piclum 10] up to NNLO)

With NLO EW-correction and m_b effects included, our **final results** read:

$$f_0^{[3]} = 0.686^{+0.002}_{-0.003}, \quad f_L^{[3]} = 0.312^{+0.001}_{-0.002}, \quad f_R^{[3]} = 0.00157^{+0.00002}_{-0.00002}.$$

Results for W-energy Distribution



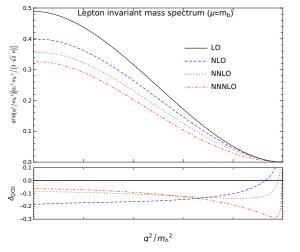
- ▶ In the bulk: QCD corrections are positive and quite sizable: pure $\mathcal{O}(\alpha_s^3)$ correction modifies the lowest order by $7 \sim 14\%$ for $E \in [94, 104]$ GeV.
- In the rightmost 1 GeV-bin: QCD corrections up to $\mathcal{O}(\alpha_s^3)$ decrease the Born-level result.

Bonus: The $b \rightarrow u e^+ \nu_e$ decay at higher orders in QCD

Every one more perturbative order higher in the $\overline{\rm MS}$ result, the term is reduced roughly by 1/2.

Bonus: lepton-pair invariant-mass spectrum in $b \rightarrow u e^+ \nu_e$

$$\Gamma(b \to u l \bar{\nu}_l) = \Gamma_0 \left(1 - 2.4131 \left(\frac{\alpha_s}{\pi} \right) - 21.27 \left(\frac{\alpha_s}{\pi} \right)^2 - 270.7 \left(\frac{\alpha_s}{\pi} \right)^3 \right).$$



- \triangleright $\mathcal{O}(\alpha_s^3)$ Leading-color inclusive part agrees with [Chen, Li, Li, Wang, Wang, Wu 23].
- $ightharpoonup {\cal O}(\alpha_s^3)$ corrections agrees with a recent approximation [Fael, Usovitsch 23] .

Summary and Outlook

We have provided the to-date most-precise theoretical prediction for top-quark decay width:

$$\Gamma_t = 1.3148^{+0.003}_{-0.005} \times |V_{tb}|^2 + 0.027 (m_t - 172.69) \text{ GeV}$$

the error of which meets the request by future colliders.

- $\ \ \, \square$ By a novel approach to complete phase-space integration over W momentum with IR-divergent integrand, we determined, in addition, W-helicity fractions, W-energy distribution at α_s^3 for the first time.
- $\ \ \, \square$ Furthermore, the lepton invariant-mass distribution in $b \to u l \bar{v}_l$ is derived up to α_s^3 , and an estimation of $\mathcal{O}(\alpha_s^4)$ correction for $\Gamma_{b \to u l \bar{v}_l}$ based on geometric series behavior is provided.
- \boxtimes The approach can be readily applied to the decay of polarized *t*-quarks at α_s^3 , as well as the mixed QCD-electroweak corrections.

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