Soft gluon resummation for associated top-quark pair production with a photon at the LHC

Michele Lupattelli



In collaboration with

Anna Kulesza, Roger Balsach

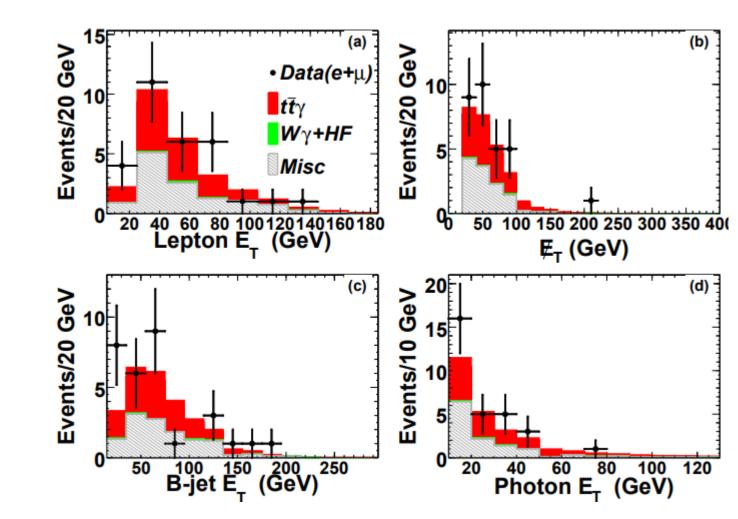
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The tty process

• First evidence at Tevatron in 2011

CDF collaboration, Phys. Rev. D 84, 031104 (2011)



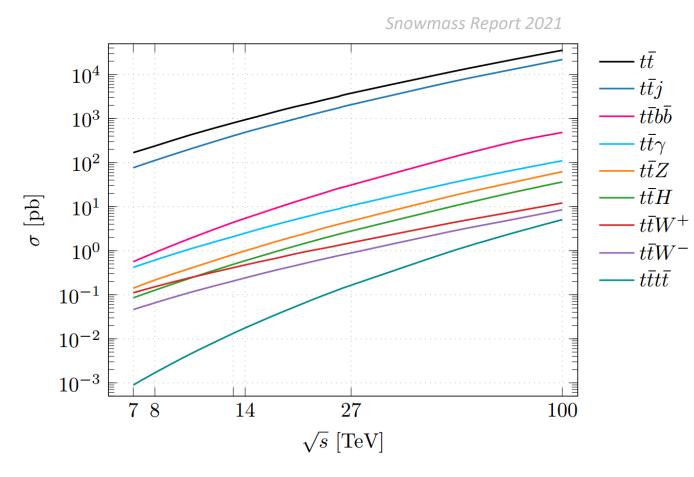
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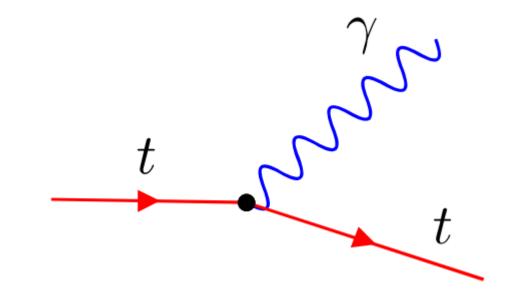


The tty process

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CDF collaboration, Phys. Rev. D 84, 031104 (2011)

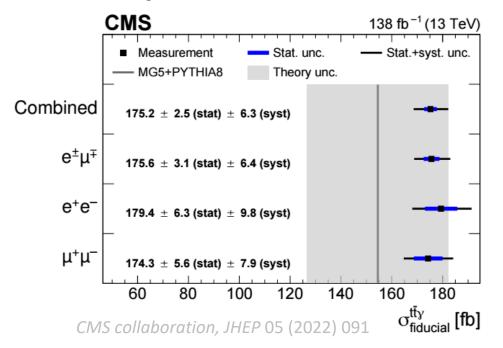
- First observation at the LHC in 2015 ATLAS collaboration, Phys. Rev. D 91, 072007 (2015)
- Probes the top-photon coupling
 - \rightarrow test of Standard Model
 - \rightarrow sensitive to new physics

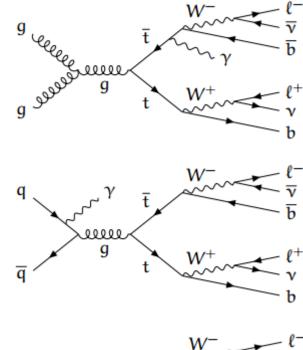


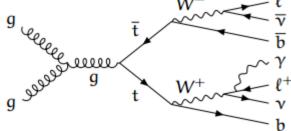
Feynman diagrams generated with FeynGame

[Harlander, Klein, Lipp, Comput. Phys. Commun. 256 (2020) 107465]

Latest experimental results

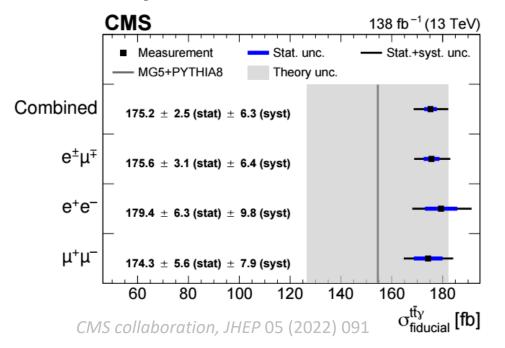


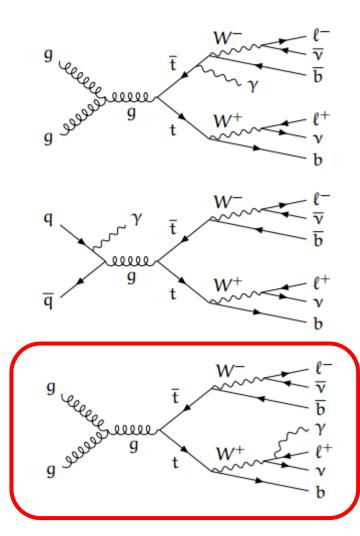




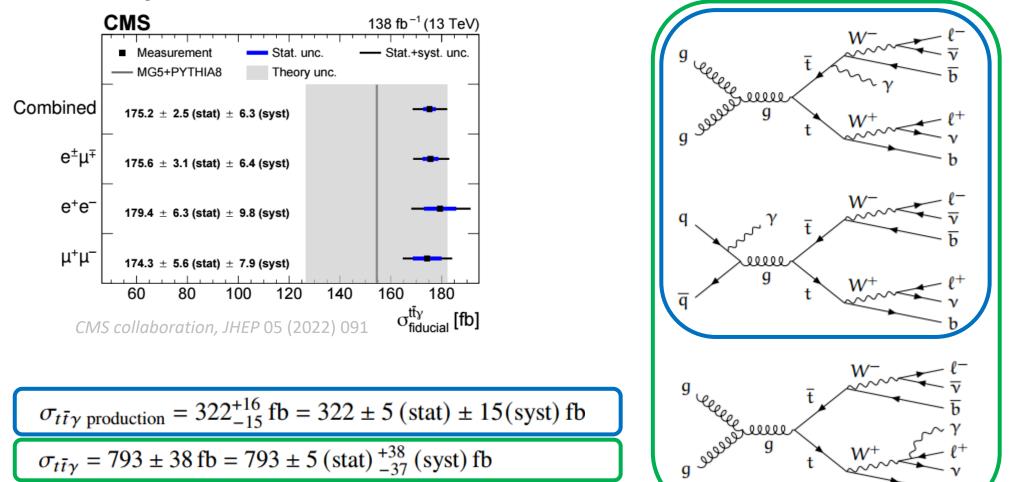


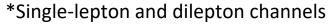
Latest experimental results





Latest experimental results





ATLAS collaboration, arXiv:2403.09452

Theoretical predictions

• NLO QCD

Duan, Ma, Zhang, Han, Guo, Wang '09, '11 | Maltoni, Pagani, Tsinikos '16

• NLO QCD + top-quark decays

Melnikov, Schulze, Scharf '11

• NLO QCD + parton shower

Kardos, Trocsanyi '15

NLO QCD + EW corrections

Duan, Zhang, Wang, Song, Li '17 | Pagani, Shao, Tsinikos, Zaro '21

• NLO QCD + off-shell effects

Bevilacqua, Hartanto, Kraus, Weber, Worek '18, '19, '20

Approximate NNLO

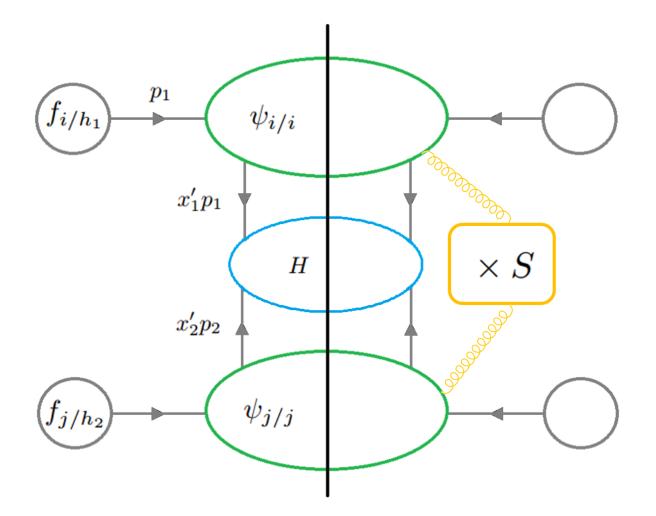
Kidonakis, Tonero '23

Complete NLO in top-quark pair production

Stremmer, Worek '24

Resummation

- Takes into account all orders in the soft-gluon emission limit
- Factorization of the amplitude close to kinematic threshold
- Factorization of the phase space in Mellin space



Invariant mass threshold resummation

$$\begin{split} \hat{\rho} &= \frac{Q^2}{s} \to 1 \qquad \qquad \alpha_s^n \left[\frac{\log^m (1-\hat{\rho})}{1-\hat{\rho}} \right]_+, \ m \le 2n-1 \\ & \text{Threshold variable} \\ \hline \frac{\mathrm{d}\sigma_{pp \to t\bar{t}\bar{B}}}{\mathrm{d}Q^2}(\rho) &= \sum_{i,j} \int_{c-\mathrm{i}\infty}^{c+\mathrm{i}\infty} \frac{\mathrm{d}N}{2\pi\mathrm{i}} \rho^{-N} \tilde{f}_{i/h_1}(N+1) \tilde{f}_{j/h_2}(N+1) \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}\bar{B}}}{\mathrm{d}Q^2}(N), \ \rho &= \frac{Q^2}{S} \\ & \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}\bar{B}}}{\mathrm{d}Q^2}(N) = \int \mathrm{d}\mathrm{PS}_3 \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}\bar{B}}(N) \bar{\mathbf{U}}_{ij \to t\bar{t}\bar{B}}(N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}\bar{B}} \left(\alpha_{\mathrm{S}} \left(\frac{Q^2}{(N+1)^2} \right) \right) \right. \\ & \left. \mathbf{U}_{ij \to t\bar{t}\bar{B}}(N+1) \right] \Delta_i(N+1) \Delta_j(N+1). \end{split}$$

$$\begin{split} & \text{Invariant mass threshold} \\ & \hat{\rho} = \frac{Q^2}{s} \to 1 \qquad \alpha_s^n \Big[\frac{\log^m(i)}{1} \\ & \text{Threshold variable} \\ \hline & \frac{\mathrm{d}\sigma_{pp \to t\bar{t}B}}{\mathrm{d}Q^2}(\rho) = \sum_{i,j} \int_{c-\mathrm{i}\infty}^{c+\mathrm{i}\infty} \frac{\mathrm{d}N}{2\pi\mathrm{i}} \rho^{-N} \tilde{f}_{i/h_1}(N + u) \\ & \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^2}(N) = \int \mathrm{d}\mathrm{PS}_3 \mathrm{Tr} \left[\mathbf{\tilde{H}}_{ij \to t\bar{t}B}(N) \mathbf{\tilde{U}}_{ij \to t\bar{t}B}(N + 1) \mathbf{\tilde{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^2}{(N+1)^2} \right) \right) \\ & \mathbf{U}_{ij \to t\bar{t}B}(N+1) \left[\Delta_i(N+1)\Delta_j(N+1) \right] \\ & \alpha_s^n \Big[\frac{\log^m(1-\hat{\rho})}{1-\hat{\rho}} \Big]_+ \qquad \longleftrightarrow \qquad \sum_k c_k \log^k N \end{split}$$

Ingredients for NNLL resummation

$$\left(\bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^2}{N^2} \right) \right) \approx \bar{\mathbf{S}}_{ij \to t\bar{t}B}^{(0)} + \frac{\alpha_{\mathrm{S}}(\mu_{\mathrm{R}}^2)}{\pi} \frac{\bar{\mathbf{S}}_{ij \to t\bar{t}B}^{(1)}}{1 - 2\lambda} \right)$$

$$\lambda = \alpha_{\rm S}(\mu_{\rm R}^2) b_0 \log(N)$$

$$\Delta_i = \exp\left(C_i(g^{(1)}\log(N) + g^{(2)} + \alpha_{\rm S}g^{(3)})\right)$$

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$$\tilde{\mathbf{H}}_{ij \to t\bar{t}B} = \tilde{\mathbf{H}}_{ij \to t\bar{t}B}^{(0)} + \frac{\alpha_{\mathrm{S}}(\mu_{\mathrm{R}}^{2})}{\pi} \tilde{\mathbf{H}}_{ij \to t\bar{t}B}^{(1)} + \dots \qquad \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^{2}} (N) = \int \mathrm{d}\mathrm{PS}_{3} \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}B} (N) \bar{\mathbf{U}}_{ij \to t\bar{t}B} (N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^{2}}{(N+1)^{2}} \right) \right) + \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^{2}} (N) = \int \mathrm{d}\mathrm{PS}_{3} \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}B} (N) \bar{\mathbf{U}}_{ij \to t\bar{t}B} (N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^{2}}{(N+1)^{2}} \right) \right) + \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^{2}} (N) = \int \mathrm{d}\mathrm{PS}_{3} \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}B} (N) \bar{\mathbf{U}}_{ij \to t\bar{t}B} (N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^{2}}{(N+1)^{2}} \right) \right) + \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^{2}} (N) = \int \mathrm{d}\mathrm{PS}_{3} \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}B} (N) \bar{\mathbf{U}}_{ij \to t\bar{t}B} (N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^{2}}{(N+1)^{2}} \right) \right) + \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^{2}} (N) = \int \mathrm{d}\mathrm{PS}_{3} \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}B} (N) \bar{\mathbf{U}}_{ij \to t\bar{t}B} (N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\alpha_{\mathrm{S}} \left(\frac{Q^{2}}{(N+1)^{2}} \right) \right) + \frac{\mathrm{d}\tilde{\sigma}_{ij \to t\bar{t}B}}{\mathrm{d}Q^{2}} (N) = \int \mathrm{d}\mathrm{PS}_{3} \mathrm{Tr} \left[\tilde{\mathbf{H}}_{ij \to t\bar{t}B} (N) \bar{\mathbf{U}}_{ij \to t\bar{t}B} (N+1) \bar{\mathbf{S}}_{ij \to t\bar{t}B} \left(\frac{Q^{2}}{(N+1)^{2}} \right) \right]$$

$$\bar{\mathbf{S}}_{ij \to t\bar{t}B}\left(\alpha_{\mathrm{S}}\left(\frac{Q^{2}}{N^{2}}\right)\right) \approx \bar{\mathbf{S}}_{ij \to t\bar{t}B}^{(0)} + \frac{\alpha_{\mathrm{S}}(\mu_{\mathrm{R}}^{2})}{\pi} \frac{\bar{\mathbf{S}}_{ij \to t\bar{t}B}^{(1)}}{1 - 2\lambda} \qquad \lambda = \alpha_{\mathrm{S}}(\mu_{\mathrm{R}}^{2})b_{0}\log(N)$$

$$\Delta_i = \exp\left(C_i(g^{(1)}\log(N) + g^{(2)} + \alpha_{\rm S}g^{(3)})\right)$$

Matching to NLO

$$\sigma_{pp \to Y}^{\text{fixed order} + \text{res}}(\rho) = \sigma_{pp \to Y}^{\text{fixed order}}(\rho) + \sum_{i,j} \int_{C_T} \frac{\mathrm{d}N}{2\pi \mathrm{i}} \rho^{-N} \tilde{f}_{i/h_1}(N+1) \tilde{f}_{j/h_1}(N+1) \left[\tilde{\sigma}_{ij \to Y}^{\text{res}}(N) - \tilde{\sigma}_{ij \to Y}^{\text{res}}(N) \right]_{\text{fixed order}}.$$

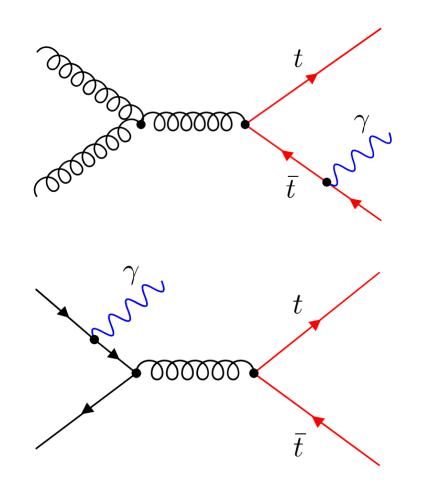
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$$\begin{split} \sigma_{pp \to Y}^{\text{fixed order} + \text{res}}(\rho) &= \sigma_{pp \to Y}^{\text{fixed order}}(\rho) \\ &+ \sum_{i,j} \int_{C_T} \frac{\mathrm{d}N}{2\pi \mathrm{i}} \rho^{-N} \tilde{f}_{i/h_1}(N+1) \tilde{f}_{j/h_1}(N+1) \\ &\left[\tilde{\sigma}_{ij \to Y}^{\text{res}}(N) - \tilde{\sigma}_{ij \to Y}^{\text{res}}(N) \right]_{\text{fixed order}} \end{split}$$

Approximate NNLO corrections (aNNLO)

- Expansion in α_s of the NNLL cross section in Mellin space
- Keep only relative order α_s^2 terms

Setup



We use the same setup as ATLAS analysis:

• Default scale choice:

$$\mu_R = \mu_F = \mu_0$$

$$\mu_0 = \frac{H_T}{2} = \frac{1}{2} \sum_i \sqrt{m_i^2 + p_{T,i}^2}$$

- Default PDF choice:
 NNPDF3.0
- Default cuts:

 $p_T^\gamma > 15~{\rm GeV}$

Results

		$\mu_0 = m_T/2$	1111 DF 5.0	
	σ [fb]	$+\delta\sigma$	$-\delta\sigma$	$\mathcal{K}^i_{t\bar{t}\gamma} = \sigma^i_{t\bar{t}\gamma}/\sigma^{\rm NLO}_{t\bar{t}\gamma}$
NLO	3.01(2)	+15.0%	-14.1%	_
NNLL	3.48(2)	+9.5%	-10.0%	1.16
aNNLO	3.35(2)	+10.8%	-11.0%	1.11

 $\mu_0 = H_T/2$ NNPDF3.0

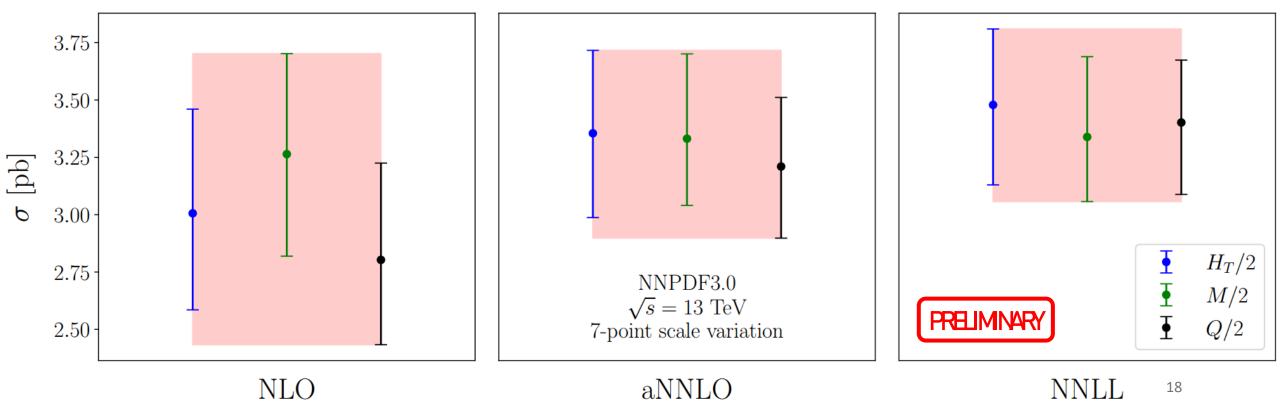
NLO predictions obtained with mg5_aMC@NLO

The higher-order soft-gluon corrections:

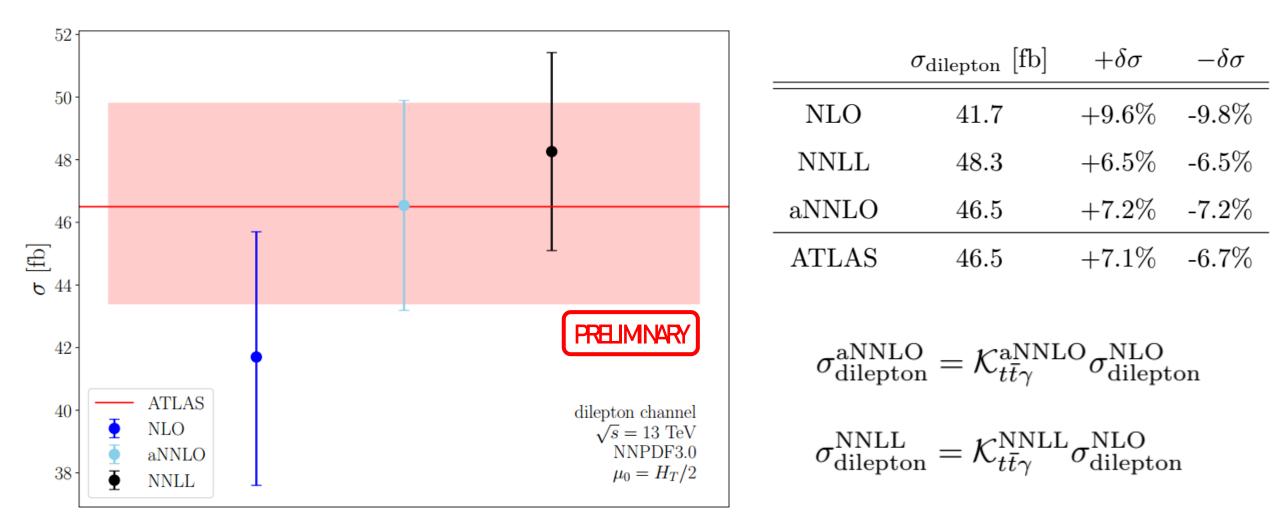
- Reduce the scale dependence by 5% (4% when keeping only NNLO)
- Their size is 16% of the NLO cross section (11% from NNLO)

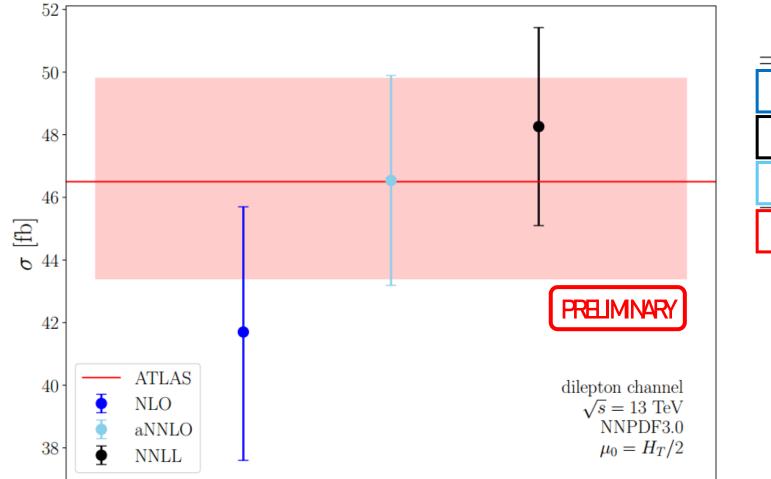
		$\mu_0 = H_T/2$	NNPDF3.0	
	σ [fb]	$+\delta\sigma$	$-\delta\sigma$	$\mathcal{K}^i_{t\bar{t}\gamma}=\sigma^i_{t\bar{t}\gamma}/\sigma^{\rm NLO}_{t\bar{t}\gamma}$
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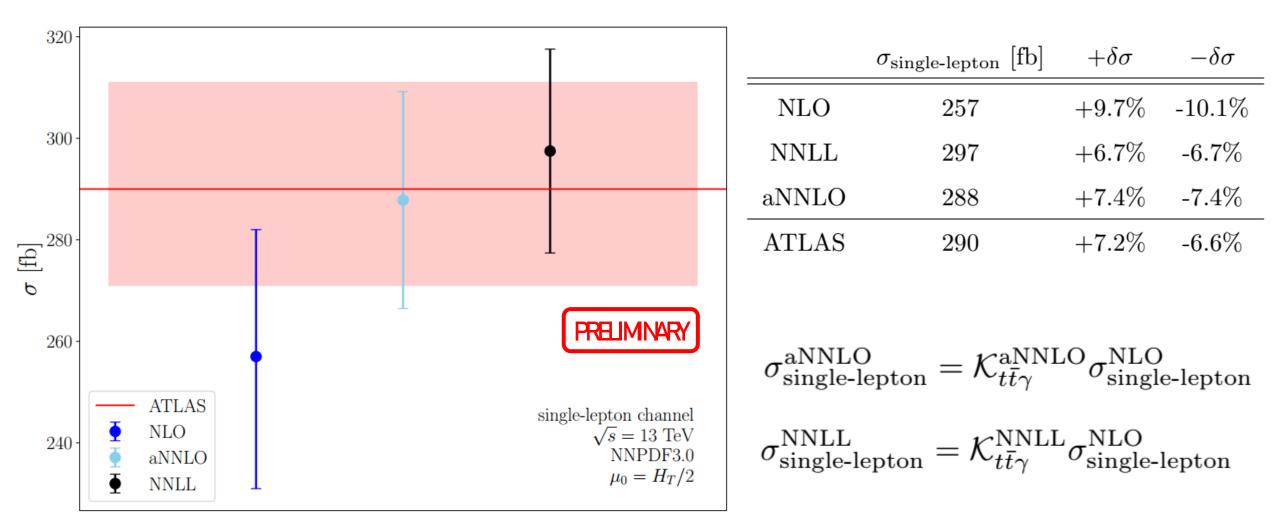
Results

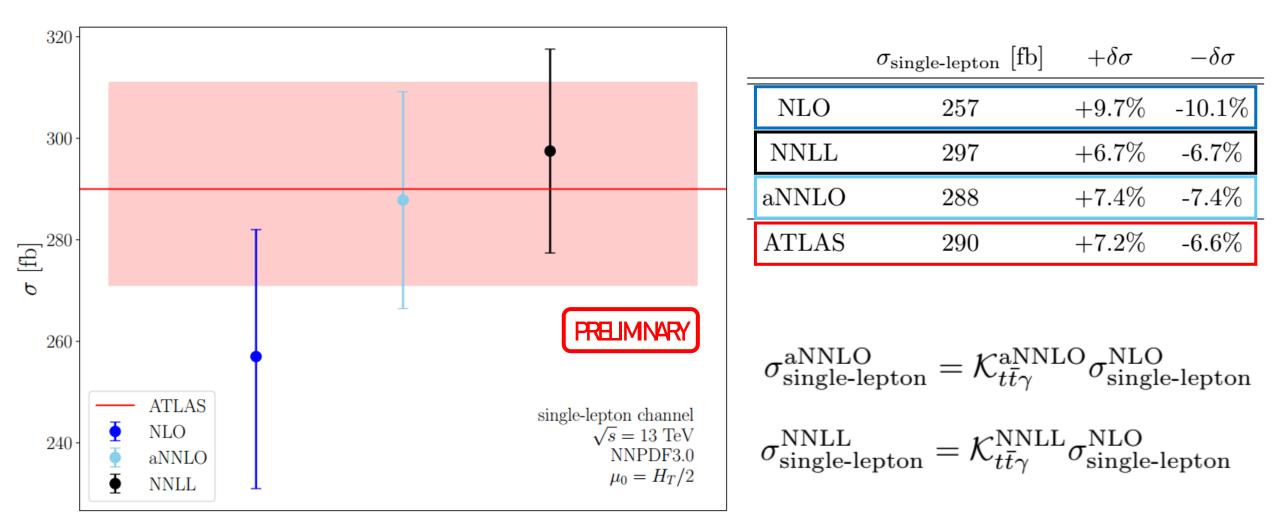




	$\sigma_{\rm dilepton}$ [fb]	$+\delta\sigma$	$-\delta\sigma$
NLO	41.7	+9.6%	-9.8%
NNLL	48.3	+6.5%	-6.5%
aNNLO	46.5	+7.2%	-7.2%
ATLAS	46.5	+7.1%	-6.7%

 $\sigma_{\rm dilepton}^{\rm aNNLO} = \mathcal{K}_{t\bar{t}\gamma}^{\rm aNNLO} \sigma_{\rm dilepton}^{\rm NLO}$ $\sigma_{\rm dilepton}^{\rm NNLL} = \mathcal{K}_{t\bar{t}\gamma}^{\rm NNLL} \sigma_{\rm dilepton}^{\rm NLO}$





Summary

- The inclusion of NNLL soft gluon corrections substantially reduces the dependence on the scales (the theoretical uncertainty from scale variation is estimated to ≈ 10%);
- Preliminary results that include NNLL soft gluon corrections bridge the gap between the theoretical predictions and the experimental results.

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