

Dependence of intrinsic kT on the collision center of mass energy using the Parton Branching Method in Drell-Yan production at NLO

42nd International Conference on High Energy Physics ICHEP 2024

Prague, Czech Republic

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Soft contributions and Sudakov form factor

- The transverse momentum dependent parton distribution functions (TMD PDFs) play an important role in the description of small transverse momentum physics as well as small x physics
- The parton branching (PB) method successful description up to the higher pT values
- The **PB method** is an **iterative procedure** using the concept of **resolvable and non-resolvable branching** and applying Sudakov form factors to describe the probability for non resolvable branchings from one evolution scale to another

$$\mathcal{A}_{a}(x,\mathbf{k},\mu^{2}) = \Delta_{a}(\mu^{2}) \mathcal{A}_{a}(x,\mathbf{k},\mu_{0}^{2}) + \sum_{b} \int \frac{d^{2}\mathbf{q}'}{\pi \mathbf{q}'^{2}} \frac{\Delta_{a}(\mu^{2})}{\Delta_{a}(\mathbf{q}'^{2})} \Theta(\mu^{2} - \mathbf{q}'^{2}) \Theta(\mathbf{q}'^{2} - \mu_{0}^{2})$$

$$\times \int_{x}^{z_{M}} \frac{dz}{z} P_{ab}^{(R)}(\alpha_{s},z) \mathcal{A}_{b}\left(\frac{x}{z},\mathbf{k}+(1-z)\mathbf{q}',\mathbf{q}'^{2}\right) ,$$
a, x
$$\mu$$
a, x
$$\mu$$

$$\mathbf{x}=\mathbf{x}_{0} \qquad \mathbf{\mu}_{0} \qquad \mathbf{b}, \mathbf{x}_{1}=\mathbf{x}_{0} \qquad \mu_{0}$$
b b, x = x_{0}
$$\mu_{0}$$
b b, x = x_{0}
$$\mu_{0}$$
b b, x = x_{0}
$$\mu_{0}$$

Intrinsic kT



$$k_T = k_{T,0} + \sum_c q_{T,c}$$

- intrinsic kT
- at the branching
- At the starting scale, parameter generated from a Gaussian as in the PB model:

$$A_{0,a}(x,{f k_T}^2,\mu_0^2)=f_{0,a}(x,\mu_0^2)(1/2\pi\sigma^2)e^{-(|k_T^2|/2\sigma^2)}$$

$$\sigma^2 = q_s^2/2$$



• At the initial state partons have not only longitudinal momentum, but also transverse momentum due to their internal (Fermi) motion -

• **Total transverse momentum** of the parton is that intrinsic transverse momentum + all the transverse momentum **qT** of the parton emitted

distribution with zero mean and a width σ expressed via parameter

- Scale dependence a much smaller sensitivity to the intrinsic-kT distribution at high Drell-Yan mass
- The intrinsic-kT interplays with the nonperturbative soft gluon contributions

Soft contributions

- z longitudinal momentum transferred at the branching, 0 < z < zM , $zM \rightarrow 1$
- **qT** the transverse momentum of the parton emitted at the branching

$$\alpha_s = \alpha_s \left(qT
ight) o q0$$

$$qT = (1-z)|q'|$$

 $z_{dyn}=1-q0/|q^{\prime}|$

- Angular ordering
 - \circ as is frozen
 - Two different regions:
 - perturbative region, with qT > q0
 - non-perturbative region of qT < q0</p>

=> qT - minimal parton transverse momentum emitted at a branching => zdyn - the dynamical resolution scale associated with the angular ordering

- Two regions of z:
 - \circ a perturbative region, with 0 < z < zdyn (qT > q0)
 - \circ a non-perturbative region with zdyn < z < zM (qT < q0)
 - Soft gluon resolution scale zM separates resolvable (z < zM) and non-</p> resolvable (z > zM) branchings

=> Define a perturbative (P) and non-perturbative (NP) ($zdyn < z < zM, zM \rightarrow 1$) Sudakov form factors

$$egin{aligned} &\Delta_a\left(\mu^2,{\mu_0}^2
ight) = exp\left(-\sum_b \int_{{\mu_0}^2}^{\mu^2} rac{d{q'}^2}{{q'}^2} \int_0^{z_{dyn}} z dz P_{ba}^{(R)}(lpha_s,z)
ight) exp\left(-\sum_b \int_{{\mu_0}^2}^{\mu^2} rac{d{q'}^2}{{q'}^2} \int_{z_{dyn}}^{z_{M2}} ec{dq'}^2 \left(\mu^2,{\mu_0}^2,{q_0}^2
ight) \cdot \Delta_a^{(NP)}\left(\mu^2,{\mu_0}^2,{q_0}^2
ight) \end{aligned}$$

 $\left(zdzP_{ba}^{(R)}(lpha_{s},z)
ight)
ight)$

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Drell-Yan pair production in hadron-hadron collisions



- - measured decay products
 - Three pT regions:
 - Non-perturbative region
 - Transition region
 - Perturbative region dominated by higher-order
 - contributions
- Low pt region significant for our analysis
 - intrinsic motion of partons
 - non-perturbative region
 - resummation of multiple soft gluon emissions
- DY production at NLO studied using the Parton Branching \bullet (PB) Method

• The production of Drell-Yan (DY) lepton pairs in hadron collisions - excellent process to study various QCD effects • Clean final state - no QCD final-state radiation, easily



Invariant mass dependance at $\sqrt{s} = 13 \text{ TeV}$



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• We used as baseline analysis the public CMS measurement Eur. Phys. J. C 83 (2023) 628

• **Detailed uncertainty breakdown**: complete treatement of experimental uncertainties + correlations between bins of the measurement

• The qs values obtained from each mass bin are compatible with each other

The most precise determination is obtained from the Z

• The sensitivity at high mass affected mainly from larger statistical uncertainties in the measurement

• The optimal qs obtained considering bins in all mass

$q_s = 1.04 \pm 0.08 \, GeV$

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DY production at different centre-of-mass energies

Analysis	√s	Collision types	Bins
<u>CMS (2022)</u> (shown before)	13 TeV	рр	25
<u>LHCb (2022)</u>	13 TeV	рр	5
<u>CMS (2021)</u>	8.1 TeV	pPb	5
<u>ATLAS (2015)</u>	8 TeV	рр	8
<u>CDF (2012)</u>	1.96 TeV	$p\overline{p}$	6
<u>CDF (2000)</u>	1.8 TeV	$p\overline{p}$	5
<u>D0 (2000)</u>	1.8 TeV	$p\overline{p}$	4
<u>PHENIX (2019)</u>	200 GeV	$p\overline{p}$	12
Total			81

• For the other measurements all uncertainties treated as being uncorrelated

Eur.Phys.J.C 84 (2024) 2, 154 Intrinsic kT-width dependence on \sqrt{s} and invariant mass

q0 = 0.01 GeV - minimal parton transverse momentum emitted at a branching



- Consistent values of qs for a large range of DY pair invariant masses
- Standard Monte Carlo event generators need a strongly increasing intrinsic-kT width with \sqrt{s}
- Strong center-of-mass energy dependence is not observed

Introducing energy dependence of the intrinsic-kT in PB

- Mimic parton-shower event generators by demanding a minimal parton transverse momentum
- q0 = 1 and 2 GeV
 - \circ qT > q0
- Non-perturbative part neglected





 Sensitivity of the DY cross section on the intrinsic-kT increases at small pair pT and with increasing of q0 value

arXiv:2404.04088

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Introducing energy dependence of the intrinsic-kT in PB



qs dependence on center-of-mass energy for the cases with q0 = 1 GeV and q0 = 2 GeV and q0 → 0 GeV
The uncertainty is estimated as a range of qs in which:

$$\chi^2\left(q_s
ight)-\chi^2_{min}<1$$

e performed a linear fit for the $q\left(q_s
ight)-log\left(\sqrt{s}
ight)$

The uncertainty bands around the fitted lines correspond to the 95% CL band

Introducing energy dependence of the intrinsic-kT in PB



- The slope increases as q0 increases
- Larger q0 means that more soft contributions are excluded
 - Larger intrinsic-kT needed to compensate missing contribution from soft gluons
 - Higher q0 values lead to an increased sensitivity to the intrinsic kTdistribution, resulting in smaller uncertainty bands

Summary

- DY production at NLO obtained with the MADGRAPH5_AMC@NLO event generator matched with the PB TMD distributions PBNLO-2018 Set2
- We study Fermi-motion of partons inside proton parameterized by a Gauss distribution of width $\sigma = q_s/\sqrt{2}$
- Proper treatment of the soft contributions in PB method leads to the intrinsic-kT width which does not depend on collision center-of-mass collision energy
 - The inclusion of soft gluons, in particularly the non-perturbative Sudakov, is crucial for providing \sqrt{s} -independent intrinsic-kT



Thank you for your attention!