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Dependence of intrinsic kT on the collision center of mass energy using the Parton Branching Method in Drell-Yan production at NLO

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Soft contributions and Sudakov form factor

- The transverse momentum dependent parton distribution functions (**TMD PDFs**) play an important role in the description of **small transverse momentum physics as well as small x physics**
- The **parton branching (PB) method** successful description up to the higher pT values
- The **PB method** is an **iterative procedure** using the concept of **resolvable and non-resolvable branching** and applying **Sudakov form factors** to describe the probability for non resolvable branchings from one evolution scale to another

$$
I_{a}(x, \mathbf{k}, \mu^{2}) = \frac{\Delta_{a}(\mu^{2}) \mathcal{A}_{a}(x, \mathbf{k}, \mu_{0}^{2})}{\Delta_{a}(\mathbf{q}^{2})} + \sum_{b} \frac{d^{2} \mathbf{q}^{2}}{\pi \mathbf{q}^{2}} \frac{\Delta_{a}(\mu^{2})}{\Delta_{a}(\mathbf{q}^{2})} \Theta(\mu^{2} - \mu^{2})
$$
\n
$$
\times \int_{x}^{z_{M}} \frac{dz}{z} P_{ab}^{(R)}(\alpha_{s}, z) \mathcal{A}_{b}(\frac{x}{z}, \mathbf{k} + (1 - z)\mathbf{q}^{2}, \mathbf{q}^{2}) ,
$$
\n
$$
I_{a} \times \sum_{\substack{\mathbf{k} \in \mathbb{Z}^{3} \\ \text{isomorphism } \mathbf{q} \text{ and } \mathbf{q} \text{ with } \
$$

1

 $- q'^2$) $\Theta(q'^2 - \mu_0^2)$

Intrinsic kT

$$
k_T = k_{T,0} + \sum_c q_{T,c}
$$

At the initial state partons have not only longitudinal momentum, but also transverse momentum due to their internal (Fermi) motion –

Total transverse momentum of the parton is that intrinsic transverse momentum + all the transverse momentum **qT** of the parton emitted

- intrinsic kT
- at the branching
- At the starting scale, parameter generated from **a Gaussian** qs in the PB model:

$$
A_{0,a}(x,\mathbf{kr}^2,\mu_0^2)=f_{0,a}(x,\mu_0^2)(1/2\pi\sigma^2)e^{-(|k_T^2|/2\sigma^2)}
$$

$$
\sigma^2 = q_s^2/2
$$

distribution with zero mean and a width expressed via parameter

- **Scale dependence** a much smaller sensitivity to the intrinsic-kT distribution at high Drell-Yan mass
- The intrinsic-kT interplays with the nonperturbative soft gluon contributions

Soft contributions

 \Rightarrow Define a perturbative (P) and non-perturbative (NP) (zdyn < z < zM, zM \rightarrow 1) Sudakov form factors

$$
\begin{aligned} &\Delta_a\left(\mu^2,{\mu_0}^2\right)=exp\left(-\sum_b\int_{\mu_0^2}^{\mu^2}\frac{dq'}{q'^2}\int_0^{z_{dyn}}zdzP_{ba}^{(R)}(\alpha_s,z)\right)exp\left(-\sum_b\int_{\mu_0^2}^{\mu^2}\frac{dq'}{q'^2}\int_{z_{dyn}}^{z_{M4}}\\ &=\left(\Delta_a{}^{(P)}\right)\left(\mu^2,\mu_0{}^2,q_0{}^2\right)\cdot\Delta_a{}^{(NP)}\left(\mu^2,\mu_0{}^2,q_0{}^2\right) \end{aligned}
$$

 $\left(\begin{matrix} \mathbb{R}^d \ 1 \end{matrix} z dz P^{(R)}_{ba}(\alpha_s, z) \right),$

- z longitudinal momentum transferred at the branching, $0 < z < zM$, $zM \rightarrow 1$
- **qT** the transverse momentum of the parton emitted at the branching

$$
\alpha_s=\alpha_s\left(qT\right)\to q0
$$

$$
qT=(1-z)|q^{\prime}|
$$

 $z_{dyn}=1-q0/|q'|$

- Angular ordering
	- αs is frozen
	- Two different regions:
		- perturbative region, with $qT > q0$
		- non-perturbative region of $qT < q0$

=> zdyn - the dynamical resolution scale associated with the angular ordering => qT - minimal parton transverse momentum emitted at a branching

- Two regions of z:
	- \circ a perturbative region, with $0 < z <$ zdyn (qT > q0)
	- \circ a non-perturbative region with zdyn $<$ z $<$ zM (qT $<$ q0)
		- Soft gluon resolution scale zM separates resolvable $(z < zM)$ and nonresolvable $(z > zM)$ branchings

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• The production of Drell-Yan (DY) lepton pairs in hadron collisions - excellent process to study various QCD effects **Clean final state** - no QCD final-state radiation, easily

- - measured decay products
	- Three pT regions:
		- Non-perturbative region
		- **Transition region**
		- **Perturbative region dominated by higher-order**
		- contributions
- **Low pT** region significant for our analysis \bullet
	- intrinsic motion of partons
	- o non-perturbative region
	- o resummation of multiple soft gluon emissions
- DY production at NLO studied using the Parton Branching \bullet (PB) Method

Drell-Yan pair production in hadron-hadron collisions

Invariant mass dependance at √s = 13 TeV

The qs values obtained from **each mass bin are compatible with each other**

The most precise determination is obtained from the **Z**

The sensitivity at high mass affected mainly **from larger statistical uncertainties** in the measurement

• The optimal qs obtained considering bins in all mass

$q_s = 1.04 \pm 0.08 \text{ GeV}$

We used as baseline analysis the public CMS measurement **Eur. Phys. J. C 83 (2023) 628**

Detailed uncertainty breakdown: complete treatement of experimental uncertainties + correlations between bins of the measurement

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DY production at different centre-of-mass energies

For the other measurements all uncertainties treated as being uncorrelated

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****q0 = 0.01 GeV - minimal parton transverse momentum emitted at a branching****

Intrinsic kT-width dependence on √s and invariant mass [Eur.Phys.J.C](https://link.springer.com/article/10.1140/epjc/s10052-024-12507-0) 84 (2024) 2, 154

- Consistent values of qs for a large range of DY pair invariant masses
- Standard Monte Carlo event generators need a strongly increasing intrinsic-kT width with \sqrt{s}
- Strong center-of-mass energy dependence is not observed

- Mimic parton-shower event generators by demanding a minimal parton transverse momentum
- **q0 = 1 and 2 GeV**
	- $qT > q0$
- Non-perturbative part neglected

Introducing energy dependence of the intrinsic-kT in PB

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• Sensitivity of the DY cross section on the intrinsic-kT increases at small pair pT and with increasing of q0 value

qs dependence on center-of-mass energy for the cases with $q0 = 1$ GeV and $q0 = 2$ GeV and $q0 \rightarrow 0$ GeV • The uncertainty is estimated as a range of qs in which:

Introducing energy dependence of the intrinsic-kT in PB [arXiv:2404.04088](https://arxiv.org/pdf/2404.04088)

$$
\chi^2(q_s) - \chi^2_{min} < 1
$$

We have performed a linear fit for the relation $log(q_s) - log(\sqrt{s})$

The uncertainty bands around the fitted lines correspond to the 95% CL band

- The slope increases as q0 increases
- **Larger q0** means that **more soft contributions**
	- Larger intrinsic-kT needed to compensate missing contribution from soft gluons
		- **Higher q0 values lead to an increased sensitivity to the intrinsic kTdistribution**, resulting in **smaller uncertainty bands**

Introducing energy dependence of the intrinsic-kT in PB [arXiv:2404.04088](https://arxiv.org/pdf/2404.04088)

Summary

- DY production at NLO obtained with the **MADGRAPH5_AMC@NLO** event generator matched with the PB TMD distributions **PBNLO-2018 Set2**
- We study Fermi-motion of partons inside proton parameterized by a Gauss distribution of width $\sigma = q_s/\sqrt{2}$
- Proper treatment of the soft contributions in PB method leads to the intrinsic-kT width which does not depend on collision center-of-mass collision energy
	- The inclusion of soft gluons, in particularly the non-perturbative Sudakov, is crucial for providing \sqrt{s} -independent intrinsic-kT

Thank you for your attention!