



# From Feynman integrals to quantum algorithms: *the LTD connection*



**German F. R. SBORLINI**

Departamento de Física Fundamental  
Universidad de Salamanca (USAL)

Prague, 20.07.2024



This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No. 101034371

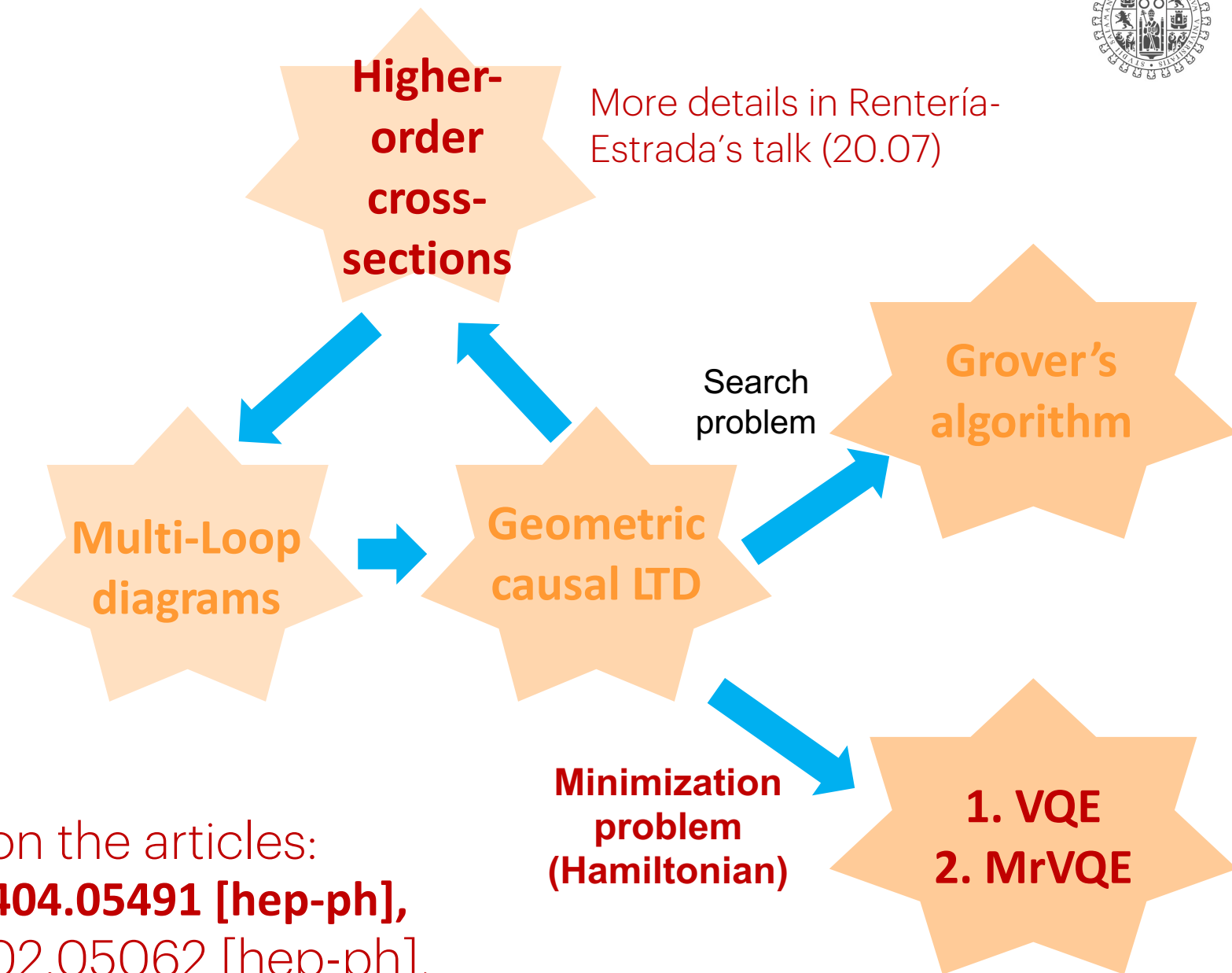


UNIVERSIDAD  
DE SALAMANCA





1. Motivation: Loop-Tree Duality
  - A. Geometrical causality
  - B. Cross-section from vacuum
2. Quantum algorithms for LTD
  - A. Grover search algorithm
  - B. Hamiltonian formulation
  - C. MrVQE
3. Conclusions



This talk is based on the articles:  
**2404.05492 [hep-ph], 2404.05491 [hep-ph],**  
**2210.13240 [hep-ph], 2102.05062 [hep-ph],**  
2001.03564 [hep-ph]



## PROBLEM

High-precision calculations require to deal with LOOPS and REAL-EMISSION. Current approaches calculate them separately.

## QUESTION

Can we do something to combine them?

## ANSWER

We can use...



## PROBLEM

High-precision calculations require to deal with **LOOPS** and **REAL-EMISSION**. Current approaches calculate them separately.

## QUESTION

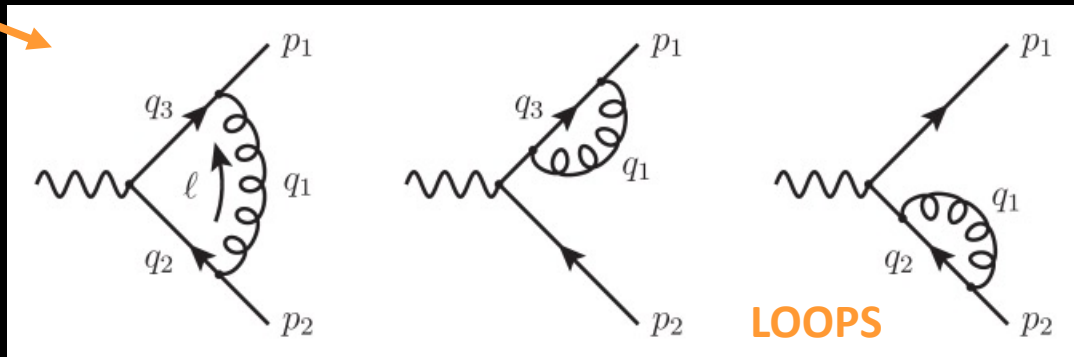
Can we do something to combine them?

## ANSWER

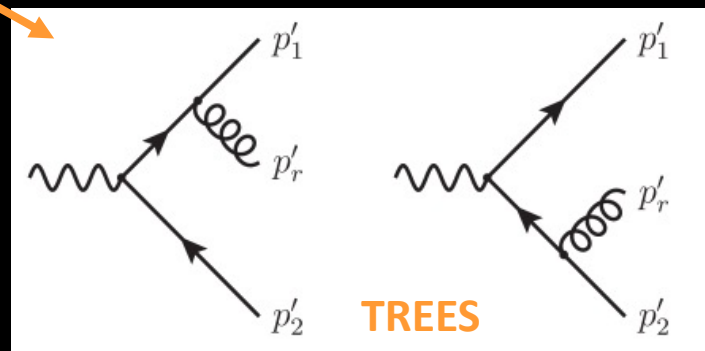
We can use...

# Loop-Tree Duality!

## NLO computations in a single slide



Loops contain a virtual particle, not observable  
**(IR/UV singularities)**



This is the real-radiation, because one extra particle is produced  
**(IR singularities)**



**UV counter-terms**



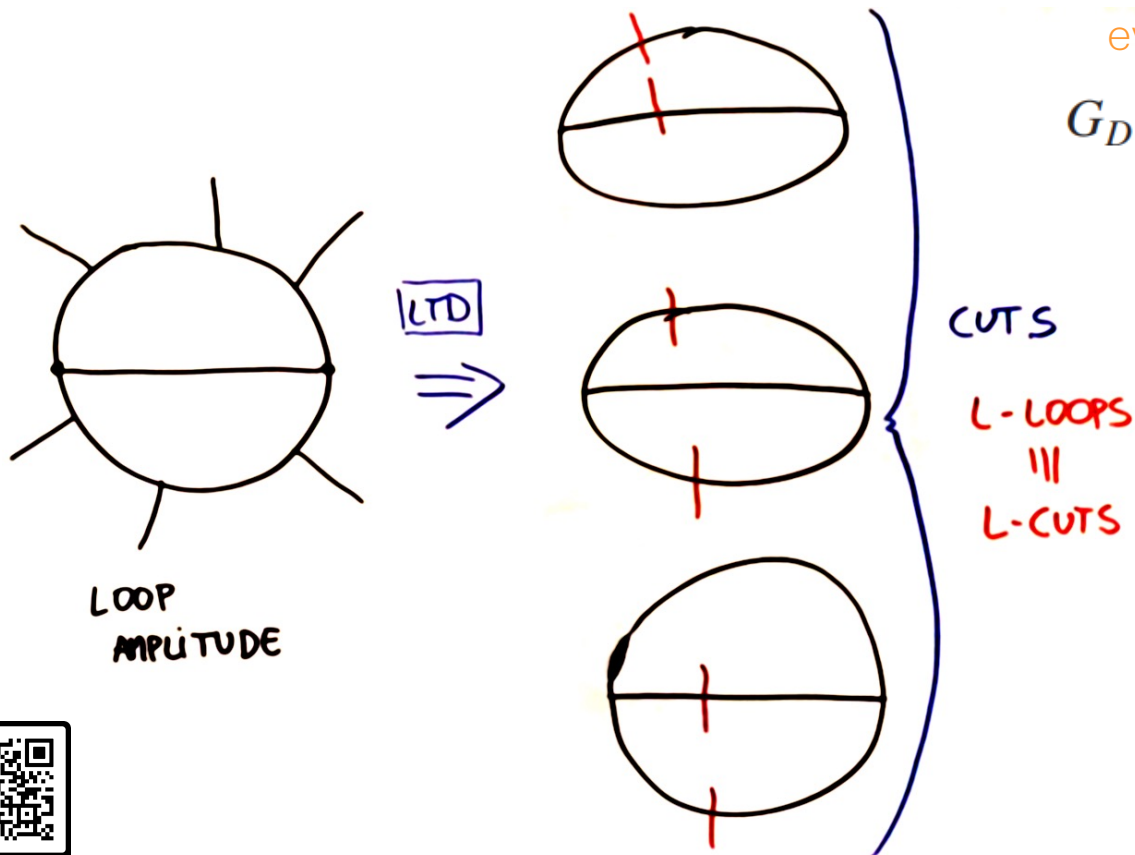
**Finite result!**

# Tackling the loops: Loop-Tree Duality



- **Main strategy:** iterate Cauchy's theorem to *open loops into trees*
- Energy component is removed by using **Cauchy's residue theorem**
- Multiloop require to iterate ("nest") the procedure (remove all the energy components)

Rodrigo & collaborators,  
 PRL 124 (2020) 21, 211602;  
 JHEP 12 (2019) 163; JHEP 01  
 (2021) 069; JHEP 02 (2021)  
 112; JHEP 04 (2021) 129



$r^{\text{th}}$  residue evaluation

Remaining sets (no residue evaluation)

$$G_D(1, \dots, r; n) = -2\pi i \sum_{i_r \in \mathcal{R}} \text{Res}(G_D(1, \dots, r-1; r, n), \text{Im}(\eta \cdot q_{i_r}) < 0)$$

Sum over all the elements of the  $r^{\text{th}}$  set

$(r-1)^{\text{th}}$  dual function

Depends on integration variables ( $q_i$ )

FEYNMAN REP.

$$\frac{1}{q} \equiv \frac{1}{q_0 - q_0^{(+)}} - \frac{1}{q_0 + q_0^{(+)}}$$

$$q_0^{(+)} = \sqrt{|\vec{q}|^2 + m^2 - i0}$$

ON-SHELL ENERGY





**Master formula** 
$$\mathcal{A}_N^{(L)}(1, \dots, L+k) = \sum_{\sigma \in \Sigma} \int_{\vec{\ell}_1, \dots, \vec{\ell}_L} \frac{\mathcal{N}_\sigma(\{q_{r,0}^{(+)}\}, \{p_{j,0}\})}{x_{L+k}} \times \prod_{i=1}^k \frac{1}{-\lambda_{\sigma(i)}} + (\sigma \leftrightarrow \bar{\sigma})$$

Set of entangled thresholds

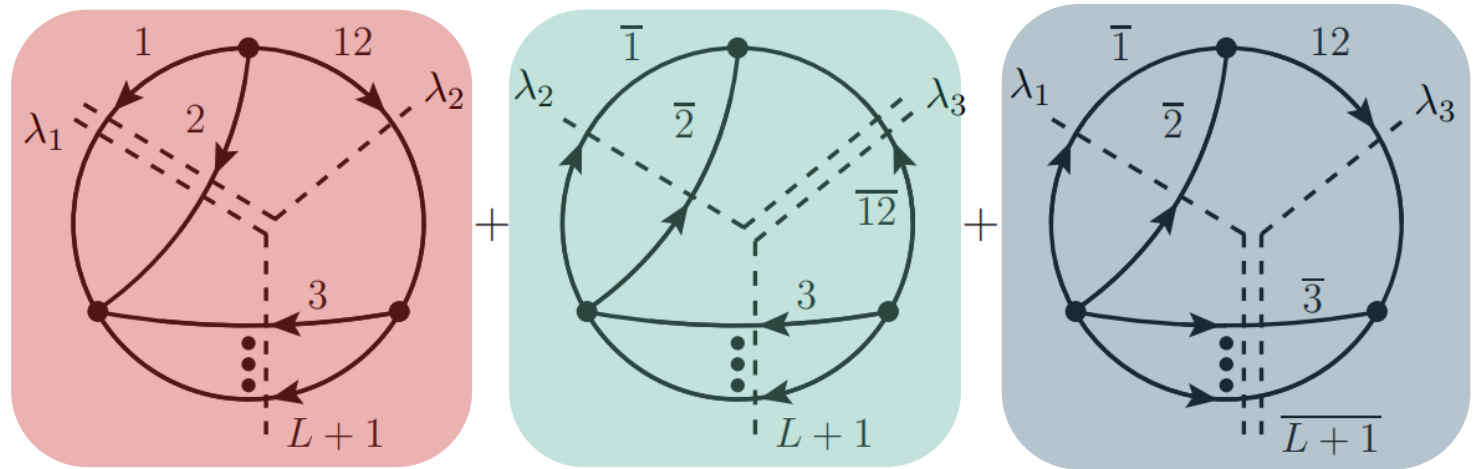
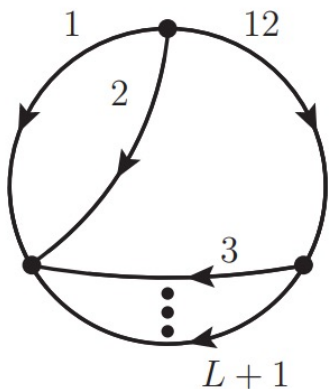
Products of  $k$  causal propagators

• Graphical interpretation in terms of **entangled thresholds**

1. Each **causal propagator** represents a **threshold** of the diagram
2. Each diagram contains **several thresholds**
3. The causal representation involves products of (**compatible**) thresholds

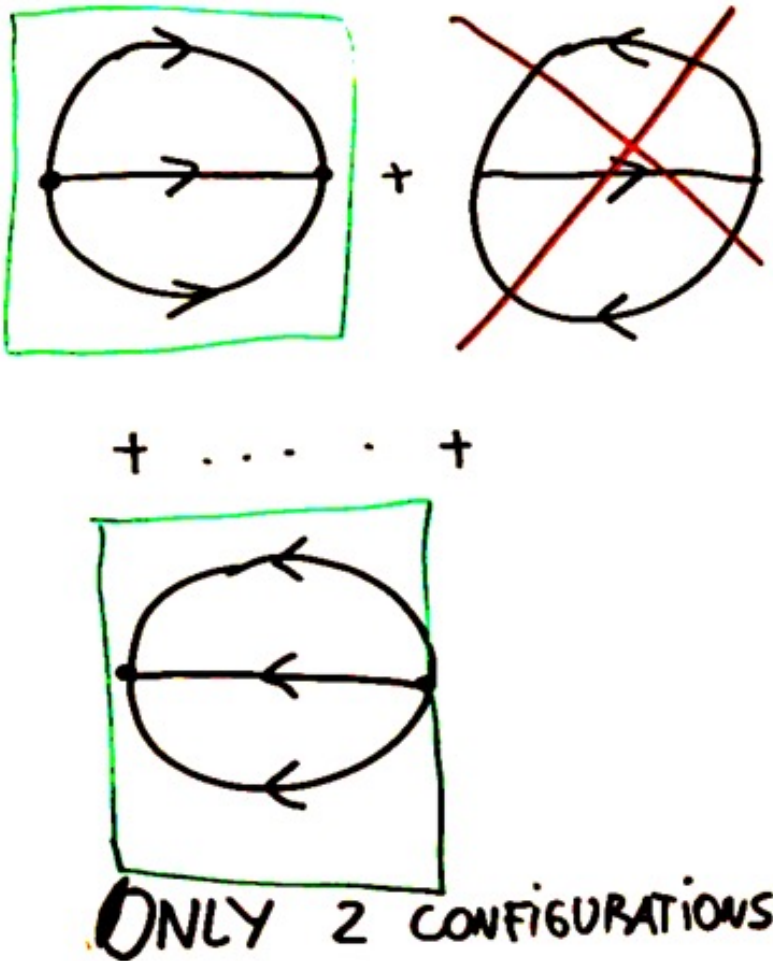
More details in: JHEP 01 (2021) 069; JHEP 04 (2021) 129; JHEP 04 (2021) 183; Eur.Phys.J.C 81 (2021) 6, 514

Causal denominators ( $\lambda$ ) are associated to **cut lines** in the diagrams: momenta flow must be adjusted to be **compatible**



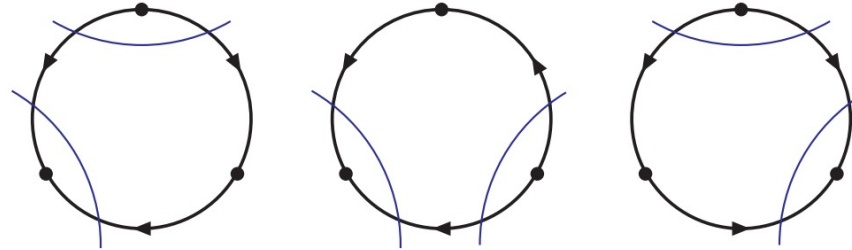
$$\mathcal{A}_{\text{NMLT}}^{(L)}(1, 2, \dots, L+2) = \int_{\vec{\ell}_1, \dots, \vec{\ell}_L} \frac{2}{x_{L+2}} \left( \frac{1}{\lambda_1 \lambda_2} + \frac{1}{\lambda_2 \lambda_3} + \frac{1}{\lambda_3 \lambda_1} \right)$$





- **Geometrical rules look into several properties of the graph**  
“If the graph can be cut in tree-level pieces (thresholds) with consistent momentum flow (entangled thresholds), then we have a causal representation”

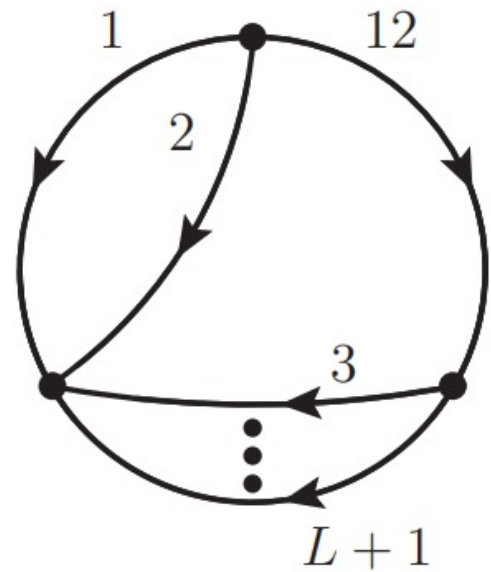
- **IMPORTANT: Causal configurations MUST FULLFIL the causal-compatible flux condition** **Bootstrapping!!**



- **From graph theory:** “Directed oriented graphs that can be split into binary partitions fulfilling (\*) have no closed directed cycles”

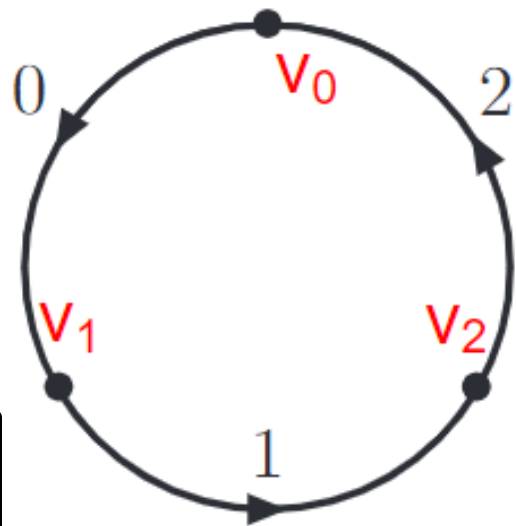
**Non-cyclical configurations = Causal flux**

More detailed explanation  
arXiv:2102.05062 [hep-ph] &  
arXiv:2105.08703 [hep-ph]

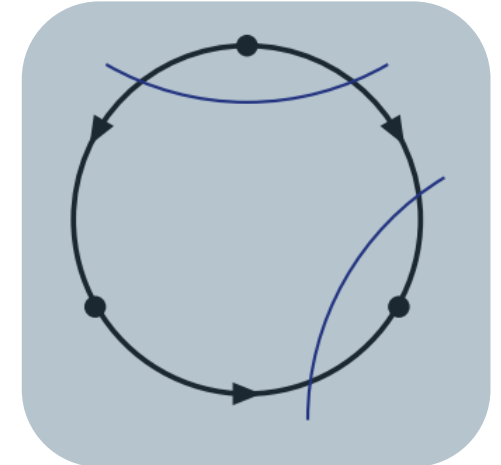
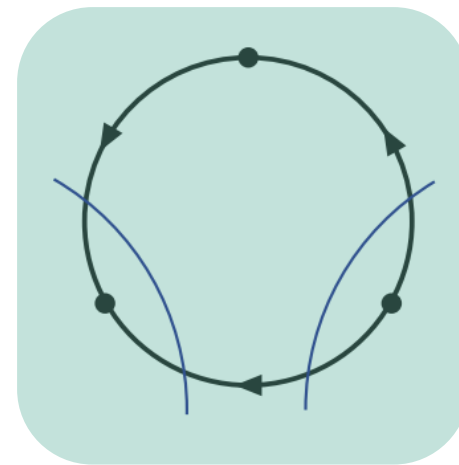
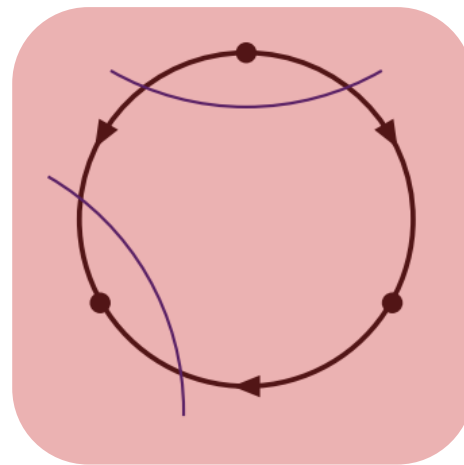


- Geometrical rules look into several properties of the graph  
 “If the graph can be cut in tree-level pieces (thresholds) with consistent momentum flow (entangled thresholds), then we have a causal representation”
- **IMPORTANT:** Causal configurations **MUST FULLFIL** the causal-compatible flux condition  $\rightarrow$  **Bootstrapping!!**

Topologically equivalent diagrams



=



$$= \int_{\vec{l}_1} \frac{2}{x_{L+2}} \left( \frac{1}{\lambda_1 \lambda_2} + \frac{1}{\lambda_2 \lambda_3} + \frac{1}{\lambda_3 \lambda_1} \right)$$



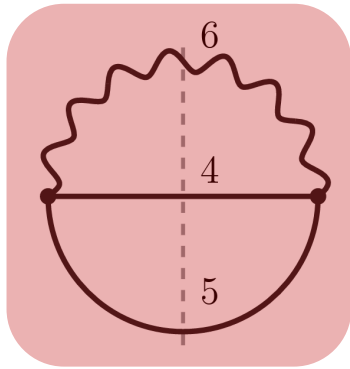




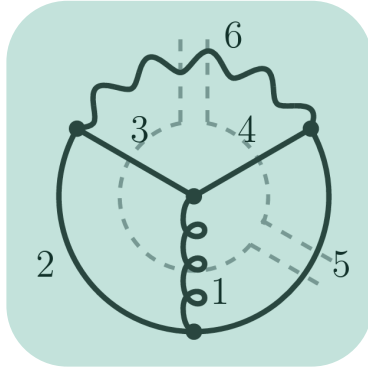
Developed by  
LTD Collaboration

- Novel technology to compute cross-sections: vacuum diagrams!
- Applying residue theorem to the causal LTD representation, real, virtual and IR counter-terms are simultaneously generated! ➔ Fully local cancellation of physical singularities!

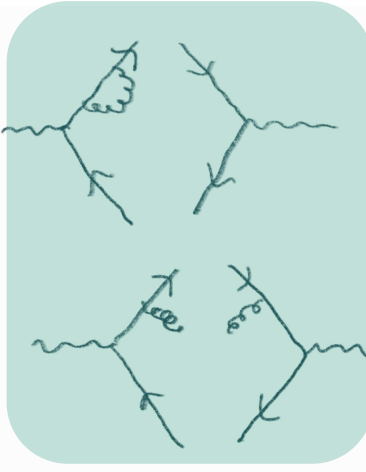
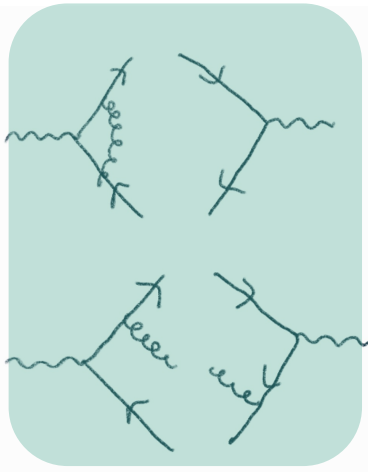
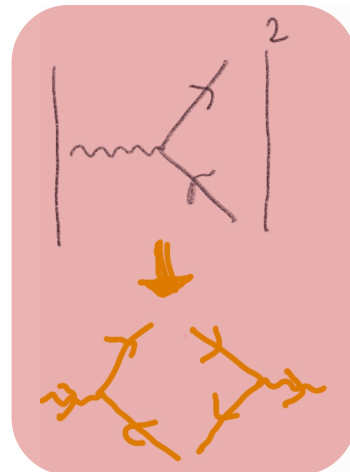
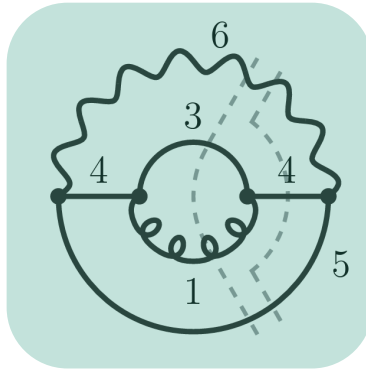
$\gamma^* \rightarrow q\bar{q} \rightarrow b\bar{b}(g)$



Tree-level  
(eq. 2-loop)



Next-to-leading  
order (eq. 3-loop)

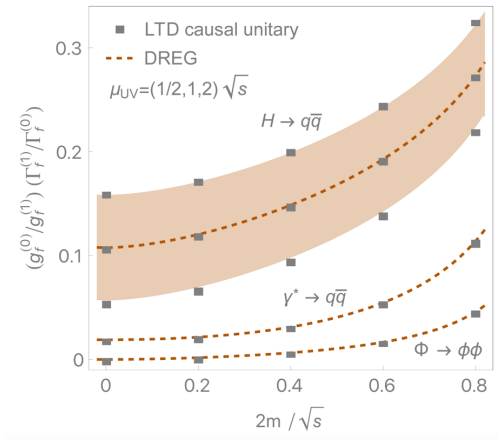


### Master formula

$$d\Gamma_a^{(k)} = \frac{d\Lambda}{2m_a} \sum_{(i_1 \dots i_n a) \in \Sigma} \mathcal{A}_D^{(\Lambda, R)}(i_1 \dots i_n a) \tilde{\Delta}_{i_1 \dots i_n \bar{a}}$$

with  $\mathcal{A}_D^{(\Lambda)}(i_1 \dots i_n a) = \text{Res} \left( \frac{x_a}{2} \mathcal{A}_D^{(\Lambda)}, \lambda_{i_1 \dots i_n a} \right)$   
the residue of the dual amplitude and the momentum conservation delta:

$$\tilde{\Delta}_{i_1 \dots i_n \bar{a}} = 2\pi \delta(\lambda_{i_1 \dots i_n \bar{a}}) \text{ where } \lambda_{i_1 \dots i_n \bar{a}} = \sum_{s=1}^n q_{i_s, 0}^{(+)} - p_{\bar{a}, 0}^{(+)}$$



More details in  
Rentería-Estrada's  
talk (20.07)

Rodrigo et al, arXiv:2404.05491 & 2404.05492 [hep-ph]



## PROBLEM

Complex topologies have many causal configurations; it takes “a lot of time” to test all the possibilities.

## QUESTION

Can we use other techniques to identify the causal terms?

## ANSWER

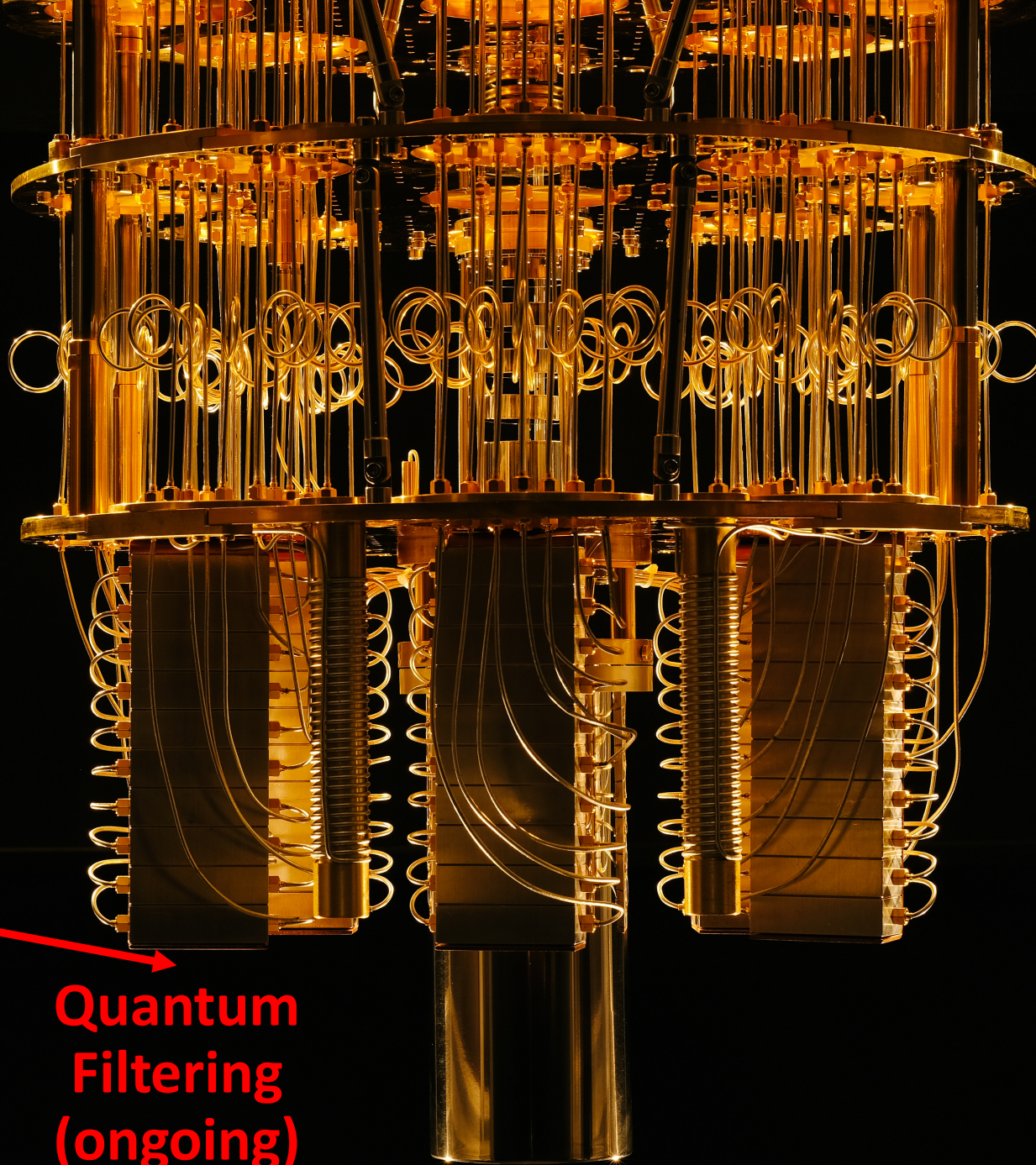
We can explore ...

## Quantum Algorithms!

Quantum  
Search  
(Grover)

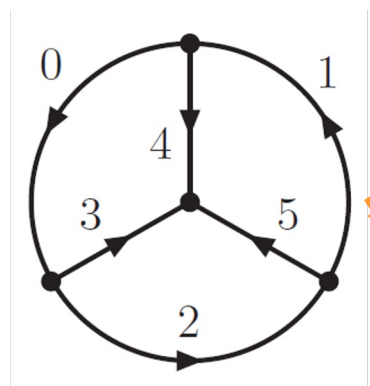
Quantum  
Minimization  
(VQE)

Quantum  
Filtering  
(ongoing)

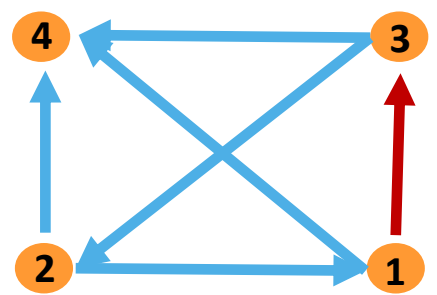




- Geometrical information is codified in the **adjacency matrix**



```
(*Configuracion momentos*)
NumeroVertices = 4; Orden = NumeroVertices - 1;
Eq[1] = {q[6] + q[4] - q[1] + p[1]};
Eq[2] = {q[1] + q[5] - q[2] + p[2]};
Eq[3] = {q[3] - q[6] + q[2] + p[3]};
Eq[4] = {-q[3] - q[4] - q[5] - (p[1] + p[2] + p[3])};
```



$$\begin{pmatrix} -1 & 0 & 0 & 1 & 0 & | & 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 & 1 & | & 0 & 0 & 1 & 0 \\ 0 & 1 & 1 & 0 & 0 & | & -1 & 0 & 0 & 1 \\ 0 & 0 & -1 & -1 & -1 & | & 0 & -1 & -1 & -1 \end{pmatrix}$$

Vertex matrix  
(from the conservation equations)

$$A = \begin{pmatrix} 0 & a_{12} & a_{13} & \dots \\ -a_{12} & 0 & a_{23} & \dots \\ -a_{13} & -a_{23} & 0 & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} = \begin{pmatrix} 0 & 0 & 1 & 1 \\ 1 & 0 & 0 & 1 \\ 0 & 1 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Adjacency matrix  
(how are vertices connected)

**Useful properties:**

$$f_3(A) \equiv \sum_{i_1 \neq i_2 \neq i_3} a_{i_1 i_2} \times a_{i_2 i_3} \times a_{i_3 i_1} \neq 0 \iff \text{There are triangles}$$

$$f_4(A) \equiv \sum_{i_1 \neq i_2 \neq i_3 \neq i_4} a_{i_1 i_2} \times a_{i_2 i_3} \times a_{i_3 i_4} \times a_{i_4 i_1} \neq 0 \iff \text{There are boxes}$$

In general, there are N-cycles iff  $f_N(A)$  is non zero

A graph with associated adjacency matrix A is acyclic if:

$$\sum_{n=1}^V \text{tr}(A^n) = 0$$




- Exploit the **adjacency matrix** to build a **Hamiltonian** → **Ground state = Acyclic graph**
- 1<sup>st</sup> approach:** Penalize oriented cycles using projectors to build the **loop Hamiltonian**

$$H_G = \sum_{\gamma \in \Gamma_{G_0}} \prod_{e \in \gamma} \pi_e^{s(e; G_0)}$$

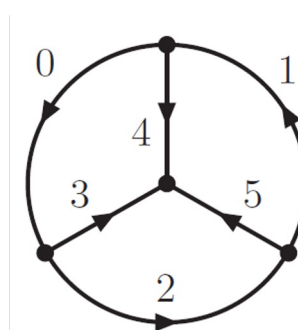
$\Gamma_{G_0}$ : Cycles (within graph)  
 $e \in \gamma$ : Edges (within cycle)  
 $\pi_e^{s(e; G_0)}$ : Projector  
 $s(e; G_0)$ : Orientation (0 or 1)

with

$$\pi_e^0 \equiv |0\rangle\langle 0|_e = \frac{1}{2}(I + Z)_e$$

$$\pi_e^1 \equiv |1\rangle\langle 1|_e = \frac{1}{2}(I - Z)_e$$

0 original direction  
1 reversed direction



Need to fix a initial orientation, but results are independent of this choice!!

- 2<sup>nd</sup> approach (BETTER):** Promote adjacency matrix to operator, and *trace over all possible cycles*

$$A \equiv \sum_{e \equiv (v_0, v_1) \in \bar{E}} \left[ \sigma_{v_0}^- \pi_e^0 \sigma_{v_1}^+ + \sigma_{v_1}^- \pi_e^1 \sigma_{v_0}^+ \right]$$

$e \equiv (v_0, v_1) \in \bar{E}$ : Edge = (origin, end)  
 $\sigma_{v_0}^\pm, \sigma_{v_1}^\pm$ : Pauli +/- operators



$$H_G = \sum_{n=1}^{M_G} \text{Tr}_V [A^n]$$

Trace over vertex space

n = cycle length

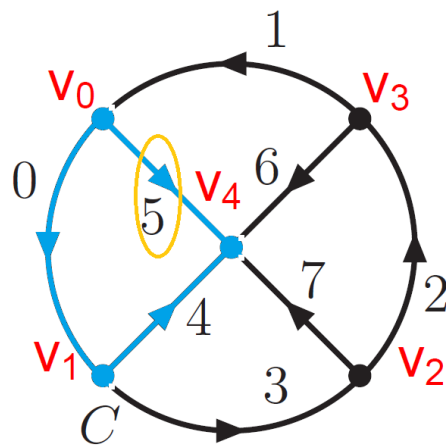
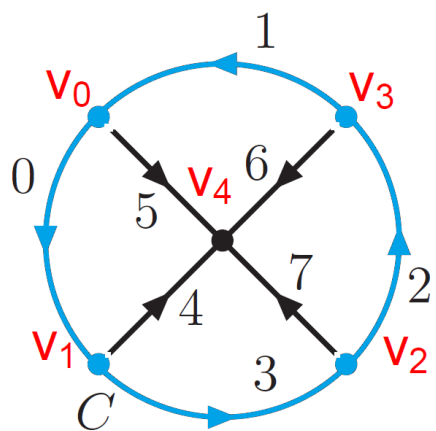
More details arXiv:2210.13240 [hep-ph]

Minimizing H, we find the acyclic graphs (0 energy)





- Exploit the adjacency matrix to build a Hamiltonian  $\rightarrow$  **Ground state = Acyclic graph**



A few comments:

- For each cycle, we have 1 penalization term
- 0 (1) corresponds to the original (reversed) direction w.r.t. the starting oriented graph
- Scaling depends on the connectivity of the graph: lower bound (1) independent on the number of vertices  $V$

$$\begin{aligned}
 H_{\text{Top.C}} = & 4\pi_0^0 \pi_1^0 \pi_2^0 \pi_3^0 + 4\pi_0^1 \pi_1^1 \pi_2^1 \pi_3^1 + 3\pi_0^1 \pi_4^1 \pi_5^0 + 5\pi_1^0 \pi_2^0 \pi_3^0 \pi_4^1 \pi_5^0 + 3\pi_0^0 \pi_4^0 \pi_5^1 \\
 & + 5\pi_1^1 \pi_2^1 \pi_3^1 \pi_4^0 \pi_5^1 + 4\pi_0^1 \pi_1^1 \pi_4^1 \pi_6^0 + 4\pi_2^0 \pi_3^0 \pi_4^1 \pi_6^0 + 3\pi_1^1 \pi_5^1 \pi_6^0 + 5\pi_0^0 \pi_2^0 \pi_3^0 \pi_5^1 \pi_6^0 \\
 & + 4\pi_0^0 \pi_1^0 \pi_4^0 \pi_6^1 + 4\pi_2^1 \pi_3^1 \pi_4^0 \pi_6^1 + 3\pi_1^0 \pi_5^0 \pi_6^1 + 5\pi_0^1 \pi_2^1 \pi_3^1 \pi_5^0 \pi_6^1 + 5\pi_0^1 \pi_1^1 \pi_2^1 \pi_4^1 \pi_7^0 \\
 & + 3\pi_3^0 \pi_4^1 \pi_7^0 + 4\pi_1^1 \pi_2^1 \pi_5^1 \pi_7^0 + 4\pi_0^0 \pi_3^0 \pi_5^1 \pi_7^0 + 3\pi_2^1 \pi_6^1 \pi_7^0 + 5\pi_0^0 \pi_1^0 \pi_3^0 \pi_6^1 \pi_7^0 \\
 & + 5\pi_0^0 \pi_1^0 \pi_2^0 \pi_4^0 \pi_7^1 + 3\pi_3^1 \pi_4^0 \pi_7^1 + 4\pi_1^0 \pi_2^0 \pi_5^0 \pi_7^1 + 4\pi_0^1 \pi_3^1 \pi_5^0 \pi_7^1 + 3\pi_2^0 \pi_6^0 \pi_7^1 \\
 & + 5\pi_0^1 \pi_1^1 \pi_3^1 \pi_6^0 \pi_7^1
 \end{aligned}$$

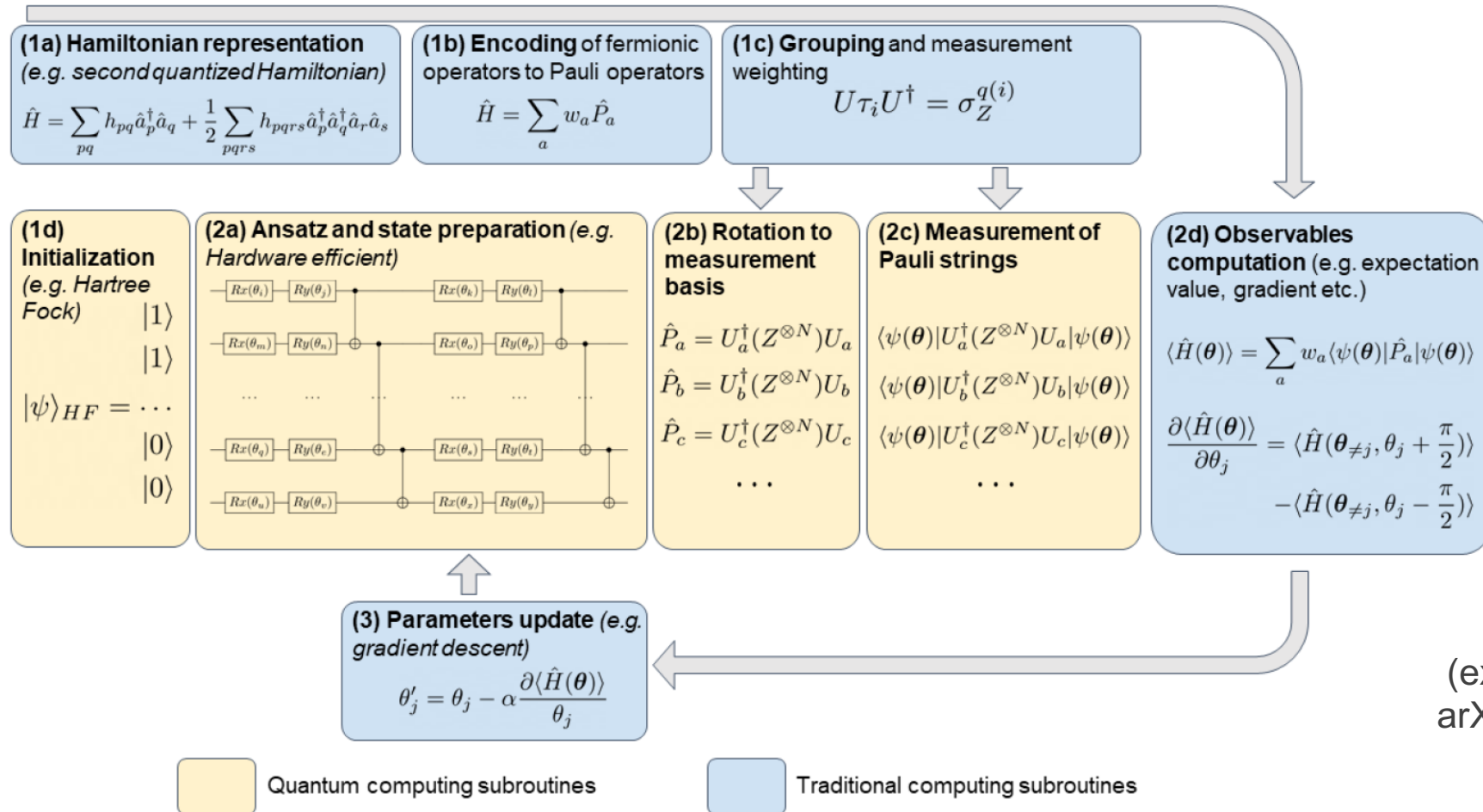
**More details arXiv:2210.13240 [hep-ph]**



# Using a Variational Quantum Eigensolver



- VQE is a hybrid quantum-classical algorithm, optimized for minimization problems
- **QUANTUM PART:** Evaluation of the Hamiltonian applied to an ansatz (parametrized quantum circuit)
- **CLASSICAL PART:** Modification of the parameters, through minimization algorithms



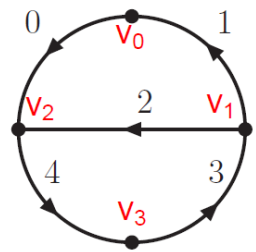
## VQE pipeline

(extracted from J. Tilly et al, arXiv:2111.05176 [quant-ph])





1. Our implementation with Qiskit: **Real Amplitudes** (ansatz) + **COBYLA** (optimizer)
2. Improved results with **multi-run VQE (MrVQE)**: set a selection **threshold**, collect **solutions** and modify the Hamiltonian with **penalization terms**



Two loop 5-point

$$\begin{aligned}
 H_{E/\{e_0\}} = & 4 I_4 \otimes I_3 \otimes I_2 \otimes I_1 + 2 I_4 \otimes I_3 \otimes I_2 \otimes Z_1 - I_4 \otimes I_3 \otimes Z_2 \otimes I_1 \\
 & - I_4 \otimes I_3 \otimes Z_2 \otimes Z_1 + I_4 \otimes Z_3 \otimes I_2 \otimes I_1 + I_4 \otimes Z_3 \otimes I_2 \otimes Z_1 \\
 & + 2 I_4 \otimes Z_3 \otimes Z_2 \otimes I_1 + Z_4 \otimes I_3 \otimes I_2 \otimes I_1 + Z_4 \otimes I_3 \otimes I_2 \otimes Z_1 \\
 & + 2 Z_4 \otimes I_3 \otimes Z_2 \otimes I_1 + 3 Z_4 \otimes Z_3 \otimes I_2 \otimes I_1 + Z_4 \otimes Z_3 \otimes I_2 \otimes Z_1
 \end{aligned}$$

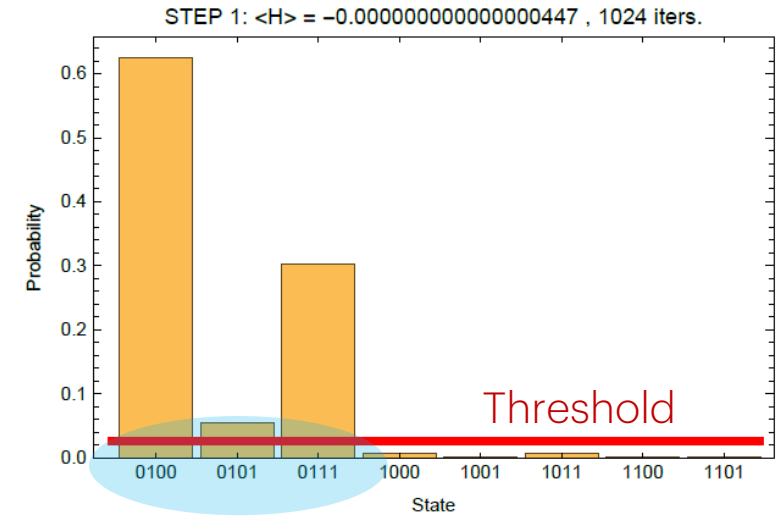
Reduced Hamiltonian

Classical methods



Complete set of solutions

$\{|0011\rangle, |0100\rangle, |0101\rangle, |0111\rangle, |1000\rangle, |1001\rangle, |1011\rangle, |1100\rangle, |1101\rangle\}$



$$\mathcal{S}_1 = \{|0100\rangle, |0101\rangle, |0111\rangle\}$$

Selected solutions (1<sup>st</sup> run)

$$\Pi^{(1)} = \sum_{|\phi_l\rangle \in \mathcal{S}_1} b_l^{(1)} |\phi_l\rangle \langle \phi_l|$$

Penalization term ...

$$H^{(1)} = H^{(0)} + \Pi^{(1)}$$

$$|\psi^{(1)}\rangle = \sum_j c_j^{(1)} |\phi_j\rangle$$

Approximated ground state found by VQE

$$\mathcal{S}_1 = \{|\phi_j\rangle \mid |c_j^{(1)}|^2 > \lambda\}$$

and subset of terms above the selection threshold

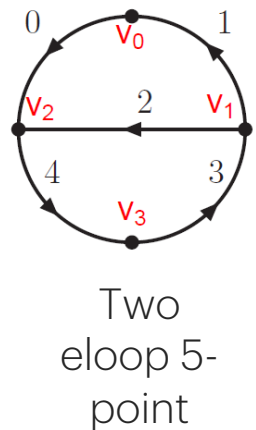
... to be added to the Hamiltonian for next run!!



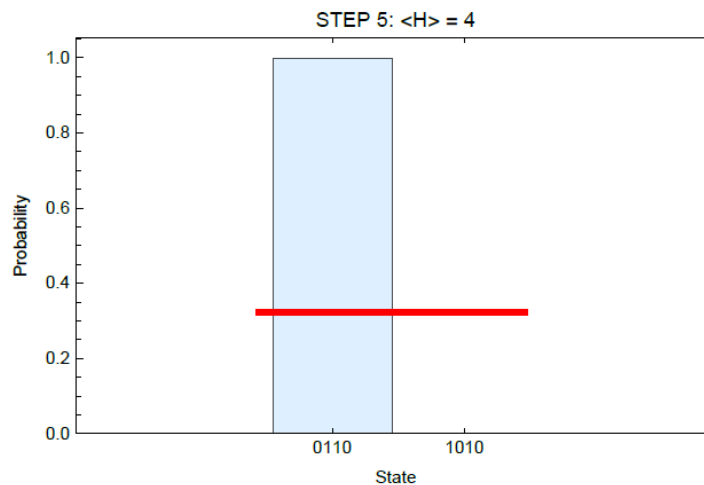
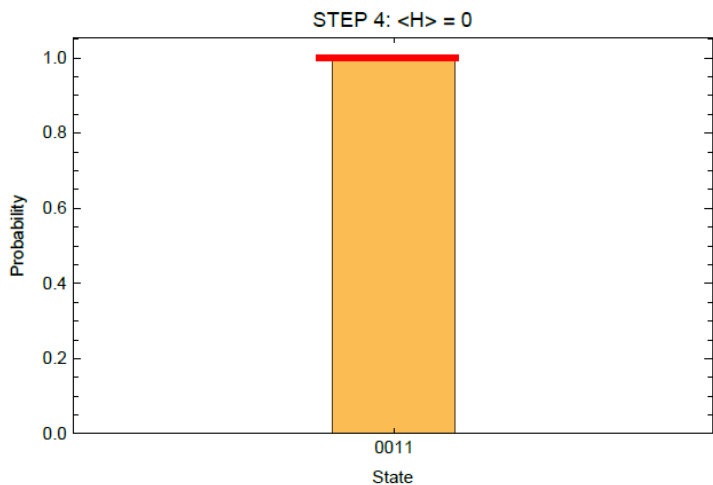
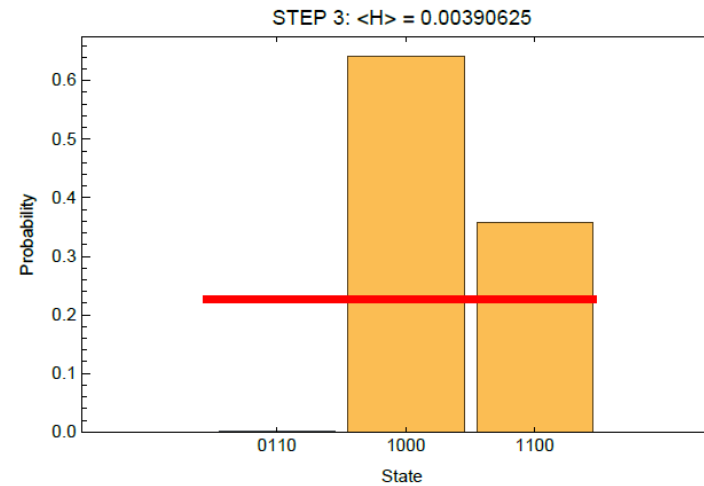
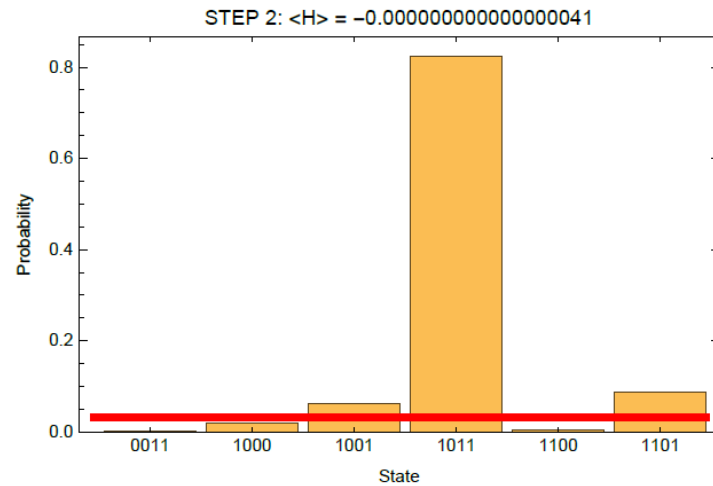
# Using a Variational Quantum Eigensolver



1. Our implementation with Qiskit: **Real Amplitudes** (ansatz) + **COBYLA** (optimizer)
2. Improved results with **multi-run VQE (MrVQE)**: set a selection **threshold**, collect **solutions** and modify the Hamiltonian with **penalization terms**



Two  
elooop 5-  
point



- We collect solutions step by step, till the algorithm converges (if  $\langle H \rangle > 1$ )
- *Problem:* it is not guaranteed that all the solutions are collected (**work in progress!!**)  
**IMPROVED!!!**







- Nested residues leads to **manifestly causal representations** in QFT
- **Geometrical rules** select entangled thresholds. **Complete reconstruction** of multiloop amplitudes!
- **LTD + causality** ➡ **New strategy to compute cross-sections from multi-loop vacuum diagrams** ➡ Several loops required (3 loops = 2->2 tree-level process)

More details in Rentería-Estrada's talk (20.07)

- **Quantum algorithms** to speed-up **causal flux selection**.
- **VQE** is a promising candidate to **calculate Feynman integrals/amplitudes** from minimization problems!
- **Optimized MrVQE** can tackle complicated (several) multiloop diagrams





THANKS!



SCAN ME!

**BACKUP**

- Causal representation obtained directly after **summing over all the nested residues**

Master formula

$$\mathcal{A}_N^{(L)}(1, \dots, L+k) = \sum_{\sigma \in \Sigma} \int_{\vec{\ell}_1, \dots, \vec{\ell}_L} \frac{\mathcal{N}_\sigma(\{q_{r,0}^{(+)}\}, \{p_{j,0}\})}{x_{L+k}} \times \prod_{i=1}^k \frac{1}{-\lambda_{\sigma(i)}} + (\sigma \leftrightarrow \bar{\sigma})$$

Set of entangled thresholds

Products of  $k$  causal propagators

- *Is it possible to do it in other way?*

- **Geometrical reconstruction**

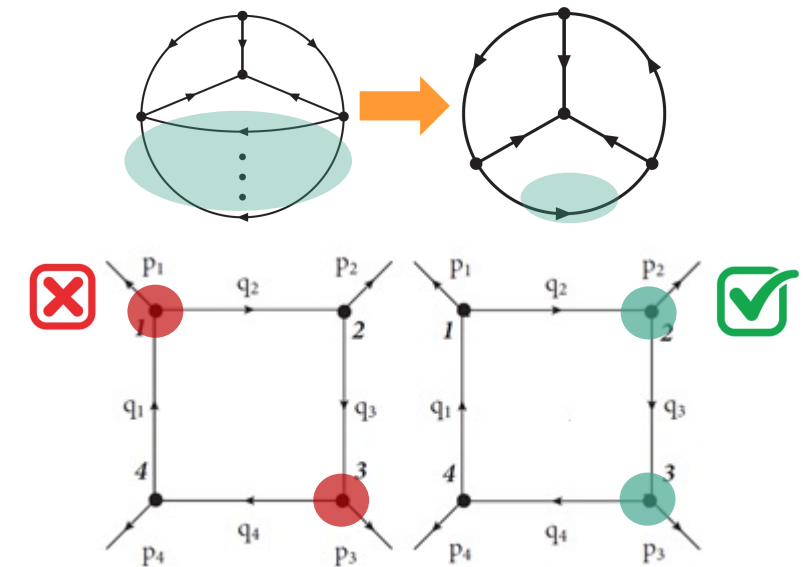
Sborlini '21

- **Algebraic reconstruction (Lotty)**

Torres Bobadilla '21

- **Important GEOMETRICAL concepts:**

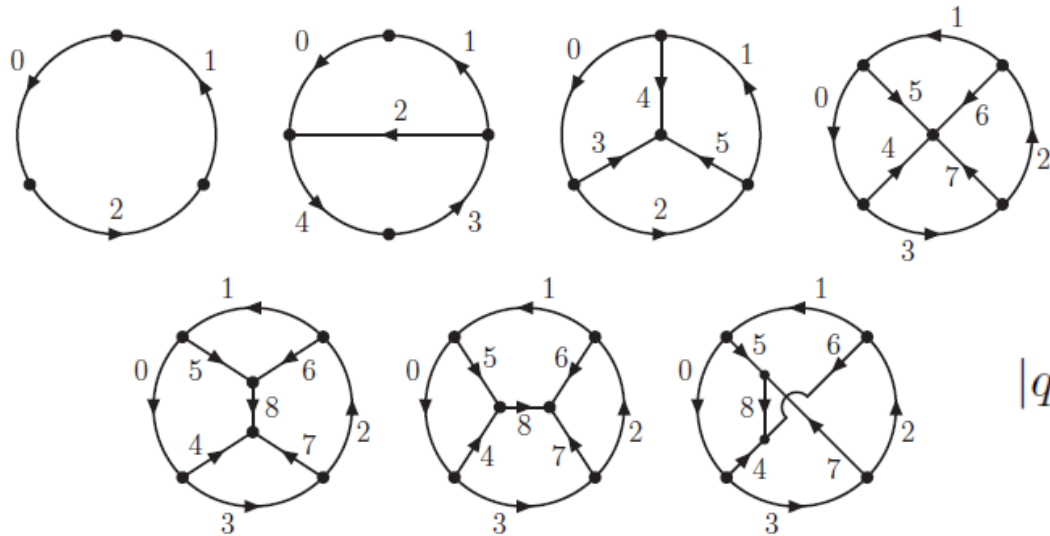
1. **Diagrams** are made of **vertices** and **multi-edges** (*bunches of propagators, connecting two given vertices*)
2. **Multi-edges** define a **basis of momenta**, that lead to the **“vertex” and adjacency matrices** **Defines the casual structure!**
3. **Binary partitions** are given by **subsets of vertices** that **splits in two** the original diagram **Connected partitions!**



# Grover's algorithm for Causal Reconstruction



- Identify momentum-orderings compatible with causality using Grover's search algorithm.
- We assign **1 qubit to each edge**, and impose logical conditions to select configurations without closed directed cycles ➔ **Non-cyclical configurations = Causal flux**
- **Important: "loop"** refers to **integration variables**; "elooop" to loops in the **graph**



Total number of orderings  
( $n = n^{\circ}$  of edges)

$$N = 2^n \quad \text{➔} \quad |q\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle$$

Quantum superposition of N flux configurations

$$|q\rangle = \cos \theta |q_{\perp}\rangle + \sin \theta |w\rangle$$

$$|w\rangle = \frac{1}{\sqrt{r}} \sum_{x \in w} |x\rangle$$

States with causal flow = "Winning states"

$$|q_{\perp}\rangle = \frac{1}{\sqrt{N-r}} \sum_{x \notin w} |x\rangle$$

States with non-causal flow = "Orthogonal states"

- We use Grover's algorithm to **enhances** the probability of the **causal states**:

$$U_w = \mathbf{I} - 2|w\rangle\langle w|$$

Oracle operator  
(changes sign of causal states)

$$U_q = 2|q\rangle\langle q| - \mathbf{I}$$

Diffusion operator  
(reflects with respect to initial state)

$$\text{➔} \quad (U_q U_w)^t |q\rangle = \cos \theta_t |q_{\perp}\rangle + \sin \theta_t |w\rangle$$

with

$$\sin^2 \theta_t \sim 1$$