

# Formal Theory

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## “Formal Theory”

What is “Formal Theory”?

~~Formal as in ... “Theory, that is not useful for any actual physics”?~~

**Theory, which develops a fundamental understanding of Quantum Field Theory and Quantum Gravity, consolidating insights from various areas of theoretical physics and mathematics.**

**Disclaimer: Often this requires studying idealized systems, that are simpler than real world setups. But we will see many avenues to connect to the latter.**

## “Formal Theory”

A (biased<sup>a</sup>) selection of topics with exciting, recent progress:

0. **Scattering Amplitudes: QFT, Strings, Gravity**  
⇒ **Formal Parallel Session at ICHEP** organized by Jusinkas, Trnka, Volovich.
1. **Quantum Gravity constraints on IR physics**
2. **Holography:**
  - (a) AdS/CFT precision holography
  - (b) Quantum Information and Quantum Gravity
  - (c) Flat-Space holography: Celestial and Carrollian holography
3. **Generalized symmetries: Non-invertible symmetries**

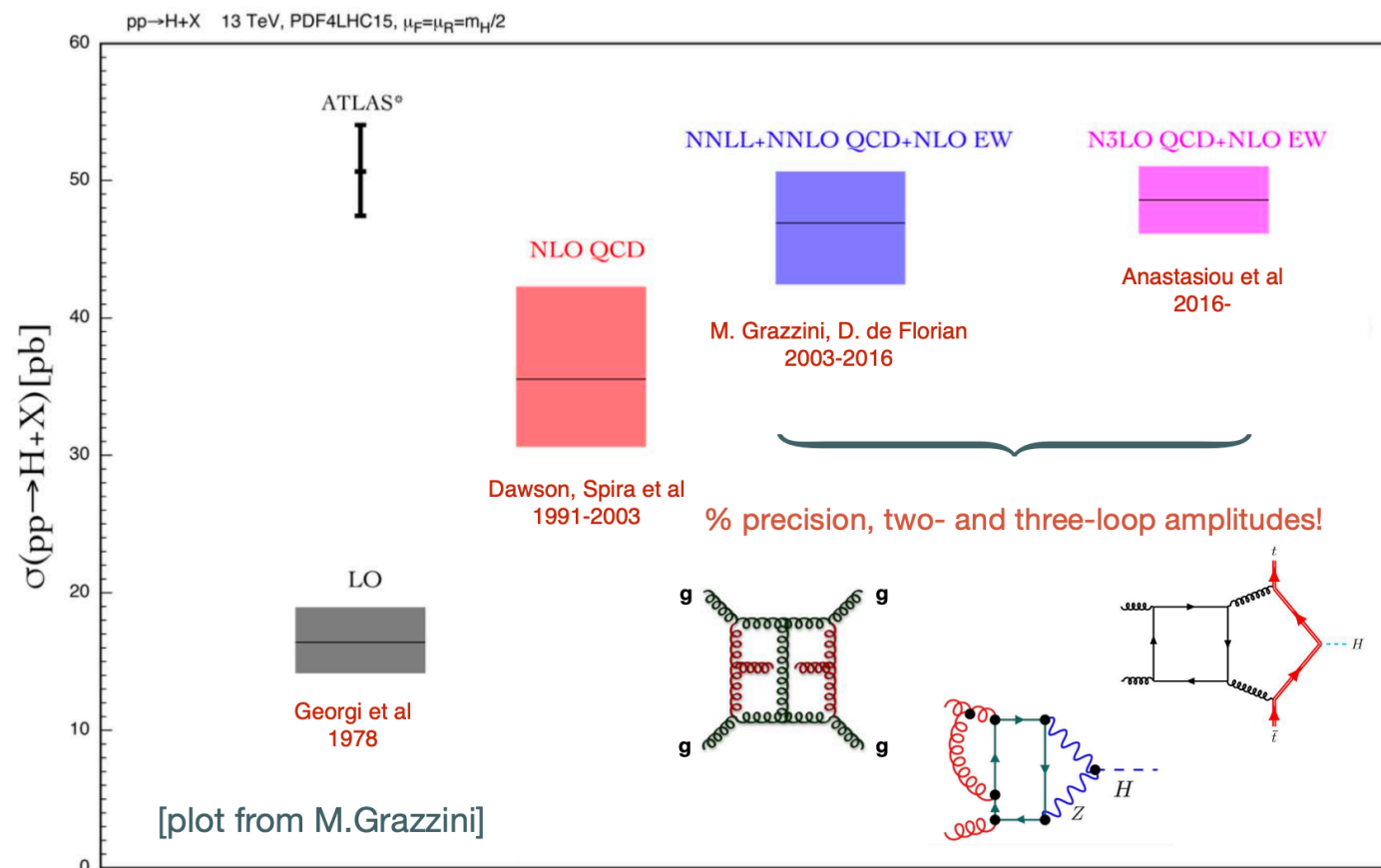
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<sup>a</sup>25 mins to cover hep-th: 2023 totals: 3834 articles + 3452 cross-lists

# 0. Scattering Amplitudes

# Scattering Amplitudes: From Formal to Precision Physics

Amplitudes are of direct relevance for current **Collider Physics**:  
 LHC requires % precision predictions: **QCD+ EW**, higher loops and legs.



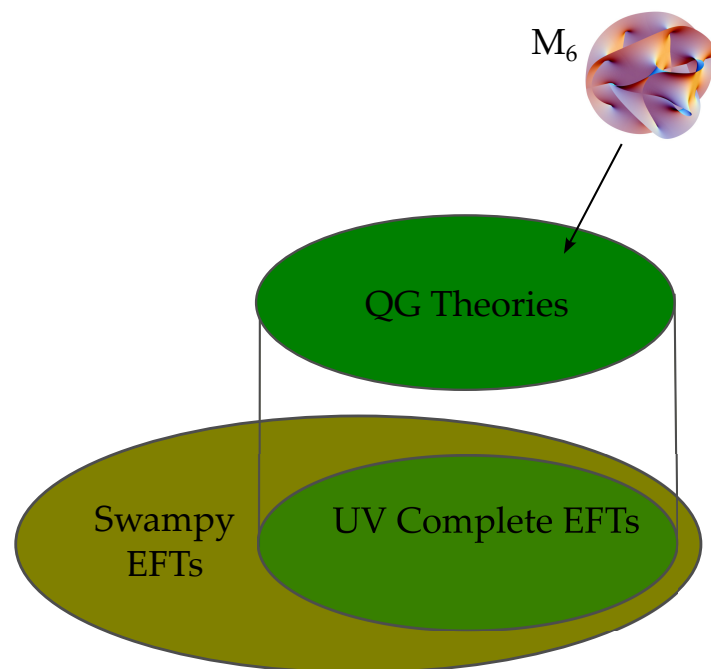
Current status of amplitudes in QCD:  $2 \rightarrow 2$  @ 3L,  $2 \rightarrow 3$  @ 2L

## Formal Session: Amplitudes Talks at ICHEP

Formal theory tools geared towards QCD-like theories (super-YM, etc) that have fed directly into phenomenologically relevant results:

- **BCFW Recursion Relations and Unitarity Methods**, Mathematical tools:  
Drummond, Parisi
- **Amplituhedron**: for  $N = 4$  Super-Yang Mills, but very powerful in producing integrands.
- **String Amplitudes**: quantum gravity, applications to holography  
Monteiro, Schlotterer, Brown<sup>2</sup>, Klisch, Lipstein
- **Double Copy**: Color-Kinematics Duality when relating gauge to gravity scattering  
Brandhuber, Carrasco, Chen, Travaglini.
- **Surface-ology**: Most recent development: promising **Loop integrals for (in principle) any colored theory** using curves on fat graphs.

# 1. Quantum Gravity (QG) Constraints on IR Physics



### Traditionally: Top-down

10d string theory  $\mathcal{T}_{10}$  compactified to 4d gives EFT, e.g. SM spectrum + X (where X is usually  $N = 1$  susy, exotics)

$$\mathbb{R}^{1,9} \rightarrow \mathbb{R}^{1,3} \times M_6$$

Defines at low energies a UV completable EFT  $(\mathcal{T}_{10}, M_6)$ .

### Modern approach: Bottom-up/Swampland Program

Question 1: Given an EFT, can it be embedded into a consistent theory of QG?

Question 2: String Universality? Do all QGs come from string theory?

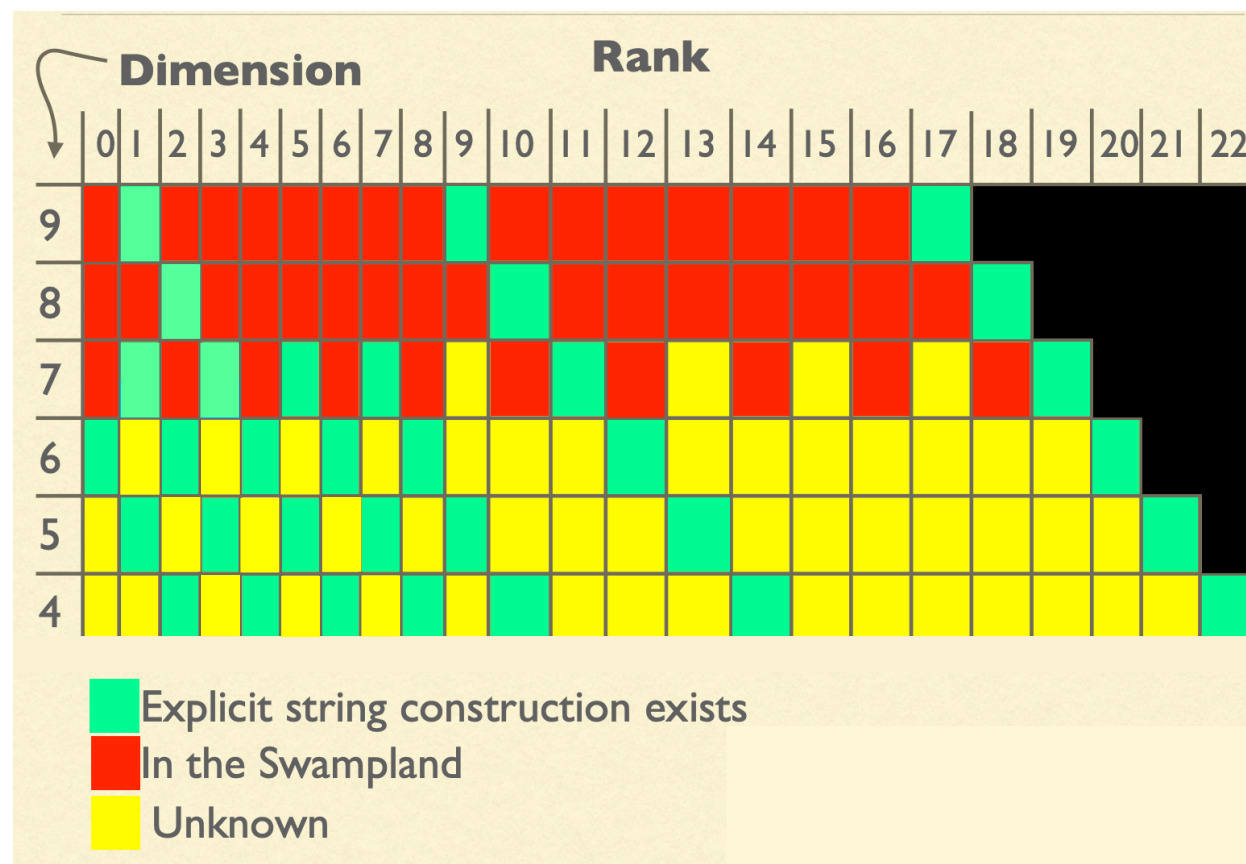
- **Distance conjecture** "infinite distance in parameter space results in infinite tower of states", i.e. breakdown of EFT.
- **Weak Gravity** "gravity is weakest force"
- **No global symmetries** "all global symmetries gauged or broken"

Some recent progress on both questions – albeit in theories with so far little pheno relevance.



## String universality: Abelian gauge theories with 16 susies

Example: 8d and 9d abelian gauge theories with 16 susies. Find agreement of top down and bottom up and confirms string universality.



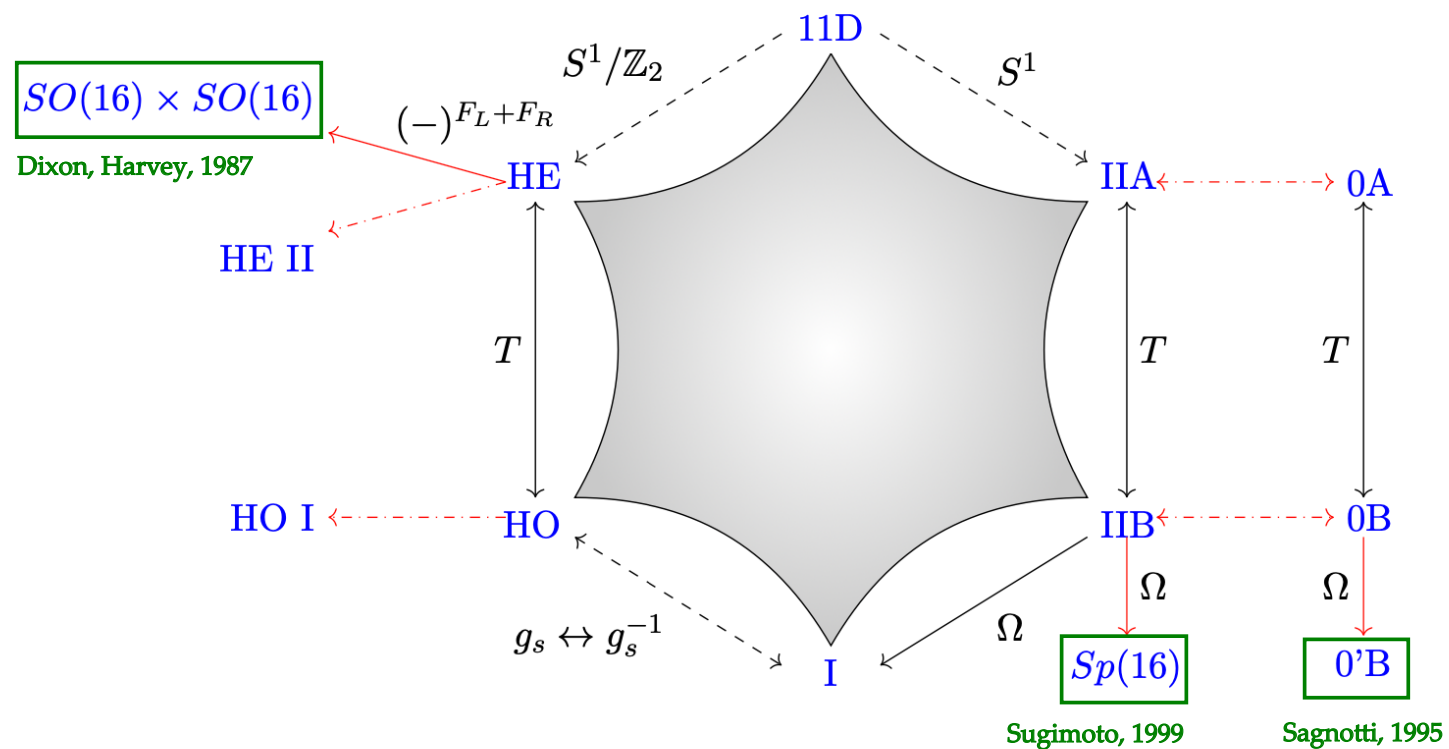
[modified from: Montero: Strings 2024 talk]

Similar arguments, with less stringent results thus far for 8 and 4 supersymmetries.

Status: Sharpening theoretical tools.

## What about no susy?

What is the status of string models with **no supersymmetry**? Start in 10d.



[from 2310.06895]

There are 3 **non-supersymmetric tachyon-free** string theories in 10d.

## 10d Non-Supersymmetric Tachyon-Free Strings

$SO(16)^2$ Het (1987)	$Sp(16)$ Sugimoto (1999)	0'B Sagnotti (1995)
$E_8 \times E_8$ Het/ $(-1)^F$	IIB + $O9^+$ + 32 $\overline{D9}$	IIB/ $(-1)^F/\Omega$
Heterotic $g, \phi, B$	Closed strings: $\mathcal{N} = 1$ susy	Closed: metric dilaton
No gravitino	Open strings: non-susy	non-susy
$SO(16)^2$ gauge group	$Sp(16)$ gauge group	$U(32)$ gauge group

Recent progress:

1. Completeness proven: these are **all** 10d non-susy heterotic theories [Boyle Smith, Lin, Tachikawa, Zheng, 2023]
2. [Basile, Debray, Delgado, Montero, 2023] showed using cobordism theory, **all local and global anomalies cancel**, i.e. not connected to identity:  
Prediction of new extended objects (branes)
3. Analysis of moduli stabilization [Sagnotti et al, (2023, 24)]
4. New  $d < 10$  non-supersymmetric compactifications [Baykara, Tarazi, Vafa, 2024]

Status: systematic progress exploring non-susy string theories.

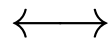
## 2. Holography

Why is Holography important? Gauge/gravity and Strong/weak dualities. Provides conceptual and computational window into quantum gravity and strongly-coupled QFTs alike.

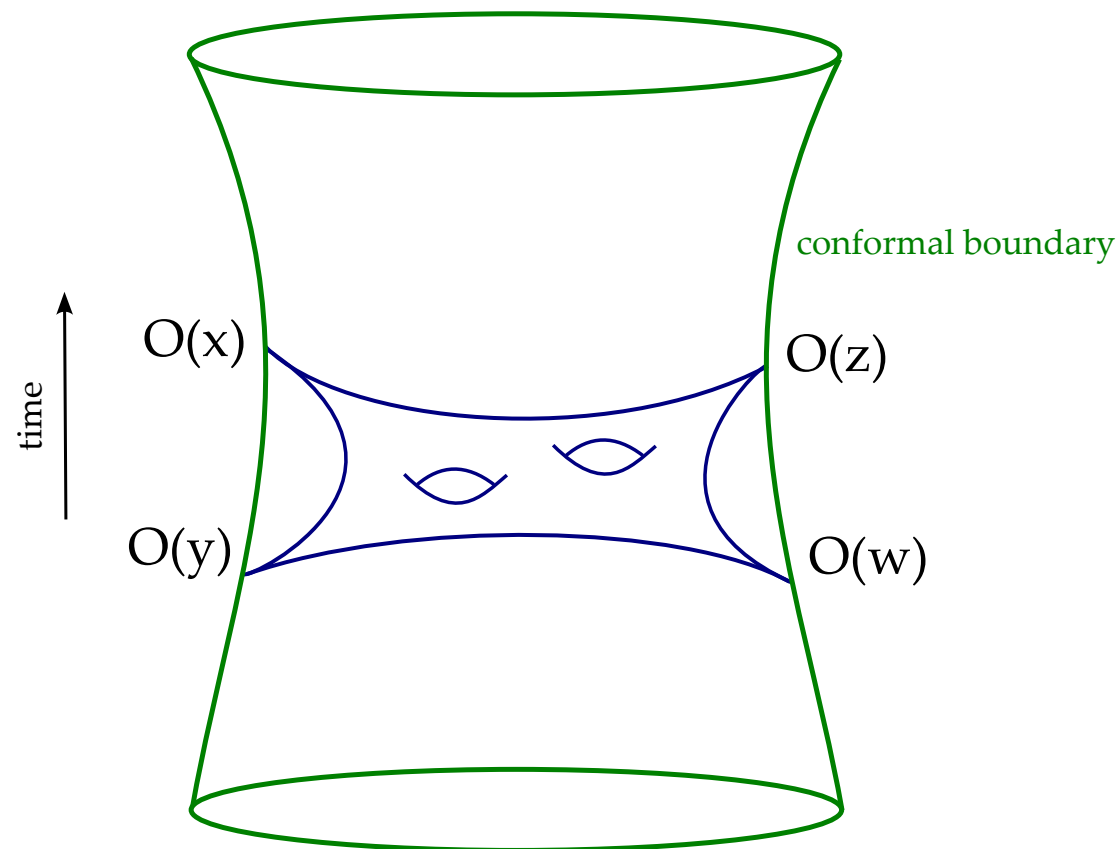
## 2.1. AdS/CFT Precision Holography

# AdS/CFT

String theory on  $\text{AdS}_{d+1} \times X$



$\text{CFT}_d$  on boundary of  $\text{AdS}_{d+1}$



Precision (in coupling constants on both sides) lab for quantum gravity and strongly-coupled CFTs alike.

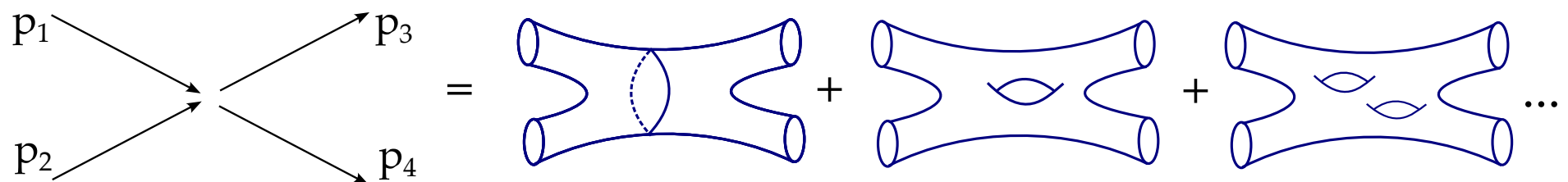
- Conformal group = isometries of AdS.
- Parameters: string coupling  $g_s = \frac{1}{N_c}$  and length  $\frac{R^2}{\alpha'} = \sqrt{\lambda} = \sqrt{g_{YM}^2 N_c}$

1. AdS<sub>3</sub>/CFT<sub>2</sub>: **exact duality proven** [Eberhardt, Gaberdiel, Gopakumar]  
 AdS<sub>3</sub> × S<sup>3</sup> × T<sup>4</sup> (with 1 unit of NSNS flux) dual to Sym<sup>N</sup>(T<sup>4</sup>) 2d CFT.

2. AdS<sub>5</sub>/CFT<sub>4</sub>: Example of deriving quantum gravity from QFT:

Reconstructing string amplitudes from QFT:

Construct string amplitudes in AdS<sub>5</sub> × S<sup>5</sup> (**HARD**). E.g. 4-graviton scattering in AdS<sub>5</sub> × S<sup>5</sup>:



This is an expansion in **string-loops**, i.e.  $1/N_c$ . Even tree-level in curved spacetimes is very difficult.

**Tree-level:** Virasoro-Shapiro amplitude has an expansion in  $\alpha'$  or 't Hooft coupling  $1/\sqrt{\lambda}$

$$A(S, T) = A^{(0)}(S, T) + \frac{A^{(1)}(S, T)}{\sqrt{\lambda}} + \dots$$

where the flat space amplitude is

$$A^{(0)}(S, T) = \frac{1}{U^2} \int d^2 z |z|^{-2S-2} |1-z|^{-2T-2} = -\frac{\Gamma(-S)\Gamma(-T)\Gamma(-U)}{\Gamma(S+1)\Gamma(T+1)\Gamma(U+1)}$$

Quantizing the string in  $\text{AdS}_5 \times S^5$  is notoriously difficult (RR-fluxes).  
Usually people leave it at supergravity level.

Using insights from

- conformal bootstrap and localization
- integrability
- number-theory

[Alday, Hansen, Silva] constructed (conjecturally) the subleading terms and wrote them in terms of a world-sheet-type integral with insertions:

$$A_4^{AdS}(S, T) \sim \int d^2z |z|^{-2S} |1-z|^{-2T} W_0(z, \bar{z}) \left( 1 + \frac{S^2}{R^2} W_3(z, \bar{z}) + \frac{S^4}{R^4} W_6(z, \bar{z}) + \dots \right)$$

$W_n$ 's are single-valued polylogarithms of weight  $n$

New progress on String Field Theory in backgrounds like  $\text{AdS}_5 \times S^5$  should soon be able to test this.

Status: AdS/CFT is now a precision lab for strongly-coupled QFTs and quantum gravity. Not real world holographic duals, but many important lessons learned from these correspondences.



## 2.2. Quantum Information and Quantum Gravity

Lessons:

1. Ordinary quantum systems, e.g. spin chains, can have emergent gravitational properties.

SYK/JT:  $N$  Majorana fermions in 1+1d with random couplings

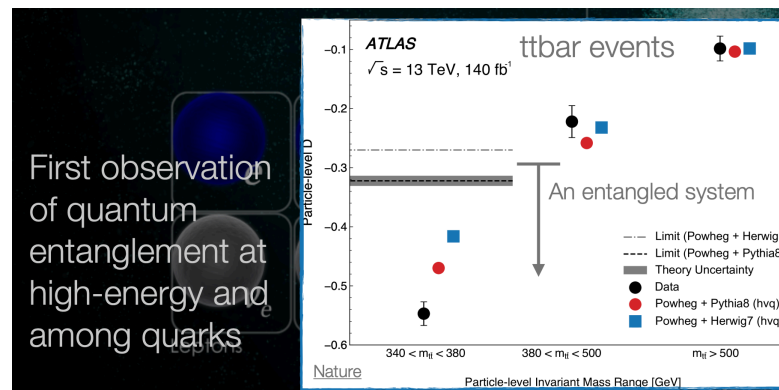
$$H_{\text{SYK}} = \sum_{i_1, i_2, \dots, i_q} J_{i_1 i_2 \dots i_q} \psi_{i_1} \psi_{i_2} \dots \psi_{i_q}$$

At large  $N$ , **emergent gravity**:

$\Rightarrow$  Emergent conformal symmetry with  $c \approx \frac{N^2}{2}$

$\Rightarrow$  Entropy is akin to black holes:  $S \approx \frac{N^2}{4} \log T$ .

2. Quantum info concepts have become key tools in hep-th:
  - (a) Spacetime emerges from quantum entanglement: ER= EPR, entangled particles are connected by wormholes.
  - (b) Holography: Entanglement entropy = Area (minimal surface)/(4G<sub>N</sub>).
  - (c) hep-ph/ex applications: **measurement of top-quark entanglement**



[ATLAS, see Dunford's ICHEP talk]

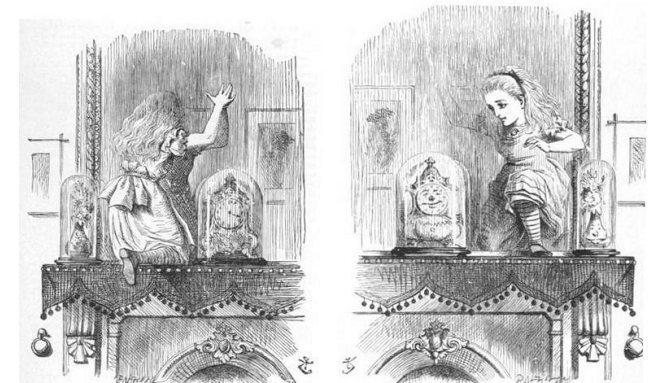
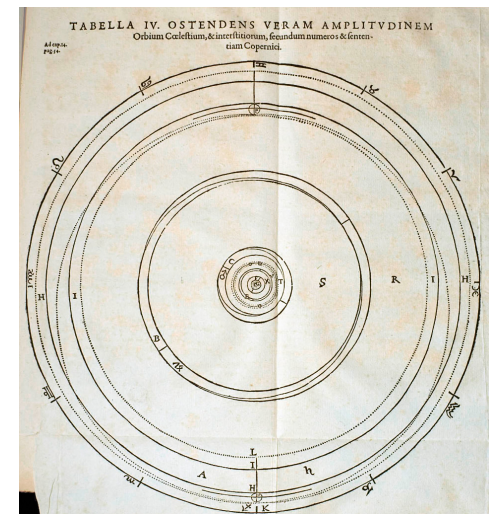
3. Full quantum gravitational path integral requires sum over all topologies of spacetime. This is well established now in AdS/CFT duality. For the future: relevant for cosmology (inflationary models e.g.).

## 2.3. Flat Space Holography

Is there a holographic dual to quantum gravity in flat space? (or asymptotically flat spacetimes (AFS)).

Two recently developed approaches for gravity in AFS spacetimes:

- **Celestial holography:**  
( $d + 2$ )-dim AFS dual to  $d$ -dim CFT on celestial sphere at future/past null-infinity  $\mathcal{I}^\pm$   
"4d gravity is dual to 2d CFT on a celestial sphere  $S^2$ "
- **Carrollian holography:**  
( $d + 1$ )-dim AFS dual to  $d$ -dim Carrollian field theory on  $\mathcal{I}^\pm$   
"4d gravity is dual to a 3d Carrollian<sup>a</sup> ( $c \rightarrow 0$ ) field theory at null infinity"

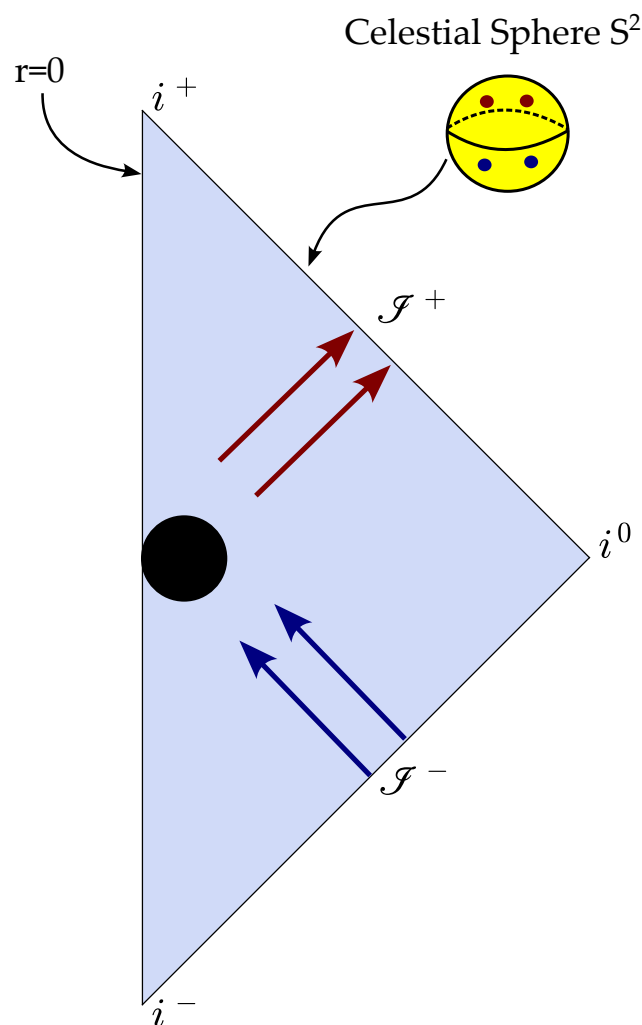


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<sup>a</sup>Peculiar causal structure, like in Lewis Carroll's "Through the Looking Glass"

# Celestial Holography

- 4d Lorentz group = conformal group on  $S^2 \subset$  null infty  $\mathcal{I} = \mathbb{R} \times S^2$ .
- Symmetries = Bondi-Metzner-Sachs (BMS) symmetries extended to  $w_{1+\infty}$
- **Scattering amplitudes in 4d become correlators in 2d celestial CFT (CCFT)!**



S-matrix between  $\mathcal{I}^-$  and  $\mathcal{I}^+$  becomes correlator in 2d CCFT:

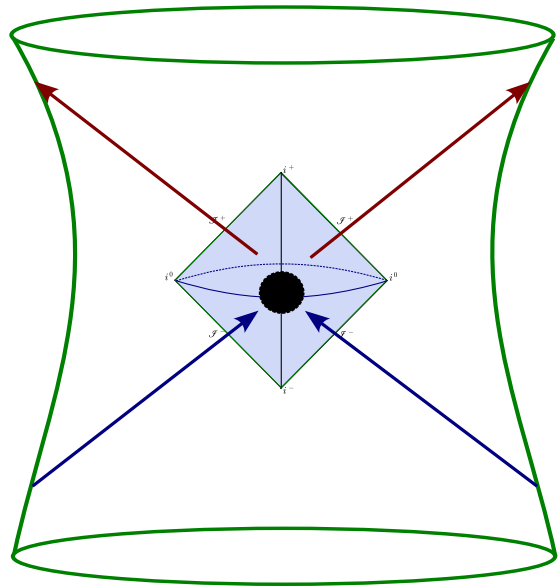
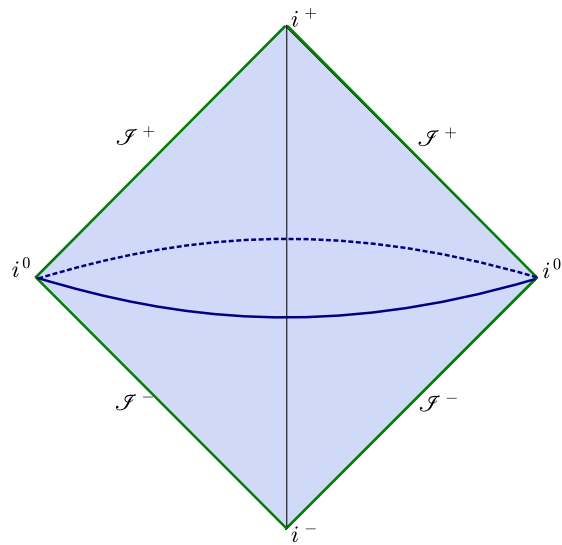
$$\langle p_n^{\text{out}} \dots | \mathcal{S} | \dots p_1^{\text{in}} \rangle \longleftrightarrow \langle \mathcal{O}_{\Delta_1}^- (x_1) \dots \mathcal{O}_{\Delta_n}^+ (x_n) \rangle$$

# **BMS symmetry** of CCFT and implies via Ward identities **soft graviton theorems**

# **Gravitational memory effects** (permanent displacement of test particles (detectors) due to the passage of gravitational waves), correspond to change of state in CCFT.

Directly related to soft theorems/BMS symmetries.

# Carrollian Holography



$d + 1$  dimensional AFS gravity dual to  $d$ -dimensional Carrollian field theory living on  $\mathcal{I} \simeq \mathbb{R} \times S^{d-1}$ .

1. **Carrollian limit**  $c \rightarrow 0$  of Poincaré group: lightcones collapse along  $t$ -axis, so spacetimes events are causally disconnected.
2. **Symmetries:**  
BMS algebra is the conformal Carrollian algebra
3. **Limit from AdS/CFT:**  $c \rightarrow 0$  of relativistic CFTs dual to flat-space local patch of AdS. Many studies of conformal to Carrollian (non-relativistic) limits. Using AdS results to get flat space results!

Status: Very active area. Many open questions (e.g. what precisely is the CCFT?), great potential to learn about flat space QG, and scattering.

### 3. Non-Invertible Symmetries

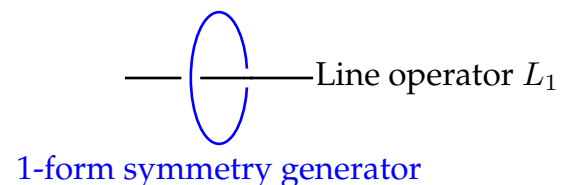
## New Structures: Generalized Symmetries

Symmetries are vital – spectrum, interactions, SSB phases, anomalies, etc.

Usually: Symmetries = **Groups**, acting on **point-like operators**.

There are new symmetries in town: Relax **acting on points** and form **group**:

- **Higher form symmetries**: [Gaiotto, Kapustin, Seiberg, Willett]  
**Example**: Line operators charged under 1-form symmetry (“center symmetry  $\mathbb{Z}_N$  of  $SU(N)$  gauge group”).



**Physics**: Confinement = 1-form symmetry preserving phase.

- **Non-invertible symmetries**: [Kaidi, Ohmori, Zheng][Choi, Cordova, Hsin, Lam, Shao][Bhardwaj, Bottini, SSN, Tiwari]  $a, b \in \mathcal{S}$  a non-invertible symmetry, have no inverse and compose as follows:

$$a \cdot b = n_{c_1} c_1 + \cdots + n_{c_k} c_k, \quad c_i \in \mathcal{S}, n_{c_i} \in \mathbb{N}$$

**Higher-form and Non-invertible symmetries are ubiquitous in 4d QFT!**



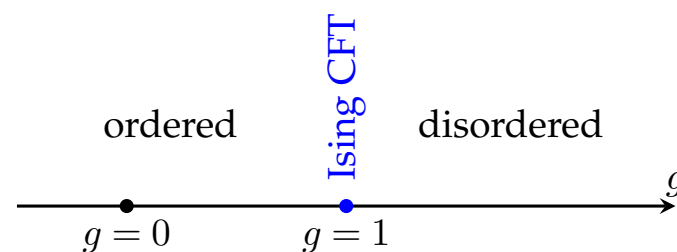
## Non-Invertible Symmetries in the Ising CFT

Transverse field **Ising model**:  $\mathcal{H} = (\mathbb{C}^2)^L$  with nearest neighbor Hamiltonian

$$H = - \sum_j \sigma_j^z \sigma_{j+1}^z - g \sum_j \sigma_j^x .$$

There is a  $\mathbb{Z}_2$  spin flip symmetry  $\eta = \prod_j \sigma_j^x$ .

- $g = 0$ : two ground states,  $|\uparrow^L\rangle$  and  $|\downarrow^L\rangle$ : "ordered phase"
- $g \gg 1$ : ground state preserves the  $\mathbb{Z}_2$ : "disordered phase"
- $g = 1$ : critical Ising CFT at  $c = 1/2$ .



**Kramers-Wannier duality  $N$ :**

$\sigma_i^x \rightarrow \sigma_j^z \sigma_{j+1}^z$  and  $\sigma_j^z \sigma_{j+1}^z \rightarrow \sigma_{j+1}^x$ , corresponds to  $g \rightarrow g^{-1}$ .

At  $g = 1$ : **symmetry of the critical Ising CFT**, which is non-invertible

$$N^2 = 1 + \eta$$

## Non-Invertible Symmetries in $d = 4$

- 4d Kramers-Wannier duality symmetries:

[Kaidi, Ohmori, Zheng][Choi, Cordova, Hsin, Lam, Shao]

$$\text{QFT} \cong \text{QFT}/D \quad \Rightarrow \quad \text{non-invertible 0-form symmetry}$$

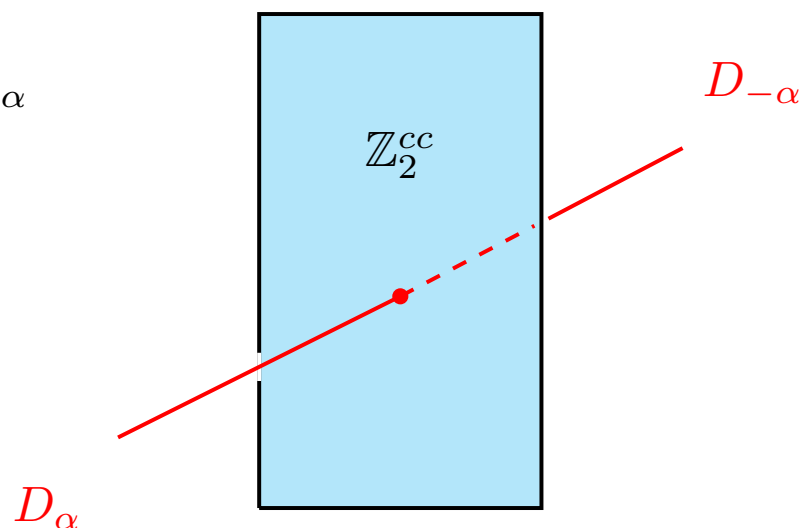
- Gauging charge conjugation [Bhardwaj, Bottini, SSN, Tiwari]

Example:  $O(2)$  gauge theory as  $U(1)/\mathbb{Z}_2^{\text{cc}}$ .

# There is a 1-form symmetry:  $D_\alpha := e^{i\alpha \int *F}$ .

#  $\mathbb{Z}_2^{\text{cc}} : D_\alpha \rightarrow D_{-\alpha}$

$$D_\alpha^{\text{inv}} = D_\alpha \oplus D_{-\alpha}$$



## Example: ABJ Anomaly as a Non-Invertible Symmetry

In 4d QED with massless charge 1 Dirac fermion

$$\mathcal{L}_{\text{QED}+\Psi} = \frac{1}{4e^2} F_{\mu\nu} F^{\mu\nu} + i\bar{\Psi} (\partial_\mu - iA_\mu) \gamma^\mu \Psi$$

the axial current  $j_\mu = \frac{1}{2} \bar{\Psi} \gamma_5 \gamma_\mu \Psi$  is not conserved due to the ABJ anomaly

$$d \star j = \frac{1}{8\pi^2} F \wedge F$$

Define an operator dressed by 3d Topological QFT that has opposite anomaly

$$\mathcal{N}_{\frac{1}{N}}(M_3) = \int [Da] \exp \left( \int_{M_3} \frac{2\pi i}{N} \star j + \frac{iN}{4\pi} ada + \frac{i}{2\pi} adA \right).$$

It is topological, but satisfies non-invertible fusion

$$\mathcal{N}_{\frac{1}{N}} \times \mathcal{N}_{\frac{1}{N}}^\dagger = \mathcal{C} = \text{condensation defect for 1-form symmetry}$$

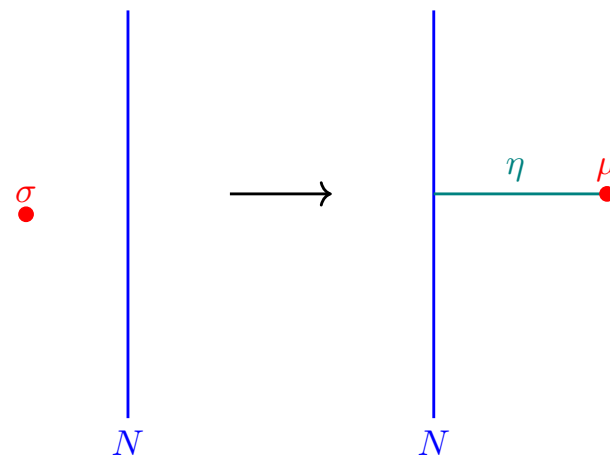
- [Choi, Lam, Shao][Cordova, Ohmori] application to pion decay
- [Cordova, Hong, Koren, Ohmori]  $Z'$  with non-invertible chiral symmetry, gives breaking by exponentially small amount, application to neutrino masses

Status: implications so far mostly known, but promising direction.

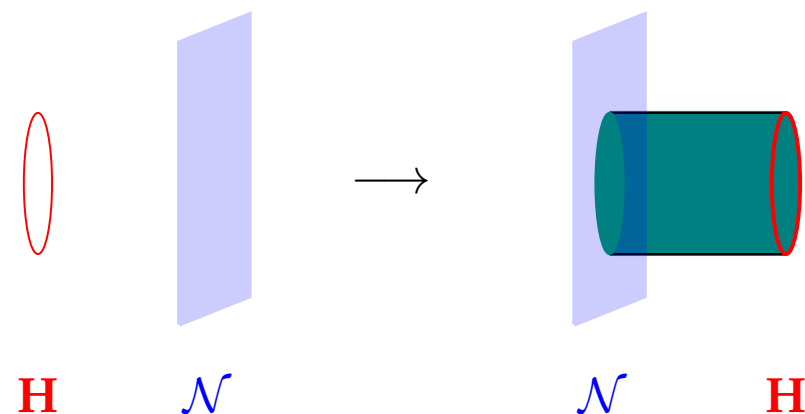
## Physical Implications of Non-Invertible Symmetries 1.

They map genuine operators to non-genuine, i.e. endpoints of extended (topological) defects, or: **order to disorder operators**.

**Example: Ising CFT**  $N^2 = 1 + \eta$



**Example: Witten effect.** 4d  $SO(3)$  SYM has non-invertible symmetry, maps 't Hooft loop to flux attached 't Hooft loop



## Physical Implications 2.: Modified Crossing Relations

Non-invertible symmetries lead to modified crossing relations for S-matrices!

Example: (1+1)d CFTs have non-invertible symmetries  $\mathcal{L}$  e.g. Ising model.

Relevant, integrable deformations can preserve some of  $\mathcal{L}$ .

$\Rightarrow$  IR are gapped vacua.  $\mathcal{L}$  constrains S-matrix of kinks through Ward ids:

$$S_{dc}^{ab}(\theta) = \sum_g \text{Diagram 1} = \sum_g \text{Diagram 2}$$

The diagrammatic equation shows the S-matrix element  $S_{dc}^{ab}(\theta)$  as a sum over states  $g$ . The first diagram shows a scattering process with incoming lines  $a$  and  $c$  and outgoing lines  $b$  and  $d$ , with angles  $\theta_1$  and  $\theta_2$ . The second diagram shows a crossing relation involving a symmetry operator  $\mathcal{L}$  (dashed blue line) and a kink  $g$  (grey circle).

[Copetti, Lucia Cordova, Komatsu] showed: crossing incompatible with symmetry/integrability/unitarity. Consistency implies **modified crossing**

$$S_{dc}^{ab}(\theta) = \sqrt{\frac{d_a d_c}{d_b d_d}} S_{ad}^{bc}(i\pi - \theta), \quad d_a = \langle \mathcal{L}_a \rangle \quad (1)$$

Modified crossing expected in (3+1)d (e.g. massive fermion-dyon scattering).

Status: Modified crossing direct implication of non-invertible symmetries.  
Compelling if in particular extendable to higher dims.

## Physical Implications 3.: Constraining Phases of Matter

Non-invertible Symmetries lead to new IR phases, and new second order Phase Transitions!

**Landau paradigm:**

A continuous (2nd order) phase transition is a symmetry breaking transition.

Gapped Phases:  $G$  spontaneously broken (SSB) to subgroup  $H$ . Phase has  $|G/H|$  vacua, which are acted upon by the broken symmetry.

**Beyond Landau:**

$\mathcal{S}$  be a non-invertible symmetry, then there are new gapped and gapless phases [Bhardwaj, Bottini, Pajer, SSN][Bhardwaj, Pajer, SSN, Warman]

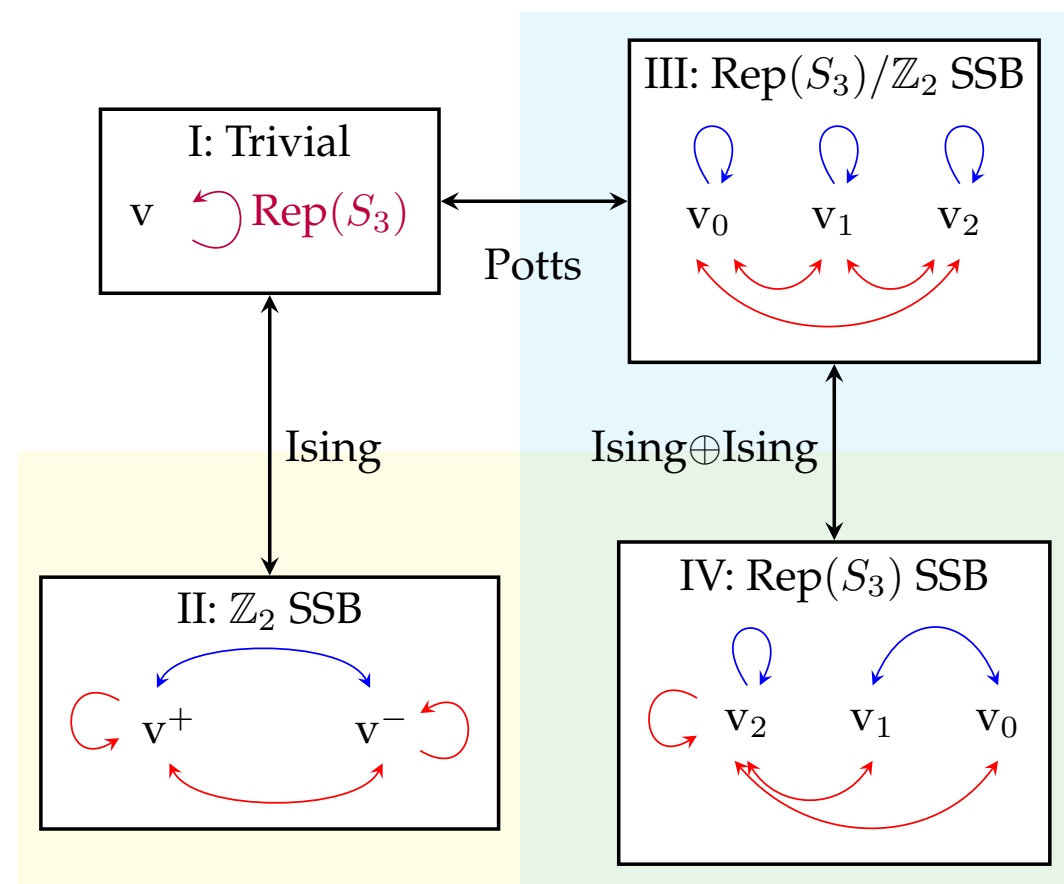
- Determined gapped (topological) phases, order parameters
- Gapless phase transitions between  $\mathcal{S}$ -symmetric gapped phases:



$\Rightarrow$  **Categorical Landau Paradigm** [Bhardwaj, Bottini, Pajer, SSN]

## Example: New Phases from Non-Invertible Symmetries

$\text{Rep}(S_3) = \{1, \sigma, E\}$  with  $\sigma^2 = 1, E^2 = 1 \oplus \sigma \oplus E$  [Bhardwaj, Pajer, SSN, Warman]



Confirmed by lattice models and numerics [Bhardwaj, Bottini, SSN, Tiwari][Chatterjee, Aksoy, Wen]

Status: huge number of predictions, currently discussing how to use "quantum experiments" (using cold atoms) to test these.

## Formal Theory – Status Report

Formal Theory is very much alive and well. Many unexpected, new directions and developments.

- ICHEP Formal Theory session: impressive progress on **Scattering Amplitudes** – string theory (Monteiro, Schlotterer), QFT/holography (Brown<sup>2</sup>, Carrasco, Drummond, Klisch, Lipstein, Parisi), gravitational waves (Brandhuber, Chen, Travaglini), in particular the latter, having direct phenomenological implications.
- **Holography** continues to be a huge source of new approaches to QFT and QG alike: precision results in AdS/CFT, new flat space holography proposals.
- **Generalized (non-invertible) symmetries** in QFTs and spin systems: revolutionarizes classification of phases. Direct (as in collaborative and computational) relevance for cond-mat.
- **Quantum Gravity constraints on EFT**: the swampland/landscape program is making tremendous progress exploring which EFTs are UV completed within QG/string theory.